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# How solvent activity and solute solvation control electrode potentials: quantitative insights from Li<sup>+</sup> insertion into TiO<sub>2</sub>

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Tuning electrolyte salt concentration has emerged as a powerful strategy to enhance the performance of rechargeable batteries by mitigating parasitic side reactions. Here, we introduce a unified and quantitative validation of an electrolyte thermodynamic framework incorporated into a modified Nernst equation that explicitly includes electrolyte-specific parameters—namely, salt and solvent activities, as well as the salt hydration number—to rationalize and predict how electrode potentials vary with electrolyte concentration in reversible ion-insertion systems. We demonstrate this generic methodology for reversible Li<sup>+</sup> insertion into anatase TiO<sub>2</sub> using aqueous LiCl electrolytes. By combining tabulated water and LiCl activities with concentration-dependent hydration numbers derived from both experimental and theoretical approaches, we achieve quantitative agreement between predicted and measured potential shifts. This work establishes a reliable and broadly applicable framework for deciphering and controlling electrolyte-concentration effects in batteries and, more generally, in electrochemical systems involving concentrated electrolytes.

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## Introduction

Electrolyte composition is a key determinant of both the performance and lifetime of rechargeable batteries.<sup>1</sup> Not only does it govern ion transport, conductivity, but it also fundamentally controls the thermodynamics and kinetics of interfacial electrochemical processes. These interfacial processes include reversible ion insertion into host materials and controlled electrodeposition of metals at current collectors. Electrolyte formulation is also essential for establishing a stable solid-electrolyte interphase (SEI) and suppressing undesirable side reactions such as degradation, dissolution, and corrosion of active materials. Furthermore, appropriate electrolyte selection mitigates dendrite growth during metal electrodeposition and reduces parasitic hydrogen or oxygen evolution reactions in aqueous systems.

Traditionally, electrolyte design has relied on molar salt concentration to balance ionic conductivity and viscosity. However, recent advances have shown that increasing salt concentrations can profoundly alter ion solvation structures and interfacial reactivities,<sup>2–4</sup> especially for aqueous electrolytes spanning dilute to “water-in-salt” (WiS) regimes. WiS electrolyte (WiSE) formulations, characterized by a low water-to-cation ratio and an anion-rich environment, have been particularly effective for expanding the electrochemical stability window (up

to ~3 V) by reducing water activity<sup>5,6</sup> and thereby suppressing both water reduction and oxidation.<sup>7</sup> This enables the reversible cycling of low-potential anodes such as Mo<sub>6</sub>S<sub>8</sub>, TiO<sub>2</sub> or Li<sub>6</sub>Ti<sub>5</sub>O<sub>12</sub> (LTO),<sup>7–9</sup> thereby opening pathways to high-energy-density aqueous batteries. Concentrated electrolytes also significantly impact the thermodynamics of ion insertion and metal electrodeposition. For instance, high lithium-ion concentrations in non-aqueous systems enable highly reversible lithium plating/stripping<sup>10</sup> and graphite intercalation.<sup>11–16</sup> This behaviour is attributed to an unusual upshift in the equilibrium potential of metal electrodeposition or ion-insertion, which reduces the driving force for side reactions (such as dendrite formation or graphite exfoliation), and thereby enhances reversibility. Comparable potential shifts have been observed for many insertion compounds (*e.g.*, LiFePO<sub>4</sub>, LiCoO<sub>2</sub>, LTO) in both non-aqueous<sup>17</sup> and aqueous electrolytes,<sup>18–20</sup> as well as for various metal ions insertion<sup>21</sup> or deposition processes.<sup>22–26</sup> These findings suggest a generic phenomenon fundamentally rooted in the physicochemical properties of the electrolyte, with applications extending well beyond the field of rechargeable batteries, as recently illustrated by electrochemical Li-mediated nitrogen reduction for sustainable ammonia synthesis,<sup>27</sup> lithium-ion battery recycling,<sup>28</sup> and facilitated electrocatalytic reduction of CO into C<sub>2</sub>H<sub>4</sub>.<sup>29</sup> All these underscore the need to quantitatively understand how electrolyte nature, composition, and concentration govern the electrode potential.

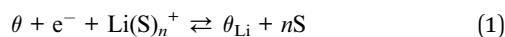
To address this, Yamada and co-workers proposed a computational framework to determine the total free energy of

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solvated  $\text{Li}^+$  ions across concentrations.<sup>17,21</sup> This method, based on molecular dynamics simulations that sum short-range coulombic interactions within the  $\text{Li}^+$  solvation shell, allows to reproduce experimental potential shifts ( $>0.3$  V) for  $\text{Li}^+$  insertion and Li metal electrodeposition. However, this computational approach is demanding and less intuitive regarding the underlying physicochemical factors driving the potential shifts. A classical physicochemical model that captures electrolyte thermodynamics and links directly to the Nernst equation offers assuredly a more accessible alternative. However, for that purpose, it is necessary to formulate a relevant Nernst expression for the electrochemical process under consideration.

If we consider the case of reversible insertion of  $\text{Li}^+$  into an active material, the latter can be described by the following cation-insertion coupled electron-transfer reaction at the electrode/electrolyte interface (assuming an ideal host lattice with equivalent, non-interacting insertion sites):



where  $\theta$  is the fraction of unoccupied insertion sites (oxidized states) converted to  $\theta_{\text{Li}}$  (reduced states) through the coupled insertion of desolvated  $\text{Li}^+$  from solvated  $\text{Li}(\text{S})_n^+$ , with  $n$  solvating solvent molecules being released during reduction and re-associated upon oxidation. The Nernst potential  $E$  (in V) is then given by:

$$E = E_{\theta_{\text{Li}}(\text{S})_n^+/\theta_{\text{Li}}}^0 + \frac{RT}{F} \ln \frac{\theta}{\theta_{\text{Li}}} + \frac{RT}{F} \ln \frac{a_{\text{Li}(\text{S})_n^+}}{a_{\text{S}}^n} \quad (2)$$

where  $E_{\theta_{\text{Li}}(\text{S})_n^+/\theta_{\text{Li}}}^0$  is the standard potential of the insertion material (in V vs. a given reference electrode, with the standard state defined for a material at half state of charge and a solute at a 1 mol  $\text{kg}^{-1}$  or 1 M under putative infinite dilution conditions),  $a_{\text{Li}(\text{S})_n^+}$  the activity of solvated  $\text{Li}^+$ , and  $a_{\text{S}}$  the activity of solvent. This expression explicitly accounts for the  $n$  solvent molecules acting as co-reactants, a contribution that is often neglected.<sup>20,30,31</sup> Omitting this term is equivalent to incorrectly assume that, regardless of concentration,  $a_{\text{S}} = 1$ , thereby leading to misinterpretations of potential shifts.

For  $\theta_{\text{Li}} = 0.5$  and using the convention that  $a_{\text{Li}(\text{S})_n^+} = a_{\text{LiX}} = \gamma_{\text{LiX}} C_{\text{LiX}}$  (with  $\gamma_{\text{LiX}}$  and  $C_{\text{LiX}}$ , the mean ionic activity coefficient and concentration of the LiX salt),<sup>32</sup> eqn (2) then leads to the following half equilibrium potential:

$$E_{1/2} = E_{\theta_{\text{Li}}(\text{S})_n^+/\theta_{\text{Li}}}^0 + \frac{RT}{F} \ln(\gamma_{\text{LiX}} C_{\text{LiX}}) - n \frac{RT}{F} \ln(a_{\text{S}}) \quad (3)$$

Applying these Nernst equations (*i.e.*, eqn (2) or (3)) correctly poses several significant challenges. The first concerns the accurate estimation of salt and solvent activities. It is worth noting that similar Nernst-type expressions have been employed to interpret the potential shifts of  $\text{Li}^+$  intercalation into graphite from molten glymes-Li salt mixtures<sup>16,33</sup> and during Li plating as a function of LiTFSI or LiFSA concentration in non-aqueous solvents.<sup>33-35</sup> However, in all these studies, the activities of  $\text{Li}^+$  and of the free solvent were approximated by their

concentrations—a simplification that becomes invalid in concentrated electrolytes. Another difficulty lies in assigning an appropriate value to  $n$ , which is expected to vary with the nature of salt and solvent, and with the electrolyte concentration. A further complication arises from the potential shift due to the liquid junction potential ( $\Delta\phi_{\text{j}}$ ) generated across the salt bridge of the reference electrode. Although numerous theoretical approaches exist to estimate this potential, none are fully satisfactory.<sup>32</sup> Nevertheless, given that many studies report maximal  $\Delta\phi_{\text{j}}$  values on the order of a few tens of mV,<sup>20,32,36-38</sup> this contribution can be regarded as negligible in the context of the present work.

To demonstrate the applicability of Nernst eqn (3) to a well-defined illustrative system, we considered here the reversible insertion of  $\text{Li}^+$  into  $\text{TiO}_2$  (anatase) from dilute to saturated aqueous LiCl electrolytes. An advantage of using aqueous LiCl electrolytes is their well-established thermodynamic properties and the very high solubility of LiCl in water, allowing to easily access WiS regimes. To experimentally determine how the solvation number  $n$  (or hydration number  $h$  in water) evolves with the electrolyte concentration, we employed a ratiometric Raman analysis of the hydration shell.<sup>39</sup> Moreover, to support the validity of our experimental analysis, the thermodynamic properties of LiCl solutions were modeled on the basis of a recently developed speciation-based solution model<sup>40-42</sup> capable of relating the solution vapor pressure (solvent activity) to the modeled solute and solvent speciated concentrations. This approach enables accurate prediction of the water activity ( $a_{\text{w}}$ ), solute activity ( $a_{\text{LiCl}}$ ), and hydration number ( $h$ ) up to saturation.

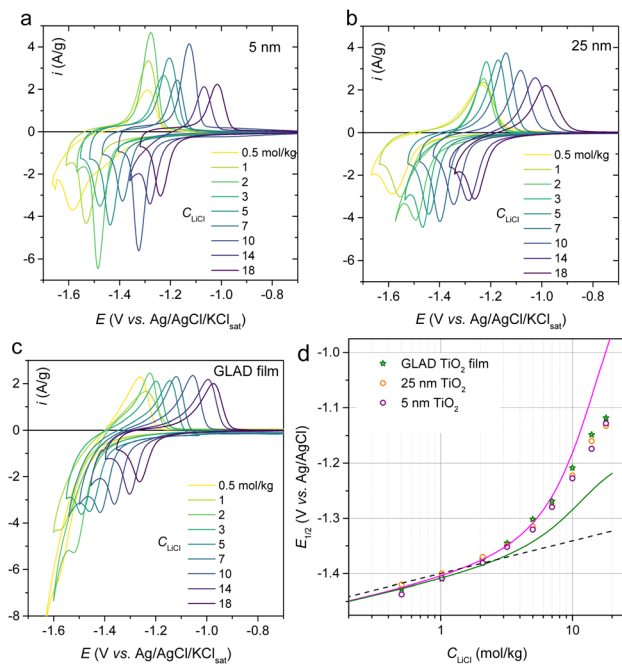
## Results and discussion

Reversible  $\text{Li}^+$  insertion equilibrium potentials into anatase were obtained from cyclic voltammograms (CVs, Fig. 1a-c) recorded in aqueous LiCl electrolytes (0.5 to 18 mol  $\text{kg}^{-1}$ ) at various anatase-based electrodes, *i.e.* (i) composite electrodes made of 80 wt% anatase nanoparticles (5 or 25 nm), 10 wt% carbon, and 10 wt% Nafion binder on titanium current collectors, and (ii) 3D nanostructured pure anatase films (1  $\mu\text{m}$ -thick) on Ti-coated glass.<sup>43,44</sup> The three-electrode cell included a Ag/AgCl (saturated KCl) reference and a Pt counter electrode.

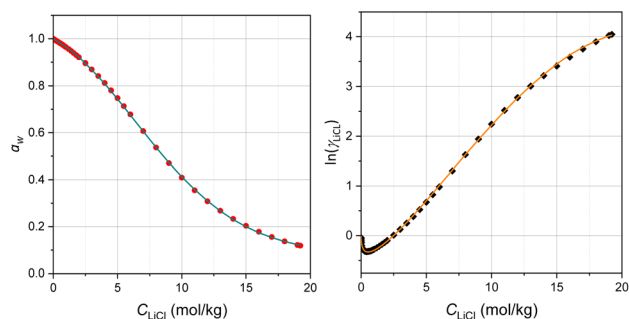
A slow scan rate of 1  $\text{mV s}^{-1}$  was used to ensure near-thermodynamic control. As shown in Fig. 1d, the extracted half-wave potentials,  $E_{1/2}$  (average anodic/cathodic peaks, which is assumed to be close to the half equilibrium potential governed by eqn (3)), evolved identically with LiCl concentration across all electrodes, regardless of their morphology, reflecting complete dependence on the electrolyte's thermodynamic properties. As a result, one can anticipate that the  $E_{1/2}$  evolution as a function of LiCl concentration can be quantitatively reproduced from Nernst eqn (3) and from the knowledge of  $\gamma_{\text{LiCl}}$  (and thereby  $a_{\text{LiCl}}$ ),  $a_{\text{w}}$ , and hydration number  $h$ .

Experimental values of  $a_{\text{w}}$  and  $\gamma_{\text{LiCl}}$  at 25 °C for LiCl molalities ranging from 0.1 mol  $\text{kg}^{-1}$  to saturation are available from compiled literature data (plots shown in Fig. 2).<sup>45</sup> Using these data, the  $E_{1/2}$  evolution was calculated with eqn (3). By adjusting the hydration number  $h$  (where  $n = h$ ) to a constant





**Fig. 1** Potential shifts in CV as a function of LiCl concentration. (a–c) CVs of anatase electrodes recorded at  $1 \text{ mV s}^{-1}$  in aqueous LiCl electrolytes with molalities ranging from  $0.5$  to  $18 \text{ mol kg}^{-1}$  (see legends on graphs). The anatase electrodes were (a and b) porous composite electrodes containing  $80 \text{ wt}\%$  of (a)  $5 \text{ nm}$  or (b)  $25 \text{ nm}$  anatase nanoparticles, and (c) nanostructured anatase films ( $1\text{-}\mu\text{m}$ -thick) deposited on Ti-coated glass. (d) Semilogarithmic plots of  $E_{1/2}$  as a function of LiCl concentration. Open dot and star symbols:  $E_{1/2}$  values extracted from the CVs in panels a–c. The dashed and solid curves correspond to plots of Nernst eqn (3) (using  $E_{\theta, \text{Li}(\text{H}_2\text{O})_h^+/\theta\text{Li}}^0 = -1.4 \text{ V}$ ) under different conditions: dashed black line, ideal case with  $\gamma_{\text{LiCl}} = 1$  and  $a_w = 1$ ; solid green curve,  $\gamma_{\text{LiCl}}$  values taken from literature (also reported in Fig. 2) and  $a_w = 1$ ; magenta solid curve, full dependence of  $\gamma_{\text{LiCl}}$  and  $a_w$  on LiCl concentration (see Fig. 2), including the solvation number  $n$  (or  $h$ ) adjusted to  $4.5$ .



**Fig. 2** Water activity and mean ionic activity coefficient of LiCl solutions. Solid dots: experimental (left) water activity and (right) mean ionic activity coefficient of LiCl as a function of concentration (from ref. 45). Solid curves: best fits of the thermodynamic speciation-based solution model to the experimental data using  $K_h = 5.46$ ,  $m = 2.21$ ,  $K_{id} = 0.45$ , and  $r_i = 0.53 \text{ nm}$ , see text.

value of  $4.5$  and using  $E_{\theta, \text{Li}(\text{H}_2\text{O})_h^+/\theta\text{Li}}^0 = -1.4 \text{ V}$ , the resulting curve (magenta) in Fig. 1d closely matches the experimental  $E_{1/2}$  at low LiCl concentrations (below  $7 \text{ mol kg}^{-1}$ ). This  $h = 4.5$  value

aligns with literature reports for dilute lithium salt solutions, where various experimental and theoretical approaches consistently estimate an average of  $4\text{--}6$  water molecules per lithium ion.<sup>†46–49</sup>

It is worth noting that if only the activity of LiCl is considered in the Nernst equation (setting the solvent activity to  $1$ , *i.e.*  $a_w = 1$ ), it leads to the green curve in Fig. 1d which fails to reproduce the experimental data, except at the lowest concentrations. This result indirectly underscores the importance of accounting for the contribution of water activity raised to the power  $h$  in Nernst eqn (3) (with  $n = h$ ). Finally, the black dashed line in Fig. 1d corresponds to the case where  $a_w = 1$  and  $\gamma_{\text{LiCl}} = 1$ , representing ideal dilute conditions in which only the LiCl concentration contributes.

From the theoretical magenta plot in Fig. 1d, it is evident that at the highest concentrations, the curve diverges from the experimental trend, increasing more rapidly. This behavior is consistent with the progressive depletion of free water molecules available to fully hydrate the solute, specifically  $\text{Li}^+$  ions.<sup>†50</sup> Indeed, at high salt concentrations, the number of water molecules becomes insufficient to maintain a constant hydration number of  $4.5$ . A simple calculation shows that at approximately  $12 \text{ mol kg}^{-1}$  LiCl, the  $55 \text{ mol}$  of water present are not enough to ensure fully hydrated ions. This thus highlights the need to determine how the hydration number  $h$  evolves with the electrolyte concentration, which we investigated using both experimental and theoretical approaches.

Given the complexity of solvation phenomena in concentrated electrolytes, it is useful to clarify the physical meaning of the hydration number  $h$  employed in this work before discussing its evolution with electrolyte concentration. The hydration number  $h$  is here defined as an average effective solvation number, representing the mean number of water molecules interacting with the solute species through sufficiently strong ion-dipole or hydrogen-bond interactions to be distinguished from bulk (free) water. This definition therefore extends beyond the strict first coordination shell and accounts for the dynamic exchange between tightly and loosely bound water molecules. The quantity  $h$  thus corresponds to an averaged value encompassing all solute-related species present in solution, including free ions and ion pairs, which naturally integrates the effects of ion pairing and aggregation without the need to assign separate hydration parameters to each species. Within this formulation, the decrease of  $h$  with increasing electrolyte concentration reflects both the consumption of solvent molecules for solvation and the progressive formation of contact and solvent-shared ion pairs in concentrated and water-in-salt regimes.

Hydration numbers of ions in aqueous media can be evaluated using various spectroscopic methods<sup>46,51</sup> such as neutron or X-ray diffraction, Raman spectroscopy, NMR, as well as by measurements of solution compressibility<sup>52</sup> or theoretical chemistry computations.<sup>53</sup> However, many of these methods have primarily determined hydration numbers under so-called infinite dilute conditions and have rarely assessed the evolution of  $h$  with solute concentration.<sup>52,54</sup> Among these techniques, solution Raman spectroscopy is particularly attractive due to its accessibility. Therefore, we have used it to experimentally determine  $h$



over a range of LiCl electrolyte concentrations, taking advantage of the OH stretching band change at 2800–3700  $\text{cm}^{-1}$  and using the ratiometric methodology proposed by Wang *et al.*<sup>39</sup> to extract hydration-shell Raman spectra, which reflects the structural features of the hydration shell.

Briefly, given that the LiCl salt makes no direct contribution to the O–H stretching band, the Raman spectra, in the 2800–3700  $\text{cm}^{-1}$  region (Fig. 3a), only relies on a linear combination of the spectra corresponding to bulk (or free) water molecules (*i.e.*  $I_{\text{bulk}}(\nu)$ , equivalent to the spectra of pure water) and water molecules in the solute hydration shell (*i.e.*  $I_{\text{hydration}}(\nu)$ ):

$$I_{\text{solution}}(\nu) = I_{\text{bulk}}(\nu) + I_{\text{hydration}}(\nu) \quad (4)$$

Upon increasing the salt concentration, the intensity below 3200  $\text{cm}^{-1}$  progressively decreases (Fig. 3a). This trend reflects a diminution of the number of water molecules involved in double donor-double acceptor hydrogen bonding, which constitute one of the dominant hydrogen-bond motifs in pure water, and is associated to large water clusters.<sup>55</sup> It clearly reflects the consumption of solvent molecules by solvation as the salt concentration increases. This spectroscopic evolution provides direct experimental support for the thermodynamic framework developed in the following section, which

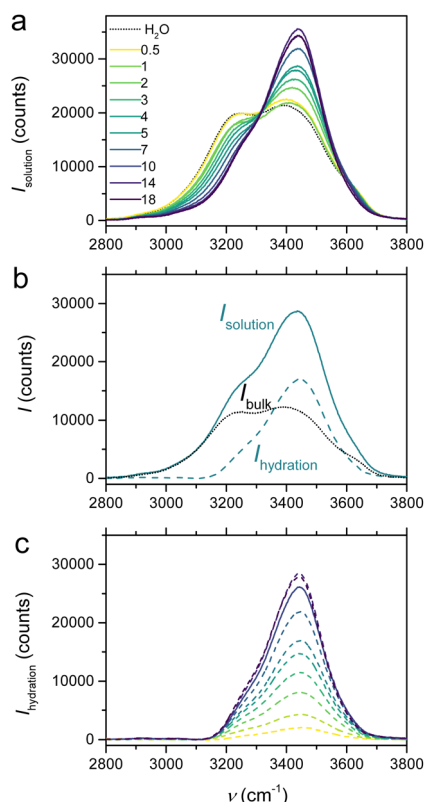


Fig. 3 Raman spectra analysis of the aqueous LiCl electrolytes in the 2800–3100  $\text{cm}^{-1}$  region: (a) solution spectra of pure water (dotted gray line) and LiCl electrolytes with concentrations ranging from 0.5 to 18  $\text{mol kg}^{-1}$  (solid yellow to dark lines); (b) deconvolution of the solution spectra (solid lines) into bulk-like (dotted lines) and hydration-shell (dashed lines) components (at 5  $\text{mol kg}^{-1}$ ); (c) hydration-shell spectra for all electrolyte concentrations (same code color as in a).

quantitatively relates the variation of solvent activity and the effective solvation number to the concentration-dependent behaviour of the electrolyte.

Raman ratio spectra calculated for each salt concentration were used to extract the proportionality coefficient linking  $I_{\text{bulk}}(\nu)$  to the spectra of pure water (see SI for details and Fig. S1). Next, the Raman spectra of both bulk and hydration-shell water molecules can be determined, as shown in Fig. 3b for the 5  $\text{mol kg}^{-1}$  LiCl electrolyte, and Fig. 3c for the hydration-shell spectra determined at all salt concentrations. These spectra account for any solvent molecules whose vibrational spectra are significantly perturbed by the solute. From their integration, the number of water molecules in the hydration shell (*i.e.*,  $n_{\text{hydration}}$ ) can be estimated, allowing in turn to access the hydration number  $h$ , being defined as  $h = n_{\text{hydration}}/n_{\text{solute}}$ .

The hydration number  $h$  is finally given by the following expression:

$$h = \frac{\int I_{\text{hydration}}(\nu) d\nu}{\frac{n_{\text{LiCl}}}{n_{\text{water}} \int I_{\text{bulk}}(\nu) d\nu + \frac{\sigma_{\text{hydration}}}{\sigma_{\text{pure}}} \int I_{\text{bulk}}(\nu) d\nu}} \quad (5)$$

where  $\sigma_{\text{hydration}}$  and  $\sigma_{\text{pure}}$  are the Raman scattering cross section corresponding to hydration shell and pure water molecules, determined as described in the SI.

The  $h$  values thus extracted (red diamonds in Fig. 4a) show a continuous decline with increasing LiCl molality, starting from an initial value of  $\sim 4.5$  at low concentrations to a minimum of about 1.7 near saturation (18  $\text{mol kg}^{-1}$ ). Injecting these  $h$  values along with the corresponding  $a_w$  and  $\gamma_{\text{LiCl}}$  values (inferred from Fig. 2) into eqn (3) yields, for each tested LiCl

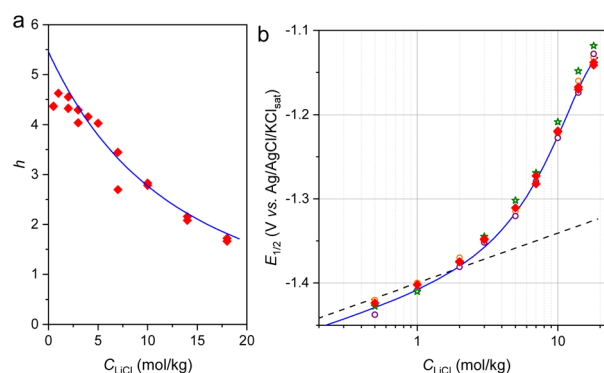


Fig. 4 Estimation of the hydration number  $h$  and its impact on  $E_{1/2}$ . (a) Average hydration number as a function of LiCl molality. Solid red diamonds: data from solution Raman spectroscopy (see SI for details). Solid blue line: theoretical curve obtained from the speciation-based solution model. (b) Semilogarithmic plots of  $E_{1/2}$  as a function of LiCl concentration. Open dot and star symbols: same experimental data as in Fig. 1d. Dashed black line: theoretical Nernst eqn (3) under ideal conditions (*i.e.*,  $a_w = 1$  and  $\gamma_{\text{LiCl}} = 1$ ). Solid red diamonds: Nernst eqn (3) using dependences of  $a_w$  and  $\gamma_{\text{LiCl}}$  on LiCl concentration (Fig. 2), plus hydration number  $h$  from Raman spectroscopy (panel a). Solid blue line: Nernst eqn (3) plotted using the values of  $a_w$ ,  $\gamma_{\text{LiCl}}$ , and  $h$  modeled from the speciation-based solution model (using  $K_h = 5.46$ ,  $m = 2.21$ ,  $K_{\text{id}} = 0.45$ , and  $r_i = 0.53$  nm, see text).



concentration, the red diamonds in Fig. 4b. The excellent overlap of the calculated and experimental  $E_{1/2}$  potentials clearly validates our approach.

To advance further, we employed the speciation-based solution model recently proposed by Wilson and coworkers<sup>40–42</sup> in order to reproduce the non-ideal variations of water and solute activities as a function of LiCl concentration, and also to provide the associated concentration-dependent hydration number. Moreover, to account not only for short-range ion–solvent interactions (*via* stepwise hydration equilibria under the law of mass action) and ion pairing, as proposed by Wilson and coworkers,<sup>40,42</sup> the model was extended to include long-range electrostatic contributions through Debye–Hückel theory.

Grounded in mechanistic principles, this model requires only a minimal set of 4 adjustable parameters: a composite hydration equilibrium constant  $K_h$ , a hydration modifier  $m$ , an ion-pair dissociation constant  $K_{id}$ , and the mean ionic radius  $r_i$  of the solute.

A primary key assumption of the model is that the experimental activity of water  $a_w$  defined by the solution vapour pressure,  $p$ , relative to the pure water at the same temperature,  $p^0$ , directly correspond to the mole fraction of free-water,  $x_w^{\text{free}}$ :<sup>40</sup>

$$a_w = \frac{p}{p^0} = x_w^{\text{free}} \quad (6)$$

Based on this assumption, the mole fraction of a 1 : 1 MX speciated solute  $i$ , *i.e.*  $x_i^{\text{spe}}$ , can be found through the water activity residual:

$$1 - a_w = x_i^{\text{spe}} \quad (7)$$

The secondary key assumption of the model is that ion hydration proceeds *via* step-wise hydration equilibria that can be aggregated, simplified, and expressed as the following relationship for the speciated (hydrated) solute concentration,  $x_i^h$  (in mole fraction):<sup>42</sup>

$$x_i^h = \frac{x_i}{x_w - hx_i + x_i} = \frac{x_i}{1 - hx_i} \quad (8)$$

where  $x_i$  and  $x_w$  are the absolute mole fractions of solute  $i$  and water  $w$ , respectively, the latter being given for a fully dissociated 1 : 1 MX salt in water by:

$$x_i = \frac{2C_i}{55.509 + 2C_i} \quad (9)$$

$$x_w = 1 - x_i = \frac{55.509}{55.509 + 2C_i} \quad (10)$$

with  $C_i$  the molality of the salt in solution, and 55.509 the number of moles of water contained in 1 kg.

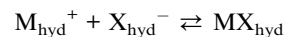
In eqn (8), the hydration number  $h$  is a function of the water mole fraction, as follows:<sup>42</sup>

$$h = K_h x_w^m \quad (11)$$

with  $K_h$  the composite hydration association constant and  $m$  the hydration modifier. This expression defines hydration as being

exponentially proportional to the mole fraction of water through the parameter  $m$ . A second-order hydration term, *i.e.*  $m = 2$ , is generally sufficient to accurately model the electrolyte behavior from dilute conditions up to saturation.<sup>42</sup>

To account for the fact that certain salts exhibit a degree of ion pairing, the following hydrated solute equilibrium has to be considered:



This equilibrium is governed by the following dissociation equilibrium constant,  $K_{id}$  (expressed here in terms of activities in the mole fraction scale):

$$K_{id} = \frac{a_M a_X}{a_{MX}} \quad (12)$$

where  $a_M$ ,  $a_X$ ,  $a_{MX}$  are the activities of the speciated cation  $M^+$ , anion  $X^-$ , and ion-pair  $MX$ .

In this context, the mole fractions of the partially dissociated and hydrated species (*i.e.*, cation  $M^+$ , anion  $X^-$ , and ion-pair  $MX$  species) must be expressed as follows:‡

$$x_M^{h,K} = \frac{\alpha C_i}{55.509 - h(1 + \alpha)C_i + (1 + \alpha)C_i} \quad (13)$$

$$x_X^{h,K} = \frac{\alpha C_i}{55.509 - h(1 + \alpha)C_i + (1 + \alpha)C_i} \quad (14)$$

$$x_{MX}^{h,K} = \frac{(1 - \alpha)C_i}{55.509 - h(1 + \alpha)C_i + (1 + \alpha)C_i} \quad (15)$$

where  $\alpha$  is the degree of dissociation of the salt.

Assuming in eqn (12) that each activity  $a$  is equal to the “true” mole fractions  $x^{h,K}$  of the partially dissociated and hydrated species, then the equilibrium constant can be written:

$$K_{id} = \frac{x_M^{h,K} x_X^{h,K}}{x_{MX}^{h,K}} = \frac{\alpha^2}{(1 - \alpha)} \frac{C_i}{55.509 - h(1 + \alpha)C_i + (1 + \alpha)C_i} \quad (16)$$

with  $\alpha$  determined by:

$$\alpha = \frac{55.509 K_{id} - \sqrt{\Delta}}{2C_i(K_{id}h - K_{id} - 1)} \quad (17)$$

and  $\Delta = (55.509 K_{id})^2 - 4K_{id}C_i(K_{id}h - K_{id} - 1)(55.509 - hC_i + C_i)$ .

From the determination of  $\alpha$ , it becomes thus possible to calculate the mole fractions  $x_M^{h,K}$ ,  $x_X^{h,K}$ , and  $x_{MX}^{h,K}$ , as well as  $x_i^{h,K}$  from:

$$x_i^{h,K} = x_M^{h,K} + x_X^{h,K} + x_{MX}^{h,K} \quad (18)$$

Using eqn (18), one can thus obtain the mole fraction of free solvent:  $x_w^{h,K} = 1 - x_i^{h,K}$ .

At this stage, the long-range electrostatic contribution to the hydrated and partially dissociated solute can be introduced into the model through the following modification of water activity by the Debye–Hückel term,  $K_w^{\text{DH}}$ :<sup>41</sup>



$$x_w^{h,K,DH} = 1 - x_i^{h,K,DH} = \frac{x_w^{h,K}}{K_w^{DH}} = \frac{1 - x_i^{h,K}}{K_w^{DH}} \quad (19)$$

with

$$K_w^{DH} = \exp\left(-\frac{2A_x I^{3/2}}{1 + r_i B_x I^{1/2}}\right).$$

Here,  $A_x = 2.917$  and  $B_x = 24.4$  are the Debye-Hückel constants (at 298 K) in the mole fraction scale,  $r_i$  is a parameter related to the average ionic size of the solute, and  $I$  is the ionic strength of the electrolyte calculated from the mole fraction of hydrated and partially dissociated solute (wherein MX contribution is null because  $z_{MX} = 0$ ):

$$I = \frac{1}{2} \sum (x_{Me}^{h,K} z_M^2 + x_X^{h,K} z_X^2 + x_{MX}^{h,K} z_{MX}^2) \quad (20)$$

The resulting modeled free water activity, *i.e.*  $x_w^{h,K,DH}$  (thus equivalent to  $x_w^{free}$  in eqn (6)), can then be used to fit either the experimental data of  $a_w$  or the mean solute activity coefficient  $\gamma_i$  (by using first eqn (21) for calculating the osmotic coefficient  $\phi$  in the absolute molal scale of  $C_i$  and then using eqn (22) to calculate  $\gamma_i$  knowing  $\phi$ ) through iterative adjustments of  $K_h$ ,  $m$ ,  $K_{id}$ , and  $r_i$ .

$$\phi = -55.509 \frac{\ln(x_w^{h,K,DH})}{2C_i} \quad (21)$$

$$-\ln \gamma_i = (1 - \phi) + \int_0^{C_i} \frac{(1 - \phi)}{C_i} dC_i \quad (22)$$

From the best fits (see Fig. 2), the following values were obtained:  $K_h = 5.46$ ,  $m = 2.21$ ,  $K_{id} = 0.45$ , and  $r_i = 0.53$  nm.

The  $K_{id}$  value is in the order of magnitude of that determined experimentally,<sup>56</sup> while the values of  $K_h$  and  $m$  are close to those previously reported by Wilson and coworkers.<sup>42</sup> Regarding the value of  $r_i$ , it aligns with expectations for the average ionic radius of hydrated ions, noting that within Debye-Hückel theory, this parameter is generally treated as flexible and non-physical. Knowing  $K_h$  and  $m$ , the evolution of  $h$  as a function of  $C_{LiCl}$  can be plotted from eqn (11) (blue plot in Fig. 4a). The calculated trend almost perfectly matches the experimental Raman data. This remarkable agreement between two independent approaches—solution Raman spectroscopy and water activity modeling—strongly supports our quantitative determination of  $h$  evolution with  $C_{LiCl}$ . On this basis, the variation of  $E_{1/2}$  was subsequently determined from eqn (3) using the values of  $a_w$ ,  $\gamma_i$ , and  $h$ , directly obtained from the model. This yields the blue curve in Fig. 4b, which nearly overlaps the experimental data. This strong agreement unambiguously validates our approach and, in particular, the proposed Nernst-type equation, which—when combined with an appropriate thermodynamic model—provides a consistent, quantitative prediction of how the insertion potential of a cation evolves with electrolyte concentration. The good overlap of both the Raman-based and model-based plots with the experimental data in Fig. 4b also indirectly confirms that the potential junction contribution to the potential shift can be considered as negligible.

With this improved understanding of the key electrolyte parameters governing the electrode potential associated with ion insertion into a material, we can better rationalize how to mitigate undesirable competing side reactions. This is clearly illustrated by the CVs in Fig. 1c. As the LiCl concentration increases, the reversible  $Li^+$  insertion wave into  $TiO_2$  progressively shifts toward more positive potentials, whereas the irreversible and undesirable hydrogen evolution reaction remains nearly unaffected, as its reduction potential changes only slightly (thermodynamically) with LiCl concentration (as attested by the electrolyte pH which remains almost constant—see Table S1 in SI). This insight largely explains the superior coulombic efficiencies reported for numerous materials in WiSEs.<sup>2,7–9,57,58</sup>

## Conclusion

In this work, we present a rigorous unification and quantitative validation of electrolyte thermodynamics in insertion systems to rationalize the anomalous evolution of the potential of an insertion material as a function of electrolyte concentration, spanning dilute regimes up to water-in-salt conditions. We demonstrate that the Nernst equation can quantitatively describe the variation of electrode potential, provided that both the salt mean activity and the solvent activity, raised to a power corresponding to the electrolyte hydration number, are properly considered. We further show how a decrease in the hydration number influences the electrode potential in the WiS regime, and how this quantity can be accessed experimentally, through Raman spectroscopy, or theoretically, *via* electrolyte modelling approaches.

By focusing on a 1 : 1 electrolyte, this study lays the groundwork for the subsequent interpretation of potential shifts observed in more complex electrolyte systems, including those exhibiting richer solution speciation, bi-salt formulations, or the presence of co-solvents. Developing a comprehensive understanding of how electrolyte composition and concentration govern electrode potential remains a key contemporary challenge, central to the rational design of advanced electrochemical devices for energy storage and beyond.

## Experimental section

### Chemicals

Anhydrous lithium chloride (free-flowing, Redi Dri™, ACS reagent,  $\geq 99\%$ ) was purchased from Sigma-Aldrich/Merck. Nafion D-520 (5 wt% in a mixture of lower aliphatic alcohols and water) was obtained from Sigma-Aldrich/Merck. Ketjenblack EC-600JD was supplied by Nanografi Nano Technology.  $TiO_2$  anatase nanopowder ( $< 25$  nm diameter) were purchased from Sigma-Aldrich, and  $TiO_2$  pure anatase nanoparticles (5 nm diameter) were obtained from NanoAmor (USA). All aqueous solutions were prepared with Milli-Q water (18.2 M  $\Omega$  cm).

### Electrodes

Anatase GLAD- $TiO_2$  electrodes were prepared as previously described.<sup>43</sup> Composite electrodes were fabricated from a slurry consisting of 80 wt% active material, 10 wt% Ketjenblack EC-



600Jd, and 10 wt% Nafion. The active material and carbon were first dispersed in the Nafion suspension. Approximately 100  $\mu\text{L}$  of Milli-Q water was then added to obtain a homogeneous ink. The suspension was stirred for 2 h and cast onto titanium sheets (0.127 mm thickness) using a manual Elcometer 3520 Baker film applicator adjusted to a 200  $\mu\text{m}$  height. The electrodes were dried at ambient temperature and pressure. The mass loading of active material was 1.7 to 2  $\text{mg cm}^{-2}$ .

### Electrochemical measurements

Cyclic voltammetry was performed in a standard three-electrode cell using a BioLogic VSP potentiostat controlled by EC-Lab software. A platinum wire served as the counter electrode, and an Ag/AgCl (saturated KCl) electrode was used as the reference. Prior to measurements, the electrolyte was degassed for 20 min, and a constant argon flow was maintained above the solution throughout the experiments. Current densities were calculated by normalizing the measured current to the mass of active material in the electrode. All potentials are reported vs. Ag/AgCl (saturated KCl).

### Electrolyte characterization

All measurements were performed at 22  $^{\circ}\text{C}$ . Conductivity and pH measurements were measured with an Accumet XL200 benchtop meter (Fisher Scientific) calibrated using standard buffer and KCl solutions before each use. Water activity was determined using an AQUALAB 4 TE chilled-mirror dew-point hygrometer (precision  $\pm 0.003$ ). Solution densities were estimated using a 5 mL volumetric flask. All data are reported in Table S1.

Raman spectroscopy was carried out using a LabRAM HR 800 high-resolution Raman microspectrometer operating over 2500–4180  $\text{cm}^{-1}$  spectral range. A He–Ne laser ( $\lambda = 633 \text{ nm}$ ) delivering 8 mW at the sample was used as the excitation source. The spectrometer was coupled to an optical microscope equipped with a 10 $\times$  objective (N.A. = 0.25) focused on a horizontal optical cuvette with a 2 mm optical path. Measurements were performed under air at 25  $^{\circ}\text{C}$ . Two sets of electrolytes were analyzed, with each spectrum acquired with an integration time of 15 s or 20 s and 20 accumulations to improve the signal-to-noise ratio. Baseline correction of the Raman spectra was performed using Labspec 6 software. Typically, a fourth-order polynomial was fitted to subtract the background. When the automatic baseline algorithm did not provide a satisfactory fit, additional baseline points were manually selected.

### Author contributions

The experiments were performed by LKB, TR, LC and AGS. The manuscript was written through contributions of LKB, TR, VB and BL. All authors have given approval to the final version of the manuscript.

### Conflicts of interest

There are no conflicts to declare.

### Data availability

The data supporting this article have been included as part of the supplementary information (SI). The data and calculations that support the findings of this study are available from the corresponding author upon request. Supplementary information: Raman analysis, Tables S1, S2, Fig. S1 and S2. See DOI: <https://doi.org/10.1039/d6sc00393a>.

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### Notes and references

† According to experimental evidence, in many aqueous electrolytes cations are considered more extensively hydrated (or solvated) than anions.

‡ Compared to Wilson's model,<sup>40–42</sup> the ion-pairing equilibrium was treated differently by expressing solute hydration in terms of partially dissociated species (*i.e.*,  $(1 + \alpha)C_i$ ) rather than fully dissociated ones (*i.e.*,  $2 C_i$ ), which avoids approximation and ensures that the dissociation constant  $K_{\text{id}}$  remains strictly constant regardless of solute concentration, thus preserving compliance with the law of mass action.

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