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## Are the Orientation and Bond Strength of the $\text{RCO}_2^- \cdots \text{M}$ Link Key Factors for Ultrafast Electron Transfers?

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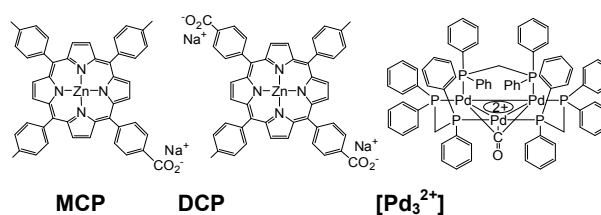
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**The photo-induced electron transfers in the straight up ionic assemblies  $[\text{Pd}_3^{2+}] \cdots \text{MCP}$  and  $[\text{Pd}_3^{2+}] \cdots \text{DCP} \cdots [\text{Pd}_3^{2+}]$  ( $[\text{Pd}_3^{2+}]^* \rightarrow \text{MCP}$  or  $\text{DCP}$ ) are ultrafast ( $< 85$  fs) indicating that it is not necessary to have a strong coordination bond or a bent geometry to obtain fast electron transfer.**

Porphyryns are heavily investigated for dye sensitized solar cell (DSSC) applications notoriously in their “push-pull” versions.<sup>1</sup> Numerous issues concerning the fast electron injection from the excited dye to the  $\text{TiO}_2$  nanoparticles were raised. For examples, the use of carboxylates vs phosphates, directly comparing their interactions with  $\text{Ti}^{4+}$ ,<sup>1c,2</sup> and the presence of J-aggregates,<sup>3</sup> were recently debated parameters. It also became clear that the DSSCs efficiency depends not only on the rate of electron injection but also on the relative amount of dyes deactivating without electron transfer.<sup>4</sup> This rate is predictably distance-dependent between the dye and  $\text{TiO}_2$  center,<sup>5</sup> which also contribute in retarding the competing charge recombination.<sup>6</sup> Moreover, recent studies also showed that the binding geometry plays a determining role on the electron transfer rate,  $k_{\text{et}}$ .<sup>7</sup> The current belief is that a strong O-Ti coordination bond and a bent conformation of the dye over the  $\text{TiO}_2$  surface promote ultrafast  $k_{\text{et}}$ . The typically reported  $k_{\text{et}}$ 's are in the fs-ps time scale,<sup>2-7</sup> going as low as  $< 50$  fs (*i.e.* the time resolution limit).<sup>8</sup> To accurately define all parameters that lead to fast electron injection, it now appears necessary to examine weakly bonded models with a “straight up” geometry; two seemingly unfavourable parameters. In such cases,  $k_{\text{et}}$  should be significantly slower (long ps). In this respect, the easily reducible unsaturated  $\text{Pd}_3(\text{dppm})_3(\text{CO})^{2+}$  cluster ( $[\text{Pd}_3^{2+}]$ ;  $\text{dppm} = \text{Ph}_2\text{PCH}_2\text{PPh}_2$ ; Chart 1;  $E^{0/1} \sim -0.50$  V vs SCE as  $\text{CF}_3\text{CO}_2^-$  salt in various solvents)<sup>9</sup> is an appropriate model as it weakly binds carboxylate ions “straight up” above the  $\text{Pd}_3$  plane strictly *via* ionic interactions (based on X-ray data using  $\text{CF}_3\text{CO}_2^-$  ions).<sup>10</sup> This geometry is favoured *via* host-guest interactions of the  $\text{RCO}_2^-$  group inside the cavity formed by the phenyls of the dppm ligands above the  $\text{M}_3$  frame.



Scheme 1. Structures of MCP, DCP and  $[\text{Pd}_3^{2+}]$  (as  $\text{PF}_6^-$  salt).

Despite the expectation for a slow rate, we now report that the photo-induced electron transfers for the  $[\text{Pd}_3^{2+}] \cdots \text{MCP}$  and  $[\text{Pd}_3^{2+}] \cdots \text{DCP} \cdots [\text{Pd}_3^{2+}]$  assemblies (MCP, DCP = 5-(4-carboxylphenyl)-10, 15, 20-tristoyl(porphyrinato)zinc(II), 5, 15-bis(4-carboxylphenyl)-15, 20-bistoyl(porphyrinato)zinc(II), respectively; as  $\text{Na}^+$  salts; Scheme 1) are in fact also ultrafast ( $< 85$  fs).

The host-guest assemblies (eqs. 1 and 2) are generated in MeOH solutions easily monitored by UV-vis spectroscopy (detail in ESI). The binding constants,  $K_{11}$  and  $K_{12}$ , were extracted using the Scatchard, Scott and Benesi-Hilderbrand plots.<sup>11,12</sup> The  $K_{11}$  and  $K_{12}$  values for the sequential formation of  $[\text{Pd}_3^{2+}] \cdots \text{DCP}$  and  $[\text{Pd}_3^{2+}] \cdots \text{DCP} \cdots [\text{Pd}_3^{2+}]$  are reasonably assumed to be the same as no slope change in the low and high concentration range of  $[\text{Pd}_3^{2+}]$  was detected. This effect is consistent with the absence of steric hindrance. The similarity of the 19300 and 22000  $\text{M}^{-1}$  constants readily attests this assumption.

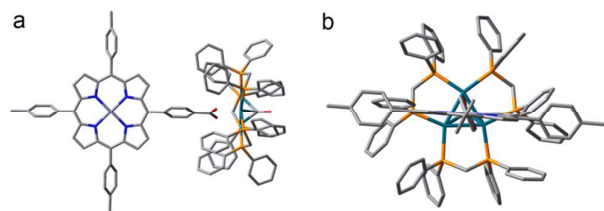


Previous X-ray data indicated a “straight up” geometry of the  $\text{CF}_3\text{CO}_2^- \cdots \text{M}_3$  unit with ionic  $\text{Pd} \cdots \text{O}$  distances of 2.57 to 3.06 Å with the  $\text{CO}_2$  unit placed closer to one of the Pd-Pd bond.<sup>10</sup> In the absence of X-ray data, the optimized geometries of the assemblies were calculated using DFT and confirmed the expected geometry (Figure 1; see ESI for  $[\text{Pd}_3^{2+}] \cdots \text{DCP} \cdots [\text{Pd}_3^{2+}]$ ). The computed “gas phase”  $\text{Pd} \cdots \text{O}$  distances ranging from 2.411 to 3.368 Å are also consistent with ionic interactions.

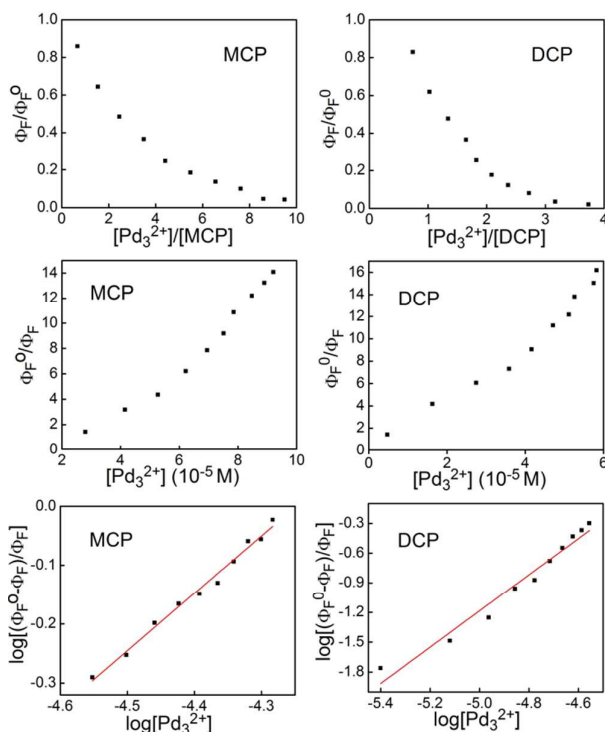
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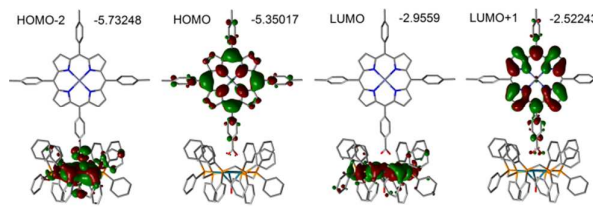


**Figure 1.** Side (a) and top (b) views of the optimized geometry (DFT) of the  $[\text{Pd}_3^{2+}] \bullet \bullet \bullet \text{MCP}$  host guest assembly (ground state).  $\text{Pd} \bullet \bullet \bullet \text{O}$  distances: 1<sup>st</sup> O; 3.368, 3.298, 2.416, 2<sup>nd</sup> O; 3.516, 3.175, 2.411 Å. The O-atoms (red) are placed more closely above a Pd-Pd bond rather than being in the center. For  $[\text{Pd}_3^{2+}] \bullet \bullet \bullet \text{DCP} \bullet \bullet \bullet [\text{Pd}_3^{2+}]$ , see ESI.



**Figure 2.** Top: graphs reporting the decrease of the relative fluorescence intensity of **MCP** and **DCP** upon addition of  $[\text{Pd}_3^{2+}]$  ( $\Phi_F$  and  $\Phi_F^0$  are the intensity in the presence and absence of  $[\text{Pd}_3^{2+}]$ , respectively). Middle: Stern-Volmer plots of the fluorescence quenching of **MCP** and **DCP** by  $[\text{Pd}_3^{2+}]$ . Bottom: graph reporting  $\log[(\Phi_F^0 - \Phi_F)/\Phi_F]$  vs  $\log[\text{Pd}_3^{2+}]$ . The red lines are the best fits allowing extracting  $n$ .

The **MCP** and **DCP** fluorescence bands exhibit quenching upon addition of  $[\text{Pd}_3^{2+}]$  (Figure 2; top) but the fluorescence lifetimes remain constant (ESI) indicating the presence of static quenching (*i.e.* the free dyes and assemblies are emissive and non-emissive, respectively). The lack of linearity in the Stern-Volmer plots (Figure 2; middle) corroborates this conclusion. In such a case, the data are analyzed using the relation  $\log[(\Phi_F^0 - \Phi_F)/\Phi_F] = \log(K_b) + (n \cdot \log[\text{Pd}_3^{2+}])$  where  $\Phi_F^0$  and  $\Phi_F$  are the fluorescence intensities in respectively the absence and presence of  $[\text{Pd}_3^{2+}]$ ,  $K_b$  is the binding constant, and  $n$  is the average number of binding sites (Figure 2; bottom).<sup>13</sup> Values of  $n = 0.97$  (**MCP**) and  $1.83$  (**DCP**) are obtained, which are consistent with the number of carboxylates on the dyes.



**Figure 3.** Selected frontier MO representations of  $[\text{Pd}_3^{2+}] \bullet \bullet \bullet \text{MCP}$  (in eV; see ESI for  $[\text{Pd}_3^{2+}] \bullet \bullet \bullet \text{DCP} \bullet \bullet \bullet [\text{Pd}_3^{2+}]$  and for the HOMO-1 (**MCP**  $\pi$ -level)).

**Table 1.**  $K_D$ ,  $V$  and  $n$  data for the assemblies (solvent = MeOH at 298 K, and MeOH/2MeTHF 1:1 at 77 K).

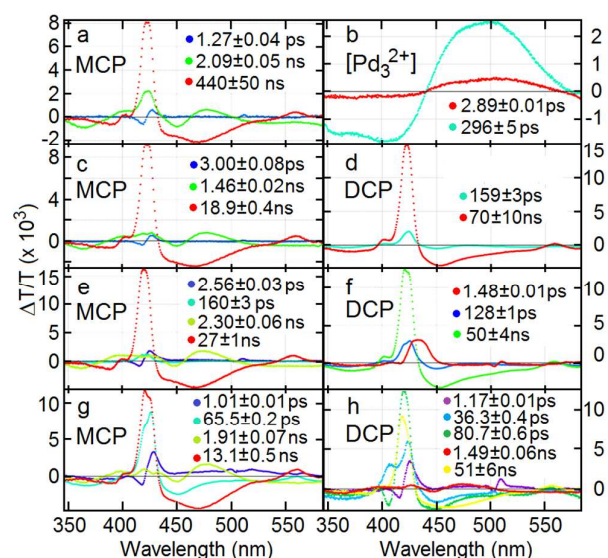
assemblies (temperature)	$n$	$K_D$ ( $\text{M}^{-1}$ )	$V$ ( $\text{M}^{-1}$ )
$[\text{Pd}_3^{2+}] \bullet \bullet \bullet \text{MCP}$ (298 K)	0.97	241	17900
$[\text{Pd}_3^{2+}] \bullet \bullet \bullet \text{DCP} \bullet \bullet \bullet [\text{Pd}_3^{2+}]$ (298 K)	1.83	2750	21300
$[\text{Pd}_3^{2+}] \bullet \bullet \bullet \text{MCP}$ (77 K)	0.91	3400	21900
$[\text{Pd}_3^{2+}] \bullet \bullet \bullet \text{DCP} \bullet \bullet \bullet [\text{Pd}_3^{2+}]$ (77 K)	2.10	3700	25800

**Table 2.** Relative percentage of complexed dyes vs the  $[\text{CO}_2^-]/[\text{Pd}_3^{2+}]$  ratio.

$[\text{CO}_2^-]/[\text{Pd}_3^{2+}]$	$K_{11}$ , $V$ (in $[\text{M}^{-1}]$ )	$K_{12}$ , $V$ (in $[\text{M}^{-1}]$ )
1:1	$[\text{Pd}_3^{2+}] \bullet \bullet \bullet \text{MCP}$ 22.1, 21.1	$[\text{Pd}_3^{2+}] \bullet \bullet \bullet \text{DCP}$ 28.0, 28.0 $([\text{Pd}_3^{2+}])_2 \bullet \bullet \bullet \text{DCP}$ 13.3, 12.7
1:2	$[\text{Pd}_3^{2+}] \bullet \bullet \bullet \text{MCP}$ 37.4, 35.8	$[\text{Pd}_3^{2+}] \bullet \bullet \bullet \text{DCP}$ 33.3, 33.3 $([\text{Pd}_3^{2+}])_2 \bullet \bullet \bullet \text{DCP}$ 32.7, 32.0
1:4	$[\text{Pd}_3^{2+}] \bullet \bullet \bullet \text{MCP}$ 55.8, 54.2	$[\text{Pd}_3^{2+}] \bullet \bullet \bullet \text{DCP}$ 21.3, 28.0 $([\text{Pd}_3^{2+}])_2 \bullet \bullet \bullet \text{DCP}$ 72.0, 58.7

In order to verify whether the quenching is dominantly static, a mixed dynamic-static model was used,<sup>14</sup> which is derived from a sphere of action quenching model:  $[1 - (\Phi_F - \Phi_F^0)]/[\text{Pd}_3^{2+}] = K_D \cdot (\Phi_F/\Phi_F^0) + (1 - W)/[\text{Pd}_3^{2+}]$  where  $W$  is the fraction of the excited-state quenching from a collisional process given by  $\exp(-V \cdot [\text{Pd}_3^{2+}])$  ( $V$  is the static quenching constant representing the volume of the sphere of action) and  $K_D$  is the dynamic quenching constant. From a plot  $[1 - (\Phi_F - \Phi_F^0)]/[\text{Pd}_3^{2+}]$  vs  $(\Phi_F/\Phi_F^0)$ ,  $K_D$  is extracted from the slope using a least-square fit (ESI) and the intercept leads to the  $W$  values as a function of  $[\text{Pd}_3^{2+}]$ . Then  $V$  is evaluated from the slopes in the  $\ln(W)$  vs  $[\text{Pd}_3^{2+}]$  plots (ESI). These values are  $K_D = 241$  and  $V = 17900$  for  $[\text{Pd}_3^{2+}] \bullet \bullet \bullet \text{MCP}$  and  $K_D = 2750$  [ $\text{M}^{-1}$ ] and  $V = 21300$  [ $\text{M}^{-1}$ ] for  $[\text{Pd}_3^{2+}] \bullet \bullet \bullet \text{DCP} \bullet \bullet \bullet [\text{Pd}_3^{2+}]$ . Because the  $V$  values are much larger than those for  $K_D$ , the static quenching dominates the overall mechanism, and unsurprisingly,  $K_{11}$  and  $K_{12}$  are similar to  $V$ . In order to confirm this conclusion, this same analysis was performed on the assemblies' fluorescence at 77 K (*i.e.* only static quenching is possible). Indeed, the same relationship (*i.e.*  $K_D < V$ ) is observed (Table 1, and all graphs are placed in the ESI). The  $K_{11}$ ,  $K_{12}$  and  $K_D$  constants are then employed to evaluate of the proportion of molecular complexes at various  $\text{CO}_2^-/[\text{Pd}_3^{2+}]$  ratios (Table 2), which is useful for the interpretation of the fs transient absorption spectra (TAS) below.

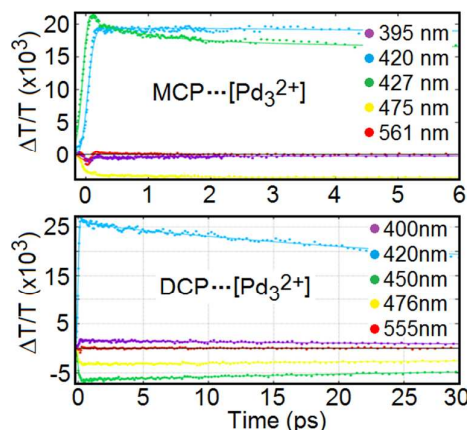
Both the **MCP** and **DCP** dyes exhibit irreversible oxidation and reduction waves in the cyclic voltammograms, CV (ESI). **MCP** is slightly harder to oxidize and easier to reduce compared to **DCP** (Table 3) which is consistent with their number of negative charges. Concurrently,  $[\text{Pd}_3^{2+}]$  exhibits oxidation and reduction potentials at  $E^{0/+1} = +0.95$  V<sup>15</sup> and  $E^{0/-1} = -0.50$  V vs SCE.<sup>9</sup>



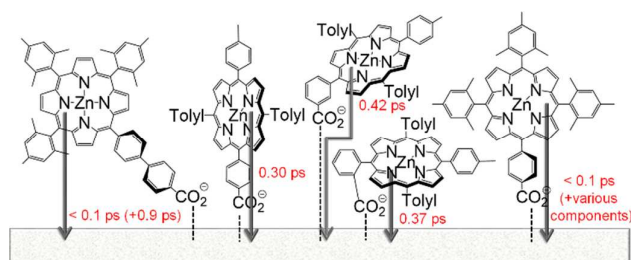
**Figure 4.** Individual components extracted from fs TAS ( $\lambda_{\text{exc}} = 600$  nm, MeOH, 298 K) with  $\text{CO}_2^-/[\text{Pd}_3^{2+}]$  ratios for MCP of 1:0 (a), 1:1 (c) 1:2 (e), and 1:4 (g) and for DCP of 1:0 (d), 1:1 (f) and 1:2 (h), and  $[\text{Pd}_3^{2+}]$  (b). The transient and bleach signals are respectively below and above the zero line.

Using the 0-0 peak position of the Q band (600 nm; 2.07 eV), the  $S_1$  driving forces for the oxidative (+1.57 V) and reductive quenching (-1.12 V vs SCE) can be extracted as well as  $\Delta G_s$  for  $\text{donor}^* \cdots [\text{Pd}_3^{2+}] \rightarrow \text{donor}^{\bullet+} \cdots [\text{Pd}_3^{\bullet+}]$  (+0.67; MCP, +0.72 V vs SCE; DCP), and for  $\text{acceptor}^* \cdots [\text{Pd}_3^{2+}] \rightarrow \text{acceptor}^{\bullet-} \cdots [\text{Pd}_3^{\bullet-}]$  (+0.38; MCP, +0.32 V vs SCE; DCP). Thermodynamically, the oxidative quenching of the dyes is more favorable than the reductive one. In order to corroborate this conclusion, DFT computations were performed in order to obtain their MO representations and energies of the frontier MOs (Figure 3 and ESI). For instance, the  $[\text{Pd}_3^{2+}] \cdots \text{MCP}$  assembly exhibits a HOMO and LUMO+1 that are typical  $\pi$ -MO for a chromophore. Concurrently, the HOMO-2 and LUMO are respectively the first filled and empty d orbitals of the  $\text{Pd}_3$  frame. Importantly, these levels are clearly energetically well isolated from one another ( $> \sim 0.4$  eV) and so the oxidation and reduction processes can only take place in the porphyrin and cluster units, respectively.

As the assemblies are non-emissive due to the oxidative quenching of the excited dyes by  $[\text{Pd}_3^{2+}]$ , TAS was used to extract  $k_{\text{et}}$ . First, these units are investigated separately to obtain their spectral and kinetic signatures (Figure 4a, b, d). The MCP and DCP TAS are characterized by the appearance of a broad transient signal (430–530 nm) and a strong bleach of the Soret band (425). The deconvolution of the data allows for the extraction of the various components and lifetimes. The strong and long ns-components (440 and 70 ns, respectively) are triplet species ( $T_1$ ) but their lifetimes are simply considered inaccurate due to the delay line of only 3.3 ns. Indeed, the previously reported  $T_1$  value for MCP in 2MeTHF is 137  $\mu\text{s}$  (without  $\text{O}_2$ ).<sup>16</sup> For the singlet species,  $S_1$ , the signals are weak. The  $\sim 2.1$  ns component in MCP is clearly the  $S_1$  emissive species ( $\tau_f = 1.97 \pm 0.08$  ns), but it is not observed for DCP ( $\tau_f = 2.04 \pm 0.10$  ns; note that a signal with a lifetime of  $\sim 1.5$  ns is detected in Figure 4f,h and is likely to be this species). This absence is due to the weakness of this signal compared to the dominant triplet species.



**Figure 5.** Monitoring of the TAS of MCP and DCP in the presence of  $[\text{Pd}_3^{2+}]$  in MeOH using  $\text{CO}_2^-/[\text{Pd}_3^{2+}]$  ratios of 1:4 and 1:2, respectively. The pulse width of the laser pulse is 85 fs and no rise time is observed.



**Figure 6.** Some examples of time scales for electron injection of various mono-carboxylates of zinc(II)porphyrin into  $\text{TiO}_2$  nanoparticles.<sup>7a,c</sup>

The ps-components are the known non-emissive species  $\text{Zn}(\text{porphyrin}) \cdots \text{solvent}$ .<sup>17</sup> No exhaustive study was performed on these species since this was not the aim of this investigation. Concurrently, the  $[\text{Pd}_3^{2+}]$  TAS exhibit two components ( $\sim 2.9$  and 296 ps) with practically identical signatures (Figure 4b). These are tentatively assigned to  $S_1$  and  $T_1$  species. The exact knowledge of these species is a minor issue in this work since these signals are not observed in the TAS of the assemblies (Figure 4). This situation is likely due to the weakness of the signal relative to those of the dyes. The TAS vs the  $\text{CO}_2^-/[\text{Pd}_3^{2+}]$  ratios allow for the reliable detection of the signals associated with the assemblies  $[\text{Pd}_3^{\bullet+}] \cdots \text{MCP}^{\bullet+}$  and  $[\text{Pd}_3^{\bullet+}] \cdots \text{DCP}^{\bullet+} \cdots [\text{Pd}_3^{2+}]$ . Upon increasing these ratios new signals with different shapes and lifetimes appear and are readily assigned to these charge separated species. Noteworthy, if the quenching was due to  $S_1$  energy transfer, the band shape of the transient signals would be constant, which is clearly not the case. For MCP, this distinctive signal is the one with lifetimes respectively of 160 and 65.5 ps (Figure 4e,g). Because its signal is expectedly stronger for the 1:4 ratio (Figure 4g), both its band shape and lifetime are more reliable (and used for the extraction of  $k_{\text{et}}$ ). For DCP, this distinctive signal is the one with a lifetime of 36.3 ps (Figure 4h).

Normally  $k_{\text{et}}$  is reliably obtained from the rise times monitored at wavelengths where their signals are stronger. However, no rise time longer than the pulse width was detected (for all monitoring wavelengths in fact) indicating that all species are formed during the laser pulse ( $< 85$  fs; Figure 5), including the charge separated  $[\text{Pd}_3^{\bullet+}] \cdots \text{MCP}^{\bullet+}$  and  $[\text{Pd}_3^{\bullet+}] \cdots \text{DCP}^{\bullet+} \cdots [\text{Pd}_3^{2+}]$  species. Therefore  $k_{\text{et}}$  must be  $> 1.2 \times 10^{13} \text{ s}^{-1}$ , which is clearly an ultrafast event.<sup>18</sup> This

time scale also compares favorably to what is commonly encountered for other zinc(II)porphyrins linked to the surface of TiO<sub>2</sub> nanoparticles (Figure 6; this list is far from exhaustive but is representative). However, this value can also be considered to be on the fast side. Similarly, the charge recombination characterized by TAS decays here found in the order of ~36 and ~66 ps, which also fall in the range of what it is currently reported.<sup>7</sup>

The ultrafast photo-induced electron transfers ( $k_{\text{et}} > 1.2 \times 10^{13} \text{ s}^{-1}$ ) occurring at the S<sub>1</sub> levels of the dyes in the structurally well-defined "straight up" ionic assemblies clearly indicate that it is not necessary to have a strong bond between the donor and acceptor and have a bent geometry.

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