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on

"Decomposition

trinitroalkyl compounds: a theoretical study from

aliphatic to aromatic nitro compounds" by G. Fayet, P.

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COMMENT

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Comment

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In a recent paper¹ Fayet et al. scrutinized primary decomposition reactions for various derivatives of 1,1,1-trinitrobutane (TNB) and nitroaromatics using DFT computations. The authors compared the calculated Gibbs free energy of activation ($\Delta^{\neq} G_{298}^{0}$) for various competing reaction channels and inferred that the C-NO₂ bond fission dominates the gas-phase thermolysis of all species studied.

For clarity, we will consider the reactions of TNB at room temperature. In contrast to isomerisation channels (reactions R2 and R3 in the manuscript),¹ the C-NO₂ bond rupture (R1) is a barrierless process, and in such case the variational transition state theory (TST) is typically used to locate the transition state (TS).² In the framework of TST, either canonical or variational,² the rate constant of a unimolecular reaction reads as:

$$\begin{aligned} k_{TST} &= \frac{k_B T}{h} exp\left(-\frac{\Delta^{\neq} G_T^0}{RT}\right) = \frac{k_B T}{h} exp\left(\frac{\Delta^{\neq} S_T^0}{R}\right) exp\left(-\frac{\Delta^{\neq} H_T^0}{RT}\right) \\ &= \frac{k_B T}{h} \frac{z^{\neq}(T)}{z^{reag}(T)} exp\left(-\frac{\Delta^{\neq} H_0^0}{RT}\right) \end{aligned}$$
(1).

In my opinion, the methodologies employed by the authors¹ to calculate the ratio of partition functions $\frac{z^{\neq}}{z^{reag}}$ (or, almost equivalently, the activation entropy $\Delta^{\neq} S^0_{298}$) and activation barrier $\Delta^{\neq} H^0_0$ for a barrierless reaction R1 (more generally, the C-NO₂ bond rupture reactions as well), both have serious shortcomings.

Let us consider first $\frac{z^{\neq}}{z^{reag}}$. As an estimation to $\Delta^{\neq} G_{298}^{0}$ the authors¹ simply proposed the Gibbs free energy of a barrierless reaction $\Delta^{r} G_{298}^{0}$, (i.e., a variational TS corresponds to the asymptote •Rad1 + •NO₂ and $z^{\neq} \approx z^{A} z^{B}$). Assuming the usual

ry factorization of the partition function of N-atomic 1- species $z = z_{tr}^{(3)} z_{rot}^{(3)} z_{vib}^{(3N-6)}$, the ratio of approximate and exact TS partition functions reads

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as:

$$\frac{z_{A}z_{B}}{z^{*}} = \frac{z_{tr}^{A(3)} z_{rot}^{A(3)} z_{vib}^{A(3N_{A}-6)} \cdot z_{tr}^{B(3)} z_{vib}^{B(3)} z_{vib}^{B(3N_{B}-6)}}{z_{tr}^{*(3)} z_{rot}^{*(3)} z_{vib}^{*(3)-7}} = \\
= \begin{cases} \left(\frac{z_{tr}^{A(3)} \cdot z_{tr}^{B(3)}}{z_{tr}^{*(3)}}\right) \left(\frac{z_{rot}^{A(3)} \cdot z_{rot}^{B(3)}}{z_{rot}^{*(3)}}\right) \\
z_{vib}^{*(3)} z_{vib}^{*(3)} z_{vib}^{*(3)} z_{vib}^{*(3)-7} z_{vib}^{*(3)-7} z_{vib}^{*(3N_{A}-6)} \cdot z_{vib}^{B(3N_{B}-6)} \\
z_{vib}^{*(3N-12)} z_$$

The ratio $\frac{z_A z_B}{z^{\neq}}$ is dominated by the first term in parentheses. For brevity, we estimate from below the second term in curly brackets (in a very conservative manner) as unity. The third term comprises contributions from approximately conserved vibrational modes and is supposed to be close to unity as well. Assuming $z_{tr}^{eff(3)} = 7.5 \cdot 10^6$ ($m_{eff} = 35$ a.e.m.) and the partition function of every TS vibrational (or hindered rotation) mode $z_{vib/hind.rot}^{\neq(1)} \leq 10$ (corresponds to a harmonic vibration with a wavenumber higher than ~20 cm⁻¹); the lower estimation of the first term in (2) yields $\frac{z_A z_B}{z^{\neq}} \geq 7 \cdot 10^4$.

More generally, the main error stems from contributions from one entirely spurious translational degree of freedom and two vibrations/hindered rotations in the TS erroneously considered to be translations as

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well. In the case of (R1), this yields at room temperature a spurious contribution to $\Delta^{\neq} G_{298}^0$ of $RT \cdot \left(ln \left(\frac{z_A z_B}{z^{\neq}} \right) - 1 \right) \ge 6$ kcal/mol. Such fictitious overestimation of $\Delta^{\neq} S_{298}^0$, inter alia, leads to an unrealistically high value of the preexponential factor $A \approx \frac{k_B T}{h} exp \left(1 + \frac{\Delta^r S_{298}^0}{R} \right) = 10^{24} \text{ s}^{-1}$.

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Table 1. Activation Barriers of Thermal Reactions of Nitroalkanes ($\Delta^{\neq} H_0^0$) Calculated Using Various Coupled Cluster Methods, Multi-Level, and DFT Procedures. In the Case of Barrierless Radical Decomposition Reactions, the C-NO₂ Bond Energies ($\Delta^{r} H_0^0$) are Listed in Parentheses. All Values are in kcal/mol.

Reactions	$\Delta^{\neq} \boldsymbol{H_0^0} \; (\Delta^r H_0^0), \text{kcal/mol}$				
	CCSD(T)/ CBS(T,Q) ^{a,b}	CCSD(T)-F12/ CBS(D,T) ^{a,b}	CCSD(T)-F12/ VDZ-F12 ^b	PBE0 ^c	M06-2X ^d
$C_3H_7C(NO_2)_3 \rightarrow \bullet C_3H_7C(NO_2)_2 + \bullet NO_2(R1)^e$			(44.0)	(35.0)	(44.0)
$C_3H_7C(NO_2)_3 \rightarrow C_3H_6=C(NO_2)_2 + HONO (R3)^e$			42.7	37.9	41.7
$C_3H_7C(NO_2)_3 \rightarrow C_3H_7C(NO_2)_2ONO (R2)^e$			59.5	59.2	62.1
$CH(NO_2)_3 \rightarrow \bullet CH(NO_2)_2 + \bullet NO_2$		(46.5)	(46.0)	(38.9)	(46.1)
$CH(NO_2)_3 \rightarrow C(NO_2)_2N(O)OH$		50.4	50.7	45.7	48.7
$CH(NO_2)_3 \rightarrow CH(NO_2)_2ONO$		61.1	60.5	60.6	62.8
$CH_3CH(NO_2)_2 \rightarrow \bullet CH_3CH(NO_2) + \bullet NO_2$		(48.5)	(47.8)	(41.1)	(47.4)
$CH_3CH(NO_2)_2 \rightarrow CH_2=CH(NO_2) + HONO$		44.4	45.0	41.3	43.6
$CH_2(NO_2)_2 \rightarrow \bullet CH_2(NO_2) + \bullet NO_2$	(49.1)	(49.3)	(48.6)	(44.3)	(48.8)
$CH_2(NO_2)_2 \rightarrow CH(NO_2)N(O)OH$	54.7	54.7	54.8	51.1	53.2
$CH_3CH_2NO_2 \rightarrow \bullet CH_3CH_2 + \bullet NO_2 (NE1)$	(61.0)	(61.1)	(60.2)	(57.5)	(62.1)
$CH_3CH_2NO_2 \rightarrow CH_2=CH_2 + HNO_2 (NE2)$	47.4	47.3	47.8	45.6	47.0

^a The two-point complete basis set extrapolations (aVTZ:aVQZ and VDZ-F12:VTZ-F12, respectively) were employed. ^b Single point CCSD(T)/aug-cc-pVTZ, CCSD(T)/aug-cc-pVQZ, CCSD(T)-F12b/cc-pVDZ-F12, and CCSD(T)-F12b/cc-pVTZ-F12 energies were calculated using the M06-2X/6-311++G(2df,p) optimized geometry. ZPE and thermal corrections to enthalpy were computed at the same level of theory. ^cPBE0/6-31+G(d,p) level of theory was used. ^dM06-2X/6-311++G(2df,p) level of theory was used. ^eThe reactions are named in accordance with ref 1.

Moreover, even higher error to $\Delta^{\neq} G_{298}^{0}$ is introduced by quantum chemical calculations of $\Delta^{\neq} H_{0}^{0.3}$ As we have previously demonstrated,⁴ some DFT functionals remarkably underbind polynitromethanes (e.g., B3LYP lowers D_{0} (C-NO₂) in trinitromethane by ~11 kcal/mol).⁴ PBE0 performs similarly and underestimates D_{0} (C-NO₂) in TNB by ~9 kcal/mol in comparison with CCSD(T)-F12 values (cf. the benchmark results in Table 1). More generally, the results of Table 1 indicate poor performance of PBE0 for decomposition and H-transfer reactions of polynitroalkanes.

Thus, the difference between a true $\Delta^{\neq} G_{298}^{0}$ and reported¹ $\Delta^{r} G_{298}^{0}$ of the reaction (R1) is at least ~15 kcal/mol (*ca.* 11 orders of magnitude in k_{TST} at room temperature). At the same time, the $\Delta(\Delta^{\neq} G_{298}^{0})$ of the reactions (R1) and (R3) was reported to be 16.4 kcal/mol.¹ However, PBE0 also underestimates the activation barrier of (R3) by ~5 kcal/mol (Table 1) and (R1) most likely remains a dominating primary reaction. The correct conclusion is therefore a result of partial error compensation between the computed activation energies of (R1) and (R3).

On the other hand, PBE0 reproduces well the activation barriers of nitro-nitrite rearrangements (Table

1). Thus, calculations of $\Delta^r H_0^0$ at PBE0 level along with incorrect estimations of $\Delta^{\neq} S_{298}^0$ render branching ratios intrinsically biased towards radical decomposition of nitroalkanes. E.g., the calculated in such manner $\Delta(\Delta^{\neq} G_{298}^0)$ of reactions NE1 and NE2 (Table 1) is less than 1 kcal/mol, while the isokinetic temperature for these channels was estimated to be higher than 600 K.⁵

Note that the M06-2X functional⁶ provides reasonable accuracy (Table 1, last column) and can be used for feasible calculations in the case of large nitroaliphatics and nitroaromatics.

Notes and references

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