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Magnetic anisotropy of iron-based metallic glassy fibers

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Flexible, uniform and surface smooth iron-based metallic glassy fibers (MGFs) were fabricated by a melt-extraction method. There was large magnetic anisotropy in the MGF. Our experimental results indicated that this large magnetic anisotropy was related to the unique magnetic domain structures in the MGF.

Iron-based and cobalt-based metallic glassy fibers (MGFs) are of commercial interest¹⁻⁴ due to their intriguing magnetic properties including giant magneto-impedance effect^{5,6}, giant stress-impedance^{7,8}, large Barkhausen effect⁹ and Matteucci effect¹⁰, which have been employed in some micro-electromechanical systems (MEMS) applications. However, various issues regarding applications of MGFs as MEMS are still unclear and have yet to be solved, such as the magnetic domain microstructures and the magnetic performance along different magnetization directions. Systematic studies of the magnetic domain structures and the intrinsic magnetic properties for MGFs would definitely shed new insights into their magnetization mechanisms. In this letter, we investigated the widely commercially used iron-silicon-based magnetic MGFs³⁻⁶, such as Fe-Si-B, to be doped with Cr to protect oxidation. The Fe74.5Si10B12Cr3.5 was demonstrated to show remarkable magnetic properties by a magnetometer and magnetic force microscopy (MFM). Interestingly, it is also found that the MGF is structural isotropy but magnetic anisotropy, which was closely correlated to the unique structure of the MGF.

Alloy ingots with a nominal composition of $Fe_{74.5}Si_{10}B_{12}Cr_{3.5}$ were prepared by arc-melting a mixture of elements with a purity level of >99.9% in a Ti-gettered high purity argon atmosphere. Minor addition of Cr is to improve the glass

forming ability and corrosion properties of the alloy. The cylindrical rods with a diameter of 8 mm were prepared by copper-mold casting under argon atmosphere. Then the fibers were prepared by the melt-extraction method. The details of MGFs preparation can be referred to as Ref. 11 and the ESI. The surface morphologies and actual fiber diameters were examined by scanning electron microscope (SEM). Magnetic measurements were performed using a Quantum Design MPMS-XL5 SQUID magnetometer. Samples were fixed in the capsule. The capsule used in magnetic measurement is made of gelatin with the size of height: 7 mm; diameter: 5 mm. The diamagnetic contribution from the sample holder was measured from 2 to 300 K, with the value of 1.4×10^{-8} emu/Oe. The fiber with a length of 5 mm was fixed on the sample holder by low temperature kapton tape. The fiber was first put parallel to the applied magnetic field to take parallel measurement, and then put perpendicular to the applied magnetic field to take the perpendicular measurement. The parallel and perpendicular measurements were conducted separately. The magnetization versus field data were collected between -5 T and 5 T, with an average field-scanning rate of 0.19 mT/s. Surface magnetic force morphology and domain structures were investigated using magnetic force microscopy (MFM) (BRUKER NANOSCOPE V SYSTEM WITH MULTIMODE 8) at room temperature.

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Fig. 1 Outer appearance of the melt-extracted fibers of $Fe_{74.5}Si_{10}B_{12}Cr_{3.5}$ alloy. Image (b) shows the magnified morphology of (a) and (c) shows the typical morphology of the cross section shape of the fiber.

Fig. 1a shows the outer appearance of the melt-extracted $Fe_{74.5}Si_{10}B_{12}Cr_{3.5}$ MGFs. All fibers show continuous and high flexible with a diameter of about 50 µm. Excellent surface smoothness and high uniformity can been seen in the magnified image of the MGFs (Fig. 1b). Spectroscopy analysis shows that there is no oxide layer on the surfaces of the MGFs. Fig. 1c shows the typical example of the micrograph of the transverse cross section of the MGFs, and high quality fibers with uniform and perfect circular cross section can be observed.



Fig. 2 Magnetic hysteresis loops measured at a maximum applied field of 5 T for the $Fe_{74.5}Si_{10}B_{12}Cr_{3.5}$ MGFs from the longitudinal and horizontal direction (a) at 2 K and (b) at 300 K,

and temperature dependence of magnetization-M(T) curves, (c) from the horizontal direction and (d) from the longitudinal direction

Fig. 2 shows the comprehensive magnetic properties of the Fe74.5Si10B12Cr3.5 MGFs at longitudinal and horizontal directions. The results of measuring the longitudinal and horizontal hysteresis M-H loops at 2K and room temperature are shown in Fig. 2a and 2b, respectively. From Fig. 2a, it can been obviously seen that the saturation magnetizations M_s of the MGFs tested along the longitudinal and horizontal direction at 2K are nearly the same (Table I), but the applied magnetic field strength to reach the saturation magnetization point (H_s) from the longitudinal direction is 0.75 kOe, which is much smaller than 8.32 kOe from the horizontal direction. Even at room temperature the MGFs possess almost the same features as shown in Fig. 2b. The data are summarized in Table I. These results indicate that the MGF is structural isotropy but magnetic anisotropy, and its easy magnetization direction is from the longitudinal side. The low coercive force H_c of the MGF listed in Table I is considered to be related with the microstructures of the MGFs. As in metallic glassy state, the MGFs lack long range order (crystalline structure), but possess homogeneous microstructures without any grain boundaries or large chemical precipitates¹². As a result, the defects are very narrow for pinning the movement and rotation of the magnetic domain walls during technical magnetization.

In order to further verify this magnetic anisotropy fact, magnetization measurements as functions of temperature and orientation were performed. The temperature dependence of magnetization-M(T) curves were measured from 2 K to 300 K parallel and perpendicular to the applied magnetic field of 1 kOe as shown in Fig. 2c and 2d, respectively. With respect to curves M(T) measured from the horizontal and longitudinal directions, the zero-field-cooled (ZFC) curves coincide much well with the field-cooled (FC) curves, but the M(T) value tested along the longitudinal direction is 65-56 emu/g, which is much larger than that along the horizontal direction, about 10-13 emu/g. These data reveal that the MGF is magnetized much more easily from the longitudinal direction, further confirming the MGF is magnetic anisotropy. It is worth noting that the MGF possessed a weak broad peak at about 225 K, as shown in the M(T) curve in Fig. 2c, and this event verified there was antiferromagnetic order in the MGF. This broad peak would disappear when the MGF reaches the saturation magnetization. The M along the horizontal direction at 10 kOe and along the longitudinal direction at 1 kOe of the MGF continuously decreases with the temperature increasing, because it is mainly related to the thermal vibrations of the atoms. Thus, it is apparent that the steep response around the original point of Fig. 2a observed along the horizontal direction includes spin-flip transition behaviour (Fig. S2, ESI). Furthermore, the M is initially proportional to the applied magnetic field H in the paramagnetic phase at 300 K along the horizontal direction (Fig. 2b). Then it will become saturated at the saturation magnetization point H_s, forming the S-shaped curve. As there is super exchange interaction at 2 K, the value

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of $H_{\rm s}$ is a little bigger than that at 300 K along the horizontal direction (Table I)

Table I. The comparison of the magnetic properties of the $Fe_{74.5}Si_{10}B_{12}Cr_{3.5}$ MGF tested from the longitudinal direction and horizontal direction. Saturation magnetization M_s , Coercive force H_c , and applied magnetic field strength to reach the saturation magnetization point H_s .

	Longitudinal	Horizontal	Longitudinal	Horizontal
	at 2K	at 2K	at 300K	at 300K
M _s (emu/g)	68.2	66	57.33	54.6
H _c (A/m)	5	10	14	33
H _s (kOe)	0.75	8.32	0.4	7.5

Magnetic force microscopy (MFM) measurements were conducted to further identify the surface magnetic domain microstructures distribution. A typical selected area of the MGF with a scanning size of 10 \times 10 μ m² by MFM was fixed (marked by A), as indicated in Fig. 3a. Fig. 3b shows the surface height morphology of the A area, which further confirms that the melt-extracted MGFs have a smooth surface. In Fig. 3c, it can be observed that there is a distinct colour contrast in the micron-sized MGFs, distinguishing features of alternate "dark" and "bright" magnetic domain structures with a width of about 1-2 µm. This suggests a strong magnetic anisotropy with the magnetic easy axis normal to the plane of the sample. As the magnetization is along the easy axis and normal to the plane of the sample, "dark" and "bright" magnetic domain structures were formed so that the demagnetization field would be reduced. These "dark" and "bright" structures also demonstrate that the magnetization direction coexist in and out the surface of the fibers, indicating that radial magnetic domain structures were formed. Even though the MGFs are amorphous and thin, the MFM results reveal its complex domain structures in the surface.





Fig. 3 Surface MFM images of the Fe_{74.5}Si₁₀B₁₂Cr_{3.5} MGFs with a scanning size of 10 μ m × 10 μ m. (a) the selected area on the MGF denoted by A, (b) the 3D surface height image, and (c) the 2D surface magnetic force image

Based on the above experimental results, the domain structures distributed in the iron-based MGF can be proposed in Fig. 4. The MGF possessed a "core-shell" domain structure consisting of large internal domain with magnetization lying along the fiber axis and covered by the external domain structure with radial magnetized domains. The difference between the magnetic domains in the core and in the shell results in the magnetic anisotropy of the MGF. In fact, many work showed that the magnetic domain structures were related to the residual stress in the MGF^{5,13,14}. In previous studies^{13,15,16}, the MGF was glass-coated and the difference in the coefficients of thermal expansion of glass and metal would lead to high radial stresses, which were maximal at the metallic core surface and decreased sharply toward the wire axis. The presence of such a residual stress field formed a complex magnetic structure. However, in our study the MGF was exposed without any coverings, thus the resultant residual stress field was mainly induced during fabrication of amorphous microfibers. As the MGF was prepared by the meltextraction method, axial tension stresses and quenching stresses related to the temperature gradient occurring in the metal on quenching could be highly expected, and formed the anisotropic "core-shell" magnetic domain structures. On the other hand, the asymmetric "shell" structure of the MGF may generate the antiferromagnetic order due to differences of the steric and electronic structures on the surface and near the surface of "shell" area, which can lead unique physical properties around these domain region. These unique domain structures lead to its magnetic anisotropy, excellent physical properties and would bring about its application as a new engineering and functional material intended for parts of inductive components (e.g. high performance magnetic sensors, electromagnetic interference shielding, MEMS, and other applications).

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Fig. 4 Schematic domain structures of metallic glassy fibers

In summary, we reported a benign approach to detect the magnetic structures of the micron-sized Fe74.5Si10B12Cr3.5 MGFs and find that there is large magnetic anisotropy, which is related to the magnetic domain structures in the MGFs. The unique domain structures distribute in the surface of the MGFs are revealed by the inhomogeneous magnetic force morphology. This would be a paradigm shift in magnetization technology because the surface domain structure characterization is a key issue for designing and building reliable and high-performance magnetic sensors and devices. As-prepared micron-sized MGFs associated with excellent flexibility, easy handling, low coercive force and high saturation magnetization, will find broad industrial applications. In addition, this work could open new opportunities of using active multi-element metallic systems to functionalize more flexible micron-sized metallic fibers with better magnetic performance.

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