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The Synergistic Mechanism of Graphene and MoS₂ for Hydrogen Generation: Insights from Density Functional Theory

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The synergistic effect of the graphene and MoS₂ was investigated by using density functional theory (DFT) calculations on the enhanced photocatalytic H₂ production activity of TiO₂/graphene/MoS₂ ternary nanoparticle. Our results indicate that it can form a weak covalent bond between the Ti atom of TiO₂ nanocluster and the nearest C atom on graphene, which not only makes the original degenerate C(2p) orbital level of the graphene (part of the conduction band energy level) split, resulting in the production of a lower level of C(2p) that is easier to accept the excited electron from the Ti(3d) orbital, but also forms a +/- sequence electric field in the interface between them. It is conducive that the electron moves from the TiO₂ cluster to the graphene. In addition, we also find that the band gap of the TiO₂ cluster can be doped by the graphene and MoS₂, and the conduction band consists predominantly of C(2p), S(3p) and Mo(4d) orbital energy level near the Fermi level. These results illustrate the excited electron will eventually accumulate in the graphene or MoS₂ film, which can effectively enhance the separate between the excited electrons and the holes in TiO₂ clusters, thereby to increase the efficiency of hydrogen evolution. Our results are consistent with the experimental results, and can provide some valuable information for the design of photocatalytic composites.

1. Introduction

Converting solar energy into hydrogen is considered to be a valid strategy to solve the energy crisis.¹⁻⁴ It is effective to use the heterogeneous photocatalytic materials for hydrogen generation, because it can strengthen the separation of photo-excited electrons and holes in the semiconductor (such as Titanium dioxide (TiO₂)), and then the hydrogen evolution and oxygen evolution reaction can be occurred at the electrons and holes on the corresponding surface of the photocatalysts, respectively.⁵⁻⁷ So, some comprehensive reviews have summarized the heterogeneous photocatalytic materials during the latest decade.⁸⁻¹⁰

In recent years, graphene-based heterogeneous photocatalytic materials are attracted considerable attentions. This is mainly due to a unique sp² hybrid carbon network resulting in many of highly significant features, such as the ultrafast electron mobility at room temperature (200000 cm²V⁻¹s⁻¹), the conductivity (10⁶ Scm⁻¹), the large surface area of the theoretical (~2600 m²g⁻¹) and the high work function (4.42 eV), etc.¹¹⁻¹³ These excellent properties of graphene showed great potential to become the component of the photocatalysts for hydrogen generation. Therefore, many researchers focus on ternary and/or multicomponent photocatalysts containing photocatalysts, graphene and cocatalyst. For example, Hou *et al* developed the CdS QDs/graphene/ZnIn₂S₄ system exhibited highly efficient hydrogen production because of the high hydrothermal stability and efficient electron transfer, and that emphasized the

importance of semiconductor heterostructure as an effective photocatalysts.¹⁴ It is worth mentioning that the interfacial electron transfer capability can be enhanced significantly and the band structure of the photocatalysts can be doped by the graphene, facilitating the charge separation efficiently at different energy levels. Meanwhile, the cocatalyst containing abundant elements like the Ni, Co and Fe can serve as electron collector and active sites for the hydrogen evolution reaction.^{9,15,16} Graphene can also play a role of the cocatalyst, and constitute a synergistic effect, with the other cocatalysts, on improving the photocatalytic efficiency.¹⁷⁻¹⁹ Recently, Yu *et al*^{17a} firstly proposed a three-component composite TiO₂/graphene/MoS₂ containing 0.25wt% graphene. The results show that graphene acts as an electron reservoirs, and MoS₂ can act as a source of active adsorption sites, to achieve synergetic highly efficient H₂ evolution with the H₂ production rate of 165.3 μmol·h⁻¹ which is 75 times that of the TiO₂/MoS₂^{17b}. Very recently, Lin *et al*¹⁹ demonstrated that ZnS loaded with 0.25 wt% graphene and 2 atom% MoS₂ reached a high H₂-evolution rate of 2258 μmol·h⁻¹ which is about 2 times that of ZnS alone. Several review articles have also summarized the latest developments related to graphene-based photocatalysts, and consider that the main reasons why that significantly enhance the photocatalytic efficiency have three main aspects: (1) it can form the space potential difference promoting the photogenerated electron-hole separation effectively; (2) it can tune the bandgap by changing the band structure of the photocatalysts; (3) it can provide catalytic reaction sites for the hydrogen evolution reaction.^{20,21}

From the existing researches, the introduction of graphene is able to lead to a significant improvement on the photocatalytic H₂ production activity. However, the roles of graphene, synergistic effects of various components and the interfacial charge transfer process in ternary photocatalysts remain unclear. And the mechanism of enhancing the photocatalytic efficiency is still unsharp at this stage. So, it is necessary to carry out the detailed investigations on electronic properties electron transfer of ternary photocatalysts. This would give valuable information for the design and synthesis of ternary photocatalysts in the future.

In this work, we investigate the electronic structure properties of the ternary photocatalysts (a case study of TiO₂/graphene/MoS₂), and characterize the synergistic effects of different components and the interfacial charge transfer process of photocatalysts by using the density functional theory (DFT) calculations. We found that it can form a weak covalent bond between the TiO₂ cluster and graphene, and the interface between graphene and TiO₂ cluster can form a +/- sequence electric field. The band gap of the TiO₂ cluster can be doped by the graphene and MoS₂. These results illustrate that the electron moves from the TiO₂ cluster to graphene, and eventually accumulate in the graphene or MoS₂ film.

2. Computational methods

All the calculations including geometry optimization and electronic structure calculations are performed by using first-principles DFT²²⁻²⁵ implemented with the Spanish Initiative for Electronic Simulation with Thousands of Atoms (SIESTA) code.^{26,27} A localized basis set composed of double zeta plus polarization (DZP) was used. The standard norm-conserving Troullier-Martins pseudopotential in the fully nonlocal (Kleinman-Bylander) form was employed. We described the exchange and correlation effects using the generalized gradient approximation (GGA) exchange-correlation function²⁸ developed by Perdew, Burke and Ernzerh.²⁹ The real space mesh cutoff was set to more than 200Ry, The Monkhorst-Pack³⁰ special k-point was set to 5×5×1, and the k-point was set to 9×9×1 on the density of state(DOS) calculation to achieve higher accuracy. We use the conjugated gradient (CG) minimization technique to obtain the optimized structure until the maximum atomic force is smaller than 0.04 eV / Å and the atomic positions were fully relaxed without constraint.

In our simulation, the interface between graphene and MoS₂ is modeled by using a supercell in the x-y plane(Figure. 1). We build the 14.76 × 14.76 Å lateral periodicity of the graphene including 72 atoms and 15.83 × 15.83 Å lateral periodicity of the MoS₂ monolayer including 75 atoms, and the lattice mismatch is 6.8%. After configuration optimization, the last lattice parameters of MoS₂/grapheme system is 15.33 × 15.32 Å in the x-y plane. The vacuum space at z axes is 20 Å in the direction normal to the interface, representing the isolated slab boundary condition. The lattice of MoS₂ was set to match to that of graphene in the supercell. The supercells are then fully relaxed for both the lattice constants and the atomic geometry. The (TiO₂)₃ configuration contains all the classical features of the well-understood R(101) surface of anatase TiO₂ which has been widely used in the literature³¹, and the results are in good agreement with the experimental data. In order to understand the photocatalytic hydrogen production synergistic mechanisms by

comparing TiO₂/graphene/MoS₂ system, we also calculated the two other systems (TiO₂/graphene and graphene/MoS₂); whose lattice parameters and calculated parameters is the same as the TiO₂/graphene/MoS₂ system. To find the most favorable adsorption configurations, the (TiO₂)₃ cluster was made initially to approach the graphene sheet in different positions, and various possible adsorbing orientations were investigated. After full relaxation, the optimized configurations obtained from the different initial states were compared, and the most stable configurations is shown in Fig. 1 with the Ti2 atom located above the C-C bond of graphene.

To reflect the combination of these three systems stability, we calculated the binding energy (E_b) of the TiO₂/graphene/MoS₂ system, which are showed as follow:

$$E_{b1} = E_{\text{tot}}[\text{TiO}_2/\text{grapheneMoS}_2] - E_{\text{tot}}[\text{grapheneMoS}_2] - E_{\text{tot}}[\text{TiO}_2] \quad (1)$$

$$E_{b2} = E_{\text{tot}}[\text{TiO}_2/\text{grapheneMoS}_2] - E_{\text{tot}}[\text{TiO}_2/\text{geaphene}] - E_{\text{tot}}[\text{MoS}_2] \quad (2)$$

$$E_{b3} = E_{\text{tot}}[\text{TiO}_2/\text{grapheneMoS}_2] - E_{\text{tot}}[\text{MoS}_2] - E_{\text{tot}}[\text{graphene}] - E_{\text{tot}}[\text{TiO}_2] \quad (3)$$

Here, $E_{\text{tot}}[\text{TiO}_2/\text{graphene/MoS}_2]$, $E_{\text{tot}}[\text{graphene/MoS}_2]$ and $E_{\text{tot}}[\text{TiO}_2/\text{geaphene}]$ are the total energies per supercell of the TiO₂/graphene/MoS₂, graphene/MoS₂ and TiO₂/graphene systems, respectively. $E_{\text{tot}}[\text{TiO}_2]$, $E_{\text{tot}}[\text{MoS}_2]$ and $E_{\text{tot}}[\text{graphene}]$ are those for isolated TiO₂ cluster, isolated MoS₂ and isolated graphene respectively. The isolated TiO₂ cluster reference energy has been computed in a 30Å × 30Å × 30Å cell.

To analyse the nature of the TiO₂/graphene/MoS₂ interaction, we also have examined electron density differences at the interface, $\Delta\rho(r)$ and the planar-averaged density difference, $\Delta\rho^{\text{avg}}(z)$ here defined as:

$$\Delta\rho(r) = \rho_{\text{tot}}[\text{TiO}_2 / \text{graphene} / \text{MoS}_2] - \rho_{\text{tot}}[\text{MoS}_2] - \rho_{\text{tot}}[\text{graphene}] - \rho_{\text{tot}}[\text{TiO}_2] \quad (4)$$

$$\Delta\rho^{\text{avg}}(z) = \int_{\Sigma(z)} dx dy \Delta\rho = \sum_{ij} \Delta x_i \Delta y_j \Delta\rho_{i,j} \quad (5)$$

$\rho_{\text{tot}}[\text{TiO}_2 / \text{graphene} / \text{MoS}_2]$ is the electron densities of TiO₂/graphene/MoS₂, $\rho_{\text{tot}}[\text{MoS}_2]$, $\rho_{\text{tot}}[\text{graphene}]$ and $\rho_{\text{tot}}[\text{TiO}_2]$ are those for the MoS₂, graphene and TiO₂ systems, respectively. The plan averaged electron density difference is obtained by integrating the electron density differences along x-y plane, where the z axis is along the surface normal direction and x/y is the cross-section of the supercell in the x-y plane at z, the i and j are one point divided lattice axis and b-axis lattice, respectively.

3. Results and discussions

3.1. Geometric structures and binding energies

To explore the roles of graphene and MoS₂ nanocluster in this ternary photocatalysts, it is necessary to consider the three systems: (a) TiO₂/graphene, (b) graphene/MoS₂ and (c) TiO₂/graphene/MoS₂. The optimized configurations are shown in Figure 1. In three arrangements, graphene keeps its plane and hexagonal atomic network, which is consistent with the previous research³². From the binding energies and the interfacial distances of the three optimized structures, Table 1, it can be found that there is no strong interaction between the TiO₂ cluster (or MoS₂)

and graphene which consists with the previous theoretical result. The binding energy of $\text{TiO}_2/\text{graphene}$ (-3.43 eV) is greater than that of $\text{MoS}_2/\text{graphene}$ (-1.19 eV). So it is easier for graphene to combine TiO_2 cluster than MoS_2 . The binding energy of $\text{TiO}_2/\text{graphene}/\text{MoS}_2$ is less than that of $\text{TiO}_2/\text{graphene}$, indicating that MoS_2 can affect the combination between TiO_2 and graphene. This conclusion is in agreement with the distance between TiO_2 and graphene below. To obtain accurate van der Waals energies, the van der Waals density functional (vdW-DF) as proposed by Dion et al.³² and recently implemented in the SIESTA code³³ was used to check the binding energy. The derivation of binding energy with vdW-DF is within 0.05eV compared with the GGA method (Table 1). It is noted that the implement of calculation van der Waals Force is ineffective at present state.

Table1: the binding energy (E_b), the interfacial distance D in the three optimized structure.

	$E_b(\text{eV})$	$E_{b, \text{vdW}}(\text{eV})$	$D (\text{\AA})$
$\text{TiO}_2/\text{graphene}$	-3.43 ^a	-3.38 ^a	2.43 ^c
$\text{graphene}/\text{MoS}_2$	-1.19 ^b	-1.16 ^b	3.80 ^d
$\text{TiO}_2/\text{graphene}/\text{MoS}_2$	-2.66 ^a	-2.62 ^a	2.40 ^c / 3.84 ^d

^a Binding energy between TiO_2 and graphene, ^b binding energy between MoS_2 and graphene, ^c the shortest distance between the TiO_2 cluster and the graphene, ^d the shortest distance between the MoS_2 nanocluster and the graphene.

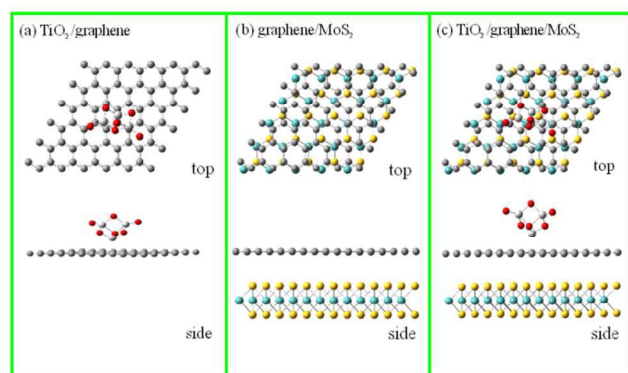


Figure 1. The top and side view of the optimized structure of the three composites: (a) $\text{TiO}_2/\text{graphene}$, (b) $\text{graphene}/\text{MoS}_2$ and (c) $\text{TiO}_2/\text{graphene}/\text{MoS}_2$. The red, white, grey, yellow and blue represent O, Ti, C, S, Mo, respectively.

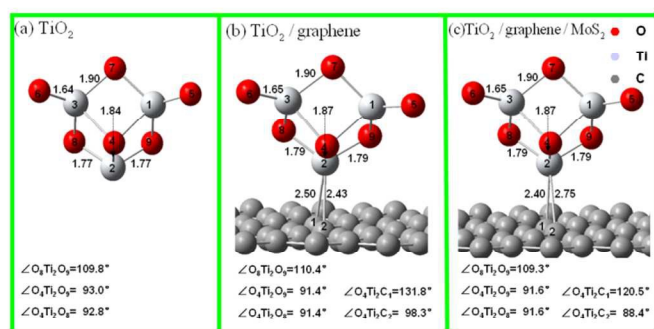


Figure 2. The bond distances and angles between $(\text{TiO}_2)_3$ and graphene in the optimized structures of (a) $(\text{TiO}_2)_3$ cluster, (b) $\text{graphene}/\text{TiO}_2$ and (c) $\text{TiO}_2/\text{graphene}/\text{MoS}_2$.

To quantitatively characterize the interaction between TiO_2 nanocluster and graphene, we draw the bond lengths and angles between $(\text{TiO}_2)_3$ and graphene among the optimized structures, which are shown in Figure 2. As seen from Figure 2a, TiO_2 nanocluster has C_2 symmetry. The shortest and longest bond lengths of Ti-O are 1.64 \AA ($\text{Ti}_3\text{-O}_6$) and 1.90 \AA ($\text{Ti}_3\text{-O}_7$) in TiO_2 nanocluster, respectively. The Ti2 atom is located at the bridge site between TiO_2 nanocluster and graphene in the $\text{TiO}_2/\text{graphene}$ and $\text{TiO}_2/\text{graphene}/\text{MoS}_2$ composites. In the $\text{TiO}_2/\text{graphene}$ system, Figure 2b, TiO_2 keeps the same symmetry as that in Figure 2a. However, compared with the TiO_2 nanocluster, the bonds near the graphene increase a little and the bonds far from graphene have no change. This indicates that graphene affects TiO_2 nanocluster mainly by the interaction with the adjacent Ti2 atom. In the $\text{TiO}_2/\text{graphene}/\text{MoS}_2$ system, Figure 2c, the structure of the TiO_2 nanocluster is consistent with that in $\text{TiO}_2/\text{graphene}$ system, which shows that the introduction of MoS_2 almost has no effect on the structure of TiO_2 nanocluster. From the distances between Ti2 atom and C_1/C_2 atoms of graphene in both Figure 2b and 2c, it can be inferred that there is weak interaction between TiO_2 and graphene.

Meanwhile, we also investigated some key bond angles formed by Ti2 atom with the surrounding O atoms and C atoms in three systems, respectively. As shown in Figure 2, there are little changes of angles $\angle \text{O}_8\text{Ti}_2\text{O}_9$, $\angle \text{O}_4\text{Ti}_2\text{O}_9$ and $\angle \text{O}_4\text{Ti}_2\text{O}_8$ in the three optimized structures. But the angles $\angle \text{O}_4\text{Ti}_2\text{C}_1$ and $\angle \text{O}_4\text{Ti}_2\text{C}_2$ are 131.8° and 98.3° in $\text{TiO}_2/\text{graphene}$ system, which differs from those (120.5° and 88.4°) in $\text{TiO}_2/\text{graphene}/\text{MoS}_2$ system. From Figure 2, the O8, O9, Ti2 and C2 atoms form an approximate plane, and the Ti2-O4 bond is substantially perpendicular to the plane. This means that the 3d orbital of Ti2 atom is hybridized. It is probable that there is an overlap between the Ti2-3d orbital and the C2-2p_z orbital which is perpendicular to the graphene plane. This helps to explain the weak interaction between graphene and TiO_2 cluster.

3.2 Charge density difference

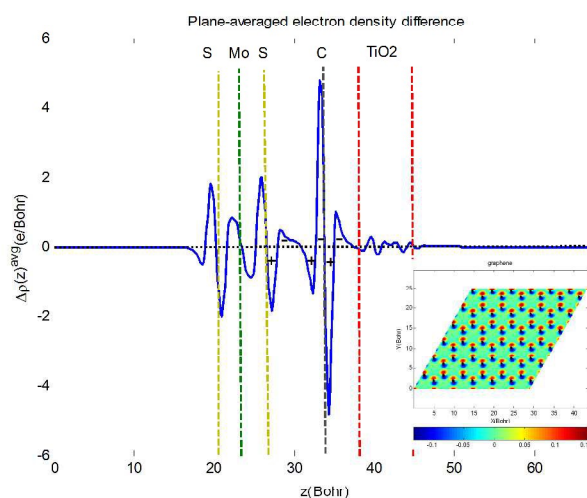
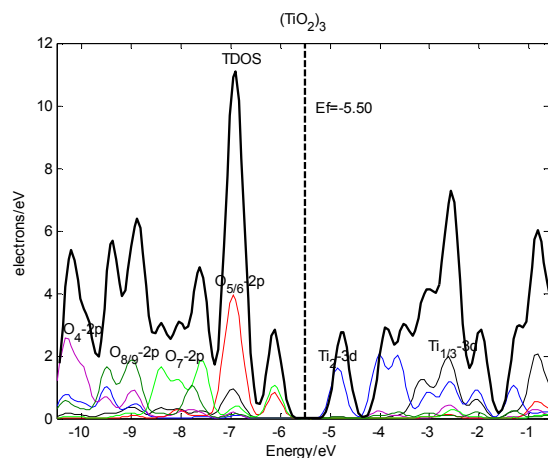


Figure 3. Plane-averaged electron density difference of the $\text{TiO}_2/\text{graphene}/\text{MoS}_2$ system along the vertical axis (z axis). The coloured

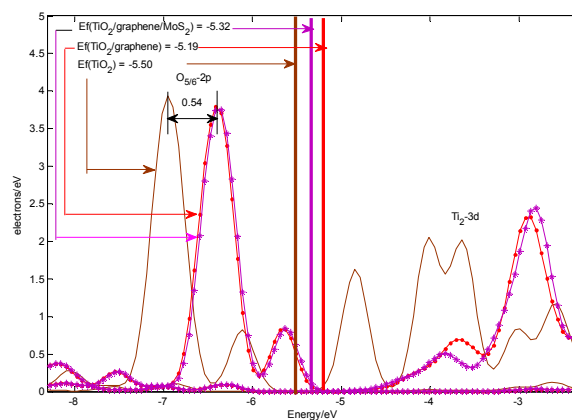
dash lines indicate the positions of each composite, and the $\Delta\rho^{\text{avg}}(z) > 0$ represents the charge accumulation. The insert is the electron density differences at the graphene surface, and these dark dots represent the C atoms on graphene

5 In order to give more detailed understandings of the interface in TiO₂/graphene/MoS₂ system, we plot the plane-averaged electron densities and electron density differences (Figure 3). From Figure 3, it can be seen clearly that graphene is polarized due to TiO₂ and MoS₂ cluster. To further verify this phenomenon, we draw the charge density difference of the surface of graphene. It is obvious that each C atom on graphene is polarized, which forms +/- sequence electric field with the intensity of about 0.1 e/Bohr³S. From the charge polarization of graphene, it is found that charge accumulation occurs on the graphene surface close to
10 MoS₂ and charge depletion happens on the other side close to TiO₂. In the interface between the graphene and the TiO₂ nanocluster, the Ti2 atom is electron withdrawing, resulting in the electron depletion in C atom of the graphene and the electron accumulation in Ti2 atom which forms a +/- sequence electric field finally. It is conducive that the electron moves from the
15 TiO₂ nanocluster to graphene. For the interface between MoS₂ and graphene, the S atom of MoS₂ nearby the graphene are also polarized to form a +/-/-/+ sequence electric field. Although the interaction between MoS₂ and graphene is very weak, a
20 significant electron accumulation is observed in the interlayer region. It indicates that charge transfer is induced between the graphene and the MoS₂ surface, which is also consistent with that of the reference³⁴. Although the large distance between graphene and MoS₂ is able to have slightly effects on the electron transfer
25 from graphene to the MoS₂ surface, it can attract more electrons due to the much more depletion of the MoS₂ surface. So, as a result, it does not preclude the charge transfer between graphene and the MoS₂ surface.

3. 3. The projected density of states (PDOS) of the composites



35 Figure 4. Projected density of states (PDOS) onto the (TiO₂)₃ cluster. The dashed line indicates the Fermi energy level (Ef).



40 Figure 5. Projected density of states (PDOS) of the O_{5/6}-2p and Ti₂-3d of the (TiO₂)₃ cluster in the three composites: (a) TiO₂, (b) TiO₂/graphene and (c) TiO₂/graphene/MoS₂. The coloured middle lines represent the Fermi level, and the same colour line represents the same system. The left nearby the Fermi level is O_{5/6}-2p PDOS, and the right is Ti₂-3d PDOS for each composites.

45 To inquire into the electronic coupling and orbital contributions in the calculated systems, the projected density of state (PDOS) for three calculated systems would be probed. It shows the total density of states (TDOS) of (TiO₂)₃ nanocluster and each valence electron orbit of Ti and O atoms in Figure 4. As from Figure 4,
50 the valence band of the TiO₂ nanocluster is composed mainly of the O_{5/6} 2p orbitals (O_{5/6} represented the 5th and 6th O atoms), and the conduction band is mainly composed of Ti₂-3d orbital. The band gap calculated by using GGA and LDA+U is only 1.36 and 1.68 eV, respectively, and is far less than the experimental
55 value (~3.2eV).³⁵ It is caused by three main reasons: (1) the well-known shortage that DFT underestimates the band gap due to the self-correlation error of electrons and inherent lack of derivative discontinuity; (2) the difference existed between a small cluster and bulk matter because of the size effect^{31,35,36}. Besides, as
60 shown in Figure 4, the electron numbers are not the same for some O-2p and Ti-3d orbitals. This is because there are two types of Ti atoms and four types of O atoms of different environments, Figure 2. The closer to the Fermi level, the more of the electrons is. Furthermore, combined with the TiO₂ cluster structure in
65 Figure 2a, on the one hand, we found that the O_{5/6} atoms not meet the eight electron configuration, so that they have high reactivity with unpaired electrons. So it is easy for O_{5/6}-2p energy levels to close to the Fermi level. On the other hand, the Ti₂-3d energy level is closest to the conduction band among all
70 the Ti-d energy levels. It can be seen that the Ti₂ atom also does not reach a stable outer electron configuration from Figure 2. Because of the atoms unreached stable configuration on the surface of the (TiO₂)₃ cluster, the number of electrons nearby the Fermi level increase. It can reduce the energy required for the
75 being excited electrons.

In order to understand the impact of graphene and MoS₂ on the O-2p/Ti-3d orbitals and the Fermi level in (TiO₂)₃ nanocluster, we also plot the O_{5/6}-2p and Ti₂-3d PDOS for the three calculated systems in Figure 5. Compared to the isolated TiO₂ nanocluster
80 system, it can be seen the graphene has more significant effects on O-2p and Ti-3d of the TiO₂ cluster than the MoS₂ in the TiO₂/graphene and TiO₂/graphene/MoS₂ system. On one hand,

because of introduction of graphene, the O-2p orbital peak level shifts to the right by 0.54 eV. The corresponding Fermi level shifts by 0.31 eV. So the graphene make the O5/6 peak energy level closer to the Fermi level. Besides, the influence of graphene on Ti2-3d peak energy level is more obvious. It can be seen clearly that the Ti2-3d peak energy level in the TiO₂/graphene system shifts to the right, and the peak changes dramatically. The main reason is that the Ti2-3d orbital interacts with the C1/C2-2p orbital of the graphene, making the Ti2-3d orbital energy level more localized at the higher level. On the other hand, the MoS₂ almost does not take any effect on shifting the O5/6 energy level. Though MoS₂ shows larger influence the impact on Ti2-3d peak energy level than that on O5/6-2p peak energy level, this influence is small compared to that resulted by graphene. However, it is interesting that the MoS₂ makes the Fermi level moves to the left by 0.14 eV. This show the O5/6-2p orbitals of the TiO₂/graphene/MoS₂ is closer to the Fermi level compared with that in the TiO₂/graphene system, which means that the electrons on O_{5/6}-2p orbitals of TiO₂/graphene/MoS₂ are more likely to be excited by the light. Apart from the O5/6 and Ti2 atoms, it can show the consistency for the analysis of other O and Ti atoms. On the whole, the graphene and MoS₂ can both make O-2p orbit closer to the Fermi level, especially the graphene. This indicates that they have the synergistic influence on making electrons of the O-2p be more easily excited to the conduction band.

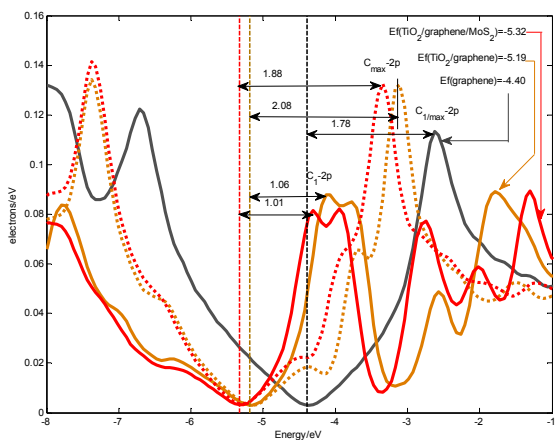


Figure 6. Projected density of states (PDOS) of the C₁-2p and C_{max}-2p of the graphene in the three composites: (a) graphene, (b) TiO₂/graphene and (c) TiO₂/graphene/MoS₂. Middle coloured dotted line represents the corresponding with the Fermi level position; the same colour line represents the same system.

To clearly understand the process of the excited electrons excited to the graphene by the covalent band formed Ti2 and the nearest C atoms (marked C₁ atom), we studied the influence of TiO₂ and MoS₂ on C-2p orbital of graphene. As can be seen from Figure 6, the TiO₂ and MoS₂ can affect C₁/C_{max}-2p (C_{max} is the C atom that has the maximum peak energy level). The peak energy levels of C-2p are the same in pure graphene system, but they show the difference in TiO₂/graphene or TiO₂/graphene/MoS₂ system with the value and shape of C₁-2p changed significantly. For the C₁ atoms in TiO₂/graphene, the main peak is splitted into two main peaks (peak energy values are -4.128 and -1.784 eV, respectively), and the energy integral value of the conduction band near the

Fermi level is larger than that of the other C atoms. The interaction between Ti2-3d and C1-2p orbitals leads to the C1-2p orbital localized and making the energy level closer to the Fermi level. For the C₁ atoms in TiO₂/graphene/MoS₂, the main peak is splitted into three peaks (peak energy values are -4.309, -2.745 and -1.303 eV, respectively). It not only reflects the interaction between C1-2p and Ti2-3d orbital, but also shows the MoS₂ also have an influence on C1-2p orbital. Meanwhile the MoS₂ also makes C1-2p orbital slightly offset to the left. Besides, the Fermi level is shifted to the left for the TiO₂/graphene or TiO₂/graphene/MoS₂ systems. The difference is 1.78 eV between the C-2p orbital peak energy level and the Fermi level for pure graphene; while the maximum/minimum difference for the TiO₂/graphene and TiO₂/graphene/MoS₂ systems are the 2.08/1.01 eV and 1.88/1.01 eV, respectively. This shows that the sp² orbital are destroyed by the TiO₂ clusters and MoS₂ splitting into lower and higher located C-2p orbital levels. Overall, the Ti1 atom makes the main distinct peak of the C-2p orbital spitting into more peaks, resulting in a lower localized C-2p orbital; while the MoS₂ can make C-2p level shift toward the Fermi level. To further explore the component of the valence and conduction band of the TiO₂/graphene/MoS₂ and the process of charge transfer, we also draw the PDOS of the TiO₂/graphene/MoS₂ Figure 7.

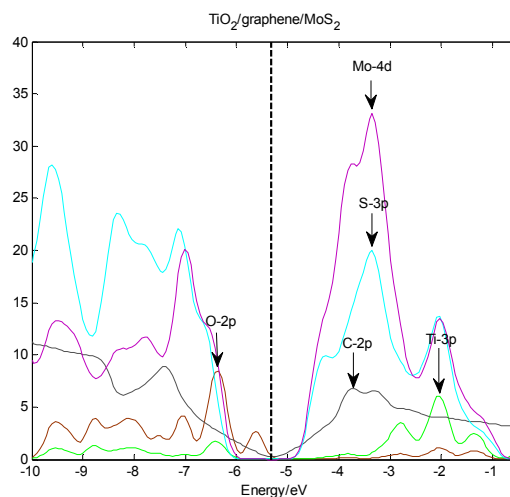
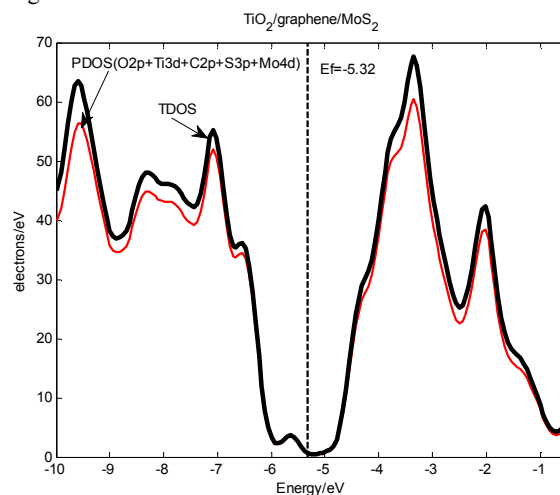


Figure 7. The Total density of states(TDOS)(left) and Projected density of states (PDOS)(right) of the TiO₂/graphene/MoS₂, middle dotted line represents the corresponding with the Fermi level position, each of the major peaks is pointed by the arrow.

From the up picture in Figure 7, the components of the valence and conduction band compose mainly of O-2p, Ti-3d, C-2p, S-2p and Mo-4d orbitals. The valence band is O-2p orbital, and the conduction band are C-2p, S-3p and Mo-4d Composition. The band gap of (TiO₂)₃ nanocluster, TiO₂/graphene and TiO₂/graphene/MoS₂ system are 1.36, 0.80 and 0.02eV, respectively. This indicates that the band gap of semiconductor photocatalysts can be adjusted by doping other material, such as the graphene and MoS₂. From Figure 7, it can be seen that the C-2p, S-3p and Mo-4d peaks exist between the O-2p and Ti-3d energy levels. This provides a sign that the excited electrons eventually appear in the C-2p, S-3p and Mo-4d energy levels. That can effectively enhance the separate between the excited electrons and the holes in TiO₂ clusters, thereby to increase the efficiency of hydrogen evolution. Our calculation results agree with the experimental results.^{17a}

4. Conclusions

In summary, DFT calculations were performed to investigate the heterogeneous TiO₂/graphene/MoS₂. The geometric structures and electrons properties were calculated to clarify the interface interaction in electronic level. And also, graphene and MoS₂ in the enhanced photocatalytic H₂ production activity of TiO₂ nanoparticle was hold key roles. Based on the analysis of the calculation results, we obtain the following conclusions.

(1) The graphene and MoS₂ can make O-2p orbital closer to the Fermi level, especially the graphene, indicating that they have the synergistic effect on making electrons of the O-2p more easily excited to the conduction band. (2) It can form a covalent bond between the Ti atoms of TiO₂ nanoclusters and the nearest C atoms on graphene, resulting in the production of a lower C-2p energy level, which is easier to accept the excited electron from the Ti-3d orbital. (3) The interface between the graphene and the Ti atoms of TiO₂ cluster can form a +/- sequence electric field. It is conducive that the electron moves from the TiO₂ cluster to the graphene. (4) The band gap of the TiO₂ cluster can be doping by the graphene and MoS₂. In addition, the conduction band consists predominantly of C-2p, S-3p and Mo-4d orbital energy level near the Fermi level, it illustrates the excited electron will eventually accumulate in the graphene or MoS₂ film.

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Notes and references

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