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Machine Learning-Assisted Optimization of Cu-Based HTLs for Lead-Free Sr_3PBr_3 Perovskite Solar Cells Achieving Over 30% Efficiency via SCAPS-1D Simulation

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Abstract

The pursuit of efficient and stable lead-free perovskite solar cells (PSCs) is critical for sustainable photovoltaic technologies. In this work, we systematically investigated Sr_3PBr_3-based PSCs incorporating five different copper-based hole transport layers (HTLs)—Cu_2O, CuI, $CuSbS_2$, CuSCN, and Cu_2BaSnS_4 (CBTS) using SCAPS-1D simulations. The device configuration FTO/SnS_2/Sr_3PBr_3/HTL/Au was optimized to evaluate the impact of HTL selection, absorber thickness, doping concentration, defect density, series resistance, and temperature on photovoltaic performance. The results demonstrate that HTL choice strongly governs charge extraction, interfacial recombination, and stability. Among the candidates, CBTS exhibited the highest efficiency, achieving a power conversion efficiency (*PCE*) of 30.78% with an open-circuit voltage (V_{OC}) of 1.32 V, short-circuit current density (J_{SC}) of 26.82 mA cm^{-2}, and fill factor (*FF*) of 87.05%. Machine learning (ML) models trained on simulation datasets provided predictive accuracies above 99.6% and, through SHAP (SHapley Additive exPlanations) analysis, revealed that acceptor density and defect density are the most influential parameters controlling device performance. This combined simulation–ML framework establishes CBTS as a highly promising non-toxic HTL and provides actionable insights for the design of stable, high-efficiency lead-free PSCs.

Keywords: Perovskite; Sr_3PBr_3; CBTS (Cu_2BaSnS_4); SCAPS-1D; Hole transport layers; Solar cell

1. Introduction

Driven by rapid technological advancements and population growth, global energy demand continues to rise, yet its reliance on fossil fuels to meet these needs poses significant environmental challenges [1]. To tackle this issue, sustainable energy alternatives, such as solar energy conversion through photovoltaic (PV) technology, are being explored[2–4]. In recent decades, there have been major advancements in the development of solar cells, including those made from GaAs, Si, CdTe, and CIGS[5–8]. Recent reports reveal that GaAs has obtained record power conversion efficiency (*PCE*) of up to 27.6%. However, the epitaxial growth process needed for GaAs layers significantly increases the cost of these solar cells [5]. In contrast,

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approximately 90% of the photovoltaic market depends on Si-based cells[6]. However, Si solar cells have several disadvantages, including being an indirect semiconductor [9] and having a low absorption coefficient. This necessitates a thick absorber layer, which increases costs due to the requirement for a substantial amount of high-quality single-crystal silicon[10,11]. It is noteworthy that materials with direct bandgaps, such as CdTe and CIGS, are commonly employed in commercial solar cell technologies [12]. However, these materials encounter difficulties because of the limited availability of elements in the Earth. Despite the reports on $Cu_2ZnSn(S,Se)_4$ (CZTSSe), which exists more extensively in the Earth's surface layer, the achieved *PCE* remains significantly lower at 14.9% compared to other established materials [13]. Conversely, the rise of organic-inorganic lead halide perovskites ($MAPbX_3$, where X = I, Br, Cl) has sparked significant attention to the advancement of photovoltaic materials[14]. Perovskite materials have demonstrated remarkable advancement, with their *PCE* enhancing dramatically from 3.8% in 2009 to an remarkable value of 26.1% by 2023. This places them as extremely promising contenders in the photovoltaic domain. Despite their outstanding performance, there are concerns regarding the poisonousness of lead and the stability of halides, which present hurdles to their durability over time and viability in the market [15–17].

In response to these concerns, scientists are enthusiastically investigating lead-free substitutes with the goal of addressing environmental issues and laying the groundwork for the next era of perovskite-based technologies[18]. More precisely, perovskites based on tin (Sn) and germanium (Ge), for example $MASnI_3$, $CsSnI_3$, $FASnI_3$, $CsGeI_3$, $MAGeI_3$, and $FAGeI_3$, have emerged as promising substitutes attributed to their lack of toxicity, narrow bandgaps, and superior carrier mobility[19]. Despite their promise, issues related to stability remain. For example, Sn^{2+} and Ge^{2+} ions are susceptible to oxidation, leading to their conversion into Sn^{4+} and Ge^{4+} when exposed to the surrounding atmosphere. Furthermore, the scarcity of Ge in the Earth's crust and its high cost further limit their practical application [20,21]. Indeed, scientists have recognized various alternative strategies. One approach entails the application of trivalent cations like Bi^{3+} and Sb^{3+} in lieu of Pb^{2+} ions. However, the achieved power conversion efficiency (*PCE*) with these replacements notably lags behind that of lead halide perovskites. This disparity could stem from their intrinsically low-dimensional structure, resulting in unfavorable optoelectronic properties [22].

Recently, researchers have increasingly employed SCAPS-1D simulations to gain deeper insights into the physical behavior of PSCs [23]. This approach has helped them fine-tune various

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components of the devices, leading to more efficient designs[24]. Moreover, machine learning has proven to be an effective approach for identifying novel methods to optimize the design and efficiency of PSCs [25]. Studies have shown that machine learning is an effective tool for predicting both the material characteristics and efficiency outcomes of PSCs.

To address these obstacles, researchers have introduced halide double perovskites (HDPs), featuring a structure resembling $A_2B^+B^{3+}X_6$, where A and B^+/B^{3+} represents cations and X indicates anions [26]. This advancement entails replacing Pb^{2+} with environmentally friendly Bi^{3+} while keeping a three-dimensional configuration [27]. This advanced approach involves substituting two Pb^{2+} ions with one B^+ ion and one B^{3+} ion. Noteworthy samples of such HDPs include $(CH_3NH_3)_2KBiCl_6$, $(CH_3NH_3)_2AgBiBr_6$, $Cs_2AgBiBr_6$, and $Cs_2NaBiCl_6$, which have recently attracted significant interest as highly favorable materials [28]. Despite their potential, HDPs encounter limitations such as weak charge carrier transport capabilities, high effective masses of charge carriers, and a high bandgap exceeding 2 eV [29]. Hence, research endeavors have concentrated on investigating antimony-based perovskite materials, including $Cs_2SbAgCl_6$, Cs_2SbAgI_6, $Cs_2SbCuCl_6$, $Cs_4SbCuCl1_2$, and $Cs_2SbAgBr_6$. These materials are acknowledged for their capacity to endure high temperatures and promote effective carrier mobility. Nonetheless, their prospects are constrained by wide bandgaps, elevated binding energies, and augmented carrier effective mass [30,31]. Consequently, there is an urgent requirement to identify novel materials that can preserve the remarkable characteristics of lead halide perovskites while mitigating the mentioned drawbacks.

As Sr_3PBr_3 is an emerging perovskite material, its crystal structure formation is influenced by the chemical interactions between Sr^{2+} cations and PBr_3 units, which stabilize the lattice through ionic bonding and maintain charge neutrality. The perovskite framework arises from the corner-sharing octahedral arrangement of $SrBr_6$ units, with P^{3+} ions occupying the interstitial sites, promoting structural stability. Among the various perovskite compositions explored, Sr_3PBr_3 has recently appeared as a promising absorber material due to its favorable optoelectronic properties and environmental safety[32,33]. This ternary halide features a stable orthorhombic crystal structure derived from alkaline earth and halide elements, which offers intrinsic thermal and structural stability under ambient conditions[32,33]. With a direct bandgap of 1.528 eV, Sr_3PBr_3 falls within the optimal range for single-junction solar cells according to the Shockley–Queisser limit, making it highly suitable for efficient solar energy conversion[32,33]. Furthermore, it exhibits a

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strong absorption coefficient across the visible spectrum and a low toxicity profile, in contrast to conventional Pb-based perovskites. Unlike Sn- and Ge-based alternatives, Sr_3PBr_3 is less susceptible to oxidation-induced instability, as it lacks divalent metal cations prone to environmental degradation (e.g., $Sn^{2+} \rightarrow Sn^{4+}$). Preliminary theoretical and experimental studies have shown that Sr_3PBr_3 can achieve long carrier diffusion lengths and low recombination rates, indicating strong potential for photovoltaic applications[34].

One critical component in PSCs that significantly influences their performance is the hole transport layer (HTL). The HTL efficiently extracts and transfers holes from the perovskite absorber to the anode, reducing recombination losses and boosting device efficiency. Copper-based HTLs, such as Cu_2O, CuI, $CuSbS_2$, CuSCN, and CBTS, have shown promise due to their favorable energy levels, high absorption coefficient, non-toxicity, good hole mobility, and chemical stability[35,36].

The device structure under investigation in this study is FTO/SnS_2/Sr_3PBr_3/Cu-based HTLs/Au. Employing fluorine-doped tin oxide (FTO) as the front electrode alongside SnS_2 as the electron transport layer (ETL) ensures good conductivity and electron extraction. The Sr_3PBr_3 serves as the active layer, with the copper-based HTLs facilitating hole transport to the gold (Au) back contact. The choice of multiple HTLs allows for a comparative analysis of their impact on device performance.

To optimize the device architecture and predict performance metrics, we employed SCAPS-1D simulation software[37,38]. SCAPS-1D is a well-established tool for modeling and simulating the electrical characteristics of PSCs, revealing the role of multiple parameters in determining device efficiency. By systematically varying the absorber's thickness, defect density, and acceptor concentration, along with adjustments to the ETL and HTL thicknesses and operating temperature, this study seeks to determine the optimal configuration that maximizes the *PCE* of Sr_3PBr_3-based PSCs.

In addition to these material considerations, the performance of any single-junction photovoltaic device is fundamentally constrained by the Shockley-Queisser (SQ) limit, which defines the maximum theoretical values of J_{SC}, V_{OC}, FF, and PCE as a function of the absorber bandgap under AM1.5 illumination. According to the theoretical framework outlined by (Morales-acevedo 2023), an absorber with a bandgap of 1.528 eV, such as Sr_3PBr_3 possesses an upper-limit V_{OC} of approximately 1.30–1.40 V and a maximum efficiency near 30%[39]. Therefore, in

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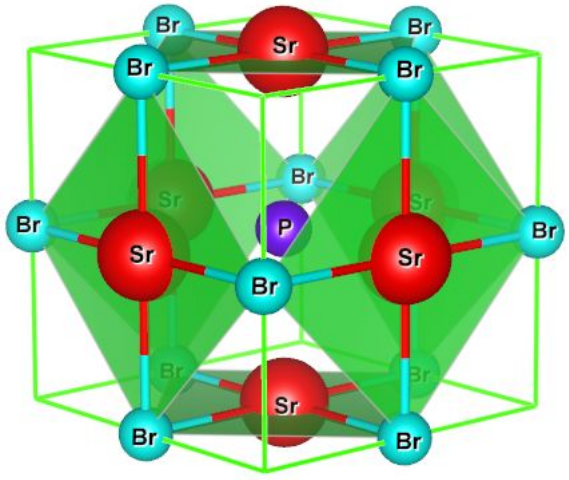
this work, all simulated device parameters were revised and analyzed to ensure consistency with these physical constraints, preventing overestimation of photovoltaic performance. This alignment with the SQ-limit framework guarantees that the simulation results remain physically meaningful and theoretically justified.

2. Device Configuration and Simulation

2.1. Simulated device configurations

In our research, we evaluated the performance of novel Sr_3PBr_3-based PSCs with five different inorganic Cu-Based HTLs-namely CBTS, Cu_2O, CuI, $CuSbS_2$, and CuSCN utilizing SCAPS-1D simulation. The configuration mimics a superstrate setup of FTO/SnS_2/Sr_3PBr_3/HTL/Au, as illustrated in Figure 1. Initial parameters for individual layers were obtained from existing literature and are detailed in Table 1. Throughout the simulations, the thermal velocities of holes and electrons were fixed at 10^7 cm/s for every layer, with a flat-band condition applied to the front contact. The simulations were performed at 300 K under AM 1.5G spectral irradiance. To more accurately replicate real operating conditions, neutral defects were incorporated at the ETL/ Sr_3PBr_3 and Sr_3PBr_3/HTL interfaces, using the parameters specified in Table 2.



(a)

(b)

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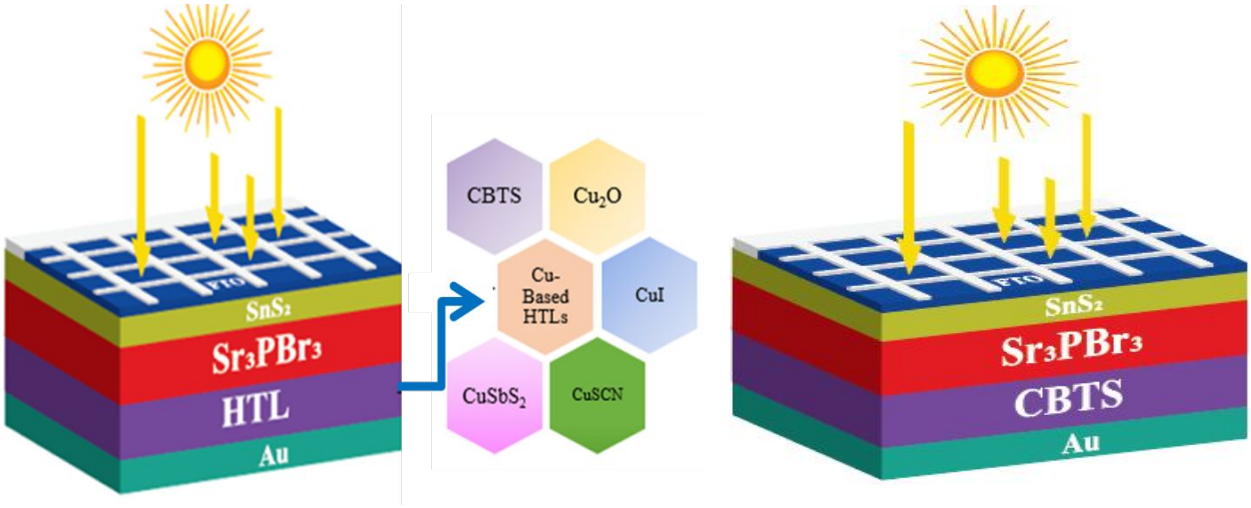

(c) (d)

Figure 1. (a) Crystal structure of the novel Sr_3PBr_3 and schematic of Sr_3PBr_3-based solar cells, (b) without HTL, (c) with various inorganic Cu-based HTLs, and (d) with the optimized HTL

Table 1. Input parameters employed in the simulation of the proposed PSCs [40–47].

Parameters	**Terms**	**FTO**	**SnS_2**	**Sr_3PBr_3** [32,33]	**Cu_2O**	**CBTS**	**CuI**	**CuSCN**	**$CuSbS_2$**
t(μm)	Thickness	0.05	0.05	1	0.05	0.05	0.05	0.05	0.05
E_g (eV)	Bandgap	3.6	2.42	1.528	2.17	1.90	3.10	3.40	1.580
χ (eV)	Electron affinity	4.5	4.24	4.160	3.2	3.60	2.1	1.9	4.2
ε_r	Relative Permittivity	10	10	5.280	7.11	5.4	6.5	10	14.6
Nc (cm^{-3})	Eff. Dos at CB	2×10^{18}	2×10^{18}	1.185×10^{19}	2.02×10^{17}	2.2×10^{18}	2.8×10^{19}	2.2×10^{18}	2×10^{18}
Nv (cm^{-3})	Eff. Dos at VB	1.8×10^{19}	1.8×10^{18}	1.660×10^{19}	1.1×10^{19}	1.8×10^{19}	1×10^{19}	1.8×10^{18}	1×10^{19}
μ_n(cm^2/(Vs))	Electrons Mobility	100	50	25	200	300	100	100	49

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$\mu_p(cm^2/(Vs))$	Holes Mobility	20	50	20	80	100	43.9	25	49
$N_A\ (cm^{-3})$	Dopant Density(Acceptor)	0	0	1×10^{17}	1×10^{18}	1×10^{18}	1×10^{18}	1×10^{18}	1×10^{18}
$N_D\ (cm^{-3})$	Dopant Density(Donor)	1×10^{18}	1×10^{17}	0	0	0	0	0	0
$N_t\ (cm^{-3})$	Defect Density	1×10^{14}	1×10^{12}	1×10^{12}	1×10^{15}	1×10^{15}	1×10^{15}	1×10^{15}	1×10^{15}

Table 2. Input parameters at the ETL/Sr_3PBr_3 and Sr_3PBr_3/ HTL interface.

Interface	Defect type	Capture cross section: electrons/ holes (cm^2)	Defect position above the highest EV (eV)	Energetic distribution	Reference for defect energy level	Total density (cm^{-2}) (integrated overall energies)
HTL/Sr_3PBr_3	Neutral	1×10^{19}	0.06	Single	Above the VB maximum	1.0×10^{10}
Sr_3PBr_3/ETL	Neutral	1×10^{19}	0.06	Single	Above the VB maximum	1.0×10^{10}

2.2. Mathematical Modeling

SCAPS-1D is an advanced simulation tool specifically designed to analyze the performance of solar cells, and it is widely utilized for evaluating various types of photovoltaic devices [48,49]. It offers the photovoltaic research community comprehensive insights into key output parameters, enabling a deeper understanding of device behavior. This capability is realized by numerically solving three fundamental equations integrated into the SCAPS-1D framework: the Poisson equation (1), the continuity equations (2) and (3), and the charge transport equations (4) and (5), which collectively govern charge carrier dynamics. These equations are expressed as follows [50,51],

$$\frac{\partial^2 \varphi}{\partial x^2} = \frac{q}{\varepsilon}\left(n(x) - p(x) + N_A(x) - N_D(x) - p_t(x) - N_t(x)\right) \quad (1)$$

Where, q denotes the elementary charge, ε represents the dielectric constant, p indicates the hole concentration, and n corresponds to the electron concentration. Furthermore, N_D refers to the

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doping concentration of donor-type impurities, while N_A designates the doping concentration of acceptor-type impurities.

$$\frac{\partial \mathrm{n}}{\partial \mathrm{t}} = \frac{1}{q}\frac{\partial J_n}{\partial \mathrm{x}} + (G_n - R_n) \quad (2)$$

$$\frac{\partial \mathrm{p}}{\partial \mathrm{t}} = -\frac{1}{q}\frac{\partial \mathrm{p}}{\partial \mathrm{x}} + (G_p - R_p) \quad (3)$$

Here, G_n denotes the electron generation rate, while G□ represents the hole generation rate. Similarly, R□ corresponds to the electron recombination rate, and R□ signifies the hole recombination rate. The symbols J□ and J□ indicate the current densities of holes and electrons, respectively [52].

$$J_n = qD_n\frac{\partial n}{\partial x} - q\mu_n n\frac{\partial \varphi}{\partial x} \quad (4)$$

$$J_p = qD_p\frac{\partial p}{\partial x} - q\mu_p p\frac{\partial \varphi}{\partial x} \quad (5)$$

Where, D□ represents the electron diffusion coefficient, while D□ corresponds to the hole diffusion coefficient. Similarly, μ□ denotes the electron mobility, and μ□ signifies the hole mobility.

3. Results and Discussion

3.1. Initial Solar Cell Performance

The original solar cells are constructed using the structure FTO/SnS_2/Sr_3PBr_3/HTL/Au, as depicted in Figure 1. Where, the inorganic HTLs like CBTS, Cu_2O, CuI, $CuSbS_2$, and CuSCN have used. The photovoltaic parameters of the new Sr_3PBr_3-based PSCs are presented in Table 3. Specifically, the power conversion efficiencies (*PCE*s) are 30.78%, 28.59%, 23.27%, 26.09%, and 27.92% for HTLs based on CBTS, Cu_2O, CuI, $CuSbS_2$, respectively and 20.80% without HTL. Table 3 shows the maximum parameters achieved under optimal temperature conditions.

Table 3. Obtained photovoltaic parameters of Sr_3PBr_3 solar cells.

Structure of PSCs	V_{OC} **(V)**	***FF*** **(%)**	J_{SC} **(mA/cm²)**	***PCE*** **(%)**
Without HTL				
FTO/SnS_2/Sr_3PBr_3/Au	1.052	86.13	22.971	20.80

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With HTL				
FTO/SnS_2/Sr_3PBr_3/CuI/Au	1.101	84.32	25.050	23.27
FTO/SnS_2/Sr_3PBr_3/ $CuSbS_2$/Au	1.111	88.82	26.429	26.09
FTO/SnS_2/Sr_3PBr_3/CuSCN/Au	1.237	84.31	26.762	27.92
FTO/SnS_2/Sr_3PBr_3/Cu_2O/Au	1.255	84.96	26.808	28.59
FTO/SnS_2/Sr_3PBr_3/CBTS/Au	1.317	87.05	26.839	30.78

3.2. Bandgap configuration of a Sr_3PBr_3-based PV apparatus

The alignment of energy band influences the current flow in a heterostructure. In our research, we examined various HTLs to determine which one provides the best *PCE*. Specifically, we utilized HTLs such as Cu_2O (Figure 2(c)), CuI (Figure 2(d)), $CuSbS_2$ (Figure 2(e)). Figure 2(a) displays the refined energy band diagram without an HTL, whereas Figure 2(b) incorporates the HTL (CBTS), which provides a superior PCE compared to the other HTLs. The results from the SCAPS simulation provide insights into the bandgap and thickness of each layer, highlighting the band bending at the interface of the CBTS and Sr_3PBr_3 layers providing more favorable band alignment, especially with different doping levels.

This bending promotes accelerated electron mobility in Sr_3PBr_3-based photovoltaic cell, enhancing the photovoltaic cell's capability for peak efficiency. The band energy diagram provides precise predictions of how holes and electrons move within the designed PSC, despite a minor offset that does not impede carrier transport as it remains below thermal energy levels.

The alignment of the energy bands plays a pivotal role in governing current flow within a heterostructure [53–55]. In this study, we systematically investigated several HTLs to identify the one that delivers the *PCE*. Specifically, HTLs including Cu_2O (Figure 2(c)), CuI (Figure 2(d)), and $CuSbS_2$ (Figure 2(e)) were evaluated. For comparison, Figure 2(a) illustrates the optimized energy band diagram of the device without an HTL, while Figure 2(b) presents the detailed insights into the bandgap configuration incorporating CBTS as the HTL, which demonstrated superior *PCE* relative to the other candidates revealing pronounced band bending at the CBTS/Sr_3PBr_3 interface. This favorable band alignment, particularly under varying doping concentrations, facilitates more efficient charge separation and transport. The observed band bending accelerates electron extraction and transport in Sr_3PBr_3-based PSCs, thereby enhancing

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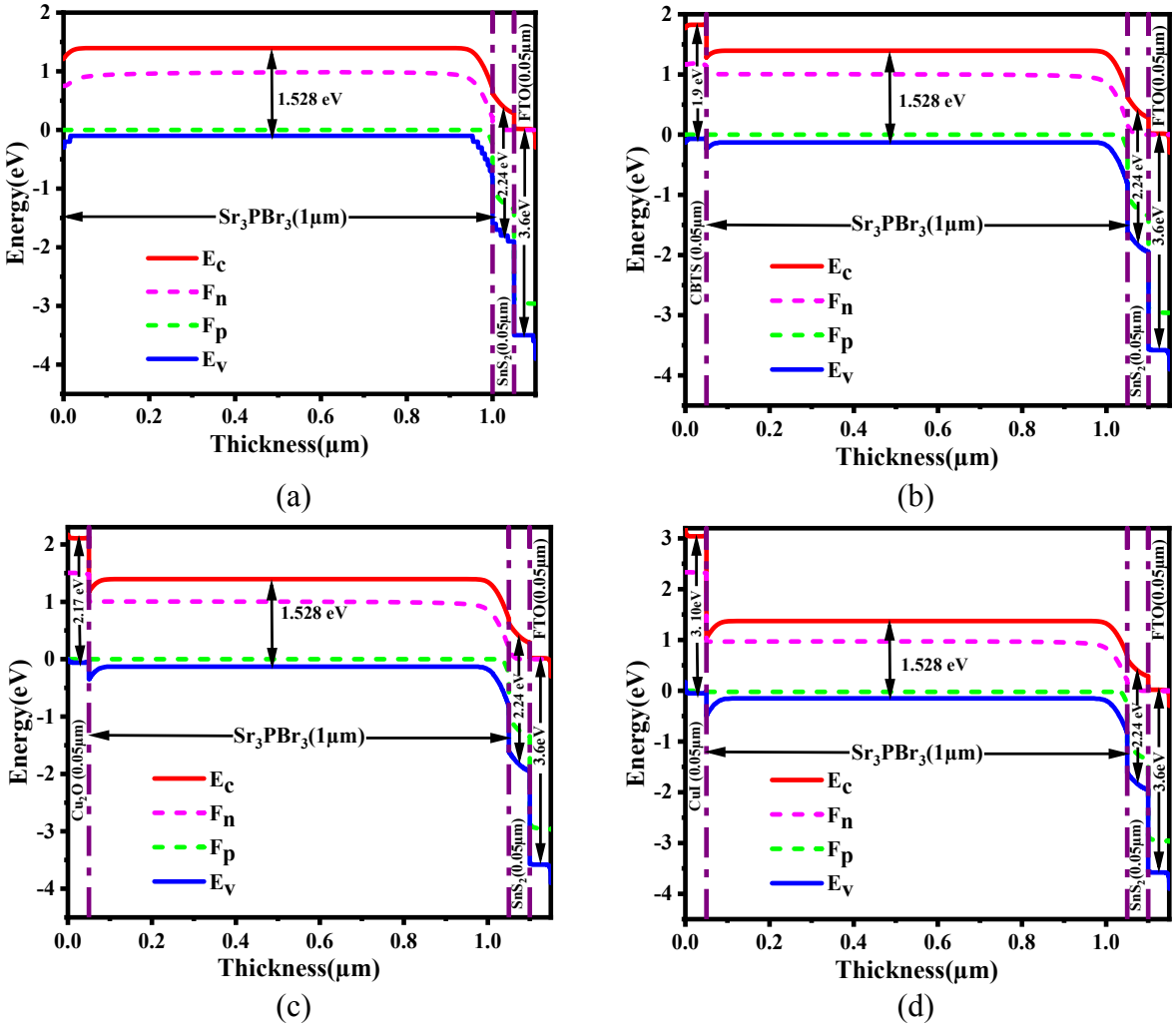
the device's capability to achieve peak performance. Furthermore, the energy band diagram accurately predicts the movement of charge carriers within the architecture, where the minimal offset present does not hinder carrier transport, as it remains well below the thermal energy threshold. Therefore, CBTS was selected as the optimized HTL and employed in all subsequent simulations and analyses throughout the remainder of this study.



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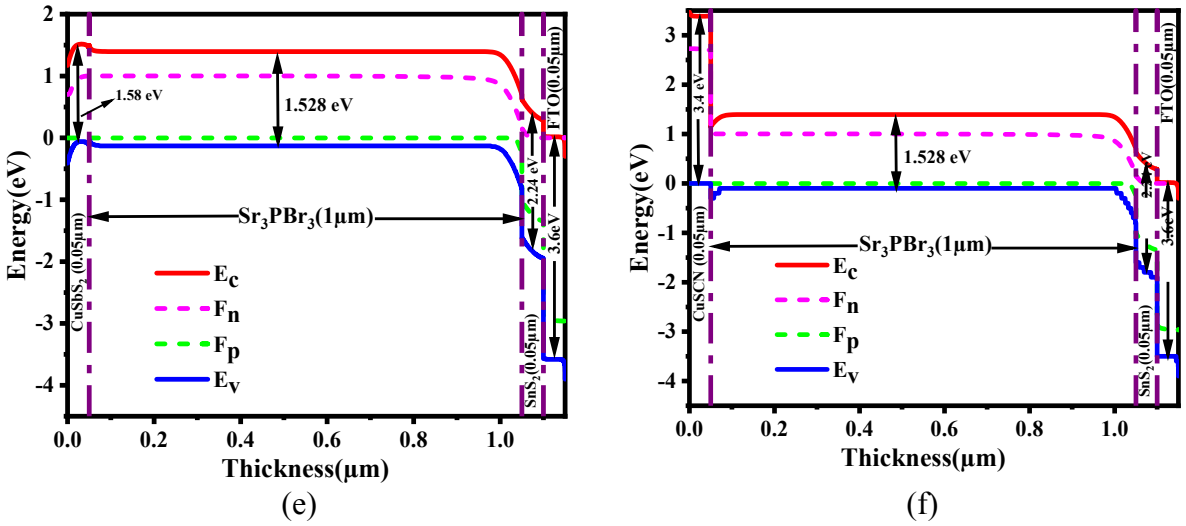

Figure 2. Band configuration of a Sr_3PBr_3-based photovoltaic cell (a) without HTL, and with HTL (b) CBTS, (c) Cu_2O, (d) CuI, (e) $CuSbS_2$, and (f) CuSCN.

3.3. Impact of Sr_3PBr_3 layer thickness and doping concentration

Figure 3 (a) illustrates the effect of varying thickness of absorber on the output characteristics of Sr_3PBr_3-based PSCs with and without CBTS HTL. As the perovskite Sr_3PBr_3 layer becomes thicker in the configurations, the J_{SC} increases, achieving a peak value of 26.839 mA/cm^2 at 1.0 μm. Subsequently, J_{SC} increases slightly with further thickness because recombination dynamics influence charge separation in the PV cell. Moreover, with a Sr_3PBr_3 thickness of 1.0 μm, overall absorption is predicted to escalate, leading to a higher exciton production rate and consequently increasing electron creation in the absorber layer due to improved photon absorption. Additionally, the changes in V_{OC} with varying Sr_3PBr_3 thicknesses from 0.1 μm to 2.0 μm are presented. The V_{OC} attains its maximum at a thickness of 1.0 μm and then remains relatively constant, a trend primarily governed by recombination dynamics at higher thicknesses. With increasing thickness of the absorber layer, there is typically a corresponding rise in the device's *FF*, potentially leveling off around a thickness of 1.0 μm. Changes in performance across different thicknesses of Sr_3PBr_3 in PV cells are evident, influenced directly by thickness variations affecting *FF*, V_{OC} and J_{SC} with peak efficiency observed around 1.0 μm which is therefore chosen as the optimized absorber thickness. Perovskite's advantage lies in its ability to

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absorb a broader spectrum of photons across a wider range of frequencies, attributed to its larger bandgap [56]. With 1.0 μm thickness, the diffusion distances are L_p = 34 μm and L_n = 34 μm.

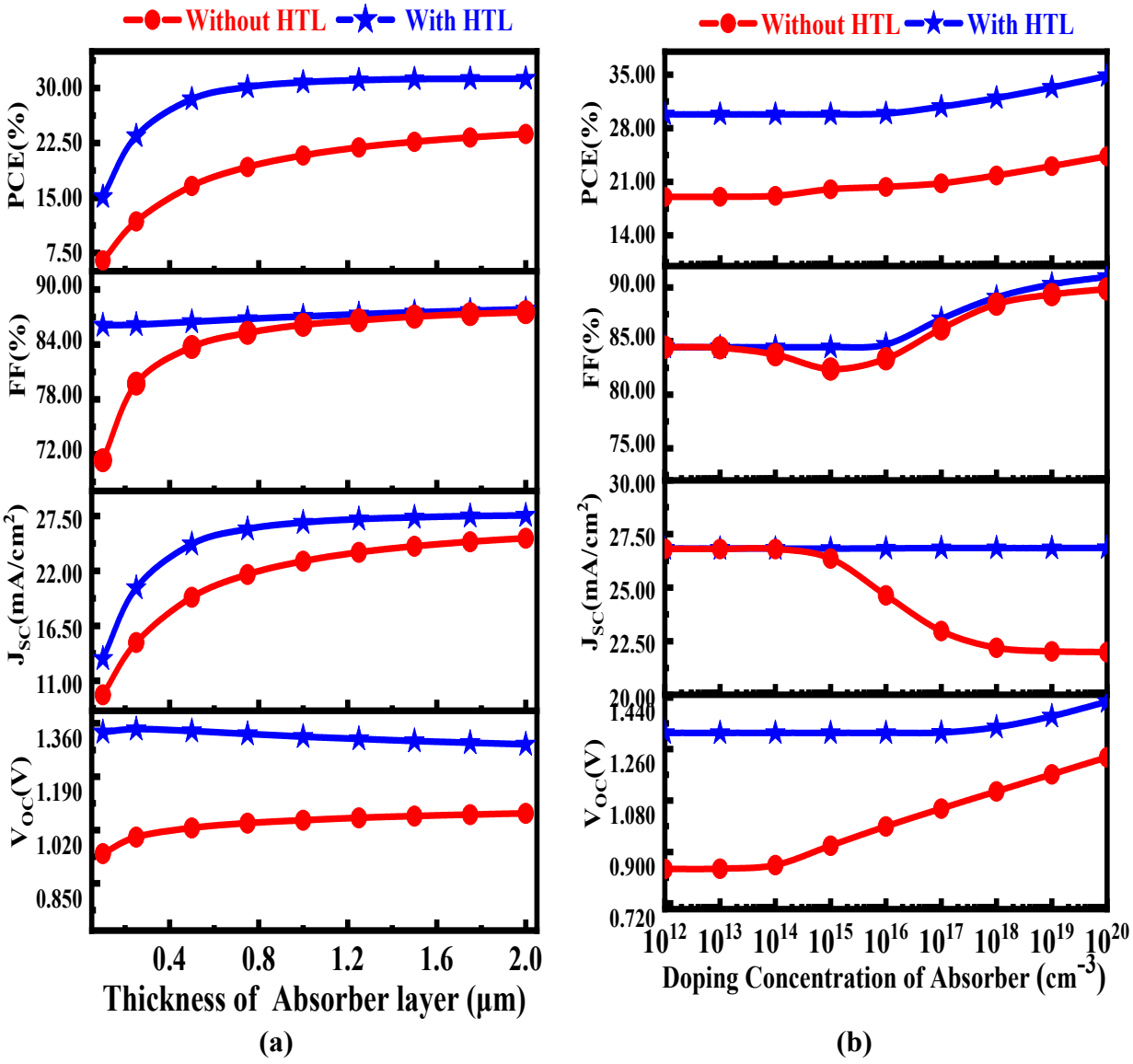

Figure 3. Impact of varying (a) the thickness of Sr_3PBr_3, and (b) the doping density in Sr_3PBr_3, without and with the CBTS HTL layer.

Enhancing the output of the PV cell involves precisely doping the absorber material[57]. Nevertheless, extreme doping may present challenges owing to its non-traditional nature, possible hindrance to carrier transport and added structural complexity in manufacturing.

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Therefore, the active layer doping density was changed from 10^{12} cm^{-3} to 10^{20} cm^{-3} in simulations to ascertain the optimal value. Figure 3(b) depicts how varying the N_A affects the Sr_3PBr_3 absorber layer without HTL and with CBTS HTL. The goal was to raise V_{OC} from 1.316 V to 1.425 V with HTL and 0.840 V to 1.241 V without HTL adjusting N_A of the Sr_3PBr_3 material from 10^{12} cm^{-3} to 10^{20} cm^{-3}. The J_{SC} rises from 26.799 to 26.841 mA/cm² with HTL and without HTL reduces J_{SC} from 26.789 to 21.988 mA/cm^2. The *FF* and *PCE* show a slight increase from 10^{12} cm^{-3} to 10^{20} cm^{-3} for both the cases.

3.4. Effect of changes in absorber thickness and density of defect on solar cell performance

The absorber layer's thickness and defect density play a crucial role in determining solar cell performance. An enhance in defects causes the film to deteriorate and form pinholes, leading to higher recombination rates that degrade the steadiness and *PCE* of PSCs [58]. Therefore, to determine the optimal defect density for a given absorber layer thickness, simulations were carried out with absorber thicknesses ranging from 0.1 µm to 2.0 µm, while adjusting the defect value from 10^{10} cm^{-3} to 10^{16} cm^{-3}. Figure 4(a–d) illustrates the variations in J_{SC}, V_{OC}, *FF*, and *PCE* for the FTO/SnS_2/ Sr_3PBr_3/CBTS/Au device configuration, providing a comprehensive depiction of how each photovoltaic parameter responds to the applied simulation conditions. It is apparent that the solar cells achieved peak V_{OC} values for absorber layer thicknesses ranging from 0.3 µm to 2.0 µm, and defect densities between 10^{10} cm^{-3} and 10^{12} cm^{-3}. This enhancement is due to the increase in quasi-Fermi level seperation as the thickness upsurges [59].

Figure 4(a) depicts how the J_{SC} values vary with active layer thickness and defect density. Upon analyzing the contour plots, a consistent pattern emerges across all solar cells regarding these variables. Remarkably, the solar cell reached its peak J_{SC} value when thickness of the active layer surpassed 1.0 µm, while maintaining a defect density below 10^{12} cm^{-3}. This rise in J_{SC} values is due to improved spectral response at greater thicknesses and minimized charge carrier recombination [60]. Figure 4 (c) illustrates the *FF*, highlighting an optimal thickness of 1.0 µm while keeping the defect density below 10^{12} cm^{-3}. Figure 4 (d) depicts how changing the defect density absorber thickness impacts the *PCE*. Increasing the absorber thickness enhances the *PCE*, while increasing the defect density decreases it.

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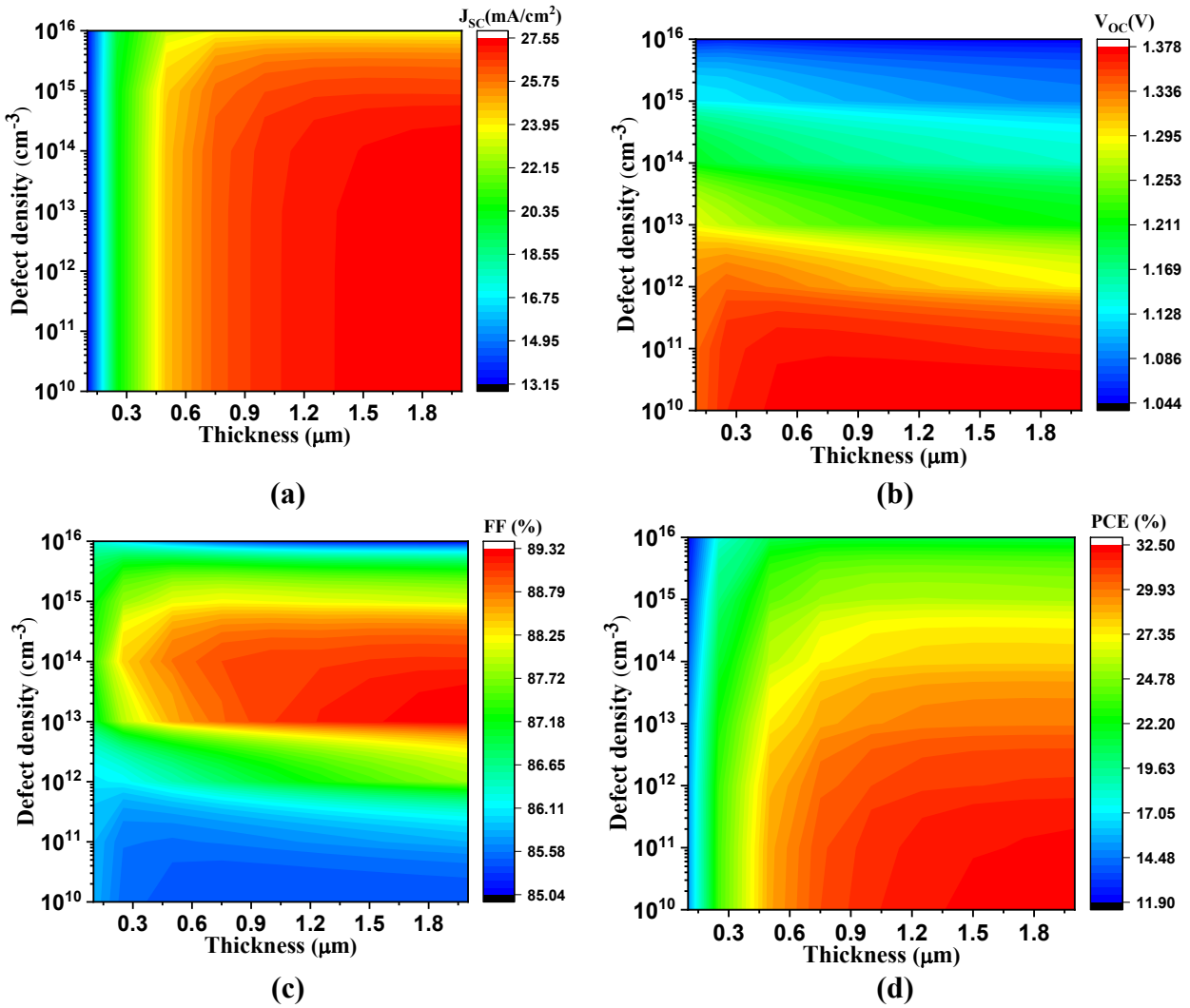

Figure 4. Contour plots illustrating the impact of photovoltaic performance metrics due to variations in active layer thickness and defect density: (a) J_{SC}, (b) V_{OC}, (c) *FF* and (d) *PCE*.

3.5. Impact of ETL thickness and Donor Density

An ETL is essential for improving light transmittance and reducing recombination inside PSCs [61]. Hence, it is crucial to adjust the characteristics of charge transport layers accordingly. In this research, we adjusted the SnS_2 thickness from 0.03 to 0.3 μm while keeping all other parameters unchanged and have shown Figure 5(a) to 5(d). Figure 5(a) illustrates the role of ETL thickness on the photovoltaic output. The variation in SnS_2 thickness has an almost negligible effect on J_{SC} , with values ranging only slightly from 26.840 mA/cm² at 0.03 μm to 26.831 mA/cm² at 0.3 μm. For V_{OC}, a comparable trend is observed. Significantly, the peak efficiency of 30.78% is

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accomplished with a SnS_2 layer thickness of 0.05 µm. Figure 5(b) demonstrates the role of varying the N_D within the ETL (SnS_2) layer on the performance of the proposed PSC device, spanning from 10^{12} to 10^{20} cm^{-3}. It has been observed that an increase in ETL doping concentration, V_{OC} remains the same for both conditions, whereas J_{SC} shows a marginal enhancement up to 10^{16} cm^{-3}, after which it exhibits a significant upsurge. Up to 10^{16} cm^{-3}, *FF* and *PCE* almost remains constant trend with doping concentration in HTL.

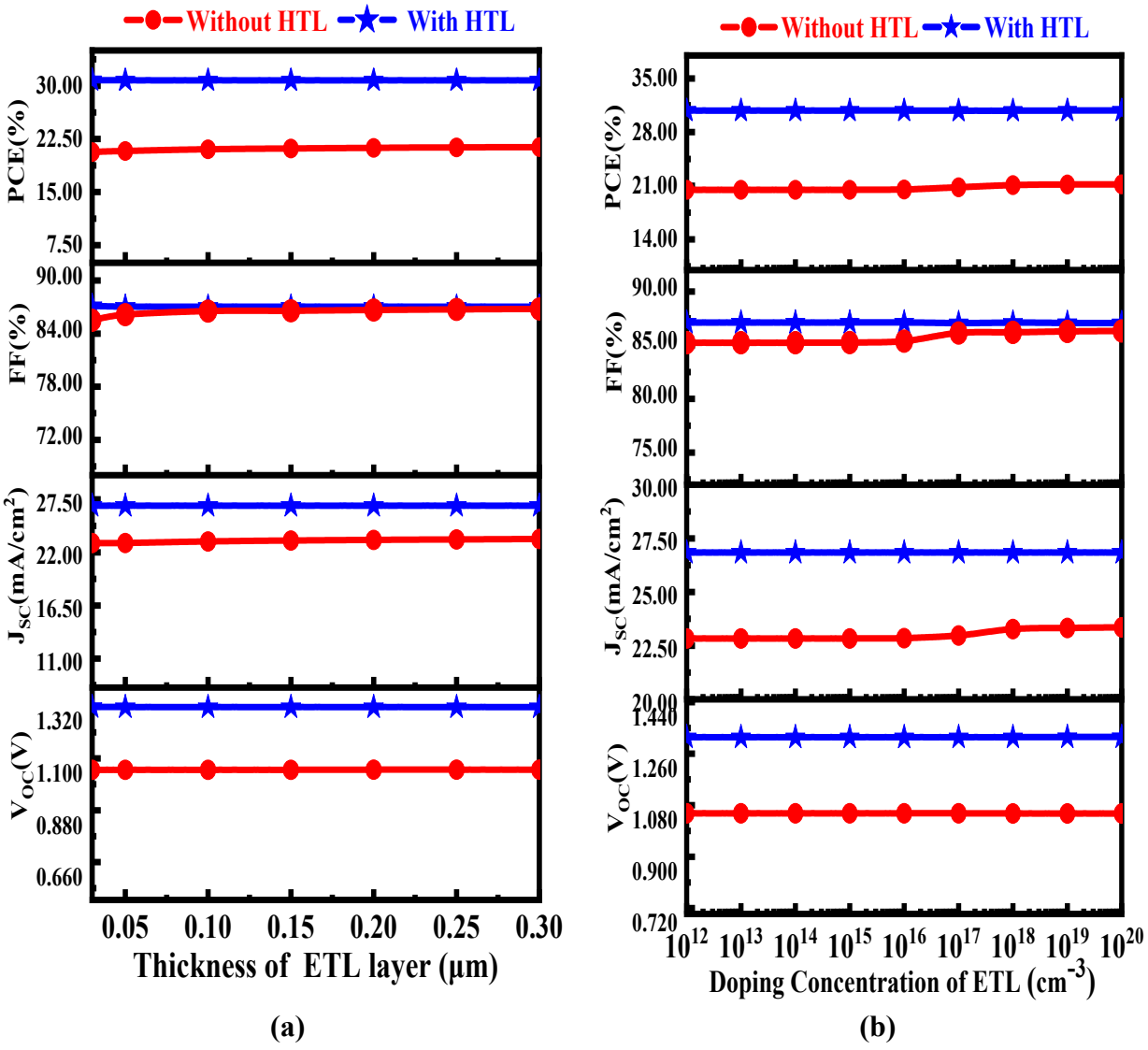

Figure 5. Effect of varying the (a) thickness of SnS_2, without and with the CBTS HTL layer, and (b) donor concentration in SnS_2, without and with the CBTS HTL layer.

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3.6. Impact of Defect Density at the ETL/Sr_3PBr_3 and HTL/ Sr_3PBr_3 Interfaces

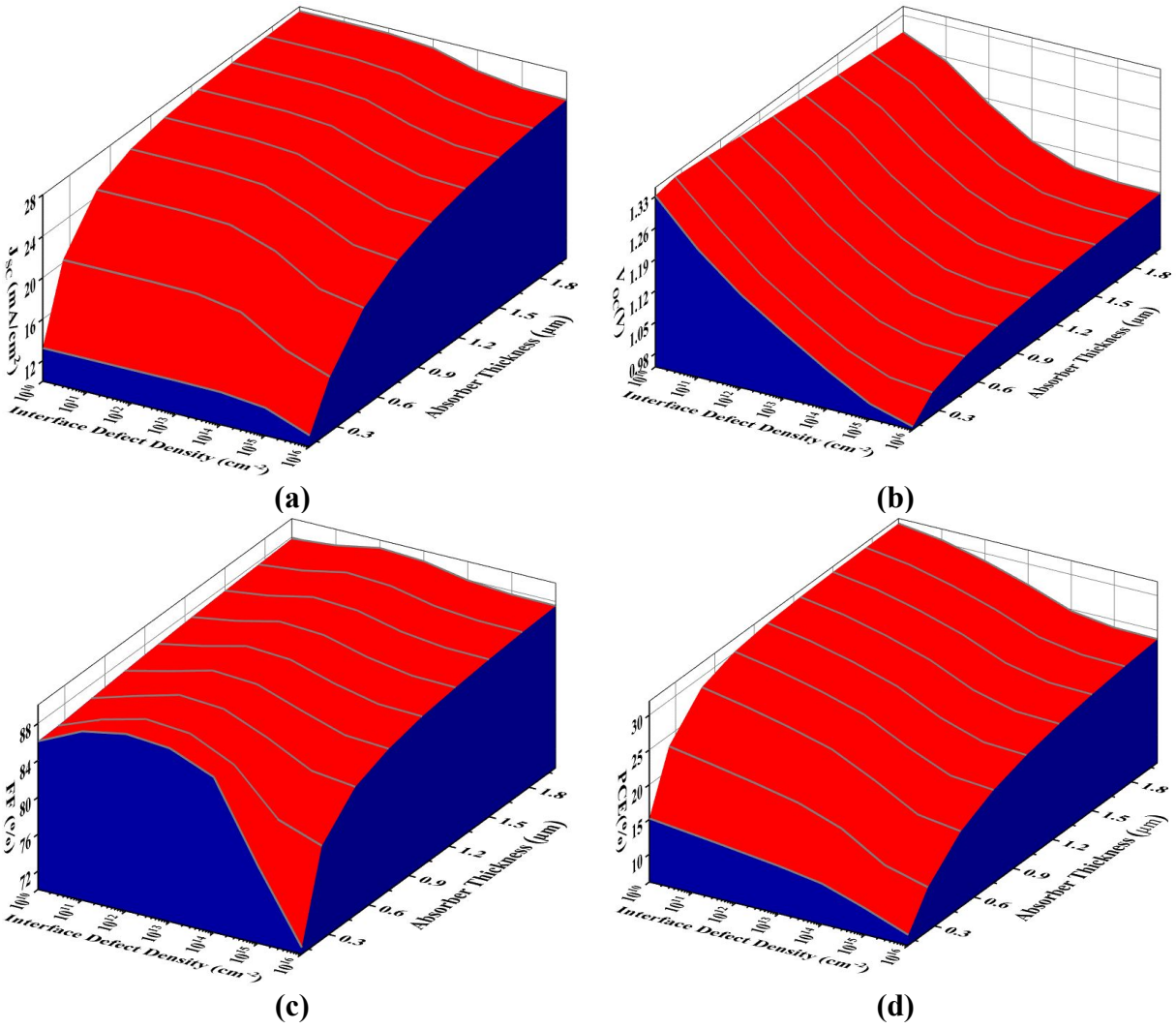

Figure 6. Simultaneous effect of changing the interface defect density (Sr_3PBr_3/CBTS) on PV output.

During the fabrication of solar cells, interface defects arise from structural defectiveness. These defects upsurge charge carrier recombination at the interface, thereby compromising the PSC's performance [62]. Therefore, investigating their impact and estimating optimal values for practical production is crucial. In our analysis, a neutral interface defect density of 10^{10} cm^{-2} was applied at both the ETL/Sr_3PBr_3 and Sr_3PBr_3/HTL interfaces with defect levels maintained at 0.6 eV across all PV cells. Here, we evaluate the impact of interface defect density (IDD) on the PV output by changing it from 10^{10} cm^{-2} to 10^{16} cm^{-2} and the thickness are ranging from 0.1 to 2.0

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μm for both interfaces. Figure 6 displays the influence of changing the IDD between CBTS and Sr_3PBr_3 affects the performance of PSC. The J_{SC}, FF, PCE, and V_{OC} of the Sr_3PBr_3-based PSCs decline from 27.50 mA/cm^2 to 11.10 mA/cm^2, 87.8% to 70.70%, 31.3% to 7.49%, and 1.341 V to 0.957 V, respectively when IDD changing from 10^{10} cm^{-2} to 10^{16} cm^{-2} at Sr_3PBr_3/CBTS interface.

Figure 7, shows the results for SnS_2/Sr_3PBr_3 interface defects. At the SnS_2/Sr_3PBr_3 interface, higher interface defect density significantly impacts the J_{SC}, which decreases from 27.52 to 13.25 mA/cm². Additionally, V_{OC} drops from 1.33 V to 0.795 V, the *FF* reduces from 87.85% to 75.52%, and the *PCE* declines from 31.30% to 7.93%.

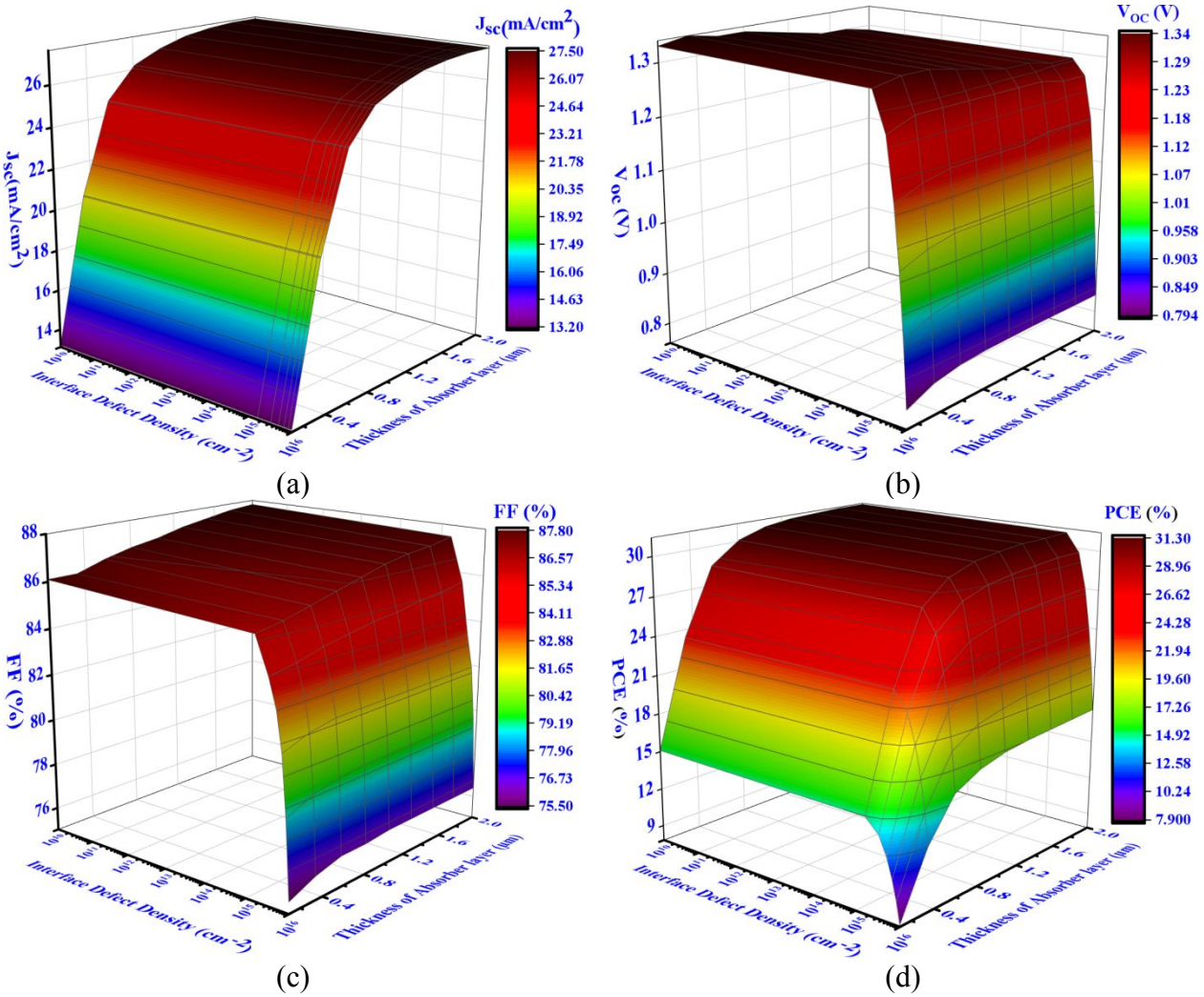

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Figure 7. The concurrent impact of changing the IDD at Sr_3PBr_3/SnS_2 interface and absorber thickness on the performance of PSC.

3.7. Influence of Series and Shunt Resistance

Parasitic resistances, such as R_S and R_{SH} are crucial factors in solar cell output. They affect the current-voltage (J-V) characteristics of the PSCs and represent sources of energy loss within the device. R_S refers to the overall resistance spanning various layers of solar cells. Conversely, R_{SH} arises from the reverse saturation current in PSCs caused by imperfections introduced during fabrication. Here, we investigate the impact of R_S and R_{SH} on solar cell parameters to see the actual behavior. Figure 8(a) depicts the output parameters, as R_S varies from 1 to 6 Ωcm² for the optimized (FTO/SnS_2/Sr_3PBr_3/CBTS/Au) solar cells. It is observed that V_{OC} and J_{SC} are unchanged across the entire range of R_S values. In the meantime, the *FF* declines significantly from 81.43% to 64.08%, and the PCE decreases from 30.78% to 26.69%. The notable reduction in *FF* is ascribed to considerable power loss (P_{loss}) within PSCs under high R_S conditions, which unfavorably impacts their overall performance [63]. As R_S increases, the corresponding rise in power loss (P_{loss}) inside the PSCs directly affects the *FF*.

Similarly, R_{SH} is varied from 10 to 10^7 Ω cm², as shown in Figure 8(b). In this scenario, J_{SC} exhibit minimal variation across all R_{SH} values. Conversely, the output of all solar cells rise as R_{SH} is increased from 10 to 10^3 Ω cm². Beyond this range, a linear relationship is observed.

Figure 8 (c) shows the series resistance dependence of photovoltaic parameters without HTL. It is observed that V_{oc} is increased with rising series resistance while other parameters are decreased. Figure 8 (d) displays the variation of photovoltaic parameters with shunt resistance without HTL and exhibited the similar trend as in Figure 8 (b) with enhanced value of all the parameters.

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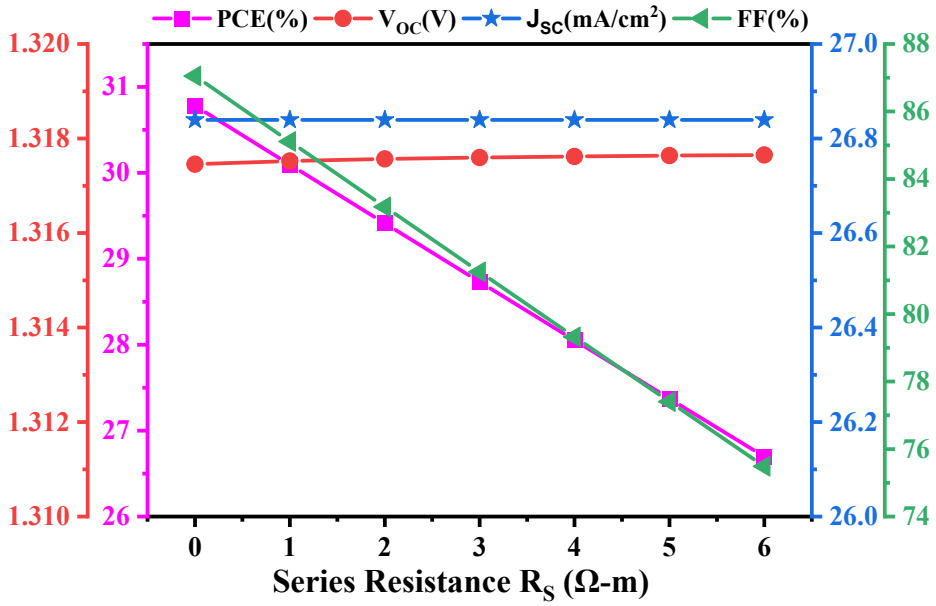
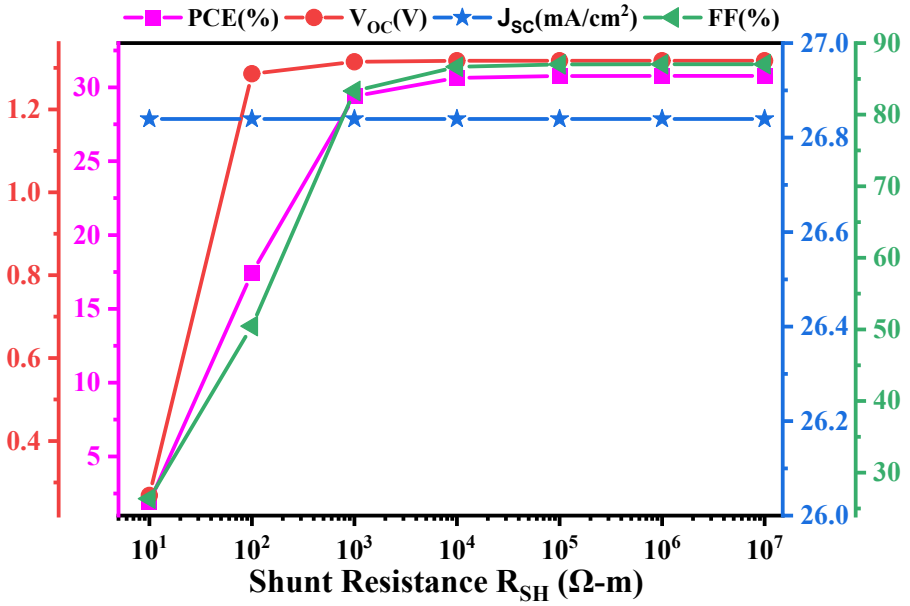


(a)



(b)

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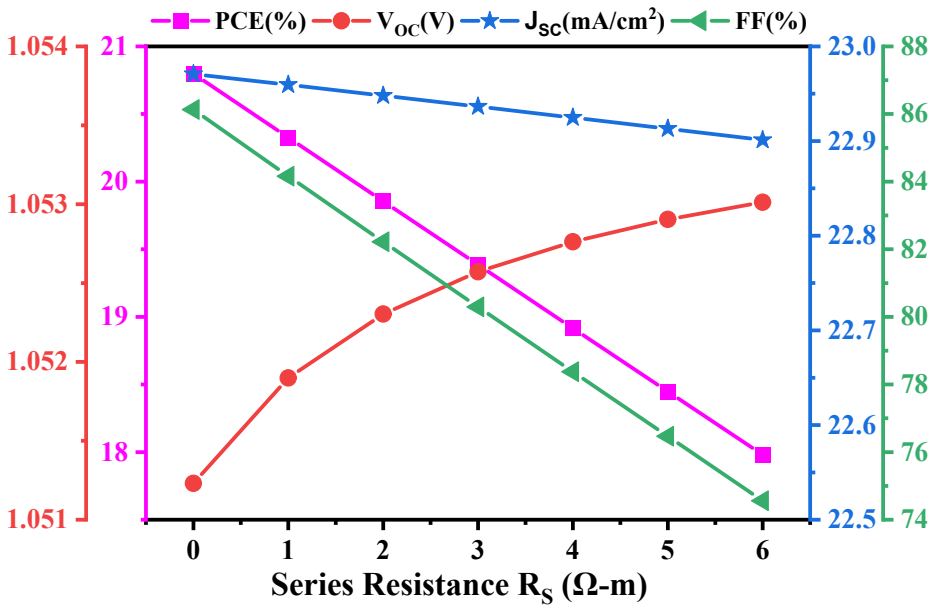
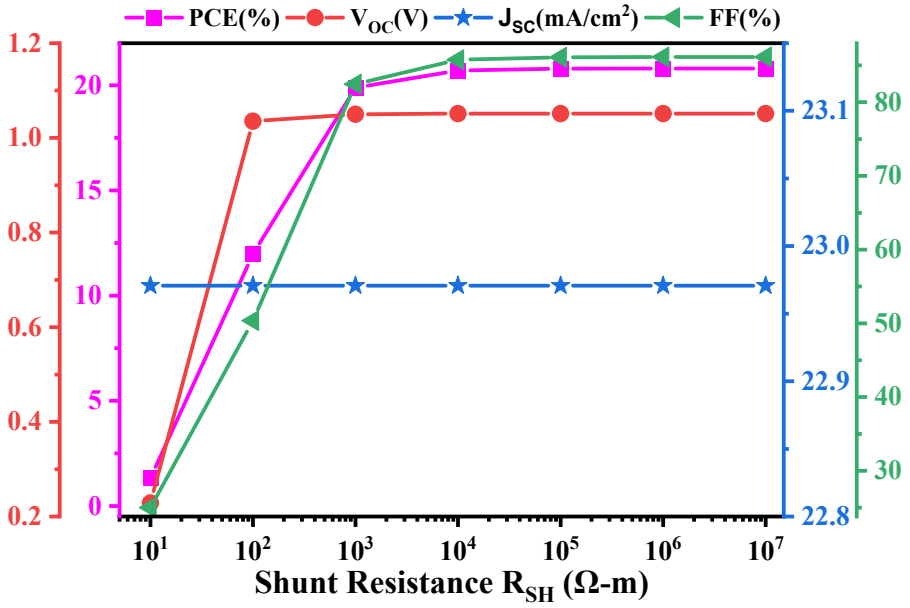


(c)



(d)

Figure 8. Influence of (a) Series resistance, and (b) Shunt resistance with CBTS HTL (c) Series resistance, and (b) Shunt resistance without CBTS HTL.

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3.8. Effect of Temperature

The operating temperature of PSCs is vital due to their exposure to ambient atmospheric conditions, necessitating extended durability. Consequently, it is essential to comprehend the degradation mechanisms of solar cells under real world conditions [64,65]. In this context, we systematically varying the working temperature of (FTO/SnS_2/Sr_3PBr_3/CBTS/Au) solar cells from 300 K to 420 K (Figure 9 (a)) to thoroughly evaluate its comprehensive impact. It is evident that as the temperature upsurges, the values of V_{OC}, *FF*, and *PCE* reduces, whereas J_{SC} maintains consistent levels. The decrease in V_{OC} values is due to the heightened vibration of thermally generated electrons at elevated temperatures, making them less stable and more possibly to recombine with holes, which in turn raises the reverse saturation current (J_O) [66]. This trend is evident from the inverse correlation between V_{OC} and J_O, as demonstrated in equation (6).

$$V_{OC} = \frac{nKT}{q}\left(\ln\left(1+\frac{J_{SC}}{J_O}\right)\right) \quad (6)$$

Where $\frac{KT}{q}$ represents the thermal voltage [67].

Additionally, the physical metrices including charge carrier mobility and carrier concentration are negatively impacted by increasing temperatures, which directly affects the transportation efficiency of charge carriers and consequently reduces the *FF* of every PSCs [68]. This combined reduction in V_{OC} and *FF* results in a lessening in *PCE* from 30.78% to 26.39% for CBTS-based solar cells. Figure 9 (b) also illustrates the *J-V* characteristics at varying temperatures, highlighting the significance of temperature effects. Figure 9 (c) and 9 (d) shows the effect of temperature on PSC's output parameters and the *J-V* properties of the PSC in different temperature, without HTL.

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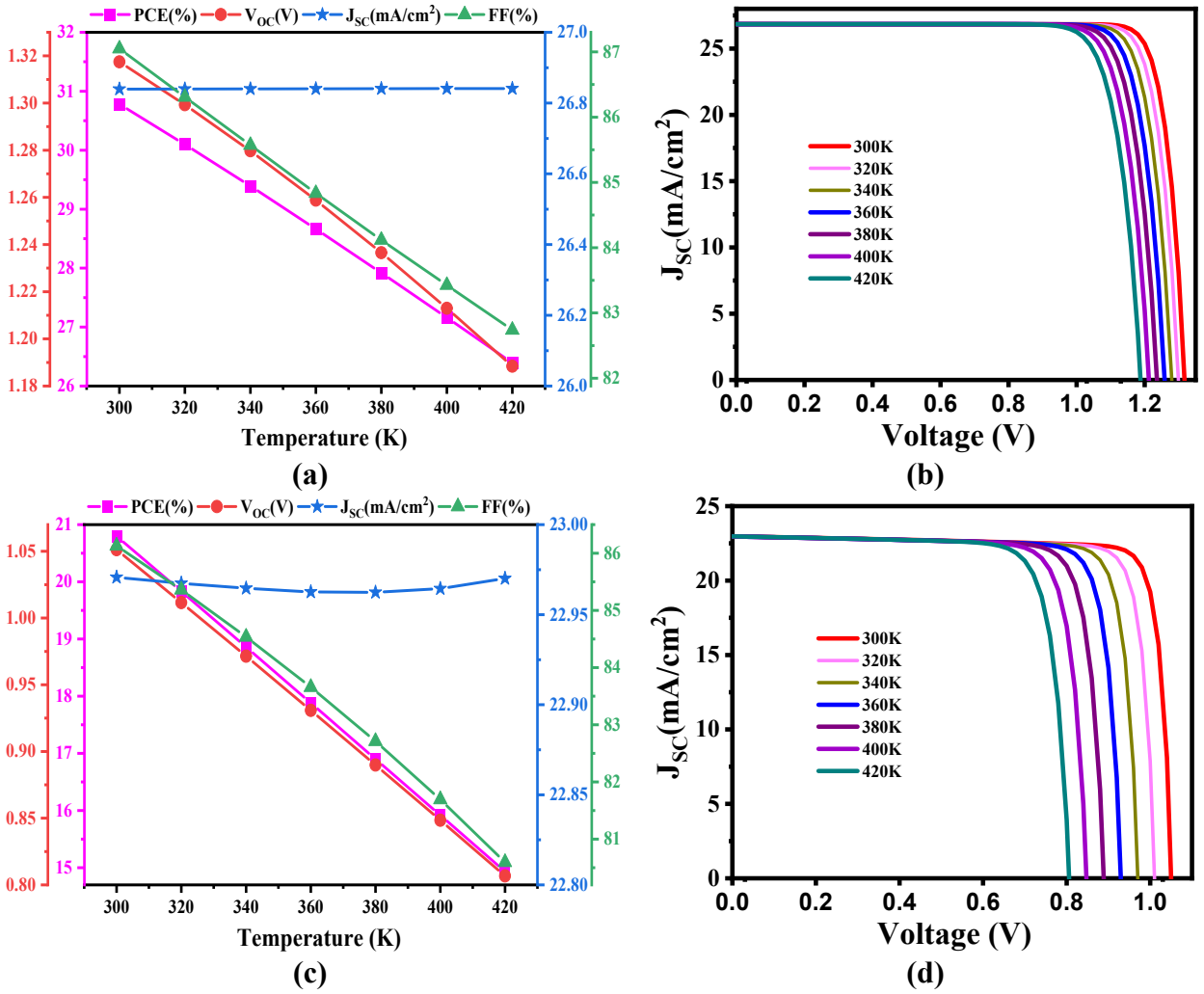

Figure 9. The effect of temperature on (a) solar cell parameters, (b) the *J-V* properties with CBTS HTL, (c) solar cell parameters, and (d) the *J-V* properties without HTL.

3.9. Current-Voltage (*J-V*) and Quantum Efficiency (*QE*) Characteristics of the Evaluated PSC Devices

Figure 10 (a) exhibits the current-voltage (*J-V*) properties of the device. Furthermore, Table 3 displays the PSC parameters without and with the HTL. The influence of the CBTS HTL in the proposed structure of Sr_3PBr_3-based solar cells is crystal clear to enhance the photovoltaic performance. Figure 11(a-b), demonstrates the *J-V* and *QE* curves of Sr_3PBr_3-based solar cells with HTL (FTO/SnS_2/Sr_3PBr_3/CBTS/Au) and without HTL (FTO/SnS_2/ Sr_3PBr_3/Au). It is evident from Table 3 that the efficiency and *FF* are improved after incorporating CBTS as the

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HTL. Under optimized device conditions, the J_{SC} with and without the HTL is measured at 26.839 mA/cm^2 and 22.970 mA/cm^2, respectively. This is attributed to the interfacial recombination of charge carriers occurring between Sr_3PBr_3 and the rear contact metal (gold). Figure 10 (b) exhibits a noticeable decline in *QE* response beyond 820 nm, verifying the band gap of Sr_3PBr_3, which is measured at 1.528 eV. It is apparent that the *QE* response of solar cells utilizing CBTS as the HTL is superior to that of PV cells without HTL. This variance in *QE* characteristics can be ascribed to the creation of a back surface field after the integration of the HTL [69,70]. Sr_3PBr_3 serves dual roles as both the absorber and p-type layer in the solar cell structure. Introducing another p-type layer such as CBTS improves the overall performance of the PSC. Therefore, CBTS as an HTL is important aimed at boosting the operational effectiveness of the proposed PSC.

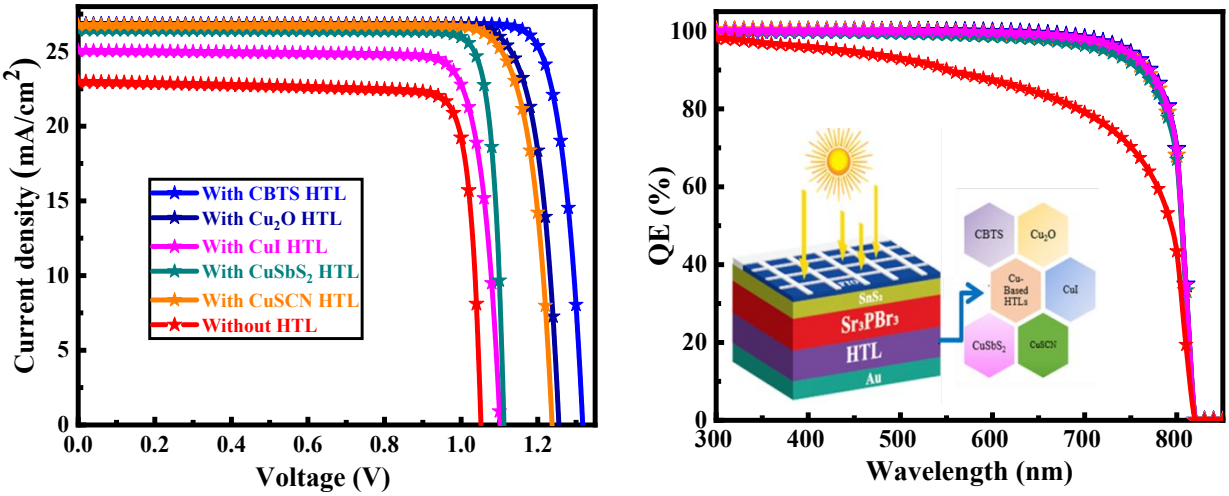

Figure 10. (a) Current-Voltage (J-V) characteristics and (b) Quantum Efficiency (QE) graph of the proposed cell, both with and without the CBTS HTL layer.

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Table 4. The device performance under both ideal and non-ideal parameter conditions

Case	V_{OC} (V)	J_{SC} (mA/cm^2)	FF (%)	PCE (%)	R_S (Ω·cm^2)	R_{SH} (Ω·cm^2)	Defect Density (cm^{-3})
Ideal	1.377	26.84	85.40	31.51	0	6000	1×10^{10}
Non-ideal	1.0311	23.25	50.62	12.15	2.0	100	1×10^{16}

Table 4 compares the simulated photovoltaic performance of the Sr_3PBr_3-based PSC under ideal and non-ideal conditions. In the ideal case, the device achieves a high PCE of 31.51%, supported by a large V_{OC} (1.377 V), high FF (85.40%), and minimal recombination losses due to negligible series resistance, extremely high shunt resistance, and low bulk defect density (1×10^{10} cm^{-3}). In contrast, the non-ideal case incorporates practical loss mechanisms—namely increased R_S (2 Ω·cm^2), reduced R_{SH} (100 Ω·cm^2), and a high defect density (1×10^{16} cm^{-3})—which significantly degrade carrier transport and increase recombination. As a result, V_{OC}, J_{SC}, and FF decline sharply, lowering the overall efficiency to 12.15%. This comparison clearly demonstrates the sensitivity of Sr_3PBr_3 PSCs to resistive and recombination losses, and validates that the simulated performance remains consistent with the SQ-limit theoretical framework.

3.10. Machine Learning

In our work, the Random Forest algorithm's internal metric was used to generate the feature importance percentages. By calculating each feature's effect on lowering impurity during data splits across the decision tree nodes, this metric evaluates each feature's contribution [71]. The *PCE* of PSCs was precisely predicted in this study using the Random Forest algorithm, which also identified the key characteristics influencing device performance. With *PCE* (%) as the desired output, the dataset included ten crucial parameters that are known to affect photovoltaic efficiency. Following data preparation and cleaning, an 80/20 train-test split was used to give the model a large amount of training data while retaining a separate portion for performance evaluation. This study emphasizes important parameters like bandgap, interface defects, and doping concentration as the most important factors influencing solar cell performance, as illustrated in Fig. 11. The analysis provides insightful direction for upcoming photovoltaic

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technology research and development by highlighting these crucial characteristics. Below are the corresponding relative importance scores for the characteristics that have an impact on *PCE*:

Table 5. Scores for different important characteristics

Parameters	Value
Acceptor Density	0.359222
Defect Density	0.322507
CB	0.121174
Electron Affinity	0.097420
VB	0.050575
Thickness	0.038474
Band Gap	0.010628

The findings highlight the significant influence of acceptor density, which is a key factor in determining how well photovoltaic devices convert energy. Similarly, the significance of defect levels emphasizes how they improve charge transport inside the device and lower recombination losses. Another important consideration is doping, which has a direct impact on carrier concentration, which is necessary to achieve the best possible device performance and have shown Fig. 5.

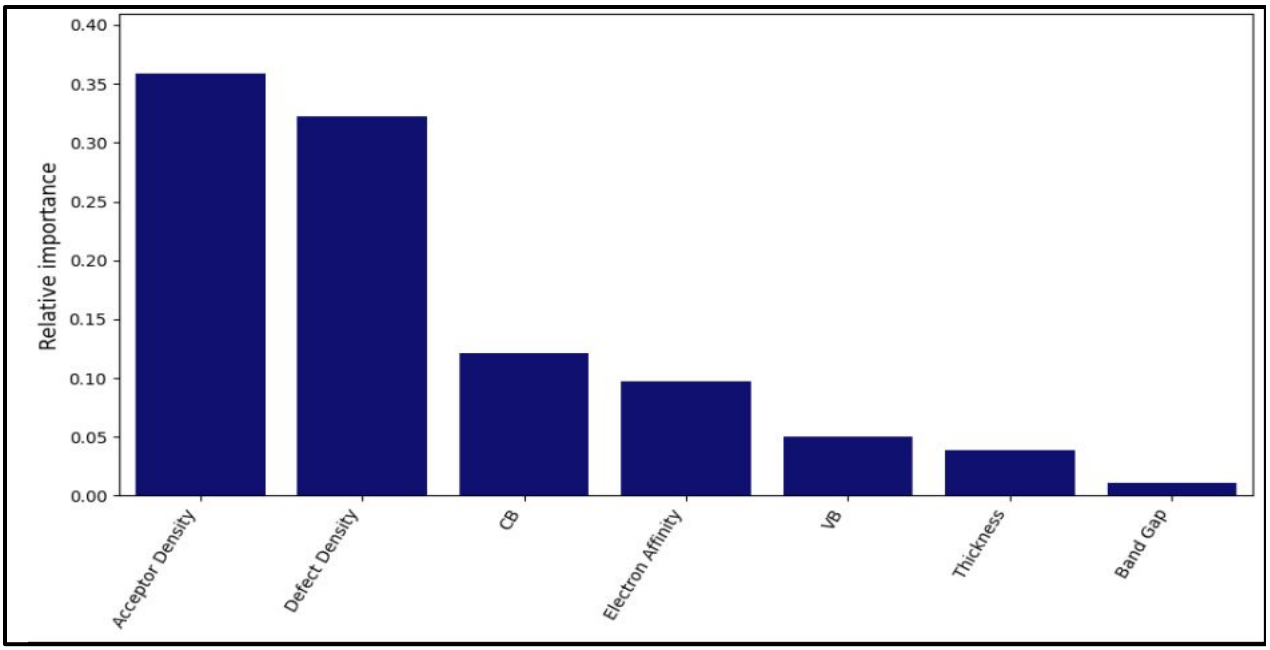

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Figure 11. Several characteristics that are most crucial in the ML mode.

The SHAP plot, which is displayed in figure 12, shows how different features affect the model's output for forecasting solar cell performance. A data point is represented by each dot, which is colored according to the feature's value, spanning from low (blue) to high (red). For instance, a high defect density (shown by the red dots on the left) significantly reduces the model's output, suggesting a detrimental effect. High acceptor density values, on the other hand, have a positive effect and raise the output. Other characteristics, such as Band Gap and Electron Affinity, also exhibit distinct effects, underscoring their importance in the prediction model.

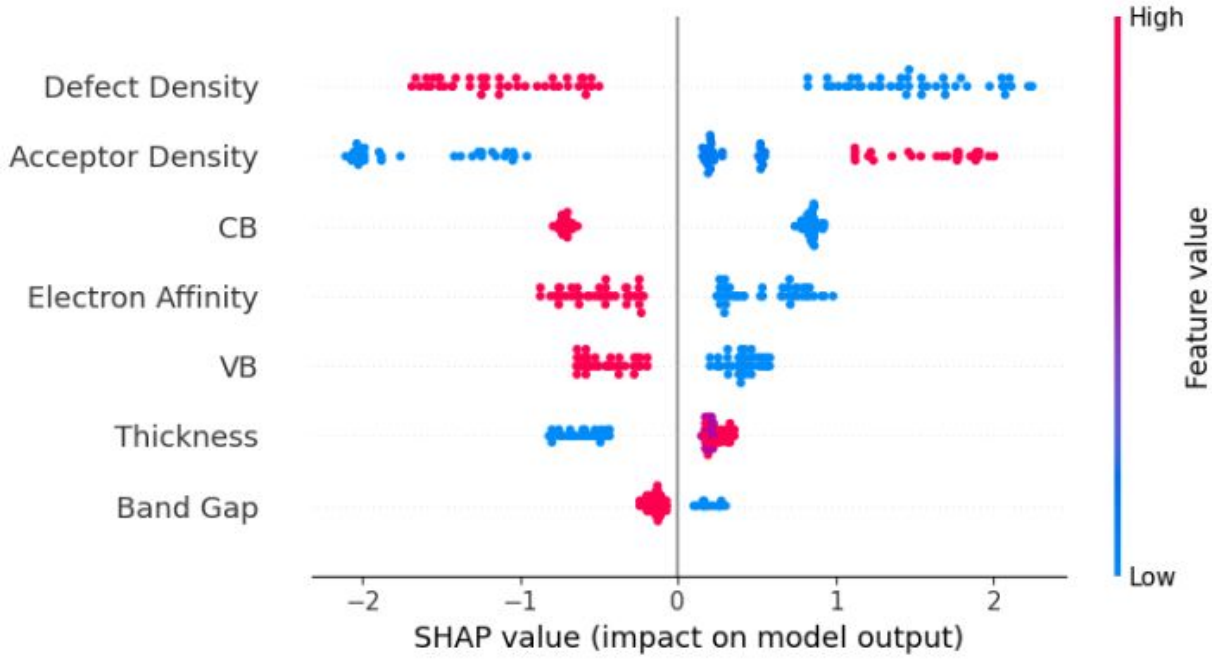

Figure 12. Random forest regression in conjunction with SHAP values was used to examine the impact of unrelated features on the device's *PCE*.

The coefficient of determination (R^2), which shows how much of the variation in the target variables the model can explain, was used to evaluate the predictive performance of the model (see Figure 13). For *PCE*, V_{OC}, J_{SC}, and *FF*, the corresponding R^2 values were 0.9943, 0.9941, 0.9944, and 0.9929, respectively. The model's high effectiveness in elucidating the majority of the variability in these significant output metrics is demonstrated by the high R^2 scores for *PCE*

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and V_{OC}. The somewhat reduced R^2 for J_{SC} indicates that although the model successfully reflects the primary patterns, a few variations might still exist, most likely as a result of experimental noise or unaccounted-for variables. The Random Forest algorithm's strong performance in forecasting the behavior of photovoltaic devices is highlighted by its average R^2 of 0.9939.

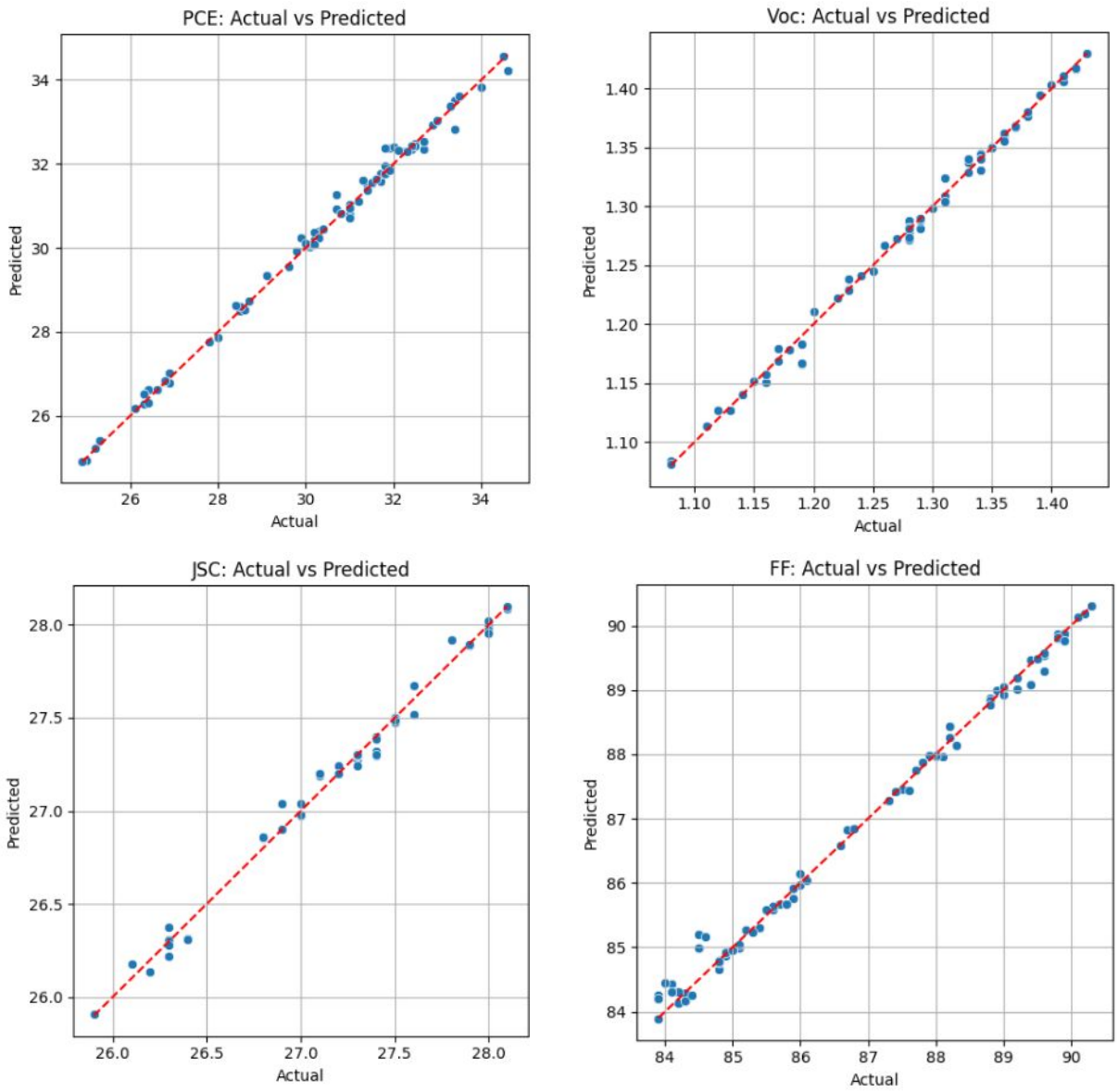

Figure 13. The assessed and actual values of *PCE*, V_{OC}, J_{SC}, and *FF* showed a linear relationship, suggesting that the model predicts these important photovoltaic parameters with little variation.

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The reported performance metrics represent theoretical upper limits obtained from SCAPS-1D simulations and machine learning predictions under idealized conditions. In practical devices, additional non-idealities such as incomplete light absorption, optical reflection, recombination at bulk and interfacial defects, series resistance, and long-term stability issues will inevitably reduce efficiency compared to the simulated values. Thus, the results presented here should be interpreted as benchmarks for guiding experimental optimization rather than directly achievable efficiencies, ensuring alignment with realistic photovoltaic expectations.

4. Conclusions

This work presents a comprehensive computational investigation of novel Sr_3PBr_3-based perovskite solar cells incorporating various inorganic Cu-based hole transport layers (HTLs), modeled using SCAPS-1D. The simulated device structure, FTO/SnS_2/Sr_3PBr_3/CBTS/Au, was systematically optimized by evaluating absorber characteristics under variations in thickness, acceptor density, and defect density, with emphasis on their interdependent effects. Performance enhancement was achieved through fine-tuning donor/acceptor densities, minimizing defect states, and optimizing thicknesses of the ETL, HTL, and absorber layers. The optimal configuration featured an absorber thickness of 1.0 μm, HTL thickness of 0.05 μm, and doping concentrations of 1.0×10^{17} cm^{-3} for both absorber and HTL. Under these conditions, the device achieved an impressive *PCE* of 30.78%, with J_{SC} = 26.839 mA/cm², V_{OC} = 1.317 V, and *FF* = 87.05%. Results confirm that thin SnS_2 layers are a promising, non-toxic alternative to conventional CdS as electron transport layers. Additionally, machine learning analysis identified acceptor density and defect density as the most critical determinants of device efficiency, achieving a predictive accuracy of 99.39% when correlating input parameters to photovoltaic performance. These findings highlight the potential of Sr_3PBr_3 absorbers and Cu-based HTLs for fabricating stable, high-efficiency, lead-free perovskite solar cells, offering a viable route toward environmentally friendly photovoltaic technologies. It should be noted that this study based on simulations supported by machine learning, reports theoretical upper limits under idealized conditions; experimental validation is therefore essential, as practical devices will face efficiency reductions from optical, recombination, and stability losses. In addition, stability assessment and large-area scalability remain key directions for future research.

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Conflicts of interest

There are no conflicts to declare.

Data availability

All data and code are available in this link: https://doi.org/10.5281/zenodo.17244830

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Data Availability Statement

All data and code are available in this link: https://doi.org/10.5281/zenodo.17244830