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**A First principle study of H<sub>2</sub>S adsorption and decomposition on Ge (100)  
surface**

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**Abstract**

We employed density functional theory (DFT) calculation to examine the adsorption configurations and possible reaction paths for H<sub>2</sub>S on Ge(100) surface. There are four reaction paths are proposed for the decomposition of adsorbed H<sub>2</sub>S on Ge(100) surface and the corresponding structural conformations are studied extensively. The present study shows that two new possible products and detailed reaction mechanism for H<sub>2</sub>S adsorption on Ge(100) surface and the results are compared with our previous study of H<sub>2</sub>S adsorption on Si(100) surface (J. Phy. Chem. C 115, 2011, 19203). The density of states (DOS) and electron density difference (EDD) analyses are used to illustrate the interaction between S-containing species and surface Ge atoms.

Key words: H<sub>2</sub>S, DOS, density functional theory, Ge(100).

## 1. Introduction

Over the past decades, high mobility semiconductors have investigated as a replacement for the Si channel due to its potential applications in high-performance metal-oxide-semiconductor (MOS) devices.<sup>1-3</sup> Recently, Germanium (Ge) has attracted much attention because it gives bulk electron and hole motilities at room temperature which are higher than the conventional silicon based materials.<sup>4, 5</sup> However, major drawbacks of the application of Ge in these devices are the poor quality of its thermal oxide compared to SiO<sub>2</sub> and the Ge-based MOS devices has high interface state density ( $D_{it}$ ) which would result in low carrier mobility in the channel. These limitations can be overcome by the passivation of Ge surface prior to the deposition of the gate dielectric which is a key issue for semiconductor surfaces.<sup>6, 7</sup>

Also, the practical use of this semiconductor in the devices requires the passivation to prevent oxidation and maintain the surface order during device processing. Sulfur (S) has been shown to be one of the best passivants of semiconductor surfaces.<sup>4</sup> Anderson et.al., reported that the Sulfur passivation of the germanium surface via both solution and vacuum deposition techniques.<sup>8</sup> A sulfur adlayer is deposited on a semiconductor surface through immersion in a solution of (NH<sub>4</sub>)<sub>2</sub>S<sup>8</sup> or on its exposure to elemental sulfur or H<sub>2</sub>S.<sup>9</sup> On the other hand, Houssa et al. proposed two approaches for the passivation of Ge surface to promote metal-oxide-semiconductor field-effect transistor (MOSFET).<sup>10</sup> One of two approaches is the exposure of the Ge surface to H<sub>2</sub>S, resulting in the formation of S-Ge bonds. Indeed, the Ge energy band gap is free of surface states after H<sub>2</sub>S exposure, resulting in the electrical passivation of the surface, which is a crucial requirement for the potential use of Ge in the MOSFET.

Similar to Si(100) and C(100) surfaces, the reconstructed Ge(100)-2 × 1 surfaces consist of dimers connected by a strong  $\sigma$ -bond and a weak  $\pi$ -bond.<sup>11</sup> The diamond surface has symmetric dimers with a bond length of approximately 1.4Å, whereas the Si and Ge surfaces both have larger asymmetric, or tilted, dimers with a bond length of 2.3-2.5Å.<sup>12-14</sup> Distortion of the dimer bonds induces a charge transfer from the down (electrophilic) to the upper (nucleophilic) surface atom.<sup>15</sup> Such zwitterionic characteristics allow the surface to undergo a nucleophilic/electrophilic reaction, often with direct analogies to the molecular systems. The adsorption of sulfur on Ge(100) surface has been studied both experimentally and theoretically.<sup>16-18</sup>

Previously, the S-passivation on Ge(100) surface has been investigated using low-energy electron diffraction (LEED),<sup>8, 19</sup> high-resolution electron-energy loss spectroscopy (EELS),<sup>20</sup> ultraviolet photoelectron spectra (UPS),<sup>21</sup> X-ray photoelectron spectra (XPS),<sup>22</sup> temperature programmed desorption (TPD),<sup>23</sup> Near edge X-ray absorption fine structure (NEXAFS)<sup>24</sup> and multiple internal reflection-Fourier transform infrared spectrometer (MIR-IR).<sup>25</sup> Kuhr and Ranke<sup>21</sup> showed in their UPS study that, H<sub>2</sub>S adsorbs dissociatively at 300K on a Ge surface; then decomposes completely to form sulfur at 550K. To understand the sulfidation of the Ge(100) surfaces and its possible passivating reactions, we choose H<sub>2</sub>S as a passivant to be adsorbed on the Ge(100) surface. In this study, we report ab initio theoretical investigations to provide a complete description of the adsorption effects and thermal decompositions of H<sub>2</sub>S on Ge(100) surface. We thereby elucidates possible mechanisms of thermal decomposition on Ge(100) that was more conformed to real surface.

## 2. Computational Details

DFT calculations are performed using the Vienna ab initio simulation package (VASP).<sup>26-28</sup> The Vanderbilt ultra-soft pseudopotential is used to describe the electron-ion interactions<sup>29</sup> and the electron- electron exchange and correlation contributions by the generalized gradient approximation (GGA) with energy truncated at 300 eV. The Brillouin zone is sampled with Monkhorst-Pack grid, and the calculations were performed with (4×2×1) Monkhorst-Pack mesh *k*-points. We optimized the structures based on the conjugate gradient-minimization scheme taking into consideration of spin polarization. The validities of all the optimized structures and determined transition states (TSs) are checked through normal-mode frequency analysis. For a real minimum structure, all frequency must be positive; TS must have one imaginary frequency corresponding to the reaction coordinate. The clean surfaces are then modeled in the form of a slab. The c(4×2) unit cell consisted of nine atomic layers separated by a 20 Å vacuum to prevent interactions between the surface adsorbates and the preceding slab. The optimized geometry of the model Ge(100) surface considered in this study is depicted in Figure 1.

In the structural optimizations, the positions of the last three Ge layers along with the H-passivated layer on the bottom are fixed upon adsorption of H<sub>2</sub>S and the remaining substrate atoms are allowed to relax with the adsorbates. To economize computing time, only the upper two atomic layers of the surface and adsorbate are relaxed in the vibrational frequency calculations. The adsorption energies are obtained by

$$E_{ads} = E_{Total} - (E_{Gas} + E_{Surface})$$

in which  $E_{Total}$ ,  $E_{Gas}$  and  $E_{Surface}$  are the respective calculated electronic energies of the adsorbed species on the surface, a gaseous molecule, and a clean surface. A negative value

for  $E_{ads}$  indicates an exothermic adsorption. The relative energies ( $E_{rel}$ ) are defined as

$$E_{rel} = E_{(LM\ or\ FS)} - E_{H_2S_{ad}}$$

in which  $E_{H_2S_{ad}}$  and  $E_{(LM\ or\ FS)}$  are the calculated adsorption energies of H<sub>2</sub>S and local minima or final products, respectively. The nudged-elastic-band (NEB) method is applied to locate the transition state structures positioned to interpolate a series of system images between the initial and final states on the potential-energy surface.<sup>30, 31</sup> A spring force between adjacent images is used to maintain constant spacing between the images, and a true force is applied to impel the images into the minimum energy path (MEP), mimicking an elastic band. Each image is optimized using the NEB algorithm based on a constrained algorithm of molecular dynamics. The highest point on the MEP corresponds to a transitional structure on the proposed reaction path and its energy, relative to that of the initial state, became the activation barrier of the reaction.

The electron density difference ( $Q_{diff}$ ) is calculated in a similar manner to the calculation of adsorption energy

$$Q_{diff} = Q_{A/Surf} - (Q_{Surf} + Q_A)$$

where  $Q_{diff}$  is the difference at each grid point in the total electron density matrix between that of the adsorbate-bonded surface ( $Q_{A/Surf}$ ) and that of the sum of the surface ( $Q_{Surf}$ ) and the single adsorbate molecule ( $Q_A$ ). According to this definition, positive and negative values correspond to increasing and decreasing electron densities, respectively.

### 3. Results and Discussion

#### 3.1. Adsorption of H<sub>2</sub>S, HS and S on Ge(100) surfaces

To ensure the reliability of the computational method, we have calculated the lattice parameters for bulk germanium using various pseudopotentials and the values are given in Table S1 of Supplementary information. As can be seen from those values PAW-PBE and the ultrasoft pseudopotential with the generalized gradient approximation (US-GGA) gives the least discrepancy between the calculated and experimental lattice parameters among the tested pseudopotentials. In order to compare with our previous study of H<sub>2</sub>S adsorption on Si(100) surface,<sup>32</sup> we considered US-GGA method for further calculations. Also, we have performed the benchmark calculations for the H<sub>2</sub>S adsorption on the clean Ge(100) surfaces with different functionals and different cutoff energy values and the calculated values are given in Table S2. It has been observed that the trends of the calculated values with different cutoff energies and functionals are almost similar and hence we have chosen the US with 300 eV cutoff energy values for our further study. Table S3 presents the comparison of calculated bond lengths and bond angles of clean Ge(100) surface with other studies<sup>33</sup> and it has been found that the calculated bond lengths are in agreement with the previous experimental reports.<sup>34-36</sup>

Previously, we found that the H<sub>2</sub>S has three stable adsorption conformations on the Si(100) surfaces with orientation of the S-H bond;<sup>32</sup> similarly here we considered the same conformation of H<sub>2</sub>S adsorbed on Ge(100) surfaces (denoted as H<sub>2</sub>S<sub>ad</sub>, H<sub>2</sub>S<sub>ad-1</sub>, and H<sub>2</sub>S<sub>ad-2</sub>) and the optimized structures are shown in Figure 2. The calculated adsorption energies and structural parameters for different conformation of H<sub>2</sub>S are summarized in Table 1. Among these three conformations, H<sub>2</sub>S<sub>ad</sub> is found to be the most stable with adsorption energy -0.49 eV and the other two conformations H<sub>2</sub>S<sub>ad-1</sub> and H<sub>2</sub>S<sub>ad-2</sub> adsorbed to the surface with

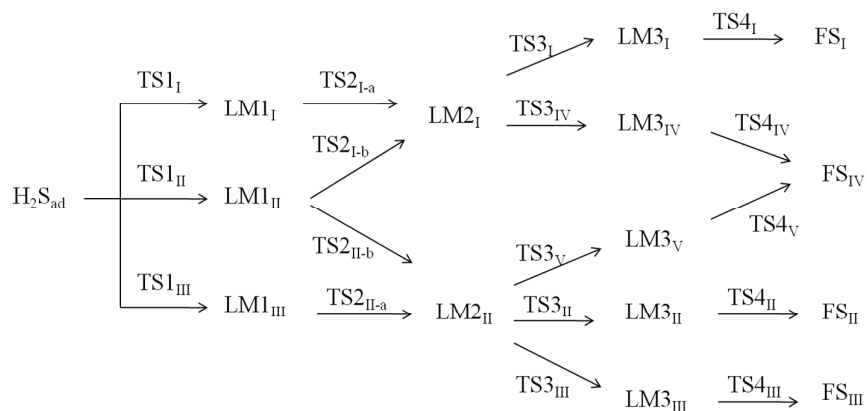


adsorption energies -0.40 and -0.30 eV, respectively. From Table 1, it has been observed that the bond length of the Ge-Ge dimer is increased ( $\sim 0.04$  Å) and the tilting angle of Ge dimer is decreased to  $16^\circ$  from  $20^\circ$  after the adsorption of  $\text{H}_2\text{S}$  on the Ge surface, indicating the decreased zwitterionic and  $\pi$ -bond characters of the Ge dimer upon the adsorption of  $\text{H}_2\text{S}$ . The calculated S-Ge bond length value for the most stable conformation ( $\text{H}_2\text{S}_{\text{ad}}$ ) is 2.61 Å, and the two remaining S-H bond lengths are 1.35 and 1.38 Å. The structural parameters of other two conformations,  $\text{H}_2\text{S}_{\text{ad-1}}$  and  $\text{H}_2\text{S}_{\text{ad-2}}$ , are similar to those of  $\text{H}_2\text{S}_{\text{ad}}$ . Figure 3 shows the DOS of surface Ge down atom and  $\text{H}_2\text{S}$  molecule before and after adsorption on Ge(100) surface, respectively. It is noticed from this figure that the  $\text{H}_2\text{S}$  molecule possesses  $C_{2v}$  symmetry with four valence states labeled as  $4a_1$ ,  $2b_2$ ,  $5a_1$ , and  $2b_1$  according to their orbital symmetries. These results are similar to that of our previous study.<sup>32</sup> The partial DOS (see Figure S1 of Supplementary information) indicates that states  $4a_1$  and  $2b_1$  are contributed mainly from  $3s$  and  $3p_z$  of S atom and  $2b_2$  and  $5a_1$  states arise from the hybridization of orbitals  $3p_x$  and  $3p_y$  of the S atom. All S  $3p$  atomic orbitals ( $2b_2$ ,  $5a_1$ , and  $2b_1$  states), especially  $3p_x$  and  $3p_y$  of adsorbed  $\text{H}_2\text{S}$  have good overlaps with d orbitals of the buckled-down Ge which is bonded to the  $\text{H}_2\text{S}$  molecule.

### 3.2. Reaction Mechanism of $\text{H}_2\text{S}$ on Ge(100)

We consider the reaction pathway for adsorption of  $\text{H}_2\text{S}$  on Ge(100) surface analogous to our previous study of the adsorption of  $\text{H}_2\text{S}$  on Si(100).<sup>32</sup> The calculated structural parameters and relative energies of intermediates and final products are listed in Table 2. The reaction energies, reaction barriers, and calculated imaginary frequencies of transition state structures are summarized in Table 3. The corresponding structures of intermediates and final products are shown in Figure 4 and the top views of their corresponding transition state structures are shown in Figure S2 of Supplementary information. The potential energy diagram (PES) for the adsorption and dissociation of  $\text{H}_2\text{S}$

on the Ge(100) surface are shown in Figure 5. According to the calculation, the possible reaction paths are summarized in the following scheme:



The possible reaction pathways described in above scheme includes first dehydrogenation, second dehydrogenation and sulfur bridged adsorption. The  $\text{H}_2\text{S}_{\text{ad}}$  species can undergo a first dehydrogenation through three paths. The first path is through  $\text{TS1}_I$ , with a 0.53 eV barrier, to produce  $\text{LM1}_I$  intermediate; i.e., one H atom of  $\text{H}_2\text{S}$  dissociated and migrated to the Ge atom of an adjacent dimer along [001] direction. It has been observed that the first dehydrogenation has higher energy barrier due to the H atom being adsorbed on the Ge(3) atom. The second path creates the intermediate  $\text{LM1}_{II}$  via  $\text{TS1}_{II}$  with a barrier of 0.49 eV; i.e., the dissociated H and HS separately adsorb on each of a dimer's Ge atom. The third path yields intermediate  $\text{LM1}_{III}$  via  $\text{TS1}_{III}$  with the smallest barrier 0.10 eV; the resulting H and HS adsorb on Ge atoms of adjacent dimers along direction [010]. Figures 6(a) and (b) show the EDD contour plots which contain the dimeric S-Ge(1) bond with S- $\text{H}_a$  and S-Ge (1) bond with S- $\text{H}_b$  bond, respectively. It can be seen from Figure 6(a) that the electron density increases between the S and Ge(1) atoms which is bonded to  $\text{H}_2\text{S}$ , indicating a strengthening of S-Ge bond. In the meanwhile, the increased electron density is observed between the H and Ge(5) atoms of a neighboring Ge dimer along [010] direction (Figure 6(a)) which interacts with each other through a hydrogen bond. This interaction results in

the lower energy barrier (0.10 eV) in the first dehydrogenation for  $\text{H}_2\text{S}_{\text{ad}}$  to form  $\text{LM1}_{\text{III}}$ . As summarized in Table 3, the reaction energies of the first dehydrogenations for the three paths are -0.89, -1.03, and -0.96 eV, respectively.  $\text{LM1}_{\text{II}}$  and  $\text{LM1}_{\text{III}}$  are comparably stable because the first dissociated H preferably adsorbs on the electron rich atoms like Ge(2) and Ge(5) atoms. The S-Ge bond length in  $\text{LM1}_{\text{I}}$ ,  $\text{LM1}_{\text{II}}$ , and  $\text{LM1}_{\text{III}}$  are shortened by  $\sim 0.3\text{\AA}$ , compared to that of in the  $\text{H}_2\text{S}_{\text{ad}}$ . In  $\text{H}_2\text{S}_{\text{ad}}$ ,  $\text{H}_2\text{S}$  adsorbed on a Ge atom via a dative bond; whereas in  $\text{LM1}_{\text{I-III}}$ , it is adsorbed on a Ge atom via a covalent bond.

Further, the local minima's  $\text{LM1}_{\text{I}}$ , and  $\text{LM1}_{\text{II}}$  can undergo dehydrogenation and form  $\text{LM2}_{\text{I}}$  via the transition states,  $\text{TS2}_{\text{I-a}}$  and  $\text{TS2}_{\text{I-b}}$  with energy barriers of 0.90 and 0.88 eV, respectively.  $\text{LM2}_{\text{II}}$  can be produced from the dehydrogenation of  $\text{LM1}_{\text{III}}$  and  $\text{LM1}_{\text{II}}$  via transition states,  $\text{TS2}_{\text{II-a}}$ , and  $\text{TS2}_{\text{II-b}}$  with barriers of 1.89 and 1.07 eV, respectively. From Table 2, it has been observed that the structural parameters of  $\text{LM2}_{\text{I}}$  and  $\text{LM2}_{\text{II}}$  are almost similar. The S-Ge bond of  $\text{LM2}$  species is much shorter than that of  $\text{LM1}$ , indicating a stronger S-Ge bond in species  $\text{LM2}$ . The energy barriers of the second dehydrogenation,  $\text{LM1} \rightarrow \text{LM2}$ , are higher than that of first dehydrogenation  $\text{H}_2\text{S}_{\text{ad}} \rightarrow \text{LM1}$ . In the reaction mechanism, there are two types of intermediates between the  $\text{LM2}$  and the final products; one is formed through the migration of H atom from  $\text{LM2}_{\text{I}}/\text{LM2}_{\text{II}}$  ( $\text{LM3}_{\text{I}}$ ,  $\text{LM3}_{\text{II}}$  and  $\text{LM3}_{\text{III}}$ ) and another one is formed through the bridging of two Ge dimer atoms via S adatom ( $\text{LM3}_{\text{IV}}$  and  $\text{LM3}_{\text{V}}$ ). In the first type, the transition states  $\text{TS3}_{\text{I}}$ ,  $\text{TS3}_{\text{II}}$ , and  $\text{TS3}_{\text{III}}$  are having the energy barriers of 1.26, 1.71, and 1.24 eV, respectively between the  $\text{LM2}_{\text{I-II}}$  and  $\text{LM3}_{\text{I-III}}$ . Also, the S adatom remains bonded with Ge atom and has a dangling bond in  $\text{LM3}_{\text{I-III}}$ . The final products,  $\text{FS}_{\text{I-III}}$  are formed from  $\text{LM3}_{\text{I-III}}$  via  $\text{TS4}_{\text{I-III}}$  with respective energy barriers of 0.48, 1.20, and 0.58 eV.  $\text{LM3}_{\text{I-III}}$  can further transform with small barriers to the final products,  $\text{FS}_{\text{I-III}}$ . In  $\text{FS}_{\text{I-III}}$ , the S adatom bridges either two neighboring dimers along [001] direction or two Ge atoms of a dimer or two neighboring dimers along [010] direction,

respectively. In LM3<sub>IV</sub> and LM3<sub>V</sub>, the S adatom bridges two Ge atoms of a dimer (Figures 4i and 4j) and their formation barriers by corresponding transition states, TS3<sub>IV</sub> and TS3<sub>V</sub> are 0.47 eV and 0.30 eV, respectively. Further, LM3<sub>IV</sub> and LM3<sub>V</sub> isomerized to form the final product FS<sub>IV</sub> (Figure 4n), with high barriers (1.13 and 1.43 eV, respectively. From Table 2, it has been noticed that the dimeric Ge-Ge bond length remains nearly at  $2.56 \pm 0.06 \text{ \AA}$  in all the intermediates and final products, except LM3<sub>IV</sub>, LM3<sub>V</sub>, FS<sub>I</sub> and FS<sub>IV</sub>. In the FS<sub>I</sub>, the Ge(1)-Ge(2) bond length increases to  $2.74 \text{ \AA}$  because the Ge-S-Ge has ring stress. Also, the Ge(1)-Ge(2) bond is cleaved in LM3<sub>IV</sub>, LM3<sub>V</sub> and FS<sub>IV</sub>, so the bond length increases to  $3.5 \text{ \AA} \sim 3.8 \text{ \AA}$ .

There are four reaction paths for the dissociation of H<sub>2</sub>S on Ge(100) surface are proposed and they are; H<sub>2</sub>S<sub>ad</sub> → LM1<sub>I</sub>/LM1<sub>II</sub> → LM2<sub>I</sub> → LM3<sub>I</sub> → FS<sub>I</sub>, H<sub>2</sub>S<sub>ad</sub> → LM1<sub>II</sub>/LM1<sub>III</sub> → LM2<sub>II</sub> → LM3<sub>II</sub> → FS<sub>II</sub>, H<sub>2</sub>S<sub>ad</sub> → LM1<sub>II</sub>/LM1<sub>III</sub> → LM2<sub>II</sub> → LM3<sub>III</sub> → FS<sub>III</sub>, and H<sub>2</sub>S<sub>ad</sub> → LM1<sub>II</sub> → LM2<sub>I</sub>/LM2<sub>II</sub> → LM3<sub>IV</sub>/LM3<sub>V</sub> → FS<sub>IV</sub>. However from the calculated results of the rate-determining steps (RDS) for the final products (FS<sub>I</sub>, FS<sub>II</sub>, FS<sub>III</sub> and FS<sub>IV</sub>) are LM2<sub>I</sub> → LM3<sub>I</sub>, LM2<sub>II</sub> → LM3<sub>II</sub>, LM2<sub>II</sub> → LM3<sub>III</sub> and LM3<sub>V</sub> → FS<sub>IV</sub>, respectively. Their barriers at the rate-determining step are 1.26, 1.71, 1.24 and 1.43 eV, respectively. Among the four, the third RDS has the lowest energy barrier. According to these calculations, FS<sub>III</sub> is thermodynamically more stable than the others and hence it can be the major product. The more ring strains in FS<sub>I</sub> and FS<sub>IV</sub> due to their inter- or intra- dimer lead to break the dimer bonds, so they are unstable products among the others.<sup>23</sup> From the above, it has been noticed that the final product FS<sub>III</sub> is favorable in terms of both kinetic and thermodynamic points.

From Table 1, it has been found that the calculated adsorption energy values for three possible adsorption conformations of H<sub>2</sub>S on Ge(100) surface are smaller than that of

Si(100) surface. When H<sub>2</sub>S adsorbs on surface, the 3p orbitals (2b<sub>2</sub>, 5a<sub>1</sub>, and 2b<sub>1</sub>) of S atom have more significant overlap with p orbitals of buckled-down Si atom<sup>32</sup> than the d orbitals of buckled-down Ge atom. The strong overlap between the p-orbitals of S and Si atoms result in the large adsorption energies for the adsorption of H<sub>2</sub>S<sub>ad</sub> on Si(100) surface.<sup>32</sup> From Table 2, it is noticed that the relative energies for all the intermediates and final products of adsorption of H<sub>2</sub>S on Si(100) surface are more negative than those on Ge(100) surface, which indicates that those are more stable on the Si surface. Based on the present and previous studies,<sup>32</sup> we summarized the total reaction energies for final products, energy barriers for RDS and RDS pathways from H<sub>2</sub>S adsorbed on Ge or Si surface and are shown in Table 4. As can be seen from this table, FS<sub>II</sub> and FS<sub>III</sub> are thermodynamically favourable products for the H<sub>2</sub>S adsorption on Ge(100) and Si(100), respectively. Their RDSs are from LM2 to LM3 via second dehydrogenation process. In the kinetic point of view, FS<sub>I</sub> and FS<sub>IV</sub> are the favorable products for Si surface and FS<sub>I</sub> and FS<sub>III</sub> are the major products for Ge surface. From the above results, it is noted that the final product FS<sub>III</sub> is thermodynamically and kinetically favorable product for the H<sub>2</sub>S adsorption on Ge surface. From the above results it has been observed that the S-passivation products on Si(100) surface are more stable than Ge(100) surface.

### 3.3. Density of States (DOS)

The electronic structure is intimately related to their fundamental physical and chemical properties. To further understand the adsorption modes of H<sub>2</sub>S on Ge(100), we analyzed DOS of most possible reaction pathway i.e., H<sub>2</sub>S<sub>ad</sub>→LM1<sub>II</sub>→LM2<sub>II</sub>→LM3<sub>III</sub>→FS<sub>III</sub>, which includes five different S conformations such as H<sub>2</sub>S, HS, S adatom and bridged S atoms. The plotted DOS for H<sub>2</sub>S<sub>ad</sub>, LM1<sub>II</sub>, LM2<sub>II</sub>, LM3<sub>III</sub> and FS<sub>III</sub> are shown in Figure 7 and it compares the different S conformations adsorbed on Ge(100) surface. In the Figure 7, the sharp band appears between -10 to -15 eV relative to the Fermi level is attributed from the

3s orbital of S atoms, which is adsorbed on the surface with different configurations. Initially, when the H<sub>2</sub>S species adsorbs on the Ge(100) surface, this band located at -15.86 eV region ( Figure 7a.) and the first and second dehydrogenation shifts this band towards the Fermi level ( from -15.86 to -13.5/-11.8 eV, see Figures 7b and 7c). This upward shift can be explained that the interaction of H<sub>2</sub>S on the surface via dative bond, whereas after the first dehydrogenation, LM1<sub>II</sub> species interact with surface via single bond. Also, the peak at -8.2 eV region (Figure 7a) is due to the p<sub>x</sub> orbital of S atom shifts close to the Fermi level after the first dehydrogenation and finally it overlaps with the p<sub>y</sub> and p<sub>z</sub> orbitals after the second dehydrogenation, which is due to the interaction of S atom with the surface by either dative bond or covalent bond. In LM2<sub>II</sub> and LM3<sub>III</sub>, S atom adsorbed in a similar way and hence, the PDOS for adsorbed S atom resembles same ( at -11.89 and -12.67 eV respectively in Figure 7c and d). Finally, the FS<sub>III</sub> possesses bridged S conformation on Ge(100) surface, so its 3s orbital band shifted downward and close to the HS conformation. Also, the broad band in the region of 0 to -5 eV, shows the bonding interaction of p orbitals of S and Ge surface atoms. These above results demonstrate that DOS spectra can comprehensively explain the interaction of adsorbed S atom and surface Ge atom

### 3.4. Vibrational Frequency Analysis for H<sub>2</sub>S Adsorption on Ge(100)

In order to guide a future surface vibrational spectroscopic study, the vibrational frequencies for adsorbed species on Ge(100) surface are analyzed. The calculated vibrational frequencies for the adsorbed species such as, H<sub>2</sub>S, LM, and FS species are summarized in Table 5. The bands observed at 2577 and 2159 cm<sup>-1</sup> for H<sub>2</sub>S species are due to asymmetric and symmetric S-H stretching modes. The strong interaction between H<sub>b</sub> and Ge (5) atoms (see Figure 6) leads to weaken the S-H<sub>b</sub> bond resulted in redshift of wave number for S-H stretching mode (2159 cm<sup>-1</sup>). The H-Ge stretching bands are observed in the range of 1900-2024 cm<sup>-1</sup>. The calculated S-H and H-Ge stretching modes of LM1<sub>II</sub> are

in agreement with the previous EELS experimental values of  $2532\text{cm}^{-1}$  and  $1967\text{cm}^{-1}$ .<sup>37</sup> For S-Ge stretching bands, LM1<sub>I-III</sub> belong to the first dissociated group which display similar S-Ge stretching vibration on the surface. There is about  $80\text{cm}^{-1}$  difference between the H<sub>2</sub>S<sub>ad</sub> and LM1<sub>I-III</sub> for S-Ge stretching mode, which is due to the change in interaction type, i.e. from dative bond to covalent bond between sulfur atom and Ge surface. The other dissociated intermediates (LM2<sub>I-II</sub> and LM3<sub>I-III</sub>) exhibits blue shift in the IR spectrum, which is due to the increase in the bond order of S-Ge after the second dehydrogenation. The S-Ge stretching bands for LM3<sub>IV-V</sub> and FS<sub>I-IV</sub> are observed at about  $300\text{-}420\text{cm}^{-1}$  which is red shifted due to the interaction of bridged sulfur atom with the Ge atom via covalent bond. As shown in Table 5, the calculated S-Ge vibration frequencies of FS<sub>I</sub> are red shifted (about  $50\text{ to }130\text{cm}^{-1}$ ) compared to other final products, which indicates that the bond of dimeric Ge become longer because of the ring stress from the Ge-S-Ge bonds.

#### 4. Conclusions

We have performed the detailed investigation of the adsorption and dissociation of H<sub>2</sub>S on Ge(100) surface using density functional theory calculations. Our results show that among the three different conformation, H<sub>2</sub>S<sub>ad</sub> is found to be adsorbed stably on the surface. Possible reaction paths are proposed for dehydrogenation between adsorbed H<sub>2</sub>S and S adatoms. In the first step, adsorbed H<sub>2</sub>S partially dissociates and produced HS and H and in the second step, HS further dissociates into S adatom. Further, the EDD contour indicate that the increased electron density between the H atom and Ge(5) atom of a neighboring Ge dimer along [010] which leads to the lower energy barrier (0.10 eV) in the first dehydrogenation for H<sub>2</sub>S<sub>ad</sub> to LM1<sub>III</sub>. The results show that the four final products have more stable conformations via LM1, LM2, and LM3 intermediates. We found that FS<sub>I</sub> and FS<sub>III</sub> are kinetically favored products whereas FS<sub>II</sub> and FS<sub>III</sub> are thermodynamically favored products for the H<sub>2</sub>S adsorption on the Ge surface. In addition, the interactions between S and Ge atoms are identified by DOS spectra. The calculated results of IR spectra are in good agreement with the available experimental EELS data.



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**Supplementary information** Table S1 lists the lattice parameters of Ge obtained with various pseudopotentials. Table S2 displays dimer bond length ( $d_D$ ) and dimer tilting angle ( $\theta$ ) for a clean Ge(100) surface, compared with results from other calculations and experimental data. Figure S1 illustrates the partial DOS of H<sub>2</sub>S adsorbed on the Ge(100) surface. Figure S2 shows a top view of all transition structures on the reaction paths of adsorbed H<sub>2</sub>S.

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**Table 1.** The calculated structural parameters and adsorption energies ( $E_{\text{ads}}$ ) for hydrogen sulphide ( $\text{H}_2\text{S}$ ) adsorption on Ge(100) surface.

System	$d(\text{S-H})$ (Å)	$d(\text{S-Ge})$ (Å)	$d(\text{Ge-Ge})$ (Å)	$\angle\text{HSH}$ (deg)	Ge-Ge tilting angle (deg)	$E_{\text{ads}}$ (eV)	<sup>a</sup> $E_{\text{ads}}$ (eV)
$\text{H}_2\text{S}_{\text{ad}}$	1.35/1.38	2.61	2.58	92.4	16	-0.49	-0.74
$\text{H}_2\text{S}_{\text{ad-1}}$	1.35/1.37	2.64	2.58	91.9	17	-0.40	-0.70
$\text{H}_2\text{S}_{\text{ad-2}}$	1.33/1.39	2.63	2.57	92.5	16	-0.30	-0.65

<sup>a</sup>The calculated adsorption energy for  $\text{H}_2\text{S}$  adsorption on Si(100) surface taken from previous study.<sup>32</sup>

**Table 2.** The calculated structural parameters and relative energies for dissociatively adsorbed H<sub>2</sub>S on Ge(100) surface.

Local minima and final products	$d(\text{Ge-Ge})^a$ (Å)	$d(\text{Ge-Ge})^b$ (Å)	$d(\text{S-Ge})$ (Å)	$d(\text{H-Ge})$ (Å)	$E_{\text{rel}}$ (eV) <sup>d</sup>
LM1 <sub>I</sub>	2.55	2.58	2.28	1.57	-0.89(-0.95)
LM1 <sub>II</sub>	2.51 (2.46 <sup>c</sup> )	2.51 (2.46 <sup>c</sup> )	2.28 (2.18 <sup>c</sup> )	1.55 (1.56 <sup>c</sup> )	-1.03(-1.43)
LM1 <sub>III</sub>	2.60	2.53	2.29	1.54	-0.96(-1.31)
LM2 <sub>I</sub>	2.52	2.58	2.10	1.55	-0.67(-1.05)
LM2 <sub>II</sub>	2.52	2.54	2.11	1.55	-0.67(-1.33)
LM3 <sub>I</sub>	2.56	2.59	2.11	1.55	-0.16(-1.13)
LM3 <sub>II</sub>	2.62	2.48	2.12	1.54	-1.03(-1.45)
LM3 <sub>III</sub>	2.54	2.58	2.11	1.56	-1.19(-1.45)
LM3 <sub>IV</sub>	3.88	2.57	2.24/2.33	1.55	-0.88(-1.46)
LM3 <sub>V</sub>	3.45	2.55	2.21/2.34	1.55	-1.12(-1.59)
FS <sub>I</sub>	2.74	2.73	2.36	1.56	-0.81(-1.38)
FS <sub>II</sub>	2.52	2.47	2.17 (2.16 <sup>c</sup> )	1.54 (1.56 <sup>c</sup> )	-1.27(-2.63)
FS <sub>III</sub>	2.51	2.53	2.33	1.55	-1.34(-2.12)
FS <sub>IV</sub>	3.56	-	2.25/2.28	1.55	-1.21(-1.99)

<sup>a</sup> Bond length of a Ge(1)-Ge(2) dimer which is bonded either with HS or S.

<sup>b</sup> Bond length of a Ge(3)-Ge(4) or Ge(5)-Ge(6) which is bonded with dissociative H.

<sup>c</sup> Taken from Reference<sup>18</sup>

<sup>d</sup> Relative energy for H<sub>2</sub>S on Si(100) surface are given in parenthesis.

**Table 3:** The calculated reaction barriers ( $^{\ddagger}E$ , eV), reaction energies ( $\Delta E$ , eV), and imaginary frequencies (IMF,  $\text{cm}^{-1}$ ) for the transition states of decomposition of  $\text{H}_2\text{S}$  on Ge(100) Surface.

Reaction Path	$^{\ddagger}E$	$\Delta E$	IMF
$\text{H}_2\text{S}_{\text{ad}} \rightarrow \text{TS1}_{\text{I}} \rightarrow \text{LM1}_{\text{I}}$	0.53	-0.89	227 <i>i</i>
$\text{H}_2\text{S}_{\text{ad}} \rightarrow \text{TS1}_{\text{II}} \rightarrow \text{LM1}_{\text{II}}$	0.49	-1.03	661 <i>i</i>
$\text{H}_2\text{S}_{\text{ad}} \rightarrow \text{TS1}_{\text{III}} \rightarrow \text{LM1}_{\text{III}}$	0.10	-0.96	300 <i>i</i>
$\text{LM1}_{\text{I}} \rightarrow \text{TS2}_{\text{I-a}} \rightarrow \text{LM2}_{\text{I}}$	0.90	0.22	492 <i>i</i>
$\text{LM1}_{\text{II}} \rightarrow \text{TS2}_{\text{I-b}} \rightarrow \text{LM2}_{\text{I}}$	0.88	0.36	419 <i>i</i>
$\text{LM1}_{\text{III}} \rightarrow \text{TS2}_{\text{II-a}} \rightarrow \text{LM2}_{\text{II}}$	1.89	0.29	1029 <i>i</i>
$\text{LM1}_{\text{II}} \rightarrow \text{TS2}_{\text{II-b}} \rightarrow \text{LM2}_{\text{II}}$	1.07	0.36	421 <i>i</i>
$\text{LM2}_{\text{I}} \rightarrow \text{TS3}_{\text{I}} \rightarrow \text{LM3}_{\text{I}}$	1.26	0.51	552 <i>i</i>
$\text{LM2}_{\text{II}} \rightarrow \text{TS3}_{\text{II}} \rightarrow \text{LM3}_{\text{II}}$	1.71	-0.36	405 <i>i</i>
$\text{LM2}_{\text{II}} \rightarrow \text{TS3}_{\text{III}} \rightarrow \text{LM3}_{\text{III}}$	1.24	-0.52	274 <i>i</i>
$\text{LM2}_{\text{I}} \rightarrow \text{TS3}_{\text{IV}} \rightarrow \text{LM3}_{\text{IV}}$	0.47	-0.32	128 <i>i</i>
$\text{LM2}_{\text{II}} \rightarrow \text{TS3}_{\text{V}} \rightarrow \text{LM3}_{\text{V}}$	0.30	-0.46	197 <i>i</i>
$\text{LM3}_{\text{I}} \rightarrow \text{TS4}_{\text{I}} \rightarrow \text{FS}_{\text{I}}$	0.48	-0.65	145 <i>i</i>
$\text{LM3}_{\text{II}} \rightarrow \text{TS4}_{\text{II}} \rightarrow \text{FS}_{\text{II}}$	1.20	-0.10	119 <i>i</i>
$\text{LM3}_{\text{III}} \rightarrow \text{TS4}_{\text{III}} \rightarrow \text{FS}_{\text{III}}$	0.58	-0.01	162 <i>i</i>
$\text{LM3}_{\text{IV}} \rightarrow \text{TS4}_{\text{IV}} \rightarrow \text{FS}_{\text{IV}}$	1.13	-0.22	919 <i>i</i>
$\text{LM3}_{\text{V}} \rightarrow \text{TS4}_{\text{V}} \rightarrow \text{FS}_{\text{IV}}$	1.43	-0.09	347 <i>i</i>

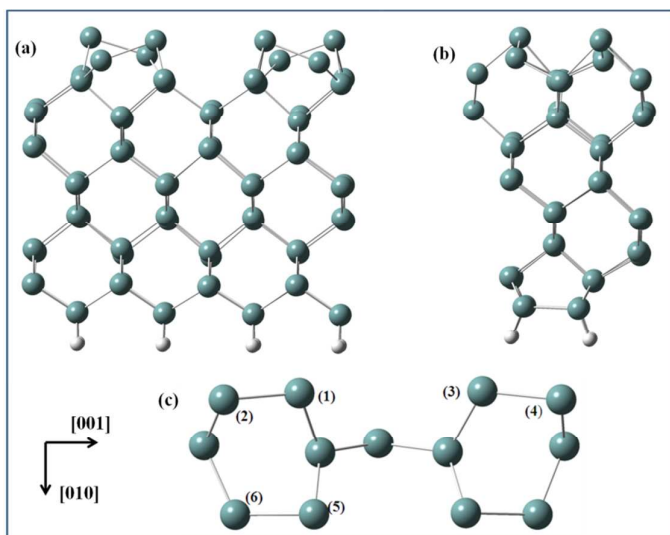
**Table 4.** The calculated total reaction energies ( $\Delta E$  in eV) for the final products, energy barrier ( $^\ddagger E$  in eV) for RDSs and their corresponding pathways for the adsorption of H<sub>2</sub>S on Ge and Si surface.

Final Products	$\Delta E$		$^\ddagger E$		RDS	
	Ge	Si	Ge	Si	Ge	Si
FS <sub>I</sub>	-0.81	-1.38	1.26	1.19	LM2 <sub>I</sub> →LM3 <sub>I</sub>	LM2 <sub>I</sub> →LM3 <sub>I</sub>
FS <sub>II</sub>	-1.27	-2.63	1.71	1.63	LM2 <sub>II</sub> →LM3 <sub>II</sub>	LM2 <sub>II</sub> →LM3 <sub>II</sub>
FS <sub>III</sub>	-1.34	-2.12	1.24	1.54	LM2 <sub>II</sub> →LM3 <sub>III</sub>	LM2 <sub>II</sub> →LM3 <sub>III</sub>
FS <sub>IV</sub>	-1.21	-1.99	1.43	1.22	LM3 <sub>V</sub> →FS <sub>IV</sub>	LM3 <sub>V</sub> →FS <sub>IV</sub>

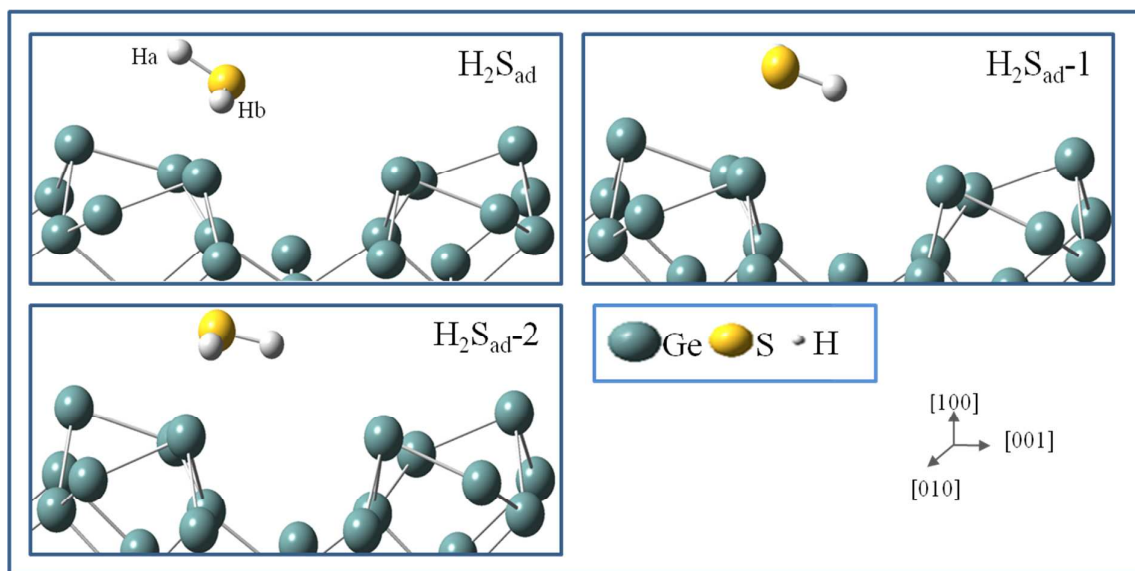
**Table 5:** The calculated vibrational frequencies of the initial state ( $\text{H}_2\text{S}_{\text{ad}}$ ), local minima's (LM), and final products (FS) for the decomposition of  $\text{H}_2\text{S}$  on Ge(100) Surfaces.

System	$\nu$ H-S ( $\text{cm}^{-1}$ )	$\nu$ H-Ge ( $\text{cm}^{-1}$ )	$\nu$ S-Ge ( $\text{cm}^{-1}$ )
$\text{H}_2\text{S}_{\text{ad}}$	2577/2159		265
LM1 <sub>I</sub>	2600	1926	342
LM1 <sub>II</sub>	2544	1967	331
LM1 <sub>III</sub>	2389	2024	348
LM2 <sub>I</sub>		1962/1915	456
LM2 <sub>II</sub>		1979/1968	461
LM3 <sub>I</sub>		1970/1921	463
LM3 <sub>II</sub>		2000/1968	448
LM3 <sub>III</sub>		1948/1927	450
LM3 <sub>IV</sub>		1913/1908	351/310
LM3 <sub>V</sub>		1934/1916	364/295
FS <sub>I</sub>		1914/1903	291/282
FS <sub>II</sub>		1998/1994	421/267
FS <sub>III</sub>		1991/1983	344/276
FS <sub>IV</sub>		1951/1923	396/333

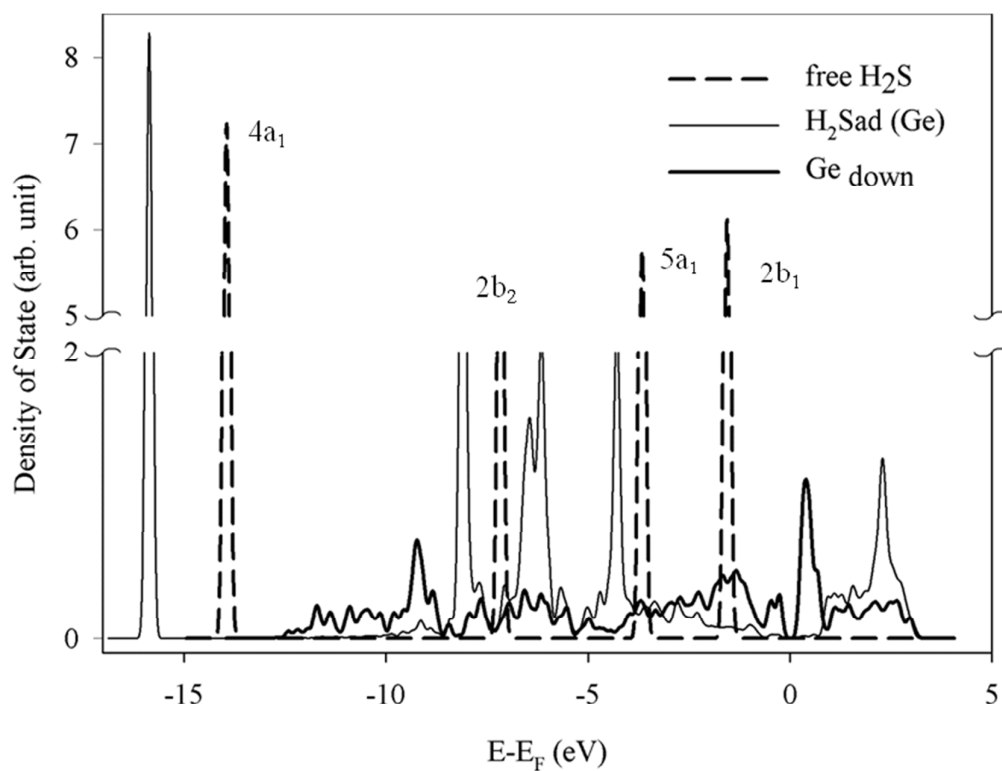




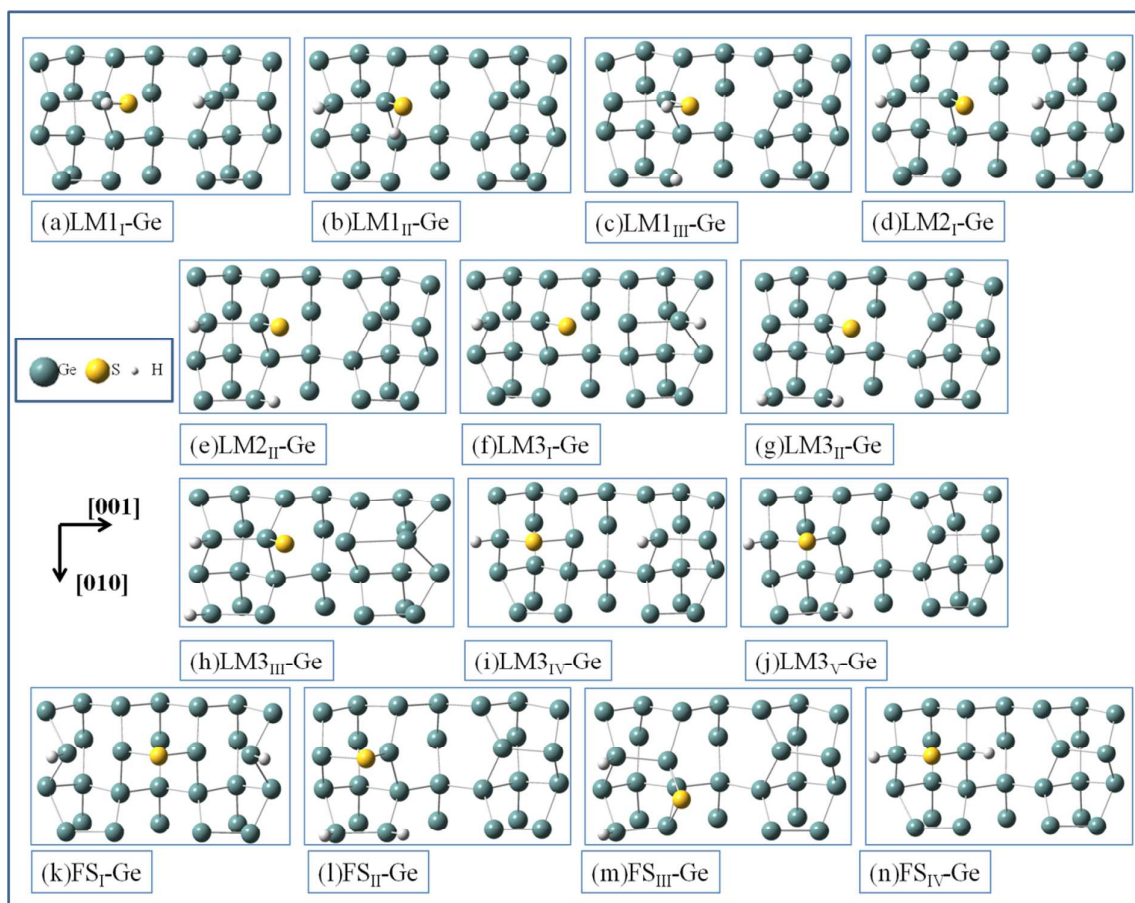
**Figure 1.** The optimized structure of Ge(100)-c(4x2) surface: (a) front view, (b) side view and (c) top view.



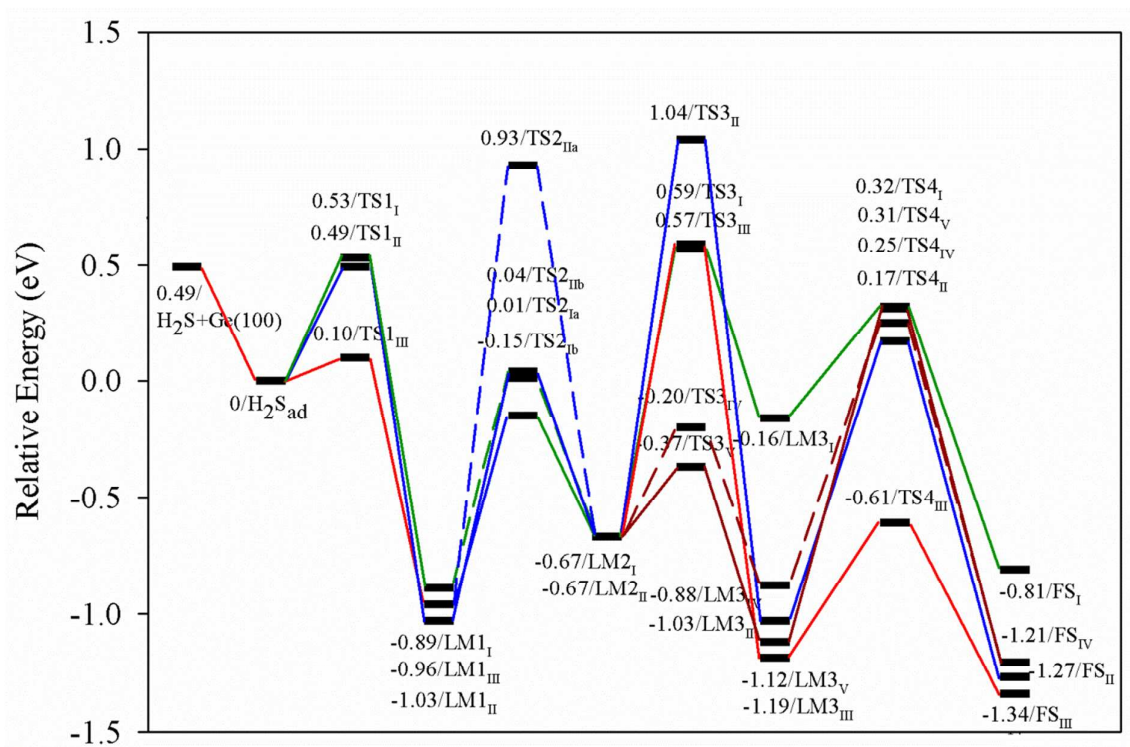
**Figure 2.** The optimized structures of different conformations of  $\text{H}_2\text{S}$  adsorbed on  $\text{Ge}(100)$  surface: (a)  $\text{H}_2\text{S}_{\text{ad}}$ , (b)  $\text{H}_2\text{S}_{\text{ad}-1}$ , and (c)  $\text{H}_2\text{S}_{\text{ad}-2}$ .



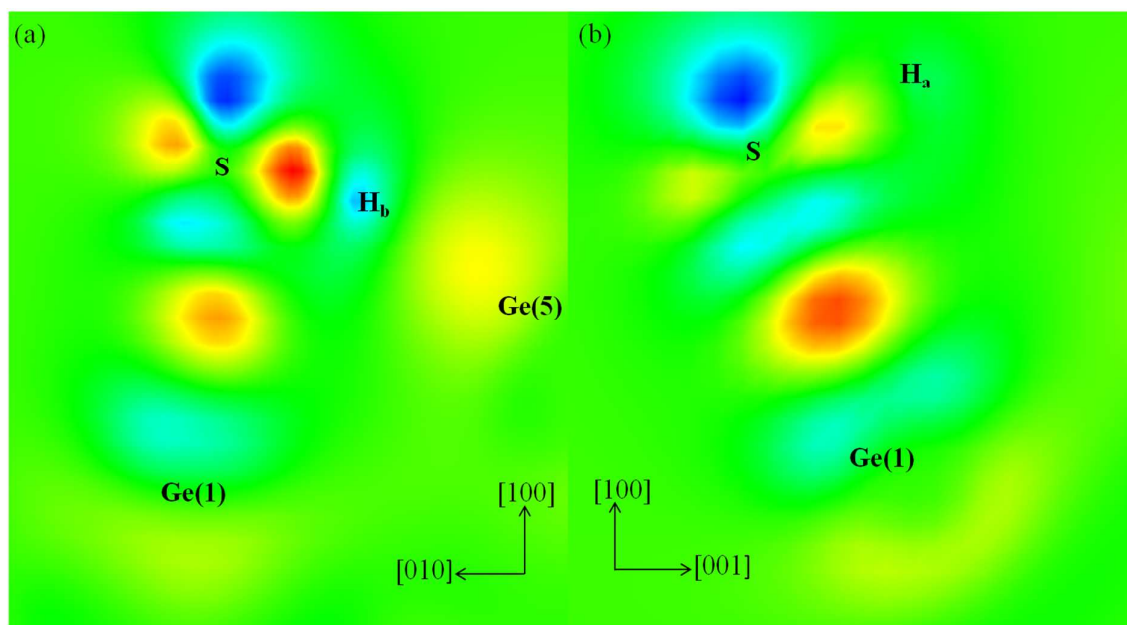
**Figure 3.** The DOS of a H<sub>2</sub>S molecule before (dashed line) and after (black solid line) adsorption on Ge (100) and DOS of Ge down atom (bond black line) bonded with H<sub>2</sub>S.



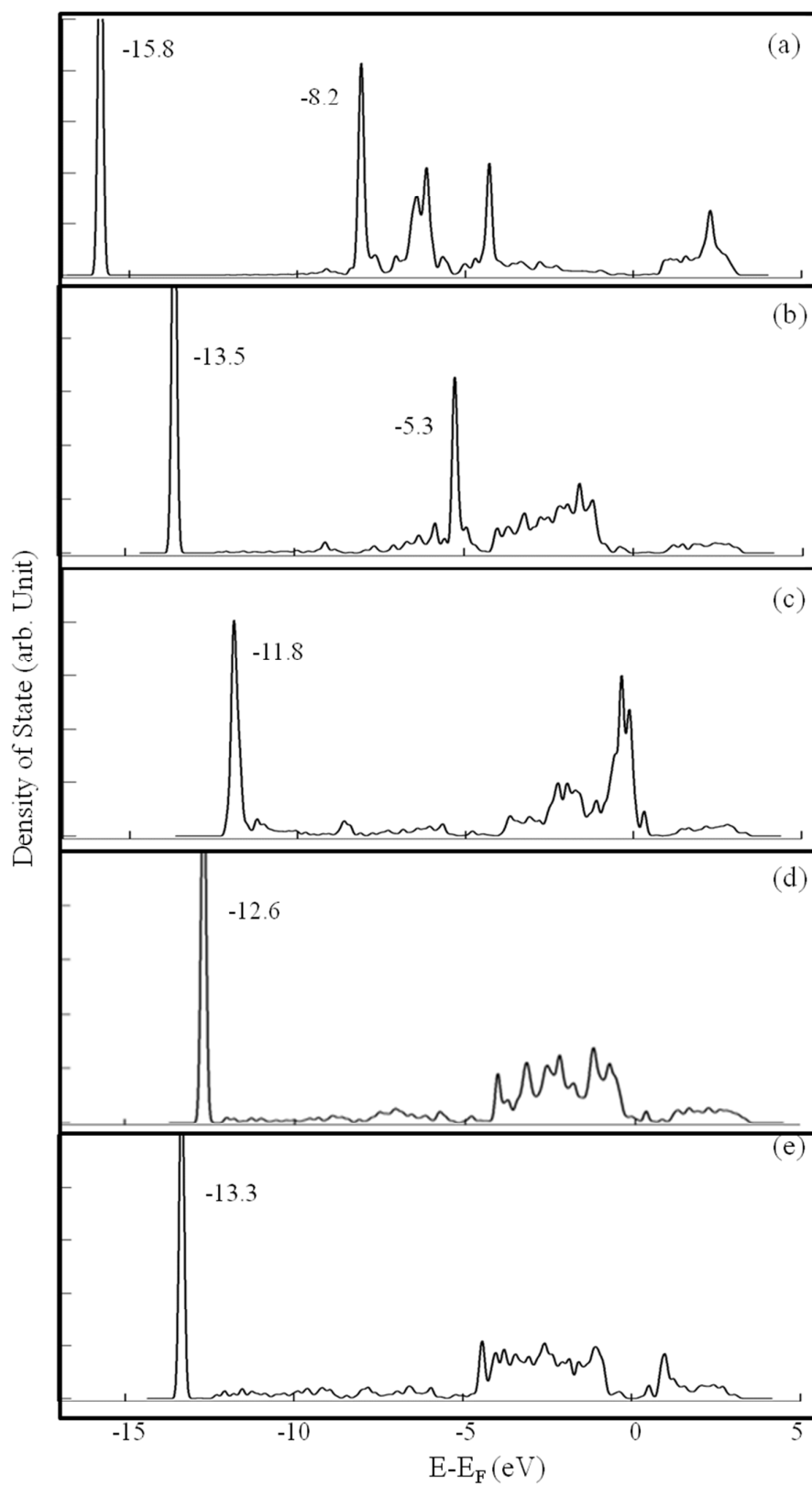
**Figure 4.** Top view of all the local minima's and the final products of the reactions for the adsorption of  $\text{H}_2\text{S}$  on  $\text{Ge}(100)$  surface.



**Figure 5.** Potential-energy diagram for adsorption and decomposition of H<sub>2</sub>S on Ge(100) surface. All potential energies of intermediates, transition structures, and final products are referred to the H<sub>2</sub>S molecule adsorbed on the surface.



**Figure 6.** EDD contours of  $\text{H}_2\text{S}_{\text{ad}}$  on a plane containing (a) a S-Ge(1) bond and a dimeric S-H<sub>a</sub> bond along [010] and (b) a S-Ge(1) bond and a dimeric S-H<sub>b</sub> bond along [001]. The red and blue colors represent the increasing and decreasing electron densities, respectively.



**Figure 7.** The total DOS of S adatom in (a)  $\text{H}_2\text{S}_{\text{ad}}$ , (b)  $\text{LM1}_{\text{II}}$ , (c)  $\text{LM2}_{\text{II}}$ , (d)  $\text{LM3}_{\text{III}}$ , and (e)  $\text{FS}_{\text{III}}$  structures.

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