

This is an *Accepted Manuscript*, which has been through the Royal Society of Chemistry peer review process and has been accepted for publication.

Accepted Manuscripts are published online shortly after acceptance, before technical editing, formatting and proof reading. Using this free service, authors can make their results available to the community, in citable form, before we publish the edited article. This Accepted Manuscript will be replaced by the edited, formatted and paginated article as soon as this is available.

You can find more information about *Accepted Manuscripts* in the **Information for Authors**.

Please note that technical editing may introduce minor changes to the text and/or graphics, which may alter content. The journal's standard <u>Terms & Conditions</u> and the <u>Ethical guidelines</u> still apply. In no event shall the Royal Society of Chemistry be held responsible for any errors or omissions in this *Accepted Manuscript* or any consequences arising from the use of any information it contains.



www.rsc.org/advances

Molecular Nitrogen in N-doped TiO₂ Nanoribbons

C. Bittencourt¹, M. Rutar^{2,3}, P. Umek², A. Mrzel², K. Vozel², D. Arčon^{2, 5}, K. Henzler⁴, P. Krüger⁶ and P. Guttmann⁴

¹Chimie des Interactions Plasma Surface, CIRMAP, University of Mons, 23 Place du Parc, B-7000 Mons, Belgium

²Jožef Stefan Institute, Jamova cesta 39, SI-1000 Ljubljana, Slovenia

³Jožef Stefan International Postgraduate School, Jamova cesta 39, SI-1000 Ljubljana, Slovenia

⁴Helmholtz-Zentrum Berlin für Materialien und Energie GmbH, Institute for Soft Matter and Functional Materials, Albert-Einstein-Str. 15, D-12489 Berlin, Germany

⁵Faculty of mathematics and physics, University of Ljubljana, Jadranska cesta 19, SI-1000 Ljubljana, Slovenia

⁶Graduate School of Advanced Integration Science, Chiba University, Chiba 263-8522, Japan.

The nitrogen doping of TiO₂ nanoribbons during the thermal transformation of hydrogen titanate nanoribbons (HTiNRs) between 400 and 650 °C in a dynamic ammonia atmosphere was investigated using X-ray photoelectron spectroscopy (XPS), transmission X-ray microscopy combined with near-edge X-ray absorption fine structure spectroscopy (NEXAFS-TXM), X-ray diffraction (XRD) and electron paramagnetic resonance measurements (EPR). Comprehensive structural characterizations have revealed that for a calcination temperature of 400 °C, the HTiNRs transform into pure monoclinic TiO₂ β-phase (TiO₂-B) whereas at higher calcination temperatures (580 and 650 °C) a mixture of TiO₂-B and anatase is obtained. XPS and EPR results clearly reveal the nitrogen doping of TiO₂ nanoribbons and that, depending on the calcination temperature, nitrogen atoms occupy interstitial and substitutional sites. Moreover, in samples calcined at 580 and 650 °C the presence of N₂-like species in the HTiNRs was detected by NEXAFS-TXM. These species are trapped in the HTiNRs structure. EPR measurements upon

light illumination have disclosed generation of photoexcited states which implies that nitrogen has an important effect on the electronic structure of N-doped TiO₂.

Introduction

Nitrogen doping titanium dioxide (TiO₂), with N at substitutional sites has been reported to be indispensable for enlarging the use of TiO₂ as a visible-light photocatalytic material.¹ The main effect of the N-doping is the narrowing of the energy band gap of TiO₂ due to the mixing of N 2p and O 2p states.² Additionally, an isolated narrow band responsible for the visible light photoactivity is formed above the valence band.² In substitutional N-doping, the N atom bonds to three Ti atoms replacing a lattice O atom in TiO₂. Otherwise, N atoms can also occupy interstitial sites where N is bound to one or more O atoms thus creating paramagnetic centres. The electronic properties of N-doped titania, the corresponding bonding configuration, and the oxidation states of nitrogen in the lattice are yet to be unambiguously determined,³⁻⁵ Recent reports base on density functional theory (DFT) calculations and are consistent with both substitutional and interstitial nitrogen doping,⁶⁻⁹ where interstitial doping is slightly lower in energy especially for surface states with low concentration of oxygen vacancies.³ For further improving the photocatalytic performance of titania, it is thus critical to identify the local bonding configurations around N atoms at the atomic level.

Nitrogen doping into the TiO₂ lattice is rather challenging as N atoms must be accommodated geometrically and electronically.^{2, 10-12} To address the fundamental issues discussed above we explored a chemical route and report here a comprehensive investigation of N-doped TiO₂ nanoribbons. As a precursor material for N-doped TiO₂ nanoribbons we used hydrogen titanate nanoribbons (HTiNRs) because upon heating HTiNRs easily transform first to the monoclinic TiO₂¹³ β -phase (TiO₂-B) and then to anatase (tetragonal phase) while retaining the morphology

of the parent nanostructure intact¹⁴. In the used chemical route, the precursor HTiNRs were calcined at different temperatures for different times in an NH₃(g)/Ar(g) flow (**Table 1**). When calcined at 400 °C the HTiNRs successfully converted to TiO₂-B, whereas calcination at 580 °C and 650 °C yielded in a mixture of two TiO₂ phases. Morphology probing transmission electron microscopy (TEM) and structure sensitive X-ray diffraction (XRD) techniques were employed to extract critical local and global bonding information. The nitrogen chemical environment was inspected with X-ray photoelectron spectroscopy (XPS), while its bond geometry was probed by transmission X-ray microscopy combined with near-edge X-ray absorption fine structure spectroscopy (NEXAFS-TXM). The nitrogen chemical environment (i.e., substitutional or interstitial) was found to depend on the calcination temperature. The N 1s XPS and the N K-edge NEXAFS spectra of the N-doped TiO₂ nanoribbons calcined at 580 and 650 °C beside anticipated covalently bonded nitrogen also surprisingly detected N₂-like species. The present results provide a coherent description of N-sites in TiO₂ nanostructures and may thus facilitate further improvements in their photocatalytic performance.

Experimental:

Synthesis of N-doped TiO₂ nanoribbons

N-doped TiO₂ nanoribbons were synthesized from hydrogen titanate nanoribbons (HTiNRs) by calcination in a dynamic NH₃(g)/Ar(g) atmosphere. The reaction was carried out in a horizontal furnace, which was purged with a mixture of NH₃(g)/Ar(g). The NH₃(g)/Ar(g) flow ratio was set to 30 mL min⁻¹/10 mL min⁻¹, respectively. In a typical experiment 225 mg of HTiNRs, prepared from sodium trititanate nanoribbons¹⁵ by an ion-exchange process¹⁴, is weighted on a quartz boat, placed into the hot zone of the furnace and heated at a ramp rate of 7 °C/min to a selected

temperature. The samples were calcined at the selected temperature for four or six hours. The experimental conditions for the transformation of HTiNRs to TiO_2 nanoribbons, nitrogen content and phase composition are summarized in **Table 1**.

Table 1. Experimental conditions for the transformation of hydrogen titanate nanoribbons to N-doped TiO₂ nanoribons in an $NH_3(g)/Ar(g)$ flow. T (°C) temperature of transformation, t (h) duration time of the transformation, N nitrogen relative atomic concentration (as determined by XPS) and phase composition.

sample	T (°C)	t (h)	N content* (wt. %)	phase composition
N400-6h	400	6	0.8	TiO ₂ -B
N580-6h	580	6	1.3	anatase, TiO ₂ -B
N650-4h	650	4	2.8	anatase, TiO ₂ -B, traces of titanium oxynitride phase

Characterization techniques

The morphology of the N-doped TiO₂ nanoribbon samples was investigated with a field emission scanning (FE-SEM, Jeol, 7600F) and a transmission (TEM, Jeol 2100) electron microscopes. The phase analysis was performed on the cut surface by X-ray powder diffraction (XRD) using a diffractometer with Cu K_{α} radiation (λ =1.5406 Å) and a Sol-X energy-dispersive detector (Endeavor D4, Bruker AXS, Karlsruhe, Germany).

The XPS measurements were performed in a VERSAPROBE PHI 5000 from Physical Electronics, equipped with a Monochromatic Al K α X-Ray source. The energy resolution was 0.7 eV. For the compensation of built up charge on the sample surface during the measurements a dual beam charge neutralization composed of an electron gun (~1eV) and an Argon ion gun (≤ 10

eV) was used. The samples for XPS measurements were prepared by pressing the sample into a pellet. A conductive double face tape was used to attach the pellet to a sample holder.

The near edge X-ray adsorption fine structure (NEXAFS) spectra were recorded with the TXM end-station installed at the undulator beamline U41-XM at the electron storage ring BESSY II, Helmholtz-Zentrum Berlin (HZB).¹⁶ The spectra were recorded at room temperature in a transmission mode by taking a sequence of images over a range of photon energies covering the investigated absorption edges with $E/\Delta E \ge 4500$. The exit slit of the monochromator was set to 5 μ m, which corresponds to a calculated spectral resolution of $E/\Delta E = 20000$. The data analysis was accomplished by using aXis2000.¹⁷ The NEXAFS spectra were normalized by using the signal intensity in the proximity of the sample to correct for variations of the photon flux with photon energy (hv) and acquisition time. For comparison and calibration molecular nitrogen N₂(g) was analyzed using Nitrogen 5.0 with 99.999% purity (Air Liquide).

For the TEM and NEXAFS–TXM analyses the samples were sonically dispersed in ethanol and a drop of the solution was deposited onto a lacey carbon film supported by a copper grid. Note that due to electron-beam damage effects on the N-doped TiO₂ nanostructures, different grids were used for the TEM and for the NEXAFS–TXM analysis.

The EPR measurements were conducted using a commercial Bruker E580 spectrometer equipped with an Oxford Cryogenics He flow cryostat that has an optical window. The spectra were measured by recording the intensity of the echo signal as a function of magnetic field. A Hahnecho pulse sequence ($\pi/2 - \tau - \pi - \tau - echo$) together with appropriate phase cycling has been used. The length of $\pi/2$ pulse was set to 20 ns and the delay between the two pulses was $\tau = 200$ ns. The sample light-illumination was performed *in-situ* at 35 K with a standard halogen lamp. A sample for EPR measurements was sealed under dynamic vacuum in a standard Suprasil quartz EPR tube.

Results and Discussion

The typical morphology of the nanostructures prepared by calcination in the NH₃(g)/Ar(g) flow, presented in **Figure 1**, show that transformation from hydrogen titanate to N-doped TiO₂ does not affect the nanoribbon morphology. The nanoribbons reach a length of several micrometres with diameters in the range of 30 to 250 nm and are thus comparable to the starting hydrogen titanate nanoribbon (HTiNRs) sample¹⁸. TEM image of **N580-6h** (**Figure S1A**) reveals a crystalline nature of the N-doped TiO₂ nanoribbons. However, in striking contrast to the parent material, a close inspection of the nanoribbons' surface reveals their porous structure (**Figure S1B**) during the transformation of HTINRs to N-doped TiO₂ nanoribbons¹⁴ in the NH₃(g)/Ar(g) atmosphere. The porous structure is characteristic of TiO₂-B nanoribbons and with further transformation to anatase disappears as we reported in [14].

The crystal structure of samples treated at different calcination temperatures was examined by powder XRD (**Figure 2**). The XRD pattern of the sample calcined at 400 °C reveals only the presence of the TiO₂-B phase (JCPDS No. 35-0088), whereas the samples treated at 580 °C and 650 °C come as a mixture of TiO₂-B and anatase phases.^{13, 19} The amount of anatase in the samples increases with increasing calcination temperature up to 650° C. At this TiO₂ already starts to transform to titanium oxynitride as indicated with the appearance of a new diffraction peak at $2\theta \sim 43^{\circ}$.²⁰



Figure 1: SEM images of A) starting hydrogen titanate nanoribbons and N-doped TiO_2 nanoribbons prepared from hydrogen titanate nanoribbons by calcination in an $NH_3(g)/Ar(g)$ flow for 6 h at B) 400 °C, C) 580 °C and D) 650 °C.



Figure 2: XRD patterns of N-doped TiO₂ nanoribbons which were obtained by calcination of hydrogen titanate nanoribbons in an NH₃(g)/Ar(g) flow at 400 °C (6h), 580 °C (6h) and 650 °C (4h). Anatase peaks are labelled with asterisks (*), TiO₂-B peaks with triangles (\blacktriangle), and titanium oxynitride phase with a diamond (\blacklozenge). Two vertical lines positioned at (110) and (020) diffractions of TiO₂-B are added for easier notion of slightly different positions of (110) diffraction of TiO₂-B and (101) diffraction of anatase, and (020) diffraction of TiO₂-B and (200) diffraction of anatase.

The XPS analysis of the samples confirmed the presence of nitrogen in all calcined samples (**Table 1**). The relative content of nitrogen increases with the calcination temperature in agreement with data published in the literature²¹⁻²³. The samples with different nitrogen contents have colours that range from white over light yellow to light green. The higher the nitrogen content, the more intense is the colour of the sample. The XPS Ti 2p spectra show a Ti $2p_{3/2}$ binding energy of 458.7 eV with a splitting of 5.7 eV between the doublets, characteristic of titanium in the oxidation state 4+ (**Figure 3A**, **S2** and **S3**).^{3, 24} The absence of peaks corresponding to titanium in lower oxidation states in the XPS spectra of **N400-6h** and **N580-6h** suggests that there is no significant net electron transfer to titanium atoms from oxygen vacancy formation.³ Conversely, in the spectrum recorded on the sample calcined at 650°C (**N650-4h**) we observed a low intensity shoulder at low binding energies characteristic of the photoelectrons generated in titanium species having lower oxidation states (Ti⁺³ at 457.3 eV and Ti⁺² 455.5 eV)²⁵ (**Figure S3**).



Figure 3: XPS spectra of N580-6h: A) Ti 2p and B) O 1s.

Figure 3B shows the typical O 1s XPS spectrum recorded on the studied samples (**Figure S2B** and **S3B**). It exhibits a major peak centred at 530.0 eV, which is characteristic of transition metal oxides. The broad peak centred at 531.2 eV was attributed to hydroxyl groups and carbon compounds contamination at the TiO₂ surface.^{2, 26} Most likely the carbon contamination originates from the preparation of HTiNRs¹⁴ where ethanol was used for the final rinsing of HTiNRs after the ion exchange.

The N 1s XPS peaks recorded on N580-6h and N650-4h, show four components whereas in the N1s peak of N400-6h only three components are observed (Figure 4). The component with binding energy at 396.1 eV is assigned to O-Ti-N linkage in the lattice, i.e. lattice oxygen atoms are replaced by nitrogen atoms (substitutional N-doping). The relative intensity of this component and the total amount of nitrogen incorporated increased with increasing calcination temperature in accordance with the results reported by Zhang et al on the N-doping of on titania nanoparticles

prepared by a sol-gel technique²¹. The component at 396.7 eV was assigned to carbon contamination²⁶ as observed in the O 1s XPS spectra, while the component at 400.1 eV to nitrogen species bonded to oxygen sites (N-O-Ti linkage); i.e. an interstitial N-doping mode formed by N-atoms bonded to one or more lattice oxygen atoms.² The contribution of the component associated to photoelectrons emitted from N atoms in O-Ti-N linkages to the N1s XPS spectrum increases at the expense of the other three components as the calcination temperature increases. DFT calculations showed that adsorbed and interstitial nitrogen atoms have a tendency to bond to oxygen vacancy sites and accept electrons,⁶ explaining the high intensity of the component associated to N atoms in the O-Ti-N linkage. The highest binding energy peak at 403.7 eV has been ascribed to molecular N₂.^{24, 27} This binding energy is significantly lower than for gas phase N₂ molecules (409.9 eV), indicating that photocreated core holes maybe screened by neighbouring atoms as was observed in physisorbed N₂ molecules on various surfaces.²⁸ In order to investigate the presence of N_2 molecules in the N-doped TiO₂ nanoribbon samples (N-580-6h and N650-4h), we next studied the Ti L-edge and the N K-edge using the NEXAFS-TXM, which allows recording NEXAFS spectra from isolated nanoribbons or selected regions of the sample.¹⁶ The features of these edges are sensitive to the local symmetry and ligand coordination, providing detailed information on crystal structure and oxidation states of solid samples and vibrational fine-structure of gas molecules.



Figure 4: Comparison of the N 1s XPS spectra for the N-doped TiO₂ NRs derived from hydrogen titanate nanoribbons calcined in an $NH_3(g)/Ar(g)$ flow at A) 400, B) 580, and C) 650 °C. The peak at 403.7 is assigned to molecular N₂. The peak appearing at 396.3 eV is attributed to the O-Ti-N linkage (substituional N-doping), the relative intensity of this peak increases at the expense of the N₂ and N-O-Ti linkage (interstitial) as the calcination temperature increases.

The TXM-NEXAFS spectra were recorded on a region of interest of an isolated nanoribbon (**Figure S4**). The Ti L-edge spectrum of the N-doped TiO₂ nanoribbons contains two sets of doublets (**Figure S5, column A**), Ti L₃ at lower energy and Ti L₂ at higher energy. The two doublets originate primarily from electronic transitions from the Ti $2p_{3/2}$ and Ti $2p_{1/2}$ levels to Ti 3d orbitals, which split into t_{2g} and e_g levels in the approximately octahedral symmetry such as in the anatase TiO₂ phase.²⁹ The features of these Ti L-edge indicate that the coordination environment of the titanium atoms in the N-doped TiO₂ samples corresponds to Ti(IV) in a tetragonal structure as suggested by the XPS measurements.¹⁶

The N K-edge spectra recorded on N580-6h and N650-4h show a very sharp peak at 401 eV (Figure 5). To examine this more carefully, we recorded detailed spectra in the 399–403 eV photon energy range with step energy of 0.05 eV and compared with the spectrum recorded on a standard sample of molecular nitrogen (Figure 6). In both cases we observed five well-resolved peaks that are characteristic of free N₂ molecules.^{24, 30, 31} The peak at 400.8 eV of the gas phase N₂ molecules was assigned to a transition from the N 1s to π^* orbital and the fine feature is a hallmark of N₂ molecule due to its vibrational structure.³² Therefore, the observed five structure

in the N K-edge undoubtedly proves the presence of N₂ molecules in the N-doped TiO₂ nanoribbons. This is further supported by the observation that the energy position of the 5 components of the peak, their area ratio, and energy separation are similar to those of gas phase N₂ molecule. This means that N-N bond order of N₂ in the N-doped TiO₂ nanoribbons is comparable to that of gas phase N₂ molecule and the interaction of N₂ with the N-doped TiO₂ NRs can be considered as very weak.³⁰ Therefore, the N1s XPS and the N K-edge NEXAFS spectra support the idea that nitrogen atoms trapped in N-doped TiO₂ nanoribbons may combine to form N₂ molecules evolving into the gas phase upon heating at elevated temperatures. This behaviour agrees well with reported DFT calculations which show that nitrogen embedded in TiO₂ is unstable energetically and spontaneously forms trapped or gaseous N₂.³ Therefore, the XPS and the NEXAFS analyses suggest that for the doping parameters used, the incorporated nitrogen atoms remain in the oxide structure as substitutional or interstitial sites and do not segregate into the ordered titanium oxynitride phase for calcination temperature in the range of 400 to 580 °C; segregation was observed at 650 °C.

RSC Advances Accepted Manuscrip



Figure 5: NEXAFS-TXM spectra of the N K-edge of the N-doped TiO_2 nanoribbons recorded with an energy step of 0.07 eV.



Figure 6: Comparison of high-energy resolution NEXAFS-TXM N K-edge spectra (step energy 0.05 eV) recorded on **650-4h** and on N_2 gas. The characteristic vibrational structures in the spectrum of the N-doped TiO₂nanoribbons indicate the presence of N_2 molecules.

The formation of paramagnetic NO_x defects and the possible changes in the electronic structure after nitrogen doping prompted us to carry out electron paramagnetic resonance (EPR). EPR has proved to be a very powerful tool for the investigation of local properties of various paramagnetic centres in TiO_x nanostructures.^{18, 33-37} Moreover, it can also distinguish between different NO_x centers through the values of electronic *g*-factors and hyperfine coupling tensors **A**.^{18, 38, 39} The X-band EPR spectrum of **N580-6h** sample measured at 35 K with the echo-detected field-sweep technique is shown in **Figure 7A**. Clearly, at least three different signals can be recognised. The

broad component at high fields can be simulated with the uniaxial *g*-factor anisotropy yielding g_{\perp} = 1.907 and g_{\parallel} = 1.939. These eigenvalues match those expected for the Ti³⁺ surface sites, (Ti³⁺)s,⁴⁰. The large homogenous broadening ($\Delta B_{1/2}$ = 4.7 mT) used in the simulation for this component is consistent with the strains and local inhomogeneities associated with such surface states and the fit could be improved if a distribution of *g*-factor values is considered. Much narrower homogenous broadening ($\Delta B_{1/2}$ = 0.9 mT) is characteristic for the second component with typical g_{\perp} = 1.974 and g_{\parallel} = 1.986, also obtained from the powder lineshape fit. These values allow us to assign this signal to the bulk Ti³⁺ sites, (Ti³⁺)_B,⁴¹. Finally, the fitting of the low-field signal at *g* = 2.0054 with two characteristic powder hyperfine satellite lines on both sides of the sharp peak yield the hyperfine coupling values of A_⊥ = 18 MHz and A_{||} = 94 MHz. Both, g-factor and hyperfine coupling values are characteristic values for NO paramagnetic centres¹⁸ and provide independent support for the formation of N-O-Ti linkage. The EPR data fully corroborates with the results of the XPS study (**Figure 4**) and thus unambiguously prove the incorporation of N into the TiO₂ lattice.

Additionally, since the nitrogen doping is also known to enhance visible light photoactivity in TiO_2 , we measured EPR spectra of illuminated samples. Interestingly, different components of the EPR spectrum respond differently to light illumination. The spectrum measured at 35 K after illumination (**Figure 7B**) still contains the same NO, $(Ti^{3+})_B$, and $(Ti^{3+})_S$ signals as their respective *g*-factors and hyperfine coupling constants coincide with those before the illumination (Figure 8A). However, while the signal assigned to surface Ti^{3+} sites remains almost the same, the signals of NO and bulk Ti^{3+} centres double in intensity after 10 min of illumination. This finding thus implies that photoexcited charges preferentially trap in the bulk of the nanoribbons where they form stable paramagnetic centres.



Figure 7: X-band electron paramagnetic resonance spectra of **N580-6h** measured at 35 K with an echo-detected field-sweep technique (grey shaded area) before A) and after B) illumination with light. The spectra fit to the sum of three different signals (solid red line) ascribed to paramagnetic NO and to bulk $(Ti^{3+})_B$ and surface $(Ti^{3+})_S$ sites, respectively.

Conclusions

We successfully prepared N-doped TiO₂ nanoribbons using hydrogene titanate nanoribbons as a precursor by a thermal treatment in a dynamic ammonia atmosphere. Two types of nitrogen species with bonding energies of 396.1 eV and 400.1 eV were observed by X-ray photoelectron spectroscopy. The two nitrogen species can be assigned to substitutional nitrogen doping and nitrogen bonded to oxygen sites (interstitial), respectively. The amount of the substitutional nitrogen increases with the increasing calcination temperature. The NEXAFS-TXM studies show that local bonding environments of oxygen and titanium atoms in N-doped TiO₂ nanoribbons are similar to the undoped TiO₂. Depending on the calcination temperature N₂ molecules can be formed and remain weakly trapped in the N-doped TiO₂ structure. The changes observed in the

EPR spectrum upon light illumination indicate that photoexcited charges preferentially trap in the

bulk of the nanoribbons where they form stable paramagnetic centres.

Acknowledgements:

This work is financially supported by the Helmholtz-Zentrum Berlin für Materialien und Energie GmbH. The research leading to these results has received funding from the European Community's Seventh Framework Programme (FP7/2007-2013) under grant agreement n.°312284. We thank HZB for the allocation of synchrotron radiation beamtime at the undulator beamline U41-TXM at the BESSY II electron storage ring, Berlin. C. Bittencourt is Research Associate of the National Funds for Scientific Research (FRS-FNRS, Belgium).

References

- 1. R. Asahi, T. Morikawa, T. Ohwaki, K. Aoki and Y. Taga, *Science*, 2001, **293**, 269-271.
- 2. H. H. Ou, S. L. Lo and C. H. Liao, *Journal of Physical Chemistry C*, 2011, **115**, 4000-4007.
- 3. H. Chen, A. Nambu, W. Wen, J. Graciani, Z. Zhong, J. C. Hanson, E. Fujita and J. A. Rodriguez, *Journal of Physical Chemistry C*, 2007, **111**, 1366-1372.
- 4. M. Zukalova, J. Prochazka, Z. Bastl, J. Duchoslav, L. Rubacek, D. Havlicek and L. Kavan, *Chemistry* of Materials, 2010, **22**, 4045-4055.
- 5. M. Batzill, E. H. Morales and U. Diebold, *Physical Review Letters*, 2006, **96**.
- 6. A. Nambu, J. Graciani, J. A. Rodriguez, Q. Wu, E. Fujita and J. F. Sanz, *Journal of Chemical Physics*, 2006, **125**.
- 7. C. Di Valentin, G. Pacchioni and A. Selloni, *Physical Review B Condensed Matter and Materials Physics*, 2004, **70**, 085116-085111-085116-085114.
- 8. C. Di Valentin, A. Tilocca, A. Selloni, T. J. Beck, A. Klust, M. Batzill, Y. Losovyj and U. Diebold, *Journal of the American Chemical Society*, 2005, **127**, 9895-9903.
- 9. C. Di Valentin, G. Pacchioni, A. Selloni, S. Livraghi and E. Giamello, *Journal of Physical Chemistry B*, 2005, **109**, 11414-11419.
- 10. X. Sun and Y. Li, *Chemistry A European Journal*, 2003, **9**, 2229-2238.
- 11. H. H. Ou, C. H. Liao, Y. H. Liou, J. H. Hong and S. L. Lo, *Environmental Science and Technology*, 2008, **42**, 4507-4512.
- 12. H. H. Ou and S. L. Lo, *Separation and Purification Technology*, 2007, **58**, 179-191.
- 13. A. R. Armstrong, G. Armstrong, J. Canales and P. G. Bruce, *Angewandte Chemie International Edition*, 2004, **43**, 2286-2288.
- 14. P. Umek, C. Bittencourt, P. Guttmann, A. Gloter, S. D. Å kapin and D. Arčon on, *Journal of Physical Chemistry C*, **118**, 21250-21257.

- 15. P. Umek, R. C. Korošec, B. Jančar, R. Dominko and D. Arčon, *Journal of Nanoscience and Nanotechnology*, 2007, **7**, 3502-3508.
- 16. P. Guttmann, C. Bittencourt, S. Rehbein, P. Umek, X. Ke, G. Van Tendeloo, C. P. Ewels and G. Schneider, *Nature Photonics*, 2012, **6**, 25-29.
- 17. A. P. Hitchock.
- 18. P. Umek, P. Cevc, A. Jesih, A. Gloter, C. P. Ewels and D. Arčon, *Chemistry of Materials*, 2005, **17**, 5945-5950.
- 19. Y. V. Kolen'ko, K. A. Kovnir, A. I. Gavrilov, A. V. Garshev, J. Frantti, O. I. Lebedev, B. R. Churagulov, G. Van Tendeloo and M. Yoshimura, *Journal of Physical Chemistry B*, 2006, **110**, 4030-4038.
- 20. M. Zukalova, J. Prochazka, Z. Bastl, J. Duchoslav, L. Rubacek, D. Havlicek and L. Kavan, *Chemistry* of Materials, **22**, 4045-4055.
- 21. Z. Zhang, X. Wang, J. Long, Q. Gu, Z. Ding and X. Fu, *Journal of Catalysis*, **276**, 201-214.
- 22. J. Wang, D. N. Tafen, J. P. Lewis, Z. Hong, A. Manivannan, M. Zhi, M. Li and N. Wu, *Journal of the American Chemical Society*, 2009, **131**, 12290-12297.
- 23. H. Irie, Y. Watanabe and K. Hashimoto, *Journal of Physical Chemistry B*, 2003, **107**, 5483-5486.
- 24. Y. Chung, J. C. Lee and H. J. Shin, *Applied Physics Letters*, 2005, **86**, 022901-022901-022901-022903.
- 25. C. D. Wagner, A. V. Naumkin, A. Kraut-Vass, J. W. Allison, C. J. Powell and J. R. Jr. Rumble, 2003.
- 26. M. Grandcolas and J. Ye, *Journal of Ceramic Processing Research*, 2012, **13**, 65-70.
- 27. P. Hoffmann and C. Pettenkofer, *Physica Status Solidi (B) Basic Research*, 2011, **248**, 327-333.
- 28. J. Stöhr, *NEXAFS Spectroscopy*, 2 edn., Springer, 2003.
- 29. C. Bittencourt, P. Krüger, M. J. Lagos, X. Ke, G. Van Tendeloo, C. Ewels, P. Umek and P. Guttmann, *Beilstein Journal of Nanotechnology*, 2012, **3**, 789-797.
- 30. M. Petravic, Q. Gao, D. Llewellyn, P. N. K. Deenapanray, D. Macdonald and C. Crotti, *Chemical Physics Letters*, 2006, **425**, 262-266.
- 31. J. Zhou, J. Wang, H. Liu, M. N. Banis, X. Sun and T. K. Sham, *Journal of Physical Chemistry Letters*, 2010, **1**, 1709-1713.
- 32. J. Zhou, J. Wang, H. Liu, M. N. Banis, X. Sun and T. K. Sham, *Journal of Physical Chemistry Letters*, **1**, 1709-1713.
- 33. P. Umek, M. Pregelj, A. Gloter, P. Cevc, Z. Jaglicic, M. Čeh, U. Pirnat and D. Arčon, *Journal of Physical Chemistry C*, 2008, **112**, 15311-15319.
- 34. P. Umek, C. Bittencourt, A. Gloter, R. Dominko, Z. Jaglicic, P. Cevc and D. Arčon, *Journal of Physical Chemistry C*, 2012, **116**, 11357-11363.
- 35. Z. V. Saponjic, N. M. Dimitrijevic, O. G. Poluektov, L. X. Chen, E. Wasinger, U. Welp, D. M. Tiede, X. Zuo and T. Rajh, *Journal of Physical Chemistry B*, 2006, **110**, 25441-25450.
- 36. P. Szirmai, E. HorvÃith, B. Náfrádi, Z. MickoviÄ[‡], R. Smajda, D. M. DjokiÄ[‡], K. Schenk, L. Forro and A. Magrez, *Journal of Physical Chemistry C*, **117**, 697-702.
- 37. R. Arroyo, G. Cardoba, J. Padilla and V. H. Lara, *Materials Letters*, 2002, **54**, 397-402.
- 38. D. Biglino, H. Li, R. Erickson, A. Lund, H. Yahiro and M. Shiotani, *Physical Chemistry Chemical Physics*, 1999, **1**, 2887-2896.
- 39. H. Yahiro, A. Lund and M. Shiotani, *Spectrochimica Acta Part A: Molecular and Biomolecular Spectroscopy*, 2004, **60**, 1267-1278.
- 40. R. F. Howe and M. Grätzel, *Journal of Physical Chemistry*, 1987, **91**, 3906-3909.
- 41. R. F. Howe and M. Grätzel, *Journal of Physical Chemistry*, 1985, **89**, 4495-4499.