The confinement inside the novel anionic Sulphobutyletherβ-Cyclodextrinnanocavity, significantly retards the torsional relaxation in Auramine O as compared to native β-CD.
Dynamics under Confinement: Torsional Dynamics of Auramine O in a Nanocavity
Prabhat K. Singh,* Aruna K. Mora, Sushant Murudkar and Sukhendu Nath

The effect of confinement on the ultrafast torsional relaxation dynamics of a well-known ultrafast molecular rotor (UMR) and a recently reported amyloid fibril sensor, Auramine O (AuO), is investigated inside the nanocavity of a novel cyclodextrin derivative, sulphobutylether β-cyclodextrin (SBE-β-CD), using sub-pisosecond fluorescence upconversion spectroscopy. The nanocavity of SBE-β-CD induces a significant increase in the emission intensity and results into a slower transient decay trace of Auramine O when compared to the native β-CD. Detailed analysis of the time-resolved emission spectral features shows that the time-dependent changes in both the mean frequency and the width of the emission spectra are considerably slower as compared to bulk water and native β-CD which suggests that the excited state torsional dynamics of AuO has been significantly effected in the nanocavity of SBE-β-CD. This effect on the torsional dynamics has been attributed to the perturbation of the water structure inside the nanocavity of SBE-β-CD which suppresses the ability of fast collective solvent reorientation motion to promote the excited-state torsional relaxation of Auramine O. The effect of ionic strength of the medium is invoked to analyze the contribution of electrostatic interaction towards the binding of AuO with SBE-β-CD and is corroborated well by computation of electrostatic potentials for the host molecules. The results also suggest that hydrophobic interaction provided by the SBE-β-CD is larger as compared to native β-CD.

Introduction

Ultrafast molecular rotors (UMR) are a special class of fluorescent probe molecules which are endowed with the ability to twist around a single bond constituting a very efficient non-radiative channel in the excited state for the molecule. This twisting motion in the excited state renders this class of molecules virtually non-emissive in low viscosity solvents but when this twisting motion is impeded by the friction of the surrounding microenvironment such as increased micro-viscosity, it leads to an increase in the emission yield of the molecule. This sensitivity towards the microscopic friction forms the basis for this class of molecules to function as viscosity sensor. Apart from sensing the micro-viscosity, molecular rotors when coupled with a specific recognition group can be a target for specific sites in the biological systems and thus can provide information about the immediate surrounding microenvironment. Therefore ultrafast molecular rotors are emerging as sensitive reporters for the micro-viscosity of different biological and chemical environments.

Auramine O (AuO) is a well-known member of this family which has been very recently reported to detect amyloid fibrils, a filamentous form of the protein aggregates responsible for several neurodegenerative disorder such as Alzheimer’s disease and Parkinson’s disease. AuO is almost non-fluorescent in bulk water but in the presence of amyloid fibrils a dramatic enhancement in emission intensity of AuO takes place. It has been established in previous studies carried out in molecular liquids that upon electronic excitation of AuO in solution, the phenyl rings undergo a barrierless rotation to form a non-emissive state of charge transfer character. The relaxed excited state then goes back to the ground state. This large amplitude torsional motion is impeded by medium friction which makes the excited-state reaction rate very sensitive to the viscosity of the environment. The emission quantum yield has been shown to vary linearly with the solvent viscosity. The excited state dynamics of AuO measured through its time resolved emission spectra was modeled by Glassbeek and coworkers as a barrierless diffusive motion of the phenyl groups on an excited state potential energy surface formed through coupling of an emissive locally excited state and a non-emissive charge transfer state. Although with the exception in the case of bulk water where the solvation dynamics appeared to control the torsional rate, this diffusive model proposed by Glassbeek and coworkers was quite successful to account for the effect of viscosity on the torsional rate in a variety of media ranging from molecular solvents such as alcohols to complex confined environments like reverse micelle.

Amyloid fibrils which are insoluble protein or polypeptide aggregates represent very complex biological media offering a wide multitude of local non-covalent interactions such as electrostatic, hydrogen bonding and hydrophobic interaction. In view of the fact that AuO has been recently shown to be very sensitive to the formation of amyloid fibrils, it becomes very relevant to study the excited state dynamics of AuO in somewhat simpler biomimetic confined environments, in particular, the ones which offer possibility of multiple non-covalent interactions.
Knowledge of the molecular dynamics of AuO on the excited state potential energy surface in such biomimetic confined environment is essential to understand the mechanism responsible for their amyloid fibril sensing ability as well as sensing other biomolecules like DNA and enzymes.

An important example of such a confined media is Sulphobutylether-β-cycloextrin (Captisol or SBE-β-CD, Scheme 1) which is a polyanionic β-cyclodextrin derivative with a sodium sulfonate salt separated from the lipophilic cavity by a butyl ether spacer group. Natural cyclodextrins (CDs) are macrocyclic molecules which comprise of elementary glucopyranose units and possess a hydrophobic cavity of varying size depending on the number of glucopyranose units. The Sulphobutylether substituent in SBE-β-CD enhances complexation by providing an extended hydrophobic cavity and an extremely hydrophilic exterior surface as compared to native β-CD. In addition it also provides improved solubility in water and better toxicity profile as compared to native β-CD. The derivative known as SBE-β-CD is an FDA approved pharmaceutical stabilizer and currently used in five commercially approved injectable products. Although there are a few literature reports that primarily focus on the structural information on the host-guest complexes of some potential drug molecules with the SBE-β-CD, to the best of our knowledge, there is no report on the molecular dynamics of the guest molecule in the ultrafast time regime in this potential lipophilic nanocavity.

In this paper, we examine the effect of confinement of nanocavity of the SBE-β-CD on the ultrafast torsional dynamics of AuO using femtosecond fluorescence upconversion technique which is uniquely suited to reveal the information on the excited state dynamics in the sub-ps time regime. Detailed analysis of the temporal evolution of the emission spectral features which include mean frequency, width of the emission spectra (FWHM) and integrated area under the emission spectra has been carried out and compared with that reported in bulk water and native β-CD. The effect of ionic strength of the medium is also performed to analyse the effect of electrostatic interaction towards the binding of AuO with SBE-β-CD.

Scheme-1: Molecular structure of Auramine O and Sulphobutylether-β-cyclodextrin

R= (H)21-n or (CH2)4-SO3Na)n where n=6.2-7.1, CAPTISOL®

Results and Discussion

Ground state absorption measurements

Ground state absorption spectra of AuO were recorded in water and in the presence of various concentrations of macrocyclic host SBE-β-CD and are shown in Figure 1. AuO shows characteristic absorption maxima at 368 nm and 431 nm in water which corresponds to the transition from the ground state to the second and first excited electronic state respectively. On addition of SBE-β-CD to the aqueous solution of AuO, a small but gradual reduction in AuO absorbance with a bathochromic shift of ~10 nm in the absorption maximum was observed. Absorption spectral changes displayed a distinct isobestic point at 381 nm and 439 nm which is suggestive of complexation equilibrium between the free and bound form of AuO. The observed bathochromic shift in the absorption spectra for the dye in SBE-β-CD could be due to the hydrophobic effect inside SBE-β-CD cavity. A small but definitive red shift (~4 nm) has been observed for the dye in the native β-CD. Besides hydrophobic effect, the observed bathochromic shift can also be ascribed to the electrostatic interaction between the cationic Auramine O and the anionic sulphobutyl ether group of SBE-β-CD. Similar bathochromic shift has been observed for the AuO in polyanionic DNA solution and SDS micelle which offers potential sites for the electrostatic interaction with AuO. The contribution of electrostatic interaction towards the binding of AuO with SBE-β-CD has been confirmed by employing the effect of ionic strength of the solution on the photophysical properties of the AuO- SBE-β-CD complex (See later).

Figure 1. Ground-state absorption spectra of Auramine-O at different SBE-β-CD concentrations: (1) 0.0 (2) 2.5 (3) 5 (4) 10.0 (5) 22.0 (6) 110.0 μM.

Steady-state emission measurements: The changes in the emission spectrum of AuO at different SBE-β-CD concentrations have been recorded and presented in figure 2. It is evident from figure 2 that the emission intensity of AuO gradually increases with the increase in the SBE-β-CD concentrations in the solution and an emission enhancement factor of ~ 8 has been obtained when the spectral changes are complete. It should be noted that native β-CD fails to provide any significant change in the emission yield of Auramine O (See Figure S1, ESI). AuO is known to be very weakly fluorescent in bulk water due to the
presence of very efficient non-radiative excited state conformational change leading to the population of a non-fluorescent intermediate state of charge transfer character (TICT) in the solution.\(^9\)\(^{11}\)

In order to evaluate the strength of binding and also the yield of AuO in SBE-\(\beta\)-CD state which further adds to the increase in the emission state fluorescence intensity \(I_f\) at any concentration of the host total dye concentration \([\text{AuO}]_0\) the same, the observed steady-state fluorescence titration curve facilitated by hydrophobic and electrostatic 15 interaction with SBE-\(\beta\)-CD owing to the presence of anionic sulfonato group together with the effect of butyl ether group which extends the hydrophobic cavity of CD leading to stronger binding strength can be ascribed to the additional electrostatic interaction with SBE-\(\beta\)-CD as compared to bulk water whereas it only 75 stronger as compared to the native \(\beta\)-CD. This difference in binding strength can be ascribed to the additional electrostatic interaction with SBE-\(\beta\)-CD owing to the presence of anionic sulfonato group together with the effect of butyl ether group which extends the hydrophobic cavity of CD leading to stronger association of the cationic AuO with SBE-\(\beta\)-CD.

**Figure 2.** (A) Steady-state emission spectra \((\lambda_{ex} = 410\) nm) of Auramine-O (5 \(\mu\)M) at different SBE-\(\beta\)-CD concentrations: (1) 0.0 (2) 2.5 (3) 5 (4) 7.5 (5) 10.0 (6) 14.8 (7) 22.0 (8) 48.0 (9) 110.0 \(\mu\)M. (B) Fluorescence titration curve for the AuO-SBE-\(\beta\)-CD system. The blue circles are the data points and the solid red line is the fitting according to the 1:1 binding model.

Thus, a significant increase in the emission intensity of AuO in the presence of SBE-\(\beta\)-CD can be ascribed to the formation of inclusion complex facilitated by hydrophobic and electrostatic interaction between AuO and SBE-\(\beta\)-CD. Due to the formation of inclusion complex, the free space available for the excited state conformational change in AuO becomes limited leading to retardation of the non-radiative phenyl group twisting process in the excited state. This leads to an enhancement in the emission yield. In addition, the reduced polarity inside the SBE-\(\beta\)-CD cavity will also lower the propensity to form a polar non-emissive TICT state which further adds to the increase in the emission yield of AuO in SBE-\(\beta\)-CD cavity.

In order to evaluate the strength of binding and also the stoichiometric compositions of the inclusion complexes, the variation in the fluorescence intensity of the AuO with the gradually varying SBE-\(\beta\)-CD concentration was analyzed using the 1:1 complexation model which gave satisfactory fitting to the experimental results. In the experimental solution, keeping the total dye concentration \([\text{AuO}]_0\) the same, the observed steady-state fluorescence intensity \(I_f\) at any concentration of the host SBE-\(\beta\)-CD can be expressed as,\(^{27,28}\)

\[
I_f = I_{\text{AuO}}^0 \left[ \frac{[\text{AuO}]_eq}{[\text{AuO}]_0} \right] + I_{\text{AuO-SBE-\(\beta\)-CD}}^\infty \left[ \frac{[\text{AuO} \cdot \text{SBE-\(\beta\)-CD}]_eq}{[\text{AuO}]_0} \right]
\]

where, \(I_{\text{AuO}}^0\) is the initial fluorescence intensity of the free dye, \(I_{\text{AuO-SBE-\(\beta\)-CD}}^\infty\) is the extrapolate fluorescence intensity when all the dye in the solution is converted to AuO-SBE-\(\beta\)-CD complex, \([\text{AuO}]_0\) is the total concentration of dye (AuO) used and \([\text{AuO}]_eq\) is the equilibrium concentration of the free AuO in the solution. For fitting of the fluorescence titration curves, the changes in fluorescence intensity \(\Delta I_f\) were estimated and plotted against the total host concentration \([\text{SBE-\(\beta\)-CD}]_0\) used. Following eq. 1, \(\Delta I_f\) is expressed as,\(^{27,28}\)

\[
\Delta I_f = \left( I_{\text{AuO}}^0 - \frac{[\text{AuO}]_eq}{[\text{AuO}]_0} \right) \left( I_{\text{AuO-SBE-\(\beta\)-CD}}^\infty - I_{\text{AuO}}^0 \right)
\]

where \([\text{AuO}]_eq\) is expressed in terms of the total host concentration, \([\text{SBE-\(\beta\)-CD}]_0\), and the equilibrium constant \(K_{eq}\) for the AuO-SBE-\(\beta\)-CD complex by the following equation.

\[
[\text{AuO}]_{eq} = \frac{[K_{eq}[\text{AuO}]_0 - [\text{SBE-\(\beta\)-CD}]_0 - 1]}{2K_{eq}([\text{SBE-\(\beta\)-CD}]_0 + [\text{SBE-\(\beta\)-CD}]_0 + 1) + 4([K_{eq}]^2[\text{AuO}]_0[SBE-\(\beta\)-CD])_0}
\]

The binding constant \(K_{eq}\) for AuO with SBE-\(\beta\)-CD as obtained from the nonlinear fitting of the titration curve in figure 2B using the equation 2 was thus estimated to be \(9.8 \pm 0.5 \times 10^4\) M\(^{-1}\). The high value of \(K_{eq}\) indicates a very strong interaction of the AuO with the nanocavity of the SBE-\(\beta\)-CD. On the contrary, the binding constant for the complexation of AuO with native \(\beta\)-cyclodextrin is reported to be \(197\) M\(^{-1}\).\(^{29}\) These values suggest that the interaction of AuO with the anionic SBE-\(\beta\)-CD is much stronger as compared to the native \(\beta\)-CD. This difference in binding strength can be ascribed to the additional electrostatic interaction with SBE-\(\beta\)-CD owing to the presence of anionic sulfonato group together with the effect of butyl ether group which extends the hydrophobic cavity of CD leading to stronger association of the cationic AuO with SBE-\(\beta\)-CD.

**Time-resolved measurements:** To understand the effect of confinement on the bond twisting process in the excited state of AuO molecule confined in the SBE-\(\beta\)-CD nanocavity, we have performed detail time-resolved fluorescence measurements using fluorescence upconversion technique. The transient emission decay traces for AuO have been measured at their emission maximum (\(\lambda_{em} = 500\) nm) in water, native \(\beta\)-CD and SBE-\(\beta\)-CD and have been presented in figure 3. It is evident from figure 3 that the emission decay trace becomes much slower in the presence of SBE-\(\beta\)-CD as compared to bulk water whereas it only becomes marginally slower in native \(\beta\)-CD. All transient decays were seen to follow the non-exponential kinetics. The ultrafast non-single exponential decay kinetics in bulk water arise from the relaxation dynamics on the barrier-less excited state potential energy surface where the radiative decay constant of AuO becomes dependent on the phenyl bond twisting angle which...
leads to the emission from different position on potential energy
surface with different time constants and causes the excited state
time to become non-single exponential.11, 14 In the literature,
such non-exponential kinetics has been reported for several other
molecules which are capable of undergoing barrier-less non-
radiative decay processes e.g., bond-twisting, cis-trans
isomerisation etc. in the excited state.28, 30

Figure 3. The transient fluorescence decay (λex = 410 nm, λem = 500 nm)
of Auramine-O in (a) water (b) 10 mM β-CD and (c) 110 μM SBE-β-CD
solution. The dotted line represents the instrument response function
(IRF).

Table 1: Fitting components of the transient emission decay traces of
AuO in different samples at the emission maximum (500 nm)

<table>
<thead>
<tr>
<th>Sample</th>
<th>τ₁/ps</th>
<th>A₁</th>
<th>τ₂/ps</th>
<th>A₂</th>
<th>τ₃/ps</th>
<th>A₃</th>
</tr>
</thead>
<tbody>
<tr>
<td>Water</td>
<td>0.12</td>
<td>0.37</td>
<td>0.35</td>
<td>0.63</td>
<td></td>
<td></td>
</tr>
<tr>
<td>Native</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>β-CD (10 mM)</td>
<td>0.17</td>
<td>0.61</td>
<td>1.3</td>
<td>0.39</td>
<td></td>
<td></td>
</tr>
<tr>
<td>SBE-β-CD (110 μM)</td>
<td>0.12</td>
<td>0.07</td>
<td>0.99</td>
<td>0.20</td>
<td>8.25</td>
<td>0.74</td>
</tr>
<tr>
<td>SBE-β-CD + 7.5 mM NaCl</td>
<td>0.10</td>
<td>0.13</td>
<td>1.05</td>
<td>0.17</td>
<td>7.39</td>
<td>0.70</td>
</tr>
<tr>
<td>SBE-β-CD + 390 mM NaCl</td>
<td>0.08</td>
<td>0.18</td>
<td>0.72</td>
<td>0.23</td>
<td>6.7</td>
<td>0.59</td>
</tr>
</tbody>
</table>

At the emission maximum (500 nm), AuO displays a bi-
exponential decay behaviour in bulk water whereas in the
presence of SBE-β-CD, the decay kinetics becomes tri-
exponential in nature. The longest component (τ₃) arises only in
the presence of SBE-β-CD and can be assigned to the bound form
of the dye in the nanocavity whereas the shorter component (τ₁)
remains invariant in the presence of host which is possibly due to
the free AuO in the solution. The intermediate component (τ₂)
is slightly lengthened in the presence of host as compared to bulk
water. This might possibly be a convoluted response of both the
slow component of bulk water and faster component of complex

Figure 4. The transient fluorescence decay of Auramine-O in 110μM
SBE-β-CD solution at different emission wavelengths: (1) 480 nm (2) 520
nm (3) 570 nm and (4) 610 nm. The dotted line represents the IRF.

Previous time-resolved fluorescence upconversion measurements
of AuO have revealed such complex wavelength dependent decay
kinetics in different media including molecular solvents like
alcohol to complex environment of reverse micelle.12, 14, 17 This
wavelength dependent decay kinetics has been attributed to the
torsional motion of the phenyl groups on a barrierless excited
state potential energy surface. The excited state gradually changes
its electronic properties as a function of the phenyl angle. At
small angle it displays LE character with large oscillator strength
while at larger angle it turns to TICT state with very low
oscillator strength (Scheme 2). Thus the radiative decay constant
of the resulting lower excited state become dependent on the
phenyl bond twisting angle which leads to wavelength dependent
decay of the excited state population. It has been shown that the
rate of torsional relaxation is controlled by the viscosity in the
case of more viscous and slowly relaxing solvents. However bulk
water appears to present a different case where, to a first
approximation, the torsional relaxation rate of AuO is controlled
by aqueous solvation dynamics rather than solvent viscosity.14

which could not be separated here. It is quite obvious from figure
3 and table 1 that the torsional relaxation in the excited state of
AuO is significantly hindered in the nanocavity of SBE-β-CD as
compared to marginal slowdown in the native β-CD. These
results are in agreement to steady-state emission measurements
and binding constant values which suggest that the association of
AuO with SBE-β-CD is much stronger as compared to the native
β-CD. Due to such a strong association of the AuO with SBE-β-
CD cavity the torsional dynamics in the excited state of AuO is
significantly hindered.
The variation in area under the emission curve (normalized to 1 at 50 fs) with time. (a) Bulk water\textsuperscript{25} (b) native \( \beta \)-CD\textsuperscript{25} and (c) SBE-\( \beta \)-CD. The symbols are the data points and the solid line is just connecting the data points (The data for water and native \( \beta \)-CD is reproduced from ref. 25 with permission).

These large changes in the integrated area under the emission spectra for AuO in different media result from the ultrafast non-radiative torsional motion in the excited AuO molecule. The fact that the relatively slower reduction takes place in the integrated area of AuO in SBE-\( \beta \)-CD solution as compared to that in bulk water and native \( \beta \)-CD is a clear manifestation of the increased retardation in the torsional motion inside the SBE-\( \beta \)-CD cavity as compared to that in native \( \beta \)-CD. The slower torsional motion in SBE-\( \beta \)-CD cavity leads to the slow rate of transfer of population from the emissive locally excited state to the non-emissive charge transfer state. It is also evident from figure 5 that the emission spectra shows a gradual red shift with time. Figure 6 shows the variation of mean frequency (first moment) of the spectrum with time. Thus a Stokes shift of \( \sim \)1200 cm\(^{-1}\) has been observed for AuO in SBE-\( \beta \)-CD solution within 10 ps time following the photoexcitation. The observed dynamic Stokes shift for AuO is reported to be due to the diffusive torsional motion of the phenyl group on a barrierless excited state potential energy surface which arises from the adiabatic coupling of the emissive locally excited state and the dark non-emissive charge transfer state. As the twisting proceeds, the excited state population moves along the barrier-less potential energy surface towards the lower energy leading to red shift with time.\textsuperscript{11, 14, 15, 17} Such dynamic Stokes shifts due to different intramolecular processes like bond-twisting, trans-cis isomerization etc. in the excited state have been reported for different chemical systems.\textsuperscript{30, 32-36} Thus, we infer that the observed dynamic Stokes shift for AuO in SBE-\( \beta \)-CD nanocavity also originates from the intramolecular torsional relaxation process operative in the excited state of AuO which causes the lowering in energy of the excited electronic state from the high energy locally excited state to the low energy charge transfer state leading to dynamic Stokes shift.

![Diagram showing potential energy surfaces](Image)

**Scheme 2.** Qualitative potential energy surface for the AuO-SBE-\( \beta \)-CD complex. On photo-excitation, the phenyl ring of AuO undergoes a torsional relaxation on the barrierless potential energy surface to form a TICT state from the initially formed LE state.

The excited state dynamics is best viewed through the time-dependent emission spectra for which wavelength-dependent fluorescence decays were transformed to time-resolved emission spectra (TRES) following the procedure proposed by Maroncelli and Fleming.\textsuperscript{31} The resulting time-dependent spectral profiles are well fitted by the lognormal function (Figure 5). From the lognormal function, different spectral parameters such as mean frequency (first moment), time-dependent integrated intensity and the spectral width can be calculated as a function of time. It is evident from figure 5 that a large reduction in the emission intensity takes place with time. The variation in the integrated area under the emission spectra (shown in the inset of figure 5) for AuO in SBE-\( \beta \)-CD is presented along with the one reported for bulk water and native \( \beta \)-CD.\textsuperscript{25} It is quite evident from figure 5 (inset) that the decrease in integrated area is reasonably slower in SBE-\( \beta \)-CD in comparison to bulk water and native \( \beta \)-CD. In a time window of 3 ps, the integrated intensity for water and native \( \beta \)-CD\textsuperscript{25} is \( \sim \)1 % and \( \sim \)2.5 % respectively whereas in the same time window, the integrated intensity is above 10 % for SBE-\( \beta \)-CD.

![Graph showing time-resolved emission spectra](Image)

**Figure 5.** Time-resolved emission spectra (TRES) of Auramine-O in 110\( \mu \)M SBE-\( \beta \)-CD solution at different times. The circles are the experimental data points and the solid lines are the experimental data points and the solid lines are the lognormal fit to the data points. Inset: The temporal profile for the dynamic Stokes shift is well described by a bi-exponential decay function with a fast component of 0.16 ps (60 %) and a relatively slow component of...

![Graph showing mean frequency with time](Image)

**Figure 6.** The variation in mean frequency with time for Auramine O in 110 \( \mu \)M SBE-\( \beta \)-CD solution. The triangles represent the data points and the solid line is a bi-exponential fit to the data points.
1.8 ps (40 %). The shorter time constant obtained in the present case is very close to that obtained in the case of bulk water and is possibly due to the free AuO present in the solution whereas the slow component of 1.8 ps can be assigned to the AuO which are associated with the nanocavity of SBE-β-CD. It is to be noted that although a nominal slowdown of the temporal decay of mean frequency for the AuO in native β-CD has been reported in literature but even in that case entire Stokes shift dynamics is complete within 1 ps whereas in the case of SBE-β-CD the Stokes shift dynamics has a slow component of 1.8 ps which dominates the decay beyond 1 ps and continues up to 10 ps. Thus the Stokes shift dynamics which is a manifestation of torsional motion of the phenyl group in the barrier-less excited state potential energy surface in AuO is significantly hindered in the SBE-β-CD as compared to the native β-CD and leads to a comparatively slower decrease in energy of the excited-state on the potential energy surface in the presence of SBE-β-CD compared to native β-CD.

Figure 7 presents the variation in the width (Full Width at Half Maximum) of the time dependent emission spectra with time. It is evident from the figure that width of the spectra gradually increases with time. The model used for describing the relaxation dynamics of AuO always predicts an increase in the width of the spectra with time and has been experimentally observed for AuO in several medium including bulk water and reverse micelles. It has also been observed that the growth in the width of the spectra follows a similar kinetics as that of the dynamic Stokes shift. In our case the variation in the width for AuO in SBE-β-CD follows a bi-exponential kinetics in contrast to the mono-exponential kinetics observed in case of bulk water. The time constants obtained in case of SBE-β-CD are 0.17 ps (44%) and 1.55 ps (56%). The shorter time constant obtained in this case is very close to the shorter component observed for the dynamic Stokes shift and is possibly due to the free AuO present in water. Further the slow component in this case is also very similar to the slow component observed for the dynamic Stokes shift and thus this longer component is assigned to originate from the inclusion complex of AuO with SBE-β-CD.

**Figure 7.** The variation in width (full width at half maximum) with time for Auramine O in 110 µM SBE-β-CD solution. The squares represent the data points whereas the solid line is a fit to the data points as mentioned in the text.

It was suggested previously that the excited state reorganisation of AuO is coupled with the solvation dynamics in case of bulk water, and the perturbed water structure at the surfactant-water interface was found to affect the rate of excited state torsional relaxation of AuO at surfactant-water interface. Thus it is logical to imagine that the slow reaction dynamics of AuO in SBE-β-CD nanocavity could be a manifestation of perturbed water structure in the nanocavity. Solvation dynamics in bulk water is driven by collective reorganisation of the H-bond network, so a significant slowdown of the torsional dynamics of AuO in SBE-β-CD cavity is indicative of a significant perturbation of the H-bond network inside the nanocavity. Important information regarding the water structure can be obtained from the molecular dynamics simulation of these systems. Although a direct MD simulation of the water structure in SBE-β-CD cavity is not available, an atomistic MD simulation of aqueous solution of β-CD and its substituted derivatives would be very relevant here. These calculation reveal that the substitution of -OH group strongly influences the dynamical properties of water, in particular those inside the cavities. It has been shown that the translational and reorientational mobility of cavity water is severely diminished with substitution of -OH group of β-CD. Reorientational dynamics of the cavity water in the tri-methyl substituted β-CD derivative was shown to be 4-5 times slower compared to unsubstituted β-CD. The enhanced degree of confinement with the increased degree of substitution of -OH group is also evident from the estimation of average structural relaxation time of cyclodextrin-water (CW) hydrogen bonds and the water-water (WW) hydrogen bonds in the cavity. Average structural relaxation time for CW hydrogen bonds of cavity water molecules in substituted β-CD are 3-4 times longer than the unsubstituted β-CD and is an order of magnitude slower than the bulk water. Interestingly the structural relaxation time of hydrogen bonds formed by the cavity water molecules in unsubstituted β-CD is similar to that of bulk water but the substitution leads to slowdown of the structural relaxation time by a factor of 4-6. Thus in a nutshell these simulations suggest that the enhanced degree of confinement of the cavity water molecules due to substitution of -OH group of β-CD has a stronger influence on the dynamics of these water molecules and the relaxation time scales of hydrogen bonds formed by them. Further the presence of sulphonato co-workers on the relevant SDS-water external interface, which also bears a sulphonated head group, revealed a strong perturbation of the orientational dynamics of water molecules at the external surface specially the first shell which is directly interacting with the sulphonated head group. Thus it is suggested that the perturbation in the water structure inside the nanocavity of SBE-β-CD suppresses the ability of fast collective solvent reorientation motion to promote the excited state reaction of AuO and leads to significant slowdown of the torsional dynamics in the nanocavity of SBE-β-CD.

**Effect of ionic strength**

Several non-covalent intermolecular interactions can contribute
towards the binding of a guest molecule to a supramolecular host. As the host in the present case is anionic due to the presence of sulphoanted groups (SBE-β-CD) and the guest (AuO) is cationic, it is highly probable that electrostatic interaction between the anionic SBE-β-CD and the cationic AuO will also contribute towards the strengthening of the association of the guest with the host. To examine this argument, we have investigated the effect of ionic strength of the medium (sodium chloride) on the complexation of AuO with SBE-β-CD. Figure 8 shows the variation of emission spectra of AuO-SBE-β-CD complex in the presence of various concentration of sodium chloride. It is observed that the emission intensity gradually decreases with increase in the concentration of salt in the solution and levels off at 100 mM of sodium chloride and remains constant beyond that. It can also be seen that the emission intensity never reaches the level observed in bulk water. The decrease in the emission intensity in the presence of salt can be ascribed to the decrease in the equilibrium association constant between the host and the guest because the increase in the ionic strength of the solution is expected to weaken the electrostatic interaction between the host and the guest. However the fact that it never reaches the bulk water behaviour suggests that along with the electrostatic interaction the hydrophobic interaction with the cavity of SBE-β-CD is also a major contributor towards the binding of AuO with SBE-β-CD. It may be reasonable to assume that in the presence of 390 mM NaCl (well beyond the saturation concentration), the remaining emission intensity from the AuO-SBE-β-CD complex is mostly an outcome of hydrophobic interaction. It would be interesting to compare this emission intensity with the AuO in presence of native β-CD in absence of salt. On comparison it turns out that this emission intensity is still much stronger compared to AuO with same concentration of native β-CD without salt. This suggests that the hydrophobic interaction in SBE-β-CD is stronger than the native β-CD. This can be ascribed to the presence of butyl ether groups which extend the hydrophobic cavity of SBE-β-CD as compared to native β-CD and results in more efficient complexation with the former.

To corroborate the steady-state emission measurement we also made ground state absorption measurements for the AuO-SBE-β-CD complex in presence of sodium chloride. Figure 9 shows the ground state absorption spectra of AuO-SBE-β-CD complex in absence and presence of sodium chloride. The absorption spectrum of AuO in bulk water has also been shown for comparison. As stated previously that the absorption maximum is red shifted by 10 nm on complexation with SBE-β-CD but on addition of salt, the absorption maximum moves towards that of bulk water however it does not reach the situation observed in bulk water. At the highest concentration of the salt studied here (390 mM NaCl) the absorption maximum appears at ~ 437 nm which is intermediate between that of bulk water and the AuO-SBE-β-CD complex. This observation is in agreement with the steady-state emission measurements which also show that the addition of salt never leads to complete recovery of the situation prevailing in bulk water.

![Figure 9. Ground-state absorption spectra of Auramine-O in (a) 110 μM SBE-β-CD (solid blue line) (b) in 110 μM SBE-β-CD with 390 mM sodium chloride (black dashed line) and (c) in water (dotted red).](image)

To understand the effect of ionic strength on the excited state dynamics of AuO in SBE-β-CD cavity, we also carried out excited state emission lifetime measurements for AuO-SBE-β-CD complex in presence of salt by monitoring the transient at their emission maximum (500 nm). Figure 10 shows the transient decay traces for the AuO-SBE-β-CD complex in absence and presence of varying concentration of sodium chloride. The transient decay trace in the bulk water has been also shown for comparison. It is observed that transient decay trace for the AuO-SBE-β-CD complex becomes gradually faster with increase in the salt concentration. The fitted parameters for the decay traces have been presented in table 1. It is evident from table 1 that the contribution of the faster components (τ₁) increases whereas the contribution of the slower component (τ₂) decreases with increase in the concentration of the salt in the solution. This is attributed to the weakening of the complexation in the presence of salt. It also becomes obvious that the transient decay trace does not reach the value observed in bulk water. This is in accordance with the ground state absorption and steady-state emission measurements in presence of salt. It can be further observed that the transient decay trace for AuO in SBE-β-CD even in presence of salt (390 mM) is still slower than the AuO in native β-CD suggesting that SBE-β-CD provides a stronger hydrophobic interaction as compared to native β-CD. This is in line with
steady-state emission measurements. These observations imply that multiple interactions with the host such as both electrostatic interaction and hydrophobic interaction contributes towards the stronger binding of AuO with the SBE-β-CD nanocavity. The stronger binding eventually leads to a large retardation of the phenyl group rotation in the excited state resulting in increased emission quantum yield and emission lifetimes as compared to native β-CD which binds to the cationic AuO with a relatively much weaker hydrophobic interaction.

Figure 10. The transient fluorescence decay (λ_{ex} = 410 nm, λ_{em} = 500 nm) of Auramine-O in (a) 110 μM SBE-β-CD solution (b) 110 μM SBE-β-CD solution with 7.5 mM sodium chloride (c) 110 μM SBE-β-CD solution with 390 mM sodium chloride (d) 10 mM β-CD and (e) water The dotted line represents the IFR.

To understand the interaction of AuO with the host better, we have computed the interaction energy (ΔH) for the complexation of AuO with native β-CD, SBE-β-CD (anionic) and SBE-β-CD (neutral) at the semiempirical PM3 level. For this purpose initially the structures of AuO, native β-CD, SBE-β-CD (neutral) and SBE-β-CD (anionic) were optimized. Subsequently the structures of 1:1 complex of AuO with all the hosts were optimized with various input geometries by placing the host along the AuO molecule. From the optimized parameters, the interaction energies of the 1:1 complexation of AuO with native β-CD, SBE-β-CD (neutral) and SBE-β-CD (anionic) is estimated to be 17.1, 64.9 and 197.9 Kcal/mole. The interaction energy obtained for the AuO-β-CD complex is comparable to the value reported previously for the β-CD inclusion complexes. The higher value of interaction energy for the SBE-β-CD (neutral) in comparison to native β-CD might be attributed to the additional hydrophobic interaction owing to the presence of butyl ether group in neutral SBE-β-CD as compared to native β-CD. The higher host-guest interaction energy for the anionic SBE-β-CD in comparison to neutral SBE-β-CD suggests the presence of additional electrostatic interaction between anionic SBE-β-CD and cationic AuO. This is further supported by the computation of electrostatic potential surfaces for the β-CD and SBE-β-CD. ESP surfaces (Fig. 11B) illustrates that the exterior of the SBE-β-CD surface is significantly negative, as expected. Further the interior of the SBE-β-CD cavity is also found to be more negative as compared to the native β-CD. This clearly explains the strong affinity of SBE-β-CD for the cationic AuO. Thus these calculations support our experimental finding that the complexation of AuO with SBE-β-CD is stronger compared to native β-CD owing to additional hydrophobic and electrostatic interaction.

Conclusions

The excited state torsional relaxation dynamics of Auramine O has been investigated inside the nanocavity of a novel cyclodextrin derivative, Sulfobutylether β-CD (SBE-β-CD) using the femtosecond fluorescence upconversion spectroscopy. AuO forms a 1:1 inclusion complex with the SBE-β-CD with a strong affinity (K_{eq}=9.8 \times 10^4 M^{-1}) as compared to native β-CD (K_{eq}=197 M^{-1}) and has been ascribed to the additional electrostatic interaction along with strong hydrophobic interaction of AuO with the former. This strong inclusion complexation affects the excited state torsional relaxation dynamics to a significant extent leading to a large emission enhancement of AuO in presence of SBE-β-CD as compared to native β-CD. The transient decay trace was also found to be significantly slowed down in case of SBE-β-CD as compared to native β-CD. The analysis of temporal evolution of emission spectra reveals the presence of dynamic Stokes shift as well as variation in the spectral width with time. The decay of the mean frequency, growth of the spectral width and the decay of the integrated area under the emission spectra is found to be slower for AuO in the nanocavity of SBE-β-CD as compared to bulk water and the native β-CD. These observations substantiate the distinct effect of confinement on the excited state torsional relaxation dynamics of AuO inside the nanocavity of SBE-β-CD. The perturbation of water structure inside the nanocavity has been suggested to be the
Experimental Section

Auramine O (AuO) was purchased from Sigma-Aldrich as the chloride salt of the dye and was purified by several sublimation steps. Sulphobutylether-β-cyclodextrin (CAPTISOL®), average degree of sulphobutyl substitution: seven; average MW 2162) (SBE-β-CD) was kindly supplied by CyDex Pharmaceutical (La Jolla, California, USA). Sodium chloride was obtained from the Sigma-Aldrich. Nanopure water (conductivity less than 0.1 µS cm⁻¹), from a Millipore Milli Q system, was used for all sample preparations. All measurements were carried out using freshly prepared solution of the dye in SBE-β-CD.

Ground state absorption measurements were performed in JASCO model V650 spectrophotometer. Steady-state fluorescence measurements were made in a Hitachi spectrofluorimeter, model F-4500. The measured spectra, F(λ), were in wavelength domain and were converted to frequency domain, I(ν), by using the following equation:

\[ \hat{I}(\nu) = \hat{\lambda}^2 I(\lambda) \]  

Time-resolved fluorescence measurements were carried out using a femtosecond fluorescence upconversion instrument (FOG 100, CDP Inc. Russia) which has been described in detail elsewhere. Briefly, a second harmonic laser pulse at 410 nm after frequency doubling of the fundamental output of a Ti-sapphire oscillator (820 nm, 50 fs, 88 MHz) in a BBO crystal was used for the sample excitation. The fluorescence light collected from the sample was overlapped with the residual fundamental laser beam (gate beam) into the BBO crystal. The gate beam passes through an optical delay rail before mixing with the fluorescence in the BBO crystal. The upconverted signal was passed through a bandpass filter to cut the excitation and the gate beam and was then dispersed in a double monochromator. The instrument response function (IRF) was measured through the cross correlation of the excitation and the fundamental laser pulse. The IRF was found to have a Gaussian intensity profile with FWHM of 210 fs. For fluorescence lifetime measurements, polarization of the excitation laser beam was set to the magic angle (54.7°) with respect to the horizontally polarized gate pulse. Each decay was collected at least for 2-3 times to see the reproducibility of the measurements. Sample was taken into a rotating cell with optical path length of 1 mm to avoid the photodecomposition of the sample.

For spectrum reconstruction, fluorescence transients were recorded at 10 nm intervals across the steady-state emission spectrum. All these fluorescence transients were fitted with a multi-exponential function convoluted with the instrument response function using the iterative convolution method. Time-resolved emission spectra were reconstructed following the method proposed by Maroncelli and Fleming. Each reconstructed spectrum was fitted using the lognormal function of the following form:

\[ I(\nu) = a \exp \left[ -\ln(2) \left( \frac{1}{b} \ln \left( 1 + \frac{2h(\nu - \nu_p)}{w} \right) \right) \right] \]  

Where a is the amplitude, \( \nu_p \) is the peak frequency, b is the asymmetry parameter and w is the width parameter.

The full width at Half maximum (FWHM) of the emission spectra is related to the width parameter, w and the asymmetry parameter, b, by the following relation:

\[ FWHM = w \frac{\sin h(b)}{b} \]  

The decay traces are fitted with a multi-exponential function of the following form,

\[ I(t) = I(0) \sum \alpha_i \exp(-t/\tau_i) \]  

All the quantum chemical calculations were performed using a set of Gaussian 03 programme. The structure of SBE-β-CD was constructed by placing four and three sulphobutylether group on the larger and smaller rim respectively of native β-CD. The geometry of Auramine O, β-CD, SBE-β-CD (neutral) and SBE-β-CD (anionic) were optimised using semiempirical methods at the PM3 level. Single point energy calculation was performed at these optimized geometry with the Hartree-Fock (HF) method using 3-21G* basis set to provide the input for the calculation of Electrostatic surface potentials. The electrostatic surface potential were obtained using the cubegen module of Gaussview 3.9 package. The relative potential values were projected onto the color coded form onto the isodensity surface by using the Gaussview 3.9 package. The interaction energy for the complexes were calculated at the semiempirical PM3 level.

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Notes and references

Radiation & Photochemistry Division, Bhabha Atomic Research Centre, Trombay, Mumbai 400 085, INDIA Fax: (+) 91-22-25505151, E-mail: prabhatk@barc.gov.in; prabhatsingh988@gmail.com
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