

PCCP

Accepted Manuscript



This is an *Accepted Manuscript*, which has been through the Royal Society of Chemistry peer review process and has been accepted for publication.

Accepted Manuscripts are published online shortly after acceptance, before technical editing, formatting and proof reading. Using this free service, authors can make their results available to the community, in citable form, before we publish the edited article. We will replace this *Accepted Manuscript* with the edited and formatted *Advance Article* as soon as it is available.

You can find more information about *Accepted Manuscripts* in the [Information for Authors](#).

Please note that technical editing may introduce minor changes to the text and/or graphics, which may alter content. The journal's standard [Terms & Conditions](#) and the [Ethical guidelines](#) still apply. In no event shall the Royal Society of Chemistry be held responsible for any errors or omissions in this *Accepted Manuscript* or any consequences arising from the use of any information it contains.

Cite this: DOI: 10.1039/c0xx00000x

www.rsc.org/xxxxxx

ARTICLE TYPE

Ligand Field Density Functional Theory for the prediction of future domestic lighting †

Harry Ramanantoanina,^a Werner Urland,^{*a,b} Amador García-Fuente,^a Fanica Cimpoesu^{*c} and Claude Daul^{*a}

Received (in XXX, XXX) Xth XXXXXXXXX 20XX, Accepted Xth XXXXXXXXX 20XX

DOI: 10.1039/b000000x

We deal with the computational determination of the electronic structure and properties of lanthanide ions in complexes and extended structures having open-shell f and d configurations. Particularly we present conceptual and methodological issues based on Density Functional Theory (DFT) enabling the reliable calculation and description of the f → d transitions in lanthanide doped phosphors. We consider here the optical properties of the Pr³⁺ ion embedded into various solid state fluoride host lattices, for the prospecting and understanding of the so-called quantum cutting process, being important in the further quest of warm-white light source in light emitting diodes (LED). We use the conceptual formulation of the revisited ligand field (LF) theory fully compatibilized with the quantum chemistry tools: LFDFT. We present methodological advances for the calculations of the Slater-Condon parameters, the ligand field interaction and the spin-orbit coupling constants, important in the non-empirical parameterization of the effective Hamiltonian adjusted from the ligand field theory. The model shows simple procedure using less sophisticated computational tools, which is intended to contribute to the design of modern phosphors and to help to complement the understanding of the 4fⁿ → 4fⁿ⁻¹5d¹ transitions in any lanthanide system.

Introduction

The analysis of the f → d electric dipole allowed transitions in lanthanide compounds has been the subject of tremendous interest amongst spectroscopists and theoreticians for a long time.¹⁻⁵ The application relevance of this issue is related to the functioning of lanthanide ions in light emitting diodes (LED) in order to generate domestic lighting. Nowadays the ban on the incandescent light bulbs being effective, artificial white light source is for instance obtained by the combination of a gallium nitride blue LED with an inorganic phosphor, *i.e.* a luminescent material such as the system yttrium aluminium garnet Y₃Al₅O₁₂ (YAG) doped with Ce³⁺ ion.¹ However, the f → d transitions obtained in this way are rather limited to the yellowish emission of Ce³⁺ combined with the blue emission of the LED. Therefore the generated white light looks always bluish cold. The concept of warm-white light is introduced in order to improve and to adjust the spectrum of known inorganic phosphors to match qualitatively the solar day light. A LED coated with two or three different phosphors aside the available Ce³⁺-doped yellow phosphor should be an appropriate choice to generate the future warm-white light. In this context the interest on the Pr³⁺ ion, which emits red light leading to the warm impression, is becoming stronger. Phosphors with Pr³⁺ ion are already in the spotlight of nowadays

experimental studies. The chemical and physical stabilities of the materials, the efficiency of the visible light emission are currently investigated⁶ even through the synthesis and preparation of the materials remains problematic due to availability and cost prices, which have to be considered. On the other hand, theoretical modelling offers a complementary vision, since its application does not require any laborious experiments and costly equipment, the computational experiments saving time by valuable preamble prospecting and offering causal insight in the property design goals.

Theoretical studies include for instance: the modelling of the solid state crystal structure doped with the Pr³⁺ ion, in order to make explicit the possibility of lattice or local disorder of the doped materials;⁷ the determination of the electronic structure of the Pr³⁺ ion surrounded by the ligands system; the prediction of the optical properties, which can be directly compared to the experimental data or can be prospectively used to the design of new phosphors. It is also of great importance to illustrate the theoretical results more conceptually toward crystal/ligand field theory,^{8,9} which as matter of principle gives a significant chemical insight, considering the theoretical support as a sort of black box. Nevertheless it is noteworthy to highlight the considerable efforts in this context made by the theoreticians¹⁰⁻¹⁵ inasmuch as the possibility to make a comparison between different approaches to the problem is now established.^{16,17,18}

Cite this: DOI: 10.1039/c0xx00000x

www.rsc.org/xxxxxx

ARTICLE TYPE

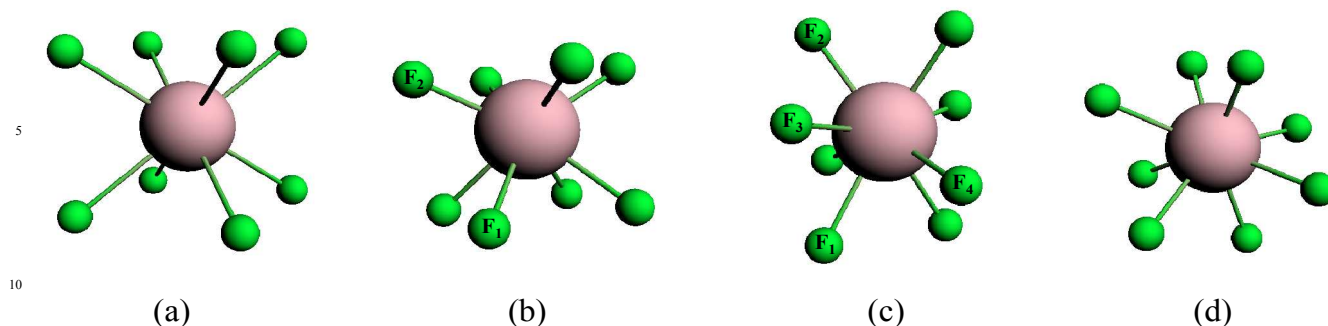


Fig. 1 Structure and spatial representation of the eight coordinated Pr^{3+} complexes being investigated in the present work, resulting from Pr^{3+} doped into: CaF_2 (a), KY_3F_{10} (b), BaY_2F_8 (c) and LaZrF_7 (d) forming a local symmetry of O_h , C_{4v} , C_2 and C_1 point groups, respectively.

Amongst various theoretical approaches to the problem, including empirical fit of experimental data to phenomenological Hamiltonian^{14,15} to *ab-initio* methodologies,^{10,13} Density Functional Theory (DFT) offers a better compromise given the acknowledged advantages of this frame in many applications. Moreover the particular procedure of controlling the spin and the orbital occupation schemes in DFT simplifies the computation and enables DFT to be the method of choice adapted to this particular problem. We have developed the Ligand Field DFT (LFDFT) approach for the analysis of the spectroscopy and magnetic properties¹⁹⁻³⁰ of an ion centre surrounded by a ligand system, using the early concept of the multideterminantal approach.³¹ This multideterminantal approach is based on the Average Of Configuration (AOC) occupation of the Kohn-Sham orbital, when we are dealing with open-shell configurations. This procedure was recently extended to account for two-open-shell f and d electrons,³² and has produced a good prediction of the f \rightarrow d transitions in Ce^{3+} and Pr^{3+} -doped compounds.^{32,33} The crystal structures with fluoride ligands are mostly suitable for inorganic phosphor candidates due to their relative stability and their ability to avoid phonons upon the f \rightarrow d transitions,⁶ ergo to dismiss radiationless transitions. While the Pr^{3+} ion is doped into certain fluoride host lattices, a particular yet interesting optical phenomenon is experimentally observed, the so-called quantum cutting process.^{6,34,35} The exhibition of such an optical effect is due to the fact that the spectral terms originating from the $4f^2$ and $4f^15d^1$ electron configurations of the Pr^{3+} are well separated on the energy scale, allowing a photon cascade emission. An overlap between both electron configurations is avoided insofar as the excitation of the materials with a high energy photon gives rise to several emitted photons leading to a quantum yield far above 100 %. The phenomenology of the quantum cutting process itself is not always acquired, since it is strongly dependent on the chemical environment of the Pr^{3+} ion. This chemical environment induces the nephelauxetic effect, which is large on the 5d orbitals of the Pr^{3+} ion but small for the 4f ones. The larger the nephelauxetic effect, the smaller is the difference between the

average energy position of the $4f^2$ and $4f^15d^1$ states of the Pr^{3+} ion, *i.e.* an overlap between those spectral terms might be possible. Accordingly host lattices with fluoride ligand would be the system of choice taking into account the spectrochemical series of the ligands.

In the extension of our previous work detailing the non-empirical prediction of the $4f^n \rightarrow 4f^{n-1}5d^1$ transitions,³² we present in this paper supplementary tools for a better parameterization of the inter-electron repulsion interaction in the ligand field theory Hamiltonian. Previously, we have extracted the Slater-Condon parameters by least squares fit from the energies of the 231 Slater-Determinants originating from the $4f^2$ (91 microstates) and $4f^15d^1$ electron configurations (140 microstates) of Pr^{3+} , where we have misrepresented the energies of some spectral terms as stressed in ref.³² In the present work, we consider the determination of the Slater-Condon parameters directly from the radial functions of the 4f and 5d Kohn-Sham orbitals. This methodological update solves the earlier problem, making now our methodology only dependent on the change of the DFT setting in the computational details and the nature³⁶ of the Kohn-Sham orbitals.

Herein we report a detailed discussion of the reliable prediction and understanding of the $4f^2 \rightarrow 4f^15d^1$ transitions considering luminescent materials like Pr^{3+} -doped fluoride host lattices in order to provide a prospective description of the optical behaviour of modern warm-white light phosphors. Various fluoride host lattices doped with Pr^{3+} ion are investigated taking into account the local symmetry of the lanthanide centre together with its coordination sphere, as well as the availability of optical spectra. Hence, a special emphasis is addressed to systems, which have been experimentally investigated before, in order to validate the theoretical approach. In this respect, we choose Pr^{3+} coordinated with eight fluoride ligands as experimentally observed in the following systems: $\text{CaF}_2:\text{Pr}^{3+}$,³⁷ $\text{KY}_3\text{F}_{10}:\text{Pr}^{3+}$,³⁸ $\text{BaY}_2\text{F}_8:\text{Pr}^{3+}$,^{38,39} and $\text{LaZrF}_7:\text{Pr}^{3+}$,⁴⁰ where the eight fluoride ligands form arrangement with O_h , C_{4v} , C_2 and C_1 symmetry, respectively (Fig.1).

Methodology

Ligand field parameters

The theory of the electronic structure of metal complexes and the concept of the ligand field theory are exhaustively explained elsewhere.^{8,9} Hereafter we are just giving a brief account of the ligand field parameters. The origin of the ligand field parameters comes from the fact that it is almost impossible to represent the full Hamiltonian operator, corresponding to a metal centre surrounding by the ligand system, to calculate the exact energies and the projected wave functions in the model space. The ligand field parameters allow to construct an effective Hamiltonian operator, which is as precise as the earlier full Hamiltonian with much smaller dimensionality. The model effective Hamiltonian is always defined in a perturbation approach of the wave function of the metal ion. Within this model, two interactions might be highlighted: first, the spherical symmetric interaction due to the inter-electron repulsion in the metal ion; secondly, the symmetry dependent ligand field interaction, which accounts for the influence of the chemical environment onto the metal ion, together with the spin-orbit coupling interaction.

The inter-electron interaction is determined from the two-electron repulsion integral,

$$\langle ab|cd\rangle = \int \Psi_a^*(r_1)\Psi_b^*(r_2)\frac{1}{r_{12}}\Psi_c(r_1)\Psi_d(r_2)dr_1dr_2 \quad (1)$$

where Ψ_i denotes the multiplet wave function that belongs to a specific configuration. Considering the Green's function expansion of the r_{12}^{-1} operator as described in most quantum chemistry textbooks,^{8,9,41} eq. 1 can be reformulated in terms of a product of an angular and a radial integrals. These two integrals are then solved using first appropriate Clebsch-Gordan coefficients for the angular one and secondly, few parameters W^k (k is an integer) because of the well known Slater's rules, for the radial one, in their definition as Condon-Shortley radial parameters:

$$W^k(ab, cd) = \int_0^\infty \int_0^\infty \frac{r_1^k}{r_{12}^{k+1}} R_a^*(r_1)R_b^*(r_2)R_c(r_1)R_d(r_2)r_1^2r_2^2dr_1dr_2 \quad (2)$$

where $R_i(r)$ represent the radial functions of atomic shells and $r_<$ and $r_>$ are the lesser and greater of the distances r_1 and r_2 of two electrons from the nuclei.⁴¹ The two-electron parameters known as Slater-Condon parameters F^k and G^k according to the nature of the interaction to be Coulomb and exchange, respectively, are obtained from eq. 2 using the following relationships:

$$F^k(ab) = W^k(ab, ab) \quad (3)$$

and

$$G^k(ab) = W^k(ab, ba) \quad (4)$$

For the problem of lanthanide ions, which are characterized by n valence electrons, one may define the F^0 , F^2 , F^4 and F^6 Slater-Condon parameters, corresponding to the ground $4f^n$ electron configuration of the lanthanide ion. Furthermore a set of new parameters appears while dealing with the excited $4f^{n-1}5d^1$ electron configuration, such as F^0 , G^1 , F^2 , G^3 , F^4 and G^5 . These parameters can be related to the $f \otimes f$ and $f \otimes d$ diadic

products, where the G^k parameters in the ground electron configuration vanish due to the *gerade* symmetry of the direct product. A parameter Δ_0 is also used to represent the energy gap between the ground and the excited electron configurations. This parameter is composed from different contributions, which implies not only the two-electron repulsions F^0 in both electron configurations but also a one-electron part, *i.e.* kinetic plus electron-nuclear terms. For the sake of clarity, (*ff*) and (*fd*) notations are added next to the parameters to emphasize the nature of the parameters to originate from the ground or excited electron configuration of the lanthanide ion, respectively. The evaluation of Eq. 2 has already been discussed earlier in the work of Morrison and Rajnak.⁴⁴ Using Hartree-Fock method, high overestimation of the parameters were obtained. Therefore it was stated that a proper evaluation of Eq. 2 requires electron correlation *via* configuration interaction between the 4f shell of lanthanides with inner shells, which is in a certain way retrieved in the frame of our multideterminantal DFT based on hybrid B3LYP functional.

The ligand field interaction, which considers the one-electron part, takes full account for the lowering of the symmetry due to the chemical environment of the lanthanide ion. There are several conventions for the interpretation of the ligand field interaction.⁴²

The classical approach is to represent the ligand field potential in the basis of spherical harmonic functions, taking into account the position of the ligands in the coordination sphere of the lanthanide ion. In this formulation, the well-known Wybourne-normalized crystal field parameters B_q^k ,⁴³ which are in general complex numbers to be placed in front of the spherical harmonic functions or can be reformulated as real if the expansion is converted to real-type harmonics, is used. Within the manifold of the merged 4f and 5d orbitals of the lanthanide ion, there are 28 parameters inside the (*ff*) part, $B_q^k(f)$, 15 aside (*dd*)-type $B_q^k(d)$ ones and 21 inter-shell $B_q^k(fd)$ terms. The count of parameters is obtained considering the q index running on the $2k+1$ components for the $k = 0, 2, 4, 6$ for the (*ff*) case, for the $k = 0, 2, 4$ for the (*dd*) case and for the $k = 1, 3, 5$ in the inter-shell block. In the one-shell case the $k = 0$ component which accounts for an overall shift of the energy scheme, is dropped by a conventional placement of the LF split with the centre of gravity in the zero-energy point. However, in two-shell problems at least one zero-order parameters should be kept to represent the f-d gap in the one-electron parameterization. In certain circumstances many of Wybourne parameters may vanish due to symmetry constrains. For instance in a ligand field of D_{4d} symmetry a limited number of B_q^k parameters remains, such as the $B_0^0(f)$, $B_0^2(f)$, $B_0^4(f)$, $B_0^6(f)$, $B_0^0(d)$, $B_0^2(d)$, $B_0^4(d)$, $B_0^6(d)$, $B_0^3(fd)$ and $B_0^5(fd)$ parameters. The labels (*f*), (*d*) and (*fd*) next to the parameters emphasize the origin of the parameters to be a perturbation of the 4f, the 5d and both 4f, 5d orbitals of the lanthanide ion, respectively. The ligand field interaction can also be described with respect to the old Angular Overlap Model (AOM),^{45,46} requiring in the present case two-open-shell f and d.

In this paper our theoretical work for the prediction of the $f \rightarrow d$ transitions involves four steps:

(i) the determination of the geometry of the Pr^{3+} -doped system using plane wave calculation involving the periodicity of the infinite crystal system;

(ii) the calculation of the Slater-Condon parameters using the radial function of the Kohn-Sham eigenfunction obtained from an AOC calculation. An AOC calculation is a molecular DFT calculation, which takes into account all interactions, *i.e.* overlap, plus electrostatic terms. The $F^k(ff)$ parameters are determined with equal occupation of two electrons in the 4f orbitals of Pr^{3+} , whereas the $F^k(fd)$ and $G^k(fd)$ are obtained from equal occupation of the 4f and 5d orbitals, *i.e.* $4f^15d^1$ electron configuration. The evaluation of the $F^k(dd)$ parameters from the $5d^2$ configuration is, in fact, not necessary, since we do not consider the very high spectral terms with this parentage. Special attention is addressed to the calculation of the $F^0(ff)$ and $F^0(fd)$ parameters, which account to the parameter $\Delta_0(fd)$, and are simply the barycentres of the energies of the Slater-Determinants originating from the $4f^2$ and $4f^15d^1$ electron configurations of the Pr^{3+} ion, respectively;

(iii) the computation of the ligand field potential from the energies of the 91 and 140 Slater-Determinants originating from the $4f^2$ and $4f^15d^1$ electron configurations of the Pr^{3+} ion, using least squares fit;

(iv) the prediction of the $f \rightarrow d$ transitions diagonalizing the ligand field theory Hamiltonian suitable for our investigation.

Computational details

The DFT calculation reported in this paper have been carried out by means of the Amsterdam Density Functional (ADF2010) program package.^{47,48,49} The LFDFT calculation has been performed using DFT calculations based on the hybrid B3LYP functional as it is implemented in the ADF program package^{47,48,49} for the exchange and correlation energy and potential. We use the B3LYP in line with our earlier work³² detailing the accuracy of the prediction in terms of the change of the exchange-correlation functional in the DFT setting. The molecular orbitals were expanded using a triple zeta STO basis sets plus two polarization functions (TZ2P+) for the Pr atom and plus one polarization function (TZP) for the F atom. Positive point charges are taken into account in order to neutralize the highly negative charged structures obtained for the $(\text{PrF}_6)^{5-}$ complex, using the *Efield* keyword available in the ADF program package.^{47,48,49} These positive point charges are placed in the position of the next-nearest neighbours to mimic the crystal hosts. The spin-orbit coupling constants were calculated for the 4f and 5d orbitals using the approach of ZORA relativistic available in the ADF program package.^{47,48,49} The LFDFT calculation were achieved following the detailed procedure already described in ref.³² where the most important step consists to the representation of the density in a totally symmetric form using the approach of the AOC type calculation. Matlab/Octave codes for the LFDFT program together with the determination of the Wybourne-normalized crystal field parameters from the ligand field matrix are available from the authors upon request. The geometrical structure of surrounding atoms around the Pr^{3+} impurity in the host lattice is obtained via periodical calculations by means of the VASP code,^{50,51} where the PBE exchange-correlation approximation⁵² is used and the interaction between valence and core electrons is emulated with the Projected Augmented Wave method.^{53,54} External as well as semi-core states are included in the valence. A plane-waves basis set with a cutoff energy of 400

eV is used. Unit cells with a number of atoms ranging from 88 to 144 were simulated, which were found to be large enough to lead to negligible interactions between the periodic images of the Pr^{3+} impurity. 3 *k*-points were included in each direction of the lattice, in order to reproduce its periodicity, and the atomic positions were allowed to relax until all forces were smaller than 0.005 eV/Å.

Results and Discussion

The conventional LFDFT procedure requires the calculation of the energies of all the Slater Determinants in the model space, *i.e.* the $4f^n$ and $4f^{n-1}5d^1$ electron configurations of the lanthanide ion in the consideration of two-open-shell f and d electrons.³² This model space is yet reasonable for the case of Ce^{3+} (24 microstates) or Pr^{3+} (231 microstates), but explodes while dealing with Eu^{2+} (33462 microstates) for instance. Phosphors with Eu^{2+} are also potential candidates for warm-white light sources. The large dimensionality is not problematic using the irreducible tensor operator since there are always few Slater Determinants from which the energies of all the spectral terms can be generated. However the LFDFT procedure considers the calculation of the Slater-Condon parameters by least squares fit, which has shown in our previous work a non-homogeneous distribution of the energy values amongst the parameters.³² Thus it was found that the $F^2(ff)$ is overestimated with respect to the $F^4(ff)$ and $F^6(ff)$, which are underestimated if compared to the experimental fitted parameters.³² Nevertheless, the results have shown sensibly good prediction of the quantum cutting process in $\text{Cs}_2\text{KYF}_6:\text{Pr}^{3+}$.³² Aiming to provide a perfect fit with the experimental results and in order to guarantee a reliable prediction of the rather delicate $f \rightarrow d$ transitions in phosphors, we readapt the methodology taking advantages of the Kohn-Sham orbitals, which do also enable information about the two-electron repulsion in the double integration given in eq. 2. The evaluation of this integration (eq. 2) can be solve within the same accuracy either numerically on a grid⁵⁵ or analytically⁵⁶ considering the radial function $R_i(r)$ as a linear combination of single- ζ Slater type orbitals (STO) functions $P_i(r, \zeta_i)$ (eq. 5) available for instance in the ADF program package.^{47,48,49}

$$R_i(r) = \sum_{u=1}^{N_{\zeta}} c_{i,u} P_u(r, \zeta_u) \quad (5)$$

where $c_{i,u}$ stands for the linear combination coefficient in front of a STO $P_i(r, \zeta_i)$ function (eq. 6),

$$P_i(r, \zeta_i) = A_i r^{n_i-1} e^{-\zeta_i r} \quad (6)$$

and A_i , n_i and ζ_i stand for the normalization factor, the principal quantum number and the screening parameter of the STO function, respectively. Within basic mathematical operations, the calculation of the Slater-Condon parameters from Gaussian type orbitals (GTO) is also possible, showing this methodology to be not restricted to basis sets options. To illustrate our model, we present in Fig. 2 the radial functions corresponding to the 4f and 5d Kohn-Sham orbitals considering both free Pr^{3+} ion and the presence of eight fluoride ligands in its coordination sphere. In this latter situation, we consider the case of eight fluoride ligands as surrogated in the cubic ligand field in $\text{CaF}_2:\text{Pr}^{3+}$.

Cite this: DOI: 10.1039/c0xx00000x

www.rsc.org/xxxxxx

ARTICLE TYPE

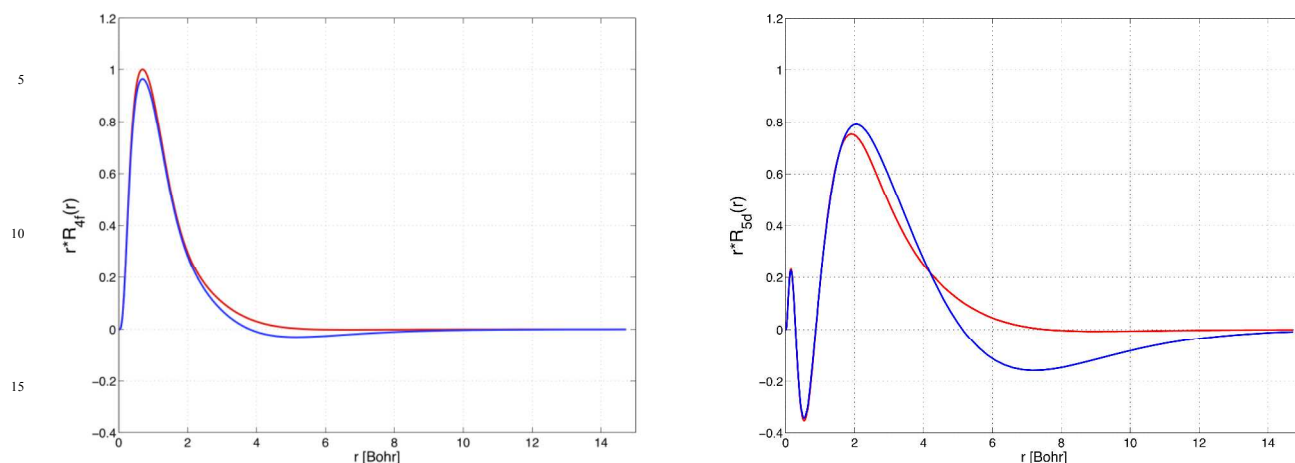


Fig. 2 Graphical representation of the radial functions of the 4f (left hand side) and 5d (right hand side) Kohn-Sham orbitals of the free Pr^{3+} ion (in red) and Pr^{3+} in $\text{CaF}_2:\text{Pr}^{3+}$ (in blue).

It is noteworthy to highlight that the graphical representations in Fig. 2 show $r.R_i(r)$ radial function, which gives more useful and handleable data not only for the numerical integration of eq. 2 but also for the fitting to STO functions (eq. 5) we need for the analytical integration of eq. 2. The change of the 4f radial function from the free to coordinated Pr^{3+} is almost negligible, in contrast to the 5d one (Fig. 2), where a drastically change of its shape is observed. This is in line to the fact that the 4f orbitals are well shielded by the outer shells, unlike the 5d ones. The 5d, which are much more diffuse establish stronger overlap with the ligand orbitals. Therefore the calculated Slater-Condon $F^k(ff)$ parameters from eq. 2 are certainly similar going from free to coordinated Pr^{3+} , although the $F^k(fd)$ and $G^k(fd)$ will be reduced. In Table 1 we present the calculated Slater-Condon parameters for the Pr^{3+} doped phosphors under consideration in the present work. Conventionally, parameters F_k and G_k (cf. Table 1) are of interest and are defined in terms of the F^k (eq. 3) and G^k (eq. 4) analogues using the conversion factor eventually given by Cowan.⁴¹ The values from Table 1 are in the magnitude of the experimental fitted parameters.⁵⁷ The small discrepancy is conceived originated from the DFT setting in the computational details.

In fine the calculation of the Slater-Condon parameters in the way presented in this work constitutes a considerable step in any conceptual formulation of the ever-modern ligand field theory along nowadays quantum chemistry tools. While working in the vicinity of our LFDFT developed approach, this methodological update will also help for the reliable determination of the magnetic properties of lanthanide or transition metal complexes being currently popular in theoretical research.⁵⁸⁻⁶⁵ The formulation of the ligand field interaction, which takes into

account the presence of the ligands in the coordination sphere of the lanthanide, is also the subject of intense work.^{32,44} Here the ligand field potential includes two distinct interactions the perturbation of the 4f and 5d orbitals of the lanthanide, respectively, leading to ligand field matrix of 12 by 12 elements. This ligand field interaction is parameterized by means of the Wybourne parameters, which in total involves 64 B_q^k 's as stated in the method section. We calculate these B_q^k parameters for the Pr^{3+} doped phosphors under consideration and we present the results in Table 2. The zero-order B_0^0 present in both 4f and 5d perturbations are dropped since they already compose the $\Delta_0(fd)$ parameter given in Table 1 leading to the formation of a traceless ligand field matrix. Moreover the display of the B_{-q}^k 's in Table 2 is redundant since B_{-q}^k can be deduced from the analogous B_q^k using their complex conjugate.

Table 1 Calculated Slater-Condon parameters in cm^{-1} for the systems $\text{CaF}_2:\text{Pr}^{3+}$, $\text{KY}_3\text{F}_{10}:\text{Pr}^{3+}$, $\text{BaY}_2\text{F}_8:\text{Pr}^{3+}$ and $\text{LaZrF}_7:\text{Pr}^{3+}$

	$\text{CaF}_2:\text{Pr}^{3+}$	$\text{KY}_3\text{F}_{10}:\text{Pr}^{3+}$	$\text{BaY}_2\text{F}_8:\text{Pr}^{3+}$	$\text{LaZrF}_7:\text{Pr}^{3+}$
$F_2(ff)$	329.4	329.5	331.6	330.2
$F_4(ff)$	42.7	42.7	42.9	42.8
$F_6(ff)$	4.5	4.5	4.6	4.7
$\Delta_0(fd)$	53473	53288	53495	53098
$G_1(fd)$	322.3	251.4	266.7	285.5
$F_2(fd)$	219.0	170.5	178.7	190.2
$G_3(fd)$	28.6	20.9	23.0	24.5
$F_4(fd)$	16.7	12.3	13.3	14.2
$G_5(fd)$	4.5	3.2	3.6	3.8

The structures of the Pr^{3+} ion surrounded by the eight fluoride ligands are obtained from plane-wave calculation taking into account the periodicity of the crystal structures by means of the VASP codes.^{50,51} The utility of the periodical calculation to obtain the local geometry around the lanthanide ion, especially in the case of the doping sites, was discussed in a previous work³³ Our precedent work³³ dealing with the doping of Ce^{3+} into $\text{Cs}_2\text{NaYCl}_6$ showed that a selective cut of the crystal host implying further molecular DFT optimization of the geometry already gives good results suitable for the LFDFT calculation. However the definition of this selective cut seems strongly dependent to the problem we are facing, knowing that the bigger is the cluster's size we cut from the crystal structure, the smaller is the difference in geometry between the molecular DFT optimization and the periodical calculation. We use the periodical calculation exclusively for geometry concern since the corresponding code do not have the leverage for the further handling in the interest of the LFDFT modelling, such as generating the AOC orbitals. In the system $\text{CaF}_2:\text{Pr}^{3+}$, the eight fluoride ligands form around the Pr^{3+} ion a cubic arrangement of O_h symmetry. The Pr^{3+} -ligand bond lengths are all equivalents and match 2.361 Å in agreement with the Shannon radii⁶⁶ of Pr^{3+} and F^- ions in such an eight coordination. In $\text{KY}_3\text{F}_{10}:\text{Pr}^{3+}$, the Pr^{3+} ion replaces the trivalent Y^{3+} ion surrounded by eight fluoride ligands forming an arrangement with C_{4v} symmetry. In this case, the eight fluoride ligands can be separated into two distinct groups: F_1 and F_2 (cf. Fig. 1). Their positions in spherical coordinates can be generated by symmetry using the following values: $d_1 = 2.439$ Å, $\theta_1 = 50.76^\circ$, $\varphi_1 = 0.00^\circ$, and $d_2 = 2.278$ Å, $\theta_2 = 118.17^\circ$, $\varphi_2 = 45.00^\circ$, respectively. Less symmetrical coordination of the Pr^{3+} is calculated considering the system $\text{BaY}_2\text{F}_8:\text{Pr}^{3+}$. Like in the precedent system, the Pr^{3+} replaces the trivalent Y^{3+} ion and is surrounded by eight fluoride ligands forming an arrangement with C_2 symmetry. In this case the ligands, which are two by two identical by symmetry, are grouped into F_1 , F_2 , F_3 and F_4 (cf. Fig. 1). In spherical coordinates, their positions can be generated by symmetry from the following coordinates: $d_1 = 2.418$ Å, $\theta_1 = 146.79^\circ$, $\varphi_1 = -2.19^\circ$, $d_2 = 2.348$ Å, $\theta_2 = 81.23^\circ$, $\varphi_2 = 16.60^\circ$, $d_3 = 2.341$ Å, $\theta_3 = 35.15^\circ$, $\varphi_3 = 126.44^\circ$, and $d_4 = 2.312$ Å, $\theta_4 = 109.35^\circ$, $\varphi_4 = 89.13^\circ$, respectively. The calculation on the system $\text{LaZrF}_7:\text{Pr}^{3+}$ yields a coordination without any symmetry operation at all (C_1 point group). This last case represents qualitatively the situation of nowadays experimental synthesis of inorganic phosphors,⁶ where highly symmetrical objects are usually not acquired. In this case, we obtain eight different coordinates and the molecular structure can be taken as a distortion of the earlier C_2 one, where in spherical coordinates we calculate: bond lengths d of 2.469, 2.453, 2.417, 2.416, 2.404, 2.369, 2.306 and 2.262 Å; polar angles θ of 43.42, 109.00, 104.97, 108.36, 179.49, 108.83, 56.60 and 37.84°; and azimuthal angles φ of -177.56, -51.05, -141.89, -29.93, 5.83, 134.44, 80.42 and -44.87°, respectively for the eight fluoride ligands.

The four Pr^{3+} -doped phosphors, we consider as examples for the application of our model in the present work, are fascinating since they are real systems, which illustrate the descent in symmetry of the ligand field potential from the O_h to C_1 point groups removing step by step high order symmetry rotation axis,

while the total number of fluoride ligands remains constant (eight). For any step in this descent in symmetry, we expect an increase of the number of the B_q^k parameters involved in the ligand field potential (Table 2). The less symmetrical is the object, the more parameters exists. For example, in the O_h ligand field occurring in the $\text{CaF}_2:\text{Pr}^{3+}$ system, only three parameters prevail: $B_0^4(f)$, $B_0^6(f)$, and $B_0^4(d)$ as shown in Table 2. The $B_4^4(f)$, $B_4^6(f)$, and $B_2^4(d)$ can be defined in terms of the previous ones by multiplying them with constant factors: $\sqrt{5/14}$, $-\sqrt{7/2}$, and $\sqrt{5/14}$, respectively.

Table 2 Calculated Wybourne-normalized crystal field parameters B_q^k in cm^{-1} for the systems $\text{CaF}_2:\text{Pr}^{3+}$, $\text{KY}_3\text{F}_{10}:\text{Pr}^{3+}$, $\text{BaY}_2\text{F}_8:\text{Pr}^{3+}$ and $\text{LaZrF}_7:\text{Pr}^{3+}$

	$\text{CaF}_2:\text{Pr}^{3+}$	$\text{KY}_3\text{F}_{10}:\text{Pr}^{3+}$	$\text{BaY}_2\text{F}_8:\text{Pr}^{3+}$	$\text{LaZrF}_7:\text{Pr}^{3+}$
$B_0^2(f)$	0	-1322	386	350
$B_1^2(f)$	0	0	0	$928+i250$
$B_2^2(f)$	0	0	$-738+i159$	$468+i100$
$B_0^4(f)$	-2973	-2752	513	-970
$B_1^4(f)$	0	0	0	$-224+i221$
$B_2^4(f)$	0	0	$-45+i97$	$-480+i188$
$B_3^4(f)$	0	0	0	$146+i770$
$B_4^4(f)$	-1777	-2015	$-815+i2008$	$-629+i84$
$B_0^6(f)$	844	2095	-564	-485
$B_1^6(f)$	0	0	0	$-159+i200$
$B_2^6(f)$	0	0	$-1027+i1408$	$-257+i588$
$B_3^6(f)$	0	0	0	$-116+i1218$
$B_4^6(f)$	-1578	1339	$582+i233$	$533+i1015$
$B_5^6(f)$	0	0	0	$-458+i378$
$B_6^6(f)$	0	0	$153+i200$	$527+i433$
$B_0^2(d)$	0	5181	2635	-12265
$B_1^2(d)$	0	0	0	$-4492+i3329$
$B_2^2(d)$	0	0	$-13261+i7361$	$9540+i7925$
$B_0^4(d)$	-47973	-43650	7286	-9432
$B_1^4(d)$	0	0	0	$-3230+i12693$
$B_2^4(d)$	0	0	$5441+i12511$	$-3395+i2568$
$B_3^4(d)$	0	0	0	$-6292+i956$
$B_4^4(d)$	-28669	3507	$15905+i17573$	$-12769+i776$
$B_0^1(fd)$	0	4342	-1553	-4823
$B_1^1(fd)$	0	0	0	$1224+i1733$
$B_0^3(fd)$	0	-6003	-1227	-3700
$B_1^3(fd)$	0	0	0	$-1165+i3171$
$B_2^3(fd)$	0	587	$740+i1162$	$-2217+i6266$
$B_3^3(fd)$	0	0	0	$-260+i1778$
$B_0^5(fd)$	0	2202	-1671	-1862
$B_1^5(fd)$	0	0	0	$3787+i1260$
$B_2^5(fd)$	0	180	$-170+i373$	$2410+i4702$
$B_3^5(fd)$	0	0	0	$-5134+i2299$
$B_4^5(fd)$	0	-9554	$2197+i982$	$-7469+i3921$
$B_5^5(fd)$	0	0	0	$4718+i4462$

Cite this: DOI: 10.1039/c0xx00000x

www.rsc.org/xxxxxx

ARTICLE TYPE

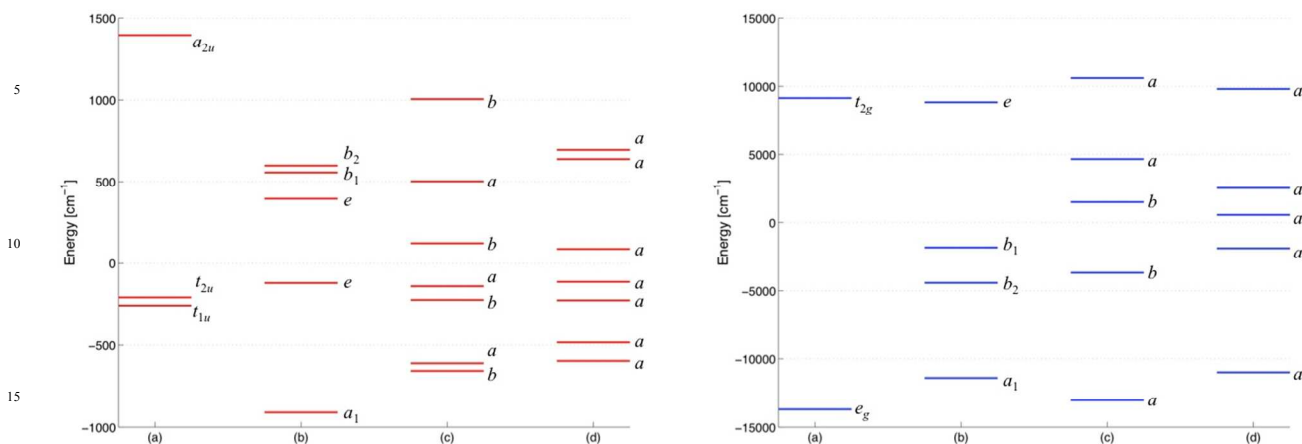


Fig. 3 Calculated energy splitting in cm^{-1} of the 4f (left hand side, in red) and the 5d (right hand side, in blue) orbitals of Pr^{3+} due to ligand field interaction in the system: $\text{CaF}_2:\text{Pr}^{3+}$ (a), $\text{KY}_3\text{F}_{10}:\text{Pr}^{3+}$ (b), $\text{BaY}_2\text{F}_8:\text{Pr}^{3+}$ (c) and $\text{LaZrF}_7:\text{Pr}^{3+}$ (d). The irreducible representations of the point group under which the local structures of the Pr^{3+} impurity belong, are also indicated.

The mixed $B_q^k(fd)$ parameters vanish in the ligand field potential due to the presence of the inversion centre symmetry operator forcing the 4f and 5d to possess opposite parity in O_h (cf. Fig. 3).

In the C_{4v} ligand field occurring in the $\text{KY}_3\text{F}_{10}:\text{Pr}^{3+}$ system, the 4f and 5d orbitals of Pr^{3+} split into a_1, b_1, b_2, e, e and a_1, b_1, b_2, e representative of the irreducible representations (irreps) of the C_{4v} point group, respectively (cf. Fig. 3).

A mixing of the 4f and 5d functions is expected when they belong to the same irreps as underlined by the appearance of the mixed $B_q^k(fd)$ parameters in the potential (Table 2). These parameters are smaller if compared to the $B_q^k(d)$ ones. One may perhaps neglect them in the potential because their small influence will be superseded by the large 4f-5d energy gap included in the $\Delta_0(fd)$ parameter. Nevertheless they will surely affect the line intensities of the calculated multiplet energy levels. The C_2 ligand field in $\text{BaY}_2\text{F}_8:\text{Pr}^{3+}$ yet differentiates two distinct irreps a and b in the potential (cf. Fig. 3). The B_q^k 's obtained for any odd q value are therefore zero because of the presence of the C_2 symmetry axis.

The C_1 ligand field met in $\text{LaZrF}_7:\text{Pr}^{3+}$ yields all the B_q^k parameters (Table 2). Here the ligand field matrix receives fully non-zero elements. This typically low symmetry coordination of Pr^{3+} ion requests special treatment since the B_q^k parameters become redundant. Indeed the values will change in terms of the orientation of the $(\text{PrF}_8)^{5-}$ structure, in line with the definition of Wybourne.⁴³ Therefore in the $\text{LaZrF}_7:\text{Pr}^{3+}$, a definition of an easy axis is not always trivial like in the systems $\text{CaF}_2:\text{Pr}^{3+}$, $\text{KY}_3\text{F}_{10}:\text{Pr}^{3+}$, $\text{BaY}_2\text{F}_8:\text{Pr}^{3+}$ where the S_4 , C_4 and C_2 axes, respectively, are set to the z -axis leading to the values in Table 2.

Any change of this orientation setting will transform the B_q^k parameters scheme from the quantities in Table 2 to different ones. Some B_q^k parameters may appear or disappear accordingly. The eigenvalue of the ligand field potential are graphically

represented in Fig. 3 considering the four systems under consideration. For clarity, the $B_0^0(f)$ and $B_0^0(d)$ are respectively taken as the origins of the energy for both one-electron perturbation of the 4f and 5d orbitals.

From Fig. 3, we realize that the energy splitting of the 4f orbitals of Pr^{3+} due to ligand field for all cases spans energy width of less than 2000 cm^{-1} . In the 5d orbitals, the energy splitting due to ligand field spans energy width of more than 20000 cm^{-1} , where the systems $\text{CaF}_2:\text{Pr}^{3+}$ and $\text{BaY}_2\text{F}_8:\text{Pr}^{3+}$ reach the strongest interaction, important in the rule of ligand field to showcase quantum cutting process in the optical behaviour. We use the formulation of the ligand field potential in terms of the Wybourne formalism for comparison purpose, especially if we want to pay attention to similar investigations mostly based on the description and data analysis of experimental findings.^{12,14,15} However it is not essential to parameterize the ligand field potential in terms of the Wybourne formalism, at least in the chemistry point of view because the Wybourne parameters are simply numbers (Table 2). Another way to illustrate the ligand field potential is obtained using the old concept of AOM following the work of Schäffer and Jørgensen⁴⁵ and Urland⁴⁶ for single-open-shell d and f electrons, respectively. In the case of the $\text{CaF}_2:\text{Pr}^{3+}$ system for instance, we obtain the AOM parameters (in cm^{-1}) $e_\sigma(f) = 562$, $e_\pi(f) = 202$, $e_\sigma(d) = 13426$ and $e_\pi(d) = 3644$, knowing that the O_h symmetry of the structure allows considering eight identical and equivalent fluoride ligands by symmetry (Fig. 1). The AOM formalism implies particularly chemical bonding since its parameters allow to determine the force of the ligands to make possible a σ or π bonding with the lanthanide ion. The AOM formalism is also used to be the method of choice relevant in the special case of low symmetry coordination, where it is known that the non-empirical calculation may sometimes fail to address.

The AOM parameters belong to one specific ligand, they are also comparable, transferable and offer chemical insight.

The final stage of the LDFFT method consists to build the ligand field theory Hamiltonian in a cumulative sum of the influence of the inter-electronic repulsion Hamiltonian, the ligand field potential and the spin-orbit coupling interaction parameterized with the values given in Table 1, Table 2 and the spin-orbit coupling constants, respectively. We obtain the spin-orbit coupling constants for the 4f and 5d orbitals from the approach of ZORA relativistic available in the ADF program package, where we calculate average spin-orbit coupling constants (in cm^{-1}) $\zeta_{4f} = 710$ and $\zeta_{5d} = 945$ considering the four systems (Fig. 1) for the 4f and 5d orbitals, respectively. These values are in the magnitude of known experimental values in some Pr^{3+} and fluoride ligands system. We calculate the multiplet energy levels corresponding to the $4f^2$ and $4f^15d^1$ electron configurations of Pr^{3+} by diagonalizing the ligand field theory Hamiltonian. The calculated multiplet energy levels are graphically represented in Fig. 4 for the four systems under consideration in the present work. They are also presented numerically in the electronically supplementary information (ESI). Different colours are used to emphasize the energy levels diagram originating from the $4f^2$ and the $4f^15d^1$ electron configurations of Pr^{3+} . The systems $\text{CaF}_2:\text{Pr}^{3+}$ and $\text{BaY}_2\text{F}_8:\text{Pr}^{3+}$ do not allow quantum cutting process considering the theoretical results, in line with the experimental observation already described for such systems.

The experimental work of Oskam et al. investigated different sites symmetry of the Pr impurity in CaF_2 . In case of the O_h site of the Pr impurity, Oskam et al. clearly deduced quantum cutting process in $\text{CaF}_2:\text{Pr}^{3+}$. Therefore in ref. 37, Fig. 7 is assigned to the excitation spectrum of $\text{CaF}_2:\text{Pr}^{3+}$ within cubic O_h symmetry of the Pr impurity, where they found two observed bands with an energy difference of 21000 cm^{-1} . In Fig. 3a (right hand side) we relate this ligand field splitting in a O_h site of the Pr impurity ($\text{CaF}_2:\text{Pr}^{3+}$) where we obtain 22844 cm^{-1} , in agreement with the observed energy by Oskam et al. Furthermore, we can also compare absolute energy values. In this context, Oskam et al. (see ref. 37 Fig. 7) deduced the position of the zero phonon line of the lower energy band at 223.3 nm (*i.e.* 44783 cm^{-1}). Fig. 4(a), this energy, corresponding to the lowest energy of the $4f^15d^1$ multiplets in $\text{CaF}_2:\text{Pr}^{3+}$ is calculated at 47484 cm^{-1} or 210.6 nm (see also ESI for $\text{CaF}_2:\text{Pr}^{3+}$). Moreover Oskam et al. expected the position of the $^1\text{S}_0$ level around 212 nm (*i.e.* 47170 cm^{-1}), we obtain 49315 cm^{-1} (202.8 nm) (Figure 4(a) and ESI). There is therefore a shift of the energy about less than 3000 cm^{-1} or 15 nm between the experimental and the theoretical results. This discrepancy between theoretical and experimental energies is due to the slight overestimation obtained for the DFT calculation, especially considering ligand field of the 4f and 5d of lanthanide.

It is noteworthy that we produce here non-empirical determination without any fitting to experimental measurements, enabling the calculation of quantum cutting process. This is because the highest energy level corresponding to the $4f^2$ electron configuration of Pr^{3+} (the $^1\text{S}_0$ spectroscopic term) is above the lowest energy levels of the $4f^15d^1$ (Fig 4). Therefore the expected cascade photon emission due to quantum cutting process will not be possible in such a situation. The optical behaviour of the systems still permits the red emission generally exhibited by the

$4f^2$ (^3P) \rightarrow $4f^2$ (^3H) transitions, *ergo* important in warm-white light phosphor candidates. In the systems $\text{KY}_3\text{F}_{10}:\text{Pr}^{3+}$ and $\text{LaZrF}_7:\text{Pr}^{3+}$ the highest energy levels $^1\text{S}_0$ from $4f^2$ are below the terms from $4f^15d^1$ leading to the possibility of quantum cutting by allowing the cascade photon emission. This cascade photon emission is experimentally characterized as the combination of the radiationless $4f^15d^1 \rightarrow 4f^2$ (^1S) transitions plus the typically blue emission in the $4f^2$ (^1S) \rightarrow $4f^2$ (^1I) transitions and the red emission (*vide supra*). The system $\text{LaZrF}_7:\text{Pr}^{3+}$ has been investigated before where the experimental evidence of the quantum cutting process was observed. In this case it was deduced that the difference in energy between the highest energy levels of the $4f^2$ configuration ($^1\text{S}_0$) and the lowest energy levels of $4f^15d^1$ was approximately around 2000 cm^{-1} , in the magnitude of the present calculation where we find an energy difference about 2638 cm^{-1} (Fig. 4(d) and ESI). The experimental accounts to the optical properties of $\text{KY}_3\text{F}_{10}:\text{Pr}^{3+}$ system suppose non quantum cutting. Of course this conclusion was not concisely precised in this way. Our non-empirical determination does not rigorously confirm the experimental observation. We attribute this discrepancy because of the ligand field interaction in the 5d orbitals, which is not strong enough to change the optical behaviour. The energy splitting of the 5d orbitals of Pr^{3+} in the $(\text{PrF}_8)^{5-}$ complex embedded in the KY_3F_{10} host enlarges if we consider a distortion of the structure following a vibronic coupling type between the degenerate triplet electronic ground state and vibrational modes in the computational details. This distorts the high symmetry C_{4v} structure into lower symmetry one, *e.g.* C_{2v} , if following a vibrational mode of b_1+b_2 symmetry. Of course if there is a distortion occurring in the system it might not be governed by only vibronic coupling type but the results of different aspects due to volume or size effect for instance (*i.e.* difference in Shannon radii) or site distortion by doping...

Experimentally, the geometry relative to the Pr^{3+} guest ion is ill-defined. It systematically follows the symmetry of the site (Wyckoff position) where the Pr^{3+} ion is placed. Theoretically in our periodical calculation, we let the doped site to relax in a totally symmetric displacement, avoiding thus any distortion of the geometry. This computational procedure does not correctly address the $\text{KY}_3\text{F}_{10}:\text{Pr}^{3+}$ problem. Taking into account, however, a possible orthorhombic distortion of the structure, $\text{KY}_3\text{F}_{10}:\text{Pr}^{3+}$ is not a quantum cutter as test made within the present model.

Conclusions

The best light source alternative after the banishment on incandescent light bulbs is certainly obtained within simple, inexpensive and energy efficient technology. LED lighting is undoubtedly the future of our domestic lighting. One usually judges the efficiency of the LED with respect to the optical manifestation of the inorganic phosphor which is used. This inorganic phosphor, currently based on the doping of lanthanide ion into stable crystal host lattices, gives suitable visible light emission necessary to the generation of warm-white light. As written, nowadays LED lighting is not only understood by the analysis of the $f \rightarrow d$ dipole allowed transitions in the lanthanide system, but also limited to the bluish cold-white light emission of the Ce^{3+} -based LED bulbs.

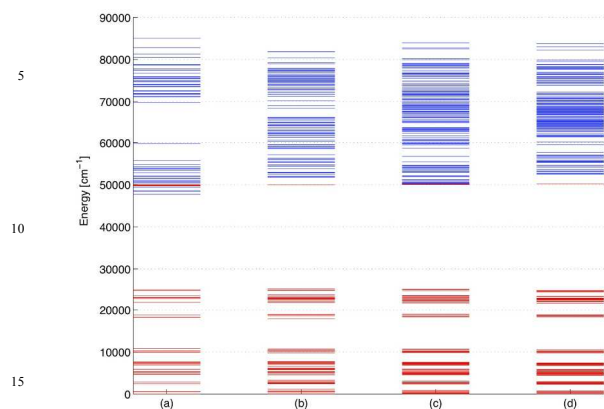


Fig. 4 Calculated multiplet energy levels in cm^{-1} obtained for Pr^{3+} doped into CaF_2 (a), KY_3F_{10} (b), BaY_2F_8 (c) and LaZrF_7 (d). The energies of the terms originating from the ground electron configuration of Pr^{3+} ($4f^3$) are in red, those for the excited configuration ($4f^25d^1$) are in blue. (cf. ESI Numerical data)

Lanthanides like Pr^{3+} and Eu^{2+} are already undeniably conceived as activators of warm-white light source, ergo solution for the improvement of LED technology. However the development and engineering of such a technology are still complex relying on the perfect understanding of the $4f^n \rightarrow 4f^{n-1}5d^1$ transitions.

In this paper, we present a fully non-empirical determination of the $4f^n \rightarrow 4f^{n-1}5d^1$ transitions in lanthanide doped into crystal host lattices, in order to understand the optical manifestation of new generation of LED phosphors. The method is particularly applied to understand the phenomenology of quantum cutting process in the optical properties of Pr^{3+} -doped into various fluoride lattices. We address the problem in a transparent and clear approach taking advantages of the facilities given by ligand field theory in conjunction with modern quantum chemistry tools. The LDFDT, once again presented in the framework of the DFT2013 meeting, shows here a useful tool providing reliable and accurate results in the perspective to design the optical properties of materials by Density Functional Theory.

Acknowledgement

This work is supported by the Swiss National Science Foundation (SNF) and the Swiss State Secretariat for Innovation and Research. Support from the UEFISC-CDI Romania research grant PCE 14/2013 is also acknowledged.

Notes and references

- ^a Department of Chemistry of the University of Fribourg, Chemin du Musée 9, 1700 Fribourg, Switzerland. Fax: +41 26 300 9738; Tel: +41 26 300 8700; E-mail: claudd.daul@unifr.ch
- ^b Institut für Anorganische Chemie der Universität Hannover, Callinstr. 9 D-30167 Hannover, Germany. E-mail: wurland@arcor.de
- ^c Institut of Physical Chemistry, Splaiul Independentei 202, Bucharest 060021, Romania. E-mail: cfanica@yahoo.com

† Electronic supplementary information (ESI) available: Numerical data for the multiplet energy levels obtained for the $4f^2$ and $4f^25d^1$ electron configuration of Pr^{3+} doped CaF_2 , KY_3F_{10} , BaY_2F_8 and LaZrF_7 .

- 1 S. Nakamura and G. Fasol, *The Blue Laser Diode*. Springer, Berlin, 1997.
- 2 T. Jüstel, H. Nikol and C. Ronda, *Angew. Chem. Int. Ed.* 1998, **37**, 3084.
- 3 C. Feldmann, T. Jüstel, C.R. Ronda and P.J. Schmidt, *Adv. Funct. Mater.* 2003, **13**, 511.
- 4 H.S. Jang, W.B. Im, D.C. Lee, D.Y. Jeon and S.S. Kim, *J. Lumin.* 2007, **126**, 371.
- 5 M. Roushan, X. Zhang, and J. Li, *Angew. Chem. Int. Ed.* 2012, **51**, 436.
- 6 Q.Y. Zhang and X.Y. Huang, *Progress in Materials Science*, 2010, **55**, 353.
- 7 D. Tu, Y. Liu, H. Zhu, R. Li, L. Liu and X. Chen, *Angew. Chem. Int. Ed.*, 2013, **52**, 1128.
- 8 J. S. Griffith, *The Theory of Transition Metal Ions*, Cambridge University Press, Cambridge 1961.
- 9 B. N. Figgis and M. A. Hitchman, *Ligand Field Theory and its Applications*, Wiley-VCH, New York 2000.
- 10 A. J. Slakeld, M. F. Reid, J. P. R. Wells, G. Sanchez-Sanz, L. Seijo and Z. Barandiaran, *J. Phys.: Condens. Matter*, 2013, **25**, 415504.
- 11 M. Laroche, J.-L. Doualan, S. Girard, J. Margerie and R. Moncorgé, *J. Opt. Soc. Am. B*, 2000, **17**, 1291.
- 12 C.-K. Duan, P.A. Tanner, V. Makhov and N. Khaidukov, *J. Phys. Chem. A*, 2011, **115**, 8870.
- 13 M.D. Faucher and O.K. Moune, *Phys. Rev. A*, 1997, **55**, 4150.
- 14 M.F. Reid, L. van Pieteron, R.T. Wegh and A. Meijerink, *Phys. Rev. B*, 2000, **62**, 14744.
- 15 P.S. Peijzel, P. Vergeer, A. Meijerink, M.F. Reid, L.A. Boatner and G.W. Burdick, *Phys. Rev. B*, 2005, **71**, 045116.
- 16 L. Hu, M.F. Reid, C.-K. Duan, S. Xia and M.Yin, *J. Phys.: Condens. Matter*, 2011, **23**, 045501.
- 17 C.-K. Duan and P. A. Tanner, *J. Phys. Chem. A*, 2010, **114**, 6055.
- 18 M.F. Reid, C.-K. Duan and H. Zhou, *Journal of Alloys and Compounds*, 2009, **488**, 591.
- 19 M. Atanasov, C. Daul and C. Rauzy, *Struct. Bond.*, 2004, **106**, 97.
- 20 M. Atanasov and C. Daul, *Chimia*, 2005, **59**, 504.
- 21 A. Borel and C. Daul, *J. Molecular Struct.: THEOCHEM*, 2006, **762**, 93.
- 22 M. Atanasov, E. J. Baerends, P. Baettig, R. Bruyndonckx, C. Daul and C. Rauzy, *Chem. Phys. Lett.*, 2004, **399**, 433.
- 23 M. Atanasov and C. Daul, *Comptes Rendus Chimie*, 2005, **8**, 1421.
- 24 M. Atanasov, C. Daul, H. U. Güdel, T. A. Wesolowski and M. Zbiri, *Inorg. Chem.*, 2005, **44**, 2954.
- 25 L. Petit, A. Borel, C. Daul, P. Maldivi and C. Adamo, *Inorg. Chem.*, 2006, **45**, 7382.
- 26 A. Borel, L. Helm and C. Daul, *Chem. Phys. Lett.*, 2004, **383**, 584.
- 27 F. Senn, L. Helm, A. Borel and C. A. Daul, *Comptes Rendus Chimie*, 2012, **15**, 250.
- 28 P. Garcia-Fernandez, F. Senn, C. A. Daul, J. A. Aramburu, M.T. Barriuso and M. Moreno, *Phys. Chem. Chem. Phys.* 2009, **11**, 7545.
- 29 F. Senn and C. A. Daul, *J. Mol. Struct.: THEOCHEM*, 2010, **954**, 105.
- 30 C. A. Daul, *J. Phys.: Conf. Ser.*, 2013, **428**, 012023.
- 31 C. Daul, *Int. J. Quantum Chem.*, 1994, **52**, 867.
- 32 H. Ramanantoanina, W. Urland, F. Cimpoesu and C. Daul, *Phys. Chem. Chem. Phys.*, 2013, **15**, 13902.
- 33 H. Ramanantoanina, W. Urland, A. Garcia-Fuente, F. Cimpoesu, C. Daul, *Chem. Phys. Lett.* 2013, **588**, 260.
- 34 J. L. Sommerdijk, A. Brill and A. W. de Jager, *J. Lumin.*, 1974, **8**, 341.
- 35 W. W. Piper, J. A. DeLuca and F. S. Ham, *J. Lumin.*, 1974, **8**, 344.
- 36 R. Stowasser and R. Hoffmann, *J. Am. Chem. Soc.* 1999, **121**, 3414.
- 37 K.D. Oskam, A.J. Houtepen and A. Meijerink, *J. Lumin.*, 2002, **97**, 107.
- 38 M. Laroche, A. Braud, S. Girard, J.-L. Doualan, R. Moncorgé, M. Thuau and L. D. Merkle, *J. Opt. Soc. Am. B*, 1999, **16**, 2269.

- 39 A.A. Pushkar, T.V. Uvarova and V.N. Molchanov, *Quantum Electronics*, 2008, **38**, 333.
- 40 E. van der Kolk, P. Dorenbos and C.W.E. van Eijk, *Optics Communications*, 2001, **197**, 317.
- 5 41 R.D. Cowan, *The theory of atomic structure and spectra*. University of California Press, Berkeley, 1997.
- 42 M. Gerloch, *Magnetism and ligand field analysis*, Cambridge University Press, Cambridge, 1983.
- 43 S. Hüfner, *Optical spectra of transparent rare earth compounds*, Academic Press, New York 1978.
- 10 44 J. C. Morrison and K. Rajnak, *Phys. Rev. A*, 1971, **4**, 536.
- 45 C. Schäffer and C. Jørgensen, *Mol. Phys.*, 1965, **9**, 401.
- 46 W. Urland, *Chem. Phys.*, 1976, **14**, 393.
- 47 E. J. Baerends, T. Ziegler, J. Autschbach, D. Bashford, A. Bércecs, F. M. Bickelhaupt, C. Bo, P. M. Boerrigter, L. Cavallo, D. P. Chong, L. Deng, R. M. Dickson, D. E. Ellis, M. van Faassen, L. Fan, T. H. Fischer, C. Fonseca Guerra, A. Ghysels, A. Giammona, S. J. A. van Gisbergen, A. W. Götz, J. A. Groeneveld, O. V. Gritsenko, M. Grüning, S. Gusarov, F. E. Harris, P. van den Hoek, C. R. Jacob, H. Jacobsen, L. Jensen, J. W. Kaminski, G. van Kessel, F. Kootstra, A. Kovalenko, M. V. Krykunov, E. van Lenthe, D. A. McCormack, A. Michalak, M. Mitoraj, J. Neugebauer, V. P. Nicu, L. Noodleman, V. P. Osinga, S. Patchkovskii, P. H. T. Philipsen, D. Post, C. C. Pye, W. Ravenek, J. I. Rodríguez, P. Ros, P. R. T. Schipper, G. Schreckenbach, J. S. Seldenthuis, M. Seth, J. G. Snijders, M. Solà, M. Swart, D. Swerhone, G. te Velde, P. Vernooijs, L. Versluis, L. Visscher, O. Visser, F. Wang, T. A. Wesolowski, E. M. van Wezenbeek, G. Wiesenekker, S. K. Wolff, T. K. Woo and A. L. Yakovlev, ADF2010.01, available at <http://www.scm.com>.
- 30 48 C. F. Guerra, J. G. Sniders, G. te Velde and E. J. Baerends, *Theor. Chem. Acc.*, **1998**, **99**, 391.
- 49 G. te Velde, F. M. Bickelhaupt, S. J. A. van Gisbergen, C. F. Guerra, E. J. Baerends, J. G. Sniders and T. Ziegler, *J. Comput. Chem.*, 2001, **22**, 932.
- 35 50 G. Kresse and J. Hafner, *Phys. Rev. B*, 1993, **47**, 558.
- 51 G. Kresse and J. Furthmüller, *Phys. Rev. B*, 1996, **54**, 11169.
- 52 J. P. Perdew, K. Burke and M. Ernzerhof, *Phys. Rev. Lett.*, 1996, **77**, 3865.
- 53 P. E. Blöchl, *Phys. Rev. B*, 1994, **50**, 17953.
- 54 G. Kresse and D. Joubert, *Phys. Rev. B*, 1999, **59**, 1758.
- 55 V. Weber, C. Daul and R. Baltensperger, *Computer Physics Communications*, 2004, **163**, 133.
- 56 F. Harris, *Computation methods of quantum chemistry Part I*, Department of Physics University of Utah, Salt Lake City, 1973.
- 45 57 G. W. Burdick and M.F. Reid, *Handbook on the Physics and Chemistry of Rare Earths*, 2007, **37**, 61.
- 58 M. Ferbinteanu, T. Kajiwara, F. Cimpoesu, K. Katagiri and M. Yamashita, *Polyhedron*, 2007, **26**, 2069.
- 59 P. Comba and M. Kerscher, *Coord. Chem. Rev.*, 2009, **253**, 564.
- 50 60 J. Paulovic, F. Cimpoesu, M. Ferbinteanu and K. Hirao, *J. Am. Chem. Soc.*, 2004, **126**, 3321.
- 61 M. Ferbinteanu, F. Cimpoesu, M. A. Girtu, C. Enachescu and S. Tanase, *Inorg. Chem.*, 2012, **51**, 40.
- 62 F. Cimpoesu, S. Dahan, S. Ladeira, M. Ferbinteanu and J.-P. Costes, *Inorg. Chem.*, 2012, **51**, 11279.
- 55 63 M. Ferbinteanu, T. Kajiwara, K.-Y. Choi, H. Nojiri, A. Nakamoto, N. Kojima, F. Cimpoesu, Y. Fujimura, S. Takaishi and M. Yamashita, *J. Am. Chem. Soc.*, 2006, **128**, 9008.
- 64 S. Tanase, M. Ferbinteanu and F. Cimpoesu, *Inorg. Chem.*, 2011, **50**, 9678.
- 60 65 M. Atanasov, P. Comba, C. A. Daul and F. Neese, *Models, Mysteries and Magic of Molecules*, Springer Berila, 2008, 411.
- 66 R. D. Shannon, *Acta Cryst. A*, 1976, **32**, 751.
- 67 J. Sugar, *J. Opt. Soc. Am.* 1965, **55**, 1058.
- 65 68 S. Khiari, M. Velasquez, R. Moncorgé, J.L. Doualan, P. Camy, A. Ferrier and M. Daif, *J. Alloys and Compounds*, 2008, **451**, 128.
- 69 H.A. Jahn and E. Teller, *Proc. R. Soc. A*, 1937, **161**, 220.
- 70 M. Atanasov, C. Daul and P.L.W. Tregenna-Piggott, *Vibronic Interactions and the Jahn-Teller Effect*, Progress in Theoretical Chemistry and Physics B23, Springer, Berlin, 2012.
- 70 71 D. Tu, Y. Liu, H. Zhu, R. Li and X. Chen, *Angew. Chem. Int. Ed.*, 2013, **52**, 1128.