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## ARTICLE

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# Separation of the Molecular Motion from Different Components or Phases using Projection MovingWindow 2D Correlation FTIR Spectroscopy for Multiphase and Multicomponent Polymers 

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#### Abstract

This study developed a new analytical method called projection moving-window 2D correlation FTIR spectroscopy (Proj-MW2D) to separate the molecular motion of groups generated from different components or phases for multiphase and multicomponent polymers. The specific implementation steps for Proj-MW2D were enumerated after the theoretical derivation and algorithm research. Two types of two-component blends, poly(Llactide)/poly(butylene succinate) and monodisperse polystyrene/monodisperse poly(ethylene-co-1-butene), were employed to validate the concept of separating the molecular motion of groups. Results showed that Proj-MW2D FTIR correlation technique successfully separated the molecular motion of the specific functional groups. Although MW2D and PCMW2D have the capacity to determine the multiple transitions of polymers, they cannot identify the origin of correlation intensity peaks without the help of other characterization methods. Proj-MW2D allows researchers to study the mechanism of the complex transition process for multiphase and multicomponent polymer systems. This method can be easily extended to three- or four-component polymers and to other spectra (e.g., Raman, X-ray, and UV).


## 1. Introduction

Polymers have a complex transition behavior. Since the discovery of polymers in organic materials chemistry, scientists have successfully developed numerous techniques to characterize polymer transition behavior. ${ }^{1}$ The advancements in characterization methods have expanded our knowledge on polymer materials. Infrared (IR) spectroscopy is a classic method for polymers. It is a vibrational spectrum that reflects the dipole vibrations of specific bonds in polymers. ${ }^{2,}{ }^{3}$ Therefore, IR spectroscopy provides molecular information at the functional group level.

Generalized 2D correlation vibrational spectroscopy, which was inspired from 2D NMR, was proposed by Noda in $1993 .{ }^{4}$ This technique uses Fourier transform and convolution to apply correlation analysis for in situ IR with an external perturbation. ${ }^{5}$ Peoples can employ it to study the interaction mechanism of functional groups in detail during polymer transitions. Given its practicality, generalized 2D correlation spectroscopy has rapidly become a new popular method of characterization. ${ }^{6-11}$ However, 2D correlation spectroscopy cannot establish the relationship of correlation intensity peaks between spectral
variables (e.g., wavenumber, Raman shift) and external perturbation variables (e.g., temperature, pressure). Thomas and Richardson ${ }^{12}$ proposed moving-window 2D correlation spectroscopy (MW2D) in 2000. The correlation intensity peaks of spectral variables directly associate with those of external perturbation variables in this technique; therefore, the transition points of polymers and the transition range can be precisely determined. ${ }^{13-19}$ Morita ${ }^{20}$ introduced perturbation-correlation moving-window 2D correlation spectroscopy (PCMW2D) in 2006. Summary findings of scientists over many years found that the combination of MW2D or PCMW2D, and generalized 2D correlation spectroscopy was the best way. ${ }^{17,18,21-26}$ In general, MW2D or PCMW2D is employed to determine the transition point and the transition range of polymers; then, generalized 2D spectroscopy is performed to study the mechanism of functional groups at a specific transition. ${ }^{17,18,21-}$
${ }^{26}$ Compared with generalized 2D spectroscopy, 2D correlation spectroscopy with moving-window (MW2D or PCMW2D) has inherent advantages on studying the multiple transitions of polymers.
Generalized 2D correlation IR spectra are sometimes complex in practical applications. Specifically, the overlapping
of correlation peaks is serious or looks "crowded." Explaining these spectra via Noda's rules is difficult because of the interference of correlation intensity peaks. ${ }^{4,5}$ This phenomenon is obvious in multiphase and multicomponent polymer systems. In some block copolymers (SBS, SIS, and SEPS) and polymer blends, the correlation intensity peaks of minor components tend to be masked when main components dominate the transitions. Therefore, caution must be implemented when using generalized 2D correlation FTIR spectroscopy to study multiphase and multicomponent polymer systems. Considering this problem, Noda ${ }^{27}$ proposed the projection technique for generalized 2D correlation spectroscopy in 2010. This method was extended from double 2D correlation analysis. ${ }^{28}$ The projection technique was specifically developed for complex systems. This technique simplifies the "crowded" 2D correlation spectra by distinguishing the corresponding correlation intensity from different components. The concept of projection has aroused great interest among scientists. ${ }^{29,}{ }^{30}$ However, generalized 2D correlation spectroscopy combined with projection still has some limitations on multiphase and multicomponent polymer systems because most polymers show multiple transitions under an external perturbation (e.g., temperature). Although MW2D or PCMW2D correlation spectroscopy has inherent capacity to determine the multiple transitions of polymers, it cannot identify the origin of correlation intensity peaks without the help of other characterization methods. To overcome these difficulties, this study developed a method called projection moving-window 2D correlation spectroscopy (Proj-MW2D) to separate the molecular motion of groups generated from different components in multicomponent polymers. The new approach not only has a capacity to precisely specify multiple transition points of polymers, but also can separate the molecular motion of groups. It will be particularly suitable to the study of the molecular movement mechanism of a specific transition for polymer systems.


Figure 1. Schematic of the proposed concept of separating the molecular motion of groups from different components for multicomponent polymers.

In this paper, through theoretical derivation, algorithm research, and infrared spectroscopy experiments, the projection technique was successfully integrated into MW2D. Figure 1 illustrates the proposed concept of separating the molecular motion of groups from different components. Proj-MW2D is a promising analytical method. This method possibly has potential outstanding and special applications in multiphase and multicomponent polymers.

## 2. Method

### 2.1. Spectral Data Matrix

The theoretical derivation of the method is outlined below.
Proj-MW2D needs to divide the entire spectral data matrix as a series of sub-matrices. Each sub-matrix is called as a window. $A$ is supposed to be an $M \times N$ matrix of spectral data with $M$ rows and $N$ columns. $v$ and $I$ represent the spectral variables (e.g., wavenumber) and external perturbation variables (e.g., temperature), respectively. Each row of $A$ corresponds to the discrete data points of a spectrum.

$$
A=\left(\begin{array}{c}
y\left(v, I_{1}\right)  \tag{1}\\
\vdots \\
y\left(v, I_{j}\right) \\
\vdots \\
y\left(v, I_{M}\right)
\end{array}\right)
$$

The sub-matrix $A_{j}$ is extracted from the main matrix $A$. At this point, the serial number of the perturbation variable $I$ ranges from $j-m$ to $j+m$. Obviously, the sub-matrix $A_{j}$ has $2 m+1$ rows, whose column number is equal to the main matrix $A$. $2 m+1$ is the window size

The sub-matrix $A_{j}$ can be expressed as

$$
A_{j}=\left(\begin{array}{c}
y\left(v, I_{j-m}\right)  \tag{2}\\
\vdots \\
y\left(v, I_{j}\right) \\
\vdots \\
y\left(v, I_{j+m}\right)
\end{array}\right)
$$

In general, the reference spectrum and the dynamic spectrum of the $j$ th sub-matrix are calculated by the following formulas.

$$
\begin{gather*}
\bar{y}(v)=\frac{1}{2 m+1} \sum_{J=j-m}^{j+m} y\left(v, I_{J}\right)  \tag{3}\\
\tilde{y}\left(v, I_{J}\right)=y\left(v, I_{J}\right)-\bar{y}(v) \tag{4}
\end{gather*}
$$

where $J$ is a row number within a sub-matrix
The mean-centered $j$ th sub-matrix can be described as follows.

$$
\widetilde{A}_{j}=\frac{1}{\sqrt{2 m}}\left(\begin{array}{c}
\widetilde{y}\left(v, I_{j-m}\right)  \tag{5}\\
\vdots \\
\widetilde{y}\left(v, I_{j}\right) \\
\vdots \\
\widetilde{y}\left(v, I_{j+m}\right)
\end{array}\right)
$$

### 2.2. Projection Matrix

A matrix $Y$ has $2 m+1$ rows and $k$ columns, and the rows $(2 m+1)$ of $Y$ should be the same as those of $\tilde{A_{j}}$. The matrix $Y$ reduces to a single vector $y$ when $Y$ has only one column.

The singular value decomposition (SVD) of matrix $Y$ is defined as follows.

$$
\begin{equation*}
Y=U_{Y} S_{Y} V_{Y}^{T}+E_{Y} \tag{6}
\end{equation*}
$$

SVD is a common method of extracting important features from a mathematical concept. ${ }^{31}$ It can be applied to an arbitrary matrix. SVD is attributed to the important field of data mining. Many applications can be related to SVD. These applications include PCA with the feature extraction, the data compression algorithm (represented by the image compression), and the latent semantic indexing for search engine.
$U_{Y}$ is a matrix of $(2 m+1) \times(2 m+1) . V_{Y}^{T}$ is a square matrix of $k \times k . S_{Y}$ is a $(2 m+1) \times k$ matrix, and its diagonal elements of $\sigma_{l}$, $\sigma_{2}, \ldots, \sigma_{k}$ are the singular values of the matrix $Y . E_{Y}$ is the residual matrix. Generally, it is small enough to be negligible. Here, we give an index $r$, which makes the value of $\left(\sigma_{1}+\sigma_{2}+\ldots+\sigma_{r}\right) /\left(\sigma_{1}+\sigma_{2}+\ldots+\sigma_{r}+\ldots+\sigma_{k}\right)$ greater than or equal to $0.9(90 \%)$, and the $1, \ldots,(2 m+1)$ rows and the $1, \ldots, r$ columns of $U_{Y}$ are used for the subsequent calculations.

Then, the projection matrix $R_{Y}$ from $Y$ can be calculated as:

$$
\begin{equation*}
R_{Y}=U_{Y} U_{Y}^{T} \tag{7}
\end{equation*}
$$

$R_{Y}$ is a $(2 m+1) \times(2 m+1)$ matrix. The used $(2 m+1) \times r$ part of $U_{Y}$ avoids $R_{Y}$ transform to an identity matrix. The corresponding null-space projection matrix is defined as:

$$
\begin{equation*}
\left(I-R_{Y}\right)=I-U_{Y} U_{Y}^{T} \tag{8}
\end{equation*}
$$

where $I$ is the $(2 m+1) \times(2 m+1)$ identity matrix, Both $R_{Y}$ and $I$ $R_{Y}$ are the symmetric matrix.

### 2.3. Generalized Projection Transformation

$\tilde{A}_{j}$ matrix can be transformed into a new matrix of $\tilde{A}_{j}^{P}$ with the projection according to the following formula:

$$
\begin{equation*}
\widetilde{A}_{j}^{P}=R_{Y} \widetilde{A}_{j} \tag{9}
\end{equation*}
$$

The corresponding null-space projection matrix is:

$$
\begin{equation*}
\widetilde{A}_{j}^{N}=\widetilde{A}_{j}-\widetilde{A}_{j}^{P} \tag{10}
\end{equation*}
$$

In the $j$ th sub-matrix, the projection synchronous and projection asynchronous 2D correlation spectra can be calculated as:

$$
\begin{gather*}
\Phi_{j}^{P}=\left(\widetilde{A}_{j}^{P}\right)^{T} \cdot \widetilde{A}_{j}^{P}  \tag{11}\\
\Psi_{j}^{P}=\left(\widetilde{A}_{j}^{P}\right)^{T} \cdot N \cdot \widetilde{A}_{j}^{P} \tag{12}
\end{gather*}
$$

The corresponding null-space projection synchronous and null-space projection asynchronous 2D correlation spectra are calculated as:

$$
\begin{gather*}
\Phi_{j}^{N}=\left(\tilde{A}_{j}^{N}\right)^{T} \cdot \tilde{A}_{j}^{N}  \tag{13}\\
\Psi_{j}^{N}=\left(\tilde{A}_{j}^{N}\right)^{T} \cdot N \cdot \widetilde{A}_{j}^{N} \tag{14}
\end{gather*}
$$

where $N$ is the Hilbert-Noda transformation matrix, ${ }^{5}$ which is a $(2 m+1) \times(2 m+1)$ matrix.

### 2.4. Positive Projection Transformation

According to Noda's proposal, ${ }^{27}$ the loading matrix $L$ of $\tilde{A}_{j}{ }^{P}$ is defined as:

$$
\begin{equation*}
L=\widetilde{A}_{j}^{T} U_{Y} \tag{15}
\end{equation*}
$$

Then, $\tilde{A}_{j}{ }^{P}$ can be expressed as:

$$
\begin{equation*}
\widetilde{A}_{j}^{P}=U_{Y} L^{T} \tag{16}
\end{equation*}
$$

In this study, we introduce a positive loading matrix $L_{+}$, which comes from $L$ by data processing. $L_{+}$retains the elements of $\geq 0$ in $L$, and the elements of $<0$ are replaced by 0 . After positive projection transformation, we obtain a new projection matrix $\tilde{A}_{j}^{P+} . \tilde{A}_{j}^{P+}$ only reserves the columns in $\tilde{A}_{j}^{P}$ that have the same change (or orientation) direction as in matrix $Y$.

$$
\begin{equation*}
\widetilde{A}_{j}^{P+}=U_{Y} L_{+}^{T} \tag{17}
\end{equation*}
$$

The corresponding null-space projection matrix is:

$$
\begin{equation*}
\widetilde{A}_{j}^{N+}=\widetilde{A}_{j}-\widetilde{A}_{j}^{P+} \tag{18}
\end{equation*}
$$

Then, the projection synchronous and projection asynchronous 2D correlation spectra in the $j$ th sub-matrix are calculated as follows:

$$
\begin{gather*}
\Phi_{j}^{P+}=\left(\widetilde{A}_{j}^{P+}\right)^{T} \cdot \tilde{A}_{j}^{P+}  \tag{19}\\
\Psi_{j}^{P+}=\left(\widetilde{A}_{j}^{P+}\right)^{T} \cdot N \cdot \widetilde{A}_{j}^{P+} \tag{20}
\end{gather*}
$$

The corresponding null-space projection synchronous and null-space projection asynchronous 2D correlation spectra are calculated as follows:

$$
\begin{gather*}
\Phi_{j}^{N+}=\left(\widetilde{A}_{j}^{N+}\right)^{T} \cdot \widetilde{A}_{j}^{N+}  \tag{21}\\
\Psi_{j}^{N+}=\left(\widetilde{A}_{j}^{N+}\right)^{T} \cdot N \cdot \widetilde{A}_{j}^{N+} \tag{22}
\end{gather*}
$$

### 2.5. Moving-window

A series of sub-matrices are obtained by sliding the window position from $j=1+m$ to $j=M-m$. The calculations are repeatedly performed using the above equations for each sub-matrix (window). Then, an overlay operation that is identical to the conventional MW2D algorithm is employed to gain a ProjMW2D correlation spectrum ${ }^{13,17}$ and a null-space Proj-MW2D correlation spectrum. For convenient explanation and calculations, our study focuses on Proj-MW2D and null-space Proj-MW2D based on auto-correlation.

### 2.6. Multiple Projection for Three to Four Components (or Phases)

The algorithm mentioned above is only applicable to polymers with two components (or phases). The multiple projection operation should be introduced to investigate threeto four-component (phases) polymers.
$\tilde{A}_{j}^{N}$ or $\tilde{A}_{j}^{N+}$ is specified as a new spectral data matrix. We select matrix $Y_{l}$ as another projection to repeat equations (6) to (10) or (15) to (18). Then, the spectral data matrix $\tilde{A}_{j}^{P I}$ or $\tilde{A}_{j}^{P I+}$ with the second projection and the spectral data matrix $\tilde{A}_{j}^{N I}$ or $\tilde{A}_{j}^{N I+}$ with null-space second projection are obtained.

$$
\begin{gather*}
\widetilde{A}_{j}^{P 1}=R_{Y 1} \widetilde{A}_{j}^{N}  \tag{23}\\
\widetilde{A}_{j}^{N 1}=\widetilde{A}_{j}^{N}-\widetilde{A}_{j}^{P 1}  \tag{24}\\
\widetilde{A}_{j}^{P 1+}=U_{Y 1} L_{1+}^{T}  \tag{25}\\
\widetilde{A}_{j}^{N 1+}=\widetilde{A}_{j}^{N+}-\widetilde{A}_{j}^{P 1+} \tag{26}
\end{gather*}
$$

This algorithm allows Proj-MW2D to analyze polymer systems with three components (or phases). Similarly, using $\tilde{A}_{j}^{P I}$ or $\tilde{A}_{j}^{P I+}$ as a new spectral data matrix for the third projection, we can use Proj-MW2D to analyze polymer systems with four components (or phases).

### 2.7. Selection of the Projection Matrix $\boldsymbol{Y}$

The projection matrix $Y$ influences the final results. The matrix $Y$ can be selected from the $j$ th sub-matrix $\tilde{A}_{j}$, which has the mean-centered processing. Generally, the matrix $Y$ contains a few columns of $\tilde{A}_{j}$. In the actual spectrum (e.g., IR), these columns contain some characteristic bands. It degenerates into a vector when $Y$ has only one column of $\tilde{A}_{j}$, which corresponds to a characteristic spectral peak. Selecting the projection matrix $Y$ from the spectral data matrix $\tilde{A}_{j}$ can satisfy most of our needs. The projection matrix $Y$ can also be chosen from other sources, such as from another experiment of IR or Raman spectra.

External perturbations (such as temperature or concentration) can be also selected as the projection matrix $Y$. In some cases, specific mathematical functions can be used as a projection matrix for model-based analyses. In Proj-MW2D correlation spectroscopy, the construction method of the projection matrix $Y$ is similar to Noda's proposal.

### 2.8. Implementation Steps for Polymers

Polymer blends with two components (or phases), such as $\mathrm{A}+\mathrm{B}$, are used as an example to show the implementation steps of Proj-MW2D in separating the molecular motion of groups. The basic steps are outlined below:

1) In situ spectral data of polymer blend samples are collected in the presence of an external perturbation (temperature, concentration, etc.). Collecting the maximum number of spectra ensured by the signal-noise ratio and allowed by the spectrometer collection rate is recommended to accurately describe the physical or chemical processes of the samples.
2) The projection matrix $Y$ is constructed. In general, $Y$ is selected from the interest characteristic bands in the in situ spectral data. For multiphase and multicomponent polymers, $Y$ is often chosen from the most characteristic bands that can represent a phase or a component. This characteristic bands can distinguish the interest component from the other components (or phases). For twocomponent polymer blends, we choose certain characteristic bands of component A (denoted as $a$ ) to construct the projection matrix.
3) Proj-MW2D correlation spectra are calculated using the positive projection transformation algorithm. The algorithm can be compiled into program using C and $\mathrm{C}++$ computer language or using Matlab directly.
4) Proj-MW2D correlation spectrum and null-space ProjMW2D correlation spectrum are obtained. For twocomponent polymer blends, the response of correlation intensities from component A appear in Proj-MW2D correlation spectra, and the response of correlation intensity from component B presents in null-space ProjMW2D correlation spectra. The reason for this result is that the response with the same variation feature as the characteristic bands (a) of component A is retained by the positive projection algorithm.
5) The transition range of components $A$ and $B$ is determined according to the response of the correlation intensity. The individual physical and chemical transitions of A or B and collaborative transitions of $\mathrm{A}+\mathrm{B}$ involved can be identified. Then, the mechanism and bands' assignment can be further analyzed.

## 3. Experimental Section

In this study, two types of two-component polymer blends were prepared using different blending methods. The molecular motion of groups generated from different components was separated using Proj-MW2D correlation analysis.

### 3.1. Materials

Poly(L-lactide) (4032D) was provided by NatureWorks LLC, USA. Poly(butylene succinate) (Bionolle 1001 MD ) was provided by Showa Highpolymer Co., Ltd., Japan. Monodisperse polystyrene (PS) and monodisperse poly(ethylene-co-1-butene) (EB, also known as fully hydrogenated polybutadiene) synthesized via anionic polymerization were supplied by Baling Petrochemical Industry Co., Ltd. (Sinopec). Cyclohexane with $99.5 \%$ (AR) purity was produced by Chemical Reagent Union Co., Ltd. The weight average molecular weight of monodisperse PS was 10,000 $\mathrm{g} / \mathrm{mol}$, and its monodisperse coefficient was 1.04 (GPC). The weight average molecular weight of monodisperse EB was $120,000 \mathrm{~g} / \mathrm{mol}$, and its monodisperse coefficient was 1.05 (GPC). The molar content of 1-butene in EB was $29 \%$ (NMR). Poly(butylene succinate) (PBS), poly(L-lactide) (PLLA), monodisperse polystyrene (PS), and monodisperse poly(ethylene-co-1-butene) (EB) were directly used without further purification. The molecular structures of these materials are shown in Scheme 1.


PBS


PLLA


EB


PS

Scheme 1. Molecular structures of PBS, PLLA, EB, and PS.

### 3.2. Preparation

### 3.2.1. PLLA/PBS blend

PBS and PLLA pellets were vacuum-dried at $70^{\circ} \mathrm{C}$ for 24 h . A PLLA/PBS blend was prepared by a twin-screw extruder through melt blending. The weight ratio of PLLA and PBS was 50:50. The temperature of melt blending ranged from $175^{\circ} \mathrm{C}$ to $185^{\circ} \mathrm{C}$, and the screw speed was 180 rpm .

### 3.2.2. PS/EB Blend

A PS/EB blend was prepared by solution blending. Monodisperse PS ( 0.714 g ) and monodisperse EB (4.286 g) were simultaneously dissolved in 250 ml of cyclohexane. After becoming transparent, the solution was stored under an air-tight seal for 4 d before use.

### 3.3. Temperature-dependent FTIR Spectroscopy

A Nicolet iS10 FTIR spectrometer equipped with a deuterated triglycine sulfate detector was used. The FTIR spectral resolution was $4 \mathrm{~cm}^{-1}$, and 20 scans were co-added for each FTIR spectrum.

### 3.3.1. PLLA/PBS

A pellet $(0.5 \mathrm{~g})$ of the PLLA/PBS blend was placed between two polyethylene terephthalate (PET) sheets. The film was directly prepared by heat-pressing the PLLA/PBS blend sandwiched between two PET sheets at $185^{\circ} \mathrm{C}$ in a hot press and then allowing the blend to naturally cool at room temperature. The heat-pressing time was 2 min , and the pressure was 20.0 MPa . Then, the film was cut into a circular shape ( $\Phi 1.3 \mathrm{~cm}$ ). The film was sandwiched between two KBr windows to prevent high-temperature flow and then placed into a homemade in situ pool (programmable heating device). The sample film was heated from $20^{\circ} \mathrm{C}$ to $120^{\circ} \mathrm{C}$ at $5^{\circ} \mathrm{C} / \mathrm{min}$ and then was stored at $120^{\circ} \mathrm{C}$ for 3 min . The molten sample was cooled from $120^{\circ} \mathrm{C}$ to $50^{\circ} \mathrm{C}$ at $5^{\circ} \mathrm{C} / \mathrm{min}$. Meanwhile, the FTIR spectra were collected from $120^{\circ} \mathrm{C}$ to $50^{\circ} \mathrm{C}$. Thirty-three FTIR spectra were collected upon cooling. The sample was protected by dried high-purity nitrogen gas ( $300 \mathrm{ml} / \mathrm{min}$ ) during the measurement.

### 3.3.2. PS/EB

A film of the PS/EB blend was spread on one side of a KBr disk by solvent casting from the solution prepared in Section 3.2.2. The film was dried in a vacuum oven at $90{ }^{\circ} \mathrm{C}$ for 120 min at -0.08 MPa . The sample film was sandwiched between two KBr disks and then placed into a homemade in situ pool. The sample was heated from $30{ }^{\circ} \mathrm{C}$ to $175{ }^{\circ} \mathrm{C}$ at $5^{\circ} \mathrm{C} / \mathrm{min}$. The FTIR spectra were collected from $30^{\circ} \mathrm{C}$ to 175 ${ }^{\circ} \mathrm{C}$ at approximately $2.1{ }^{\circ} \mathrm{C}$ increments. The number of the spectra collected was 68 . The sample was also protected by dried high-purity nitrogen gas.

### 3.4. Proj-MW2D Correlation Analysis

Proj-MW2D correlation FTIR spectra were processed and calculated using 2DCS 6.1 software, developed by one of the authors using Visual C++. Before the analysis, the IR spectra were treated using a linear baseline correction. The window size was chosen as $2 m+1=11$ to produce high-quality ProjMW2D FTIR spectra and null-space Proj-MW2D FTIR spectra. The $1 \%$ correlation intensity of the Proj-MW2D and null-space Proj-MW2D correlation spectra was regarded as noise and was cut off. In the 2D spectra, red and blue areas represent positive and negative correlation intensities, respectively.


Figure 2. Temperature-dependent FTIR spectra of PLLA/PBS in the $1850 \mathrm{~cm}^{-1}$ to $720 \mathrm{~cm}^{-1}$ region upon cooling ( $120^{\circ} \mathrm{C}$ to $50^{\circ} \mathrm{C}$ ). The region from $1635 \mathrm{~cm}^{-1}$ to $1520 \mathrm{~cm}^{-1}$ is truncated. The bands within $1850 \mathrm{~cm}^{-1}$ to $1635 \mathrm{~cm}^{-1}$, which is the overlapping of PBS and PLLA, are assigned to the $\mathrm{C}=\mathrm{O}$ stretching vibration. The bands at $806 \mathrm{~cm}^{-1}$ is the $\mathrm{CH}_{2}$ in-plane bending in the $\mathrm{OC}\left(\mathrm{CH}_{2}\right)_{2} \mathrm{CO}$ of the PBS component, and it is chosen as the projection vector. The FTIR spectra at 806 $\mathrm{cm}^{-1}$ are enlarged for clarity. For clear illustration, not all the collected FTIR spectra are shown.

## 4. Results and Discussion

### 4.1. Separation of the Molecular Motion of $\mathbf{C = O}$ Groups from the Crystalline Phase of the PBS Component during PLLA/PBS Crystallization upon Cooling

### 4.1.1. Temperature-dependent FTIR of PLLA/PBS Crystallization

PBS and PLLA are completely biodegradable plastics. ${ }^{32-35}$ The temperature-dependent FTIR spectra of PLLA/PBS in the $1850 \mathrm{~cm}^{-1}$ to $720 \mathrm{~cm}^{-1}$ region upon cooling are shown in Figure 2. The region within $1635 \mathrm{~cm}^{-1}$ to $1520 \mathrm{~cm}^{-1}$ is an entirely horizontal baseline; thus, the spectra in this range are truncated. For the sake of clarity, Figure 2 does not show all the collected spectra. The bands' assignments are listed in Table 1. This study focuses on the bands within the $1850^{-1}$ to $1635 \mathrm{~cm}^{-1}$ region. These bands, which are the overlapping of PBS and PLLA, are assigned to the $\mathrm{C}=\mathrm{O}$ stretching vibration of PBS and PLLA. ${ }^{36-41}$ Some minor changes are observed within $1850 \mathrm{~cm}^{-1}$ to $1635 \mathrm{~cm}^{-1}$ when the temperature decreases from $120{ }^{\circ} \mathrm{C}$ to $50{ }^{\circ} \mathrm{C}$. The peak intensity at $1759 \mathrm{~cm}^{-1}$ gradually increases, whereas the bands at $1714 \mathrm{~cm}^{-1}$ slightly move to the lower wavenumber. These results show that temperaturedependent FTIR cannot separate the molecular motion from the $\mathrm{C}=\mathrm{O}$ groups in the PBS crystalline phase. Our target is to use Proj-MW2D to separate the $\mathrm{C}=\mathrm{O}$ motion generated from the PBS crystalline phase during PLLA/PBS crystallization.

Table 1. Band Assignments of PLLA, PBS, PS, and EB

| $\begin{aligned} & \text { Wavenumber } \\ & \left(\mathrm{cm}^{-1}\right) \end{aligned}$ | Assignment |  |
| :---: | :---: | :---: |
|  | PLLA | PBS |
| 1759 | $\mathrm{C}=\mathrm{O}$ stretching vibration | -- |
| 1714 | -- | $\mathrm{C}=\mathrm{O}$ stretching vibration |
| $\begin{aligned} & 955 \text { (960- } \\ & 930) \end{aligned}$ | $\mathrm{CH}_{3}$ rocking+C-C stretching | $\mathrm{C}-\mathrm{O}$ symmetric stretching |
| $\begin{gathered} 919(924- \\ 920) \end{gathered}$ | $10^{3}$ helical chain skeletal deformation | $\mathrm{CH}_{2}$ skeletal deformation |
| 806 | -- | $\mathrm{CH}_{2}$ in-plane bending in $\mathrm{OC}\left(\mathrm{CH}_{2}\right)_{2} \mathrm{CO}$ |
|  | PS | EB |
| 2925 | $\mathrm{CH}_{2}$ antisymmetric stretching vibration of the main chain | $\mathrm{CH}_{2}$ antisymmetric stretching vibration |
| 2856 | $\mathrm{CH}_{2}$ symmetric stretching vibration of the main chain | $\mathrm{CH}_{2}$ symmetric stretching vibration |
| 1462 | -- | $\mathrm{CH}_{2}$ bending vibration |
| 1453 | $\mathrm{CH}_{2}$ bending vibration of the main chain | -- |
| 757 | $=\mathrm{CH}$ out-of-plane bending vibration of benzene rings | -- |
| 720 | -- | $\mathrm{CH}_{2}$ rocking of the ethylene units |
| 698 | $=\mathrm{CH}$ out-of-plane bending vibration of benzene rings | -- |

### 4.1.2. Selection of Projection Vector

Determining the projection matrix (or vector) is crucial. Because we want to sepaerate the molecular motion of $\mathrm{C}=\mathrm{O}$ groups from PBS crystalline phase, the projection matrix (or vector) should be the characteristic bands of PBS and the bands can accurately describe the functional group movements of the PBS crystalline phase during crystallization. In addition, the bands should not overlap with any bands of the PLLA component. In Figure 2, also refer to the literature, ${ }^{38,42}$ we find three characteristic bands of PBS at 955,919 , and $806 \mathrm{~cm}^{-1}$. The bands at $955 \mathrm{~cm}^{-1}$ is assigned to the $\mathrm{C}-\mathrm{O}$ symmetric stretching of PBS. ${ }^{42}$ The bands at $919 \mathrm{~cm}^{-1}$ is attributed to the $\mathrm{CH}_{2}$ skeletal deformation of PBS, ${ }^{38,42}$ and that at $806 \mathrm{~cm}^{-1}$ is attributed to the $\mathrm{CH}_{2}$ in-plane bending in $\mathrm{OC}\left(\mathrm{CH}_{2}\right)_{2} \mathrm{CO} .{ }^{42}$ These three bands can be judged as the crystallization bands of PBS because their intensities gradually increase as the temperature decreases from $120^{\circ} \mathrm{C}$ to $50^{\circ} \mathrm{C}$. The crystallization behavior of PBS in PLLA/PBS is clearly described by 955,919 , and 806 $\mathrm{cm}^{-1}$ upon cooling. However, according to the literature, ${ }^{43,} 44$ the bands in the $960 \mathrm{~cm}^{-1}$ to $930 \mathrm{~cm}^{-1}$ and $924^{-1}$ to $920 \mathrm{~cm}^{-1}$ regions are also assigned to the $\mathrm{CH}_{3}$ rocking $+\mathrm{C}-\mathrm{C}$ stretching and $10^{3}$ helical chain skeletal deformation of the PLLA component, respectively. Therefore, in Figure 2, only $806 \mathrm{~cm}^{-1}$ is fully attributed to PBS. The bands at $806 \mathrm{~cm}^{-1}$ does not overlap with any bands of PLLA, and it accurately describes the crystallization behavior of PBS. The molecular structure of PBS also shows that the corresponding $\mathrm{CH}_{2}$ groups of $806 \mathrm{~cm}^{-1}$ directly connect to the $\mathrm{C}=\mathrm{O}$ groups, thereby reflecting the real $\mathrm{C}=\mathrm{O}$ movements of PBS. In this study, the bands at $806 \mathrm{~cm}^{-1}$ is certainly chosen as the projection. To simplify the calculations, we only select the spectral intensity at $806 \mathrm{~cm}^{-1}$ as the
projection. Thus, the projection vector (not a matrix) is the spectral intensity at $806 \mathrm{~cm}^{-1}$.


Figure 3. Key: (a) Proj-MW2D correlation FTIR spectrum of PLLA/PBS ( $110^{\circ} \mathrm{C}$ to $60^{\circ} \mathrm{C}$ ); (b) null-space Proj-MW2D correlation FTIR spectrum of PLLA/PBS ( $110{ }^{\circ} \mathrm{C}$ to 60 ${ }^{\circ} \mathrm{C}$ ); (c) conventional MW2D correlation FTIR spectrum ( $110^{\circ} \mathrm{C}$ to $60^{\circ} \mathrm{C}$ ). The positive projection transformation algorithm is used for the calculations. The molecular motion of the $\mathrm{C}=\mathrm{O}$ groups involved in the crystallization of the PBS component is retained in the Proj-MW2D correlation spectrum, and the response temperature is approximately $78^{\circ} \mathrm{C}$. The $\mathrm{C}=\mathrm{O}$ motion of the PLLA component ( $98^{\circ} \mathrm{C}$ ) appears in the null-space Proj-MW2D correlation spectrum.


Figure 4. Key: (a) synchronous PCMW2D correlation FTIR spectrum of PLLA/PBS ( $110^{\circ} \mathrm{C}$ to $60^{\circ} \mathrm{C}$ ) upon cooling; (b) asynchronous PCMW2D correlation FTIR spectrum of PLLA/PBS ( $110^{\circ} \mathrm{C}$ to $60^{\circ} \mathrm{C}$ ) upon cooling

### 4.1.3. Proj-MW2D Analysis of the Molecular Motion of the $\mathbf{C}=\mathbf{O}$ Groups

Figure 3 is the Proj-MW2D correlation FTIR spectra of PLLA/PBS upon cooling, and the spectral intensity at 806 $\mathrm{cm}^{-1}$ is the projection vector. Figures 3(a) and 3(b) show the Proj-MW2D correlation spectrum and null-space Proj-MW2D correlation spectrum, respectively.

The molecular motion of the $\mathrm{C}=\mathrm{O}$ groups with the same motion feature as the projection vector at $806 \mathrm{~cm}^{-1}$ is retained in the Proj-MW2D correlation spectrum [Figure 3(a)]. Results
show that the correlation intensity appears at approximately $1714 \mathrm{~cm}^{-1}$, which represents the $\mathrm{C}=\mathrm{O}$ motion involved in PBS crystallization. The response temperature is approximately 78 ${ }^{\circ} \mathrm{C}$. This temperature is the crystallization temperature of PBS in PLLA/PBS. In the null-space Proj-MW2D correlation spectrum [Figure 3(a)], the $\mathrm{C}=\mathrm{O}$ motion that is completely orthogonal (not identical) with the projection vector at 806 $\mathrm{cm}^{-1}$ is observed. Figure 3(b) represents the molecular motion of the $\mathrm{C}=\mathrm{O}$ groups of PLLA because $806 \mathrm{~cm}^{-1}$ is attributed to the PBS component. In Figure 3(b), the corresponding correlation intensity is observed at around $1759 \mathrm{~cm}^{-1}$, and the
response temperature is $98{ }^{\circ} \mathrm{C}$. Therefore, the crystallization temperature of PLLA is $98{ }^{\circ} \mathrm{C}$ under the experimental conditions described in the present study.

Aside from separatting the molecular motion of functional groups from the different components, Proj-MW2D correlation analysis can also be used to assign the overlapping bands. In Figures 3(a) and 3(b), it can be easily distinguished that the crystalline bands of PBS and PLLA are 1714 and $1759 \mathrm{~cm}^{-1}$, respectively. This result is similar to previous reports. ${ }^{37,}{ }^{39,}{ }^{41-44}$ However, these reports used conventional FTIR, and their experimental methods were more complex than the ones used in the present study. In addition, these reports used the pure substance to confirm the assignment of PLLA, whereas we directly specified the assignment at 1714 and $1759 \mathrm{~cm}^{-1}$ from the PLLA/PBS blend. Therefore, ProjMW2D correlation analysis also provides a simple way to determine the spectral assignments for multiphase and multicomponent polymers.

Figure 3(c) shows the conventional MW2D correlation FTIR spectrum of the PLLA/PBS blend for comparison. The conventional MW2D correlation FTIR spectrum is the combination of the Proj-MW2D [Figure 3(a)] and null-space Proj-MW2D correlation FTIR spectra [Figure 3(b)]. However, in-depth analysis shows that the conventional MW2D method can only determine two transition temperatures of PLLA/PBS upon cooling. In addition, it cannot specify the origin of the molecular motion of the $\mathrm{C}=\mathrm{O}$ groups in the $1850 \mathrm{~cm}^{-1}$ to $1635 \mathrm{~cm}^{-1}$ region. The PCMW2D correlation FTIR spectra of the PLLA/PBS blend upon cooling are illustrated in Figure 4, including synchronous [Figure 4(a)] and asynchronous spectra [(Figure 4(b)]. Similiarly, PCMW2D can determine two transitions of PLLA/PBS at 98 ${ }^{\circ} \mathrm{C}$ and $78{ }^{\circ} \mathrm{C}$, respectively. However, similar to conventional MW2D, PCMW2D also cannot specify the origin of the $\mathrm{C}=\mathrm{O}$ motion at 1759 and $1714 \mathrm{~cm}^{-1}$. In Figure 4(a), the negative correlation peaks of blue areas represent the increase in the spectral intensities at 1759 and $1714 \mathrm{~cm}^{-1}$ in the cooling direction.

### 4.2. Extraction of the Molecular Motion of $\mathbf{C H}_{\mathbf{2}}$ Groups from Ethylene Units in the EB Component during Microcrystalline Melting in PS/EB

### 4.2.1. Temperature-dependent FTIR of PS/EB

The monodisperse PS and EB are thermodynamically incompatible. The molar content of the ethylene unit in EB was $71 \%$ (measured by NMR); thus, the ethylene units can form some microcrystals. ${ }^{17,45-47}$ The monodisperse PS used in this study is atactic PS. Figure 5 shows the temperaturedependent FTIR spectra of PS/EB in the $3055 \mathrm{~cm}^{-1}$ to 630 $\mathrm{cm}^{-1}$ region upon heating. The bands' assignments are listed in Table 1. The bands within the $3055 \mathrm{~cm}^{-1}$ to $2800 \mathrm{~cm}^{-1}$ region (Figure 5) show serious overlapping, particularly at 2925 and $2856 \mathrm{~cm}^{-1}$. The bands at 2925 and $2856 \mathrm{~cm}^{-1}$ correspond to the antisymmetric and symmetric stretching vibrations of $\mathrm{CH}_{2}$, respectively. ${ }^{48}$ However, not only are these
bands assigned to the $\mathrm{CH}_{2}(\mathrm{a}+\mathrm{b}+\mathrm{c})$ of the ethylene units and the 1-butene units in the EB component but also to the $\mathrm{CH}_{2}$ (d) of the main chain in the PS component. ${ }^{48}$ The target is to separate the molecular motion of $\mathrm{CH}_{2}$ groups of the ethylene units from the EB component upon heating. The temperaturedependent FTIR spectra cannot achieve this goal. Furthermore, no other characteristic means can be used to separate the molecular motion of $\mathrm{CH}_{2}$ groups from the EB component.


Figure 5. Temperature-dependent FTIR spectra of PS/EB in the $3055 \mathrm{~cm}^{-1}$ to $630 \mathrm{~cm}^{-1}$ region upon heating ( $30^{\circ} \mathrm{C}$ to $175^{\circ} \mathrm{C}$ ). The bands at 2925 and 2856 $\mathrm{cm}^{-1}$, which are serious overlapping bands, are assigned to the antisymmetric and symmetric stretching vibrations of $\mathrm{CH}_{2}$. The bands at $720 \mathrm{~cm}^{-1}$ is the $\mathrm{CH}_{2}$ rocking of the ethylene units of the EB component. The spectral intensity within the $730 \mathrm{~cm}^{-1}$ to $712 \mathrm{~cm}^{-1}$ region ( $720 \mathrm{~cm}^{-1}$ as the center) is selected as the projection matrix. The FTIR spectra within the $790 \mathrm{~cm}^{-1}$ to $670 \mathrm{~cm}^{-1}$ region are enlarged for clarity. For clear illustration, not all the collected FTIR spectra of PS/EB are shown.

### 4.2.2. Selection of the Projection Matrix

The projection matrix (or vector) must be determined to extract the molecular motion of $\mathrm{CH}_{2}$ groups of the ethylene units. According to the literature, ${ }^{45,}{ }^{48-52}$ the bands at 1462 , 1453, and $720 \mathrm{~cm}^{-1}$ are attributed to $\mathrm{CH}_{2}$ in the $1630 \mathrm{~cm}^{-1}$ to $630 \mathrm{~cm}^{-1}$ region. The bands at 1462 and $1453 \mathrm{~cm}^{-1}$, which still show serious overlapping, are attributed to the $\mathrm{CH}_{2}$ bending vibrations. ${ }^{48-51}$ These two bands are actually the corresponding bending vibration of 2925 and $2856 \mathrm{~cm}^{-1}$. Therefore, they cannot be used as the projection matrix (or vector). The bands at $720 \mathrm{~cm}^{-1}$, which is the $\mathrm{CH}_{2}$ rocking of the ethylene units of the EB component, is the characteristic bands of the ethylene units. ${ }^{45,48,52}$ This bands does not overlap with any bands of the PS component. The peak intensity of $720 \mathrm{~cm}^{-1}$ gradually weakened to disappear when the temperature increases from $30{ }^{\circ} \mathrm{C}$ to $175{ }^{\circ} \mathrm{C}$. Thus, the intensity variation at $720 \mathrm{~cm}^{-1}$ indicates the microcrystalline melting of the ethylene units of the EB component. Therefore, $720 \mathrm{~cm}^{-1}$ is particularly
suitable for the projection matrix (or vector). With a somewhat different approach from PLLA/PBS, we select the spectral intensity within the $730 \mathrm{~cm}^{-1}$ to $712 \mathrm{~cm}^{-1}\left(720 \mathrm{~cm}^{-1}\right.$ as the
center) region as the projection matrix rather than the projection vector.


Figure 6. Key: (a) Proj-MW2D correlation FTIR spectrum of PS/EB ( $40^{\circ}{ }^{\circ} \mathrm{C}$ to $165^{\circ} \mathrm{C}$ ); (b) null-space Proj-MW2D correlation FTIR spectrum of PS/EB ( $40^{\circ}{ }^{\circ} \mathrm{C}$ to $165^{\circ} \mathrm{C}$ ); (c) conventional MW2D correlation FTIR spectrum ( $40^{\circ} \mathrm{C}$ to $165^{\circ} \mathrm{C}$ ). The positive projection transformation algorithm is used for the calculations. The $\mathrm{CH}_{2}$ motion during the microcrystalline melting of the ethylene units from the EB component is extracted into the Proj-MW2D correlation spectrum, and the temperature is $53^{\circ} \mathrm{C}$. The molecular motion of $\mathrm{CH}_{2}$ groups of the main chain from the PS component is separated in the null-space Proj-MW2D correlation spectrum. However, the temperature is considerably high, namely, at $99^{\circ} \mathrm{C}$ and $144^{\circ} \mathrm{C}$, respectively.


Figure 7. Key: (a) Proj-MW2D correlation FTIR spectrum of PS/EB ( $40^{\circ} \mathrm{C}$ to $165^{\circ} \mathrm{C}$ ); (b) null-space Proj-MW2D correlation FTIR spectrum of PS/EB ( $40{ }^{\circ} \mathrm{C}$ to $165{ }^{\circ} \mathrm{C}$ ). The positive projection transformation algorithm is used for the calculations. The spectral intensity within the $710 \mathrm{~cm}^{-1}$ to $687 \mathrm{~cm}^{-1}$ region ( $698 \mathrm{~cm}^{-1}$ as the center), which is assigned to the out-of-plane bending vibration of the $=\mathrm{CH}$ of benzene rings of the PS component, is selected as the projection matrix. Once again, the molecular motion of $\mathrm{CH}_{2}$ groups generated from the PS component is successfully extracted. The molecular motion of $\mathrm{CH}_{2}$ groups of the ethylene units from the EB component is separated in the null-space Proj-MW2D correlation spectrum.

### 4.2.3. Proj-MW2D Analysis of the Molecular Motion of $\mathbf{C H}_{\mathbf{2}}$ Groups

The temperature-dependent FTIR spectra in the $3055 \mathrm{~cm}^{-1}$ to $2800 \mathrm{~cm}^{-1}$ were analyzed via the positive Proj-MW2D algorithm. Figures 6(a) and 6(b) are the Proj-MW2D and
null-space Proj-MW2D correlation FTIR spectra of the PS/EB blend upon heating, respectively. The molecular motion of $\mathrm{CH}_{2}$ groups with the same feature as the projection matrix $730 \mathrm{~cm}^{-1}$ to $712 \mathrm{~cm}^{-1}$ ( $720 \mathrm{~cm}^{-1}$ as the center) is retained in Figure 6(a). The $\mathrm{CH}_{2}$ motion in the ethylene units from the EB component is successfully separated. The correlation
intensity peaks appear at 2925 and $2856 \mathrm{~cm}^{-1}$, and the temperature is $53{ }^{\circ} \mathrm{C}$, representing the $\mathrm{CH}_{2}$ motion during the microcrystalline melting of the ethylene units from the EB component. ${ }^{17,46}$

In the null-space Proj-MW2D correlation spectra [Figure $\mathbf{6 ( b )}]$, the $\mathrm{CH}_{2}$ motion that is completely orthogonal (not identical) to the projection matrix is determined. That is, the molecular motion of $\mathrm{CH}_{2}$ groups is not related to the ethylene units of the EB component [Figure 6(b)]. The correlation intensity peaks appear at 2925 and $2856 \mathrm{~cm}^{-1}$ but at high temperatures of $99{ }^{\circ} \mathrm{C}$ and $144{ }^{\circ} \mathrm{C}$, respectively. Two transition temperatures represent the molecular motion of $\mathrm{CH}_{2}$ groups of the main chain from the PS component. According to the literature, ${ }^{53-62} 99^{\circ} \mathrm{C}$ is the glass transition temperature of PS. ${ }^{53-57}$ The temperature at $144{ }^{\circ} \mathrm{C}$ represents a liquidliquid transition. Previous studies attributed this transition to the viscous flow of PS. ${ }^{58-62}$

Figure 6(c) shows the conventional MW2D correlation FTIR spectrum, which is a combination of the Proj-MW2D and null-space Proj-MW2D correlation spectra. The correlation intensity peaks appear at 2925 and $2856 \mathrm{~cm}^{-1}$, and the corresponding temperatures are $53^{\circ} \mathrm{C}, 99^{\circ} \mathrm{C}$, and $144^{\circ} \mathrm{C}$. Although three transition temperatures can be observed, the conventional method cannot separate the molecular motion of $\mathrm{CH}_{2}$ groups from EB or PS component without the help of other characterization methods. The conventional MW2D correlation FTIR spectra cannot indicate the source component of the transitions at the functional group level.

### 4.2.4. Separation of the Molecular Motion of $\mathbf{C H}_{\mathbf{2}}$ Groups from the Main Chain in PS

The above results suggest that the characteristic bands of the PS component can be employed as the projection matrix to extract the $\mathrm{CH}_{2}$ motion of the main chain from PS component into the Proj-MW2D correlation FTIR spectrum. In Figure 5, the bands at 757 and $698 \mathrm{~cm}^{-1}$ are assigned to the out-of-plane bending vibration of the $=\mathrm{CH}$ of benzene rings. ${ }^{51}$ In the present study, the spectral intensity within the $710 \mathrm{~cm}^{-1}$ to 687 $\mathrm{cm}^{-1}$ region ( $698 \mathrm{~cm}^{-1}$ as the center) is selected as the projection matrix. Within the entire temperature-dependent FTIR spectrum (Figure 5), the bands within $710 \mathrm{~cm}^{-1}$ to 687 $\mathrm{cm}^{-1}$ are characteristic of the PS component. The spectral intensities of these bands are relatively large and do not coincide with any of bands of the EB component. Therefore, $710 \mathrm{~cm}^{-1}$ to $687 \mathrm{~cm}^{-1}$ ( $698 \mathrm{~cm}^{-1}$ as the center) is a suitable projection matrix for extracting the $\mathrm{CH}_{2}$ motion from PS component.

The Proj-MW2D and null-space Proj-MW2D correlation FTIR spectra are shown in Figures 7(a) and 7(b), respectively. The molecular motion with the same motion feature as the projection matrix ( $710 \mathrm{~cm}^{-1}$ to $687 \mathrm{~cm}^{-1}$ ) is retained in Figure 7(a). Once again, the molecuar motion of $\mathrm{CH}_{2}$ groups from the PS component is successfully separated. These results confirm that the correlation intensity peaks of 2925 and $2856 \mathrm{~cm}^{-1}$ at $99{ }^{\circ} \mathrm{C}$ and $144^{\circ} \mathrm{C}$ are the $\mathrm{CH}_{2}$ motion of the main chain from the PS component. Figure 7(a) is
almost the same as Figure 6(b). Different from Figure 6, Figure 7(b) shows that the molecular motion of $\mathrm{CH}_{2}$ groups of the ethylene units from the EB component is extracted in the null-space Proj-MW2D correlation FTIR spectrum. Similar to Figure 6(a), Figure 7(b) shows that the correlation intensity peaks of 2925 and $2856 \mathrm{~cm}^{-1}$ appear at $53{ }^{\circ} \mathrm{C}$. The use of $710 \mathrm{~cm}^{-1}$ to $687 \mathrm{~cm}^{-1}$ ( $698 \mathrm{~cm}^{-1}$ as the center) as the projection matrix can also separate the molecular motion of $\mathrm{CH}_{2}$ groups generated from the PS component.

## 5. Conclusions

Proj-MW2D was developed to separate the molecular motion of groups generated from different components or phases for multicomponent and multiphase polymers. The specific implementation steps for Proj-MW2D were enumerated after the theoretical derivation and algorithm research.

Two types of two-component blends, namely, PLLA/PBS and PS/EB, were employed to validate the concept of separating the molecular motion of groups. The molecular motion of $\mathrm{C}=\mathrm{O}$ groups from the crystalline phase in the PBS component was successfully extracted through the ProjMW2D correlation FTIR spectra of PLLA/PBS upon cooling. The characteristic bands at $806 \mathrm{~cm}^{-1}$ of PBS was used as the projection vector. Meanwhile, the molecular motion of $\mathrm{C}=\mathrm{O}$ groups from the PLLA component was separated in the nullspace Proj-MW2D correlation FTIR spectrum. The PBS crystalline phase showed the $\mathrm{C}=\mathrm{O}$ motion at $78{ }^{\circ} \mathrm{C}$. The molecular motion of $\mathrm{C}=\mathrm{O}$ groups was generated from the PLLA component at $98^{\circ} \mathrm{C}$.

The molecular motion of $\mathrm{CH}_{2}$ groups during the microcrystalline melting of the ethylene units from the EB component was successfully extracted at $53^{\circ} \mathrm{C}$ using the ProjMW2D correlation FTIR spectra of PS/EB upon heating. The projection matrix was the spectral intensity of $730 \mathrm{~cm}^{-1}$ to 712 $\mathrm{cm}^{-1}$ ( $720 \mathrm{~cm}^{-1}$ as the center). The molecular motion of $\mathrm{CH}_{2}$ groups of the main chain in the PS component was separated in the null-space Proj-MW2D correlation FTIR spectrum at 99 ${ }^{\circ} \mathrm{C}$ and $144^{\circ} \mathrm{C}$.

This study successfully used the Proj-MW2D correlation technique to separate the molecular motion of groups for multiphase and multicomponent polymers. This method allows researchers to study the mechanism of the complex transition process for multiphase and multicomponent polymer systems.

In the future, Proj-MW2D can be extended to three- or fourcomponent polymers, to other materials (e.g., small organic, biological, inorganic, and non-metallic materials), and to other spectra (e.g., Raman, X-ray, and UV). In addition, the projection technique was directly applied to MW2D correlation spectroscopy in this study. Nevertheless, it can be incorporated into PCMW2D methods and other forms of moving-window methods.

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## Notes and references

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