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# Diffusion-Induced Enhanced Photoresponsivity, Detectivity in $\text{Ag}_2\text{S}$ and $\text{In}_2\text{Se}_3$ Heterostructure for UV-Visible Photodetector: An Experimental and Computational Analysis

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## Abstract:

For recent optical communication and imaging systems, photodetectors are playing a pivotal role. Metal chalcogenide-based photodetectors are widely used for visible light photodetection. In this regard, the  $\text{In}_2\text{Se}_3$  and  $\text{Ag}_2\text{S}$  combined heterostructure is a promising candidate for Visible light photodetection. The annealing-induced  $\text{Ag}_2\text{S}$  diffusion into the  $\text{In}_2\text{Se}_3$  layer resulted in a high-performance photo detectivity of  $7.32 \times 10^9$  Jones. It shows the highest photocurrent during both the rise (62.35 nA) and decay (67.74 nA) phases, coupled with a strong  $I_{\text{on}}/I_{\text{off}}$  ratio of 3.96 (rise) and 3.18 (decay). Its rise and fall times ( $\tau_r = 7.15$  s,  $\tau_d = 6.35$  s) are moderate and well-balanced, suggesting efficient charge separation and recombination kinetics. The bandgap of the annealed film increased with the reduction of structural disorder as probed from the UV-Visible spectroscopy and well supported by the DFT result. The amorphous to polycrystalline phase transformation induced surface morphology change, reduced the contact angle, thus decreasing the hydrophobicity. The refractive index decreased with an increase in optical transmission and skin depth, while optical density reduced upon annealing. The X-ray photoelectron spectroscopy revealed the oxidation states of the elements, while energy dispersive X-ray analysis presented the elemental composition in the films. The heterostructure formation and its mixing upon annealing are viewed from the cross-sectional FESEM images, and the presence of the planes was noticed through HRTEM images. The observed optical properties, along with the enhanced photodetection, pave the way toward the construction of novel III-VI metal chalcogenide-based heterojunctions for high-performance and broadband photodetectors.

**Keywords:**  $\text{Ag}_2\text{S}/\text{In}_2\text{Se}_3$  heterostructure; Photodetector; Optoelectronics; Bandgap; Surface wettability.



## 1. Introduction

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The recent requirement demands cost-effective and less toxic material-based high-speed, highly sensitive photodetectors, based on thin films [1]. Thin films are very useful for visible light photo detection. The semiconducting thin films have a wide range of photodetection applications based on their exceptional optoelectronic properties [2-4]. The factors that influence the photodetector performance are response time, detectivity, responsivity, and wavelength range of detection [5,6]. For various applications in instrumentation and industrial purposes, the photo response in the visible and ultraviolet (UV) regimes has gained attention in recent times. It is basically used for image sensing, chemical analysis, radiation detection, environmental monitoring, optical communications, medical diagnostics, astronomical studies, etc. [7]. Photodetectors (PDs), are the most valuable optoelectronic components in recent days. Its principle is based on the transformation of optical energy into electrical one by photoelectric effect [8]. The ultimate aim of recent methods is to have highly efficient, larger spectral range, high-speed photodetectors from less expensive materials, as well as a low-cost fabrication method. For this purpose, the present study relies on the metal-based chalcogen materials.

The binary combinations from the IIIA-VIA group elements formed compounds basically In-based alloys such as  $\text{In}_2\text{S}_3$ ,  $\text{InSe}$ , and  $\text{In}_2\text{Se}_3$  have great importance and are widely used for photodetection applications [9-11]. The important aspects of these combinations include controlled morphology, tunable bandgap, flexibility, good stability, 2D structure, carrier mobility, and remarkable fundamental properties. The 2D materials also have very good light-matter interaction with excellent optoelectronic properties. However, in recent times, the heterostructure films have been widely used for photodetector applications in a broad wavelength scale. For a wider spectral range, the stacking of  $\alpha\text{-In}_2\text{Se}_3$  with any 2D semiconducting alloy is a good choice for achieving speedy photo response. The  $\alpha\text{-In}_2\text{Se}_3/\text{Si}$  heterojunction acts as a highly efficient heterojunction-based photodetector used for photoelectric imaging and object recognition [12]. The  $\alpha\text{-In}_2\text{Se}_3/\text{Ta}_2\text{NiSe}_5$  heterojunction-based photodetector works both under bias voltage state and self-driven state. The detection waveband covers between 405-1550 nm with a response time of only 25  $\mu\text{s}$  [13]. The  $\alpha\text{-In}_2\text{Se}_3/\text{WSe}_2$  heterostructure-based photodetector has detectivity and photoresponsivity of  $4.34 \times 10^{14}$  Jones and  $4.61 \times 10^5 \text{ A W}^{-1}$  [14]. The self-powered  $\gamma\text{-In}_2\text{Se}_3/\text{p-Si}$  heterojunction-based photodetector has enhanced photoresponsivity of 47.9 mA/W, photosensitivity of 282, and photo detectivity of  $8.45 \times 10^{10}$  Jones. The heterostructure explores humidity and light-intensity-dependent photo response [15]. The partially gated  $\text{MoS}_2/\alpha\text{-In}_2\text{Se}_3$  heterojunction photodetector shows excellent photo response performance and a non-volatile photo memory



effect. This is suitable for applications like non-volatile memory and other optoelectronic applications involving designing 2D ferroelectric devices [16]. The photo response in the 650-900 nm range is shown by a  $\text{WSe}_2/\text{In}_2\text{Se}_3$  heterojunction-based photodetector with good optoelectronic performance at room temperature. The direct band gap multilayer material-based 2D vdW heterojunctions are suitable for future optoelectronic devices [17]. The photo response of  $\text{In}_2\text{Se}_3/\text{MoS}_2$  heterojunction photodetector is up to 1310 nm. Here, the photocurrent generation mechanism shifts from photogating to photoconductive effect [18]. The observation of no persistent photoconductivity with fast transient response is observed from p-GaN/ $\alpha$ - $\text{In}_2\text{Se}_3$  heterojunction-based UV/Near-IR dual-band photodetector. The formation of  $\alpha$ - $\text{In}_2\text{Se}_3$  along with a wide bandgap semiconductor-based heterostructure is used for a special type of optoelectronic device [19].

The different energy treatments on such a heterostructure bring modifications for better performances that include ion irradiation, laser irradiation, annealing, etc [20-22].  $\text{In}_2\text{Se}_3/p$ -Si heterojunction photodetectors show enhanced performance at increased argon ion fluence. It brings out higher photosensitivity, photoresponsivity, and photo detectivity with faster rise and decay time. It shows good response to radiation hardening and is used for developing space-based photodetectors [20]. The excitation wavelength of 638 nm at 10 nW laser power and 1500s by  $\text{In}_2\text{Se}_3/\text{PtSe}_2$  photodetector resulted in a fast response broadband photodetector. Such a photodetector is suitable for low-energy consumption and high-capacity optoelectronic devices [21]. The thermal annealing on SnS-based thin film resulted in the specific detectivity of  $6.8 \times 10^{10}$  Jones, which is nearly two orders of magnitude greater than that of the unannealed film [22]. In our earlier work, the annealing effect on the  $\text{Ag}_2\text{S}$ -based  $\text{In}_2\text{Se}_3$  heterostructure film of 800 nm resulted in remarkable photodetection ability with  $2.01 \times 10^{-1}$  A/W responsivity and detectivity of  $7.32 \times 10^9$  Jones. The nA to mA current increase with 250 °C annealing dramatically enhanced the photo response [23]. The current study relies on the low-thickness-based  $\text{Ag}_2\text{S}/\text{In}_2\text{Se}_3$  heterojunction thin film annealed at various temperatures to probe the optimized performance. The annealed heterostructure was subjected to various experimental tools like X-ray diffraction (XRD), X-ray photoelectron spectroscopy (XPS), cross-sectional field emission scanning electron microscopy (FESEM), and photo response study. The experimental results is supported by the density functional theory (DFT) based calculation.

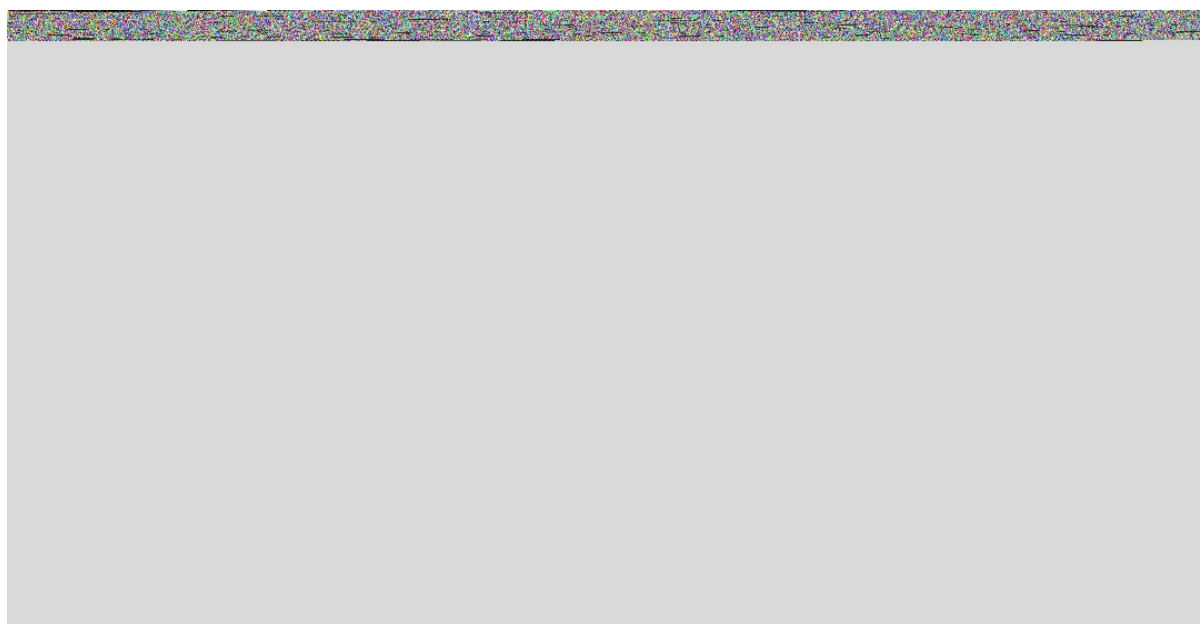
## 2. Experimental details

### 2.1. Thin film deposition and thermal annealing

The formation of a bilayer structure of  $\text{Ag}_2\text{S}$  and  $\text{In}_2\text{Se}_3$  was carried out by the thermal evaporation technique. The HHV Smart coat 3.0 unit was employed for coating of the two



layers from high-purity (99.999% Sigma Aldrich)  $\text{Ag}_2\text{S}$  and  $\text{In}_2\text{Se}_3$  on a cleaned glass substrate. The rate of deposition was slow ( $5\text{\AA}/\text{sec}$ ) for uniformity and homogeneity. The attached thickness controller determines the thickness deposited, and the substrate holder was rotated slowly inside the chamber. The  $\sim 150\text{ nm}$  of  $\text{Ag}_2\text{S}$  layer was deposited onto  $\sim 450\text{ nm}$  of  $\text{In}_2\text{Se}_3$  layer under high vacuum state ( $10^{-5}$  Torr). The preparation was done under normal conditions only. The coated films were then cut into pieces for annealing at different temperatures. The diffusion of  $\text{Ag}_2\text{S}$  inside the  $\text{In}_2\text{Se}_3$  layer was carried out by heat treatment. The hot air oven was used for thermal annealing at  $100\text{ }^\circ\text{C}$ ,  $150\text{ }^\circ\text{C}$ ,  $200\text{ }^\circ\text{C}$  and  $250\text{ }^\circ\text{C}$  temperatures. The time duration of annealing was 2 hrs inside the oven.



**Scheme 1:** Annealing of the films at different temperatures

## 2.2. Experimental characterizations

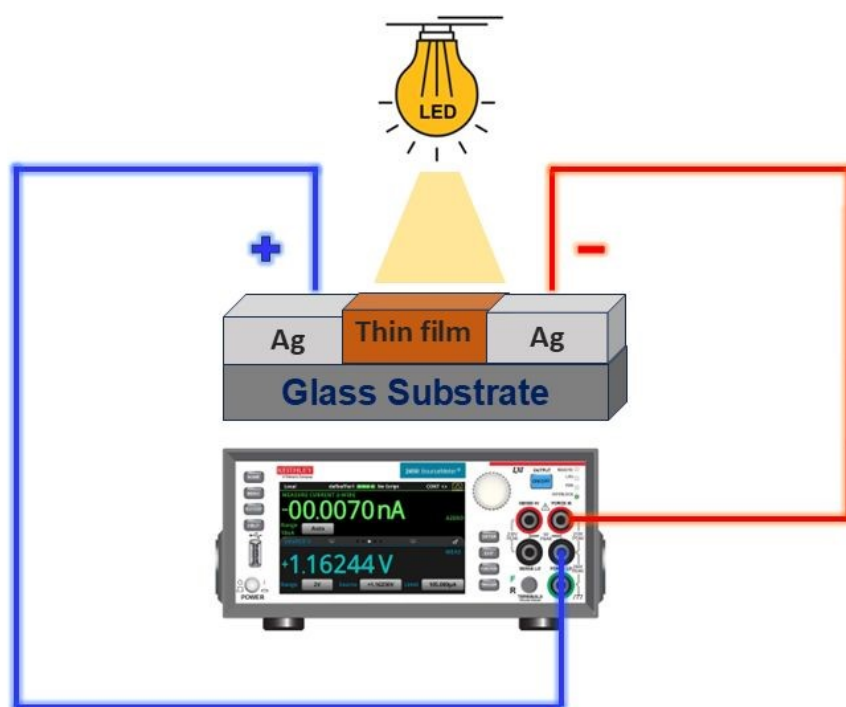
The thermally annealed films were subjected to different types of experiments, like XRD, FESEM, XPS, Energy dispersive X-ray analysis (EDS), Contact angle, UV-Visible spectroscopy, and IV measurement. The formation of the heterostructure was first checked through cross-sectional FESEM by a JEOL-JSM-7610F unit. The structure of the as-deposited and annealed bilayer was determined through a Bruker D8 Advance XRD unit. The scanning was from  $10\text{--}80^\circ$  with a  $0.03^\circ/\text{sec}$  scan rate. Further structural data were collected from HRTEM. The TEM and SAED images were taken by the JEOL HRTEM unit at different scales. The film composition after preparation and annealing was tested by the EDS unit at a high vacuum state ( $9 \times 10^{-5}$  Torr vacuum) at 16 kV voltage. The surface morphology pictures were recorded from a JEOL FESEM unit at different scale bars with high magnifications. The optical data were recorded by a JASCO-770 spectrometer with 1 nm resolution from 500 to



2500 nm wavelength range. The surface wettability test was done by the contact angle meter (DME-211 Plus make). The chemical state of the elements was probed through XPS by Axis Ultra, Kratos Analytical instrument. It has Al  $K_{\alpha}$  energy source (1485.6 eV). The spectra were collected at  $3 \times 10^{-9}$  Torr.

### 2.3. Photodetection study setup:

For the photodetection experiment, the film's edges were carefully scratched on two sides, leaving a central film area ( $0.25 \text{ cm}^2$ ) over the substrate. Silver paste was then applied to the two opposite edges to establish conductive contacts. The specimen was connected to a Source Measure Unit (SMU-Keithley 2450) using two probes. Photodetection measurements were carried out under ambient conditions using a bulb (9W white LED) as the light source, chosen for its broad coverage of the visible spectrum and its reduced risk of excessive local heating.



**Scheme 2:** Photo response study through IV and IT measurement

## 3. Results and Discussion

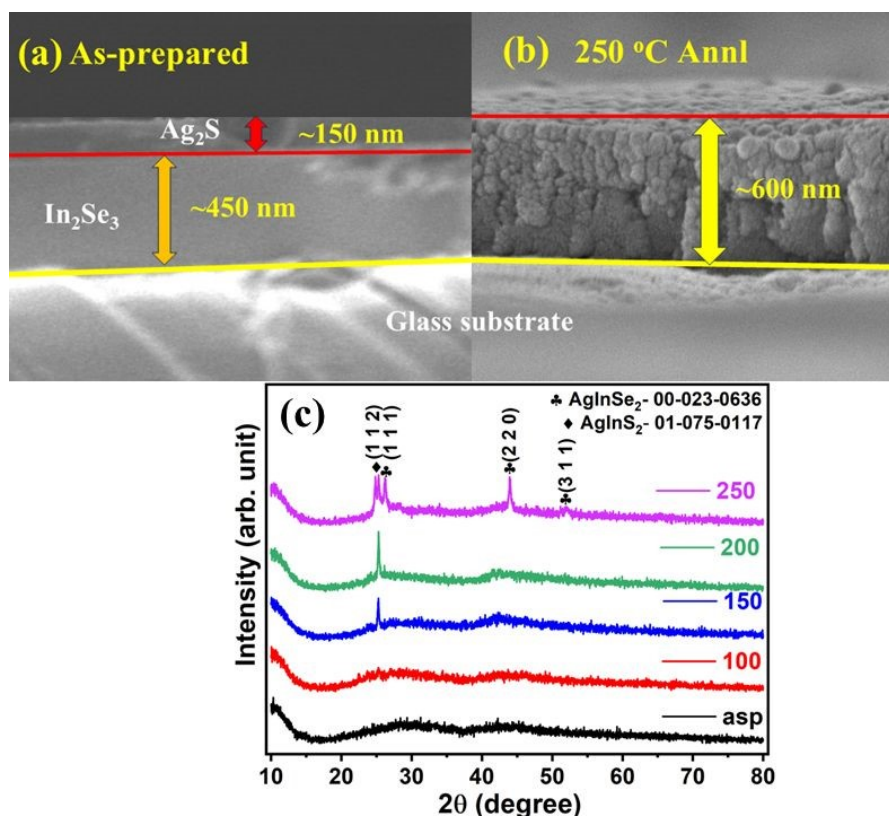
### 3.1. Cross-sectional image and XRD study

To verify the film thickness and bilayer structure formation, the cross-sectional FESEM images were collected from the bilayer and annealed film. Fig. 1(a) presents the bilayer images with a top layer of  $\sim 150 \text{ nm}$  ( $\text{Ag}_2\text{S}$ ) and a bottom layer of  $\sim 450 \text{ nm}$  ( $\text{In}_2\text{Se}_3$ ). Approximately the same thickness was deposited during thermal evaporation, as recorded by the thickness monitor. After annealing, the thickness remained nearly the same, as shown in Fig. 1(b). The disappearance of the distinct layer infers the formation of a single layer by annealing, through





which there is intermixing of the two layers. The different structure can be seen in the annealed film than in the as-prepared one.



**Fig. 1.** Cross-sectional image of (a) as-deposited, (b) annealed film, (c) XRD pattern

XRD analysis, as depicted in Fig. 1(c), reveals a significant transformation in the structural properties of  $\text{Ag}_2\text{S}/\text{In}_2\text{Se}_3$  thin films upon thermal annealing. The as-prepared films demonstrate an amorphous structure, lacking the long-range atomic order characteristic of crystalline materials. In contrast, the application of heat treatment promotes the development of a polycrystalline structure, indicating the formation of distinct crystalline grains within the material. This enhancement in crystallinity is directly proportional to the annealing temperature, suggesting that higher thermal energy facilitates more ordered atomic arrangements and grain growth [23]. The increased annealing temperature also encourages interdiffusion between the  $\text{Ag}_2\text{S}$  and  $\text{In}_2\text{Se}_3$  phases, potentially leading to the formation of new or mixed phases. The XRD patterns of the annealed thin films exhibit several characteristic diffraction peaks, signifying the presence of multiple crystalline phases. Specifically, peaks observed at approximately  $26.21^\circ$ ,  $43.88^\circ$ , and  $51.92^\circ$  are attributed to the cubic  $\text{AgInSe}_2$  phase, correlating with the (1 1 1), (2 2 0), and (3 1 1) crystallographic planes, respectively (ICDD: 00-023-0636). Additionally, a peak detected around  $25.43^\circ$  corresponds to the (1 1 2) plane of the tetragonal  $\text{AgInS}_2$  phase (ICDD: 01-075-0117). The crystallite size, denoted as  $D$ ,



represents the extent of coherent diffraction within the thin films for each specific diffraction peak, essentially reflecting the grain size in a polycrystalline material. This critical structural parameter can be quantitatively determined using Scherrer's equation, which correlates the size of the crystallites to the broadening of the XRD peaks.

The crystallite size, often symbolized as  $D$ , refers to the dimensions of the coherently diffracting domains within a thin film, is expressed as,

$$\text{crystallite size } (D) = \frac{0.94\lambda}{\beta \cos\theta} \quad (1a)$$

Where Bragg's angle is  $\theta$ , and the full width at half maximum (FWHM) is  $\beta$ . Here,  $\lambda = 1.54 \text{ \AA}$  [24].

Lattice strain ( $\epsilon$ ) quantifies the deformation in the dimensions of a crystal lattice relative to its unstrained state. This strain can manifest in two primary forms: uniform and non-uniform strains. Uniform strain involves a consistent expansion or contraction of the unit cell across the entire crystalline volume. While uniform strain does not contribute to the broadening of diffraction peaks, it causes a shift in the peak positions and an alteration in the calculated lattice parameters [25]. This shift is a direct indicator of the average change in interplanar spacing. The lattice strain has been calculated by using the equation [26],

$$\text{lattice strain } (\epsilon) = \frac{\beta \cot\theta}{4} \quad (1b)$$

Dislocation density ( $\delta$ ) is a fundamental microstructural parameter defined as the total length of dislocation lines per unit volume of a crystalline material. While dislocations are non-equilibrium defects, meaning they are not inherently stable according to thermodynamic principles, their observable density holds significant importance in influencing a material's macroscopic properties [27]. These linear crystallographic defects play a crucial role in plastic deformation, strengthening mechanisms, and various other physical phenomena in solids [28].

$$\text{Dislocation density } (\delta) = \frac{1}{D^2} \quad (1c)$$

The quantity of crystallites within a material, often denoted as the number of crystallites ( $N_c$ ), is primarily governed by intrinsic structural characteristics, specifically the crystallite size ( $D$ ) and the extent of agglomeration or clustering among these crystalline domains. These structural factors directly influence how many individual crystallites are present in a given volume or area of the material. A smaller average crystallite size generally implies a larger number of crystallites for a constant material volume and can be calculated as,

$$N_c = \frac{d}{D^3} \quad (1d)$$





The calculated structural parameters are summarized in Table 1. The crystallite size ( $D$ ) increases progressively with annealing temperature. This enhancement can be attributed to the additional thermal energy that promotes recrystallization and re-nucleation, thereby improving the degree of crystallinity. Densification promotes the development of small nuclei during annealing, which eventually expand and unite to create bigger, more structured crystallites. As a result, numerous crystallites are generated, leading to an increased effective surface area for light absorption. This structural modification favors stronger light-matter interaction, thereby reducing the optical bandgap through enhanced absorption relative to transmission [29,30].

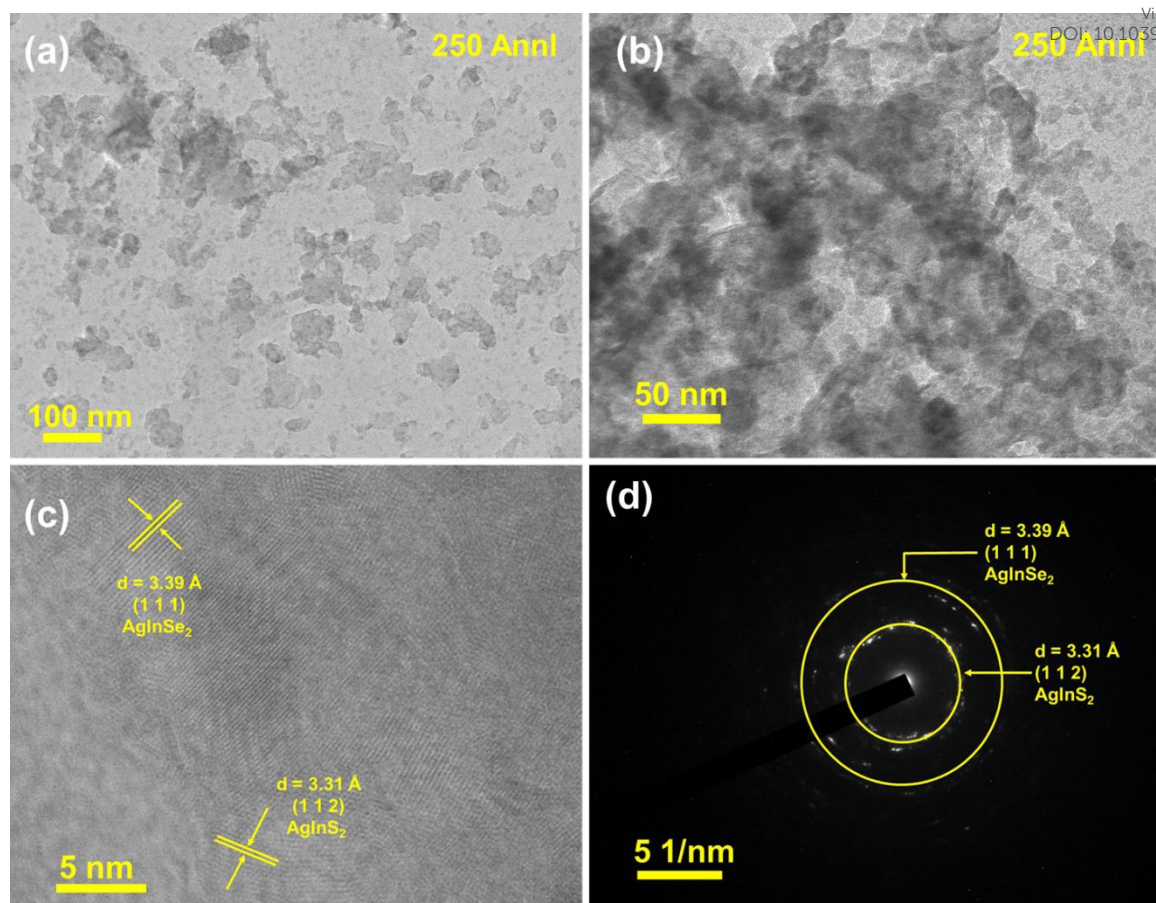
**Table 1:** Structural parameters of the  $\text{Ag}_2\text{S}/\text{In}_2\text{Se}_3$  films.

Structural parameters	Asp	150 °C	150 °C	200 °C	250 °C
$D$ (nm)	Amorphous	Amorphous	15.820	17.320	18.156
$\delta$ ( $\times 10^{15} \text{ m}^{-2}$ )	----	---	0.004	0.003	0.006
Lattice strain ( $\epsilon$ )	----	-----	0.010	0.009	0.007
$N_c$ ( $\text{nm}^{-2}$ )	----	----	0.151	0.115	0.100

### 3.2.TEM Study:

TEM pictures of the prepared films at two magnifications (100 nm and 50 nm, Fig. 2(a, b)) reveal that the material consists of well-dispersed nanostructures with uniform morphology. The nanoscale features are clearly distinguished, indicating good crystallinity and homogeneity across the sample. High-resolution TEM (HRTEM) analysis, presented in Fig. 2(c), provides further insight into the structural characteristics of the material. The lattice fringes exhibit two distinct interplanar spacings: 3.39 Å, which can be assigned to the (111) plane of the  $\text{AgInSe}_2$  phase, and 3.31 Å, corresponding to the (112) plane of  $\text{AgInS}_2$ . The coexistence of these phases suggests a mixed chalcogenide system, where both selenium and sulfur contribute to the crystalline framework. The SAED pattern, shown in Fig. 2(d), displays a series of sharp diffraction rings, confirming the polycrystalline nature of the material. The diffraction features are in close agreement with the XRD results, thereby validating the presence of multiple crystalline phases.

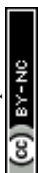




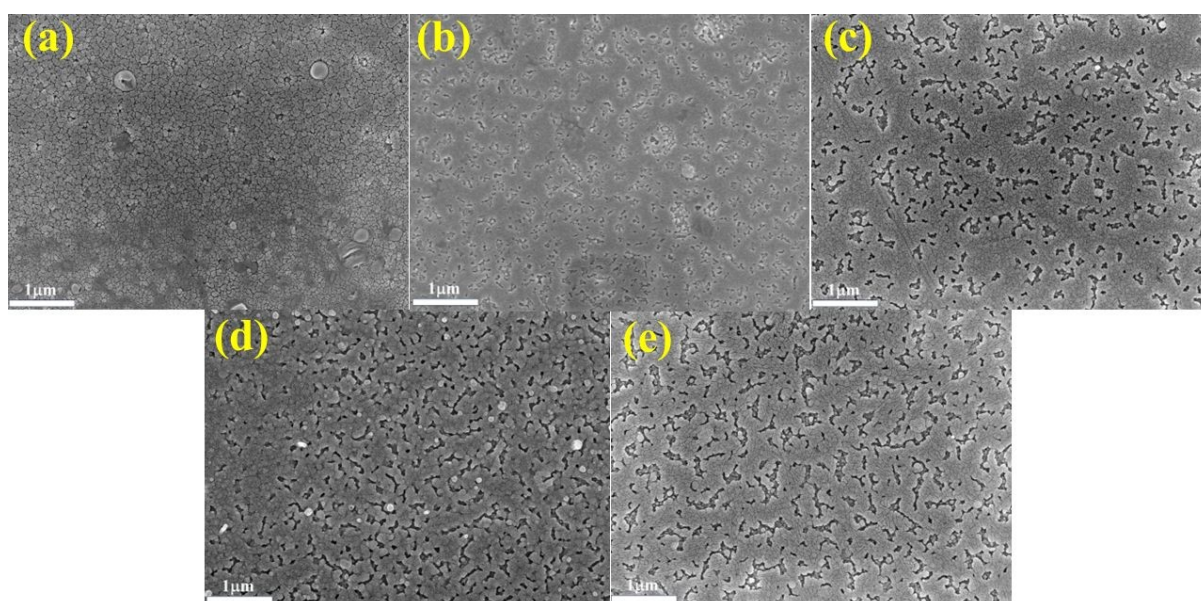
**Fig. 2.** TEM images at (a) 100 nm, (b) 50 nm, (c) HRTEM view, and (d) SAED pattern of 250 °C annealed  $\text{Ag}_2\text{Se}/\text{In}_2\text{Se}_3$  thin film.

### 3.3. FESEM and EDS study

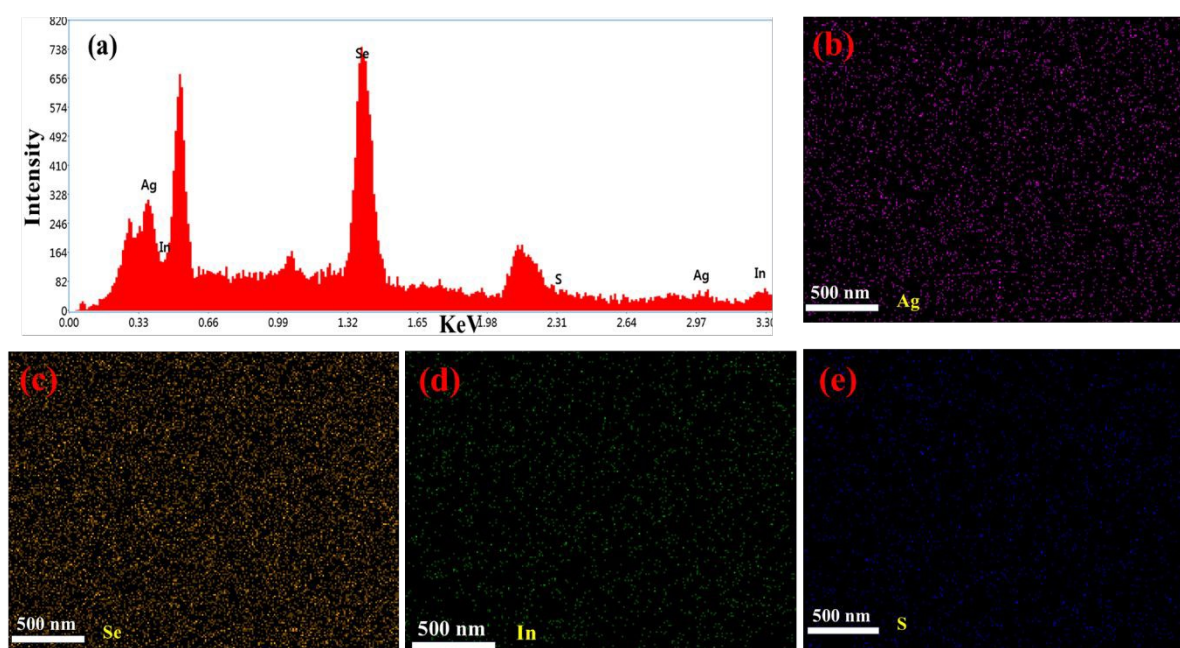
The morphology of film structure before thermal treatment and after thermal annealing is presented through the  $1\mu\text{m}$  FESEM images (Fig. 3). The as-prepared and 100 °C annealed film structures do not have as much porosity as the annealed one beyond 100 °C (Fig. 3a,b). The annealing induces porosity as seen in Fig. 3(c-e). The films consist of nanostructured particles. The degree of porosity is enhanced with the annealing conditions. The 100 nm scale film morphology shows the porosity in a better way, as presented in Fig. S1. The enhanced porosity is also seen from the TEM image (Fig. 2a, b). This increase in porosity originates from diffusion-induced phase transformation during annealing, where atomic intermixing between  $\text{Ag}_2\text{S}$  and  $\text{In}_2\text{Se}_3$  leads to grain boundary reorganization and void formation. Such porosity enhances the active surface area, facilitating improved light absorption and faster surface charge exchange processes, which together contribute to the enhanced photodetection performance observed at higher annealing temperatures.



The existence of elements after film preparation and annealing is confirmed through EDS data. The 250 °C annealed film shows the different elemental peaks, inferring the existence of related elements. The unassigned peaks refer to the carbon and gold peaks. Carbon tape and gold coating were used to avoid charging of the film, as the films were coated on a glass substrate. Fig. 4(a) shows the peaks for Ag, In, S, and Se, thus confirming the existence of elements in the film after annealing. The EDS pictures of the pristine and other annealed samples are depicted in Fig. S2. The elemental distribution of elements infers the uniformity of the film as illustrated in Fig. 4(b-e). The Atomic% values obtained from the EDS peaks of Ag, In, S and Se for each thin film sample shown in Table S1.



**Fig. 3.** FESEM view of (a) as-prepared, (b) 100 °C, (c) 150 °C, (d) 200 °C, (e) 250 °C annealed film.

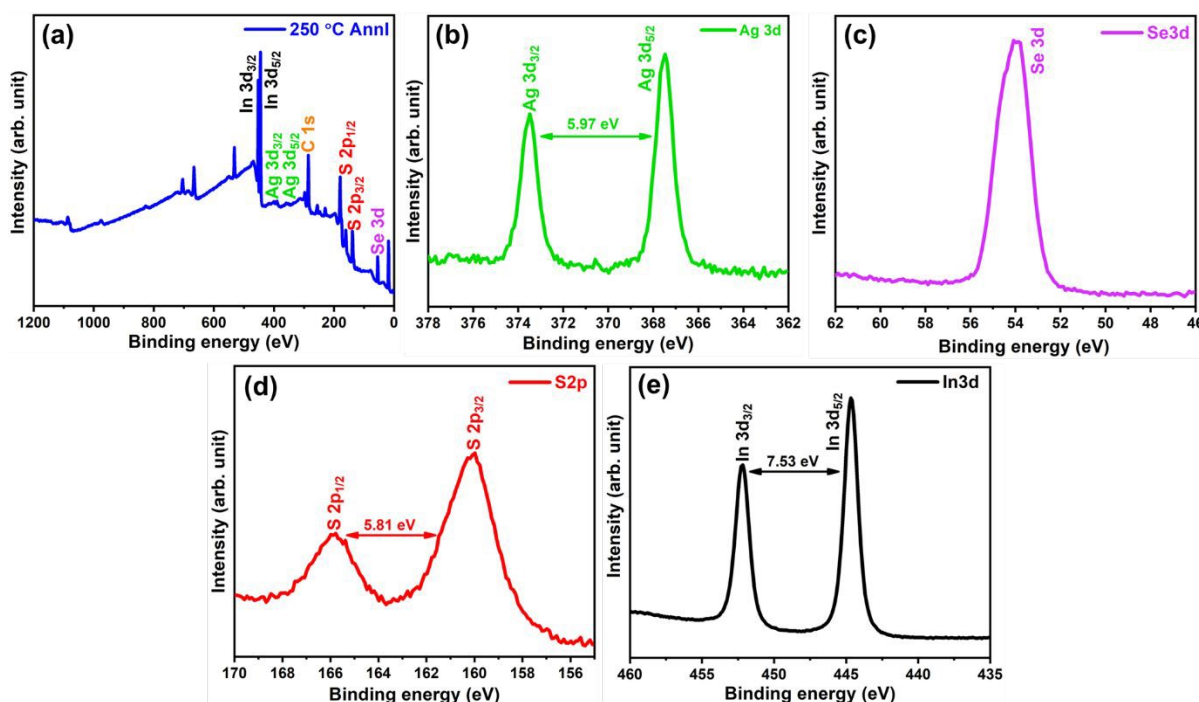




**Fig. 4.** (a) EDS spectra, elemental mapping of (b) Ag, (c) Se, (d) In, and (e) S in a 250 °C annealed film.

### 3.4. XPS Study:

XPS serves as an indispensable tool for characterizing the elemental composition of materials, particularly in the study of  $\text{Ag}_2\text{S}/\text{In}_2\text{Se}_3$  films. The survey spectra of thin films annealed at 250 °C, presented in Fig. 5(a), clearly display characteristic photoelectron peaks. These peaks correspond to C 1s, In 3d, S 2p, Ag 3d, and Se 3d, providing an initial identification of the elements present within the film. A more precise understanding of the elemental states can be derived from the analysis of their respective core-level spectra.



**Fig. 5.** XPS spectra of (a) survey spectra, (b) Ag 3d, (c) Se 3d, (d) S 2p, and (e) In 3d spectra of 250 °C annealed  $\text{Ag}_2\text{S}/\text{In}_2\text{Se}_3$  thin film.

Further detailed analysis of the core-level spectra reveals specific chemical environments for each element. As depicted in Fig. 5(b), the Ag 3d core level spectrum exhibits a doublet, characteristic of spin-orbit splitting, with peaks observed at approximately 367.46 eV (Ag 3d<sub>5/2</sub>) and 373.50 eV (Ag 3d<sub>3/2</sub>) with a binding energy of 5.97 eV [31]. These binding energies are indicative of metallic silver (Ag (0)). Furthermore, the Se 3d peak, presented in Fig. 5(c) at around 53.95 eV, signifies the spin-orbit coupling characteristics typical for selenium in this particular thin film. Similarly, the S 2p peak, shown in Fig. 5(d), presents a doublet structure at around 160.12 eV (S 2p<sub>3/2</sub>) and 165.83 eV (S 2p<sub>1/2</sub>), confirming the presence of sulfur in a specific chemical state within the film and having the binding energy of 5.81 eV. Finally, the



In 3d spectrum, illustrated in Fig. 5(e), displays two distinct peaks at approximately 444.61 eV and 452.21 eV, corresponding to the In 3d<sub>5/2</sub> and In 3d<sub>3/2</sub> components, respectively. These values are consistent with indium in a compound form. These detailed core-level analyses collectively provide critical insights into the chemical states and bonding environments of the constituent elements in the Ag<sub>2</sub>S/In<sub>2</sub>Se<sub>3</sub> thin films [32-34].

### 3.5. Contact angle study

The surface wettability plays a vital role in determining the surface-related properties for different applications. The solid-liquid interface is an important aspect in applied sciences. It is evaluated based on the contact angle measurement, which provides an idea about the type of surface, whether hydrophilic or hydrophobic. The tangent (angle) of a liquid drop with a solid contact at the base determines the contact angle value ( $\theta_c$ ). The value of  $\theta_c$  defines the nature of the surface. The angle value greater than 90° is categorized as a hydrophobic one, while the value less than 90° is regarded as a hydrophilic one [35]. The decrease in angle in hydrophilic surfaces makes the sample a super-hydrophilic one, whereas the increase in angle in hydrophobic surfaces takes it to a super-hydrophobic one. The obtained  $\theta_c$  are illustrated in Fig. 6, which infers a decrease in the contact angle. This infers the decrease in hydrophobicity, which indicates the decrease in surface roughness. This suggests the water-repelling ability gradually decreased with annealing. The reason for this is the formation of pores in the annealed film with increased temperatures (FESEM image). Young's relation is used to estimate the surface energy associated with the sample by taking the surface tension of water (71.99 mN/m) [36].

$$\gamma_{se} = \frac{\gamma_w (1 + \cos \theta_c)^2}{4} \quad (2a)$$

The evaluated values in Table 2 indicate the increased surface free energy with heat treatment. The value of  $\gamma_{se}$  was 7.791 mN/m for the bilayer film, which increased to 12.289 mN/m after 250 °C annealing. The work of adhesion ( $W_{sl}$ ) was calculated from Young–Dupré equation [37,

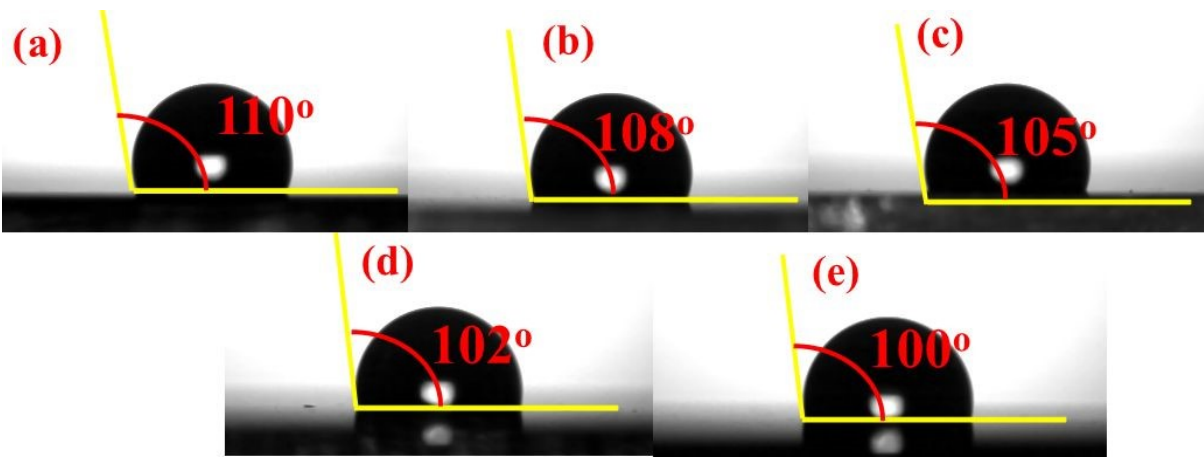
$$W_{sl} = \gamma_w (1 + \cos \theta) \quad (2b)$$

The  $W_{sl}$  for all the films was found to be increasing from 47.367 to 59.489 mN/m upon heat treatment (Table 2). The retention of the hydrophobic nature of the films is helpful for different applications due to their self-cleaning nature, greater chemical stability, and less degradation. For such important properties, there are potential applications in various sectors involving heat transfer applications, biomedical devices, water treatment, and so on [38]. In the context of photodetection, the gradual decrease in contact angle with annealing signifies enhanced surface





activity and oxygen adsorption–desorption capability, which directly influence carrier dynamics and photocurrent generation efficiency in  $\text{Ag}_2\text{S}/\text{In}_2\text{Se}_3$  films.



**Fig.6.** Contact angle images for (a) as-prepared, (b)100 °C, (c) 150 °C, (d) 200 °C, (e) 250 °C annealed film.

**Table 2.** Parameters associated with contact angle

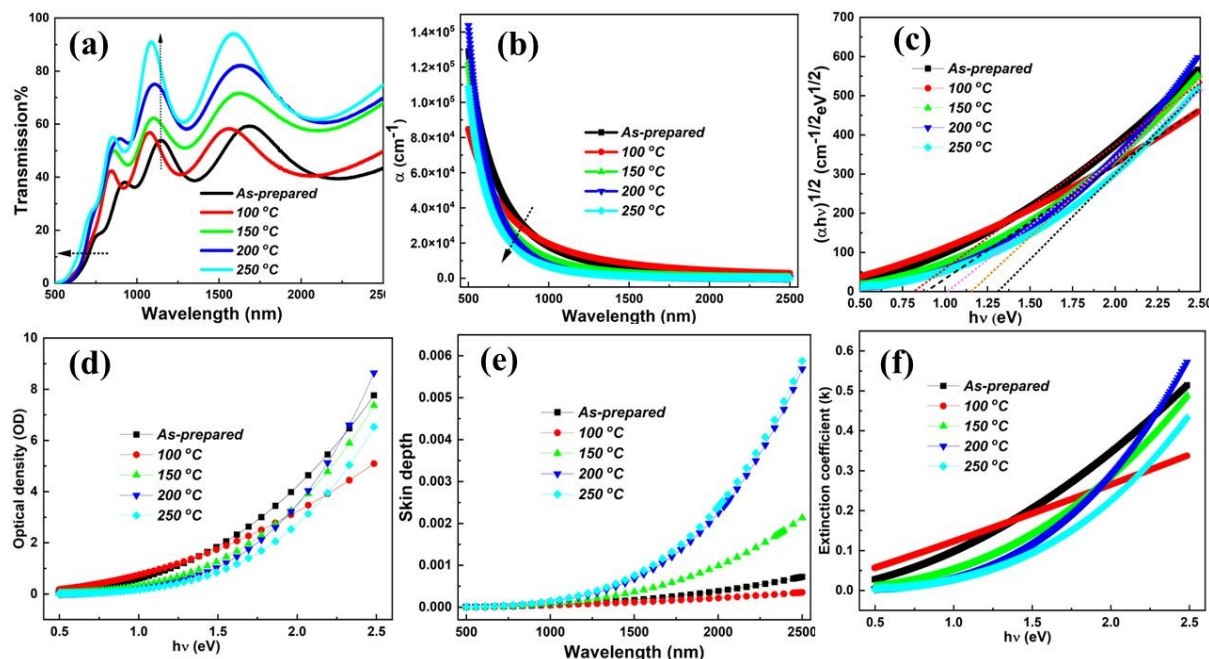
Parameters/irradiation (min)	As-prepared	100°C	150°C	200°C	250°C
$\theta_c$ (degree)	110°	108°	105°	102°	100°
$\gamma_{se}$ (mN/m)	7.791	8.593	9.886	11.291	12.289
$W_{sl}$ (mN/m)	47.367	49.743	53.357	57.022	59.489

### 3.6. Optical data (UV-Visible study)

The optical response of the films under study is reflected through their transmittance curve, as shown in Fig. 7(a). The systematic increase in the transmission power through annealing is well observed from Fig.7(a). The unannealed film transmittance of ~60% is enhanced to ~95% at 1750 nm. Such increased value shows the improved transparency in the film by annealing. The appearance of wave-type fringes is the signature of uniform and homogeneous film quality [39]. The change in gap states by annealing resulted in an increased magnitude of the transmittance. The concentration of the defects, along with disorder, changed by such annealing [40]. The magnitude of transmittance is less below 1000 nm, which is favorable for solar cell applications. The higher transparency after 1000 nm wavelength range is the result of free electron interaction with the incident wave. It results in the polarization change, which is good for temperature sensors as well as IR lenses [41]. Simultaneously, the absorption edge shift to the lower wavelength side indicates the reduction in its energy gap as calculated through its absorption coefficient  $(\alpha) = \left(\frac{1}{t}\right)\ln\left(\frac{1}{T}\right)$  [42]. Here, T stands for transmittance, and ‘t’ is the film



thickness. Fig. 7(b) infers the reduction in absorption value upon annealing. The physical quantity is essential for evaluating other necessary constants. The stability in the absorption coefficient at higher wavelengths signifies the normal dispersion behaviour.



**Fig. 7.** (a) Transmission (b) absorption coefficient (c) bandgap (d) optical density (e) skin depth (f) extinction coefficient changes with annealing conditions.

The bandgap ( $E_g$ ) depends on the value of ' $\alpha$ ' along with the incident photon energy through the Tauc relation, [43],  $(\alpha hv) = C(hv - E_g)^w$  (3)

The kind of transition between the valence and conduction bands depends on the exponent ' $w$ ' value. The direct and indirect allowed transition follows  $w=1/2$  and 2. The direct and allowed forbidden transition follows  $w=3/2$  and 3. The present data follows an indirect allowed transition as illustrated in Fig. 7(c). The linear fitting of the data evaluates the  $E_g$  value from its X-intercept, and the slope determines the Tauc parameter. The individual plot is illustrated in Fig. S3. The increased  $E_g$  value for the annealed film in comparison to the bilayer film shows the intermixing effects of the two layers influencing the optical bandgap [44]. The bilayer film with  $E_g$  value 0.816 eV showed an increased band gap to 1.311 eV after 250 °C annealing. Such a change is explained through Davis and Mott, the so-called "density of state model" [45]. The enhanced  $E_g$  value is because of the improved structural change by annealing. This is reflected in the XRD data. Such changes are due to the decreased disorder as measured through the Urbach energy ( $E_u$ ) and Tauc constant ( $B^{1/2}$ ). The  $E_u$  value is determined from the Urbach relation [46]  $\alpha = \alpha_0 \exp\left(\frac{hv}{E_u}\right)$ . The absorption coefficient at  $E_g$  point is represented by  $\alpha_0$ .



The linear fit of the absorbance data at the low absorption regime with photon energy evaluates the  $E_u$  value. From the calculation, the  $E_u$  value for the bilayer structure was found to be 543 meV, which was reduced to 432 meV for the 250 °C annealed film. The other values are mentioned in Table 3. Such reduced strength is the signature of increased structural order or decreased disorder [46]. The other parameter  $B^{1/2}$  is inversely related to  $E_u$ , which shows the increased value as mentioned in Table 3 [47]. The annealing-induced saturated bond formation resulted in the reduction of defect states [48]. This implies an improvement in structure in terms of microstructural order inside the film.

**Table 3.** Optical parameters of the investigated  $Ag_2S/In_2Se_3$  films.

Optical parameter	As-prepared	100°C	150 °C	200 °C	250 °C
Optical band gap (eV)	0.816	0.903	1.017	1.162	1.311
Tauc parameter ( $cm^{-1/2}eV^{1/2}$ )	243	267	298	321	344
Urbach Energy $E_u$ (meV)	543	512	489	465	432
$\sigma$ ( $\times 10^{-2}$ )	4.761	5.048	5.286	5.559	5.983
$S_{e-p}$	14.001	13.205	12.610	11.991	11.141
$(\eta_{opt})$	1.633	1.645	1.661	1.676	1.691
$(\epsilon_{\infty})$	12.852	12.117	11.303	10.445	9.715
$(n_0)$	3.585	3.481	3.362	3.232	3.117

The other parameters, the steepness parameter ( $\sigma$ ) as well as electron-phonon interaction strength ( $S_{e-p}$ ) and are associated with  $E_u$  [49].

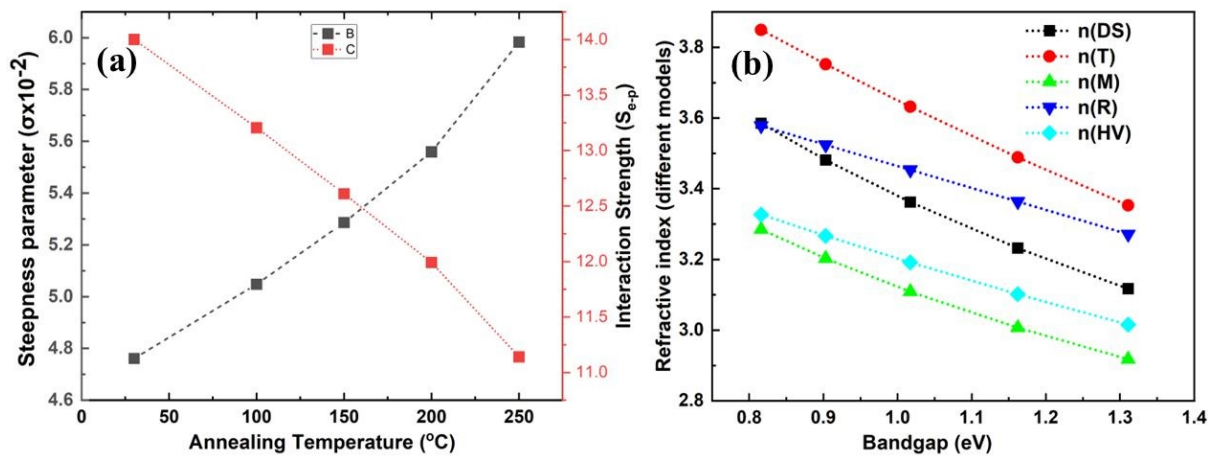
$$\sigma = \frac{K_B T}{E_u} \text{ and } S_{e-p} = \frac{2}{3\sigma} \tag{4}$$

In Fig. 8(a), the increased value of  $\sigma$  from  $4.761 \times 10^{-2}$  to  $5.883 \times 10^{-2}$  (Table 3) suggests a steeper absorption edge with less disorder. On the other hand,  $S_{e-p}$  reduced from 14.001 to 11.141 with annealing. It refers to the weakening of electron-phonon coupling. The reduced  $S_{e-p}$  indicates the minimised lattice vibrations. It leads to an increment in lattice stability and structural order. It ultimately accounts for the improvement in optical quality with less disorder.

Optical density (OD) presents the absorption strength of the films. It is determined from the relation,  $OD = \alpha \times t$ . The variation in OD with ‘ $h\nu$ ’ at various annealed conditions is illustrated in Fig.7(d). Its value decreases for the annealed film at higher temperatures. The skin depth or penetration depth ( $\delta$ ) is another important parameter for optical quality determination. It is equal to  $1/e$  times the photon density at the surface state. It is associated with semiconducting materials conductivity. The value is determined by the reciprocal of ‘ $\alpha$ ’ ( $\delta = 1/\alpha$ ) [49]. The increased skin depth with photon energy, as well as annealing temperature, results from the

low absorbance, which enhances the transmittance (Fig. 7e) [42]. The extinction coefficient ( $k = \frac{\alpha\lambda}{4\pi}$ ) measures the loss of light during crossing through the sample by scattering and absorption. The change in  $k$  various annealing states is presented in Fig. 7(f). It presents a lower value for the annealed film over the as-prepared bilayer film. The 'k' values increased with photon energy and decreased with annealing temperature.

These optical property trends with annealing confirm that the improvement in transparency, bandgap, and optical constants arises from enhanced crystallinity, reduced defect states, and the formation of well-ordered  $\text{AgInS}_2/\text{AgInSe}_2$  phases, all of which play a crucial role in optimizing the light absorption and charge transport behavior of the thin film photodetector.



**Fig. 8.** (a) Variation of  $S_{e-p}$  and ' $\sigma$ ' with annealing temperature (b) Refractive index at different bandgap values through various models.

The high-frequency dielectric parameter is evaluated by  $\epsilon_L = n^2$ . In the present study,  $\epsilon_L$  increased with irradiation time (Table 3). It is attributed to the electrical polarizability of the material and is connected with free charge carriers affecting polarization [50]. The optical conductivity ( $\eta_{opt}$ ) is determined by,  $\eta_{opt} = \left(\frac{c}{n_0}\right)^{1/4}$ , with  $c = 25.54$  and are given in Table 3 [50]. It measures the affinity of an atom for bond formation. In the present case, it increased from 1.633 (un-annealed) to 1.691 for 250°C annealed film (Table 3).

The refractive index ( $n$ ) evaluates the dispersion of light waves and other nonlinear optical phenomena. The calculation of ' $n$ ' from  $E_g$  by using the Dimitrov and Sakka empirical relation establishes a useful connection between the electronic band state and the optical dispersion of the films [52]

$$\frac{n^2-1}{n^2+2} = 1 - \sqrt{\frac{E_g}{20}} \quad (5a)$$



With increasing annealing temperature, the refractive index ( $n$ ) exhibited a notable decrement (Table 3 and Fig. 7(b)). In the present study,  $n$  values were found to vary between 3.58 and 3.11, suggesting their suitability for applications in integrated photonic devices, ultrafast optical modulators, infrared (IR) sensors, and related optoelectronic systems [53]. The opposite variation of “ $E_g$ ” with “ $n$ ” follows

Moss’s rule, i.e.,  $E_g n^4 \sim \text{constant}$  [54]. The  $E_g n^4 \sim \text{constant}$  leads to as,  $n_M = \sqrt[4]{\frac{95}{E_g}}$ , (5b)

with ‘ $k$ ’ = 95 eV. The rectified Ravindra’s formula  $n_R = 4.084 [0.62 \times E_g]$  is used [55] for the evaluation of ‘ $n$ ’ below 4 eV materials. Likewise, for low  $E_g$  films, Herve-Vandamme [56] relation is used, which is in accordance with oscillator theory, expressed as,  $n^2 = 1 + \left(\frac{A}{E_g + B}\right)^2$ . (5c)

With,  $A \cong 13.6 \text{ eV}$ , ( $\sim$ hydrogen ionisation energy), and  $B = 3.4 \text{ eV}$ . Considering these values, the above

relation becomes,  $n_{[HV]} = \sqrt{1 + \left(\frac{13.6}{E_g + 3.47}\right)^2}$ . (5d)

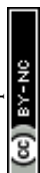
The exponential relation of ‘ $n$ ’ for semiconductors with band gaps is given by Tripathy [57], the relation,  $n_T = 1.73 \times [1 + 1.9017 \times e^{-0.539 \times E_g}]$  (5e)

The values obtained from different models are expressed in Fig. 8(b). It presents the inverse relation with  $E_g$ .

### 3.7. Photo response study

The I–V features for all films under illuminated and dark states are illustrated in Fig. 9(f, g), while Fig. S4 presents corresponding plots with a linear Y-axis. These curves exhibit clear symmetry for each thin film, with the photocurrent consistently surpassing the dark current, indicating effective light-induced charge separation and material interaction. [58] From the linear plotting in Fig. S4, the peak current values under light ( $I_L$ ) and dark ( $I_D$ ) states were extracted and illustrated in Table 4. The photosensitivity of each film was calculated by the formula in equation (6a). All thin films displayed significant photosensitivity, with the sample annealed at 250°C exhibiting the highest photodetection capability. This enhancement is linked to improved crystallinity at higher annealing temperatures, which boosts photoresponsivity and detectivity. [59]

Photodetector performance under various light intensities is basically driven by the formation of electron-hole pairs proportional to the power of incident light. Incorporating an  $\text{Ag}_2\text{S}$  layer on  $\text{In}_2\text{Se}_3$  increases carrier concentration in the  $\text{In}_2\text{Se}_3$ , enhancing the overall response. Additionally, surface interactions involving oxygen play a critical role: in the dark, oxygen molecules adsorb onto the surface, capturing electrons ( $\text{O}_2 + e^- \rightarrow \text{O}_2^-$ ), while under illumination, photogenerated holes trigger desorption of oxygen, freeing electrons to conduction band and elevating photocurrent [23]. Key performance indicators such as





photoresponsivity (R) and detectivity (D) were calculated by standard formulas and are also detailed in Table 4[1,60-62]

$$\text{Photo Sensitivity} = \frac{I_L - I_D}{I_L} \times 100\% \quad (6a)$$

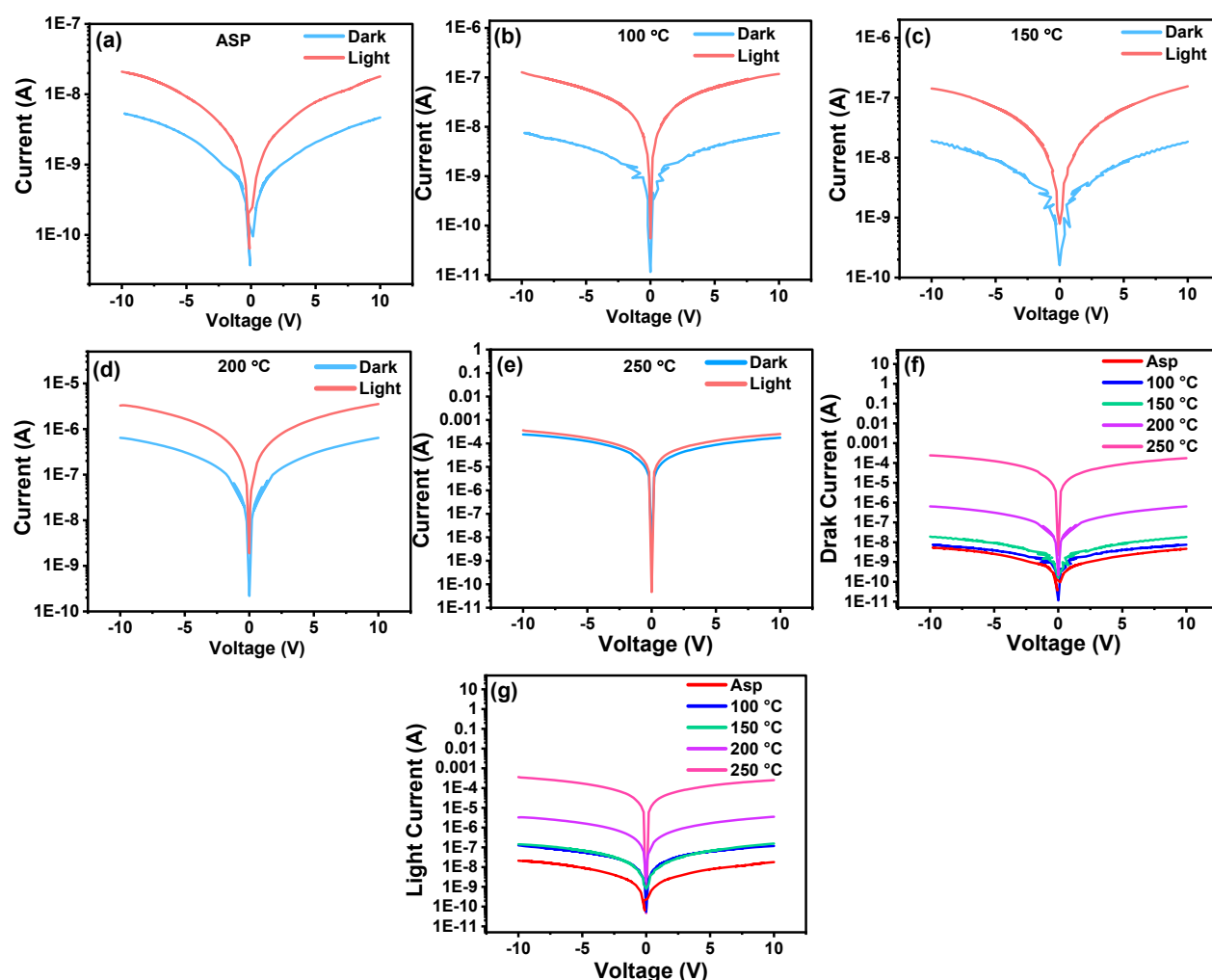
$$\text{Responsivity (R)} = \frac{I_L}{A \times P_{in}} \quad (6b)$$

$$\text{Detectivity (D}^*) = R \sqrt{\frac{A}{2eI_D}} \quad (6c)$$

Here, A = Active surface area = 0.25 cm<sup>2</sup>,

P<sub>in</sub> = Incident light's power density = 20 mWcm<sup>-2</sup>,

e = 1.602 × 10<sup>-19</sup> C = electronic charge



**Fig. 9.** The current–voltage (I–V) characteristics of Ag<sub>2</sub>S/In<sub>2</sub>Se<sub>3</sub> thin films, presented on a logarithmic scale under both dark and illuminated conditions. Individual plots are displayed for (a) the as-prepared sample, and those annealed at (b) 100°C, (c) 150°C, (d) 200°C, and (e) 250°C. Additionally, combined I–V curves for all samples are provided under (f) dark and (g) light conditions.



**Table 4.** The figure of merit of all the as-prepared and annealed Ag<sub>2</sub>S/In<sub>2</sub>Se<sub>3</sub> thin films

Sample	I <sub>L</sub> (A)	I <sub>D</sub> (A)	I <sub>L</sub> - I <sub>D</sub> (A)	Photo sensitivity (%)	Responsivity (R) (AW <sup>-1</sup> )	Detectivity (D*) (Jones)
Asp	1.79 × 10 <sup>-8</sup>	4.66 × 10 <sup>-9</sup>	1.32 × 10 <sup>-8</sup>	73.96	3.58 × 10 <sup>-6</sup>	3.09 × 10 <sup>8</sup>
100°C	1.17 × 10 <sup>-7</sup>	7.5 × 10 <sup>-9</sup>	1.09 × 10 <sup>-7</sup>	93.61	2.34 × 10 <sup>-5</sup>	2.89 × 10 <sup>7</sup>
150°C	1.55 × 10 <sup>-7</sup>	1.85 × 10 <sup>-8</sup>	1.36 × 10 <sup>-7</sup>	88.06	3.10 × 10 <sup>-5</sup>	1.95 × 10 <sup>8</sup>
200°C	3.53 × 10 <sup>-6</sup>	6.42 × 10 <sup>-7</sup>	2.89 × 10 <sup>-6</sup>	81.84	7.07 × 10 <sup>-4</sup>	7.32 × 10 <sup>9</sup>
250°C	2.51 × 10 <sup>-4</sup>	1.75 × 10 <sup>-4</sup>	7.66 × 10 <sup>-5</sup>	30.43	5.03 × 10 <sup>-2</sup>	1.44 × 10 <sup>7</sup>

The I–V characteristics of Ag<sub>2</sub>S/In<sub>2</sub>Se<sub>3</sub> thin films, under both dark and illuminated states, demonstrate clear symmetry and a consistent increase in current under illumination across all samples. The logarithmic plots (Fig. 9a–e) show enhanced photocurrent (I<sub>L</sub>) compared to dark current (I<sub>D</sub>), with the 250°C annealed sample exhibiting the highest absolute current levels. However, Fig. S4 (with a linear y-axis) reveals that while the 250°C sample produces the highest I<sub>L</sub> and I<sub>D</sub> values, the contrast between light and dark currents, indicative of sensitivity, is reduced compared to other annealed conditions. Quantitative analysis from Table 4 supports this observation. The highest photosensitivity (93.61%) is recorded for the 100°C sample, followed closely by the 150°C and 200°C samples. The 250°C sample, despite showing the largest I<sub>L</sub> (2.51 × 10<sup>-4</sup> A), demonstrates the lowest relative photosensitivity (30.43%) due to a high baseline dark current. Nevertheless, its responsivity (5.03 × 10<sup>-2</sup> A/W) and detectivity (1.44 × 10<sup>7</sup> Jones) remain significant, highlighting its utility in applications requiring high absolute photocurrent rather than contrast.

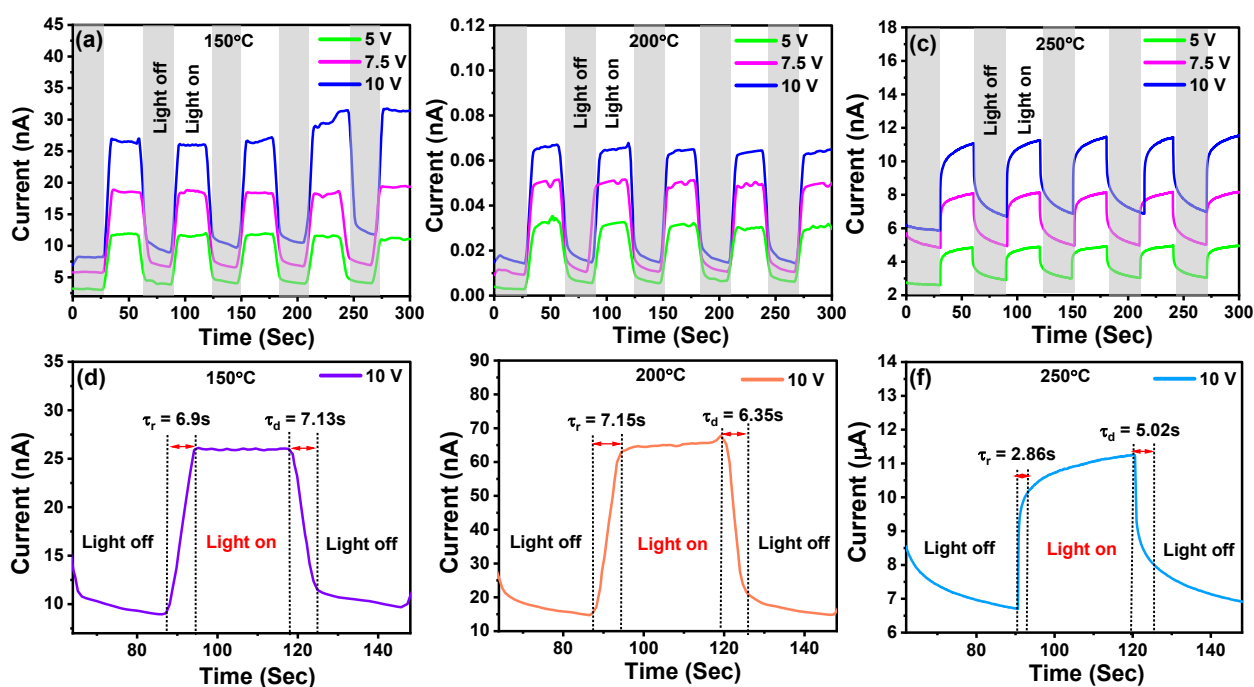
Overall, moderate annealing temperatures (100–250°C) improve photodetection performance by optimizing light-induced carrier generation and suppressing dark current. The 200°C sample strikes a balance, offering enhanced current response, high responsivity (7.07 × 10<sup>-4</sup> A/W), and the best detectivity (7.32 × 10<sup>9</sup> Jones), attributed to improved crystallinity and reduced surface recombination.

A comparative analysis of the films annealed at different temperatures reveals that the photoresponsivity (R) correlates strongly with the films’ structural and optical properties. As the annealing temperature increases from 100 °C to 200 °C, enhanced crystallinity, improved interfacial uniformity, and reduced defect density collectively enhance carrier mobility and light absorption, leading to higher R and D values. However, further annealing to 250 °C introduces excess charge carriers and defect-assisted conduction pathways, increasing dark current and reducing photosensitivity. Thus, the 200 °C annealed film achieves the optimal



balance between crystallinity, defect control, and interfacial quality, resulting in the highest responsivity and detectivity among all samples.

However, the rise in dark current at 250 °C indicates that excessive annealing increases carrier concentration and defect-assisted conduction, thereby reducing photosensitivity. This demonstrates that optimal photodetector performance arises not from maximum current output but from a balanced combination of crystallinity, defect density, and interfacial integrity conditions best achieved at 200 °C. Table S2 compares the key photodetection parameters of the present  $\text{Ag}_2\text{S}/\text{In}_2\text{Se}_3$  device with other 2D/ $\text{In}_2\text{Se}_3$ -based systems reported in the literature. The results demonstrate that the optimized  $\text{Ag}_2\text{S}/\text{In}_2\text{Se}_3$  heterostructure exhibits comparable responsivity and detectivity, while offering a broader visible spectral response due to efficient interfacial charge transfer and reduced recombination losses.



**Fig. 10.** Temporal photoresponse curves illustrating the change in current over time at bias voltages of 5, 7.5, and 10 V for annealed thin films at (a) 150°C, (b) 200°C, and (c) 250°C. The rise and fall times of the photocurrent within a selected time interval are shown for the films annealed at (d) 150°C, (e) 200°C, and (f) 250°C, respectively.

Fig. 10(a-c) displays I-T curves for  $\text{Ag}_2\text{S}/\text{In}_2\text{Se}_3$  films annealed at 150°C, 200°C, and 250°C under bias voltages of 5, 7.5, and 10 V, showing distinct on/off light cycles. All samples exhibit a sharp increase in current upon illumination, followed by a quick drop when the light is off, driven by enhanced carrier drift velocity. The increased photocurrent with higher bias voltages is attributed to stronger electric fields, which improve charge separation, reduce recombination,



and expand the depletion region, leading to more efficient carrier generation and collection. Variations in maximum current among samples are likely influenced by defect states. After the light is off, the photocurrent decreased due to recombination between trapped and free carriers, with additional effects from non-radiative processes and carrier transport variations. For consistent analysis, a specific peak within the 60-150 second interval is selected for all samples Fig. 10(d–f). From this peak, key parameters such as light-on ( $I_{on}$ ), light-off ( $I_{off}$ ), and  $I_{on}/I_{off}$  ratios during rise and decay phases are extracted and presented in Table 2 to evaluate photodetector performance.

Furthermore, the transient photoresponse (I-T) curves in Fig. 10(a-c) show that the photocurrent quickly returns to its baseline after each illumination cycle, indicating the absence of persistent photoconductivity. This behavior reflects efficient carrier recombination and minimal trap-assisted retention, confirming the fast and reversible photoresponse characteristics of the  $Ag_2S/In_2Se_3$  heterostructure.

**Table 5.** Photodetector parameters of the  $Ag_2S/In_2Se_3$  thin films for both rise and decay conditions.

Sample	During Rise				During Decay			
	$I_{on}$ (nA)	$I_{off}$ (nA)	$I_{on}/I_{off}$	$\tau_r$ (s)	$I_{on}$ (nA)	$I_{off}$ (nA)	$I_{on}/I_{off}$	$\tau_d$ (s)
150°C	25.88	9.08	2.85	6.9	25.92	11.26	2.3	7.13
200°C	62.35	15.72	3.96	7.15	67.74	21.25	3.18	6.35
250°C	10.14	6.71	1.51	2.86	11.24	8	1.4	5.02

Rise time ( $\tau_r$ ) and decay time ( $\tau_d$ ) are essential indicators of a photodetector’s responsiveness. The rise time measures how quickly the current increases (from 10% to 90%) when light is applied, while the decay time reflects how fast it returns to baseline (from 90% to 10%) after light is removed. Shorter  $\tau_r$  and  $\tau_d$  values indicate faster response and recovery, which are critical for high-speed light detection. [5] These parameters, extracted from Fig. 10(d–f), are presented in Table 5 to evaluate the dynamic performance of the  $Ag_2S/In_2Se_3$  photodetectors.

Fig. 10 and Table 5 present the temporal photo response behavior of  $Ag_2S/In_2Se_3$  films annealed at 150°C, 200°C, and 250°C under different bias voltages (5, 7.5, and 10 V). All samples show a repetitive and reversible change in current upon cyclic illumination, indicating stable photodetector operation. The response amplitude increases with higher bias voltage, especially evident in the 200°C sample (Fig. 10b), which exhibits the highest photocurrent. The



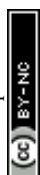
extracted photodetector parameters (Table 5) reveal that the 200°C sample demonstrates the most favourable characteristics overall. It shows the highest photocurrent during both the rise (62.35 nA) and decay (67.74 nA) phases, coupled with a strong  $I_{\text{on}}/I_{\text{off}}$  ratio of 3.96 (rise) and 3.18 (decay). Its rise and fall times ( $\tau_r = 7.15$  s,  $\tau_d = 6.35$  s) are moderate and well-balanced, suggesting efficient charge separation and recombination kinetics. The optimal dynamic behavior at 200 °C arises from a well-balanced microstructure where enhanced crystallinity and controlled interdiffusion reduce trap-assisted recombination, thereby enabling faster carrier transport and recovery compared to under or over-annealed films.

Fig. S5 illustrates the current–time (I-T) stability behavior of  $\text{Ag}_2\text{S}/\text{In}_2\text{Se}_3$  thin films annealed at 150 °C, 200 °C, and 250 °C under continuous illumination for 1 hour (3600 sec.). All three samples exhibit a steady photocurrent response with minimal fluctuation, confirming good temporal stability. The photocurrent initially increases slightly before reaching saturation, which can be attributed to the gradual stabilization of charge trapping and de-trapping processes at the film–electrode interface. The 200 °C annealed sample shows the most stable and consistent current output, with negligible drift over time, indicating robust charge transport and minimal defect-mediated recombination. In contrast, the 250 °C film exhibits higher absolute current but with a slow saturation trend due to increased defect density. Overall, the I-T stability analysis confirms that the  $\text{Ag}_2\text{S}/\text{In}_2\text{Se}_3$  heterostructures maintain stable photocurrent characteristics over prolonged operation, particularly for the optimally annealed 200 °C sample.

In contrast, the 250°C sample exhibits the fastest response ( $\tau_r = 2.86$  s) and decay ( $\tau_d = 5.02$  s) but suffers from reduced current contrast ( $I_{\text{on}}/I_{\text{off}} \approx 1.5$ ), indicating lower sensitivity. Meanwhile, the 150°C sample shows a moderate current level and contrasts with longer response times ( $\tau_r = 6.9$  s,  $\tau_d = 7.13$  s), suggesting slower carrier dynamics. Overall, the 200°C annealed film stands out as the optimal candidate, combining high photocurrent, strong on/off contrast, and reasonable response times, ideal attributes for practical photodetector applications. Although wavelength-dependent photoresponse measurements could not be performed due to the unavailability of a monochromatic light source, the use of a broad-spectrum white LED ensures uniform excitation across the visible region. The strong absorption and rapid transient response observed confirm the film's potential for broadband visible photodetection. Future work will include detailed spectral responsivity analysis using tunable light sources.

#### 4. Computational Study

The structural stability, electronic structure, and optical behavior of the heterostructure layer

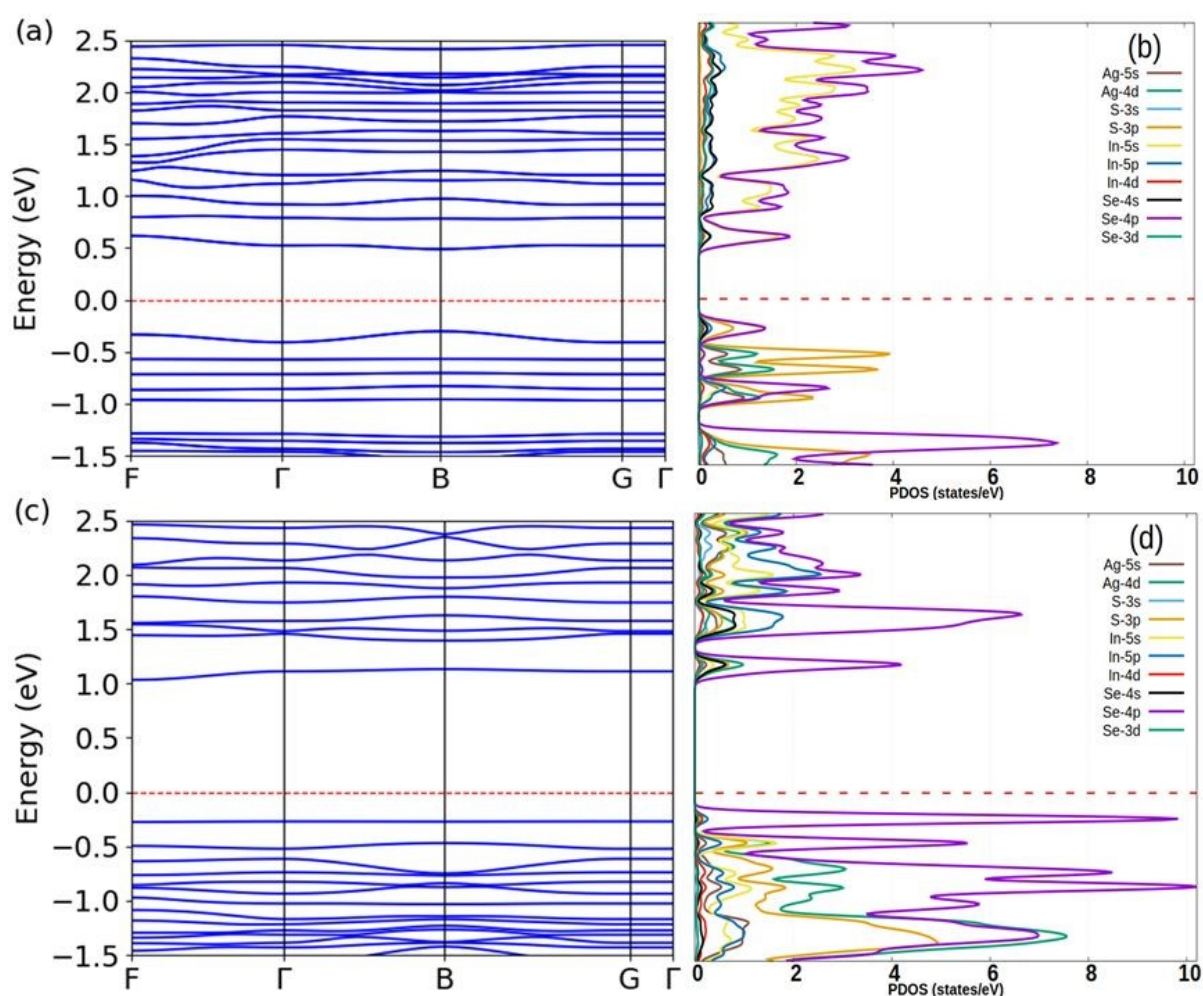




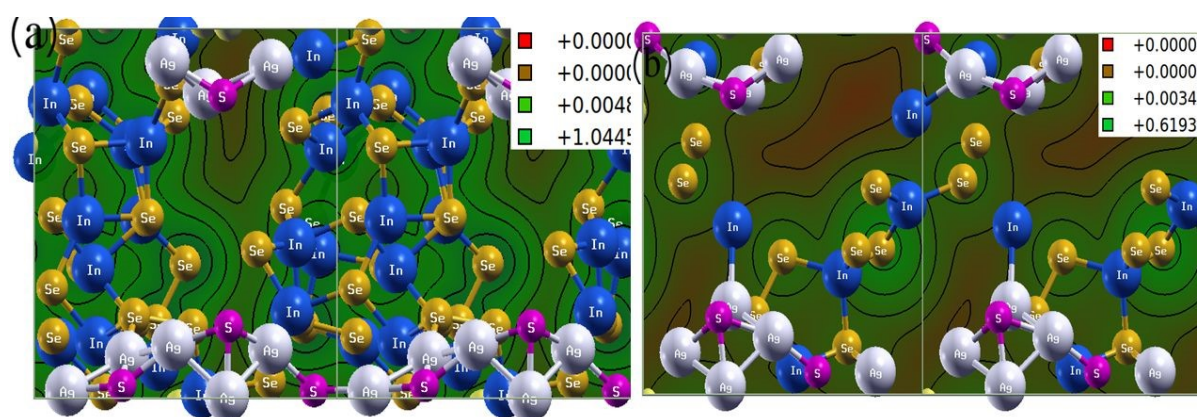
of  $\text{Ag}_2\text{S}$  and  $\text{In}_2\text{Se}_3$  with different proportions (1:3 and a single layer of Ag-S-Se-In) have been calculated by the full potential linearized augmented plane wave (FP-LAPW) within the framework of DFT through MedeAVASP package [63,64]. DFT was used to obtain an understanding of atomistic-scale interactions in determining the various physical properties of a material. The heterostructures were constructed along the (1 0 0) plane of  $\text{Ag}_2\text{S}$  and the (2 -1 0) of the  $\text{In}_2\text{Se}_3$  with 15Å of vacuum region given along the z-direction by using VESTA visualisation software [65]. The correlation and exchange energies were evaluated within GGA by Perdew-Burke-Ernzerhof (PBE) parametrization [66]. Finally, the Hubbard correction value,  $U_d=4$  for Ag and  $U_p=10$  for In, Se, and S atoms, was added to enhance the localization and short-range interactions. The self-consistent calculations were performed with a self-consistent field (SCF) energy convergence threshold of  $10^{-6}$  eV and a planewave cutoff of 258.689 eV. The Brillouin zone was integrated with a 2x2x1 Monkhorst-Pack grid, and the Gaussian scheme with a broadening of 0.003 Ry was used in calculation. The band structure was calculated along F- $\Gamma$ -B-G- $\Gamma$  high symmetry directions, and the corresponding density of states is depicted in Fig. 11(a-c). The conduction band is dominated by Se-4p, In-5s atoms in a 1:3 layer, but in the 1:1 case, mostly by Se-4p atoms, and the valence band is occupied by Se-4p, S-3p in a 1:3 layer, whereas in a 1:1 layer, it is occupied by Se-4p near the Fermi level. This indicates that stronger interaction or bonding is more likely in Se, In, and S atoms. The band gap ( $E_g$ ) for a 1:3 proportion heterostructure layer of  $\text{Ag}_2\text{S}$  and  $\text{In}_2\text{Se}_3$  is 0.78987 eV, whereas for a single layer of Ag-S-Se-In is 1.37339 eV. The contour plot of electron charge density shows that more charge is accumulated around the Se atom in both these cases, which shows that the Se atom plays a crucial role in the materials' properties [Fig. 12]

A similar simulation was implemented for the optical parameter calculations, such as real and imaginary parts of dielectric constant, real and imaginary parts of conductivity, reflectivity, refractive index, attenuation coefficient, extinction coefficient, and absorption, respectively, within 0-20 eV photon energy.



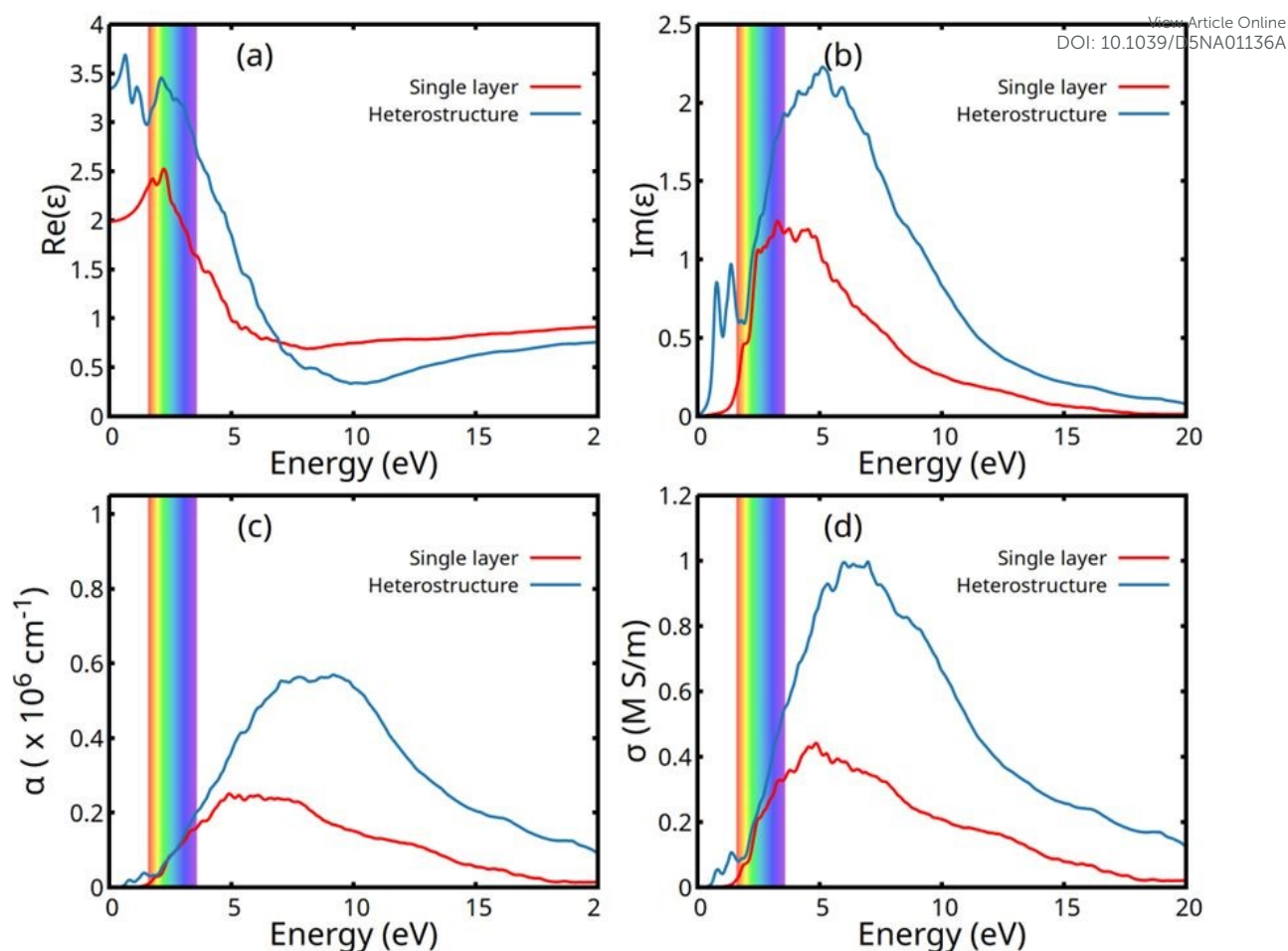


**Fig. 11.** Calculated band structures and corresponding density of state of (a,b) 1:3 proportion layer of  $\text{Ag}_2\text{S}$  and  $\text{In}_2\text{Se}_3$  heterostructure (c,d) single layer of Ag-S-Se-In.



**Fig. 12.** Contour plot for the charge density of (a) 1:3 proportion layer of  $\text{Ag}_2\text{S}$  and  $\text{In}_2\text{Se}_3$  heterostructure and (b) single layer of Ag-S-Se-In.



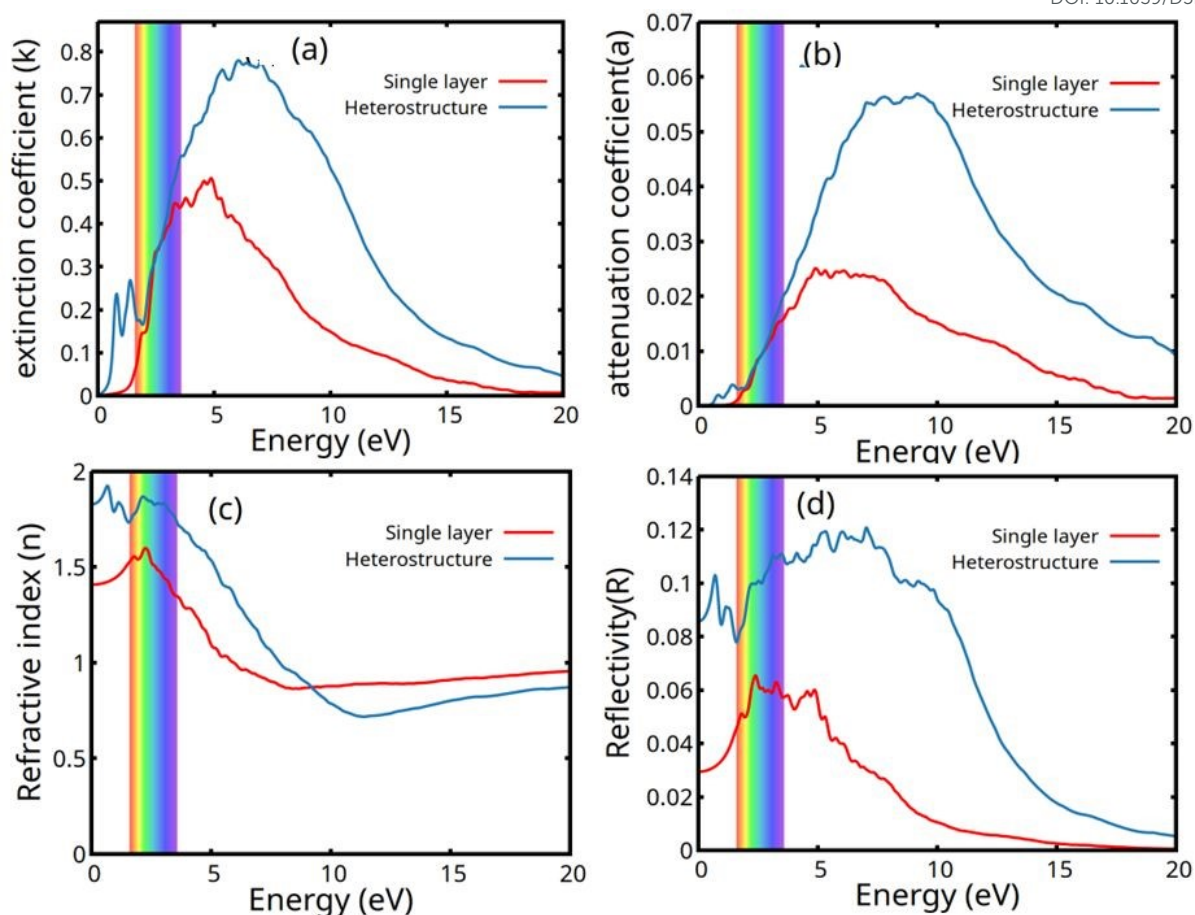


**Fig. 13.**  $\text{Ag}_2\text{S}$  and  $\text{In}_2\text{Se}_3$  heterostructure 1:3 and 1:1 proportion (a) real part, (b) imaginary part of dielectric function, (c) absorption coefficient, (d) optical conductivity.

The optical parameters of a material are crucial to investigate an excellent candidate for photovoltaic and optoelectronic applications. When the matter interacts with the electromagnetic radiation, it gets polarized by the applied electric field. It affects its interaction with electromagnetic waves. Fig. 13(a, b) illustrates the real and imaginary parts of the dielectric parameter. The imaginary component, indicating the material's absorption of electromagnetic radiation, denotes the energy loss per unit volume. This aspect is directly linked to electronic transitions between unoccupied and occupied states within the sample. The real part presents peaks  $\sim 1.13$  eV and  $1.78$  eV for hetero and single layers of Ag-S-Se-In, respectively. Then it decreased with increased photon energy. Within the zero-frequency range, the static dielectric constant for the heterolayer is  $3.34$ , whereas for a single layer, it is  $1.98$ . Fig. 13(c) depicts the maximum absorption coefficient values at photon energy of  $9.1$  eV for hetero and  $4.8$  eV for a single layer of Ag-S-Se-In.







**Fig. 14.** Ag<sub>2</sub>S and In<sub>2</sub>Se<sub>3</sub> heterostructure 1:3 and 1:1 proportion (a) extinction coefficient, (b) attenuation coefficient (c) refractive index, (d) reflectivity.

Optical conductivity, extinction coefficient, attenuation coefficient, and refractive index are displayed in Fig. 14(a-d) for both the hetero layer and the single layer of Ag-S-Se-In. Fig 14(d) indicates around 8% of reflectance for heterolayer and 3% of reflectance for single layer at 0 eV.

It should be noted that the computational model represents an atomically thin Ag-S-In-Se heterostructure, designed to approximate the local interfacial environment between Ag<sub>2</sub>S and In<sub>2</sub>Se<sub>3</sub> layers. This simplified configuration allows accurate insight into charge transfer, electronic coupling, and band alignment at the junction, which govern the experimentally observed optoelectronic behavior, even though it does not replicate the full experimental film thickness.



## 5. Conclusion

The diffusion of Ag<sub>2</sub>S into In<sub>2</sub>Se<sub>3</sub> heterostructure-based films was studied for a Visible light photodetector. The annealing resulted an increased photo detectivity of  $7.32 \times 10^9$  Jones for 200 °C annealed film. The intermixing of the layers resulted in the transformation of amorphous to crystalline phase and enhanced the photo current during both the rise (62.35 nA) and decay (67.74 nA) phases. The strong  $I_{on}/I_{off}$  ratio of 3.96 (rise) and 3.18 (decay) was observed, resulting in its rise and fall times as  $\tau_r = 7.15$ s, and  $\tau_d = 6.35$ s. These are moderate and well-balanced, suggesting efficient charge separation and recombination kinetics suitable for visible light photodetection. The crystallinity increase is reflected by the increase in the crystallinity size from 15.820 nm to 18.156 nm upon annealing. The developed AgInSe<sub>2</sub> and AgInS<sub>2</sub> phases resulted from the intermixing, which reduced the disorder, as well as being well visualised from SAED and HRTEM images. The increase in porosity in the films, as noticed from FESEM morphology, resulted in a decrease in hydrophobicity. The contact angle decreased from 110 ° to 100 ° upon 250 °C annealing. The increased transmittance is accompanied by reduced optical density with enhanced skin depth. The optical gap increased from 0.816 to 1.311 eV upon 250 °C annealing, which is well supported by the DFT study. The steepness parameter increased from  $4.761 \times 10^{-2}$  to  $5.983 \times 10^{-2}$ , resulting from the decreased Urbach energy from 543 meV to 432 meV upon annealing. The decreased refractive index and extinction coefficient is well supported by the DFT result. The overall optical and photo response behaviours of the films suggest their wider application as Visible light photodetectors and are useful for various optoelectronic devices.

## Supporting Information

FESEM images at 100 nm scale for all films, EDS spectra, individual bandgap plot and IV data for dark and light current of all the films.

## Conflicts of interest

There are no conflicts to declare.

## Data Availability Statement

All the data related to the present work are included in the manuscript and supporting information and will be available upon reasonable request.

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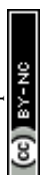
Computing (HPC) facility provided by the Center of Excellence in High Energy and Condensed Matter Physics, Department of Physics, Utkal University, India.

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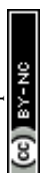
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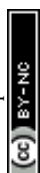
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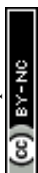


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## **Data Availability Statement**

View Article Online  
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**Title of Manuscript:** *Diffusion-Induced Enhanced Photoresponsivity, Detectivity in  $\text{Ag}_2\text{S}$  and  $\text{In}_2\text{Se}_3$  Heterostructure for UV-Visible Photodetector: An Experimental and Computational Analysis*

**Corresponding Authors:** Dr. Ramakanta Naik

- All data needed to support the conclusions in the paper are presented in the manuscript and the Electronic Supplementary Material. Additional data related to this paper may be requested from the corresponding author upon request.
- The data are plotted in with the help of Origin software.
- Analysis and quantification of images were performed using ImageJ software.

