

Cite this: *J. Mater. Chem. C*,
2026, 14, 1009Received 14th September 2025,
Accepted 10th December 2025

DOI: 10.1039/d5tc03428k

rsc.li/materials-c

Atomic-layer precision etching of SiO₂ using sequential molecular adsorption and plasma activation

Rakshith Venugopal,^a Nian Ran,^{*b} Robert Blick,^{id ac} Robert Zierold^{id *a} and Jun Peng^{id *a}

As device architectures in electronics, photonics, and quantum technologies reach atomic dimensions, precise and controllable material processing becomes essential. However, achieving atomic-layer precision in materials etching, even in silicon dioxide (SiO₂), remains a major challenge for next-generation nanofabrication. Here, we present a cyclic process that integrates sequential sulfur hexafluoride (SF₆) molecular adsorption with argon (Ar) plasma activation, enabling a stable etch-per-cycle (EPC) of ~1.4 Å per cycle and 100% synergy between modification and removal steps. Mechanistic studies combining experiments, *ab initio* molecular dynamics, and density functional theory reveal that etching proceeds *via* a combination of reversible physisorption and defect-mediated chemisorption. Moreover, detailed morphology characterization over multiple cycles reveals a directional and uniform etching effect. This work introduces a scalable, contamination-free, precise etching strategy using standard reactive ion etching (RIE) equipment and commercially available gases, offering a robust and transferable platform for next-generation nanofabrication.

1. Introduction

Atomic layer processing technologies, most notably atomic layer deposition (ALD) and atomic layer etching (ALE), have emerged as key enablers for continued scaling of device architectures in electronics, photonics, and quantum technologies.^{1–4} These techniques offer true atomic-level control over material growth and removal, facilitating extreme dimensional precision, high aspect ratio features, and compatibility with three-dimensional integration. Their utility has extended into

the fabrication of next-generation nanoelectronic systems, including quantum devices,^{5–9} all while remaining compatible with wafer-scale processing and high-volume manufacturing.^{10–12} The concept of ALE was first introduced in a 1988 patent describing the cyclic removal of atomic layers from crystalline diamond *via* alternating nitrogen dioxide exposure and inert gas ion bombardment.¹³ At that time, limited demand for such precision constrained further development. However, as Moore's law approaches its physical limits, ALE has gained renewed interest, particularly in advanced nanofabrication at the 10 nm technology node and beyond, where atomic-scale controllability, uniformity, and low damage are essential for enabling high-density integrated circuits, 3D NAND flash memory, and next-generation devices that require highly precise fabrication.^{11,14} In 2015, Lee and George¹⁵ reported the first isotropic thermal ALE of Al₂O₃, inspiring further exploration of atomic layer level precision etching techniques. Although etching of SiO₂ with atomic layer precision, one of the most basic semiconductor materials, has been tested by different etching strategies, it still remains challenging to keep control of accuracy and precision.¹⁶ Strategies using fluorocarbons (PFCs) or hydrofluorocarbons (HFCs) such as C₄F₈, CHF₃, and CF₄ have been leading the way. PFCs or HFCs are used to modify the SiO₂ surface to form a thin carbon-carbon film and subsequently Ar plasma bombardment is performed to remove the carbon-carbon polymer layer together with the SiO₂ thin layer. However, these etching approaches suffer from several drawbacks: polymer accumulation can destabilize the etch rate and uniformity and residual carbon contamination can affect both the processing chamber and device surfaces. These issues likely contribute to the significant variability in reported EPC values, ranging from 1.9 to 27 Å per cycle (Table S1). Consequently, other approaches, such as using trimethylaluminum and HF,^{17,18} plasma-involved processes with mixed precursors, or infrared-activated etching,¹⁹ are being actively pursued.

In this context, SF₆ is a promising alternative due to its clean decomposition chemistry and the absence of polymer-forming

^a Center for Hybrid Nanostructures, University of Hamburg, Luruper Chaussee 149, 22607 Hamburg, Germany. E-mail: jun.peng@uni-hamburg.de, robert.zierold@uni-hamburg.de

^b State Key Laboratory of High-Performance Ceramics and Superfine Microstructures, Shanghai Institute of Ceramics, Chinese Academy of Sciences, 200050 Shanghai, China. E-mail: rannian@mail.sic.ac.cn

^c DESY Photon Science Deutsches Elektronen-Synchrotron DESY, 22607 Hamburg, Germany



byproducts. In conventional etching processes, the use of SF₆ primarily relies on its plasma-phase dissociation into highly reactive fluorine radicals, which enables extremely fast etching rates, up to 1 μm min⁻¹,^{20–24} which introduce challenges in process control for precise etching. Here, we report a self-limiting, highly synergistic etching process for SiO₂ using sequential SF₆ gas exposure and Ar plasma etching near room temperature. Our process achieves a stable etch per cycle of ~1.4 Å per cycle, with 100% synergy confirmed by decoupled half-cycle tests. Through systematic parametric studies, we identified both a temperature window and an ICP power window characteristic of ALE-like behavior. Mechanistic insights, supported by *ab initio* molecular dynamics and density functional theory calculations, reveal that the adsorption process in the etching proceeds *via* a combination of reversible physisorption and defect-mediated chemisorption. Notably, the process exhibits excellent directionality and maintains high uniformity across large areas, leveraging commonly available gases and existing commercial RIE equipment. This work introduces a practical, scalable, contamination-free, and potentially SF₆-recyclable etching strategy, which can be an important step toward precise nanofabrication in next-generation semiconductor and photonic devices.

2. Experimental

2.1. Etching and characterization

Si wafers with 285 nm wet thermally grown SiO₂ (SIEGERT WAFER GmbH) were cut into small pieces of dimensions approximately 1 cm × 1 cm or 4.5 cm × 4.5 cm for most etching experiments. In addition, SiO₂ layers deposited by ALD and plasma-enhanced chemical vapor deposition (PECVD) on Si substrates (SIEGERT WAFER GmbH) were etched as control samples. Details of the respective deposition processes can be found in our previous work.^{25,26} Etching was carried out in a commercial ICP-RIE system (SENTECH SI 500). The film thickness was measured *ex situ* using an ellipsometer (SENTECH SENpro), with thickness values extracted using a Cauchy model for SiO₂. The film thickness was measured before and after etching. For pillar patterning, a positive electron beam resist (ARP661.09) was applied, and features were defined using an electron beam lithography system (Raith). The resulting sample was then subjected to a deep ion etching in the RIE system: at a continuous flow of 60 sccm SF₆, a plasma was generated with an ICP power of 300 W and a radio frequency (RF) bias of 60 W. The etching was carried out for 70 seconds. Then, the samples were cleaned with an Ar plasma at an ICP power of 300 W for 120 seconds in the same RIE system. The surface morphology was characterized using scanning electron microscopy (SEM, Zeiss Crossbeam 550), and surface roughness was measured by atomic force microscopy (AFM, Dimension).

2.2. MD calculations

Ab initio molecular dynamics (AIMD) simulations were performed using a non-spin-polarized configuration, an electronic

energy convergence cut-off of 10⁻⁴ eV, a Γ -centered 1 × 1 × 1 *k*-point mesh, and a time step of 3 fs. The Perdew–Burke–Ernzerhof (PBE) functional within the generalized gradient approximation (GGA) was employed. Following structural relaxation, AIMD simulations were conducted at 300 K in the NVT ensemble using a Nosé–Hoover thermostat for a total simulation time of 30 ps.

2.3. DFT calculations

All DFT calculations were carried out using the projector augmented wave (PAW) method^{27,28} in the Vienna *Ab initio* Simulation Package (VASP).²⁹ The exchange–correlation functional was described using GGA with the PBE parameterization. Convergence criteria were set to 10⁻⁵ eV for total energy and -0.05 eV Å⁻¹ for forces. The plane-wave cutoff energy was set to 520 eV. The PAW pseudopotentials were VRHFIN = Si: s²p², VRHFIN = O: s²p⁴, and VRHFIN = F: s²p⁵. Gaussian smearing (ISMEAR = 0) with a smearing width SIGMA = 0.05 eV was applied. The Brillouin zone was sampled using a 2 × 2 × 1 Monkhorst–Pack grid for calculations involving surface defects on the SiO₂ (001) surface. For electronic structure analysis, a denser 4 × 4 × 1 *k*-point grid was applied. The initial SiO₂ unit cell (*a* = *b* = 4.91656 Å, *c* = 5.43163 Å, space group *P*3₂21, no. 154, and mp-6930) was constructed according to a Materials Project database.³⁰ The detailed computational model is shown in Fig. S1. Long-range van der Waals interactions were accounted for using Grimme's DFT-D2 dispersion correction.

3. Results and discussion

3.1. Etching process and synergy characterization

The standard etching procedure follows a cyclic mode analogous to conventional ALE, as illustrated in Fig. 1a. Each cycle comprises four steps: (i) SF₆ exposure: SF₆ gas is introduced into the reaction chamber at a flow rate of 20 sccm for 5 seconds, allowing a full modification on the surface. (ii) Purge: the SF₆ gas supply is stopped and the chamber is purged for 30 seconds. (iii) Ar plasma etching: an inductively coupled plasma (ICP) of Ar is ignited at 100 W for 60 seconds to perform the etching. (iv) Cleaning purge: the ICP is turned off, and the chamber is purged again for 30 seconds. In each cycle, steps (i) and (ii) constitute the modification half-reaction, while steps (iii) and (iv) constitute the removal half-reaction. This alternating half-cycle process yields a stable EPC of ~1.4 Å per cycle (Fig. 1b). Throughout the process, a continuous Ar flow of 100 sccm is maintained as a carrier gas, while the reactor is kept at room temperature and a working pressure of 1 Pa. Comparable EPC values, as shown in Fig. S2, were obtained for SiO₂ deposited by ALD and PECVD using the same ALE recipe, indicating that the etching is predominantly surface-reaction-driven, with internal material defects exerting a negligible influence. Fig. 1c and Table S1 present a comparison of EPC values with previously reported ALE studies on SiO₂ over the past decade,^{19,31–42} highlighting the superior precision of our “sequential SF₆ gas and Ar plasma” etching strategy. Notably,



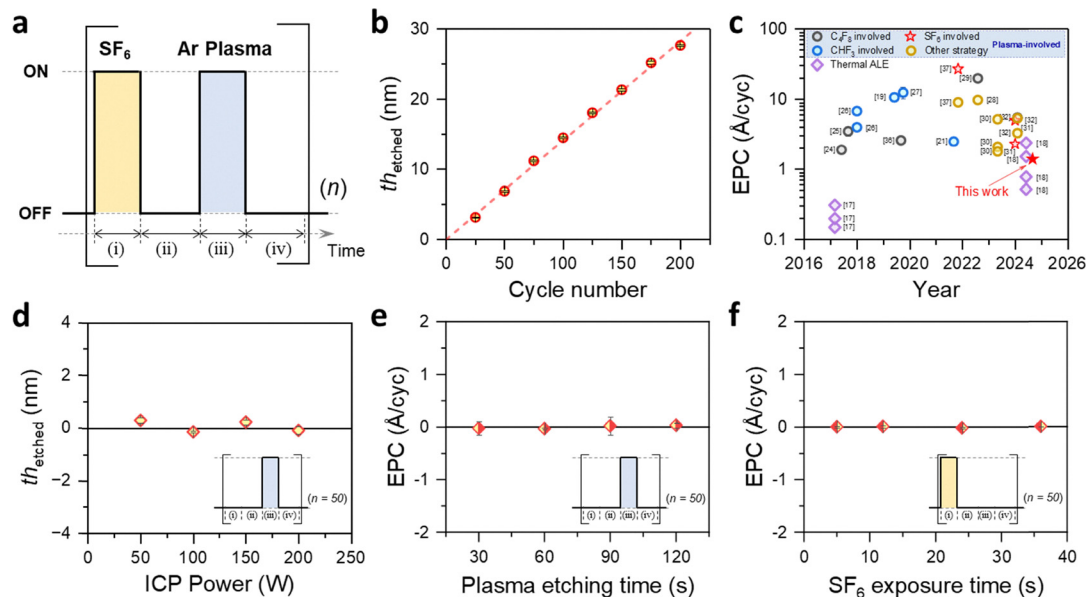


Fig. 1 Process configuration and synergy validation. (a) Schematic of the cyclic etching process comprising n sequential cycles. Each cycle includes four steps: (i) SF_6 exposure, (ii) purge, (iii) Ar plasma etching and (iv) cleaning purge. (b) The EPC determined to be $\sim 1.4 \text{ \AA}$ per cycle based on a linear fitting with $R^2 \approx 0.999$ of etched thickness, th_{etched} , versus cycle number of the standard loop recipe. (c) Comparison of the EPC from this work with the state-of-the-art SiO_2 ALE in the past decade. The numbers adjacent to data points correspond to literature references. (d) A blank period replaces step (i) in the standard recipe. The ICP power is changed in step (iii), indicating that Ar plasma cannot etch SiO_2 by sputtering within this power range. (e) and (f) Synergy validation by replacing either step (i) or (iii) in the standard process with a blank period. No etching is observed when either half-reaction is applied independently, confirming that etching only occurs when both steps are combined, indicative of 100% synergy.

similar EPC values were reproduced using the same recipe on a second, comparable RIE system, demonstrating the reproducibility and robustness of the method. Our process delivers EPC performance on a par with thermal ALE but without the need for specialized equipment.

To gain a better understanding of the etching mechanism, control studies were performed by omitting either step (i) or (iii) from the standard etching cycle. First, the SF_6 gas input in step (i) was withdrawn, and the ICP power range in step (iii) was widened. No significant etching was observed (Fig. 1d). This observation indicates that material removal in the standard process does not originate from Ar plasma-induced physical sputtering alone. Rather, it results from the synergistic interaction between the two half-reactions. In the field of ALE, the degree of synergy is quantitatively evaluated using the ALE synergy factor, S , defined as

$$S = (\text{EPC} - (\alpha + \beta)) / \text{EPC} \times 100\%.$$

Here, α and β denote the apparent etch contributions arising from the isolated modification (steps i + ii) and removal (steps iii + iv) half-cycles, respectively, which correspond to parasitic etching outside the ideal self-limiting ALE regime.^{10,43} As shown in Fig. 1e and f, repeating either half-reaction alone produces negligible etching, confirming that both α and β are effectively zero. Consequently, the synergy factor is calculated to be $S = 100\%$, indicating a fully synergistic process. This value exceeds that of previous ALE reports employing alternating fluorocarbon plasma and Ar ion bombardment, which achieved $\sim 80\%$ synergy in 10 nm logic device fabrication.⁴³ Note that the cyclic configuration that separates SF_6 exposure and

Ar plasma etching steps is critical for maintaining the precise etching behavior. As shown by the tests in Table S2, mixing the SF_6 and the Ar during the plasma step, either cyclically or continuously, yields significantly higher EPC values, resembling a reactive ion etching (RIE) process.

3.2. Etching characteristics

Comparative analysis of various etching strategies suggests that our “sequential SF_6 gas and Ar plasma” method follows strictly neither RIE nor ALE characteristics but is closer to a combination of both. To validate this assumption, a series of systematic experiments were conducted to assess the process behavior at varying parameters. Fig. 2a shows the effect of wafer holder temperature on the etching rate. Between room temperature and approximately $40 \text{ }^\circ\text{C}$, the EPC remains stable, defining a distinct temperature window for optimal etching. Beyond this range, up to $160 \text{ }^\circ\text{C}$, the EPC gradually decreases, which is likely due to increased thermal energy causing SF_6 molecules to desorb more easily from the surface, thereby reducing their availability for effective surface modification. In addition to the temperature window, a working pressure window was also identified (Fig. 2b). As the pressure increased to 0.8 Pa , the EPC increased to 1.4 \AA per cycle sharply; then, until 8 Pa , the EPC remained at a plateau. However, with further increases in pressure, the EPC begins to rise again. This behavior may be attributed to the insufficient purge time at higher pressures, allowing excessive SF_6 molecules to persist in the chamber. Note that there is no sputtering etching within the testing pressure if the step (i) SF_6 exposure is absent (Fig. S3). Upon plasma ignition, these



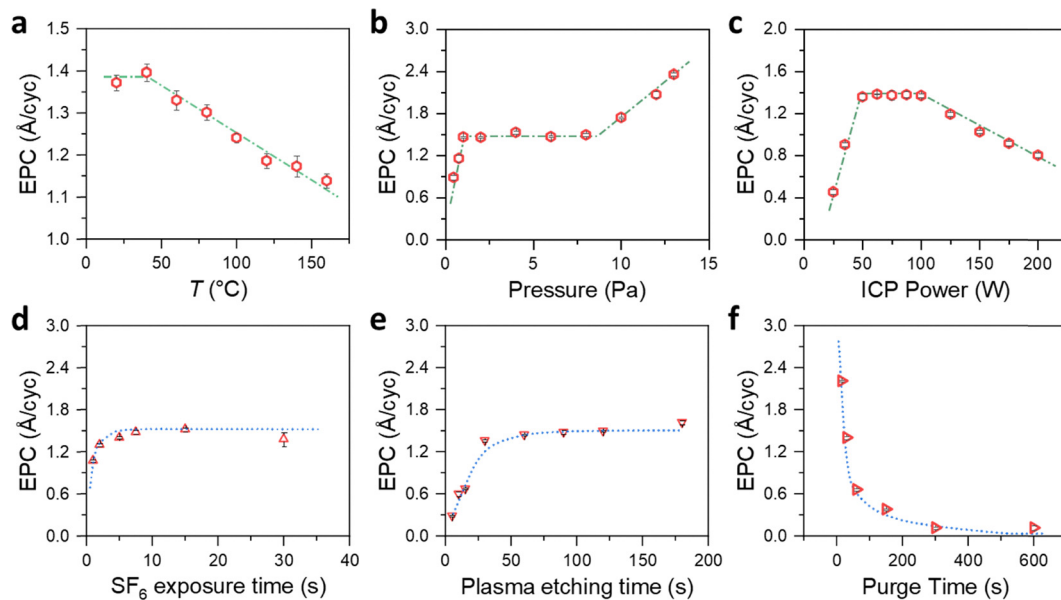


Fig. 2 Etching characteristics under single-parameter variations. (a)–(c) Process windows are identified in which the EPC remains stable when varying (a) wafer holder temperature, (b) chamber pressure, and (c) ICP power. These plateaus indicate the self-limiting behavior. As the (d) SF_6 exposure time and (e) the plasma etching duration increase, the EPC gradually approached saturation. The phenomena in (d) and (e) exhibit conventional ALE-like self-limitation characteristics. (f) However, as the purge time in step (ii) prolongs, the EPC drops dramatically and then approaches a no-etching state gradually, suggesting physical absorption of SF_6 in step (i).

residuals generate additional F radicals for etching. Fig. 2c presents the EPC dependence on ICP power. In general, higher ICP power results in increased gas ionization and ion density (Fig. S4). However, the EPC curve, which first increases and then decreases with increasing ICP power, mirrors the trend of Ar plasma ion energy as measured on identical equipment (Fig. S4). This correlation suggests that ion energy, rather than ion density, is the dominant factor influencing etching in the tested 0–700 W range. The relatively flat EPC result within the 50–100 W range likely corresponds to an “etching window,” where the incident plasma ion energy is sufficient to remove the SF_6 -modified surface layer without damaging the underlying SiO_2 . Below this energy threshold, incomplete removal of the modified surface leads to reduced EPC. As shown in Fig. 2d and e, as the SF_6 exposure time in step (i) and the plasma etching time in step (ii) increase—corresponding to an increase in dose—the EPC saturates at an EPC of ~ 1.4 Å per cycle. The above behaviors are consistent with the self-limiting saturation characteristic expected in ALE processes. However, when the purge time following the SF_6 exposure (step (ii)) is extended, the EPC does not remain constant as expected. Instead, it drops quickly and ultimately approaches zero etching, as shown in Fig. 2f. This result indicates that the modification half-reaction is reversible, deviating from the typical ALE working principle derived from the ALD concept, which assumes irreversible surface reactions.⁴⁴ Furthermore, the shape of the EPC–purge time curve implies that the adsorption behavior in the modification step is governed predominantly by reversible type I physisorption,⁴⁵ rather than by chemisorption.

However, an etching process, whose modification half-reaction is dominated by reversible saturated physisorption, exhibits an

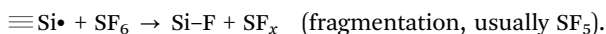
ICP power window (Fig. 2c) comparable to the temperature window typically observed in ALE/ALD processes governed by chemisorption. This phenomenon is counterintuitive. To clarify this behavior, simplified SiO_2 surface models terminated with either Si or O atoms (Fig. S1) were constructed to examine the adsorption behavior of SF_6 . *Ab initio* molecular dynamics (AIMD) simulations (Fig. S5 and Video S1, S2) reveal that Si-terminated surfaces can chemisorb and partially dissociate SF_6 molecules, whereas O-terminated surfaces remain inert and exhibit negligible adsorption. DFT calculations (Figure S6) further show that SF_6 adsorption on Si-terminated surfaces has an adsorption energy E_{ads} of -5.98 eV, indicative of strong chemisorption, with F atoms forming stable Si–F bonds with exposed Si atoms. In contrast, adsorption on O-terminated surfaces exhibits a much weaker E_{ads} equal to -0.22 eV, without the formation of stable chemical bonds. These results confirm that the modification half-reaction involves a combination of chemisorption and physisorption, consistent with experimental observations: the ICP power window (Fig. 2c) reflects chemisorption behavior, while the EPC–purge time dependence (Fig. 2f) represents physisorption dynamics. Moreover, an etching process using only pulsed SF_6 plasma (Table S3) yielded an EPC of 0.5 Å per cycle, approximately one-third of that of the standard ALE process, indicating that chemisorption contributes only partially to the overall adsorption. This demonstrates that, although our process achieves atomic-level etching precision, it does not conform to a purely conventional ALE mechanism.

3.3. Etching mechanism exploration

Next, the SiO_2 surface defect sites as potential active adsorption centers for SF_6 are explored. At room temperature, the SiO_2



surface hosts various stable or metastable defects, with the most prevalent being silicon dangling bonds ($\equiv\text{Si}\cdot$), oxygen-centered radicals ($\equiv\text{Si-O}\cdot$), doubly coordinated silicon radicals ($\equiv\text{Si}\cdot$), and strained siloxane rings (Si_2O_2).^{46–48} Density functional theory (DFT) simulation was used to construct and relax these defect structures under vacuum at room temperature. Among them, three were found to be thermodynamically stable: $\equiv\text{Si}\cdot$, $\equiv\text{Si-O}\cdot$ and Si_2O_2 rings, as shown in Fig. 3a. F atoms were sequentially introduced to these defect sites, and the corresponding adsorption energy E_{ads} was calculated. All three sites exhibit negative E_{ads} values, indicating spontaneous and energetically favorable formation of Si–F bonds. In the etching cycles, F atoms can originate from two potential primary sources: first, the spontaneous dissociation of SF_6 at defect sites during step (i), and second, the plasma-induced dissociation of physically adsorbed SF_6 during step (iii). For the former origin, a thermodynamically favorable chemical reaction occurs only if E_{ads} exceeds the dissociation energy E_{diss} of SF_6 into SF_5 and F, which is calculated to be 3.03 eV. For example, the chemisorption at $\equiv\text{Si}\cdot$ defect sites can be expressed as



As shown in Fig. 3b, among the studied defects, only the Si_2O_2 ring possesses a sufficiently negative E_{ads} to overcome E_{diss} , thereby providing enough energy and enabling spontaneous dissociation of SF_6 and subsequent F adsorption. The other defects exhibit insufficient E_{ads} , which aligns with the experimental observation that the overall surface behavior is dominated by physisorption (Fig. 2f). For the latter sources, plasma-activated Ar^+ ions dissociate physically adsorbed SF_6 , generating free F radicals, which can spontaneously react with defect sites to form stable Si–F bonds. Crystal orbital Hamilton population (COHP) analysis⁴⁹ was used to compare the related bond strengths E_{b} as a function of F adsorption. As shown in Fig. 3c, after adsorbing F atoms, the Si–O bond strengths $E_{\text{b,Si-O}}$ at the defect sites decrease. The Si–O bonds at $\equiv\text{Si-O}\cdot$, Si_2O_2 rings with single F adsorption, and $\equiv\text{Si}\cdot$ with two F adsorptions exhibit lower Si–O bond strength $E_{\text{b,Si-O}}$ than the Si–O bond strength in defect-free regions $E_{\text{b,Si-O,normal}}$. This implies that, under Ar plasma bombardment, the weaker Si–O bonds near Si–F sites are more preferentially broken than the Si–O bond in defect-free regions. Therefore, when the incident Ar ion energy is larger than $E_{\text{b,Si-O}}$ but smaller than $E_{\text{b,Si-O,normal}}$, the etching is highly controllable and only limited to the surface,

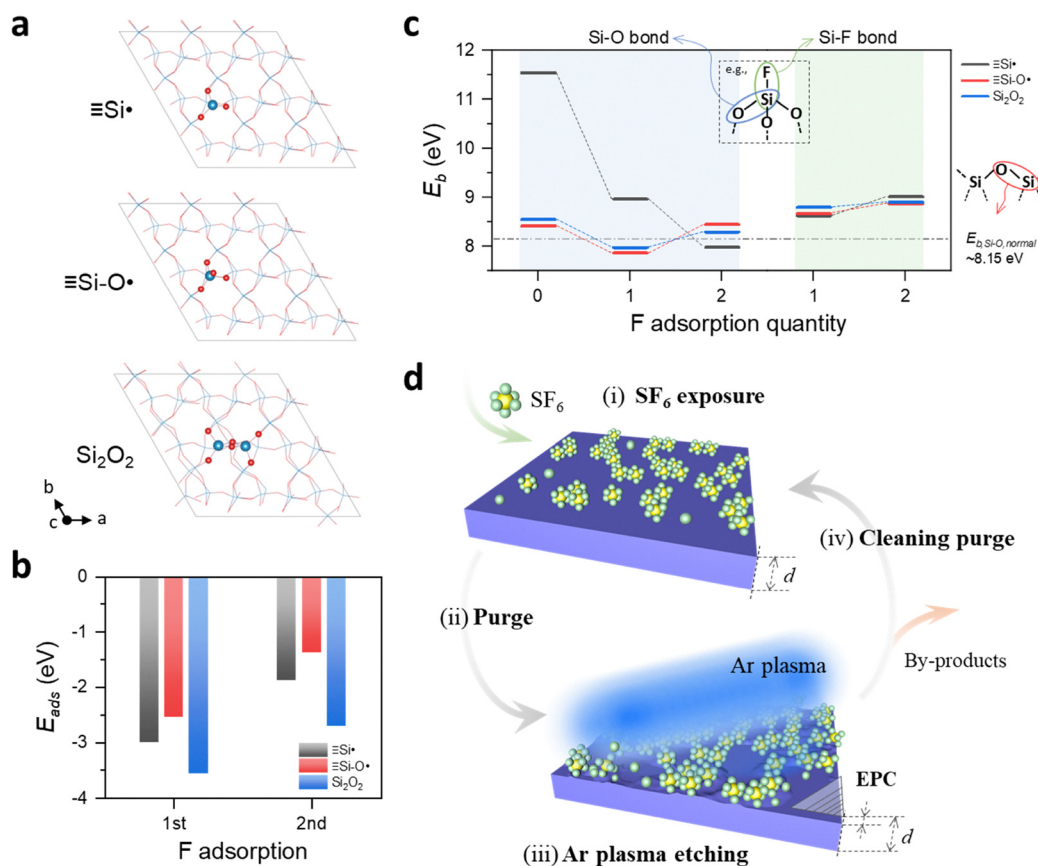


Fig. 3 Mechanistic insights into the etching process. (a) Three dominant surface defects on the SiO_2 surface. (b) The adsorption energy E_{ads} for the different surface defects of F atoms. All negative E_{ads} values indicate that the reaction occurs spontaneously. (c) Comparison of bonding strength, E_{b} , for Si–O and Si–F with the F adsorption quantity. The black dashed line represents the Si–O bond strength in the defect-free regime on the surface, $E_{\text{b,Si-O,normal}}$. Insets: A $\equiv\text{Si}\cdot$ defect with one F adsorbed as an example of Si–O and Si–F studied; Si–O bond in the defect-free surface regime. (d) Schematic of the proposed etching mechanism using sequential SF_6 gas exposure and pulsed Ar plasma. Unlike conventional ALE, this process combines both chemical adsorption at reactive defect sites and reversible physical adsorption in the step (i) SF_6 exposure, contributing to etching during step (iv) Ar plasma.



similar to pure ALE processes. This supports the existence of an “etching window” (Fig. 2c), in which the energy of incident Ar^+ ions falls between $E_{\text{b,Si-O}}$ and $E_{\text{b,Si-O,normal}}$ —sufficient to break the Si–O bonds that are modified and weakened by Si–F bonds but insufficient to sputter the unmodified and defect-free surface regions. Note that, when the ion energy exceeds $E_{\text{b,Si-O,normal}}$, the surface will be etched by physical sputtering.

Based on these findings, the etching mechanism for this four-step cyclic sequence is proposed in Fig. 3d. In step (i) SF_6 exposure, most SF_6 molecules are physically adsorbed onto the SiO_2 surface, and only a small amount of SF_6 molecules can be dissociated at high-reactivity defect sites, *e.g.*, Si_2O_2 rings, forming Si–F bonds. None of the adsorption produces an etching effect. And the physical adsorption remains the dominant interaction. In step (ii) Purge, excess SF_6 is removed, and the adsorbed SF_6 amount on the surface can be tuned by adjusting the purge duration or chamber pressure. In step (iii) Ar plasma etching, the plasma dissociates physisorbed SF_6 and F radicals are produced, which react with surface defects to form Si–F bonds. Simultaneously, Ar^+ ions provide energy to cleave weakened Si–O bonds adjacent to the Si–F bonds. The fluorinated Si atoms then further react with radicals to form volatile byproducts.^{50,51} Finally, the volatile by-products are evacuated in step (iv) cleaning purge, preparing the surface for the next cycle. It is worth noting that in most conventional etching schemes utilizing SF_6 , the entire gas flow is converted into plasma, and fluorine radicals generated through

ionization or fragmentation act as the primary etching species. In contrast, in our process, only the absorbed SF_6 molecules are fragmented to participate in etching, while the unadsorbed fraction remains intact and is purged as stable SF_6 molecules during the modification half-reaction. These unreacted molecules can potentially be captured, compressed, and recycled, significantly reducing process emissions. Although the use of SF_6 , a potent greenhouse gas, cannot be entirely eliminated, our etching approach achieves lower gas consumption and improved environmental compatibility, thereby offering a more sustainable pathway for precision plasma etching.^{52–54}

3.4. Directional etching

Validating the directionality—directional or isotropic—of this etching process determines how this process can be used in advanced patterning applications. To assess this, a SiO_2 substrate patterned with cylindrical pillars of diameter 600 nm was used. Optical microscopy images (Fig. 4a) were captured after various etching cycles. The observed color variations in the optical images reflect changes in the SiO_2 film thickness due to material removal. The corresponding atomic force microscopy (AFM) images (Fig. 4b) and detailed interface profile (Fig. 4c) show, however, no significant difference after each etching process. These results reveal that the surface morphology of the pillars remains essentially unchanged throughout the etching cycles. Specifically, the sidewalls of the pillars retain their shape, indicating that the etching is anisotropic and proceeds

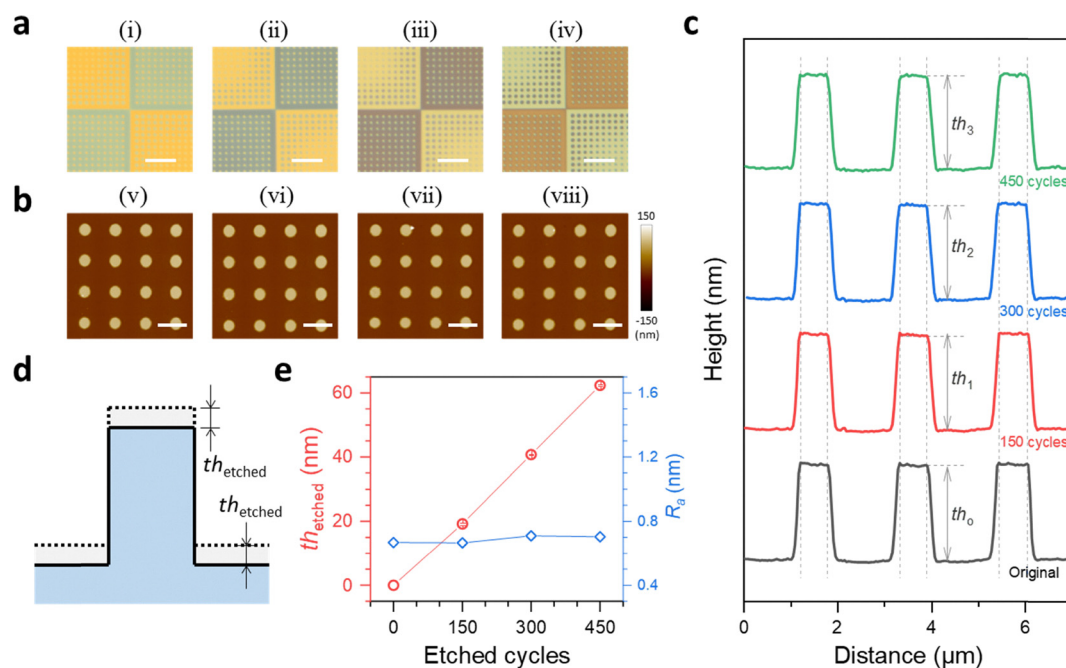


Fig. 4 Directional etching. Pillar structures were used to evaluate the anisotropy of the etching process. (a) Optical microscopy images and (b) the corresponding AFM images of the original pillar sample (i) and (v), the same position after 150 etching cycles (ii) and (vi), the same position after 300 etching cycles (iii) and (vii), and the same position after 450 etching cycles (iv) and (viii). The color variation is due to the different thicknesses of SiO_2 after etching. The scalebars for (a) and (b) are 10 μm and 2 μm , respectively. (c) Height profiles of three consecutive pillars in their original state and after 150, 300, and 450 etching cycles. The corresponding pillar heights th_0 , th_1 , th_2 , and th_3 are 91.4 ± 1.17 , 91.3 ± 0.86 , 90.7 ± 0.91 , 89.6 ± 1.00 nm, respectively, which remain the same during the etching. (d) Sketch of the directional etching result. (e) Corresponding etched thickness th_{etched} and roughness R_a during the test as a function of cycle number. The stable and low R_a suggests a damage-free and uniform etching process.



primarily in the vertical direction (Fig. 4d). In contrast, isotropic etching should reduce both the height and diameter of the pillars uniformly, as sketched in Fig. S8. Quantitatively, the total etched depth was approximately 62 nm (Fig. 4e) corresponding to the determined 1.4 Å per cycle, while the pillar diameter remained constant. If the process would have been isotropic, the pillar diameter would have decreased by an estimated 27%. This preservation of lateral dimensions strongly supports the directional nature of the etch. The observed anisotropy is likely driven by a bias voltage generated unintentionally during Ar plasma pulses, potentially due to the self-bias effect and capacitive coupling from the ICP power supply. Although no deliberate RF bias was applied, a passive bias was detected during plasma ignition (Fig. S9). This bias establishes an electric field in the plasma sheath, accelerating charged particles, primarily Ar⁺ ions, perpendicularly toward the substrate. As a result, the etching becomes angle-selective: ions arriving at high incident angles exhibit lower kinetic energy, making them less likely to erode sidewalls or lateral features.^{55,56} Additionally, the low operating pressure ensures a large mean free path, further promoting unidirectional ion trajectories and reinforcing vertical etch selectivity. This directional etching behavior was also consistently observed in hole-patterned samples (Fig. S10), indicating excellent repeatability in different nanostructured geometries. Surface characterization after etching confirmed the gentle nature of the process, with an unchanged post-etch surface roughness (R_a) of approximately 0.7 nm. Moreover, large-area uniformity tests on a 4.5 cm × 4.5 cm wafer showed an etching uniformity of 3.96% after 150 cycles (Fig. S11). It is well below the ±5% threshold typically required by the semiconductor industry for dry etching processes and is comparable to values reported in other ALE studies,^{57–59} highlighting the potential of this process to achieve high-fidelity, uniform etching over wafer-scale dimensions.

4. Conclusions

In summary, we have developed a highly precise and reproducible etching process for SiO₂ by combining sequential SF₆ gas exposure with Ar plasma at room temperature. The process achieves a stable EPC of ~1.4 Å per cycle, with 100% synergy confirmed between the modification and removal half-reactions. Parameter studies revealed well-defined process windows for temperature and ICP power, while saturation behavior and physisorption-dominated reversibility were confirmed through dose-dependent experiments. Mechanistic insights obtained *via* AIMD and DFT simulations highlight the role of surface defects as active sites for SF₆ dissociation and F adsorption. These findings support a model of the etching process involving both reversible physical adsorption and defect-mediated chemical adsorption. Moreover, directional etching was experimentally verified using pillar and hole structures. Uniformity testing over a 4.5 cm × 4.5 cm area demonstrated high reproducibility with only 3.96% variation after 150 cycles. Importantly, the process relies on commercial RIE equipment and widely available gases, making it scalable and

industrially compatible. Although demonstrated here for SiO₂, the concept could be extended to other materials that exhibit selective reactivity between neutral gas-phase species and plasma-generated radicals, opening pathways for broader applications in nanoscale fabrication.

Author contributions

J. P. proposed the concepts and designed the experiments; V. R. and J. P. performed the experiments and analysis; R. N. performed the MD and DFT calculation; R. Z. supervised the study; R. B. provided infrastructure to conduct the experiments; J. P., V. R., and R. Z. co-wrote the manuscript. All authors discussed the experimental and theoretical results and commented on the manuscript. All authors have approved the final version of the manuscript.

Conflicts of interest

There are no conflicts to declare.

Data availability

The data supporting the findings of this work are available within the article and its supplementary information (SI). Supplementary information is available. See DOI: <https://doi.org/10.1039/d5tc03428k>.

All other relevant data supporting the findings of this study are available from the corresponding author on request.

Acknowledgements

This work was funded by the Deutsche Forschungsgemeinschaft (DFG, German Research Foundation) – Projektnummer 192346071 – SFB 986 “Tailor-Made Multi-Scale Materials Systems”. The authors acknowledge the use of the Hamburg-CRR cleanroom facility, jointly operated by the University of Hamburg (UHH), DESY, and the Max Planck Institute for the Structure and Dynamics of Matter (MPSD), which was instrumental to the experimental work presented in this study. We also acknowledge support from the BMBF ForLab initiative. Nian Ran acknowledges funding from the National Natural Science Foundation of China (NSFC, grant no. 22403103) and the Shanghai Sailing Program (grant no. 23YF1454900). The authors thank Dr Paul Plate, Dr Michael Höfner, and Dr Marcel Schulze (SENTECH Instruments GmbH) for valuable discussions and for providing Ar plasma characterization data used in this work.

References

- 1 V. Cremers, R. L. Puurunen and J. Dendooven, *Appl. Phys. Rev.*, 2019, **6**, 021302.
- 2 A. Fischer, A. Routzahn, S. M. George and T. Lill, *J. Vac. Sci. Technol., A*, 2021, **39**, 030801.



- 3 T. Faraz, F. Roozeboom, H. C. M. Knoop and W. M. M. Kessels, *ECS J. Solid State Sci. Technol.*, 2015, **4**, N5023.
- 4 J. Peng and R. Zierold, in *Encyclopedia of Condensed Matter Physics*, ed. T. Chakraborty, Academic Press, Oxford, 2nd edn, 2024, pp. 716–728, DOI: [10.1016/B978-0-323-90800-9.00206-7](https://doi.org/10.1016/B978-0-323-90800-9.00206-7).
- 5 N. P. de Leon, K. M. Itoh, D. Kim, K. K. Mehta, T. E. Northup, H. Paik, B. S. Palmer, N. Samarth, S. Sangtawesin and D. W. Steuerman, *Science*, 2021, **372**, eabb2823.
- 6 W. Lu, Y. Lee, J. C. Gertsch, J. A. Murdzek, A. S. Cavanagh, L. Kong, J. A. del Alamo and S. M. George, *Nano Lett.*, 2019, **19**, 5159–5166.
- 7 A. P. M. Place, L. V. H. Rodgers, P. Mundada, B. M. Smitham, M. Fitzpatrick, Z. Leng, A. Premkumar, J. Bryon, A. Vrajitoarea, S. Sussman, G. Cheng, T. Madhavan, H. K. Babla, X. H. Le, Y. Gang, B. Jäck, A. Gyenis, N. Yao, R. J. Cava, N. P. de Leon and A. A. Houck, *Nat. Commun.*, 2021, **12**, 1779.
- 8 Z. Chen, I. M. Grace, S. L. Woltering, L. Chen, A. Gee, J. Baugh, G. A. D. Briggs, L. Bogani, J. A. Mol, C. J. Lambert, H. L. Anderson and J. O. Thomas, *Nat. Nanotechnol.*, 2024, **19**, 986–992.
- 9 C. Y. Chen, Z. Sun, R. Torsi, K. Wang, J. Kachian, B. Liu, G. B. Rayner, Z. Chen, J. Appenzeller, Y.-C. Lin and J. A. Robinson, *Nat. Commun.*, 2024, **15**, 4016.
- 10 K. J. Kanarik, T. Lill, E. A. Hudson, S. Sriraman, S. Tan, J. Marks, V. Vahedi and R. A. Gottscho, *J. Vac. Sci. Technol., A*, 2015, **33**, 020802.
- 11 K. J. Kanarik, S. Tan and R. A. Gottscho, *J. Phys. Chem. Lett.*, 2018, **9**, 4814–4821.
- 12 X. Chen, H.-R. Park, M. Pelton, X. Piao, N. C. Lindquist, H. Im, Y. J. Kim, J. S. Ahn, K. J. Ahn, N. Park, D.-S. Kim and S.-H. Oh, *Nat. Commun.*, 2013, **4**, 2361.
- 13 M. N. Yoder, *US Pat.*, US4756794A, 1988.
- 14 T.-Y. Lee, P.-T. Chen, C.-C. Huang, H.-C. Chen, L.-Y. Chen, P.-T. Lee, F.-C. Chen, R.-H. Horng and H.-C. Kuo, *Nanoscale Adv.*, 2025, **7**, 2796–2817.
- 15 Y. Lee and S. M. George, *ACS Nano*, 2015, **9**, 2061–2070.
- 16 T. Lill, *Atomic layer processing: semiconductor dry etching technology*, 2021.
- 17 J. W. DuMont, A. E. Marquardt, A. M. Cano and S. M. George, *ACS Appl. Mater. Interfaces*, 2017, **9**, 10296–10307.
- 18 D. Catherall, A. Hossain and A. Minnich, arXiv, 2024, preprint, arXiv:2405.05491, DOI: [10.48550/arXiv.2405.05491](https://doi.org/10.48550/arXiv.2405.05491).
- 19 N. Miyoshi, K. Shinoda, H. Kobayashi, M. Kurihara, Y. Kouzuma and M. Izawa, *J. Vac. Sci. Technol., A*, 2021, **39**, 052601.
- 20 R. Dussart, T. Tillocher, L. Becerra, P. Lefauchaux and L. J. Overzet, *Jpn. J. Appl. Phys.*, 2025, **64**, 05SP01.
- 21 V. Bliznetsov, H. M. Lin, Y. J. Zhang and D. Johnson, *J. Micromech. Microeng.*, 2015, **25**, 087002.
- 22 S. F. Yoon, *Microelectron. Eng.*, 1991, **14**, 23–40.
- 23 R. Hsiao and J. Carr, *Mater. Sci. Eng., B*, 1998, **52**, 63–77.
- 24 X. Man, N. Bao, Y. Hao, Y. Feng and X. Ma, *Phys. Status Solidi A*, 2020, **217**, 2000223.
- 25 S. Haugg, C. Hedrich, R. H. Blick and R. Zierold, *Nanomaterials*, 2021, **11**, 3357.
- 26 C. Hedrich, D. Deduytsche, R. R. Petit, T. Krekeler, J. Peng, M. Ritter, J. Dendooven, C. Detavernier, R. H. Blick and R. Zierold, *Surf. Interfaces*, 2025, **57**, 105696.
- 27 J. P. Perdew, K. Burke and M. Ernzerhof, *Phys. Rev. Lett.*, 1996, **77**, 3865–3868.
- 28 P. E. Blöchl, *Phys. Rev. B:Condens. Matter Mater. Phys.*, 1994, **50**, 17953–17979.
- 29 G. Kresse and J. Furthmüller, *Phys. Rev. B:Condens. Matter Mater. Phys.*, 1996, **54**, 11169–11186.
- 30 A. Jain, S. P. Ong, G. Hautier, W. Chen, W. D. Richards, S. Dacek, S. Cholia, D. Gunter, D. Skinner, G. Ceder and K. A. Persson, *APL Mater.*, 2013, **1**, 011002.
- 31 S. S. Kaler, Q. Lou, V. M. Donnelly and D. J. Economou, *J. Phys. D: Appl. Phys.*, 2017, **50**, 234001.
- 32 R. J. Gasvoda, A. W. van de Steeg, R. Bhowmick, E. A. Hudson and S. Agarwal, *ACS Appl. Mater. Interfaces*, 2017, **9**, 31067–31075.
- 33 K. Koh, Y. Kim, C.-K. Kim and H. Chae, *J. Vac. Sci. Technol., A*, 2017, **36**, 01B106.
- 34 S. Dallorto, A. Goodyear, M. Cooke, J. E. Szornel, C. Ward, C. Kastl, A. Schwartzberg, I. W. Rangelow and S. Cabrini, *Plasma Processes Polym.*, 2019, **16**, 1900051.
- 35 S. Dallorto, M. Lorenzon, J. Szornel, A. Schwartzberg, A. Goodyear, M. Cooke, M. Hofmann, I. W. Rangelow and S. Cabrini, *J. Vac. Sci. Technol., B*, 2019, **37**, 051805.
- 36 S. Y. Kim, I.-S. Park and J. Ahn, *Appl. Surf. Sci.*, 2022, **589**, 153045.
- 37 Y. S. Lee, S. J. Kim, J. J. Lee, C. H. Cho, I. H. Seong and S. J. You, *J. Phys. D: Appl. Phys.*, 2022, **55**, 365203.
- 38 Y. Kim, H. Kang, C. Kim and H. Chae, *ACS Sustainable Chem. Eng.*, 2023, **11**, 6136–6142.
- 39 A. Osonio, T. Tsutsumi, B. Mukherjee, R. Borude, N. Kobayashi and M. Hori, *Jpn. J. Appl. Phys.*, 2023, **62**, 121001.
- 40 J. Kim, H. Kang, Y. Kim, M. Jeon and H. Chae, *Plasma Processes Polym.*, 2024, **21**, 2300216.
- 41 R. J. Gasvoda, Y. G. P. Verstappen, S. Wang, E. A. Hudson and S. Agarwal, *J. Vac. Sci. Technol., A*, 2019, **37**, 051003.
- 42 N. Miyoshi, H. Kobayashi, K. Shinoda, M. Kurihara, K. Kawamura, Y. Kouzuma and M. Izawa, *J. Vac. Sci. Technol., A*, 2021, **40**, 012601.
- 43 K. J. Kanarik, S. Tan, W. Yang, T. Kim, T. Lill, A. Kabansky, E. A. Hudson, T. Ohba, K. Nojiri, J. Yu, R. Wise, I. L. Berry, Y. Pan, J. Marks and R. A. Gottscho, *J. Vac. Sci. Technol., A*, 2017, **35**.
- 44 V. Miikkulainen, M. Leskelä, M. Ritala and R. L. Puurunen, *J. Appl. Phys.*, 2013, **113**, 021301.
- 45 K. S. W. Sing, *Pure Appl. Chem.*, 1985, **57**, 603–619.
- 46 A. D. Kulkarni, D. G. Truhlar, S. Goverapet Srinivasan, A. C. T. van Duin, P. Norman and T. E. Schwartzentruber, *J. Phys. Chem. C*, 2013, **117**, 258–269.
- 47 V. A. Radzig, in *Defects in SiO2 and Related Dielectrics: Science and Technology*, ed. G. Pacchioni, L. Skuja and D. L. Griscom, Springer, Netherlands, Dordrecht, 2000, pp. 339–370, DOI: [10.1007/978-94-010-0944-7_12](https://doi.org/10.1007/978-94-010-0944-7_12).
- 48 P. Norman and T. E. Schwartzentruber, Presented in part at the LAMMPS Users' Workshop and Symposium, 2011.



- 49 R. Dronskowski and P. E. Bloechl, *J. Phys. Chem.*, 1993, **97**, 8617–8624.
- 50 R. Knizikevičius, *Vacuum*, 2009, **83**, 953–957.
- 51 Y. A. Mankelevich, E. N. Voronina, T. V. Rakhimova, A. P. Palov, D. V. Lopaev, S. M. Zyryanov and M. R. Baklanov, *Eur. Phys. J. D*, 2017, **71**, 126.
- 52 H. Pedersen, S. T. Barry and J. Sundqvist, *J. Vac. Sci. Technol., A*, 2021, **39**, 051001.
- 53 S. Shi, Y. Li, Z. Cui, Y. Yan, X. Zhang, J. Tang and S. Xiao, *Chem. Eng. J.*, 2023, **470**, 144166.
- 54 G. S. Oehrlein, S. M. Brandstadter, R. L. Bruce, J. P. Chang, J. C. DeMott, V. M. Donnelly, R. Dussart, A. Fischer, R. A. Gottscho, S. Hamaguchi, M. Honda, M. Hori, K. Ishikawa, S. G. Jaloviar, K. J. Kanarik, K. Karahashi, A. Ko, H. Kothari, N. Kuboi, M. J. Kushner, T. Lill, P. Luan, A. Mesbah, E. Miller, S. Nath, Y. Ohya, M. Omura, C. Park, J. Poulouse, S. Rauf, M. Sekine, T. G. Smith, N. Stafford, T. Standaert and P. L. G. Ventzek, *J. Vac. Sci. Technol., B*, 2024, **42**, 041501.
- 55 M. Teichmann, J. Lorbeer, F. Frost and B. Rauschenbach, *Nanoscale Res. Lett.*, 2014, **9**, 439.
- 56 H. Winter, *Phys. Rep.*, 2002, **367**, 387–582.
- 57 K. J. Kanarik, T. Lill, E. A. Hudson, S. Sriraman, S. Tan, J. Marks, V. Vahedi and R. A. Gottscho, *J. Vac. Sci. Technol., A*, 2015, **33**, 020802.
- 58 M. Huff, *Micromachines*, 2021, **12**, 991.
- 59 X. Luo, Z. Zhao, Y. Wang, P. Gao, X. Ma, M. PU, X. Li and Y. Guo, *US Pat.*, US12272530B2, 2025.

