

This is an *Accepted Manuscript*, which has been through the Royal Society of Chemistry peer review process and has been accepted for publication.

Accepted Manuscripts are published online shortly after acceptance, before technical editing, formatting and proof reading. Using this free service, authors can make their results available to the community, in citable form, before we publish the edited article. This Accepted Manuscript will be replaced by the edited, formatted and paginated article as soon as this is available.

You can find more information about *Accepted Manuscripts* in the **Information for Authors**.

Please note that technical editing may introduce minor changes to the text and/or graphics, which may alter content. The journal's standard <u>Terms & Conditions</u> and the <u>Ethical guidelines</u> still apply. In no event shall the Royal Society of Chemistry be held responsible for any errors or omissions in this *Accepted Manuscript* or any consequences arising from the use of any information it contains.



www.rsc.org/advances

Atomistic insight into the oxidation of monolayer transition metal dichalcogenides: from structures to electronic properties

Hongsheng Liu, Nannan Han, Jijun Zhao*

Key Laboratory of Materials Modification by Laser, Ion and Electron Beams (Dalian University of Technology), Ministry of Education, Dalian 116024, China

ABSTRACT: Monolayer transition metal dichalcogenides (TMDs) stand out in twodimensional (2D) materials due to their potential applications in future microelectronic and optoelectronic devices. In experiments, field effect transistors (FET) based on MoS₂ monolayer are sensitive to the environmental gases, especially O₂. Thus, the oxidation of monolayer TMDs becomes a critical concern. By first-principles calculations, we reveal that perfect single-layer sheet of TMDs keep intact when exposed in O_2 due to the weak physical adsorption of O_2 . However, O₂ can be chemically adsorbed on monolayer TMDs (including MoS₂, MoSe₂, MoTe₂, WS₂, WSe₂, and WTe₂) with single vacancies of chalcogen, which are the most common defects in realistic TMD materials. The adsorption configurations and dissociation behavior of O₂ molecule at vacancy sites as well as the possible diffusion behavior of oxygen adatom on TMD monolayer surface are explored. Oxidation significantly influences the electronic properties of defective MoS₂ monolayer, while other defective TMD monolayers (especially MoTe₂ and WTe₂) suffer less from oxidation. Our theoretical results provide valuable atomistic insight into the oxidation of TMD monolayers and are useful for the future design of TMD-based 2D devices.

Introduction

The boom of graphene has initiated tremendous research interests in the other two-dimensional (2D) atomic crystals¹⁻³. Among these 2D materials, single-layer and few-layer sheets of transition metal dichalcogenides (TMDs) stand out due to their extraordinary physical properties as well as great potentials in technological applications⁴⁻⁷. MoS₂, a typical example of TMDs, forms a hexagonal layered crystal structure with weak interlayer van der Waals (vdW) interactions. Each layer of MoS2 is composed of one layer of molybdenum atoms sandwiched between two layers of sulfur atoms. The intra-layer bonding is covalent and strong, resulting in absence of dangling bonds at clean MoS₂ surface and thus no intrinsic surface states. Most excitingly, monolayer MoS₂ is a direct gap semiconductor with a quasiparticle band gap of 2.8~2.9 eV^{8,9} and an optical gap of 1.8 eV¹⁰, which are important for the applications in microelectronic devices. A monolayer MoS₂ field effect transistor (FET) was constructed by Kis's group¹¹, which exhibited a reasonable charge carrier mobility (~ 200 $\text{cm}^2 \text{V}^{-1} \text{S}^{-1}$) and a room-temperature current on/off ratio of 1×10^8 . In addition, 2D TMD materials show great potentials for many other applications, such as phototransistors^{12,13}, analog small-signal amplifier¹⁴, chemical sensor¹⁵⁻¹⁷, thermoelectric nanodevice¹⁸, ultrasensitive photodetector¹⁹, and solar $cell^{20}$.

Compared to the gapless graphene, MoS_2 and other 2D semiconducting TMD sheets have finite gaps (e.g., 2.8~2.9 eV for $MoS_2^{8,9}$, 2.31 eV for $MoSe_2^{21}$, 1.77 eV for $MoTe_2^{21}$, 2.91 eV for WS_2^{21} , 2.51 eV for WSe_2^{21} , 1.79 eV for WTe_2^{21}) but lower carrier mobility. Moreover, the performance of MoS_2 -based FET can be significantly affected by the environmental gases like oxygen. For instance, Qiu et al. found that bi-layer MoS_2 FET is very sensitive to oxygen and water ambient.²² Chemisorption of oxygen and water reduced the conductance by up to about

RSC Advances Accepted Manuscript

100 times; but this effect could be reversibly recovered by a simple vacuum annealing at 350 K. Park et al. also explored the influence of oxygen environment on tri-layer MoS_2 FET.²³ The current level of the MoS_2 -based FET and the carrier mobility of MoS_2 channel were substantially reduced in the O₂ environment, whereas the current can recover to ~80% of the original level after thermally annealing at 350 K in vacuum. They also revealed that polymethyl methacrylate protective layer coated on MoS_2 can effectively protect the MoS_2 FET from suffering the oxygen disturbance.

Since most applications in microelectronics and optoelectronics require sufficiently high electric conductivity and carrier mobility, the effect of oxygen exposure on the electronic properties of TMD-based devices has become a critical concern for their performance under normal conditions. Recently, Tongay and co-workers pointed out that O2 molecules are physisorbed on perfect monolayer TMDs and monolayer TMDs with sulfur-vacancies through DFT calculations.^{24,25} Here we performed systematical *ab initio* calculations to explore oxidation effect on the electronic properties of MoS₂, MoSe₂, MoTe₂, WS₂, WSe₂ and WTe₂ single layers. We show that perfect TMD monolayer sheets are inert to oxygen adsorption. However, chemical adsorption and dissociation of O2 molecule occurs at the vacancy sites of TMD monolayers exothermically. According to the computed band structures, the electronic properties of MoS₂ monolayer with vacancy defects suffers greatly from oxidation, while oxidation effect is relatively less prominent for MoTe₂ and WTe₂ sheets. Our theoretical results not only give atomistic insights into the oxidation configuration and electronic properties of TMD monolayers under ambient environments but also provide usefull guidance for the design of future microelectronic and optoelectronic devices.

Computational methods

All calculations were carried out using the Vienna Ab initio Simulation Package (VASP) based on the density functional theory (DFT).²⁶ The electron-ion interactions were described by the projector augmented wave (PAW) potentials.²⁷ To treat the exchange-correlation interaction of electrons, we chose the Perdew-Burke-Ernzerhof (PBE) functional within the generalized-gradient approximation (GGA).²⁸ A kinetic energy cutoff of 500 eV for the planewave basis and a convergence criterion of 10^{-4} eV for the total energies were carefully tested and adopted for all DFT calculations.

Due to the chemical inertness of perfect MoS₂ surface, oxidation of bulk MoS₂ mainly occurs at the edge-plane surfaces and defect sites on the basal-plane surface.²⁹⁻³¹ Monolayer or fewlayer MoS₂ materials fabricated from both micromechanical cleavage and CVD methods always contain some stable defects such as single and double sulfur vacancies, vacancy complexes of one molybdenum and three nearby sulfur, vacancy complexes of molybdenum and three nearby disulfur pairs, and antisite defects by exchanging molybdenum and sulfur atom.³²⁻³⁵ Among them, sulfur single vacancies are most frequently observed, while other kinds of defects are only occasionally found.^{34,35} Moreover, single sulfur vacancy can arise under electron irradiation.^{33,36} Therefore, here we only considered oxidation at the single chalcogen vacancy (as the major structural defects) of six semiconducting monolayer TMDs, including MoS₂, MoSe₂, MoTe₂, WS₂, WSe₂ and WTe₂.

In our simulation models, one chalcogen single vacancy (SV) was created in a (4×4) supercell of TMD monolayer, ensuring the distance between two single vacancies larger than 12 Å. Such supercell size corresponds to a vacancy density of 7.1×10^{13} cm⁻², comparable to the

RSC Advances Accepted Manuscript

experimental density (in the magnitude of 10^{13} cm⁻²) in the MoS₂ monolayer from micromechanical cleavage.³⁴ Our test calculations for bigger (6 × 6) supercells of MoS₂ and WS₂ (corresponding to a SV density of ~3.2 × 10^{13} cm⁻²) show no significant difference from the results of (4 × 4) supercells. The 2D systems were modeled by a slab model with a vacuum space of more than 12 Å to avoid interactions between adjacent periodic images. During geometry optimization, the spacing of **k** point grids was chosen to be 0.03 Å⁻¹ and a convergence criterion of 0.02 eV/Å for force was adopted. In order to obtain more accurate energies and electronic properties, denser **k** point meshes with uniform spacing of 0.015 Å⁻¹ were used. Bader analysis³⁷ was preformed to evaluate the charge transfer between TMD monolayers and O₂ molecule. Spin polarized DFT calculations were carried out for all geometry optimization.

The vacancy formation energy E_f^{SV} is defined as:

$$E_f^{SV} = E_{vac} + \mu_X - E_{TMD}, \tag{1},$$

where E_{vac} and E_{TMD} are the energies of the defective (with vacancy) and pristine supercells, respectively. The chemical potential μ_X of chalcogen is taken as the energy of one chalcogen atom in its solid phase (i.e. orthorhombic sulfur, α -Se and α -Te)³⁸. The absorption energy E_{ads} of an oxygen molecule on chalcogen vacancy of TMD monolayer is defined as:

$$E_{ads} = E_{vac} + E_{O2} - E_{oxi},\tag{2},$$

where E_{oxi} is the energy of the oxidized TMD supercell, E_{O2} is the energy of an isolated O₂ molecule in gas phase.

6

Results and discussions

Oxygen adsorption and dissociation

Firstly, all six perfect TMD monolayers were well relaxed along with the lattice constants. The calculated 2D lattice constants are 3.183 Å, 3.318 Å, 3.547 Å, 3.182 Å, 3.317 Å and 3.551 Å for MoS₂, MoSe₂, MoTe₂, WS₂, WSe₂ and WTe₂ monolayers, in agreement with previous results.^{21,39} We started from the perfect monolayers of TMDs and placed an oxygen molecule on the supercell of two representative systems, i.e., MoS_2 and $MoTe_2$. DFT-D2 method^{40,41} was used to properly take into account the long-range vdW interactions. After optimization, no chemical bond is ever formed between O2 and TMD layer, with the nearest distance between O atom and S/Te atom of about 3.4~3.5 Å. The calculated absorption energies of oxygen molecule are 58 meV and 10 meV for MoS₂ and MoTe₂ monolayers, respectively. The charge density difference between O₂ and perfect MoS₂ monolayer was also calculated and presented in Figure S1 of the Supporting Information. From the charge density difference, no chemical bond occurs between O₂ and perfect MoS₂ monolayer. Previous DFT calculations showed that O₂ can be physisorbed on the MoS₂ surface with a binding energy of 79 meV,²⁴ which is slightly larger than the absorption energy in our work. The small difference in the adsorption energy may originate from the different kinetic energy cutoff in DFT calculations. It is known that the gaseous O₂ molecule adopts triplet spin state, which is retained when O₂ molecule is adsorbed on perfect TMD monolayers. Such weak vdW interactions clearly indicate that perfect TMD monolayers cannot be oxidized owing to the chemically inert surface with saturated chalcogen ions. Further calculations of electronic band structures show that weak physical adsorption of O₂ has virtually no influence on the electronic properties of TMD monolayers. Hence, in the following discussions we will concentrate on the oxidation effect of defective TMD monolayers with single chalcogen vacancies.



Figure 1. Top views (upper) and side views (lower) of the atomic configurations of (a) MoS_2 monolayer with sulfur SV; (b) MoS_2 monolayer with O_2 molecule adsorption on sulfur SV; (c) oxidized MoS_2 monolayer after O_2 dissociation. Black circle in (a) represents sulfur SV.

For MoS₂, MoSe₂, MoTe₂, WS₂, WSe₂ and WTe₂, the structures of defective monolayer sheets with single chalcogen vacancy are nearly identical (Figure 1a). Creating a single chalcogen vacancy in the supercell of these TMD sheets does not induce large deformation of the original lattice, that is, the maximum displacement being 0.132 Å and 0.243 Å for chalcogen and metal atoms, respectively. Moreover, our spin polarized calculations show that O₂ molecule adsorbed on the vacancy site of defective TMD monolayers becomes non-magnetic due to formation of TM-O bonds. Table 1 lists the formation energies of a chalcogen single vacancy (E_f^{SV}) for all TMD monolayers being explored. For MoX₂ and WX₂ (X=S, Se, Te), the formation energy E_f^{SV} decreases with the increasing of atomic number of chalcogen, which can be understood by the reduction of electronegativity of chalcogen (thus weakening of Mo-X or W-X bond). Among them, WS₂ is most reluctant to form sulfur SV, while creation of sulfur SV is relatively easy from thermodynamic point of view.

Table 1. Formation energy of chalcogen single vacancy E_f^{SV} , adsorption energy E_{ads} and charge transfer (*Q*) of O₂ molecule on TMD monolayer (positive *Q* means that O₂ gains electrons).

TMD monolayer	$E_f^{SV}(eV)$	E_{ads} (eV)	<i>Q</i> (e)	
WS ₂	2.856	2.082	1.083	
WSe ₂	2.688	2.433	1.136	
WTe ₂	2.311	2.805	1.218	
MoS ₂	2.681	1.843	0.988	
MoSe ₂	2.609	2.155	1.044	
MoTe ₂	2.396	2.537	1.118	

For MoS₂, MoSe₂, MoTe₂, WS₂, WSe₂ and WTe₂ monolayers with SV, the adsorbed O₂ molecule forms covalent bonds with three surrounding TM atoms at the vacancy site, i.e., one O atom with two TM atoms, another O atom with one TM atom (see Figure 1b). The charge

RSC Advances Accepted Manuscript

density difference between O_2 and defective MoS_2 monolayer is shown in Figure S1b of the Supporting Information. Chemical bond between O atoms and Mo atoms and the weakening of O-O bond can be seen, indicating that O_2 adsorbs on defective MoS_2 monolayer chemically. This is quite different from previous DFT results, in which O_2 was predicted to be physisorbed on the sulfur vacancy sites of MoS_2 monolayer^{24,25}. As summarized in Table 2, the O-TM bond lengths for all oxidized TMD monolayer sheets are in the range of 2.10~2.12 Å, which are typical covalent bond lengths for transition metal oxides⁴². The O-O bond lengths (d_{O-O}) of O_2 molecules before and after adsorption on TMD monolayers are also presented in Table 2. Compared with the gas-phase molecule, d_{O-O} in the chemisorbed O_2 is elongated by 0.2~0.3 Å. Such weakening of O-O bond is a natural consequence of formation of O-TM bonds, which are expected to affect the electronic properties of the TMD sheets to certain extent.

Table 2. O-O bond lengths of O₂ molecules in vacuum and chemisorbed on defective TMD monolayers.

	In vacuum	MoS ₂	MoSe ₂	MoTe ₂	WS_2	WSe ₂	WTe ₂
$d_{\text{O-O}}(\text{\AA})$	1.23	1.44	1.45	1.48	1.46	1.48	1.53
$d_{ ext{O-TM}}(ext{\AA})$		2.12 2.06	2.12 2.08	2.12 2.07	2.11 2.05	2.10 2.04	2.10 2.04

The adsorption energies for O_2 on chalcogen vacancies are summarized in Table 1. For all TMD monolayers, the E_{ads} values are positive, indicating that oxidation of defective ultrathin TMD sheets is an exothermic process. Among all these TMD monolayers, the E_{ads} of MoS₂ is smallest, which corresponds to the weakest interaction between O_2 and vacancy. For MoX₂ (X =

S, Se, Te), the order of E_{ads} is MoS₂ < MoSe₂ < MoTe₂. This can be understood by the following picture. In the oxidized MoX₂, O₂ molecule bonds to the unsaturated Mo atoms at the vacancy site and gains electrons from Mo. Specifically, one Mo atom bonds with five X atoms and one O atom. Hence, there is a competition between O and X atoms to gain electrons from the Mo atom. The smaller electronegativity of X atom, the more electrons are transferred from Mo to O and thus the stronger Mo-O bond. Indeed, Bader charge analysis (see Table 1) confirmed this conjecture, i.e., O₂ molecule gains 0.988 electrons from defective MoS₂ sheet, 1.044 electrons from MoSe₂, and 1.118 electrons from MoTe₂, respectively. The same trend can be also applied to the WX₂ systems. In addition, with the same chalcogen atom (X), O₂ gains more electrons from WX₂ than MoX₂ and interacts more strongly with WX₂ since W has larger reducibility than Mo.



Figure 2. Dissociation process of O₂ on sulfur SV of MoS₂ monolayer.

RSC Advances Accepted Manuscript

To understand the oxidation behavior of defective TMD monolayers, we investigated dissociation of an O_2 molecule at the chalcogen SV with climbing-image nudged elastic band (cNEB) method⁴³. For all TMD monolayers considered, the dissociation processes of O_2 molecule at vacancy site remain the same; thus only that for MoS₂ is depicted in Figure 2 as a representative. The structure of oxidized monolayer MoS₂ after O_2 dissociation as the final state is also presented in Figure 1c, in which one oxygen atom at the vacancy site replaces the original missing sulfur and the other one at the top of the adjacent sulfur atom bonds with the sulfur atom perpendicular to the monolayer sheet. Naturally, the two dissociated oxygen atoms at the chalcogen SV are non-magnetic; thus spin polarization was not considered during the transition state search.

The energy barriers and heats of reaction for O_2 dissociation at vacancy site for all the TMD monolayers are summarized in Table 3. For all systems considered, dissociation of oxygen molecule at the chalcogen vacancy is exothermic with heat of reaction ranging from 1.49 eV to 2.61 eV. Oxygen molecule is most reluctant to dissociate at the vacancy of MoS₂ monolayer with an energy barrier of 0.93 eV, while the barrier remains high for WS₂ (0.86 eV) and MoSe₂ (0.71 eV). The dissociation barrier reduces to 0.58 eV for WSe₂ and 0.25 eV for MoTe₂, whereas there is even no barrier for O_2 dissociation in the case of WTe₂. Therefore, O_2 molecule adsorbed at the chalcogen vacancy sites of the defective MoTe₂ and WTe₂ monolayers would easily dissociate into two oxygen atoms at ambient conditions.



Figure 3. Diffusion behavior of single oxygen atom on perfect MoS₂ monolayer.

After the adsorbed O₂ molecule dissociates on a defective TMD monolayer sheet, one oxygen atom would stay in the original vacancy location and saturate the dangling bond, and another oxygen atom may either stay or migrate to the nearby region of the TMD sheet. The most energetically favorable absorption position for a single oxygen atom on perfect TMD monolayer is the top of chalcogen atom, in accordance with previous studies for oxygen absorbed MoS₂, MoSe₂, MoTe₂ and WS₂ monolayers^{44,45}. To gain further insight into the final structures of oxidized TMD monolayers, we explored the diffusion behavior of an oxygen atom on perfect TMD monolayer. The diffusion path from one top site to another nearby top site is shown in Figure 3, and the corresponding barriers are summarized in Table 3. The barriers for oxygen diffusion on MoS₂ and WS₂ are very high (over 2.5 eV), and it becomes lower on MoTe₂ and WTe₂ (around 1.0 eV). Nevertheless, oxygen adatoms are hard to diffuse on TMD monolayers.

RSC Advances Accepted Manuscript



Figure 4. Simulated STM images of (a) TMD monolayers with SV defects with the configuration corresponding to Figure 1a; (b) oxidized MoS_2 , $MoSe_2$, WS_2 and WSe_2 monolayers with the configuration corresponding to Figure 1b; (c) oxidized $MoTe_2$ and WTe_2 with the configuration corresponding to Figure 1c.

From the above discussions, we can deduce the most probably atomic structures of oxidized TMD monolayers with vacancy defects. For oxidized MoS₂, MoSe₂, WS₂ and WSe₂ monolayers with SV defects, their final structures should correspond to Figure 1b, with oxygen molecules adsorbing on the vacancy sites. For oxidized MoTe₂ and WTe₂ monolayers, the most likely structures could be described by Figure 1c, with some oxygen adatom sitting on the top of chalcogen atom and some oxygen adatoms residing at the vacancy site to replace the missing chalcogen atom. To help recognize the configuration of oxidized TMD monolayers in experiments, the STM images of six defective and oxidized TMD monolayers were simulated at bias of -1 V and displayed in Figure 4. In the STM images, each bright point represent a chalcogen atom. As shown in Figure 4a, a SV defect can be clearly seen in the STM image with absence of a bright point in the perfect triangular lattice of bright points, in agreement with previous results.⁴⁶ To our surprise, the chemisorbed oxygen molecule cannot be seen in the STM image (Figure 4b), which hints that some other methods should be used to detect the chemisorbed oxygen molecule. After the dissociation of oxygen molecule, a big bright point

14

appear in the STM image (Figure 4c) that corresponds to the oxygen adatom at the top of one chalcogen atom, while the other oxygen atom at the vacancy site is still invisible. The different oxidization configurations of defective TMD sheets would lead to different effects on the electronic properties, as we will discuss in the following.

Table 3. Energy barrier (E_a) and heat of reaction (E_{dis}) for dissociation of O_2 on chalcogen SV of TMD monolayer sheets. Energy barrier for diffusion of oxygen atom on perfect TMD monolayers (E_{dif}).

	Oxidized TMD monolayers					
	MoS_2	MoSe ₂	MoTe ₂	WS_2	WSe ₂	WTe ₂
E _a (eV)	0.93	0.71	0.25	0.86	0.58	0.00
E _{dis} (eV)	2.61	1.59	1.54	2.54	1.49	1.49
$E_{dif}(eV)$	2.53	1.55	0.90	2.68	1.66	1.01

Oxidation effect on electronic band structures

Hereafter we will discuss the electronic properties of six monolayer TMD systems including the perfect sheets, defective sheets and oxidized sheets. Figure 5 displays the band structures of pristine, defective, and oxidized MoX₂ (X = S, Se, Te) monolayers in various circumstances. From our DFT calculations, the perfect MoX₂ monolayer sheets are direct gap semiconductors with band gaps of 1.67 eV, 1.44 eV and 1.09 eV for MoS₂, MoSe₂ and MoTe₂ respectively (Figure 5a, e, i), in good agreement with previous theoretical values^{21,39}. For the defective MoX₂ monolayer, some impurity states in the gap region are induced by chalcogen SV and the top of valence band is substantially disturbed by the SV defect (see Figure 5b, f, g), transforming the defective MoX₂ monolayer sheets into deeply p-type doped semiconductors. The same effect was also found in previous DFT calculations.^{34,45,46} Interestingly, the SV-induced impurity states in the gap region are completely removed upon adsorption of oxygen molecule (Figure 5c, g, k), simply because of the saturation effect by formation of TM-O bonds at the vacancy sites. However, the valence and conduction bands are still disturbed to certain extent with regard to the pristine MoX₂ monolayers.



RSC Advances Accepted Manuscript

Figure 5. Band structures of perfect, defective and oxidized MoX_2 monolayer sheets. (a), (e) and (i) present band structures of perfect MoS_2 , $MoSe_2$ and $MoTe_2$ monolayer sheets, respectively. (b), (f) and (j) show band structures of MoS_2 , $MoSe_2$ and $MoTe_2$ monolayer sheets with chalcogen SV (corresponding to the structure in Figure 1a), respectively. (c), (g) and (k) show band structures of defective MoS_2 , $MoSe_2$ and $MoTe_2$ monolayer sheets adsorbed by O_2 molecule (corresponding to the structure in Figure 1b), respectively; (d), (h) and (l) show band structures of defective MoS_2 , $MoSe_2$ and $MoTe_2$ monolayer sheets adsorbed by O_2 molecule (corresponding to the structure in Figure 1b), respectively; (d), (h) and (l) show band structures of defective MoS_2 , $MoSe_2$ and $MoTe_2$ monolayer sheets adsorbed with two separated oxygen atoms (corresponding to the structure in Figure 1c). The Fermi level is set to zero (denoted by red dotted line).

After adsorption of oxygen molecule on the defective MoS₂ monolayer, a nearly flat band arises at the top of valence band (Figure 5c), which undoubtedly has a disadvantageous influence on the electronic transport properties. This flat band mainly originates from O₂, since it disappears after dissociation of O₂ at the vacancy site (Figure 5d). A very sharp peak arises at the valence band maximum (VBM) conresponding to the flat band in the partial density of states (PDOS) shown in Figure S2a of the Supporting Information, which confirms that the flat band originates from O₂. Similar to MoS₂, O₂ adsorption on single MoSe₂ and MoTe₂ layers also induces an isolated band at the top of valence band, but with more pronounced dispersion (Figure 5g, k). In the PDOS of MoSe₂ (Figure S2b of the Supporting Information), there is also a sharp peak from O atoms near VBM, but with a small shoulder, which promotes the dispersion of top valance band. For MoTe₂, no sharp peak of O atoms arises in the PDOS. (Figure S2c) The different dispersion behaviors of the top valence bands can be related to the interaction strengths between O₂ molecule and TMD sheets. After chemisorption on the TMD layers, the electronic

level of the O_2 molecule will be broadened due to the formation of O-TM bonds. Intuitively, the stronger O-TM interaction, the broader molecular level of O_2 . As presented in Table 1, the interaction between O_2 and vacancy in MoS₂ is weakest among all TMD monolayers explored. Thus, O_2 molecule induces a rather flat band when it is adsorbed at the vacancy site of MoS₂ layer. On the contrary, for MoSe₂ and MoTe₂, the electronic states of O_2 at the top of valence band are broadened and show certain dispersion due to relatively stronger adsorption interaction.

After the dissociation of O_2 molecule, the features of valence and conduction bands of the oxidized TMD monolayer sheets are partly recovered compared with the perfect monolayers (Figure 5d, h, l). According to above discussions, O_2 molecule is difficult to dissociate at the vacancy sites of MoS₂ and MoSe₂ monolayers at ambient conditions. In order to reduce the oxidation effect on the electronic properties of MoS₂ and MoSe₂ monolayer sheets, appropriate thermal annealing is needed to promote the dissociation of adsorbed O_2 molecule. In contrast, most O_2 molecules would dissociate on the defective MoTe₂ monolayer at room temperature. Therefore, MoTe₂ suffers least from oxidation among three MoX₂ monolayers.



Figure 6. Band structures of perfect, defective and oxidized WX_2 monolayer sheets. (a), (e) and (i) present band structures of perfect WS_2 , WSe_2 and WTe_2 monolayer sheets, respectively. (b), (f) and (j) show band structures of WS_2 , WSe_2 and WTe_2 monolayer sheets with chalcogen SV (corresponding to the structure in Figure 1a), respectively; (c), (g) and (k) show band structures of defective WS_2 , WSe_2 and WTe_2 monolayer sheets adsorbed by O_2 molecule (corresponding to the structure in Figure 1b), respectively; (d), (h) and (l) show band structures of defective WS_2 , WSe_2 and WTe_2 monolayer sheets adsorbed by O_2 molecule (corresponding to the structure in Figure 1b), respectively; (d), (h) and (l) show band structures of defective WS_2 , WSe_2 and WTe_2 monolayer sheets adsorbed with two separated oxygen atoms (corresponding to the structure in Figure 1c), respectively. The Fermi level is set to zero (denoted by red dotted line).

From our calculations, the perfect WX_2 monolayer sheets are also direct gap semiconductors with band gaps of 1.81 eV (WS₂), 1.55 eV (WSe₂), and 1.08 eV (WTe₂), respectively (Figure 6a, e, i), in accordance with previous theoretical results.^{21,39} Similar to the MoX₂ systems, existence of SV defects leads to some impurity states in the gap region (Figure 6b, f, j). However, the valence and conduction bands are less affected by defects; only dispersion of the top valence band is slightly disturbed. When the W atom at chalcogen vacancy is saturated by O₂ adsorption, the impurity states disappear (Figure 6c, g, k). Different from the MoX₂ systems, no sharp peak of O atoms exists in the PDOS of WX₂ (X=S, Se, Te). (Figure S1d, e and f) Thus, no flat band occurs at the top of valence band after O₂ adsorption on the defective WX₂ monolayers. The large dispersion and parabolic shape are well preserved at both the top of valence band and the bottom of conduction band (Figure 6c, g, k). Meanwhile, modification of the conduction bands by O₂ adsorption is more pronounced than that for MoX₂ systems, that is, isolated bands at the bottom of conduction band emerge. After O₂ dissociation, the dispersion of bottom conduction

RSC Advances Accepted Manuscript

band becomes sharp and its shape resemble to that of perfect WX_2 monolayer sheets (Figure 6d, h, l). In other words, O_2 dissociation is indeed helpful to reduce the influence on the electronic properties due to oxidation. Generally speaking, the electronic properties of oxidized WS_2 and WSe_2 sheets should be described by the band structures shown in Figure 6c and Figure 6g respectively, since almost no O_2 molecule is able to dissociate at the vacancy site. On the other hand, all O_2 molecules would dissociate after adsorbing on the defective WTe_2 monolayer. Eventually, oxidation has only little influence on the electronic properties of WTe_2 monolayer, as demonstrated by the band structure shown in Figure 61.

Table 4. Effective masses of hole and electron carriers (in unit of electron mass m_e) for the perfect TMD monolayers without oxidation and defective TMD monolayers adsorbed by an O₂ molecule or two separated oxygen atoms.

	Perfect sheet		Defective sheet adsorbed with O ₂		Defective sheet adsorbed with two O atoms	
	VBM	CBM	VBM	CBM	VBM	CBM
MoS ₂	-0.574	0.472	-6.494	0.587	-2.833	0.452
MoSe ₂	-0.649	0.554	-0.855	0.528	-0.712	0.493
MoTe ₂	-0.694	0.578	-0.848	0.477	-0.818	0.468
WS_2	-0.410	0.308	-0.535	0.349	-0.457	0.299
WSe ₂	-0.439	0.340	-0.516	0.331	-0.498	0.314
WTe ₂	-0.417	0.316	-0.509	0.220	-0.527	0.269

To further understand the oxidation effect on electronic and transport properties of TMD monolayers, we calculated the effective masses of charge carriers at VBM and conduction band minimum (CBM) for the perfect TMD monolayers without oxidation and the defective monolayers with oxidation. The effective mass m^* can be easily derived from band dispersion $E_k(k)$ using the following formula⁴⁷:

$$m^* = \frac{\hbar^2}{d^2 E_k(k) / dk^2}$$
(3),

where \hbar is the reduced Planck's constant. The theoretical results are summarized in Table 4. First of all, the effective masses of charge carriers for perfect single MoS₂ layer from our band structure calculations coincide well with the previous values (-0.637 m_e for hole and 0.483 m_e for electron).⁴⁸

For the oxidized MoS₂ with O₂ adsorption (corresponding to the structure in Figure 1b), the hole effective mass at VBM is as large as $-6.494 m_e$, while the increase of electron effective mass at CBM is not prominent compared with the pristine MoS₂ monolayer. Even after O₂ dissociation, the hole effective mass is still very large ($-2.833 m_e$). In contrast, for oxidized MoSe₂ and MoTe₂ with O₂ adsorption, the hole effective mass at VBM increases by only about 0.2 m_e and that of electron carrier at CBM even decreases slightly compared with the corresponding perfect monolayers. O₂ dissociation, which easily occurs in oxidized MoTe₂ monolayer, weakens the oxidation effect on the effective mass of hole. It is known that the carrier mobility is inversely proportional to its mass.⁴⁷ Thus, MoTe₂ monolayer suffers least from oxidation among the three MoX₂ monolayers in terms of effective masses and mobilities of carriers.

As for WS₂ and WSe₂ monolayers, the hole effective mass at VBM increases by about 0.1 m_e and the electron effective mass at CBM only changes slightly upon O₂ adsorption. After the dissociation of O₂, the effective masses of hole and electron of oxidized TMD monolayers are partly recovered. For oxidized WTe₂ monolayer, in which the adsorbed O₂ molecules tend to dissociate at the vacancy sites, the final effective mass of hole (electron) are larger (smaller) than that of perfect sheet by 0.11 m_e (0.047 m_e). The overall carrier mobilities of oxidized WTe₂ layer would be only slightly downgraded with regard to the perfect one.

Implication to device applications

Our theoretical results provide important guidance for the selection of 2D TMD materials for future microelectronic devices. First of all, all perfect TMD monolayers are inert to O₂ molecule, whereas oxidation mainly occurs at the chalcogen single vacancies, which are the most common defects in MoS₂ monolayer materials^{34,35}. In order to avoid oxidation and the consequent effects on electronic properties, it is thus desirable to reduce the defect density during the fabrication of TMD sheets. Existence of chalcogen SV induces some impurity states in the gap region and transforms these defective TMD monolayers into deeply p-type doped semiconductors. Interestingly, all SV-induced impurity states in the gap region are completely removed after oxidation, because of the saturation effect by formation of TM-O bonds at the vacancy sites. However, the valence bands and conduction bands are disturbed with regard to the original monolayer sheets to different extents, depending on the specific composition of TMD. The different oxidation effects on the electronic properties of defective TMD monolayers can be related to the interaction strengths between O₂ molecule and the vacancy site of TMD sheets.

24

Dissociation of O₂ at the vacancy site further weakens the oxidation effect on electronic properties of TMD monolayers. Fortunately, O₂ can easily dissociate on the oxidized MoTe₂ and WTe₂ layers. But dissociation of O₂ is rather difficult in the other four oxidized TMD monolayer sheets due to large energy barrier (0.58~0.93 eV). Among them, MoS₂ monolayer suffers most from oxidation, with a flat band at the top of valence band. In the realistic MoS₂-based devices like FET, a protective layer is usually required.²³ Compared to MoS₂, oxidation effect is less prominent in the WS₂ and WSe₂ monolayer materials. Recently, ambipolar FET based on few-layered WS₂ sheet and p-type FET composed of monolayer WSe₂ sheet have been successfully fabricated in experiments.^{49,50} Finally, due to the easy dissociation of O₂, the band structures of MoTe₂ and WTe₂ monolayer sheets show some robustness to oxidation. We anticipate MoTe₂- or WTe₂- based FET devices to be experimentally fabricated in near future.

Conclusions

Using first-principles calculations, we systematically investigated the oxidation behaviors of several semiconducting TMD monolayers (MoS₂, MoSe₂, MoTe₂, WS₂, WSe₂ and WTe₂) with single chalcogen vacancies and discussed the oxidation effects on the electronic properties. Without defects, all TMD monolayers are inert to O₂ gaseous molecules; whereas oxidation as an exothermic process mainly occurs at the chalcogen vacancies with adsorption energies of 1.8~3.9 eV. For all semiconducting TMD monolayers explored, chalcogen vacancy can induce some impurity states in the gap region, which can be removed by oxidation. After adsorption, O₂ molecule may further dissociate exothermically by overcoming an energy barrier. Due to small or even zero energy barrier, O₂ molecule would easily dissociate after adsorption on defective

MoTe₂ and WTe₂ monolayers at ambient conditions. For the other four oxidized TMD monolayers, dissociation of O_2 is more difficult. Moreover, it is difficult for oxygen adatoms to diffuse on TMD monolayers because of the large diffusion barriers (0.90~2.68 eV). Simulated STM images indicate that chemisorbed oxygen molecules and oxygen adatoms in the vacancy sites are invisible and that oxygen adatoms on the top of chalcogen atoms result in bright spots in the STM images.

Most importantly, the oxidation effects on the electronic properties of different TMD monolayers are different. Oxidation can badly affect the electronic properties of MoS_2 monolayer. Among MoX_2 and WX_2 systems, $MoTe_2$ and WTe_2 monolayers suffer least from oxidation due to the easy dissociation of O_2 , which is helpful to reduce the influence on the electronic properties due to oxidation. No doubt, our results give valuable atomic insight into oxidation of the TMD monolayers. These theoretical findings will not only help understand the oxidation effects on TMD-based 2D materials in current experiments but also provide useful guidance for designing microelectronic devices in the future.

AUTHOR INFORMATION

Corresponding Author

* Corresponding author. E-mail: zhaojj@dlut.edu.cn

ACKNOWLEDGMENT

This work was supported by the National Natural Science Foundation of China (11134005).

REFERENCES

- K. S. Novoselov, D. Jiang, F. Schedin, T. J. Booth, V. V. Khotkevich, S. V. Morozov and A. K. Geim. Proc. Natl. Acad. Sci. U. S. A. 2005, 102, 10451-10453
- J. N. Coleman, M. Lotya, A. O'Neill, S. D. Bergin, P. J. King, U. Khan, K. Young, A. Gaucher, S. De, R. J. Smith, I. V. Shvets, S. K. Arora, G. Stanton, H.-Y. Kim, K. Lee, G. T. Kim, G. S. Duesberg, T. Hallam, J. J. Boland, J. J. Wang, J. F. Donegan, J. C. Grunlan, G. Moriarty, A. Shmeliov, R. J. Nicholls, J. M. Perkins, E. M. Grieveson, K. Theuwissen, D. W. McComb, P. D. Nellist and V. Nicolosi. *Science* 2011, 331, 568-571
- 3. M. Xu, T. Liang, M. Shi and H. Chen. Chem. Rev. 2013, 113, 3766-3798
- Q. H. Wang, K. Kalantar-Zadeh, A. Kis, J. N. Coleman and M. S. Strano. *Nat. Nanotechnol.* 2012, 7, 699-712
- 5. M. Chhowalla, H. S. Shin, G. Eda, L.-J. Li, K. P. Loh and H. Zhang. Nat. Chem. 2013, 5, 263-275
- S. Z. Butler, S. M. Hollen, L. Cao, Y. Cui, J. A. Gupta, H. R. Gutiérrez, T. F. Heinz, S. S. Hong, J. Huang,
 A. F. Ismach, E. Johnston-Halperin, M. Kuno, V. V. Plashnitsa, R. D. Robinson, R. S. Ruoff, S. Salahuddin, J. Shan, L. Shi, M. G. Spencer, M. Terrones, W. Windl and J. E. Goldberger. *ACS Nano* 2013, 7, 2898-2926
- S. Balendhran, S. Walia, H. Nili, J. Z. Ou, S. Zhuiykov, R. B. Kaner, S. Sriram, M. Bhaskaran and K. Kalantar-zadeh. *Adv. Funct. Mater.* 2013, 23, 3952-3970
- 8. D. Y. Qiu, F. H. da Jornada and S. G. Louie. Phys. Rev. Lett. 2013, 111, 216805
- 9. H. Shi, H. Pan, Y.-W. Zhang and B. Yakobson. Phys. Rev. B 2013, 87, 155304
- 10. K. F. Mak, C. Lee, J. Hone, J. Shan and T. F. Heinz. Phys. Rev. Lett. 2010, 105, 136805
- 11. RadisavljevicB, RadenovicA, BrivioJ, GiacomettiV and KisA. Nat. Nanotechnol. 2011, 6, 147-150

- Z. Yin, H. Li, H. Li, L. Jiang, Y. Shi, Y. Sun, G. Lu, Q. Zhang, X. Chen and H. Zhang. ACS Nano 2011, 6, 74-80
- H. S. Lee, S. W. Min, Y. G. Chang, M. K. Park, T. Nam, H. Kim, J. H. Kim, S. Ryu and S. Im. *Nano Lett.* 2012, **12**, 3695-3700
- 14. B. Radisavljevic, M. B. Whitwick and A. Kis. Appl. Phys. Lett. 2012, 101, 043103
- H. Li, Z. Yin, Q. He, H. Li, X. Huang, G. Lu, D. W. Fam, A. I. Tok, Q. Zhang and H. Zhang. *Small* 2012, 8, 63-67
- J. Feng, L. Peng, C. Wu, X. Sun, S. Hu, C. Lin, J. Dai, J. Yang and Y. Xie. Adv. Mater. 2012, 24, 1969-1974
- F. K. Perkins, A. L. Friedman, E. Cobas, P. M. Campbell, G. G. Jernigan and B. T. Jonker. *Nano Lett.* 2013, 13, 668-673
- M. Buscema, M. Barkelid, V. Zwiller, H. S. van der Zant, G. A. Steele and A. Castellanos-Gomez. *Nano Lett.* 2013, 13, 358-363
- 19. O. Lopez-Sanchez, D. Lembke, M. Kayci, A. Radenovic and A. Kis. Nat. Nanotechnol. 2013, 8, 497-501
- 20. M. Bernardi, M. Palummo and J. C. Grossman. Nano Lett. 2013, 13, 3664-3670
- 21. Y. Ding, Y. Wang, J. Ni, L. Shi, S. Shi and W. Tang. *Physica B: Condensed Matter* 2011, **406**, 2254-2260
- 22. H. Qiu, L. Pan, Z. Yao, J. Li, Y. Shi and X. Wang. Appl. Phys. Lett. 2012, 100, 123104
- 23. W. Park, J. Park, J. Jang, H. Lee, H. Jeong, K. Cho, S. Hong and T. Lee. *Nanotechnology* 2013, 24, 095202

- 24. S. Tongay, J. Zhou, C. Ataca, J. Liu, J. S. Kang, T. S. Matthews, L. You, J. Li, J. C. Grossman and J. Wu. *Nano letters* 2013, **13**, 2831-2836
- 25. S. Tongay, J. Suh, C. Ataca, W. Fan, A. Luce, J. S. Kang, J. Liu, C. Ko, R. Raghunathanan, J. Zhou, F. Ogletree, J. Li, J. C. Grossman and J. Wu. *Sci Rep* 2013, **3**, 2657
- 26. G. Kresse and J. Furthmüller. Phys. Rev. B 1996, 54, 11169-11186
- 27. G. Kresse and D. Joubert. Phys. Rev. B 1999, 59, 1758-1775
- 28. J. P. Perdew, K. Burke and M. Ernzerhof. Phys. Rev. Lett. 1996, 77, 3865-3868
- 29. S. M. Davis and J. C. Carver. Applications of Surface Science 1984, 20, 193-198
- 30. W. Jaegermann and D. Schmeisser. Surf. Sci. 1986, 165, 143-160
- 31. J. Lince and P. Frantz. Tribol. Lett. 2001, 9, 211-218
- 32. J. S. Ha, H.-S. Roh, S.-J. Park, J.-Y. Yi and E.-H. Lee. Surf. Sci. 1994, 315, 62-68
- H.-P. Komsa, S. Kurasch, O. Lehtinen, U. Kaiser and A. V. Krasheninnikov. *Phys. Rev. B* 2013, 88, 035301
- 34. H. Qiu, T. Xu, Z. Wang, W. Ren, H. Nan, Z. Ni, Q. Chen, S. Yuan, F. Miao, F. Song, G. Long, Y. Shi, L. Sun, J. Wang and X. Wang. *Nat. Commun.* 2013, 4, 2642
- 35. W. Zhou, X. Zou, S. Najmaei, Z. Liu, Y. Shi, J. Kong, J. Lou, P. M. Ajayan, B. I. Yakobson and J. C. Idrobo. *Nano Lett.* 2013, **13**, 2615-2622
- H.-P. Komsa, J. Kotakoski, S. Kurasch, O. Lehtinen, U. Kaiser and A. V. Krasheninnikov. *Phys. Rev. Lett.* 2012, **109**, 035503
- 37. G. Henkelman, A. Arnaldsson and H. Jónsson. Comput. Mater. Sci. 2006, 36, 354-360

- 38. J. Donohue. The Structures of the Elements; John Wiley & Sons, Inc. Press: New York, 1974.
- 39. Q. Yue, Z. Shao, S. Chang and J. Li Nanoscale Res. Lett. 2013, 8, 425
- 40. S. Grimme. J. Comput. Chem. 2006, 27, 1787-1799
- 41. T. s. Bučko, J. r. Hafner, S. b. Lebègue and J. n. G. Ángyán. J. Phys. Chem. A 2010, 114, 11814-11824
- 42. J. B. Parise, E. M. McCarron Iii and A. W. Sleight. Mater. Res. Bull. 1987, 22, 803-811
- 43. G. Mills, H. Jónsson and G. K. Schenter. Surf. Sci. 1995, 324, 305-337
- 44. J. He, K. Wu, R. Sa, Q. Li and Y. Wei. Appl. Phys. Lett. 2010, 96, 082504
- 45. Y. Ma, Y. Dai, M. Guo, C. Niu, J. Lu and B. Huang. Phys. Chem. Chem. Phys. 2011, 13, 15546-15553
- 46. J. Fuhr, A. Saúl and J. Sofo. Phys. Rev. Lett. 2004, 92, 026802
- 47. C. Kittel. Introduction to Solid State Physics; John Wiley & Sons, Inc. Press: USA, 2005.
- 48. W. S. Yun, S. W. Han, S. C. Hong, I. G. Kim and J. D. Lee. Phys. Rev. B 2012, 85, 033305
- 49. D. Braga, I. Gutierrez Lezama, H. Berger and A. F. Morpurgo. Nano Lett. 2012, 12, 5218-5223
- 50. H. Fang, S. Chuang, T. C. Chang, K. Takei, T. Takahashi and A. Javey. Nano Lett. 2012, 12, 3788-3792