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## ARTICLE

## Time-resolved photoelectron imaging of the isolated deprotonated nucleotides

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Using time-resolved photoelectron spectroscopy, the excited state dynamics of gas-phase mass-selected nucleotide anions have been monitored following UV excitation at 4.66 eV. The spectra reveal that the dynamics of the 2'-deoxyguanosine 5'-monophosphate anion (dGMP<sup>-</sup>) are very similar to those of the adenosine nucleotide (dAMP<sup>-</sup>) and insensitive to a solvation environment. Comparison of our results with other literature suggests that nucleotides of the two purine bases share a common relaxation pathway, whereby the initially populated <sup>1</sup>ππ\* states relax to the ground electronic state without involvement of any other intermediary electronic states. In the analogous pyrimidine nucleotides of thymine and cytosine, dTMP<sup>-</sup> and dCMP<sup>-</sup>, no such unified mechanism is observed. Photoexcited dTMP<sup>-</sup> behaves much like the isolated nucleobase thymine, exhibiting rapid relaxation to the ground electronic state, although with a minor long-lived channel. On the other hand, isolated dCMP<sup>-</sup> is longer lived than its cytosine nucleobase, and hence it appears that the presence of the sugar and phosphate in the nucleotide arrangement leads to a modification of the available relaxation pathways. Nucleotides are the basic monomer building blocks of DNA and our results present important new benchmark data to develop an understanding of the molecular mechanism by which photodamage can be mediated when DNA is exposed to UV light.

### Introduction

The photophysics of DNA are a subject of great interest, as ultraviolet (UV) light is capable of inducing harmful strand breaks and mutations.<sup>1-3</sup> Despite the constant irradiation from solar light, photoinduced mutations are rare, and thus it appears that nature has selected the structure of DNA to have in-built photoprotective mechanisms. Part of this photoprotection takes the form of ultrafast non-radiative relaxation processes which enable potentially damaging electronic energy to be rapidly dispersed as harmless thermal energy. A major goal in the photochemistry and photobiology communities has been to develop a molecular level understand of these relaxation mechanisms.

The UV absorption spectrum of DNA is dominated by the excitation of the four nucleobases: adenine (Ade), guanine (Gua), thymine (Thy) and cytosine (Cyt). These are aromatic species and possess a strong UV absorption around 4.7 eV (264 nm).<sup>4</sup> In all four nucleobases, this UV absorption populates a <sup>1</sup>ππ\* excited state and considerable efforts have been devoted to understanding the photophysics of the isolated (gas phase) nucleobases and related chromophores. All the nucleobases

possess decay mechanisms that lead to rapid relaxation to the ground electronic state. In the purine nucleobases (Ade and Gua), this relaxation occurs on a < 1 ps timescale, whilst in the pyrimidine nucleobases (Thy and Cyt) relaxation is somewhat slower, though still < 10 ps. A combination of experimental and theoretical work has provided great insights towards determining the plausible mechanisms by which excited state flux undergoes relaxation; in all cases the relaxation proceeds non-adiabatically by exploiting conical intersections (CIs) where an electronically excited state intersects the ground or an intermediate excited state.<sup>1-3</sup>

Although gas phase studies have been extremely important in deciphering the photophysics of isolated nucleobases,<sup>5-13</sup> this reductionist approach has been limited due to the involatility of larger DNA components. In biological systems, DNA is a polymer of nucleotides, each composed of a nucleobase bound to a deoxyribose molecule and a phosphate group (see Fig. 1 for a representative example). To progress, the major question is: to what extent are the relaxation mechanisms in isolated nucleobases still active in their nucleotides, oligonucleotides and ultimately DNA? Solution phase studies have gone a long way to answering these questions, as transient absorption and

fluorescence upconversion experiments can easily be compared between single nucleobases and nucleotides or oligonucleotides.<sup>14-20</sup> These results have shown that, to a first approximation, the relaxation mechanisms of solvated nucleobases and nucleobases within nucleotides are very similar. On the other hand, interactions between stacked nucleobases within a DNA strand can lead to the formation of excimer states, which may result in remarkably different relaxation dynamics.<sup>21-25</sup> However, even in these large systems, internal conversion to the ground state remains a dominant decay route. In general, mechanistic assignment of solution phase data is complicated by selection rules and solvent broadening, and additionally it is a much more challenging environment for computational chemistry to simulate than the gas phase. In the gas phase, these solvent-induced complications are not present, but to date, have been limited in size to the study of Cyt, Ade and Thy.

We have recently demonstrated how a combination of electrospray ionization (ESI) and time resolved photoelectron spectroscopy (TRPES) could be used to build up complexity and size from an isolated nucleobase towards DNA. We demonstrated this using the (deprotonated) nucleotide of adenine, 2'-deoxyadenosine 5'-monophosphate (dAMP<sup>-</sup>, see Fig. 1).<sup>26</sup> Using ESI, transportation of deprotonated nucleotides to the gas phase is facile, and TRPES provides rich details of the path electronically excited state population takes as the nucleotides decay back to the ground electronic state. Deprotonation occurs on the phosphate group of the nucleotides; hence, the excess charge also resides at this moiety.<sup>27</sup> In the dAMP<sup>-</sup> experiments, and those presented herein, a femtosecond UV pulse at 4.66 eV (266 nm) prepares an electronically excited state on the mass selected nucleotide anions, and a second, time delayed, femtosecond pulse at 3.10 eV (400 nm) detaches an electron. Fig. 1 outlines the electronic states involved in these experiments. The kinetic energy of the electron,  $eKE$ , is analysed using velocity map imaging (VMI). A key feature of this experimental scheme is that, although in principle the total photon energy (7.76 eV) is sufficient to detach the electron from many parts of the anion, including the charge carrying phosphate group, by exploiting the  ${}^1\pi\pi^* \leftarrow S_0$  resonance, we can selectively detach electrons only from the nucleobase moiety.<sup>28</sup> This allows us to think of the experiment as photoelectron spectroscopy on a neutral subgroup, whilst the charged phosphate group remains a spectator.

In our previous dAMP<sup>-</sup> experiments, we found that the excited state dynamics of the nucleotide could be modelled using biexponential kinetics. Using a global fitting technique, the energetic components of these two pathways could be deconvoluted, and there was clear evidence of a rapid (< 60 fs) initial relaxation within the initially populated  ${}^1\pi\pi^*$  state, followed by a slower (290 fs) relaxation to the ground electronic state. These results compared extremely well to similar experiments performed on isolated 9-methyladenine (Ade-9Me; methylation occurs on the N atom where the deoxyribose group would attach).<sup>5, 6, 29</sup> Theoretical studies have posited two plausible relaxation mechanisms for dAMP<sup>-</sup> and

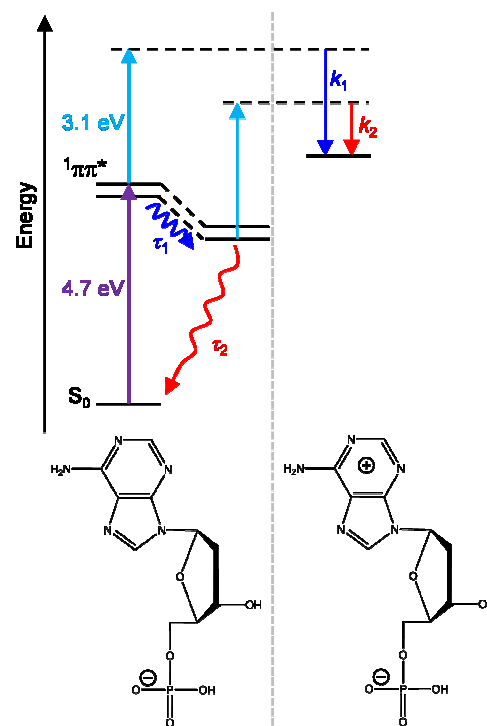


Fig. 1 The pump-probe scheme employed in the experiments, along with a schematic for the relaxation process observed in dAMP<sup>-</sup>. A 4.66 eV photon prepares the molecule in the  ${}^1\pi\pi^*$  excited state, which subsequently relaxes on this excited state before converting back to the ground state ( $S_0$ ). Below, the geometry of dAMP<sup>-</sup> and the subsequent neutral are shown.

Ade-9Me following excitation to the optically brightest  ${}^1\pi\pi^*$  state (the  ${}^1L_a$  state).<sup>30-41</sup> Firstly, a two-state mechanism may be active whereby the  ${}^1\pi\pi^*$  state initially converts to a lower lying  ${}^1n\pi^*$  state, and this then relaxes to the ground state through a CI (termed  ${}^1S_6$  in the Boeyens classification scheme).<sup>42</sup> Alternatively, the relaxation may occur solely on the  ${}^1\pi\pi^*$  state, where first rapid wavepacket motion occurs on this state, followed by internal conversion to the ground state through a CI (termed  ${}^2E$ ). Through comparisons with both Ade-9Me and solution phase dAMP<sup>-</sup> data, we concluded that the single-state mechanism was more likely to be operative in gas phase dAMP<sup>-</sup>.<sup>26</sup> This mechanism is sketched in Fig. 1.

Here, we apply this experimental technique to study the remaining deprotonated deoxynucleotides, 2'-deoxyguanosine 5'-monophosphate (dGMP<sup>-</sup>), 2'-deoxythymidine 5'-monophosphate (dTMP<sup>-</sup>), and 2'-deoxycytidine 5'-monophosphate (dCMP<sup>-</sup>). The purine base within dGMP<sup>-</sup> is especially interesting as no TRPES has ever been performed on isolated Gua, due to its low volatility,<sup>43</sup> so it is much less well understood than Ade derivatives. dTMP<sup>-</sup> and dCMP<sup>-</sup> contain the pyrimidine nucleobases, which follow different mechanisms from the purine bases, but appear to follow the same general theme. We show that dGMP<sup>-</sup> behaves in a similar manner to dAMP<sup>-</sup>, and likely shares a very similar excited state relaxation pathway. dTMP<sup>-</sup> shows similar dynamics to isolated thymine, and hence very likely follows a similar relaxation route. Finally dCMP<sup>-</sup> is shown to behave rather differently from isolated Cyt,

suggesting that the addition of the sugar/phosphate backbone significantly alters the dynamics of this nucleobase.

## Experimental

The experiment has previously been described in detail,<sup>44-46</sup> and only a brief overview shall be presented here. Isolated nucleotides were produced using ESI at  $-2.5$  kV with  $\sim 1$  mM solutions of the sodium salts of the nucleotides (Sigma-Aldrich, >99% purity) in methanol. Anions were accumulated in a ring electrode trap, where they are thermalized to  $\sim 300$  K, and then ejected into a collinear time-of-flight mass spectrometer at a 50 Hz repetition rate. Mass selected ions were intersected by femtosecond pump (4.66 eV, 266 nm) and probe (3.10 eV, 400 nm) laser pulses in a perpendicular velocity map imaging arrangement.<sup>47</sup> Detached electrons were imaged using a position sensitive detector. Around  $5 \times 10^4$  laser shots were accumulated for each pump-probe time delay ( $t$ ). The 4.66 eV light creates additional background counts, so for each delay an equal number of photoelectron images were accumulated without ions, so that this background could be subtracted. Likewise, to remove one-colour signal produced from the pump pulse, the photoelectron image at a delay of  $-1$  ps (probe pulse precedes pump pulse) was subtracted from all images. Images were deconvoluted using the polar onion peeling algorithm<sup>48</sup> to derive the photoelectron spectra and photoelectron angular distributions. The latter reveal no time-resolved dynamics and will not be discussed further. The spectrometer was calibrated using the known spectrum of iodide, and has a resolution  $\Delta eKE/eKE \approx 5\%$ .

Femtosecond laser pulses were derived from a commercial Ti:Sapphire oscillator and regenerative amplifier, delivering fundamental pulses centred at 1.55 eV (800 nm) and with  $\sim 50$  fs duration. The 4.66 eV pump pulse was produced using a combination of second harmonic and sum frequency generation with the fundamental, in a series of two type I  $\beta$ -barium borate (BBO) crystals. The 3.10 eV probe pulse was produced by second harmonic generation in a further type I BBO crystal. The two pulses were delayed relative to each other using a motorized delay line. Both pulses were combined collinearly using a dichroic mirror, and were loosely focused into the interaction region using a curved metal mirror. The intensity of both beams were below  $10^{11}$  W cm<sup>-2</sup>. The cross-correlation of the two beams was estimated by fitting the rising edge of the data to be approximately 120 fs, providing a temporal resolution of  $\sim 60$  fs.

## Results and Analysis

Figures 2 to 4 show the results and analysis following excitation at 4.66 eV (266 nm) and probing at 3.10 eV (400 nm) for dGMP<sup>-</sup>, dTMP<sup>-</sup>, and dCMP<sup>-</sup>, respectively. Their structures are shown in (a) of each figure. The time-resolved photoelectron spectra shown in (b) of these figures have a similar overall appearance with ultrafast decay that is mostly complete by 2 ps. Additionally, a component at high  $eKE$

(between 1 and 2 eV) can be distinguished from a peak at lower  $eKE$  which appears to decay on a longer time-scale for the three systems. The timescales and detailed shapes of the features are however different.

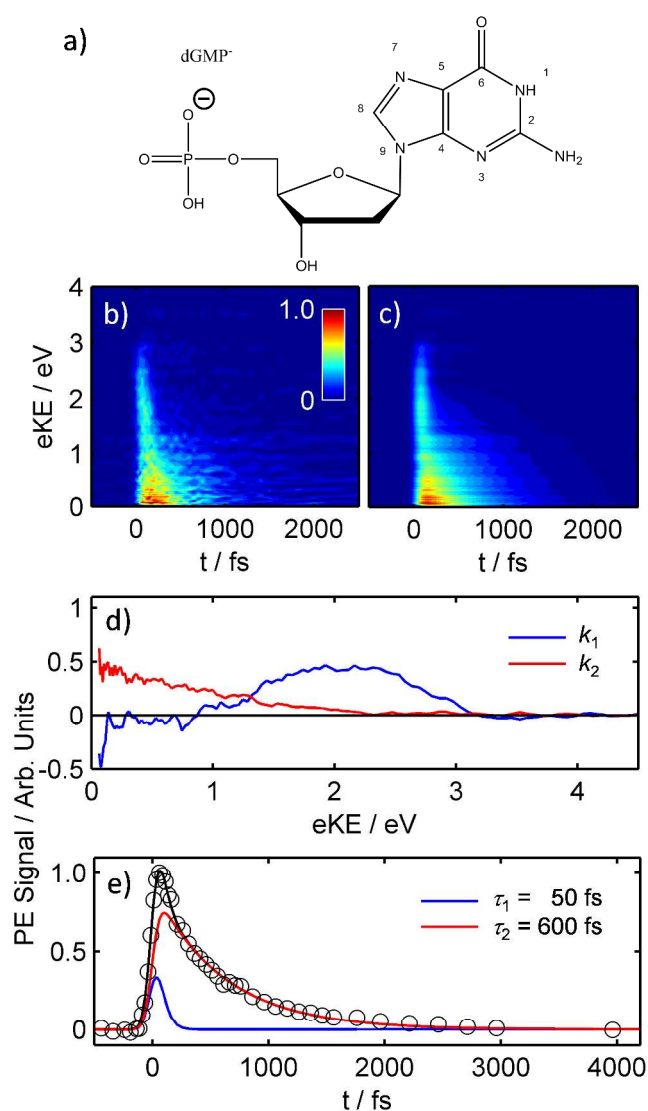
All three nucleotides were analysed using a global fitting routine, in an identical manner to previous TRPES studies on nucleobases,<sup>5</sup> and our recent work on the dAMP<sup>-</sup> nucleotide.<sup>26</sup> Global fitting operates by fitting spectra at all time delays simultaneously to the same set of time constants for a sum of  $i$  exponential decays:

$$S(eKE, t) = \sum_i G(t) * k_i(eKE) e^{-\frac{t-t_0}{\tau_i}}$$

where  $k_i(eKE)$  is the decay associated spectrum (DAS) corresponding to the exponential decay with a time constant of  $\tau_i$ .  $t_0$  gives the temporal position of overlap between pump and probe pulses, and  $G(t)$  is the Gaussian instrument response function (120 fs). The Levenberg-Marquardt algorithm (implemented in the Matlab optimization toolbox) is used to minimize the sum of the squares of the difference between  $S(eKE, t)$  and the experimental dataset.

The shapes of the DAS are highly informative in guiding the interpretation of the spectra. In particular, any negative component corresponds to an exponential rise in that energy region, concomitant with a decay in the positive regions. This can be interpreted as a flow of population from the positive region into the negative region, which may be observed due to either a change in electronic state, or relaxation within a single electronically excited state.<sup>49</sup> It is important to note that global fitting *cannot* precisely capture large amplitude motions in the time resolved photoelectron spectrum; if motion of population occurs which does not behave as one state decaying into another, then the fitting algorithm is only able to approximate this motion.<sup>50</sup> In instances where the motion is relatively fast compared to the instrument response function, however, the global fitting approximation is valid.

In all the photoelectron spectra and DAS presented, there is a small amplitude of signal with an  $eKE$  that extends beyond the maximum energy expected from just the pump (4.66 eV) and probe (3.10 eV) photons. This signal can be attributed to two-photon probing of the system, giving a combined probe energy of 6.20 eV. The two-photon spectra resemble the one-photon spectra observed at lower  $eKE$ , but blue shifted by 3.10 eV. In principle, the presence of the two-photon peak extends the probing range significantly, although its weak intensity makes concrete assignments difficult. The ability to significantly extend the probe window is noteworthy in that it allows the adiabatic onset of photoelectron signal to be observed, and hence the Franck-Condon detection window. In the DAS, the same dynamics are assigned to the two-photon probe signal as the one-photon probe signal, indicating that we are capturing all the dynamics of the system, and there are no further processes lying energetically below the one-photon probe limit.



**Fig. 2** a) Structure of  $\text{dGMP}^-$ ; b) TRPES data from  $\text{dGMP}^-$ , presented as a false-colour plot; c) global fit to the TRPES data, using two exponential decay functions; d) decay associated spectra showing the spectra of the two fit components used to produce the fit; and e) integrated photoelectron signal as a function of delay time. Open circles give experimental data, coloured lines give the two exponential decay functions, and black line gives the total fitted integral.

### $\text{dGMP}^-$

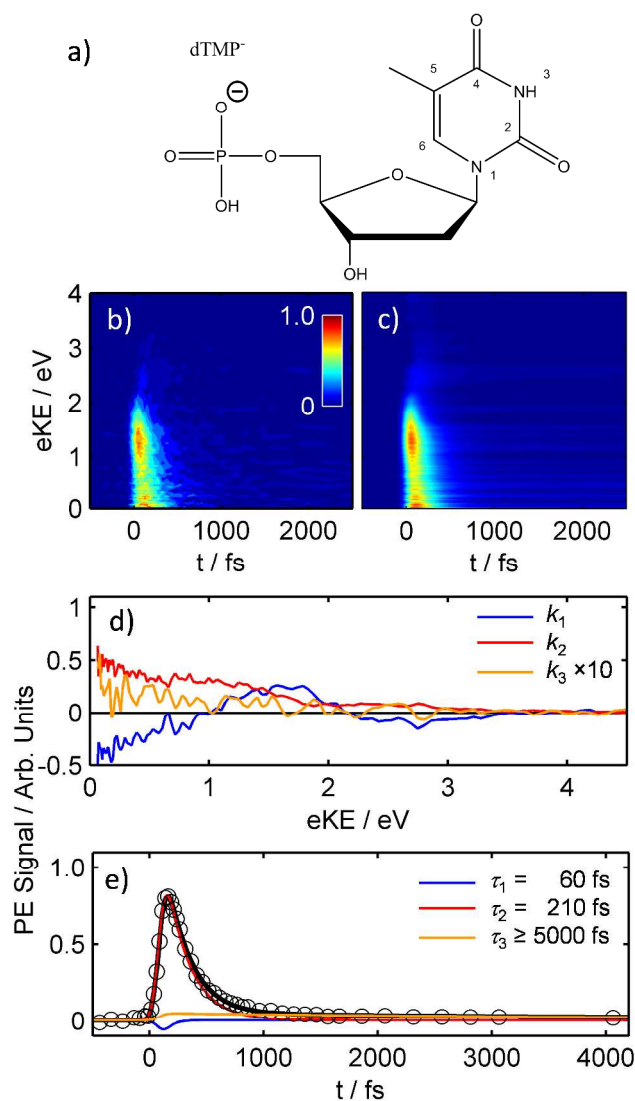
Background subtracted time resolved photoelectron spectra of  $\text{dGMP}^-$  are presented as an intensity plot in Fig. 2b. Immediately following excitation, a feature is seen extending to  $\sim 3$  eV, which rapidly decays in  $< 200$  fs. At lower  $eKE$ , a second feature is observed, which has a slightly delayed onset compared to the higher  $eKE$  feature (see Fig. S3 in the supporting information). This lower energy feature decays on a much slower timescale, taking  $\sim 1000$  fs to substantially decay. The total integrated electron signal is shown as a function of delay time in Fig. 2e; the biexponential decay kinetics are readily visible.

The global fitting algorithm disentangles the spectra of the two decay components, and provides quantitative lifetimes. Fig. 2c shows the intensity plot of a global fit to the  $\text{dGMP}^-$  data, using two decay components. The residuals (see Fig. S2 in the supporting information) indicate that an acceptable fit has been found. The lifetimes extracted for the two decay components are  $\tau_1 = 50 \pm 60$  fs, and  $\tau_2 = 600 \pm 170$  fs; corresponding exponential decay functions are shown as coloured lines in Fig. 2e. Fig. 2d, show the spectral profile of the two DAS,  $k_1$  and  $k_2$ :  $k_1$  is a broad peak from  $\sim 1$  eV  $< eKE < 3.1$  eV, and is predominately negative for  $< 1$  eV; while  $k_2$  takes the form of a continually decreasing slope from  $0$  eV  $< eKE < 2.2$  eV. The  $k_1$  high energy cutoff of 3.1 eV agrees well with the adiabatic detachment energy for the nucleobase within  $\text{dGMP}^-$ , which was found to be 4.65 eV.<sup>28</sup> The negative intensity of  $k_1 < 1$  eV is indicative of a population transfer from the higher  $eKE$  region to the lower  $eKE$  region. Note that the extracted  $\tau_1$  lifetime is within our temporal resolution.

### $\text{dTMP}^-$

Fig. 3b and c shows the background subtracted time-resolved photoelectron data and the global fit to this, respectively, for  $\text{dTMP}^-$ . Unlike the purine nucleotides, two decay components are not adequate to fit the data: close inspection of Fig. 3e shows a small offset from the baseline at long time delays. The global fit was thus performed with three decay components, with the lifetime of the third component arbitrarily set to  $\tau_3 = 5$  ps. The data are reminiscent of those for the purine nucleobases, with a  $\tau_2 = 211 \pm 127$  fs feature decreasing from  $0$   $eKE$ , but remaining universally positive, and a faster  $\tau_1 = 58 \pm 111$  fs feature that is positive around 1.5 eV, but negative for  $eKE < 1$  eV. The significantly larger errors extracted for  $\text{dTMP}^-$  compared to  $\text{dGMP}^-$  are a consequence of the additional timescale required in the fit, reducing the steepness of  $\chi^2$  in phase space. Inspection of Fig. 3b (and Fig. S3 in the supporting information) show the higher energy feature shifting to lower  $eKE$  as a function of time, a clear indication of sequential dynamics. The DAS for the picosecond component,  $k_3$ , is mainly located at low  $eKE$ . However, the feature is too weak (amplified 10-fold in Fig. 3d for clarity) to make any definitive comments about its shape. A strong two-photon feature is observed for  $eKE > 2.1$  eV, which is the adiabatic cut off for a one-photon probe. The spectra show a clear two-photon feature which effectively extends the available probing window, although we note that convolution of the one- and two-photon probe peaks may complicate the data analysis. The initial rapid shift in  $eKE$  is especially visible in the two-photon signal.

The DAS of the  $\tau_1$  fs component ( $k_1$ ) has considerable negative amplitude at  $eKE < 1$  eV, and in fact the total integral of it is negative (see Fig. 3e). The overall negative integral indicates that we are observing population flowing from a region that has a lower ionisation cross-section into a region where the cross-section has increased. The two-photon feature for  $k_1$  also shows a significant negative component. The large amount of negative signal in this feature points to a

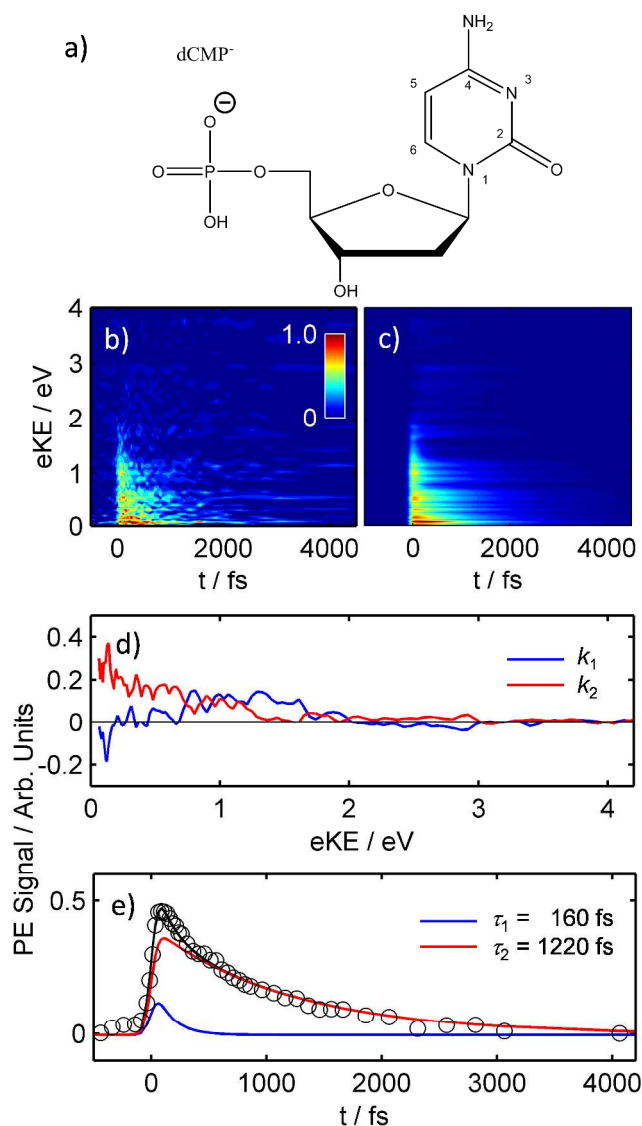


**Fig. 3** a) Structure of  $\text{dTMP}^-$ ; b) TRPES data from  $\text{dTMP}^-$ , presented as a false-colour plot; c) global fit to the TRPES data, using three exponential decay functions (one of which was arbitrarily set to 5 ps); d) decay associated spectra showing the spectra of the two fit components used to produce the fit. The long lived component,  $k_3$ , has been expanded ten-fold for clarity. e) integrated photoelectron signal as a function of delay time. Open circles give experimental data, coloured lines give the three exponential decay functions, and black line gives the total fitted integral.

considerable population flow in the system; in the 2D spectra, this flow is clearly manifested as a shift to later onset time for the  $eKE < 1$  eV feature and the two-photon feature, as can also be seen in Fig. S3 in the supporting information.

#### dCMP<sup>-</sup>

Fig. 4b and c shows the background subtracted time-resolved photoelectron data and the global fit to this, respectively, for dCMP<sup>-</sup>. As indicated by the larger residuals (see Fig. S2 in the supporting information), the signal to noise ratio for dCMP<sup>-</sup> is worse than for the other three nucleotides. This is partly because the ion signal of dCMP<sup>-</sup> was lower, presumably



**Fig. 4** a) Structure of  $\text{dCMP}^-$ ; b) TRPES data from  $\text{dCMP}^-$ , presented as a false-colour plot; c) global fit to the TRPES data, using two exponential decay functions; d) decay associated spectra showing the spectra of the two fit components used to produce the fit; and e) integrated photoelectron signal as a function of delay time. Open circles give experimental data, coloured lines give the two exponential decay functions, and black line gives the total fitted integral.

because it does not electrospray as easily, and partly due to reduced absorption and ionization cross sections. As with the purine nucleobases, a single decay component is insufficient to capture the dynamics, and so a biexponential model has been used. Unlike  $\text{dTMP}^-$ , there is no long-time offset, although the relatively poorer signal-to-noise may obscure its presence. The two decay components have lifetimes of  $\tau_1 = 160$  fs and  $\tau_2 = 1.22$  ps, and the DAS spectral profile resembles the other nucleotides, with the fast component found at higher  $eKE$ , and the slower decay increasing towards  $eKE = 0$ , although no evidence of sequential dynamics (negative DAS) is observed. We have not assigned confidence intervals to the lifetimes extracted for dCMP<sup>-</sup>, because the support plane analysis

revealed that the  $\chi^2$  landscape of the fit is exceedingly shallow (see Fig. S1 in the supporting information) and so we cannot make any precise judgements of the errors. The poorer quality of these data mean they are somewhat less reliable than for the other nucleotides, and so conclusions drawn from them are by necessity more tenuous. We can, however, definitively state that the excited state population in dCMP<sup>-</sup> is considerably longer lived than in the other nucleotides, but also that it shares similar DAS.

## Discussion

### Dynamics of dGMP<sup>-</sup> excited at 4.66 eV

Studying the dynamics of the isolated purine nucleobase Gua have presented a particular challenge, given that Gua tends to undergo thermal decomposition upon heating in a molecular beam apparatus.<sup>43, 51</sup> Additionally, Gua can form multiple stable conformers in the gas phase,<sup>52-54</sup> that further complicates experiments. Time resolved ion yield (TR-IY) experiments with mass selective product detection (which are sensitive only to the product of interest) are the only reported excited state dynamics experiments on isolated Gua.<sup>9, 10</sup> Photoexcitation at 4.7 eV, resulted in biexponential kinetics, with lifetimes of 148 fs and 360 fs for the two components.<sup>9</sup>

Gua has, however, been the focus of several *ab initio* calculations. These reveal that, at the ground state equilibrium geometry, the S<sub>1</sub> state is the <sup>1</sup>L<sub>a</sub> (<sup>1</sup>ππ\*) state. At higher energy lie the <sup>1</sup>nπ\* and <sup>1</sup>L<sub>b</sub> (<sup>1</sup>ππ\*) states, which are close in energy.<sup>55-63</sup> Calculations of CIs reveal a similar pattern to Ade: three CIs have been located which may be active in the relaxation of Gua.<sup>3, 57, 59, 60, 62, 64</sup> The first two have a geometry puckered at the C<sup>2</sup> position (see Fig. 2a for atom label scheme), analogous to the <sup>2</sup>E CI in Ade. Both of these CIs, which differ only in the geometry of the amino moiety, are between the <sup>1</sup>L<sub>a</sub> and S<sub>0</sub> states. The third CI has a geometry with a C<sup>6</sup> pucker, again like in Ade, and this CI is associated with a state of mixed <sup>1</sup>nπ\* and <sup>1</sup>ππ\* character. There are also <sup>1</sup>πσ\* states associated with the N<sup>1</sup>-H and N<sup>(amino)</sup>-H bonds, which possess CIs that intersect the ground state,<sup>60</sup> although most recent theory does not consider these states. Minimum energy path calculations suggested that the rapid decay observed was attributed to direct decay from the <sup>1</sup>L<sub>a</sub> state through the <sup>2</sup>E CIs,<sup>59</sup> and the longer lived component was attributed to dynamics via the <sup>1</sup>nπ\*/<sup>1</sup>L<sub>b</sub> mixed state. More recently, nonadiabatic dynamics simulations have been performed which indicate that dynamics occur almost exclusively on the first <sup>1</sup>ππ\* state, with almost all population relaxing through the two <sup>2</sup>E CIs,<sup>62</sup> although an earlier simulation suggested that the C<sup>6</sup> pucker CI may also be involved.<sup>65</sup>

In the solution phase, Gua itself is rarely studied as it has poor solubility, although dGMP<sup>-</sup> has been studied in a number of experiments with transient absorption<sup>1, 16, 66</sup> and fluorescence upconversion.<sup>14, 16, 19, 67 14, 16, 19, 64</sup> As for gas phase Gua excited at around 4.7 eV, solvated dGMP<sup>-</sup> is observed to relax through biexponential kinetics.<sup>1</sup> Recent measurements give the short

lifetime,  $\tau_1 = 120$  fs and the longer lifetime,  $\tau_2 = 680$  fs.<sup>16</sup> Computational studies show that solvation dramatically blue-shifts the first <sup>1</sup>nπ\* state, to the extent where it is considered unlikely to have any substantial influence on the dynamics.<sup>65, 68</sup> Instead, *ab initio* results have proposed that all dynamics occur upon the <sup>1</sup>L<sub>a</sub> state, and any population on the <sup>1</sup>L<sub>b</sub> state converts to the <sup>1</sup>L<sub>a</sub> in < 20 fs.<sup>68</sup> The dynamics then proceed purely through CIs between the <sup>1</sup>L<sub>a</sub> and S<sub>0</sub> states, although there are indications that the presence of the solvent influences which path – and ultimately which CI – is accessed to relax back to the ground electronic state.

Comparing the lifetimes obtained in our experiments on isolated dGMP<sup>-</sup> (50 fs and 600 fs) with those for solvated dGMP<sup>-</sup> (120 fs and 680 fs) reveals a striking similarity: the two agree within error despite the radically different environments.<sup>16</sup> This strongly suggests that a similar decay pathway is occurring in both environments, despite the strong effect of the solvent on the energy of the <sup>1</sup>nπ\* state.<sup>65, 68</sup> Analogous to dAMP<sup>-</sup>,<sup>26</sup> our data suggests that the dynamics in dGMP<sup>-</sup> are occurring predominately on the <sup>1</sup>L<sub>a</sub> excited state, with the <sup>1</sup>nπ\* state not actively participating. This is in agreement with some theoretical predictions.<sup>57, 61, 62</sup> We therefore assign the initial motion observed in the time resolved photoelectron spectrum to relaxation of the excited system away from the Franck-Condon geometry, and the slower decay to internal conversion using a <sup>2</sup>E CI. If this assignment is correct, then by extension we propose that the previous biexponential kinetics observed in isolated Gua also correspond to dynamics solely upon a single excited state,<sup>9</sup> and excited state dynamics of the purine nucleobases may follow a similar mechanism. The DAS can then be correlated to the photodetachment spectra of different regions along the <sup>1</sup>L<sub>a</sub> decay coordinate.

The similarity between the dGMP<sup>-</sup> data presented here and those obtained previously for dAMP<sup>-</sup> is striking.<sup>26</sup> The shapes of the DAS for both the fast and slow components are extremely similar, although the fast feature in dAMP<sup>-</sup> has a positive component around half the width of that in dGMP<sup>-</sup>, and extends to a lower *eKE* (~ 2.1 eV compared to ~ 3.1 eV). The first observation is perhaps a reflection of a flatter potential energy surface in the Franck Condon region for dGMP<sup>-</sup>. The lower *eKE* cut-off is simply a reflection of the lower adiabatic binding energy of the nucleobase in dGMP<sup>-</sup> compared to dAMP<sup>-</sup> (4.65 eV vs. 5.65 eV), which is partly due to the configuration of the nucleotide, and partly due to an innately low binding energy in Gua.<sup>28, 69</sup> In the temporal domain, the lifetime of dGMP<sup>-</sup> is around twice as long than for dAMP<sup>-</sup>, although still ultrafast. Given the similarities of the global fits, it seems probable that both dAMP<sup>-</sup> and dGMP<sup>-</sup> follow a similar relaxation mechanism. The agreement of dGMP<sup>-</sup> dynamics with solution phase experiments, as well as our previous work on dAMP<sup>-</sup>, suggest that this mechanism follows a pathway involving only one electronic state, with no explicit internal conversion to a <sup>1</sup>nπ\* state. Precise mechanistic details will of course differ, as the differing excited state lifetimes indicate, however, to a first approximation, a unified mechanism for

relaxation of purine nucleotides seems feasible. Moreover, it appears that a solvent environment has a relatively small impact on the dynamics of the purine nucleotides. This makes our current findings particularly relevant as they can be directly compared to solution-phase work. In the DNA polymer however, the environment does cause major changes, in particular for Ade, where base-stacking leads to the formation of excimer states that are much longer lived.<sup>1, 2, 21</sup> Nevertheless, the direct internal conversion mechanism for dAMP<sup>-</sup> and dGMP<sup>-</sup> remains active even in the much larger systems and is a key mechanism for transferring electronic energy into harmless heat.<sup>2, 21, 70</sup>

The combination of ESI and photoelectron imaging allows us, for the first time, to record the time resolved photoelectron spectrum of a Gua moiety in the gas phase (albeit in its nucleotide). Moreover, the gas phase spectroscopy of the dGMP<sup>-</sup> provides a useful comparison with solution phase work on the same species where the charge is screened and certain electronic states are dramatically shifted (e.g. <sup>1</sup>nπ\* states).

### Dynamics of dTMP<sup>-</sup> excited at 4.66 eV

A number of experimental studies have focused on the decay dynamics of isolated Thy, when excited at 4.7 eV (266 nm), which are of particular interest as photoinitiated cyclo-addition between two Thy groups is a leading form of DNA damage.<sup>71</sup> The dominant conformer in molecular beams is the biological keto tautomer.<sup>72, 73</sup> TR-IY experiments using multiphoton probes showed that isolated Thy has a much longer lived excited state than the other nucleobases; three decay constants could be associated with the relaxation, with lifetimes of approximately 100 fs and 5 ps for the first two decays.<sup>9, 10, 74</sup> The third feature has a lifetime of 2.8 ns, and has been assigned to intersystem crossing onto the <sup>3</sup>ππ\* state, by measurement of the vibrational spectrum of the transient species.<sup>74</sup> TRPES, using a 4.96 eV (250 nm) pump and a 6.20 eV (200 nm) probe, were fit using three time constants of < 50, 530 and 6400 fs.<sup>5</sup> Presumably owing to the differing photoexcitation and single-photon ionisation energies, the TRPES data are not in perfect agreement with the TR-IY data. Nevertheless, a pattern for the early dynamics in which an ultrafast decay, followed by a decay of ~5 ps, emerges.

Numerous computational studies have been performed in an attempt to disentangle the initial dynamics of Thy. In the vertical excitation region, the S<sub>1</sub> is an optically dark state of <sup>1</sup>nπ\* character, whilst the S<sub>2</sub> is the optically bright <sup>1</sup>ππ\* state.<sup>75</sup> CIs have been located between the <sup>1</sup>ππ\* and both <sup>1</sup>nπ\* and S<sub>0</sub> states: the former has a boat conformation, puckered at C<sup>6</sup> and C<sup>3</sup>, while the latter is puckered at the C<sup>5</sup> atom (see Fig. 3a for atom labels).<sup>64, 75-79</sup> From these computational studies, a number of possible explanations for the dynamics of Thy have emerged. Domcke and coworkers<sup>76</sup> calculated a barrierless reaction path to the <sup>1</sup>ππ\*/S<sub>0</sub> CI from the vertical Franck-Condon region, responsible for the fast time constant, whilst the slower time constant was assigned to trapping of excited state flux on the optically dark <sup>1</sup>nπ\* state. On the other hand, Merchan *et al.* attributed the slower dynamics to trapping on the <sup>1</sup>ππ\* state,

without active contribution from the <sup>1</sup>nπ\* state.<sup>78</sup> Interpretation of the time resolved photoelectron spectra, using the *ab initio* multi spawning (AIMS) technique, concluded that the initial femtosecond component results primarily from motion on the <sup>1</sup>ππ\* state, without any change of electronic character.<sup>11</sup> The slower component was then suggested to be borne mainly from a barrier crossing at the <sup>1</sup>ππ\*(S<sub>2</sub>)/<sup>1</sup>nπ\*(S<sub>1</sub>) CI. A key observation, however, was that multiple CI pathways are available, rather than just a single decay mechanism.

In the solution phase, Hare *et al.* found that a wavepacket bifurcation occurs in excited Thy, which may decay via two mechanisms.<sup>2, 80</sup> The excited state population mostly (~90%) undergoes direct relaxation via a <sup>1</sup>ππ\*/S<sub>0</sub> CI with sub-picosecond timescales,<sup>1, 17-19</sup> while the remaining population becomes trapped on the <sup>1</sup>nπ\* state for 30 ps. They found that solvated dTMP displayed similar early time dynamics, but the slow component has a lifetime of 127 ps, indicating a much larger barrier to deactivation of the <sup>1</sup>nπ\* state after ribosyl substitution. Similar to the purine bases, calculations show that solvation destabilizes the <sup>1</sup>nπ\* state significantly more than the <sup>1</sup>ππ\* state,<sup>81</sup> so comparisons between gas and solution phases may not be valid in mechanisms where the <sup>1</sup>nπ\* state plays a significant role.

The dynamics of deprotonated nucleotides presented here do not directly correlate with those of isolated nucleobases, although some similarities are observed. By analogy with the purine nucleotides presented above, as well as with reference to the sizable body of work describing the potential energy surfaces of Thy,<sup>75-77, 79, 82</sup> the initial τ<sub>1</sub> = 60 fs time constant can likely be assigned to motion of the wavepacket on the S<sub>2</sub> <sup>1</sup>ππ\* state, from the vertical Franck-Condon region towards some local minimum. Assignment of the two slower features is less straightforward, and we shall use the identity of the long lived component, τ<sub>3</sub>, to guide this discussion. As this feature manifests as a flat offset on the timescale of our experiment, of too small intensity to accurately fit, we can only assign it a lifetime ≥ 5 ps, which could in principle stretch into the nanosecond regime. If this feature is indeed very long lived (nanoseconds), then it is likely a triplet feature, as seen in neutral Thy.<sup>74</sup> The spectrum of k<sub>3</sub> should in principle aid with this assignment – the detachment energy of the triplet state can be calculated, and compared with the experimental spectrum. Even without direct spectral knowledge of k<sub>3</sub>, calculations should be able to inform as to whether or not detachment from the triplet state is energetically feasible. If the τ<sub>3</sub> decay is associated with the triplet state decay, then presumably the 210 fs component corresponds to ground state relaxation via a <sup>1</sup>ππ\*/S<sub>0</sub> CI. In this scenario, the <sup>1</sup>nπ\* state does not appear to play an active role, and we must presume that the picosecond channel observed in both isolated Thy and solvated dTMP has been quenched by either the presence of the charge or the unique geometry of the nucleotide. This scenario is shown schematically in Fig. 5a.

A second possible assignment of τ<sub>3</sub> is to the internal conversion from the <sup>1</sup>nπ\* state. In this scenario, the lifetime is likely to be on the order of picoseconds, as in neutral Thy. The



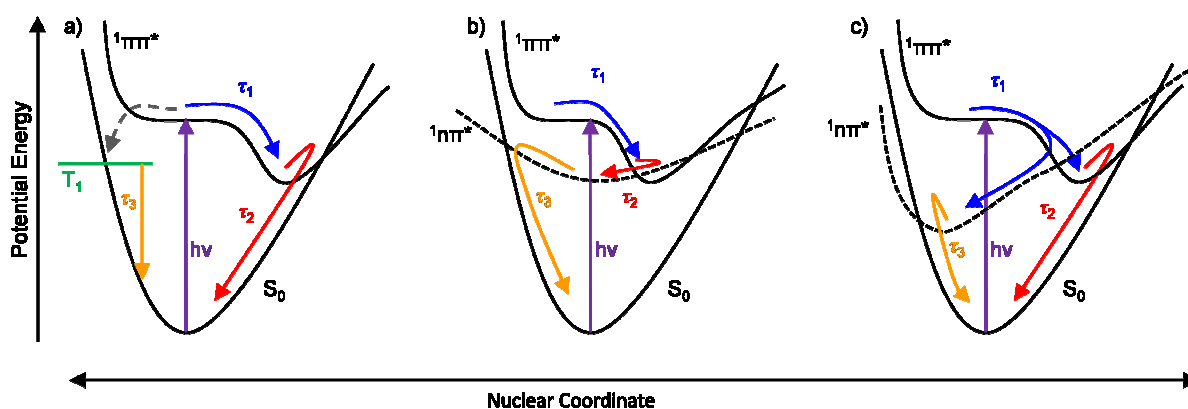


Fig. 5 Three possible mechanisms to describe the dynamics observed in dTMP<sup>-</sup>. a) The two fast components correspond to motion upon the <sup>1</sup>ππ\* state, followed by internal conversion, akin to dGMP<sup>-</sup>. The long lived state is the result of intersystem crossing to a triplet state. b) The two fast time constants correspond to motion on the <sup>1</sup>ππ\* state, followed by conversion to the <sup>1</sup>nπ\* state, and the slow time constant is relaxation of the <sup>1</sup>nπ\* state. c) The majority of population proceeds via a <sup>1</sup>ππ\* → S<sub>0</sub> route, whilst a small portion becomes trapped upon the <sup>1</sup>nπ\* state.

$\tau_2 = 210$  fs component would then correspond to <sup>1</sup>ππ\* → <sup>1</sup>nπ\* internal conversion. Fig. 5b outlines this relaxation scheme. It is important to note that the 60 fs  $\tau_1$  component has negligible absolute amplitude, and would be difficult to detect in non-differential measurements such as ion yield experiments.<sup>9, 10</sup> In these, only the 210 fs  $\tau_2$  component would then be observed. However, this lifetime is longer than that observed in experiments on neutral Thy (~100 fs), although variations in the potential landscape due to the addition of the sugar/phosphate backbone and differences in initial temperature may well account for this. Note that in our experiments, the internal temperature is on the order of 300 K, whereas in a molecular beam, this is much lower: typically on the order of 10s K. The very small amplitude of  $k_3$  appears to disagree with this model. But, our spectral window is limited, and the 3.1 eV probe may be too low for the majority of detachment from the <sup>1</sup>nπ\*. If this was the case, a long lived component should be visible from two-photon detachment. No significant amplitude is seen in the DAS of the long lived component at  $eKE > 2.1$  eV, although it may simply be the case that detachment signal from the <sup>1</sup>nπ\* state is very weak due to a small cross section.

Finally, it is feasible that  $\tau_3$  is due to bifurcated population trapped on the <sup>1</sup>nπ\* state, as has been observed in solution,<sup>80</sup> and suggested by photoelectron spectroscopy and *ab initio* calculations on Thy.<sup>11</sup> Fig. 5c gives a sketch of this mechanism. In this scenario, the remaining dynamics are likely similar to those described for the triplet scenario, with  $\tau_1$  being associated with rapid initial motion, and  $\tau_2$  with internal conversion from the <sup>1</sup>ππ\* state to the ground state. We note that solvation significantly destabilizes the <sup>1</sup>nπ\* state (as in the purine bases), and hence this mechanism is likely perturbed strongly in the gas phase.<sup>71, 77</sup> The short term dynamics are otherwise rather comparable with those in solution, perhaps again indicating a less active role of the <sup>1</sup>nπ\* state than previously thought.

The above comparisons have all been between dTMP<sup>-</sup> and neutral Thy, and hence one may query whether substitution at the N<sup>1</sup> site has a great effect on the dynamics. The ion yield experiments of Brutschy and co-workers showed that almost

identical dynamics are obtained for Thy and 1-methylthymine,<sup>74</sup> suggesting that the hydrogenated nucleobase is an acceptable model for the dynamics of Thy in more complex environments, however as this was a non-differential measurement, there may be additional masked features. Overall, it appears that although the specifics of the relaxation mechanism for dTMP<sup>-</sup> are debatable, our results represent important new information to which theoretical models can be benchmarked. Specifically, the spectral shape of the DAS should provide a basis for assigning the decay components to specific processes by comparison with calculated time-resolved photoelectron spectra.

#### Dynamics of dCMP<sup>-</sup> excited at 4.66 eV

The dynamics of cytosine (Cyt) in the gas phase have, until recently, been poorly understood. Early femtosecond experiments, using either TR-IY<sup>9, 10</sup> or TRPES<sup>5</sup> showed poor agreement. The ion yield experiments, performed with a 4.7 eV pump, gave biexponential decay kinetics with a fast component limited by the instrument response, and a slower component with a lifetime of 1.9 – 3.2 ps.<sup>9, 10</sup> The TRPES, performed with a 4.96 eV (250 nm) pump, resulted in triexponential kinetics, with lifetimes of <50 fs, 820 fs and 3.2 ps (the latter lifetime fixed to the TR-IY results by Kang *et al.*).<sup>5</sup>

The origin of this confusion is almost certainly due to multiple tautomers present in the molecular beam. Microwave<sup>83</sup> and high resolution IR<sup>84, 85</sup> spectroscopies reveal that there are (at least) three tautomers of Cyt present within a molecular beam. These are the biological keto tautomer, and the non-biological enol and imino tautomers, which are present in a ratio of approximately 4:4:1, respectively (from microwave experiments).<sup>83</sup> In nucleotides, the sugar/phosphate backbone is attached on the N<sup>1</sup> site, and hence the enol tautomer is unavailable in biological systems, whilst the imino tautomer interferes with base pairing, and so the primary biological tautomer is the keto form.<sup>86</sup>

Three recent experimental publications have sought to untangle the dynamics within Cyt to isolate the behaviour of

each tautomer, and in particular to ascertain the nonadiabatic relaxation of the keto tautomer. Kosma *et al.* noted that the onset of absorption to the  $S_1$  of the enol and imino tautomers is energetically higher than in the keto tautomer:<sup>87</sup> the keto-Cyt onset is  $\sim 3.90$  eV (318 nm), whilst the enol-Cyt origin is  $\sim 4.46$  eV (278 nm).<sup>88, 89</sup> Thus, by restricting pump wavelengths to  $\lambda \geq 280$  nm, these authors could selectively observe only dynamics of the keto-Cyt tautomer. Dynamics were tracked using TR-IY, pumped using either 4.28 or 4.43 eV (290 or 280 nm) radiation, and probed using 3 photons at 1.55 eV (800 nm). Three components were observed in the decay, with lifetimes of  $<100$  fs, 1.1 (1.2) ps and  $>150$  (55) ps for excitation with 4.28 (4.43) eV. The proposed assignment was ultrafast motion away from the Franck-Condon region following photoexcitation to the  $^1\pi\pi^*$  state, followed by internal conversion directly to the ground state. The long lived component was tentatively assigned to an excited state tautomerization,<sup>87</sup> recent theory has shown this to be feasible.<sup>90</sup>

Ho *et al.* used chemical substitution to measure tautomer-specific behaviour:<sup>12</sup> methylation at the  $N^1$  position (where the ribose group attaches, see Fig. 4a) blocks keto-enol tautomerization, and also better represents the biological nucleotide. TR-IY experiments on Cyt-1Me with a range of pump energies,  $4.13 \text{ eV} \leq h\nu \leq 4.77 \text{ eV}$  ( $260 \text{ nm} \leq \lambda \leq 300 \text{ nm}$ ), and a 1.55 eV probe, gave biexponential decay kinetics composed of an initial spike limited by the instrument response, and a decay with a lifetime of 0.4 - 1.3 ps (with shorter lifetimes at higher pump energies). Intriguingly, when pumped at 4.00 eV, the lifetime increases to 3.2 ps. The authors did not observe a slower component, and assigned the slower component observed by Kosma *et al.* to dynamics of fragments. Power dependency studies showed the initial spike to be primarily multi-photon in nature, although a one-photon component could not be ruled out.<sup>12</sup>

Very recently, Lobsiger *et al.* have measured linewidths in resonance-enhanced multiphoton ionisation experiments in order to determine the lifetime of keto-Cyt near the  $S_1$  origin and shed light on the femtosecond measurements.<sup>13</sup> It was found that the lifetime of Cyt is *not* ultrafast at the absorption onset, and rapid dynamics only appear after  $\sim 500 \text{ cm}^{-1}$  (60 meV) of excess energy is imparted in the  $S_1$  state. The predominant vibrational progression was in an out-of-plane mode, and this was taken as evidence that relaxation occurs via a non-planar CI, as long as there is sufficient energy to access it.

Theoretical studies have gone a considerable way towards understanding the nonadiabatic dynamics of keto-Cyt. It is generally agreed that the  $S_1$  state is of  $^1\pi\pi^*$  character in the Franck-Condon region. The next two excited states are of  $^1n\pi^*$  character, with the non-bonding orbital residing on either the  $N^3$  ( $^1n_N\pi^*$ ) or the O ( $^1n_O\pi^*$ ) atom.<sup>91-93</sup> Several CIs with the ground state have been located, the three most accessible being a  $^1\pi\pi^*/S_0$  CI from puckering on the  $C^6$  atom, a  $^1n_N\pi^*/S_0$  CI from puckering at the  $N^3$  atom, or a semiplanar  $^1n_O\pi^*/S_0$  CI.<sup>64, 78, 91-96</sup> Computations utilizing either minimum energy paths or surface hopping techniques have disagreed as to whether it is the  $C^6$

pucker<sup>78, 92</sup> or the semiplanar<sup>97, 98</sup> CI which dominates, however very recently high level calculations have suggested that the  $^1\pi\pi^*/S_0$  CI is the primary relaxation route, and that earlier calculations over-emphasized the importance of the  $^1n_O\pi^*$  state.<sup>91</sup> This result is also supported by calculations on 5-fluorocytosine, which prevents facile state switching to the  $^1n\pi^*$  states.<sup>99, 100</sup> A mechanism utilizing a puckered CI accords well with the findings that vibrational energy in out-of-plane modes is required before rapid relaxation can occur.<sup>13</sup>

In the solution phase, the dynamics of Cyt and CMP (the RNA nucleotide of Cyt) when excited at 4.66 eV are very much reminiscent of those of Thy and dTMP.<sup>2, 80</sup> Bifurcation of the population occurs between the optically bright  $^1\pi\pi^*$  state and a dark state, presumably of  $^1n\pi^*$  character. The dynamics of the rapid component are sub-picosecond,<sup>1, 17-19</sup> whilst the slow, dark, component decays in 12 ps.<sup>80</sup> As with Thy, moving to the nucleotide results in an increase in dark state lifetime, in this case to 34 ps, which indicates an increased barrier to relaxation. The dark state yield in CMP is considerably higher than in Cyt, at 41% and 9%, respectively,<sup>80</sup> perhaps indicative of a decrease in rate of decay purely upon the  $^1\pi\pi^*$  state, and a change in dynamics upon building up to the nucleotide.

Our results show that dCMP<sup>-</sup> is remarkable in that it is the only isolated nucleotide which possesses a significant component with a lifetime  $>1$  ps. In particular, it is the only example where the hot nucleotide has longer lived dynamics than the isolated, cold nucleobase; Cyt-1Me has a lifetime of 0.5 ps when pumped at 4.59 eV.<sup>12</sup> There are two basic explanations for the increased lifetime of dCMP<sup>-</sup> compared to Cyt-1Me. Firstly, the general mechanism may be the same, relaxation mediated by the  $^1\pi\pi^*/S_0$  CI, but the nucleotide may not reach the CI as easily as the methylated nucleobase. The CI geometry is thought to involve a puckering of the ring at the  $C^6$  atom,<sup>64, 78, 91-96</sup> and it may be that inclusion of the sugar and phosphate impedes the required out of plane motion. In this scenario, all observed dynamics are borne from the  $^1\pi\pi^*$  state, and the explanation follows that of the other nucleobases, with initial motion from the Franck-Condon region, followed by slower relaxation via the CI. The DAS appear to support such a mechanism, as they are qualitatively very similar to those of the other three nucleotides.

Alternatively, the mechanism for relaxation of dCMP<sup>-</sup> and Cyt-1Me need not be the same. The large difference in  $^1n\pi^*$  utilization between Cyt and CMP in solution suggests that a significant modification to the electronic landscape of the chromophore may occur upon addition of the backbone.<sup>80</sup> Involvement of either of the  $^1n\pi^*$  states cannot be ruled out. In fact, one could surmise that a bifurcation occurs, and that the two relaxation lifetimes reflect a split between relaxation using the  $^1\pi\pi^*$  and  $^1n\pi^*$  states. The lack of observed negative DAS tenuously supports this hypothesis.

Comparing the pyrimidine nucleotides, dTMP<sup>-</sup> and dCMP<sup>-</sup>, less agreement in the excited state dynamics is observed than with the purine nucleotides. They both appear to follow a similar structure of rapid dynamics at higher *eKE*, followed by a slower decay in the lower *eKE* region. dTMP<sup>-</sup> employs

definite sequential dynamics, while for dCMP<sup>-</sup> this is less definitive. Looking at the dCMP<sup>-</sup> time resolved spectra, normalized by each *eKE* slice (Fig. S3 in the supporting information), a small motion of maximum intensity towards later delays appears to be visible. This suggests that perhaps dCMP<sup>-</sup> also engages initially in rapid wavepacket motion, followed by exponential decay, as in the other nucleotides. On the other hand, the longer lifetimes observed in dCMP<sup>-</sup> indicate that relaxation must be more hindered in this system. The increase of excited state lifetime of dCMP<sup>-</sup> compared to Cyt-1Me<sup>12</sup> suggests that the presence of the sugar/phosphate backbone is very significant for the dynamics of this system. This is in opposition to the purine bases, where the dynamics seem to be solely driven by the nucleobase and are not affected by the backbone. dTMP<sup>-</sup> dynamics appear also to be mainly determined by the Thy nucleobase, but further work is required to assign the identity of the third component in the dTMP<sup>-</sup> dynamics. The spectral information provided through the DAS provides important information to aid in assigning the dynamics and, in combination with excited state calculations that consider the final states accessible through ionisation, a clearer picture of the potential energy surfaces and mechanisms may be generated. This is a prerequisite for studying processes that occur in the DNA polymer such as the photoinduced cyclo-addition of a pair of Thy nucleobases.<sup>101</sup>

Going forward, dCMP<sup>-</sup> perhaps presents the greatest mystery of the nucleotides studied. It is the only nucleotide which shows dynamics slower than the corresponding isolated nucleobase, tentatively suggesting a change in relaxation mechanics as the system is built up. Future improvements to the instrument should improve the signal-to-noise observed in the dispersed spectra, so that dCMP<sup>-</sup> can be studied more robustly, and a more concrete result obtained. Nonetheless, these measurements provide important benchmark data for theoretical simulations of dCMP<sup>-</sup> relaxation mechanisms.

### Neutral-in-Anion Photoelectron Spectroscopy

A novel aspect of these experiments is that, although the overall nucleotide system is anionic, the phosphate group carries the excess charge and so the probed nucleobase moiety is essentially neutral. This “neutral-in-anion” scheme carries a number of interesting experimental features. The proximity of a charge to the nucleobase site lowers the electron binding energy, which means that longer wavelength lasers can be used to produce photoelectron spectra. In nucleotides, the detachment energy from the base moiety is shifted by ~3 eV relative a neutral nucleobase.<sup>28</sup> This shift is on the order of that expected by purely Coulombic interactions between a base and a charged phosphate group separated. It corresponds to a separation of ~5 Å, which is in reasonable agreement with the distance in calculated nucleotide anion structures.<sup>102</sup> The charge also influences the energies of other molecular orbitals. For example, calculations on dAMP<sup>-</sup> suggest that the  $\pi$  orbitals are shifted by a similar amount to the  $\pi^*$  orbitals.<sup>26</sup> In solution, where the charge is screened, the transition energies are similar between gas- and solution-phases, but slightly red-shifted. Non-

bonding orbitals are affected differently and lead to transition energies of the  $1n\pi^*$  states that are blue shifted relative to the isolated Ade. Experimentally, the spectral profile of the action (absorption) spectra of dGMP<sup>-</sup> and dAMP<sup>-</sup> looks similar to the solution phase spectra. dGMP<sup>-</sup> shows no appreciable shift while dAMP<sup>-</sup> is blue shifted 0.12 eV in the gas-phase.<sup>103</sup> These small shifts in absorption spectra are in line with our observations that solvation has a small effect on the excited states of the purine nucleotides. The spectrum of isolated dTMP<sup>-</sup> is similar to that in solution but slightly broadened. Interestingly, the dCMP<sup>-</sup> action spectrum has a somewhat different shape to that in solution and is also the system that shows the most deviations from its base dynamics.

For time-resolved experiments, where very large datasets are required, background noise produced by deep UV lasers can be problematic, and so the possibility to use a near UV probe significantly enhances signal-to-noise. On the other hand, if background noise can be reduced, then using deep UV on a neutral-in-anion system allows for much deeper probing than in just the neutral alone. Caution must be applied to using a deep UV probe for nucleotides, however, as such a probe will produce a very large one-photon signal from detachment at the phosphate, which may mask the pump-probe signal. Additionally, such a probe would also be resonant with the excited states of the nucleobase. At 3.10 eV, there are no resonances, which makes the interpretation of the time-resolved spectra more straightforward.

In photoelectron spectroscopy of anions, the photodetachment cross section follows Wigner's law,<sup>104, 105</sup> and is expected to be zero into a given open channel at the threshold of that channel. The consequence is that at *eKE* = 0, no photoelectron signal is expected. In all the spectra presented here, this is not the case and a strong photodetachment signal is observed at *eKE* = 0. A possible explanation for this is the fact that the electron is detached from a *neutral* region of the molecule that is spatially well separated from the anionic phosphate group. Hence, photodetachment produces a cation locally (and a zwitterion overall) for which the Wigner law predicts a finite cross section at threshold. These Coulomb interactions may relax the Wigner law for large anions with Zwitterionic final states.

In our experiments, we have assumed that the presence of a nearby charge does not significantly perturb the structure and dynamics of the nucleobase moiety. The similarities in dynamics in the purines between neutral nucleobases and anionic nucleotides suggest that this assumption may hold. On the other hand, each nucleotide has a different and fluctuating gas phase conformation,<sup>27</sup> so the charge may not remain a spectator in all systems. Further theory is required to explore the effect of the charge on the dynamics and excited states. A second concern of the experimental scheme is that detachment leads to a zwitterionic final state, where an electron resides on the phosphate, and a hole is on the nucleobase. Resonance enhanced two-photon photodetachment showed that this zwitterion is the lowest lying neutral state for dGMP<sup>-</sup>,<sup>28</sup> whilst in the other nucleotides, the neutral species without a charged

phosphate is lower in energy. These considerations will have no effect on comparisons between nucleotides, as the same class of final state is accessed in all of them. Likewise, the nature of the final state should not affect dynamics comparisons between anions and neutral species, although different Franck-Condon factors may complicate detailed spectral comparisons.

## Conclusions

Using combination of ESI mass spectrometry and TRPES, the relaxation dynamics of excited deprotonated nucleotides have been measured. For dGMP<sup>-</sup>, this represents the first time TRPES has been applied to a gas phase Gua moiety. The relaxation dynamics of dGMP<sup>-</sup> are analogous to those of dAMP<sup>-</sup>, suggesting that both purine nucleotides relax in a similar manner with biexponential decay kinetics assigned to wavepacket motion on the <sup>1</sup>ππ\* state followed by internal conversion on the <sup>1</sup>ππ\* to the S<sub>0</sub> through a CI. dTMP<sup>-</sup> showed behaviour reminiscent of isolated Thy, though not identical. A long lived component was seen in the relaxation of Thy, in addition to two rapid components. Three feasible mechanisms were presented to interpret these dynamics of which rapid wavepacket motion on the <sup>1</sup>ππ\* state, followed by internal conversion to the ground state is the most consistent with our data. The long lived component is probably due to population trapped in the <sup>1</sup>nπ\* state, or possibly in a triplet state. Finally, dCMP<sup>-</sup> showed slower dynamics than the other nucleotides. For this the dynamics likely involve motion entirely on the <sup>1</sup>ππ\* surface, as with the other nucleobases. Intriguingly, dCMP<sup>-</sup> is longer lived than isolated Cyt, suggesting that the presence of the sugar/phosphate backbone inhibits the system reaching an appropriate CI geometry.

These results present benchmark data for the dynamics of isolated nucleotides. Advances in *ab initio* techniques will allow complex simulations to be directly compared to experimental results, providing high levels of detail about the relaxation mechanisms. These results also present a first step towards understanding the dynamics of more complex isolated DNA systems, such as oligonucleotides with stacked or paired bases, which can also be generated using ESI. Finally, our results show that a direct comparison between solution phase and gas phase experiments on these systems provides deeper insights by considering the effects solvation has on the location of excited states.

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## Notes and references

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