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DPT tautomerisation of the G·A_{syn} and A*·G*_{syn} DNA mismatches: A QM/QTAIM combined atomistic investigation

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Abstract. By applying a combined QM and QTAIM atomistic computational approach we have established for the first time that the G·A_{syn} and A*·G*_{syn} DNA mismatches (rare tautomers are marked with an asterisk), causing spontaneous transversions with substantially various probability, radically differ from each other in their ability to tautomerise through the double proton transfer (DPT). The A*·G*_{syn} mismatch tautomerises quite easily ($\Delta\Delta G_{TS}\approx 4\cdot kT$ at room temperature) into the A·G*_{syn} mismatch through the asynchronous concerted mechanism, whereas the G·A_{syn} base mispair does not tautomerise *via* the DPT at all, since there is no local minimum corresponding to the tautomerised G*·A*_{syn} mismatch on the potential energy surface. It was established that the A·G*_{syn} base mispair is dynamically unstable H-bonded complex with an extremely low lifetime of 2.17·10⁻¹³ s. Consequently, the obtained results allow us to believe that spontaneous or forced dissociation of both the G·A_{syn} and A*·G*_{syn} DNA mismatches by the DNA-polymerase occurs with the preservation of the tautomeric status of the bases.

Keywords. Rare tautomers of the adenine and guanine \cdot *Syn-* and *anti*-conformations \cdot Sweeps of the energetic, electron-topological, geometric, polar and NBO parameters along the IRC \cdot The double proton transfer \cdot Cooperativity of the H-bonds \cdot B3LYP and MP2 levels of QM theory \cdot QTAIM analysis

Introduction

The penetration into the underlying physico-chemical mechanisms that provide the fantastically high precision of the enzymatic DNA biosynthesis in the living cell is an important fundamental problem in structural biology and related disciplines, which is now far from its final solution [1,2,3,4,5,6,7,8]. It is now found for certain that the root cause of spontaneous mutations, in particular point mutations, that inevitably accompany this very important biological process, responsible for the reproduction of the genetic information from generation to generation, is the formation of the mispairs of nucleotide bases adopting the geometry that closely mimics the structure of the Watson-Crick base pairs in the base-pair recognition pocket of the high-fidelity replicational DNA polymerase [9,10,11,12]. Meanwhile, the asymmetry of the spontaneous point mutations in DNA, i.e., the dependence of their frequency of occurrence on the appurtenance of the bases to the template DNA strand or to the incoming deoxynucleotide triphosphate, is one of their least understood features [13,14]. It was experimentally registered (see ref. [15] and bibliography presented there) that the adenine guanine $(A \cdot G)$ and $G \cdot A$ transversions demonstrate noticeable asymmetry: mistakes happen more often when the A is the base of the incoming deoxynucleotide triphosphate, while the G – of the template strand (here and below in the $X \cdot Y$ mispair the X base belonging to the template strand is in left and the Y base of the incoming deoxynucleotide triphosphate - in the right).

In our recent works [16,17] we have postulated for the first time that the incorrect purine-purine $G \cdot A_{syn}$ DNA base pair [18,19,20] is responsible for the highly probable $G \cdot A$ transversion, while the $A^* \cdot G^*_{syn}$ base mispair [16,17] – for the low probable $A \cdot G$ transversion (see Scheme 1). The assumption has also been made that both these mispairs are formed in the base-pair recognition pocket of the replicational DNA polymerase through a common intermediate - the so-called long $A \cdot G/G \cdot A$ Watson-Crick base pair, which is stabilized by the N6H···O6 and N1H···N1 hydrogen bonds (H-bonds) and the N2H···HC2 dihydrogen bond [21].

Thus, the $A^* \cdot G^* \rightarrow A^* \cdot G^*_{syn}$ conformational transition is preceded by the double proton transfer (DPT) tautomerisation of the A·G base pair [21], that takes off the steric constraints for its implementation; the G·A_{syn} base mispair is formed through the conformational transformation of the G·A base mispair without changing of its tautomeric status [22]. Within this model, the asymmetry of the transversions is explained by the significantly lower Gibbs free energy of the G·A_{syn} base pair than of the A*·G*_{syn} base pair [16,17]. It should be remarked that this simplest physico-chemical model is based on the experimentally observed regularity - in the recognition pocket of the DNA polymerase only the base in the incoming deoxynucleotide triphosphate can perform *anti*→*syn* conformational switch; while for the template base this switch is prohibited [10,13,14].

At the same time the biologically important question remains unanswered - what is the atomistic mechanism of the DPT tautomerisation of the $G \cdot A_{syn}$ and $A^* \cdot G^*_{syn}$ base mispairs causing spontaneous transversions in DNA? An answer to this question provides information about the cooperativity of the intermolecular H-bonds that stabilize these base pairs and also allows us to understand the way in which the dissociation of the $G \cdot A_{syn}$ and $A^* \cdot G^*_{syn}$ DNA base mispairs by the polymerase machinery occurs – with or without the alteration of the tautomeric status [23,24,25] of the bases that are part of these pairs.

This work is devoted to the searching for the answers to the just formulated questions and discussion of the possible biological consequences that can be derived from them. We have managed to demonstrate that the studied $G \cdot A_{syn}$ and $A^* \cdot G^*_{syn}$ DNA base mispairs exhibit unique physico-chemical properties: while the first base pair does not tautomerise due to the absence of the local minimum corresponding to the $G^* \cdot A^*_{syn}$ base mispair on its potential energy surface, the $A^* \cdot G^*_{syn}$ base mispair tautomerises with the constant of tautomerisation, greatest among all investigated earlier incorrect base pairs [16,17].

It was found for the first time that whereas all three investigated $G \cdot A_{syn}$, $A^* \cdot G^*_{syn}$ and $A \cdot G^*_{syn}$ base mispairs are thermodynamically stable structures, since the values of their Gibbs free energies of interaction are less than zero, the $A^* \cdot G^*_{syn}$ base mispair is dynamically unstable complex with an extremely low lifetime of $2.17 \cdot 10^{-13}$ s, that is by several orders less than the time spent by the DNA-polymerase machinery to forcibly dissociate a base pair into the monomers (several ns [8]).

In the result of the investigation it was found that the $A^* \cdot G^*_{syn} \leftrightarrow A \cdot G^*_{syn}$ Löwdin's [23,27,28] tautomerisation *via* the DPT is a concerted (i.e., this reaction involves no stable intermediates) and asynchronous (i.e., both protons involved in the H-bonds move with a time gap) process.

It was established based on the profiles of the H-bond energies that the $A^* \cdot G^*_{syn}$ base mispair is stabilized by the cooperative [29,30,31] O6H···N6 (9.71) and N1H···N7 (6.99) H-bonds, while the $A \cdot G^*_{syn}$ base mispair – by the cooperative N6H···O6 (7.23) and N7H···N1 (9.59 kcal·mol⁻¹) H-bonds.

At the same time, a new approach to the definition of the cooperativity of the intermolecular Hbonds in the H-bonded complexes that do not principally tautomerise through the DPT because of the lack of the local minimum corresponding to the tautomerised complex has been proposed for the first time. By using this approach it was found that the N6H····O6 (5.08) and N1H····N7 (6.12 kcal·mol⁻¹) Hbonds in the G·A_{syn} base mispair are cooperative.

Computational Details

Gaussian'09 suite of programs has been used to provide all quantum-chemical calculations [32]. Geometries and harmonic vibrational frequencies of the G·A_{svn}, A*·G*_{svn} and A·G*_{svn} DNA base mispairs and the transition state $(TS_{A^* \cdot G^* syn \leftrightarrow A \cdot G^* syn})$ of their mutagenic tautomerisation via the DPT were obtained employing Density Functional Theory (DFT) [33] with the B3LYP hybrid functional [34], which includes Becke's three-parameter exchange functional (B3) [35] combined with Lee, Yang and Parr's (LYP) correlation functional [36] in connection with Pople's 6-311++G(d,p) basis set in vacuum. The DFT method has been recommended in the literature for the description of the tautomerisation phenomena of the H-bonded nucleobase pairs in the ground electronic state [37,38,39,40] and for the study of the vibrations of the constituents of the nucleic acids [41,42,43,44,45,46,47], since it has shown a good balance between computational cost and accuracy and therefore can be considered as the shortest way to MP2 results [27,28,48,49,50]. A scaling factor of 0.9668 [51,52] has been used in the present work at the B3LYP quantum-mechanical (QM) level of theory to correct the harmonic frequencies of all the studied structures. We then performed single point energy calculations using correlated MP2 level of theory [53] with the 6-311++G(2df,pd) and 6-311++G(3df,2pd) Pople's [54,55,56] and cc-pVTZ and cc-pVQZ Dunning's cc-type [57,58] basis sets for the B3LYP/6-311++G(d,p) geometries to consider electronic correlation effects as accurately as possible.

The correspondence of the stationary points to the $G \cdot A_{syn}$, $A^* \cdot G^*_{syn}$ and $A \cdot G^*_{syn}$ base mispairs (see Scheme 1) or to the $TS_{A^* \cdot G^*syn \leftrightarrow A \cdot G^*syn}$, located by means of Synchronous Transit-guided Quasi-Newton (STQN) method [59,60], on the potential energy landscape has been checked by the absence or the presence, respectively, of one and only one imaginary frequency corresponding to the normal mode that identifies the reaction coordinate.

Since the stationary points $(A^* \cdot G^*_{syn}, A \cdot G^*_{syn})$ base mispairs and the $TS_{A^* \cdot G^* syn \leftrightarrow A \cdot G^* syn})$ were located, the IRC path was established by following the IRC in the forward and reverse directions from the TS using the Hessian-based predictor-corrector (HPC) integration algorithm [61,62,63] with tight convergence criteria. These calculations eventually ensure that the proper reaction pathway leads either to the expected reactants or products.

The electronic interaction energies E_{int} have been computed at the MP2/6-311++G(2df,pd)//B3LYP/6-311++G(d,p) level of theory as the difference between the total energy of the base pair and the energies of the isolated monomers, neglecting the deformational energy. In each case the interaction energy was corrected for the basis set superposition error (BSSE) [64,65] through the counterpoise procedure [66,67]. The Gibbs free energy G values for all structures were obtained at room temperature (T=298.15 K) in the following way:

$$G=E_{el}+E_{corr},$$
(1)

where E_{el} – the electronic energy, E_{corr} – the thermal correction.

The lifetime τ of the A·G*_{syn} base mispair was calculated using the formula $1/k_r$, where the value of the reverse k_r rate constant for the A*·G*_{syn} \leftrightarrow A·G*_{syn} tautomerisation reaction was obtained as [68,69,70,71]:

$$k_r = \Gamma \cdot \frac{k_B T}{h} e^{-\frac{\Delta \Delta G_r}{RT}},\tag{2}$$

where $\Gamma = 1 + \frac{1}{24} \left(\frac{hv_i}{k_B T} \right)^2$ – Wigner's tunneling correction [72], k_B – Boltzmann's constant, T = 298.15 K –

absolute temperature, h – Planck's constant, $\Delta\Delta G_r$ – the Gibbs free energy of activation for the DPT reaction in the reverse direction, R – universal gas constant, v_i – the magnitude of the imaginary frequency at the TS_{A*-G*syn} \leftrightarrow A·G*syn.

Bader's quantum theory "Atoms in molecules" (QTAIM) was applied to analyse the electron density distribution [73], using program package AIMAll [74] with all the default options. Wave functions were obtained at the level of theory used for geometry optimisation. The presence of a bond critical point (BCP) [73], namely the so-called (3,-1) BCP and a bond path between the H-bond donor and acceptor, as well as the positive value of the Laplacian at this BCP ($\Delta \rho \ge 0$), were considered as criteria for the H-bond formation [75,76,77].

The energies of the intermolecular H-bonds E_{HB} in the A*·G*_{syn} and A·G*_{syn} base mispairs were evaluated by the empirical logansen's formula [78]:

$$E_{HB} = 0.33 \cdot \sqrt{\Delta \nu} - 40 \tag{3}$$

where Δv - the magnitude of the redshift (relative to the free molecule) of the stretching mode of the Hbonded groups. The partial deuteration of the CH, NH, OH and NH₂ groups involved in the H-bonding was applied to eliminate the effect of vibrational resonances [77].

The energy of the N6H····O6 H-bond in the $TS_{A^*\cdot G^*syn\leftrightarrow A\cdot G^*syn}$ was estimated by the Nikolaienko-Bulavin-Hovorun formula [79]:

$$E_{N6H\cdots O6} = -2.03 + 225 \cdot \rho$$

where ρ is the electron density at the (3,-1) BCP of the N6H···O6 H-bond.

The energies of all intermolecular H-bonds E_{HB} under the investigation of the sweeps of their energies were evaluated by the empirical Espinosa-Molins-Lecomte (EML) formula [80,81] based on the electron density distribution at the (3,-1) BCPs of the H-bonds:

(4)

$E_{HB}=0.5 \cdot V(r),$

(5)

where V(r) - the value of a local potential energy density at the (3,-1) BCPs.

To determine the deformation energy necessary to apply for the investigated base mispairs to acquire the Watson-Crick sizes, we have adjusted the glycosidic angles at the G, A* and A bases in the $G \cdot A_{syn}$, $A^* \cdot G^*_{syn}$ and $A \cdot G^*_{syn}$ base pairs to those corresponding to the A and G bases in the A·T and G·C Watson-Crick base pairs, respectively, while at the A_{syn} , G^*_{syn} and G^*_{syn} bases - to those corresponding to the T and C bases and the distance between the appropriate glycosidic hydrogens and then frozen them using "opt=modredundant" key word.

To determine the cooperativity of the intermolecular H-bonds in the G·A_{syn} base mispair, that does not tautomerise *via* the DPT, we have realized a new approach. The lengths of the N6H and N1H atomic groups – the donors of the N6H····O6 and N1H····N7 H-bonds - has been forcedly stretched (each of them separately) with the step=0.01 Å and then the energy E_{HB} calculated by the EML formula [80,81] has been acquired at the each step at the (3,-1) BCPs of the H-bonds. As a result we obtained the dependency of the energy E_{HB} on the lengths N6H/N1H, that allows us to determine the cooperativity of the N6H····O6 and N1H····N7 H-bonds in the G·A_{syn} base mispair.

Results and Discussion

We have divided this section into three paragraphs aiming to simplify the perception of the material, concretely, the first one is dedicated to the thermodynamical/dynamical stability of the investigated $G \cdot A_{syn}$, $A^* \cdot G^*_{syn}$ and $A \cdot G^*_{syn}$ base mispairs and to their physico-chemical parameters; in the second part the physico-chemical mechanism of the $A^* \cdot G^*_{syn} \leftrightarrow A \cdot G^*_{syn}$ tautomerisation *via* the DPT would be discussed; and finally in the third part the evolution of the selected physico-chemical characteristics of the base mispairs and of the intermolecular specific contacts stabilizing them along the IRC of the $A^* \cdot G^*_{syn} \rightarrow A \cdot G^*_{syn}$ DPT tautomerisation would be covered. Obtained by us results are numerically represented in Tables 1-3 and graphically depicted in Scheme 1 and Figures 1-2, S1-S6.

1. Thermodynamical and dynamical stability of the $G \cdot A_{syn}$, $A^* \cdot G^*_{syn}$ and $A \cdot G^*_{syn}$ DNA base mispairs and their main physico-chemical characteristics

1.1. The A*·G*_{syn} and A·G*_{syn} base mispairs

Firstly, we have analysed the thermodynamical and dynamical (vibrational) stability of the investigated $A^* \cdot G^*_{syn}$ and $A \cdot G^*_{syn}$ base mispairs, containing A base in the *anti*- and G base in the *syn*-orientations, or structural and energetical characteristics of the interbase interactions stabilizing them.

It was found that the planar $A^* \cdot G^*_{syn}$ ($\Delta G = \Delta E = 0.00$), stabilized by the O6H…N6 (9.71) and N1H…N7 (6.99) H-bonds, and the planar $A \cdot G^*_{syn}$ ($\Delta G = 1.89$ and $\Delta E = 2.20$ kcal·mol⁻¹), stabilized by the

N6H···O6 (7.23) and N7H···N1 (9.59 kcal·mol⁻¹) H-bonds, structures with C_s symmetry and planar heterocycles of the flexible [41,83,84] A and G bases, are thermodynamically stable, since their commensurable Gibbs free energies of interaction are less than zero (ΔG_{int} =-11.47 (A*·G*_{syn}) and -12.96 (A·G*_{syn}) kcal·mol⁻¹) (Tables 1-3, Fig. 2). In contrast to the A·A* [21], A·A*_{syn} [31] and A·G [71] base mispairs, we have not established electron-topological and spectroscopic manifestations of the third specific C2H···HC8 contact as in the A*·G*_{syn}, A·G*_{syn} and TS_{A*·G*syn+A·G*syn} stationary structures, so along the IRC of the A*·G*_{syn} \leftrightarrow A·G*_{syn} tautomerisation *via* the DPT. Moreover, the electronic and Gibbs free interaction energies for the A*·G*_{syn} (ΔG_{int} =-11.47 and ΔE_{int} =-23.00) and A·G*_{syn} (ΔG_{int} =-12.96 and ΔE_{int} =-26.38) base mispairs are close, but less than those for the G·C Watson-Crick DNA base pair (ΔG_{int} =-15.97 and ΔE_{int} =-29.28 kcal·mol⁻¹ obtained at the MP2/6-311++G(2df,pd)//B3LYP/6-311++G(d,p) level of QM theory) [27]. Notably, the intermolecular Hbonds in the A*·G*_{syn} and A·G*_{syn} base mispairs consists more than half (72.6 and 63.8 %, respectively) of their electronic interaction energies, making a major contribution into their stabilization. The $\Sigma E_{HB}/l\Delta E_{int}$! electronic-energy relationship can be considered as an invariant characteristic of all H-bonded nucleobase pairs [48,27,28,77], but not exceptionally of these base pairs.

A computational survey of the potential energy surface of the planar $A^* \cdot G^*_{syn}$ base mispair reveals that it tautomerises into the planar $A \cdot G^*_{syn}$ base mispair through the slightly non-planar $TS_{A^*G^*syn\leftrightarrow A\cdot G^*syn}$ (C₁ symmetry; $\angle C6N6(A)C6C5(G_{syn})=11.3^\circ$), stabilized by the N1-H-N7 covalent bridge and the N6H····O6 H-bond (12.62 kcal·mol⁻¹), with imaginary frequency $v_i=940.7i$ cm⁻¹ (Tables 1 and 2; Fig. 2).

Transition state optimizations and intrinsic reaction coordinate (IRC) calculations for the $A^* \cdot G^*_{syn} \leftrightarrow A \cdot G^*_{syn}$ tautomerisation *via* the DPT showed that there is additional TS at the IRC=-5.23 Bohr connecting the $A^* \cdot G^*_{syn}$ base mispair (IRC=-8.75 Bohr) and the $A^+ \cdot G^-_{syn}$ zwitterionic intermediate (IRC=-3.76 Bohr) (Figs. 2 and S1). We established that the localized $A^+ \cdot G^-_{syn}$ zwitterionic intermediate is formal and cannot be considered as thermodynamically stable structure, since the value of the Gibbs free energy for the reverse barrier of the $A^* \cdot G^*_{syn} \rightarrow A^+ \cdot G^-_{syn}$ transition are negative (-1.31 kcal·mol⁻¹ calculated at the MP2/6-311++G(2df,pd)//B3LYP/6-311++G(d,p) level of theory). This means that there is a single TS (IRC=0.00 Bohr) for the $A^* \cdot G^*_{syn}$ tautomerisation *via* the DPT, that governs the reaction rate and therefore the $A \cdot G^*_{syn}$ product selectivity (Fig. 2).

By comparing the value of the zero-point energy (3.85 kcal·mol⁻¹) of the corresponding vibrational mode, which frequency becomes imaginary at the $TS_{A^*\cdot G^*syn \leftrightarrow A \cdot G^*syn}$, with the electronic activation energy for the reverse reaction of the $A^*\cdot G^*_{syn} \rightarrow A \cdot G^*_{syn}$ DPT tautomerisation (2.08 kcal·mol⁻¹), it was revealed that the $A \cdot G^*_{syn}$ base mispair is dynamically unstable structure [27,28,85]. These

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observations do not depend on the used QM levels of theory (for more details see Table 3). Moreover, the Gibbs free energy of activation for the reverse reaction of the $A^* \cdot G^*_{syn} \rightarrow A \cdot G^*_{syn}$ DPT tautomerisation is negligibly small (*<kT*). This results in extremely short lifetime of the $A \cdot G^*_{syn}$ base mispair that is equal to 2.17 · 10⁻¹³ s (Table 3). Therefore any of the 6 low-frequency intermolecular vibrations [86] (9.9, 21.7, 63.2, 73.1, 103.7 and 111.5 cm⁻¹) can develop during this period of time, since their periods are noticeably greater than this time interval. This observation additionally indicates that the $A \cdot G^*_{syn}$ base mispair is not dynamically stable structure and the tautomeric status of the bases included in the $A^* \cdot G^*_{syn}$ base mispair does not change upon its spontaneous or forced dissociation by the DNA-polymerase machinery: in both cases this process occurs according to the scheme $A^* \cdot G^*_{syn} \xrightarrow{dissociation} A^* + G^*_{syn}$.

1.2. The $G \cdot A_{syn}$ base mispair

Lets now discuss the results of the investigation of the G·A_{syn} tautomerisation *via* the DPT into the G*·A*_{syn} base mispair, involving the G base in the *anti*- and the A base in the *syn*-orientations respectively to the sugar residue. The G·A_{syn} base mispair is non-planar (\angle N1C6(G)C6C5(A_{syn})=-21.0°) structure with C₁ symmetry (Scheme 1). The electronic and Gibbs free energies of interaction between the G and A_{syn} bases in the G·A_{syn} base mispair constitute -17.00 and -2.80 kcal·mol⁻¹, respectively.

Providing the calculations for the tautomerised $G^* \cdot A^*_{syn}$ base mispair and the TS of the $G \cdot A_{syn} \leftrightarrow G^* \cdot A^*_{syn}$ tautomerisation, it was revealed that there is no potential-energy minima associated with these structures. This means that the $G \cdot A_{syn}$ base mismatch, stabilized solely by the two N6H···O6 (5.08) and N1H···N7 (6.12 kcal·mol⁻¹) H-bonds, which total energy makes 65.9 % of the electronic interaction energy, does not tautomerise at all and the tautomerised $G^* \cdot A^*_{syn}$ base mismatch almost instantly converts into the $G \cdot A_{syn}$ base mismatch without any barrier.

In order to clarify such important physico-chemical characteristic of the intermolecular H-bonds as their cooperativity, we have proposed a new approach based on the initial forced stretching of the N6H and N1H atomic groups – donors of the N6H…O6 and N1H…N7 H-bonds, respectively, with subsequent sequential fixation of their length and geometry optimization (Fig. 1). As a result, we found out that the N6H…O6 and N1H…N7 H-bonds are cooperative (Fig. 1).

We established that comparably small amount of the electronic deformational energy ΔE_{def} [16,17,21] should be applied for the G·A_{syn} base mispair (R(H-H)=10.399 Å; α_1 =51.6°; α_2 =38.5°) to aquire the glycosidic sizes of the A·T (R(H-H)=10.132 Å; α_1 =54.2°; α_2 =54.7°) (ΔE_{def} =3.00 kcal·mol⁻¹) and G·C (R(H-H)=10.213 Å; α_1 =53.1°; α_2 =55.1°) (ΔE_{def} =3.61 kcal·mol⁻¹ obtained at the MP2/6-

311++G(2df,pd)//B3LYP/6-311++G(d,p) level of theory) Watson-Crick DNA base pairs, that represents an imperative property for the structural isosterisity with Watson-Crick base pairs and faithful DNA replication (see ref. [21] and bibliography there). Notably, we revealed the formation of third C8H…N2 H-bond at the adjusting of the geometry of the G·A_{syn} base mispair to the sizes of the A·T (0.48 kcal·mol⁻¹) and G·C (0.41 kcal·mol⁻¹) Watson-Crick DNA base pairs.

2. Atomistic mechanism of the A*·G*_{syn}↔A·G*_{syn} tautomerisation *via* the DPT

In order to penetrate into the essence of the energetical, structural, polar, electron-topological and NBO reorganizations of the $A^* \cdot G^*_{syn}$ base mispair and intermolecular interactions stabilizing it along the IRC we have carried out the calculations of the electronic energy, the first derivative of the electronic energy with respect to the IRC, the dipole moment of the base pair, the distances and the angle of the intermolecular H-bonds, the electron density, the Laplacian of the electron density, ellipticity and the energy at the (3,-1) BCPs of the intrapair H-bonds, the NBO charges of the hydrogen atoms involved in the tautomerisation, the glycosidic angles and the distance between the glycosidic hydrogens at each step along the IRC of the $A^* \cdot G^*_{syn} \leftrightarrow A \cdot G^*_{syn}$ tautomerisation. In such a way we obtained the scanning or, in other words, sweeps or profiles [27,28,30,31,48]) of these characteristics *vs.* the IRC (Figs. S1-S6), revealing surprising physico-chemical details of the $A^* \cdot G^*_{syn} \leftrightarrow A \cdot G^*_{syn} \oplus A \cdot$

Analysing the spans of the χ -like conjunctions on the d_{O6H/HN6}, d_{N1H/HN7}, ρ and $\Delta\rho$ profiles of the intermolecular H-bonds (Figs. S3a, S3b and S4b), we have arrived at the conclusion that the A*·G*_{syn} \leftrightarrow A·G*_{syn} tautomerisation proceeds through the *asynchronous concerted mechanism*. Additionally, basing on the changes of the electron density and geometry of the intermolecular H-bonds along the IRC of the A*·G*_{syn} \leftrightarrow A·G*_{syn} tautomerisation, we have allocated 9 key points, that characterise the course of this reaction [69,70,71] (Figs. 2, S3a, S3b and S4b).

Thus, the $A^* \cdot G^*_{syn} \leftrightarrow A \cdot G^*_{syn}$ DPT tautomerisation is initiated by the sequential migration of the first proton at the O6 oxygen atom of the G^*_{syn} enol tautomer in the $A^* \cdot G^*_{syn}$ base mispair (the key point 1; IRC=-8.75 Bohr) along the O6H····N6 H-bond in the direction of the N6 nitrogen atom of the A* imino tautomer. This transition passes through the key points 2 ($\Delta \rho_{N6\cdots H}=0$; IRC=-5.48 Bohr), 3 (ρ_{N6} . $_{H}=\rho_{H-O6}$; IRC=-5.23 Bohr), that coincides with the aforementioned formal TS of the A*·G*_{syn} $\rightarrow A^+ \cdot G^-_{syn}$ transition, 4 ($\Delta \rho_{H\cdots O6}=0$; IRC=-5.09 Bohr) and is accompanied by the breakage of the H-O6 covalent bond and the formation of the new N6-H covalent bond. Consequently, the reaction is arrived at the A⁺·G⁻_{syn} non-stable intermediate. Then, the second mobile proton, localized at the N1⁺ nitrogen atom of the A⁺ protonated base starts out to move towards the N7⁻ nitrogen atom of the G_{syn}⁻ deprotonated base

across the key points 5 ($\Delta \rho_{H...N7}=0$; IRC=-0.19 Bohr), 6 - the TS_{A*·G*syn $\leftrightarrow A \cdot G^*$ syn (IRC=0.00 Bohr), 7 ($\rho_{N1-H}=\rho_{H-N7}$; IRC=0.06 Bohr) and 8 ($\Delta \rho_{N1...H}=0$; IRC=0.31 Bohr), eventually reaching the A·G*_{syn} base mispair (the key point 9; IRC=3.73 Bohr). At this the rupture of the N1-H covalent bond and the formation of the H-N7 covalent bond are observed.}

It is interesting to note that the values of the electron density ρ at the (3,-1) BCPs of the Hbonds in the nine key points involving stationary structures, usually treated as a measure of these interactions [75,76,77], range from 0.038 a.u. up to the 0.167 a.u. (Tables 1 and 2). The values of the Laplacian of the electron density $\Delta \rho$ at the (3,-1) BCPs are positive for all intrapair H-bonds and lie within a wide range from 0.000 a.u. up to the 0.160 a.u. (Tables 1 and 2), demonstrating that the Hbonds are attractive closed-shell ineractions. The value of the ellipticity ε varies in the range (1.56÷5.73)·10⁻².

The classical geometrical criteria are satisfied for all canonical H-bonds in the nine key points along the IRC: the $d_{H\dots N6/N1/N7}$ (1.261÷1.855 Å) and $d_{H\dots O6}$ (1.285÷1.714 Å) distances are less than the sum of corresponding Bondi's [87] van der Waals radii (2.75 and 2.72 Å, respectively). The values of the angle of the H-bonding (161.0÷179.0°) are also the trends of the H-bonds (Tables 1 and 2). These observations indicate that the classical geometrical criteria are fulfilled along the IRC, even in the non-stationary points.

The established nine key points allow us to provide a partitioning of the pathway of the $A^* \cdot G^*_{syn} \rightarrow A \cdot G^*_{syn}$ DPT tautomerisation into three distinct regions: the reactant, TS and product regions, separated by the key points 2 (location of the reaction force minimum) and 8 (location of the reaction force maximum), in which the extrema of the first derivative of the electron energy with respect to the IRC - dE/dIRC – are reached [88] (Figs. 2 and S1).

It was established that *the reactant region*, where structural changes such as bond stretching, angle bending and mutual adjustment of the bases within the $A^* \cdot G^*_{syn}$ base pair take place preparing the complex for the DPT tautomerisation, is located between the key points 1 and 2 (-8.75÷-5.48 Bohr). Examination of the barrier heights shows that the electronic energy necessary to bring the donor and acceptor atoms as close as possible to each other and to acquire such mutual deformation and orientation, that eventually lead to the DPT reaction, that is the electronic energy difference between the key points 2 and 1, is equal to 2.31 kcal·mol⁻¹, representing 51.7 % of the TS electronic energy. The DPT reaction actually occurs at *the TS region*, limited by the key points 2 and 8 (-5.48÷0.31 Bohr), where the A* and G*_{syn} bases lose their chemical individuality and where the structural and electronic rearrangements occur accompanying bond breakages and formations. *The product region*, where the structural relaxation of the reaction complex starts to reach the final A·G*_{syn} base mispair, is located

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between the key points 8 and 9 ($0.31\div3.73$ Bohr), respectively. It was estimated, that comparably small amount of the electronic energy (2.47 kcal·mol⁻¹) releases during the transformation of the key point 8 to the key point 9, that is the energy difference between the key points 8 and 9, representing 55.2 % of the TS energy. Observed data evidence that almost equal amount of energy is spent as at the rebuilding and reorganization of the A* and G*_{syn} bases within the A*·G*_{syn} base pair actually before the initiation of the chemical reaction, so at the relaxation of the key point 8 to the final A·G*_{syn} base mispair.

In the course of the investigation we established the electronic deformational energies of the A*·G*_{syn} (R(H-H)=10.411 Å; α_1 =50.3°; α_2 =37.5°) and A·G*_{syn} (R(H-H)=10.211 Å; α_1 =53.0°; α_2 =39.6°) base mispairs [16,17,21], which are inversely proportional to the probability of their incorporation into the DNA double helix: ΔE_{def} (A·T)=3.18 and ΔE_{def} (G·C)=3.72 for the A*·G*_{syn} base mispair and ΔE_{def} (A·T)=4.68 and ΔE_{def} (G·C)=5.54 kcal·mol⁻¹ for the A·G*_{syn} base mispair obtained at the MP2/6-311++G(2df,pd)//B3LYP/6-311++G(d,p) level of QM theory. The presented above values indicate that these DNA base mispairs most closely resemble the geometry of the A·T Watson-Crick base pair and that the substantial part of the deformational energy is spent on the adjusting of the glycosidic angles.

3. Evolution of the selected physico-chemical characteristics of the DNA base mispairs and of the intermolecular specific contacts stabilizing them along the IRC of the $A^* \cdot G^*_{syn} \rightarrow A \cdot G^*_{syn}$ DPT tautomerisation

Notably, the electron energy profile of the $A^* \cdot G^*_{syn} \leftrightarrow A \cdot G^*_{syn}$ DPT tautomerisation is doublewell (Fig. S1a) and zero vibrational level of the vibrational mode in the $A \cdot G^*_{syn}$ base mispair (2690.4 cm⁻¹), which frequency becomes imaginary in the $TS_{A^* \cdot G^*syn \leftrightarrow A \cdot G^*syn}$ of the DPT tautomerisation, lies above the barrier of the DPT tautomerisation.

It is important to note that the O6H···N6 and N1H···N7 H-bonds in the A*·G*_{syn} base mispair exist between the key points 1 and 2 and between the key points 1 and 5, reaching their maximal energies at the key points 2 and 5, for which $\Delta \rho=0$, respectively. The N6H···O6 and N7H···N1 H-bonds in the A·G*_{syn} base mispair exist within the range from the key point 4 to the key point 9 and from the key point 8 to the key point 9, reaching their maximal energies at the key points 4 and 8, for which $\Delta \rho=0$, respectively (Fig. S3d).

Analysis of the dependencies of the H-bond energies on the IRC presented in Figure S3d allows us to reach a clear conclusion that the O6H···N6/N6H···O6 H-bonds and the N1H···N7/N7H···N1 H-bonds are cooperative with each other [27,28,29] and mutually strengthen each other in the $A^*\cdot G^*_{syn}$ ($dE_{N1H···N7}/dE_{O6H···N6}=0.92$ at the IRC=-8.75 Bohr) and $A\cdot G^*_{syn}$ ($dE_{N6H···O6}/dE_{N7H···N1}=0.21$ at the IRC=3.73 Bohr) base mispairs, respectively.

The graphs of the electron density ρ and the Laplacian of the electron density $\Delta\rho$ at the (3,-1) BCPs of the N6-H and H-O6 bonds intersect exactly at the key point 3, while of the N1-H and H-N7 bonds – at the key point 7 (Figs. S3a and S3b). Moreover, the shapes of these profiles are similar to each other, indicating a strong correlation between them [77]. In contrast, the graph of the ellipticity ε at the (3,-1) BCP of the N7-H bond crosses with the graphs for the N1-H and N6-H bonds between the key points 5 and 6 (Fig. S3c) and with the O6-H bond – nearby the key point 8, indicating the extreme sensitivity of the intermolecular bonds to the dynamical behavior of the base pair and the modulation of their energies by the low-frequency intermolecular vibrations of the base pair that tautomerises [86]. The values of the ρ (0.038÷0.312 a.u.), $\Delta\rho$ (-2.003÷0.161 a.u.) and ε (0.012÷0.057) parameters lie within a wide range of values and are in line with the results, presented in our recent works [27,28,30,31].

We have established significant alterations of the absolute value of the dipole moment μ of the studied base pairs within the range of values 6.42÷10.79 D along the IRC of the A*·G*_{syn} \leftrightarrow A·G*_{syn} tautomerisation (Fig. S2). It is evident from the profile of the dipole moment that the tautomerised A·G*_{syn} base mispair is more polar than the starting complex – the A*·G*_{syn} base mispair and a significant increasing of the polarity of the system is observed in the vicinity of the TS_{A*·G*syn} \leftrightarrow A·G*_{syn}, that can be explained by the zwitterionic structure of the base pair.

The obtained sweeps of the NBO charges [89] of the hydrogen atoms localized between the N6 and O6, N1 and N7 atoms along the IRC of the $A^* \cdot G^*_{syn} \leftrightarrow A \cdot G^*_{syn}$ tautomerisation *via* the DPT are presented in Fig. S5. Interestingly, the graph for the NBO charge of the hydrogen atom involved in the O6H…N6 H-bond lies above the graph for the N1H…N7 H-bond at the reactant region and then, changing places at the TS region, the graph for the N6H…O6 H-bond starts to be settled down under the graph for the N7H…N1 H-bond at the product region. The sharp peaks on the sweeps of the NBO charges are observed nearby the key points 4 and 5-8, reflecting the proton movements (Fig. S5).

We revealed the compression, or so-called "breathing", of the starting $A^* \cdot G^*_{syn}$ base mispair throughout the tautomerisation process (Figs. 2, S4a and S6a). It is believed that this phenomenon occurs due to the contraction of the distance between the O6 oxygen and N6 nitrogen atoms at the key point 3 (by 0.289 Å) and between the N1 and N7 nitrogen atoms at the key point 6 (by 0.304 Å) (Fig. S4a). Moreover, the H-bond angles and glycosidic parameters vary at the compression of the $A^* \cdot G^*_{syn}$ base mispair along the IRC of the $A^* \cdot G^*_{syn} \leftrightarrow A \cdot G^*_{syn}$ tautomerisation, that to a large extent is observed at the TS region: $\angle O6HN6$ (161.1÷173.5°); $\angle N1HN7$ (172.9÷179.0°) (Fig. S4a); R(H-H) (9.966÷10.399 Å); α_1 ($\angle N9H(A)H(G_{syn})$) (50.4÷53.8°) and α_2 ($\angle N9H(G_{syn})H(A)$) (36.4÷40.2°) (Fig. S6).

Concluding Remarks

B3LYP and MP2 levels of QM theory, QTAIM analysis, NBO theory and the methodology of the sweeps of the energetical, electron-topological, geometrical, polar and NBO parameters were used to study the structure, barrier heights, thermodynamics, kinetics, electron-topological, polar and NBO properties of the $A^* \cdot G^*_{syn} \leftrightarrow A \cdot G^*_{syn}$ tautomerisation *via* the DPT.

It was found that the $A^* \cdot G^*_{syn}$ DNA base mispair is stabilized by the cooperative O6H…N6 (9.71) and N1H…N7 (6.99) H-bonds, while the $A \cdot G^*_{syn}$ base mispair – by the cooperative N6H…O6 (7.23) and N7H…N1 (9.59 kcal·mol⁻¹) intermolecular H-bonds.

By elaborating a novel approach to the definition of the cooperativity of the intermolecular Hbonds in the H-bonded complexes, which in principle can not tautomerise, it was found that the N6H···O6 (5.08) and N1H···N7 (6.12 kcal·mol⁻¹) intermolecular H-bonds in the G·A_{syn} base mispair are cooperative.

The 9 key points were detected and completely investigated along the IRC of the $A^* \cdot G^*_{syn} \leftrightarrow A \cdot G^*_{syn}$ tautomerisation *via* the DPT.

We have established for the first time that the $G \cdot A_{syn}$ and $A^* \cdot G^*_{syn}$ DNA base mispairs causing spontaneous transversions radically differ from each other in their ability to tautomerise through the DPT. The $A^* \cdot G^*_{syn}$ base mispair (C_s) tautomerises into the $A \cdot G^*_{syn}$ base mispair (C_s) *via* the asynchronous concerted DPT through the non-planar transition state (C_1), whereas the $G \cdot A_{syn}$ base mismatch (C_1) does not tautomerise at all, since there is no local minimum corresponding to the tautomerised $G^* \cdot A^*_{syn}$ base mispair on its electronic (potential) energy surface. The $A^* \cdot G^*_{syn}$ base mispair tautomerises quite easily: the Gibbs free energy of activation for the forward reaction of the $A^* \cdot G^*_{syn} \rightarrow A \cdot G^*_{syn}$ DPT tautomerisation constitutes only $4 \cdot kT$ at room temperature.

It was established that whereas all three investigated $G \cdot A_{syn}$, $A^* \cdot G^*_{syn}$ and $A \cdot G^*_{syn}$ DNA base mispairs are thermodynamically stable structures, since the values of their Gibbs free energies of interaction are less than zero, the $A^* \cdot G^*_{syn}$ base mispair is dynamically unstable H-bonded complex with an extremely low lifetime (2.17 $\cdot 10^{-13}$ s) and any of its 6 low-frequency intermolecular vibrations can develop during this span of time. Consequently, the dissociation of both the $G \cdot A_{syn}$ and $A^* \cdot G^*_{syn}$ DNA base mispairs occurs spontaneously or forcedly by the DNA-polymerase machinery, with the preservation of the tautomeric status of the bases, that form them.

Obtained results allow us to make the following biologically important assumptions. The dissociation of the $G \cdot A_{syn}$ and $A^* \cdot G^*_{syn}$ base mispairs, responsible for the asymmetric transversions, is carried out by the DNA polymerase machinery without changing of the tautomeric status of the bases that form them, i.e. $G \cdot A_{syn} \rightarrow G + A_{syn}$ and $A^* \cdot G^*_{syn} \rightarrow A^* + G^*_{syn}$. In the last case it means that the

 $A \cdot G^*_{syn}$ mismatch "escapes from the hands" of the replication machinery due to its immediate transformation into the $A^* \cdot G^*_{syn}$ base mispair. This finding is important, in our point of view, in terms of the theory of the spontaneous point mutations, namely - of the process of their emergence and consolidation in the DNA biosynthesis in a living cell.

Electronic Supplementary Material. Sweeps of the electronic energy, the first derivative of the electronic energy with respect to the IRC, the dipole moment of the base pair, the electron density, the Laplacian of the electron density, the ellipticity and the energy at the (3,-1) BCPs of the intrapair H-bonds, the distances and the angle of the intermolecular H-bonds, the NBO charges of the hydrogen atoms involved in the intermolecular interactions and the glycosidic parameters (angles and distance between the glycosidic hydrogens) of the base mispair along the IRC of the A*·G*_{syn} \leftrightarrow A·G*_{syn} tautomerisation *via* the DPT obtained at the B3LYP/6-311++G(d,p) level of QM theory in vacuum are available free of charge *via* the Internet at http://rsc.org.

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$F A = O_{syn}, A = O_{syn}, O = A_{syn}$ and $F = S_{syn} + A = G^{syn} + O = O = O = O = O = O = O = O = O = O$										
Complex	AH…B	ρ^{a}	$\Delta \rho^{\rm b}$	$100 \cdot \varepsilon^{c}$	$d_{A \cdots B}^{d}$	$d_{H \cdots B}^{e}$	$\angle AH \cdots B^{f}$	Δd_{AH}^{g}	$\Delta v^{\rm h}$	$E_{\mu\nu}^{i}$
	H-bond									
A*•G* _{syn}	O6H…N6	0.057	0.103	4.56	2.645	1.664	161.0	0.048	905.4	9.71
	$N1H \cdots N7$	0.038	0.094	5.72	2.894	1.855	179.5	0.028	489.3	6.99
A·G* _{syn}	N6H···O6	0.042	0.135	2.64	2.745	1.714	172.4	0.030	520.3	7.23
	$N7H \cdots N1$	0.051	0.093	5.73	2.801	1.739	175.2	0.051	884.3	9.59
TS _{A*·G*syn↔A·G*syn}	N6H···O6	0.065	0.155	2.33	2.614	1.553	172.4	-	-	12.62*
G·A _{syn}	N6H···O6	0.029	0.104	2.85	2.886	1.874	169.9	0.016	276.9	5.08
	N1H···N7	0.032	0.087	5.50	2.958	1.926	175.6	0.022	383.8	6.12

Table 1 Electron-topological, structural, vibrational and energetical characteristics of the intermolecular H-bonds revealed in the $A^* \cdot G^*_{\text{syn}}$, $A \cdot G^*_{\text{syn}}$, $G \cdot A_{\text{syn}}$ and $TS_{A^* \cdot G^*_{\text{syn}}}$ obtained at the B3LYP/6-311++G(d,p) level of theory in vacuum

^aThe electron density at the (3,-1) BCP, a.u.

^bThe Laplacian of the electron density at the (3,-1) BCP, a.u.

^cThe ellipticity at the (3,-1) BCP

^dThe distance between the A (H-bond donor) and B (H-bond acceptor) electronegative atoms of the H-bonds, Å

^eThe distance between the H and B atoms of the H-bonds, Å

^fThe H-bond angle, degree

^gThe elongation of the H-bond donating group AH upon H-bonding, Å

^hThe redshift of the stretching vibrational mode v(AH), cm⁻¹

¹The energy of the H-bonds, calculated by Iogansen's [78] and Nikolaienko-Bulavin-Hovorun (marked with an asterisk) [79] formulas, kcal mol⁻¹

Table 2 Electron-topological and structural characteristics of the intermolecular bonds revealed in the 9 key points and the polarity of the latters along the IRC of the $A^* \cdot G^*_{syn} \leftrightarrow A \cdot G^*_{syn}$ DPT tautomerisation obtained at the B3LYP/6-311++G(d,p) level of theory in vacuum

Complex	AH…B H-bond/	ρ	Δho	100·e	$d_{A \cdots B}$	$d_{H \cdots B}$	∠AH…B	μ
	A-H-B covalent bond							
Key point 1 (-8.75 Bohr):	O6H···N6	0.057	0.103	4.56	2.645	1.664	161.0	6.43
A*•G* _{syn}	$N1H \cdots N7$	0.038	0.094	5.72	2.894	1.855	179.5	Ç
Key point 2 (-5.48 Bohr):	O6H···N6	0.117	0.000	3.66	2.464	1.384	164.1	7.37
$\Delta \rho_{\text{N6}\cdots\text{H}}=0$	N1H···N7	0.052	0.103	5.04	2.776	1.723	176.6	
Key point 3 (-5.23 Bohr):	O6-H-N6	0.162	-0.275	3.29	2.461	1.261	166.0	8.63
ρ _{N6-H} =ρ _{H-O6}	N1H···N7	0.053	0.099	4.94	2.773	1.713	176.7	
Key point 4 (-5.09 Bohr):	N6H···O6	0.135	0.000	1.56	2.465	1.285	166.4	9.39
Δρ _{H···O6} =0	$N1H \cdots N7$	0.054	0.097	4.88	2.771	1.707	176.8	2
Key point 5 (-0.19 Bohr):	N6H···O6	0.067	0.152	2.31	2.610	1.542	172.4	10.23
$\Delta \rho_{\rm HN7}=0$	$N1H \cdots N7$	0.112	0.000	3.50	2.588	1.418	173.9	ſ
Key point 6 (0.00 Bohr):	N6H···O6	0.065	0.155	2.33	2.614	1.553	172.4	9.49
$TS_{A^* \cdot G^* syn \leftrightarrow A \cdot G^* syn}$								C
Key point 7 (0.06 Bohr):	N6H···O6	0.065	0.156	2.34	2.615	1.556	172.4	9.24
$\rho_{\text{N1-H}} = \rho_{\text{H-N7}}$	N1-H-N7	0.154	-0.217	2.98	2.584	1.293	174.7	
Key point 8 (0.31 Bohr):	N6H···O6	0.062	0.160	2.36	2.618	1.568	172.5	8.39
$\Delta \rho_{N1\cdots H}=0$	N7H···N1	0.111	0.000	4.12	2.588	1.419	175.0	
Key point 9 (3.73 Bohr):	N6H···O6	0.042	0.135	2.64	2.745	1.714	172.4	7.75
A·G* _{syn}	N7H···N1	0.051	0.093	5.73	2.801	1.739	175.2	č

Notes: For footnote definitions see Table 1. μ – the dipole moment of the complex, D.

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Table 3 Energetic and kinetic characteristics of the $A^* \cdot G^*_{syn} \leftrightarrow A \cdot G^*_{syn}$ tautomerisation via the DPT in vacuo (T=298.15 K) obtained at the
different levels of QM theory for the geometry calculated at the B3LYP/6-311++G(d,p) level of QM theory

Level of OM theory	ΛG^{a}	٨F ^b	AAG°	AAEd	$\Lambda\Lambda G^{e}$	$\Delta\Delta E^{\rm f}$		τ^{g}
Level of Qivi theory	20	$\Delta \mathbf{L}$		ΔΔLTS	ΔΔΟ	kcal·mol ⁻¹	cm ⁻¹	ι
MP2/6-311++G(2df,pd)	1.83	2.14	2.03	4.21	0.20	2.08	726.4	$1.25 \cdot 10^{-13}$
MP2/6-311++G(3df,2pd)	1.71	2.02	2.16	4.35	0.45	2.33	815.5	$1.92 \cdot 10^{-13}$
MP2/cc-pVTZ	2.67	2.98	3.07	5.25	0.40	2.27	795.2	$1.75 \cdot 10^{-13}$
MP2/cc-pVQZ	1.89	2.20	2.42	4.60	0.52	2.40	840.2	$2.17 \cdot 10^{-13}$

^aThe Gibbs free energy of the A·G*_{syn} base mispair relatively to the A*·G*_{syn} base mispair ($\Delta G_{A*\cdot G*syn}=0.00$), kcal·mol⁻¹

^bThe electronic energy of the A·G*_{syn} base mispair relatively to the A*·G*_{syn} base mispair ($\Delta E_{A*\cdot G*_{syn}}$ =0.00), kcal·mol⁻¹

^cThe Gibbs free energy of activation for the forward reaction of the $A^* \cdot G^*_{syn} \rightarrow A \cdot G^*_{syn}$ tautomerisation, kcal·mol⁻¹

^dThe activation electronic energy for the forward reaction of the $A^* \cdot G^*_{syn} \rightarrow A \cdot G^*_{syn}$ tautomerisation, kcal·mol⁻¹

^eThe Gibbs free energy of activation for the reverse reaction of the $A^* \cdot G^*_{syn} \rightarrow A \cdot G^*_{syn}$ tautomerisation, kcal·mol⁻¹

^fThe activation electronic energy for the reverse reaction of the $A^* \cdot G^*_{syn} \rightarrow A \cdot G^*_{syn}$ tautomerisation

^gThe lifetime of the $A \cdot G^*_{syn}$ base mispair, s

The frequency of the vibrational mode in the $A \cdot G^*_{syn}$ base mispair, which becomes imaginary in the $TS_{A^* \cdot G^* syn \leftrightarrow A \cdot G^* syn}$ of the $A^* \cdot G^*_{syn} \leftrightarrow A \cdot G^*_{syn}$ DPT tautomerisation, is equal to 2690.4 cm⁻¹ and the zero-point vibrational energy associated with this normal mode is equal to 3.85 kcal·mol⁻¹ or 1345.2 cm⁻¹ obtained at the B3LYP/6-311++G(d,p) level of QM theory.

The Gibbs free energy of the G·A_{syn} base mispair relatively to the A*·G*_{syn} base mispair ($\Delta G_{A*\cdot G*syn} = \Delta E_{A*\cdot G*syn} = 0.00$) is equal to $\Delta G_{G\cdot Asyn} = -8.91$ kcal·mol⁻¹ and the electronic energy – is equal to $\Delta E_{G\cdot Asyn} = -9.70$ kcal·mol⁻¹ obtained at the MP2/6-311++G(2df,pd)//B3LYP/6-311++G(d,p) level of QM theory *in vacuo*.

Graphical Abstract.

The $A^* \cdot G^*_{syn} \leftrightarrow A \cdot G^*_{syn}$ DPT tautomerisation proceeds through the asynchronous concerted mechanism; the $G \cdot A_{syn}$ DNA base mispair does not tautomerise at all.

