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PAPER

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Synthesis, structure and magnetic properties of phenylhydroxamate-based coordination clusters†§

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The strategic recombination of preformed coordination clusters in the presence of polymodal bridging ligands has successfully led to the characterisation of five new compounds of structural and magnetic interest. Indeed using the dinuclear complex $[M_2(H_2O)(piv)_4(Hpiv)_4]$ ($M=Co, Ni$; $Hpiv$ =pivalic acid) as starting material and reacting it with phenylhydroxamic acid (H_2pha) has yielded the four tetrametallic coordination clusters $[Co_4(Hpha)_2(piv)_6(Hpiv)_4]$ (**1**), $[Ni_4(Hpha)_2(piv)_6(Hpiv)_2(DMF)_2]$ (**2**), $[Co_4(Hpha)_2(piv)_6(EtOH)_2(H_2O)_2]$ (**3**), $[Ni_4(Hpha)_2(piv)_6(EtOH)_2(H_2O)_2]$ (**4**) and the hexanuclear complex $[Co_6(Hpha)_4(piv)_8(EtOH)_2] \cdot EtOH$ (**5**). All the compounds have been structurally characterised revealing a particular binding mode for the hydroxamate ligand. The study of their magnetic properties has been performed and the modelling of these properties has been done using the appropriate hamiltonians for each compound. The experimental data and their modellings show non-zero spin ground states for compounds **4** and **5**.

Introduction

Molecule-based magnetic materials have attracted considerable interest for the past two decades.¹ Indeed the existence of molecular compounds that can retain their magnetisation below a blocking or a critical temperature or else switch their magnetic state under external stimuli has motivated both chemists and physicists to further explore this area of research and its potential applications in data storage, quantum computing or spintronics.² One particular focus has been made on discrete polymetallic compounds –coordination clusters– where fundamental quantum magnetic phenomena can be observed.³ Convincing results have been obtained and ever-growing efforts are currently being made towards the comprehension of such phenomena, the structuring of the molecular systems and of course the synthesis of novel coordination clusters.^{4,5}

One strategy to achieve the synthesis of paramagnetic coordination clusters lies in the use of specific ligands that can display multiple binding modes through distinct coordination sites and thus promote magnetic exchange between several metal ions. Indeed ligands containing O- and/or N-donors such as carboxylates, poly-ols, poly-azoles or phosphonates to name a few, have efficiently led to numerous paramagnetic clusters⁶ including single-molecule magnets (SMM).^{4a,c,d} To this end, hydroxamic acids constitute excellent candidates for the

preparation of coordination clusters as demonstrated by V. Pecoraro's group work on metallocrowns which includes examples of SMMs.^{7,8} Another known strategy for the preparation of magnetic coordination compounds is based on the use of pre-formed clusters. Used as reactants, specific clusters can be assembled into higher nuclearity complexes or into extended structures, they can be used to incorporate co-ligands or additional metal ions or they can simply represent handy and/or reactive precursors in comparison to commercial sources. There are now several examples of pre-formed carboxylato clusters that have been used either as starting materials or building blocks to generate macromolecular, supramolecular or multidimensional assemblies.⁹ Similarly arrays of topologically SMM-like molecules have been obtained from in-situ preparation of the polymetallic building blocks and further linking with organic or inorganic spacers.¹⁰

Aiming at the synthesis of novel magnetic coordination clusters we have combined the two strategies mentioned above in studying the reactivity of phenylhydroxamic acid, **H₂pha**, with the preformed bimetallic complex, $[M_2(H_2O)(piv)_4(Hpiv)_4]\{M_2piv\}$ ($Hpiv$: trimethylacetic acid, $M=Co^{II}, Ni^{II}$). Surprisingly the **H₂pha** ligand has scarcely been used in the synthesis of coordination clusters and only a limited number of clusters containing paramagnetic early transition metal ions have been reported.¹³ Both the cobalt(II) and nickel(II) pivalate dimers have been successfully used as starting materials to build coordination

clusters.^{9h,9i, 9k,11,12,14,15,16} In particular, the cobalt(II) dimer alone is known to recombine in solution yielding homo- and hetero-valent species with nuclearities ranging from three to fourteen cobalt ions.¹¹ Its use in the presence of additional bridging ligands has also led to single-molecule magnets^{14e} as well as remarkable nanosized clusters.^{14f} Moreover the unique and appealing magnetic properties of the cobalt(II) ion has in recent years motivated numerous research on the synthesis of cobalt-based coordination clusters.^{17,18} Besides, cobalt-based coordination clusters are also receiving considerable attention in water oxidation catalysis.¹⁹

Herein we report the synthesis, the structural characterisation and the magnetic properties of five coordination clusters we have obtained reacting **H₂pha** with either {Co₂piv} or {Ni₂piv}. Coordination of the **H₂pha** ligand has led to a recombination of the dinuclear building blocks which, upon varying solvent and base, has afforded the four tetrametallic clusters [Co₄(Hpha)₂(piv)₆(Hpiv)₄] (**1**), [Ni₄(Hpha)₂(piv)₆(Hpiv)₂(DMF)₂] (**2**), [Co₄(Hpha)₂(piv)₆(EtOH)₂(H₂O)₂] (**3**), [Ni₄(Hpha)₂(piv)₆(EtOH)₂(H₂O)₂] (**4**) and the hexanuclear complex [Co₆(Hpha)₄(piv)₈(EtOH)₂]·EtOH (**5**).

Experimental

Synthesis

All reagents were used as purchased with no further purification. [Co₂(H₂O)(O₂CCMe₃)₄(HO₂CCMe₃)₄], {Co₂piv}, was prepared according to the literature procedure.¹¹ Elemental analysis (%) calculated for C₄₀H₇₈Co₂O₁₇ (*M_r*=948.9 g mol⁻¹): C 50.63, H 8.29. Found: C 50.46, H 8.25. [Ni₂(H₂O)(O₂CCMe₃)₄(HO₂CCMe₃)₄], {Ni₂piv}, was prepared according to the literature procedure.¹² Elemental analysis (%) calculated for C₄₀H₇₈Ni₂O₁₇ (*M_r*=948.4 g mol⁻¹): C 50.65, H 8.29. Found: C 50.44, H 8.24.

Hydrate phenylhydroxamic acid, **H₂pha·1.5H₂O**, was prepared by a modification of the literature procedure.²⁰ Sodium hydroxide (43.7 g, 1.093 mol) dissolved in 100 mL of water was added to an aqueous solution (100 mL) of hydroxylamine hydrochloride (38 g, 0.546 mol). The mixture was stirred for 15 min and then added to an ethanolic solution (200 mL) of ethyl benzoate (41 g, 0.273 mol). The resulting solution was stirred for 72 hrs at 40°C under inert atmosphere and left to cool to room temperature. pH was adjusted to 2 with 37% HCl and a yellow precipitate was obtained after evaporation of the solvents under vacuum. The crude product was dissolved in 250 mL of absolute ethanol and the solution filtered to remove the NaCl precipitate that was formed. Further evaporation of the solvent under vacuum has yielded a white powder that was recrystallised from hot water. Yield: 29g (78%). RMN ¹H (400 MHz, DMSO, 298 K, δ): 7.45 (2 H, t, ¹J = 7.4 Hz, Ph); 7.51 (1 H, t, ¹J = 7.2 Hz, Ph); 7.76 (2 H, d, ¹J = 7.76 Hz, Ph); 11.23 (1 H, s, NH). RMN ¹³C (400 MHz, DMSO, 298 K, δ): 126.82; 128.33; 131.09; 132.73; 164.20. Elemental analysis (%) calculated for C₇H₁₀NO_{3.5} (*M_r*=164.16 g.mol⁻¹): C, 51.2; H, 6.1; N, 8.5. Found: C, 51.1; H, 5.9; N, 8.5. ATR/FT-IR (cm⁻¹): 3283s; 3030w; 2693br; 1642m; 1622w; 1601s; 1554s; 1488m; 1452m; 1433m; 1326w; 1314s; 1302sh; 1185w; 1160m; 1077w; 1039w; 1020w; 930w; 896m; 796m;

786m; 704w; 687s; 674m; 618w; 513s; 486s; 426s; 368s; 332m; 291sh; 282m.

[Co₄(Hpha)₂(piv)₆(Hpiv)₄] (**1**).

H₂pha·1.5H₂O (0.041 g, 0.25 mmol) and {Co₂piv} (0.237 g, 0.25 mmol) were dissolved in MeOH (20 mL) and the mixture was stirred overnight. The resulting purple solution was left to slowly evaporate at room temperature. After eight weeks purple crystals of **1** were collected by filtration and dried in air. Yield: 0.073 g (38% based on Co). Elemental analysis (%) calculated for C₆₄H₁₀₆Co₄N₂O₂₄ (*M_r*= 1523.2 g mol⁻¹): C, 50.4; H, 7.0; N, 1.8; Co, 15.5. Found: C, 50.3; H, 7.0; N, 1.8; Co, 15.3. ATR/FT-IR (cm⁻¹): 3372w; 2960(m); 2931w; 2872w; 2579br; 1686m; 1614s; 1604sh; 1572sh; 1519m; 1480s; 1460w; 1410s; 1357s; 1315m; 1223sh; 1199s; 1157w; 1078w; 1025w; 936sh; 914w; 897w; 872w; 796sh; 786m; 766w; 712m; 697w; 683m; 603m; 574w; 538m; 482m; 412m; 393m; 371w; 302w.

[Ni₄(Hpha)₂(piv)₆(Hpiv)₂(DMF)₂] (**2**).

H₂pha·1.5H₂O (0.041 g, 0.25 mmol) and {Ni₂piv} (0.237 g, 0.25 mmol) were dissolved in DMF (15 mL) followed by the addition of triethylamine (0.036 mL, 0.25 mmol). The mixture was stirred overnight and the resulting green solution left to slowly evaporate at room temperature. After six weeks green crystals of **2** were collected by filtration and dried in air. Yield: 0.086 g (47% based on Ni). Elemental analysis (%) calculated for C₆₀H₁₀₀N₄Ni₄O₂₂ (*M_r*=1464.2 g mol⁻¹): C, 49.2; H, 6.9; N, 3.8; Ni, 16.0. Found: C, 48.5; H, 6.8; N, 4.4; Ni, 16.2. ATR/FT-IR (cm⁻¹): 3326w; 2956m; 2926w; 2906sh; 2869w; 1689sh; 1664m; 1615s; 1586w; 1572w; 1555sh; 1524m; 1481s; 1461w; 1448w; 1416s; 1371w; 1354w; 1324w; 1260w; 1228m; 1208m; 1158m; 1100m; 935w; 915w; 894w; 873w; 797w; 787m; 771w; 713m; 685s; 606m; 587w; 566w; 538w; 504w; 455w; 414m; 401w; 371w; 318m; 275w; 229m.

[Co₄(Hpha)₂(piv)₆(EtOH)₂(H₂O)₂] (**3**).

H₂pha·1.5H₂O (0.041 g, 0.25 mmol) and {Co₂piv} (0.237 g, 0.25 mmol) were dissolved in EtOH (20 mL) followed by the addition of tetrabutylammonium hydroxide (1M in MeOH, 0.25 mL, 0.25 mmol). The mixture was stirred overnight and the resulting purple solution left to slowly evaporate at room temperature. After eight weeks pink crystals of **3** were collected by filtration and dried in air. Yield: 0.071 g (43% based on Co). Elemental analysis (%) calculated for C₄₈H₈₂Co₄N₂O₂₀ (*M_r*= 1242.9 g mol⁻¹): C, 46.4; H, 6.6; N, 2.2;. Found: C, 46.3; H, 6.7; N, 2.3. ATR/FT-IR (cm⁻¹): 3345,84(f); 2959m; 2928w; 2904w; 2868w; 1674w; 1599s; 1566sh; 1539w; 1521w; 1482s; 1457w; 1417s; 1374sh; 1361m; 1324w; 1259w; 1226s; 1207sh; 1157w; 1088m; 1051m; 1026w; 910m; 893m; 805sh; 786m; 708w; 684w; 603m; 556w; 491m; 409m; 392m.

[Ni₄(Hpha)₂(piv)₆(EtOH)₂(H₂O)₂] (**4**).

H₂pha·1.5H₂O (0.041 g, 0.25 mmol) was dissolved in EtOH (10 mL) followed by the addition of triethylamine (0.072 mL, 0.5 mmol) and of a solution of {Ni₂piv} (0.237 g, 0.25 mmol) in dichloromethane (5 mL). The mixture was stirred overnight and the resulting green solution left to slowly evaporate at room temperature. After ten days green crystals of **4** were collected by filtration and dried in air. Yield: 0.101 g (62% based on Ni).

Elemental analysis (%) calculated for $C_{48}H_{82}N_2Ni_4O_{20}$ ($M_r =$ least-squares refinement against F^2 using SHELXL software.²¹

Table 1 Crystallographic data for compounds **1-3** and **5**

	1	2	3	5
Formula ^a	$C_{64}H_{106}Co_4N_2O_{24}$	$C_{60}H_{100}Ni_4O_{22}$	$C_{48}H_{82}Co_4N_2O_{20}$	$C_{74}H_{114}Co_6N_4O_{27}$
$F.W.$ [g mol ⁻¹]	1523.2	1464.2	1242.9	1845.3
Crystal system	Triclinic	Triclinic	Triclinic	Monoclinic
Space group	$P-1$	$P-1$	$P-1$	$C2/c$
a [Å]	12.1420(5)	11.8753(4)	10.9832(2)	26.774(2)
b [Å]	13.3720(5)	13.2048(4)	12.7170(2)	13.963(1)
c [Å]	14.1654(6)	13.8854(5)	13.1082(1)	25.509(2)
α [°]	62.550(2)	62.403(1)	114.643(1)	90
β [°]	80.916(2)	82.791(1)	98.680(1)	102.481(5)
γ [°]	71.263(2)	69.928(1)	106.364(1)	90
V [Å ³]	1932.70(14)	1810.99(10)	1519.21(4)	9311(1)
Z	1	1	1	4
T [K]	200(2)	200(2)	200(2)	200(2)
λ [Å]	0.71073	0.71073	0.71073	0.71073
ρ_{calc} [g cm ⁻³]	1.309	1.343	1.358	1.349
μ (MoK α) [mm ⁻¹]	0.914	1.095	1.141	1.117
Measured reflections	35172	33162	23018	49173
Unique reflections	8772	10630	8871	13563
R_{int}	0.0271	0.0264	0.0188	0.0401
Reflections $I > 2\sigma(I)$	7532	8011	7250	9251
Parameters	540	420	391	590
Restraints	219	0	155	102
R^2 [$I > 2\sigma(I)$]	0.0528	0.0446	0.0309	0.0398
wR^2 [$I > 2\sigma(I)$]	0.1549	0.1185	0.0795	0.0971
GO F	1.042	1.023	1.015	1.001
Largest residuals [eÅ ⁻³]	-0.652 ; 1.175	-0.835 ; 0.765	-0.405 ; 0.536	-0.383 ; 0.442

^a Including solvate molecules. ^b $R_1 = \sum ||F_o| - |F_c|| / \sum |F_o|$; $wR_2 = \left[\frac{\sum (w(F_o^2 - F_c^2))^2}{\sum (w(F_o^2))^2} \right]^{1/2}$

1241.9 g mol⁻¹): C, 46.4; H, 6.6; N, 2.2; Ni, 18.9. Found: C, 46.3; H, 6.7; N, 2.4; Ni, 18.8. ATR/FT-IR (cm⁻¹): 3442w; 2956m; 2929w; 2902w; 2869w; 1603s; 1569w; 1539w; 1523w; 1482s; 1485w; 1417s; 1372w; 1362m; 1324w; 1225s; 1207sh; 1156m; 1090m; 1051m; 1026w; 916w; 893w; 801sh; 787m; 708w; 687w; 607m; 561w; 495m; 419m; 399w.

[Co₆(Hpha)₄(piv)₈(EtOH)₂]-EtOH (**5**).

H₂pha·1.5H₂O (0.041 g, 0.25 mmol) was dissolved in EtOH (10 mL) followed by the addition of triethylamine (0.072 mL, 0.5 mmol) and of a solution of {Co₂piv} (0.237 g, 0.25 mmol) in dichloromethane (5 mL). The mixture was stirred overnight and the resulting purple solution left to slowly evaporate at room temperature. After six days pink crystals of **5** were collected by filtration and dried in air. Yield: 0.095 g (69% based on Co). Elemental analysis (%) calculated for $C_{74}H_{114}Co_6N_4O_{27}$ ($M_r =$ 1845.3 g mol⁻¹): C, 48.1; H, 6.2; N, 3.0; Co, 19.1. Found: C, 47.8; H, 6.3; N, 2.9; Co, 18.9. ATR/FT-IR (cm⁻¹): 3353w; 3261br; 2963m; 2929w; 2901w; 2871w; 1623sh; 1613s; 1573s; 1523w; 1481s; 1458w; 1418s; 1374m; 1359m; 1319w; 1226s; 1154s; 1091w; 1043w; 1025w; 912s; 888w; 787s; 707sh; 685s; 595w; 494w; 422m; 387w.

Physical measurements

Crystallographic data were collected on a Bruker Kappa-APEX II CCD diffractometer for **1**, **2**, **3** and **5** (MoK α , $\lambda = 0.71069$ Å). Crystals were mounted on a Hamilton cryoloop using Paratone-N oil and placed in the cold flow produced with an Oxford Cryocooling device. Partial hemispheres of data –preselected with the APEX 2 software²¹– were collected using ϕ and ω scans (25 s/frame for **1-2** 40 s/frame for **3** and **5**). Integrated intensities were obtained with SAINT and were corrected for absorption with SADABS.²² The structures were solved by direct methods and completed by iterative cycles of ΔF syntheses and full-matrix

Crystallographic data and refinements parameters for **1-3** and **5** are given in Table 1. Crystallographic details are available in CIF format.[‡] CCDC numbers 915714-915717.

The X-ray powder diffraction diagram of **4** was collected on a Philips X'pert Pro diffractometer using CuK α 1 monochromatised radiation ($\lambda = 1.54060$ Å) and equipped with a Pixcel detector. The comparison of the experimental powder pattern of **4** with the calculated one for **3** shows that the compounds are isostructural (Figure S1, ESI).

Magnetic measurements in dc mode were performed on a Quantum Design MPMS SQUID on a crushed crystalline sample restrained in a plastic film, drops of Paratone-N oil were added to prevent torquing of the crystallites.

Data were corrected for the diamagnetism contributions of the samples using Pascal constants. The sample holder and Paratone-N oil diamagnetism were measured and subtracted from the raw data. To model the magnetic properties we used a homemade Mathematica code which establishes the hamiltonian matrix and calculates the partition function allowing the derivation of the physical properties. The best parameters were found using the Neldear-Mead algorithm.²³

¹H and ¹³C NMR spectra were collected on a 400 MHz Bruker Avance spectrometer at 298 K.

ATR/FT-IR spectra were collected on a Bruker TENSOR 27 equipped with a simple reflexion ATR diamond plate of the Harrick MPV2 series.

Results and discussion

Structures and synthesis

The tetranuclear compounds **1-4** [M_4 (Hpha)₂(piv)₆L_b-L_w]₂ (**1**: M=Co^{II}, L_b=L_w=Hpiv; **2**: M=Ni^{II}, L_b=DMF, L_w=Hpiv; **3**: M=Co^{II}, L_b=EtOH, L_w=H₂O; **4**: M=Ni^{II}, L_b=EtOH, L_w=H₂O) are structurally closely related and all adopt a centrosymmetric butterfly topology with two edge-sharing {MO₆} octahedra constituting the butterfly's body positions and two {MO₆} octahedra defining the wingtip positions, each connected to the body via a vertex (Figure 1a).

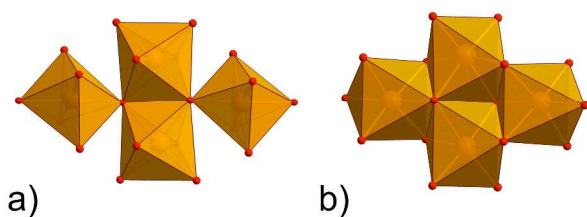


Fig. 1 Polyhedra representation of the butterfly topologies: a) $\{M_4(\mu_3-O)_2\}$ and b) $\{M_4(\mu_3-O)_2(\mu-O,N,X)_4\}$.

In each cluster, two (Hpha)⁻ ligands and six pivalate groups are found bridging the metal ions. The two body sites (M(2) and M(2)#) are bridged to one wingtip position (M(1)) by a μ_3 -oxygen atom from the hydroxamate group and M(1) and M(2) are further bridged by one or two pivalate ligands alternatively (Figure 2). The oxygen atom from the carbonyl group of the phenylhydroxamate ligand binds M(1). In **1** the coordination spheres of the cobalt ions are completed by pivalic acid both on the body and wingtip positions ($L_b=L_w$) where in **2**, DMF and pivalic acid molecules are found on the body and wingtip nickel ions, respectively.

In the isostructural compounds **3** and **4** ethanol act as terminal ligand for the metal ions located on the body positions and water is found on the wingtip metal centres. In compounds **1** to **4** the metal ions are all six-coordinate with distorted octahedral geometries. Charge balance considerations, elemental analysis and BVS calculations support the presence of cobalt(II) ions in **1** and **3**. There are numerous butterfly-type clusters described in the literature although only one example is known with hydroxamate as a constitutive ligand.²⁴ Yet the coordination mode of the ligand differs from the mode observed here with a μ hydroxamate oxygen atom instead of the μ_3 binding mode seen in compounds **1** to **4**. Regarding cobalt(II) and nickel(II) butterflies clusters, distinguishing $\{M_4(\mu_3-O)_2\}$ from $\{M_4(\mu_3-O)_2(\mu-O,N,X)_4\}$ (Figure 1) one can notice that the former structural type—to which **1** to **4** belong—is less common and almost systematically obtained in the presence of bridging carboxylate ligands.^{25,26,27,28} There are no clear signs of intermolecular interactions in the crystal packing of **1** and **2** however the presence of water molecules as terminal ligands in **3** and **4** leads to intermolecular H-bonds with pivalate and hydroxamate oxygen atoms from adjacent clusters (O-H)⋯O distances of 2.72 and 2.88 Å resulting in the formation of chains in the solid (Figure S2, ESI).

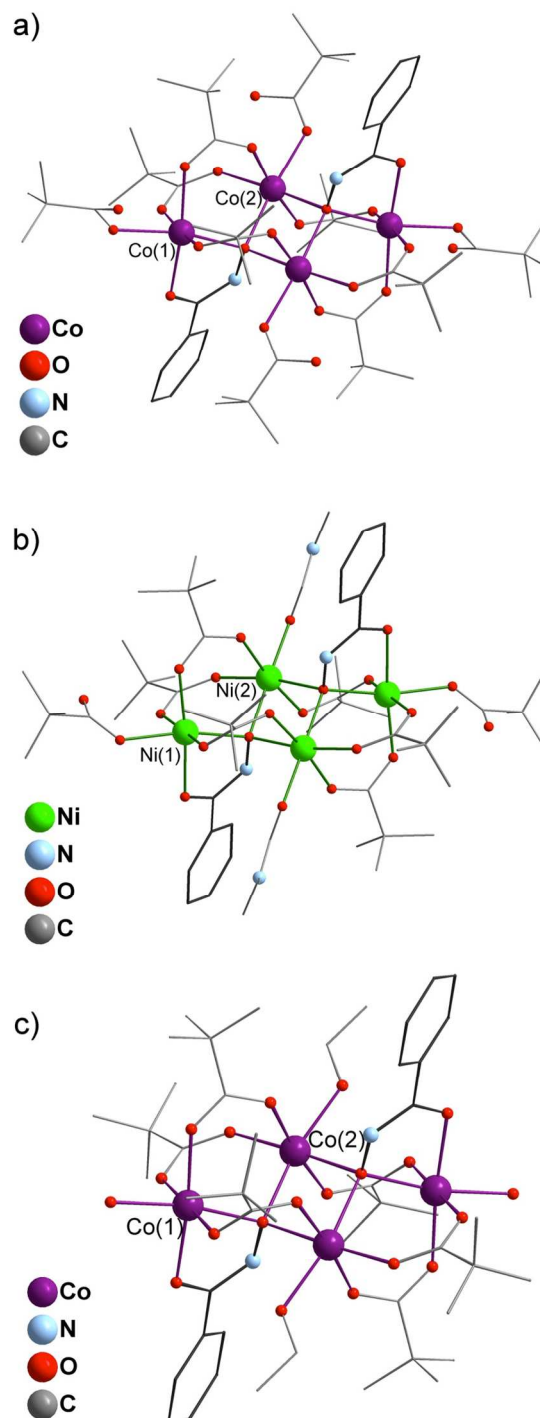


Fig. 2 Structure representations of a) **1**, b) **2** and c) **3**. H atoms have been omitted for clarity and (Hpha)⁻ ligands represented in black.

Compound 5 is a hexametallc cluster made of four bridging hydroxamate ligands and eight bridging pivalate groups. The inorganic skeleton can be described as a butterfly with two additional cobalt centres extending the butterfly wings, formally resulting in the oxo-cluster $\{Co_6(\mu_3-O)_4(\mu-O)_2\}$ (Figure 3).

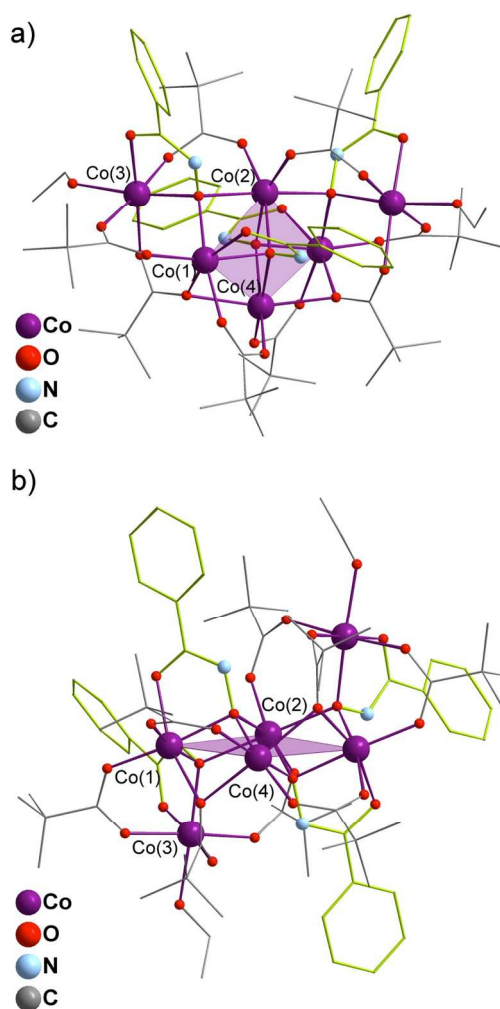


Fig. 3 Structure representations of **5**, a) top view and b) side view. H atoms have been omitted for clarity and (Hpha)⁻ ligands highlighted in lime.

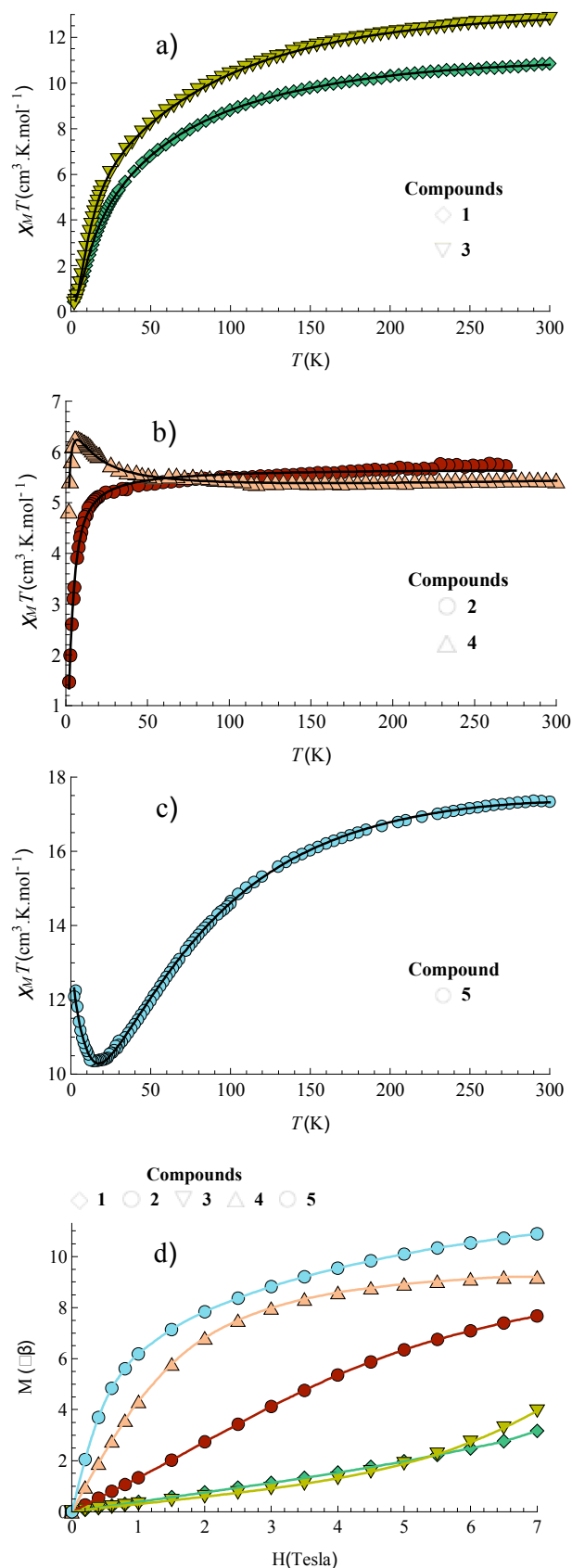
5 The core butterfly is defined by Co(2) and Co(4) on the body positions, Co(1) and its symmetry equivalent being located on the wingtips. The additional cobalt atoms (Co(3) and Co(3)#) are attached to two adjacent edges of the butterfly and situated on each side of its metallic plane (Figure 3b). The four (Hpha)⁻ ligands have similar coordination modes all with μ_3 hydroxamate oxygen atoms but with two distinct environments. One consist in bridging Co(1) with both Co(2) and Co(4) within the internal butterfly and the second has a μ_3 -oxo group binding the external cobalt atoms Co(3) to the butterfly's body- and wingtip-located metal ions Co(1) and Co(2). The connectivity within the cluster is further supported by two kinds of bridging pivalate groups. Six carboxylates adopt a μ_2 : η^1 : η^1 coordination mode connecting Co(1) with Co(3), Co(2) with Co(3) and Co(1) with Co(4) (and symmetry equivalents respectively). The remaining two pivalate groups show a μ_3 : η^2 : η^1 binding mode linking Co(3) with Co(1) and Co(4) (and symmetry equivalents resp.). Oxygen atoms from the carbonyl function of the phenylhydroxamate ligands complete the coordination spheres of Co(1) and Co(3). The latter is also coordinated to a terminal ethanol molecule. All the cobalt ions in **5** are six-coordinate with distorted octahedral geometries. Charge balance considerations, elemental analysis and BVS calculations

support cobalt ions in the +II oxidation state. To the best of our knowledge the topology of **5** has never been encountered to date in early transition metal ions coordination clusters.

30 In analogy to the observations made by Winpenny et al.¹¹ it is likely that the tetrametallic cluster observed for **1** to **4** results from the condensation of two dimeric sub-unit that occurs when substituting the bridging water molecule in the metal-pivalate starting material by the hydroxamate group, hence forming a μ_3 -oxo bridge. Similarly the substitution of the bridging water molecules in the cobalt(II) dimer by the hydroxamate ligands should be driving the condensation of the metal ions in **5**, however arguing whether an intermediate cluster of lower nuclearity is involved in a multi-step self-assembly mechanism seems rather inappropriate with no further solution studies. Structurally all compounds display a μ_3 -hydroxamate ligand which is a striking feature since there are only a limited number of hydroxamate-based coordination compounds displaying such a bonding mode.²⁹ This in turn would explain the relatively high nuclearity of the clusters described here. Indeed the phenylhydroxamate-based coordination clusters reported to date are dinuclear or trinuclear species with μ hydroxamate oxygen atoms only.^{13,30} There is only one recent example of high nuclearity compound based on phenylhydroxamate displaying a μ_3 oxygen atom, a $\{\text{Co}_{16}\}$ cluster.^{13d} This cluster was prepared solvothermally. Compounds **1** to **5** represent then some of the highest nuclearities obtained so far in the presence of phenylhydroxamic acid. However, the incomplete deprotonation of the ligand has most definitely restrained the nuclearity of the formed species. The synthesis of **1** when compared to **2-5** tends to indicate that in the absence of base proton exchange between the carboxylate groups and the hydroxamic acid is sufficient to generate the μ_3 -hydroxamate bridge. A look at the synthesis of **4** and **5** shows that the addition of two equivalents of base per ligand does not lead to the formation of the hydroxamate bridge but results, in the case of cobalt(II), in a slight increase of the hydroxamate to metal ratio.

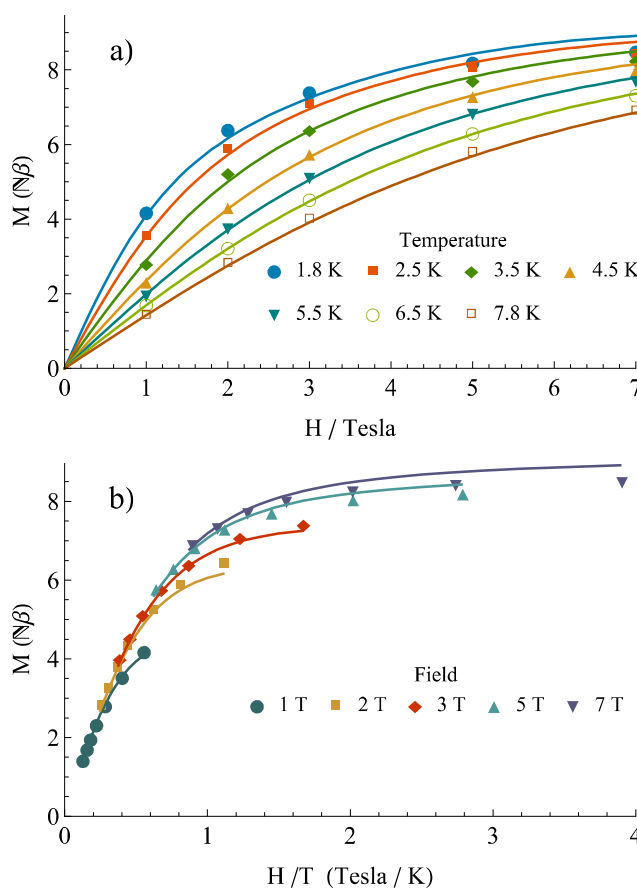
Magnetic properties

DC magnetic susceptibility measurements were performed for **1** to **5** at 1 kOe in the 2-300 K temperature range (Figure 4). At 300 K, $\chi_M T$ values for **1** and **3** are equal to 13.1 cm³Kmol⁻¹ and 12.7 cm³Kmol⁻¹, values compatible with the expected ones for four non-interacting cobalt(II) ions.³² Upon cooling, the $\chi_M T$ products steadily decrease to reach values of 0.6 cm³Kmol⁻¹ and 0.44 cm³Kmol⁻¹ at 2.5 K. At 300 K, **5** has a $\chi_M T$ value of 20.6 cm³Kmol⁻¹ which is in agreement with the expected one for six non-interacting cobalt(II) ions.³²



5 **Fig. 4** a-c) Plots of $\chi_M T$ vs. T measured from 300 to 2 K at 1 kOe for **1** (◆), **2** (●), **3** (▼), **4** (▲) and **5** (●), the solid black lines are the best-fit curves; d) plots of M vs. H measured from 0 to 7 T at 2 K for **1** (◆), **2** (●), **3** (▼), **4** (▲) and **5** (●), the solid lines are a guide for the eyes.

The product continuously decreases when lowering the temperature, reaches a minimum of $10.4 \text{ cm}^3 \text{Kmol}^{-1}$ at ca. 16 K then increases to a local maximum of $12.25 \text{ cm}^3 \text{Kmol}^{-1}$ at 3 K and finally drops to $12.1 \text{ cm}^3 \text{Kmol}^{-1}$ at 2 K. Besides spin-orbit coupling, the rise of $\chi_M T$ at low temperature for **5** tends to indicate the co-existence within the cluster of antiferro- and ferromagnetic interactions. At room temperature the $\chi_M T$ values for **2** and **4** are equal to 5.8 and $5.43 \text{ cm}^3 \text{Kmol}^{-1}$ respectively and are relatively higher than the expected value for four non-interacting nickel(II) ions ($4.8 \text{ cm}^3 \text{Kmol}^{-1}$, assuming $g=2.2$). The $\chi_M T$ product for **2** is almost constant from 300 to 22 K and then rapidly falls close to zero at 2 K. This behaviour would be consistent with weak antiferromagnetic interactions along with zero-field splitting. The $\chi_M T$ curve for **4** shows no significant variation down to 100 K, slowly increases to $6.3 \text{ cm}^3 \text{Kmol}^{-1}$ at 6 K and then drops to 5.3 at 2K, which indicates weak ferromagnetic coupling as well as zero-field splitting.

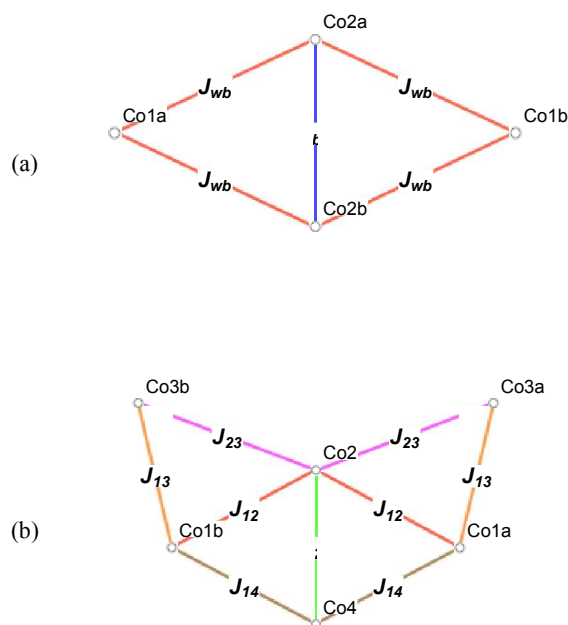


30 **Fig. 5** a) M vs. H measured from 0 to 7 T and 1.8 K to 7.8 K for **4**; b) Isofield M vs H/T from 1 T to 7 T for **4**, the solid lines are the best-fit curves.

The magnetisation curves measured at 2 K support the observations made from the $\chi_M T$ vs. T plots. For **1**, **2** and **3**, the M vs. H curves tend to confirm the existence of moderate to weak antiferromagnetic exchange interactions and no signs of saturation are observed at 7 T (Figure 4d). The magnetisation curve of **4** saturates at $8.5 \text{ N}_\mu\text{B}$ which would confirm the existence of ferromagnetic interactions leading to an $S = 4$ ground state. For **4** M vs. H at various temperature and isofield data have also been measured between 0 and 7 T and between 1.8 K and 8

K (Figure 5). The separation of the isofield curves indicates the presence of zero field splitting or (and) the existence of populated low lying excited states. At 7 T, the magnetisation of **5** reaches 11 N β but no saturation is observed which is coherent with the

presence of anisotropic cobalt(II) ions. Since the pioneering work of Lines,³¹ modelling the magnetic properties of high-spin Co(II) polynuclear complexes is still an open challenge due to the presence of unquenched orbital angular momentum and several approaches have been proposed.³² In the cobalt butterfly compounds **1** and **3** there are two different pathways between the cobalt ions of the wing and the two cobalt ions of the body. To limit overparametrisation only one wing-body interaction has been considered according to scheme 1 (a).



Scheme 1: Topology of interaction in **1** and **3** (a) and in **5** (b)

In this context, the appropriate Hamiltonian to model the magnetic properties of **1** and **3** is: $\mathcal{H}_{tot} = \mathcal{H}_{int} + \mathcal{H}_{so} + \mathcal{H}_{Ze}$ with

$$\begin{aligned} \mathcal{H}_{int} &= -J_{bb} S_{Co2a} \cdot S_{Co2b} \\ &\quad - J_{wb} (S_{Co1a} \cdot S_{Co2a} + S_{Co1a} \cdot S_{Co2b} + S_{Co1b} \cdot S_{Co2a} + S_{Co1b} \cdot S_{Co2b}) \\ \mathcal{H}_{so}^v &= \sum_{i=1}^4 -\frac{3}{2} \alpha \lambda L_{Co_i} S_{Co_i} \\ \mathcal{H}_{Ze} &= \sum_{i=1}^4 \left[\Delta \left(L_{z Co_i}^2 - \frac{1}{3} L_{Co_i}^2 \right) \right. \\ &\quad \left. + \left(-\frac{3\alpha}{2} L_{v Co_i} + g_e S_{v Co_i} \right) \beta \cdot H \right] \end{aligned}$$

with $v = x, y, z$

and where J_{bb} is the coupling constant between the two cobalt ions in the butterfly body positions and J_{wb} the coupling constant between the body and the wing cobalt ions. For the \mathcal{H}_{so} Hamiltonian, λ is the spin-orbit coupling constant, α the orbital

reduction factor and Δ the axial distortion parameter. L et S are respectively the orbital and spin operators with $L = 1$ and $S = 3/2$ in the T-P isomorphism approach³³. To simplify, all the cobalt ions are assumed to be identical and the values of α , λ and Δ are averaged values. However, even within this idealised model, the size of the matrix (20736 x 20736) is too large to have the problem treated exactly in a reasonable amount of time. To overcome this situation we used the perturbational approach developed by Lloret et al^{32b} where the $\chi_M T$ product is expressed

$$\chi_M T = F_{tetra} [G_i(T, J_{ij}, \lambda, \alpha, \Delta), T, J_{ij}]$$

where F_{tetra} is the thermal variation of $\chi_M T$ calculated for a tetranuclear cluster with the same topology than **1** and **3** but with a fictitious $S_{eff} = 1/2$ to mimic the ground state Kramer doublet of a Co(II) ion. To establish the F_{tetra} analytical law it is necessary to take into account a 25/9 scaling factor in the Heisenberg Hamiltonian H_{int} ($S \Rightarrow 5/3 S_{eff}$) leading to the following Hamiltonian:

$$\begin{aligned} \mathcal{H}_{tot} &= \mathcal{H}_{int} + \mathcal{H}_{Ze} \\ \mathcal{H}_{int} &= -\frac{25}{9} J_{bb} S_{Co2a}^{eff} \cdot S_{Co2b}^{eff} \\ &\quad - \frac{25}{9} J_{wb} (S_{Co1a}^{eff} \cdot S_{Co2a}^{eff} + S_{Co1a}^{eff} \cdot S_{Co2b}^{eff} + S_{Co1b}^{eff} \cdot S_{Co2a}^{eff} \\ &\quad \quad + S_{Co1b}^{eff} \cdot S_{Co2b}^{eff}) \end{aligned}$$

$$\mathcal{H}_{ZE} = \sum_{i=1}^4 G_{Co_i}(T, J_{ij}) \beta S_{Co_i}^{eff} H$$

In F_{tetra} the Landé factors of S_{eff} are replaced by the $G_{Co_i}(T, J_{ij}, \lambda, \alpha, \Delta)$ fictitious Landé factors introduced to take into account the influence of the excited states due to spin-orbit coupling and exchange interaction. These G_{Co_i} factors are calculated according to

$$G_{Co_i}^2(T, J_{ij}, \lambda, \alpha, \Delta) = \frac{4k}{N\beta^2} (\chi_M T)_{Co} + f(J_{ij}, T)^2.$$

$(\chi_M T)_{Co}$ is the thermal variation of $\chi_M T_i$ for an isolated anisotropic cobalt ion calculated by full diagonalisation of the hamiltonian matrix and $f(J_{ij}, T)$ is a perturbation term due to excited levels (see SI).

It is extremely delicate to obtain reliable values for a model depending on five parameters. To avoid purely mathematical fits, at least 20 sets of starting parameters were used. All the optimisations lead to a dominant antiferromagnetic J_{wb} interaction around 5 cm^{-1} but to a relatively wide range of J_{bb} values. As J_{wb} determines the magnetic behaviour of the compound and in fact creates an effective ferromagnetic interaction between the two Co(II) ions in the body positions, a reliable determination of the J_{bb} value by magnetic measurements is difficult. Nevertheless, if for **1** the J_{bb} value does not have a great influence on the quality of the fit, there is an improvement by a factor 10 of the agreement factor for **3** when a ferromagnetic J_{bb} is taken into account. It is noteworthy that the need of a ferromagnetic J_{bb} interaction to model the magnetic properties of **3** is coherent with the ferromagnetic interaction observed in the isostructural Ni(II) compound **4**. The final fit for **1** is obtained setting J_{bb} to zero and the least square fit of the magnetic data gave $J_{wb} = -4.8 \text{ cm}^{-1}$, $\lambda = -143. \text{ cm}^{-1}$, $\alpha = 0.78$, $\Delta = 429 \text{ cm}^{-1}$, with $R = 3.1 \cdot 10^{-5}$. The least square fit for **3** gave $J_{wb} = -3.7 \text{ cm}^{-1}$, $J_{bb} = 6.2 \text{ cm}^{-1}$, $\lambda = -152. \text{ cm}^{-1}$, $\alpha = 0.99$, $\Delta = 466 \text{ cm}^{-1}$, with $R = 6.4 \cdot 10^{-6}$ for **3**. Almost equivalent fits are obtained with negative Δ values. The significant antiferromagnetic J_{wb} interaction is not unexpected as

the bridging angles Co(1)-O(1)-Co(2) between the Co(II) ions of the body and the Co(II) ions of the wings are in the 108°-120° range leading to an overlap between the magnetic orbitals. On the other hands, the angles Co(2)-O(1)-Co(2) between the body ions are equal to 99° and 97.4° for **1** and **3** respectively. These values are close to the situation of accidental orthogonality of the magnetic orbitals and should lead to weak ferromagnetic J_{bb} interactions or to J_{bb} values close to zero since such angles also fall in the range where AF contribution can compensate the ferromagnetic one. The spin-orbit coupling constant λ is in a reasonable range for both compounds but the orbital reduction parameter α seems too high for **3**. However it is impossible to obtain a good fit with an α value smaller than 0.9.

The appropriate Hamiltonian to model the magnetic properties of

5 is (scheme 1(b)): $\mathcal{H}_{tot} = \mathcal{H}_{int} + \mathcal{H}_{so} + \mathcal{H}_{Ze}$ with

$$\begin{aligned} \mathcal{H}_{int} &= -J_{12}(S_{Co1a} \cdot S_{Co2} + S_{Co1b} \cdot S_{Co2}) \\ &\quad -J_{13}(S_{Co1a} \cdot S_{Co3a} + S_{Co1b} \cdot S_{Co3b}) \\ &\quad -J_{14}(S_{Co1a} \cdot S_{Co4} + S_{Co1b} \cdot S_{Co4}) \\ &\quad -J_{23}(S_{Co2} \cdot S_{Co3a} + S_{Co2} \cdot S_{Co3b}) - J_{24} S_{Co2} \cdot S_{Co4} \\ \mathcal{H}_{so}^v &= \sum_{i=1}^6 -\frac{3}{2} \alpha \lambda L_{Co i} S_{Co i} \\ \mathcal{H}_{Ze} &= \sum_{i=1}^6 \left[\Delta \left(L_{Z Co i} - \frac{1}{3} L_{Co i}^2 \right) \right. \\ &\quad \left. + \left(-\frac{3\alpha}{2} L_{V Co i} + g_e S_{V Co i} \right) \beta \cdot H \right] \end{aligned}$$

Where J_{ij} are the coupling constants, the other symbols have the same meaning than in the previous formula. As for **1** and **3**, overparametrisation is minimised assuming that all the cobalt ions are identical and thus that the values of α , λ and Δ are averaged values. Nevertheless, due to the size of the hamiltonian matrix (2 985 984 x 2 985 984) it is almost impossible to fit the magnetic data using this model and as for the butterfly compounds we used the perturbational approach developed by Lloret et al³⁴ where the $\chi_M T$ product is expressed has:

$$\chi_M T = F_{hexa} [G_i(T, J_{ij}, \lambda, \alpha, \Delta), T, J_{ij}]$$

where F_{hexa} is the thermal variation of $\chi_M T$ for an hexanuclear cluster with the same topology than **5** but with an effective spin

$S^{eff} = 1/2$ to mimic the ground state Kramer doublet of a Co(II) ion. For **5**, it is not possible to derive an analytical law for the $S = 1/2$ hexanuclear complex and the thermal variation of $\chi_M T$ is obtained by full diagonalisation of the hamiltonian matrix established with the following hamiltonian. $\mathcal{H}_{tot} = \mathcal{H}_{int} + \mathcal{H}_{Ze}$

$$\begin{aligned} \mathcal{H}_{int} &= -\frac{25}{9} J_{12} (S_{Co1a}^{eff} S_{Co2}^{eff} + S_{Co1b}^{eff} S_{Co2}^{eff}) \\ &\quad -\frac{25}{9} J_{13} (S_{Co1a}^{eff} S_{Co3a}^{eff} + S_{Co1b}^{eff} S_{Co3b}^{eff}) \\ &\quad -\frac{25}{9} J_{14} (S_{Co1a}^{eff} S_{Co4}^{eff} + S_{Co1b}^{eff} S_{Co4}^{eff}) \\ &\quad -\frac{25}{9} J_{23} (S_{Co2}^{eff} S_{Co3a}^{eff} + S_{Co2}^{eff} S_{Co3b}^{eff}) - \frac{25}{9} J_{24} S_{Co2}^{eff} S_{Co4}^{eff} \\ \mathcal{H}_{Ze} &= \sum_{i=1}^6 G_{Co i} (T, J_{ij}) \beta S_{Co i}^{eff} H \end{aligned}$$

The mathematical form of the $G_{Co i}(T, J_{ij})$ factors for **5** are given in supplementary informations.

In a first attempt and to limit overparametrisation we fixed λ and

α to sensible values ($\lambda = -160 \text{ cm}^{-1}$ and $\alpha = 0.85$) and find the best values for the other parameters. After this first step we blocked the J_{ij} and Δ parameters to their best values and optimised the λ and α parameters. In an iterative way we repeated these two steps procedure several times until the process converges to stable values. The best fit curve shown in figure 4c is obtained for $J_{12} = -1.3 \text{ cm}^{-1}$, $J_{13} = -0.28 \text{ cm}^{-1}$, $J_{14} = 5.9 \text{ cm}^{-1}$, $J_{23} = -9.8 \text{ cm}^{-1}$, $J_{24} = -0.19 \text{ cm}^{-1}$ and $\Delta = 348. \text{ cm}^{-1}$, $\lambda = -159 \text{ cm}^{-1}$ and $\alpha = 0.84$.

The magnetic behaviour of **5** is determined by two main interactions: J_{14} and J_{23} . The existence of one ferromagnetic interaction, J_{14} , was expected and is consistent with the rise of $\chi_M T$ below 16 K. Magnetically, this polynuclear complex is made up of two weakly coupled trinuclear parts. The first part is comprised of the Co(2) ion and the two Co(3) ions. In this trimetallic entity a large antiferromagnetic interaction is expected between Co(2) and the Co(3) ions due to a large Co(2)-O(1)-Co(3) angle of 120° which allows a very good overlap between the magnetic orbitals. The second part is composed by the Co(4) ion and the two Co(1) ions. The modelling gives a ferromagnetic interaction between Co(4) and the Co(1) ions. This is not unexpected since the values of the bridging angles between Co(1) and Co(4) are equal to 90.1° for Co(1)-O(9)-Co(4) and 95.5° for Co(1)-O(5)-Co(4). These angles are close 90° and lead to accidental orthogonality of the magnetic orbitals and consequently give ferromagnetic interaction between the Co(II) ions. The weak values of J_{12} and J_{13} are more difficult to explain, but for J_{12} the two bridging angles Co(1)-O(1)-Co(2) and Co(1)-O(5)-Co(2) are equal to 99.3° and 103.3° respectively and these values are probably not sufficiently large to give an ample overlap between the magnetic orbitals. For J_{13} the bridging angle Co(1)-O(1)-Co(3) is equal to 109° and relatively large antiferromagnetic interaction is expected. However, these two ions are also bridged by two carboxylate groups and it is well established that such bridges can drastically reduce the antiferromagnetic interaction due to countercomplementarity effects.³⁴

With the lack of first order spin-orbit coupling the two Ni(II) complexes **2** and **4** are easier to model. The Hamiltonian for **2** is the following :

$$\begin{aligned} \mathcal{H} &= -J_{bb} S_{Ni2a} \cdot S_{Ni2b} - J_{wb1} (S_{Ni2a} \cdot S_{Ni1a} + S_{Ni2b} \cdot S_{Ni1b}) \\ &\quad - J_{wb2} (S_{Ni2a} \cdot S_{Ni1b} + S_{Ni2b} \cdot S_{Ni1a}) \\ &\quad + g_{Ni} \beta H (S_{Ni2a} + S_{Ni2b} + S_{Ni1a} + S_{Ni1b}) \end{aligned}$$

where J_{bb} is the body-body interaction and J_{wb1} and J_{wb2} are the two different wing-body interactions. To simplify we have in this Hamiltonian taken the same isotropic g_{Ni} factor in the Zeeman term for all the Ni(II) ions. To avoid overparametrisation we have set $J_{wb1} = J_{wb2} = J_{wb}$ and the least square fit of the magnetic data gave $J_{bb} = 0.96 \text{ cm}^{-1}$, $J_{wb} = -1.3 \text{ cm}^{-1}$ and $g_{Ni} = 2.38$ with $R = 1.6 \cdot 10^{-4}$. An almost equivalent fit is obtained setting J_{bb} to zero ($J_{wb} = -1.28 \text{ cm}^{-1}$, $g_{Ni} = 2.38$ $R = 2.1 \cdot 10^{-4}$) showing that as for compounds **1** and **3** the J_{wb} governs the magnetic behaviour of **2** and that J_{bb} is weak or null. It is worthy of note that, as in the cobalt butterfly complexes, in spite of the large Ni(1)-O(1)-Ni(2) angles (111° and 121°) the antiferromagnetic J_{wb} value is weak. This is most probably due to the non planar conformation of the butterfly structure and also to the presence of carboxylate bridges which could bring a countercomplementarity effect.³⁴

The experimental evidence of a magnetic ground state and the strong likelihood of anisotropy in **4** lead us to add a ZFS term in the Hamiltonian to model the magnetic data. Indeed attempts to fit the magnetisation data by taking into account only exchange interactions without ZFS parameter failed. The anisotropy and Landé factors of all the nickel ions are assumed to be identical and the D_{Ni} and g_{Ni} values are averaged values. To limit overparametrisation we also set $J_{wb1} = J_{wb2} = J_{wb}$ and take isotropic g_{Ni} Landé factor. In this approximation, the appropriate Hamiltonian for compound **4** is:

$$\mathcal{H} = -J_{bb} S_{Ni2a} \cdot S_{Ni2b} - J_{wb} (S_{Ni2a} \cdot S_{Ni1a} + S_{Ni2b} \cdot S_{Ni1b} + S_{Ni2a} \cdot S_{Ni1b} + S_{Ni2b} \cdot S_{Ni1a}) + \sum_{j=1}^4 D_{Ni} (S_{Ni(j)z}^2 - 2/3) + g_{Ni} \beta H (S_{Ni2a} + S_{Ni2b} + S_{Ni1a} + S_{Ni1b})$$

Since only low temperature data give information on the ZFS in the $\chi_M T$ plot we simultaneously fitted the $\chi_M T$ and isofield data to get a more reliable value for D_{Ni} .³⁵ The least square fit of both set of data gives $J_{bb} = 7.0 \text{ cm}^{-1}$, $J_{wb} = 0 \text{ cm}^{-1}$, $D_{Ni} = 3.2 \text{ cm}^{-1}$, $g = 2.27$, with $R = 1.8 \cdot 10^{-4}$. The ferromagnetic interaction between the Ni(II) ions in **4** is not surprising, this compound is isostructural to **3** and the bridging angle Ni(2)-O(1)-Ni(2) between the body ions should be close to 97° , a value in the range of the expected ones to observe a ferromagnetic interaction between Ni(II) ions.³⁶ By contrast, in **2** the bridging angle Ni(2)-O(1)-Ni(2) is larger and it is equal to 99.6° : in the probable range where the weak antiferromagnetic contribution due to the overlap of the magnetic orbitals compensates the ferromagnetic contribution due to exchange integrals leading to a very weak or null value for J_{bb} . The null value of the wing-body interactions is more puzzling but this interaction is already quite small in **2** and is also probably due to the non-planar conformation of the butterfly structure and also to the possible existence of countercomplementarity effects due to the presence of carboxylate bridges.³⁴

Conclusions

This preliminary study shows great potential in the strategic use of pre-formed coordination clusters to react with hydroxamate ligands since it has led to five new coordination clusters. The use of ortho-substituted phenylhydroxamic acid to favour the formation of hydroximate bridges or else the use of polyhydroxamic acid should lead to higher nuclearity clusters. It would also be relevant to study the reactivity of the phenylhydroxamic acid towards coordination clusters of different shape and/or nuclearity. These new synthetic routes are currently explored in our group. The magnetic properties of the clusters described in this paper are less exciting than expected. Indeed none of the compounds show slow relaxation of the magnetisation at low temperature. Yet compounds **4** and **5** present magnetic ground states with some ferromagnetic interactions hinting at the possible preparation of SMMs following the synthetic strategy described here. We also would like to emphasize that it is possible to obtain a satisfactory model of the magnetic properties of coordination clusters containing metal ions with unquenched orbital angular momentum in spite of

the overparametrisation risk. However obtaining reliable results requires caution. It can be fulfilled with a robust optimisation method, multiple tries using several sets of starting parameters and common sense, of course.

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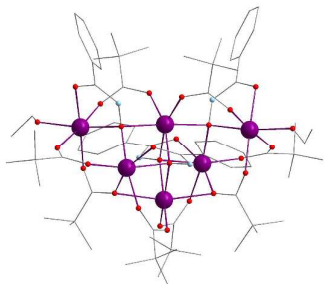
Notes and references

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- † Electronic Supplementary Information (ESI) available: XRD powder data for **3**, additional structural data for **3**, additional magnetic data (χ_M vs. T and $1/\chi_M$ vs. T for compounds **1** to **5**) and informations on the modelling; BVS calculations for **1**, **3** and **5**. See DOI: 10.1039/b000000x/
- ‡ Dedicated to Pr Marius Andruh in celebration of his 60th birthday
- ‡ These data can be obtained free of charge via <http://www.ccdc.cam.ac.uk/cgi-bin/catreq.cgi> (or from the Cambridge Crystallographic Data Centre, 12, Union Road, Cambridge CB2 1EZ, UK; fax: (+44) 1223-336-033; or deposit@ccdc.cam.ac.uk).
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The successful clustering of nickel(II) or cobalt(II) dinuclear complexes into tetra- and hexa-metallic species has been made possible with the use of the phenylhydroxamate bridging ligand. The magnetic properties of these complexes have been studied and modelled.