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2	<b>Mechanisms of Enhanced Sulfur Tolerance on</b>			
3	Samarium (Sm)-doped Cerium Oxide (CeO <sub>2</sub> ) from			
4	First Principles			
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## ABSTRACT

The role of samarium (Sm) 4f states and Sm-perturbed O 2p states in determining the sulfur 3 tolerance of Sm-doped CeO<sub>2</sub> was elucidated by using the density functional theory (DFT)+U 4 calculation. We find that the sulfur tolerance of Sm-doped CeO<sub>2</sub> is closely related to the modification of 5 O 2p states by the strong interaction between Sm 4f and O 2p states. In particular, the availability of 6 unoccupied O 2p states near the Fermi level is responsible for enhancing the sulfur tolerance of Sm-7 8 doped CeO<sub>2</sub> compared to the pure CeO<sub>2</sub> by increasing the activity of surface lattice oxygen toward sulfur adsorbate, by weakening the interaction between Sm-O, and by increasing the migration 9 10 tendency of subsurface oxygen ion toward the surface.

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15 Keywords: Sulfur poisoning, solid oxide fuel cell, Samarium, CeO<sub>2</sub>, f electron

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# 2 Introduction

One of the main challenges in the development of anodes for solid oxide fuel cells (SOFCs) is the prevention of sulfur poisoning of Ni-based anodes because sulfur can block the active sites of the Ni anode and deteriorate the electrochemical performance of SOFCs.<sup>1, 2</sup> Recently, modifications of Nibased anodes with ceria (CeO<sub>2</sub>)-based materials has shown promise in reducing sulfur poisoning owing to the affinity of CeO<sub>2</sub> for sulfide at high temperature.<sup>3-6</sup> By coating ceria onto conventional nickel/yttria-stabilized zirconia (Ni/YSZ), improved SOFC cell performance was obtained<sup>7, 8</sup> because of the sulfur adsorption by ceria.

Furthermore, samarium (Sm) doping of CeO<sub>2</sub> may lead to enhanced sulfur tolerance and cell 10 11 performance. Sm-doped ceria (SDC)-coated Ni/YSZ may serve as a more effective sulfur adsorbent than ceria, and has resulted in stable cell performance in the presence of 500 ppm of H<sub>2</sub>S over 800 h.<sup>9,10</sup> 12 Approximately 25 to 50% improvement of the cell performance has been achieved with SDC-coated 13 14 Ni/YSZ than with ceria-coated Ni/YSZ and conventional Ni/YSZ owing to the enhanced sulfur adsorbing capability of Sm and enlarged triple-phase boundary (electrode-electrolyte-gas phase) areas 15 that reduce electrode polarization in the former.<sup>9, 11, 12</sup> The underlying mechanism governing the sulfur 16 17 tolerance of Sm-doped CeO<sub>2</sub>-based materials still remains unclear.

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Recently, the origin of sulfur poisoning on Ni-based anode materials has been investigated using 18 density functional theory (DFT). Zeng et al.<sup>13</sup> found that the smaller ionic radius of cation dopant  $X^{3+}$ 19 results in the higher sulfur tolerance on Ni/XSZ (X<sub>2</sub>O<sub>3</sub> stabilized zirconia), which is attributed to the fact 20 that the dopant size directly influences the Ni-XSZ interaction and in turn affects the Ni-S interaction. 21 Malyi et al.<sup>14</sup> revealed that the hybridization of Ni 3d and S 3p states serves an important role in the 22 Ni-S interaction on Ni surfaces. Lu et al.<sup>15</sup> investigated the interaction of atomic sulfur with the ceria 23 surfaces and found that sulfur tends to be in the  $CeO_2(111)$  after replacing oxygen, but on the  $CeO_2(110)$ 24 surface as  $SO_2^{2-}$ . Mayernick et al.<sup>16</sup> examined H<sub>2</sub>S adsorption and dissociation on Lanthanum (La)- and 25

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Terbium (Tb)-doped CeO<sub>2</sub>(111) by using DFT+U and ab initio thermodynamics, and indicated that substoichiometric oxygen vacancies play an important role in sulfur adsorption process occurring via H<sub>2</sub>S adsorption and dissociation where a strong endergonic molecular adsorption of H<sub>2</sub>S is the rate limiting step.

In this study, we investigated the role of Sm in determining the sulfur tolerance of Sm-doped 5  $CeO_2$  (Sm<sub>x</sub>Ce<sub>1-x</sub>O<sub>2</sub>) using spin-polarized DFT+U calculations. We examined the oxidation of H<sub>2</sub>S on 6  $Sm_xCe_{1-x}O_2$  and stoichiometric CeO<sub>2</sub> by calculating the reaction energetics and energy barriers of 7 various reaction pathways. H<sub>2</sub>S oxidation proceeds via H<sub>2</sub>S decomposition (H<sub>2</sub>S<sup>\*</sup>  $\rightarrow$  HS<sup>\*</sup> + H<sup>\*</sup>  $\rightarrow$  S<sup>\*</sup> + 8  $H^* + H^*$ ) and SO<sub>2</sub> formation (S<sup>\*</sup> + O<sup>\*</sup> + O<sup>\*</sup>  $\rightarrow$  SO<sup>\*</sup> + O<sup>\*</sup>  $\rightarrow$  SO<sub>2</sub><sup>\*</sup>) (asterisk (\*) denotes adsorbed species 9 10 hereafter) as shown in Supporting Information, Figure S1. We also present the electronic properties of 11 Sm-doped  $CeO_2$  and stoichiometric  $CeO_2$  to elucidate the origin of sulfur tolerance in  $CeO_2$ -based materials. 12

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# 14 Computational Methodology

Spin-polarized DFT+U calculations were performed using the Vienna ab-initio Simulation 15 Package (VASP)<sup>17</sup> within with the generalized gradient approximation (GGA-PBE<sup>18</sup>). The projector 16 17 augmented wave (PAW) method with a planewave basis set was employed to describe the interaction between core and valence electrons<sup>19, 20</sup>. A plane-wave expansion with a kinetic energy cutoff of 400 eV 18 was employed with a Gaussian smearing with a width of 0.2 eV. The integration of the Brillouin zone 19 was conducted using a  $4 \times 4 \times 1$  Monkhorst-Pack grid<sup>21</sup> for the  $2 \times 2$  CeO<sub>2</sub>(111) slab model, where three 20 21 Ce atom layers and six O atom layers are present as shown in Figure 1A (this nine layer slab model is be close to converged with respect to slab thickness and has been used in previous studies<sup>22-24</sup>). 22

The Sm-doped CeO<sub>2</sub> was modeled by introducing a Sm-doping atom to substitute Ce atoms in the CeO<sub>2</sub> (111) slab. Here a Sm dopant was located on the top surface, rather than on the subsurface since Sm dopant tends to segregate on the CeO<sub>2</sub> surface, rather than migrate into the bulk<sup>25, 26</sup>.

The geometries were optimized until the forces were reduced below 0.05 eV/Å. A Hubbard-U
term, U = 5 eV, was added to properly describe the localized nature of the 4*f* electrons of Ce.<sup>27, 28</sup>
Furthermore, the climbing image-nudged elastic band (CI-NEB)<sup>29</sup> method was employed to calculate
the activation energy for H<sub>2</sub>S decomposition and SO<sub>2</sub> formation along the examined pathway.

The charge difference density (Δn(r)) plots were obtained by subtracting the charge densities of
the separated adsorbate and the CeO<sub>2</sub> (or Sm-doped CeO<sub>2</sub>) surface from the charge density of the
adsorbate-surface system as follows:

$$\Delta n(\mathbf{r}) = n_{surf+adsorbate}(\mathbf{r}) - [n_{surf}(\mathbf{r}) + n_{adsorbate}(\mathbf{r})]$$
(1)

8

9 where  $n_{surf+adsorbate}(\mathbf{r})$ ,  $n_{surf}(\mathbf{r})$ , and  $n_{adsorbate}(\mathbf{r})$  are the electron charge distributions of the surface with 10 adsorbed species, isolated surface, and isolated adsorbate, respectively. The source of the isolated 11 adsorbate and surface structures used for the charge difference density calculations have been directly 12 obtained from the optimized structures of the bound systems, rather than the optimized free adsorbate 13 and surface systems.

# 14 The oxygen vacancy formation energy per oxygen atom $(E_{form})$ is calculated as:

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 $E_{form} = \left(E_{vac} - E_{perfect} + n(\frac{1}{2}E_{O_2})\right) / n$ 

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where  $E_{vac}$  and  $E_{perfect}$  are the total energies of oxygen vacant and perfect surface, respectively. *n* is the number of oxygen vacancy and  $E_{O2}$  is the total energy of oxygen molecule.

19 The adsorption energy  $(E_{ads})$  of an adsorbate is defined as  $E_{ads} = E_{substrate+adsorbate} - E_{substrate} - 20$ 20  $E_{adsorbate}$ , where  $E_{substrate+adsorbate}$ ,  $E_{substrate}$ , and  $E_{adsorbate}$  are the total energies of a substrate and adsorbate 21 (e.g., CeO<sub>2</sub> with adsorbed sulfur), a substrate (e.g., CeO<sub>2</sub>), and a gas phase adsorbate (e.g., sulfur). A

(2)

negative adsorption energy indicates that adsorption is exothermic (stable) with respect to the free gas phase adsorbate.

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#### 4 **Results and Discussion**

The modified electronic properties of Sm-doped CeO<sub>2</sub> were examined by analyzing the charge 5 6 density difference caused by the replacement of Ce by Sm (See Figure 1). Because of the higher electron density of Sm ([Xe]4f<sup>6</sup>6s<sup>2</sup>) than that of Ce ([Xe]4f<sup>1</sup>5d<sup>1</sup>6s<sup>2</sup>), strong charge accumulation is 7 observed at the location of Sm (up to 3.60  $e/Å^3$ ) in the charge density difference plot. Importantly, the 8 electron charge of the neighboring lattice oxygen atoms of Sm becomes depleted because of the 9 replacement of Ce<sup>4+</sup> by Sm<sup>3+</sup>, resulting in a relatively weak interaction of lattice oxygen with the Sm-10 11 doped CeO<sub>2</sub> surface. The weaker Ce–O interaction on Sm-doped CeO<sub>2</sub> is supported by the oxygen vacancy formation energy as shown in Table 1. Less energy is required to create one or two oxygen 12 vacancies in Sm-doped CeO<sub>2</sub> compared to that required for stoichiometric CeO<sub>2</sub>. This observation may 13 14 imply that the lattice oxygen with depleted electron charge helps promote sulfur oxidation by lowering the activation energy for SO<sub>2</sub> formation ( $S^* \rightarrow SO_2^*$ ). In the ensuing discussion, we demonstrate the 15 enhanced SO<sub>2</sub> formation by analyzing the reaction pathways and estimating the activation energies 16 using the CI-NEB<sup>29</sup> method. 17

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Figure 1. (A) Side view of the 2×2 CeO<sub>2</sub> structure consisting of three Ce and four O layers. (B) Charge density difference ( $\Delta n(\mathbf{r})$  in units of  $e/Å^3$ ) of Sm-doped CeO<sub>2</sub> upon Sm replacement defined by  $[n(\mathbf{r})_{\text{Sm-doped CeO_2}} - n(\mathbf{r})_{\text{CeO_2}}]$ , where  $n(\mathbf{r})_{\text{Sm-doped CeO_2}}$  and  $n(\mathbf{r})_{\text{CeO_2}}$  are the charge densities of Sm-doped CeO<sub>2</sub> and stoichiometric CeO<sub>2</sub> surfaces. Ce<sub>sm</sub> indicates the Ce atom replaced by Sm in (B).

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**Table 1**. Oxygen vacancy formation energies depending on the number of oxygen vacancies.
Superscripts *a*, *b*, *c*, *d*, and *e* indicate oxygen vacancy locations of (O1), (O2), (O1, O2), (O1, O4), and

9	(01, 03)	, respectively,	as depicted	in Figure 1B.
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	One vacancy	Two vacancies
	(eV/O atom)	(eV/O atom)
CeO <sub>2</sub>	3.33 <sup><i>a</i></sup>	3.01 <sup><i>c</i>,<i>d</i></sup>
Sm-doped CeO <sub>2</sub>	$1.79^a, 1.97^b$	$2.65^c, 2.32^e$

<sup>10</sup> 

Prior to detailed examination of the reaction pathways and corresponding activation energy, an
 overall reaction pathway from H<sub>2</sub>S adsorption to SO<sub>2</sub> desorption was examined as shown in the

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Supporting Information, Figure S2, indicating that the  $H_2S$  decomposition reaction is exothermic whereas the SO<sub>2</sub> desorption reaction is endothermic. The reaction limiting step of the sulfur removal mechanism may lie in the last step of SO<sub>2</sub> desorption. In-depth understanding of the effect of Sm on the reaction pathways can be achieved by comparing the activation energies required for each reaction step on both surfaces.

Figure 2 shows the lowest-energy reaction pathways of H<sub>2</sub>S decomposition on stoichiometric 6 CeO<sub>2</sub> and Sm-doped CeO<sub>2</sub>, demonstrating that Sm doping promotes H<sub>2</sub>S decomposition by lowering the 7 8 activation energy required for each step. Sm doping significantly lowers the activation energy barrier of the  $HS^* + H^* \rightarrow S^* + H^* + H^*$  step from 0.64 to 0.12 eV, and furthermore, turns an endothermic reaction 9 into an exothermic one at the last step of H<sub>2</sub>S decomposition. The H<sub>2</sub>S molecule binds weakly on the top 10 11 site of Ce or Sm where the sulfur is bound on Ce or Sm with respective adsorption energies of -0.21and -0.22 eV for stoichiometric CeO<sub>2</sub> and Sm-doped CeO<sub>2</sub>, which agrees with previous DFT<sup>30, 31</sup> and 12 experimental<sup>32</sup> studies. Notably, the adsorption configuration of  $S^* + H^* + H^*$  on CeO<sub>2</sub>, where sulfur is 13 located at the top site of Ce, is ~0.15–0.19 eV more stable than those in which sulfur is located on the 14 bridge or hollow sites of Ce (Supporting Information, Figure S3); the most stable configuration of  $S^*$  + 15  $H^* + H^*$  is where sulfur interacts directly with lattice oxygen (Figure 2A). Marrocchelli and Yildiz<sup>30</sup> also 16 showed that the  $HS^* + H^* \rightarrow S^* + H^* + H^*$  reaction on stoichiometric CeO<sub>2</sub>(111) is exothermic when 17 atomic sulfur is bound to the lattice oxygen, and an energy barrier of 0.69 eV is present between the two 18 steps, which may be attributed to an intermediate step where sulfur is bound to the top site of Ce as 19 shown here. 20

The CeO<sub>2</sub> and Sm-doped CeO<sub>2</sub> surfaces differ notably in terms of the stabilization of atomic sulfur and hydrogen on each surface; Sm doping considerably enhances the stability of the S<sup>\*</sup> + H<sup>\*</sup> + H<sup>\*</sup> step. This result agrees well with the empirical observation that Sm doping enhances the sulfuradsorbent properties of CeO<sub>2</sub> in recent SOFC studies employing a Ce<sub>0.8</sub>Sm<sub>0.2</sub>O<sub>2</sub>-coated Ni/YSZ anode.<sup>9</sup>, <sup>10</sup> In the presence of 100 ppm of H<sub>2</sub>S, the polarization resistance of the Ce<sub>0.8</sub>Sm<sub>0.2</sub>O<sub>2</sub>-coated Ni/YSZ

anode (40%) increased to a lesser extent than that of the CeO<sub>2</sub>-coated Ni/YSZ anode (90%);<sup>9</sup> this
demonstrates that Sm enhances the capture of sulfur by CeO<sub>2</sub>, where it is stored as Ce<sub>2</sub>O<sub>2</sub>S.<sup>9</sup> The current
DFT results also demonstrate that Sm-doped CeO<sub>2</sub> can better decompose the H<sub>2</sub>S molecule into atomic
sulfur compared to stoichiometric CeO<sub>2</sub>, which may lead to a higher storage of Ce<sub>2</sub>O<sub>2</sub>S on Sm-doped
CeO<sub>2</sub> than on stoichiometric CeO<sub>2</sub>.

6



8 Figure 2. The lowest-energy reaction pathways (A) and relative energy diagrams (B) of  $H_2S$ 9 decomposition on stoichiometric CeO<sub>2</sub> (dashed lines) and Sm-doped CeO<sub>2</sub> (solid lines). *a* and *b* indicate 10 the location of Sm depending on Sm-doped CeO<sub>2</sub> (a) and (b), respectively. Empty rectangles represent 11 transition states (TS) in eV. The absence of TS between two steps indicates the absence of an energy

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barrier. Case(I) and Case(II) represent the S\* + H\* + H\* configurations where sulfur is located on the top
of Ce and lattice oxygen, respectively.

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The enhanced decomposition of H<sub>2</sub>S on Sm-doped CeO<sub>2</sub> may be attributed to the modified 4 electronic structure of Sm-doped CeO<sub>2</sub> due to Sm doping. Figure 3A shows a comparison of the total 5 density of states (DOS) between stoichiometric CeO<sub>2</sub> and Sm-doped CeO<sub>2</sub>. For the stoichiometric CeO<sub>2</sub> 6 case, we find that a gap is present between the top of the O 2p valence band and the bottom of the 7 unoccupied Ce 4f conduction band, which agrees well with previous DFT studies.<sup>27, 33, 34</sup> On the other 8 hand, for the Sm-doped CeO<sub>2</sub> case, our calculation predicts that a new gap state between the O 2*p*-Ce 4*f* 9 gap is formed near the Fermi level. This new gap state is related to Sm 4f and Sm-perturbed O 2p states, 10 which may be responsible for the significantly increased stability of sulfur radicals and in turn enhanced 11 H<sub>2</sub>S decomposition on Sm-doped CeO<sub>2</sub> compared to that on stoichiometric CeO<sub>2</sub>. 12

To unravel the relative role of Sm 4f and Sm-perturbed O 2p states in determining the surface 13 reactivity toward the sulfur species, we calculated the sulfur adsorption energy at various surface sites 14 and projected density of state (PDOS) onto Sm 4f and O 2p states. First, our calculation shows the slight 15 enhancement of the sulfur adsorption energy  $(E_{ad})$  at the Sm top site  $[E_{ad} = -0.88 \text{eV}]$  on Sm-doped 16 CeO<sub>2</sub>, compared to the Ce top site on stoichiometric CeO<sub>2</sub> [ $E_{ad} = -0.70$ eV]. On the other hand, we find a 17 significant improvement of the sulfur adsorption energy at the O top site on Sm-doped CeO<sub>2</sub> by 0.82eV 18 in comparison with stoichiometric  $CeO_2$  case, indicating that the surface O site is the critical factor to 19 boost the sulfur stability in Sm-doped CeO<sub>2</sub> (Supporting Information, Table S1). Figure 4 shows the 20 charge difference density upon sulfur adsorption on lattice oxygen in Sm-doped CeO<sub>2</sub> and 21 stoichiometric CeO<sub>2</sub> cases. We find the larger electronic charge accumulation in the bonding region 22 between S 2p an O 2p orbital in Sm-doped CeO<sub>2</sub> than that in stoichiometric CeO<sub>2</sub> cases, suggesting the 23 formation of much stronger covalent bond in Sm-doped CeO<sub>2</sub>. 24



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Figure 3. (A) Total density of states (TDOS) of stoichiometric  $CeO_2$  (dashed lines) and Sm-doped  $CeO_2$ (solid lines). (B) Projected density of states (PDOS) of surface oxygen 2p state (four O atoms) of stoichiometric  $CeO_2$  and Sm-doped  $CeO_2$ . (C) PDOS of Sm 4*d* and 4*f* states. Spin-up and -down are indicated by positive and negative values, respectively. The Fermi energy is referenced at 0 eV.



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Figure 4. Charge difference density (Δn(r) in units of e/Å<sup>3</sup>) of stoichiometric CeO<sub>2</sub> and Sm-doped CeO<sub>2</sub>
 upon sulfur adsorption on lattice oxygen. Two arrows indicate strong charge depletion upon sulfur
 adsorption.

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The dramatic enhancement of the surface activity at the lattice O site is related to the formation 5 6 of unoccupied O 2p states caused by the strong mixing between Sm 4f and O 2p states. According to our DFT+U calculation on the PDOS of Sm 4f and O 2p states (Figure 3B and 3C), a significant overlap 7 between Sm 4f and O 2p states are observed in the  $-1.7 \text{ eV} \le E - E_f \le 1.0 \text{ eV}$ , suggesting a strong 8 hybridization of Sm 4f-O 2p states in Sm-doped CeO<sub>2</sub>. Notice no overlap between Ce 4f-O 2p states in 9 stoichiometric CeO<sub>2</sub>. Correspondingly, the PDOS of occupied O 2p states in Sm-doped CeO<sub>2</sub> noticeably 10 shifts toward the Fermi level compared to the stoichiometric CeO<sub>2</sub> case. In particular, the onset of low 11 binding energy tail shifts above the Fermi level, which results in the formation of unoccupied O 2p 12 states in the 0 eV  $\leq E - E_f \leq 1.0$  eV. The different oxidation state of dopants and Bader charge analysis 13 also support the creation of empty unoccupied O 2p states. Our DFT+U calculation shows that the 14 oxidation state of Sm dopant is about +3 in Sm-doped CeO<sub>2</sub>, which is lower than that of Ce atom (Ce<sup>4+</sup>) 15 (Note that the magnetization of Sm 4f states in Sm-doped CeO<sub>2</sub> is ~ 4.9e. Together with the fact that 16 there is no electron in Sm 6s state, this indicates that the oxidation state of Sm is close to +3 in Sm-17 doped CeO<sub>2</sub>). As a result, the electronic charge of  $O^{2-}$  perturbed by Sm 4f states is substantially 18 decreased by 0.060*e* (from Bader charge analysis) compared to the  $O^{2-}$  case in stoichiometric CeO<sub>2</sub>, 19 leading to the formation of empty O 2*p* states. 20

Moreover, the unoccupied O 2p states above the Fermi level may weaken the binding energy between Sm and the lattice oxygen due to reduced electron density of oxygen, promoting the formation of oxygen vacancies (see comparison of oxygen vacancy formation energies in Table 1). A similar phenomenon has been reported for copper adsorbed-CeO<sub>2</sub>(111) where the formation of oxygen

- 1 vacancies was promoted by formation of unoccupied states near the Fermi level upon copper
- 2 adsorption.<sup>27</sup>
- 3



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**Figure 5.** (A) PDOS of surface oxygen 2p (four O atoms) and dopant (Ce, Pr, Pm, Sm, and Eu) 4f (one dopant atom) states where spin-up and -down states are summed. The Fermi energy is referenced at 0 eV. (B) Sulfur adsorption energies on top sites of dopants (rhombus), top sites of surface lattice oxygen (square), and bridge sites of dopant-oxygen (circles). *x*-axis represents the portion of unoccupied lattice oxygen 2p states below 1 eV in PDOS. Sulfur on the top of dopants is restrained in *x* and *y* directions.

To better understand the importance of unoccupied O 2*p* states in determining the activity of surface lattice oxygen, the Sm atom in Sm-doped CeO<sub>2</sub> was replaced with three 4*f* shell dopants (i.e., praseodymium (Pr) ([Xe]4f<sup>3</sup>6s<sup>2</sup>), promethium (Pm) ([Xe]4f<sup>5</sup>6s<sup>2</sup>), and europium (Eu) ([Xe]4f<sup>7</sup>6s<sup>2</sup>)), and the sulfur adsorption and the electronic properties of Pr-, Pm-, and Eu-doped CeO<sub>2</sub> were examined.

From the PDOS of O 2*p* states and dopant 4*f* states (See Figure 5A), we find that the number of unoccupied O 2*p* states strongly depends on the degree of overlap between O 2*p* states and dopant 4*f* states. That is, as the number of 4*f* electrons of dopants (Ce = 2, Pr = 3, Pm = 5, Sm = 6, and Eu = 7) increase, the hybridization of Sm 4*f*-O 2*p* states becomes stronger, leading to the increase of the number of unoccupied O 2*p* states of 4*f* shell dopant-perturbed oxygen ion (Ce = 0, Pr = 1.47, Pm = 3.04, Sm = 3.85, and Eu = 5.34 % in terms of the percentage of unoccupied O 2*p* state). TDOS and PDOS of Pr-, Pm-, Sm-, and Eu-doped CeO<sub>2</sub> are compared in Supporting Information, Figure S4.

In Figure 5B, the sulfur adsorption energy is displayed as a function of the number of 12 unoccupied O 2p states. We find the improvement of sulfur adsorption energy at the O-associated 13 14 surface sites (dopant-oxygen bridge sites) as the number of unoccupied O 2p states increases [Ce (-2.58) < Pr(-2.81) < Pm(-3.08) < Sm(-3.26) < Eu(-3.67 eV)]. This trend is inversely proportional 15 to the oxygen vacancy formation energy depending on the 4*f* shell dopants (Ce (3.33) > Pr (2.96) > Pm16 (2.11) > Sm(1.79) > Eu(1.17 eV), which is attributed to the increased depletion of electron density 17 from oxygen as the shift of O 2p states above the Fermi level. This demonstrates that the availability of 18 unoccupied O 2p states plays a key role in enhancing the stability of sulfur and in turn H<sub>2</sub>S 19 decomposition on the surface. 20

Figure 6 represents the lowest-energy reaction pathways of gas phase  $SO_2$  production from adsorbed sulfur on stoichiometric CeO<sub>2</sub> and Sm-doped CeO<sub>2</sub>. In this reaction scheme, adsorbed sulfur uses surface lattice oxygen to form  $SO^*$  and  $SO_2^*$  according to the following hypothesis, where adsorbed sulfur (S<sup>\*</sup>) is electrochemically oxidized by lattice oxygen that is replenished from oxygen ion flux from a cathode.<sup>35-38</sup>

$$S^* + 2O_{\text{SDC}} \rightarrow * + SO_{2(\text{gas})} + 2V\ddot{o} + 4e^-$$
(1)

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Here,  $O_{SDC}$ , \*, and  $V\ddot{o}$  represent a lattice oxygen, the SDC surface, and an oxygen vacancy in SDC, respectively. Although the H<sub>2</sub>S reactions on ceria involve the formation of H<sub>2</sub>O<sub>(gas)</sub> and H<sub>2(gas)</sub>, a DFT calculation model based on Eq. (1) may provide enough insight into the role of Sm in determining the sulfur tolerance of SDC given that the resistance to sulfur poisoning is mainly attributed to the SO<sub>2</sub> desorption reaction.<sup>35-37</sup>

9 Our calculation predicts that Sm doping enhances the adsorbed SO<sub>2</sub> formation by lowering the 10 activation energy for the SO<sup>\*</sup>  $\rightarrow$  SO<sub>2</sub><sup>\*</sup> + *Vö* step from 0.90 to 0.62 or 0.50, depending on the location of 11 Sm (Figure 6B). The lowered energy barriers are related to the significant reduction of oxygen vacancy 12 formation energy due to the weakened interaction between lattice oxygen and Sm dopant by the 13 substantial existence of unoccupied O 2*p* states. Note that the oxygen vacancy formation energy is 14 reduced by 1.54 (or 1.36) eV on Sm-doped CeO<sub>2</sub> compared to pure ceria as shown in Table 1.

15 The energy for the desorption of SO<sub>2</sub> on stoichiometric  $CeO_2(111)$  is 2.05 eV, which agrees with the results of a previous DFT study in which the desorption energy was calculated at 1.96 eV. However, 16 Sm doping lowers the desorption energy of SO<sub>2</sub> to 2.00 and 1.69 eV. The relatively more stabilized 17 18 oxygen vacant Sm-doped CeO<sub>2</sub> surface compared to the oxygen vacant CeO<sub>2</sub> surface also contributes to enhancing sulfur removal by promoting SO<sub>2</sub> desorption. Notably, the significantly lowered SO<sub>2</sub> 19 desorption energy of 1.69 eV on Sm-doped CeO<sub>2</sub> is due to the fact that the subsurface oxygen ion near 20 21 Sm tends to migrate onto the surface during the SO<sub>2</sub> desorption. The calculated distances of migrated subsurface oxygen ions toward the surface in stoichiometric CeO<sub>2</sub> and Sm-doped CeO<sub>2</sub> are 0.35 and 22 1.07 Å, respectively (See Supporting Information, Figure S5), which implies that Sm doping may 23 promote the migration process of oxygen ion formed at the subsurface toward the surface during the 24

SO<sub>2</sub> desorption, resulting in the further stabilization of oxygen vacant Sm-doped CeO<sub>2</sub> surface compared to stoichiometric CeO<sub>2</sub> and in turn the improvement of SO<sub>2</sub><sup>\*</sup> desorption energetics. This is also confirmed by higher exothermicity by 0.18 eV on Sm-doped CeO<sub>2</sub> than on stoichiometric CeO<sub>2</sub> for the complete migration process of oxygen ion from subsurface to surface.

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Figure 6. The lowest-energy reaction pathways (A) and energy diagrams (B) for SO<sub>2</sub> formation on
stoichiometric CeO<sub>2</sub> (dashed lines) and Sm-doped CeO<sub>2</sub> (solid lines). *a* and *b* indicate the location of
Sm depending on Sm-doped CeO<sub>2</sub> (A) and (B), respectively. Empty rectangles represent transition
states (TS) in eV. The absence of TS between two steps indicates the absence of an energy barrier.

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#### 12 Conclusions

13 In summary, we find that the strong interaction between Sm 4f and O 2p states on Sm-doped 14 CeO<sub>2</sub> increases the availability of unoccupied O 2p states near the Fermi level. The increased

unoccupied O 2p states in turn enhance the sulfur tolerance of Sm-doped CeO<sub>2</sub> compared to the pure ceria by boosting the stability of sulfur species in H<sub>2</sub>S decomposition owing to the increased activity of surface lattice oxygen toward sulfur adsorption, by reducing the oxygen vacancy formation energy in SO<sub>2</sub> production due to the weakened interaction between Sm and O, and by further stabilizing the oxygen vacant surface in SO<sub>2</sub> desorption process attributed to the increased migration tendency of subsurface oxygen ion toward the surface. This work also hints on the importance of properly engineering the activity of oxygen ion in CeO<sub>2</sub>-based materials by adding dopant.

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# 9 Supporting Information

Schematic diagrams of H<sub>2</sub>S decomposition and SO<sub>2</sub> formation pathways, adsorption energies of H<sub>2</sub>S,
HS, and S species, additional adsorption configurations of S<sup>\*</sup> + H<sup>\*</sup> + H<sup>\*</sup> on CeO<sub>2</sub>(111), TDOS and
PDOS of Pr-, Pm-, Sm-, and Eu-doped CeO<sub>2</sub>, and configurations of diffused subsurface oxygen atoms.
This materials is available free of charge via the Internet at http://pubs.acs.org.

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