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# A Novel Electrochemical Exfoliation Route to Tailor Graphene Bandgap through Silicon Incorporation: Semi-metallic to Semiconducting Transition

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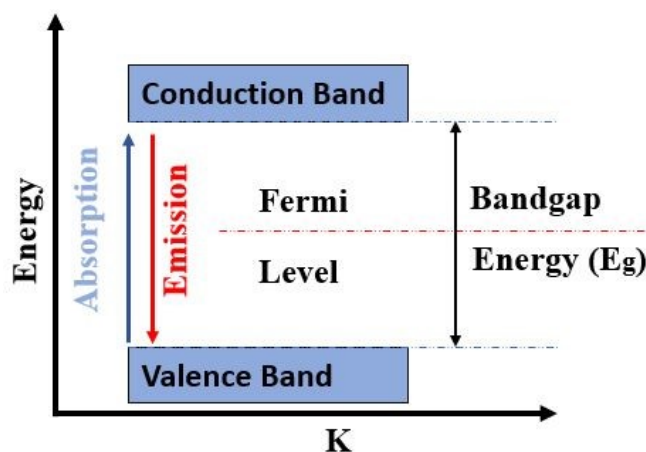
## ABSTRACT

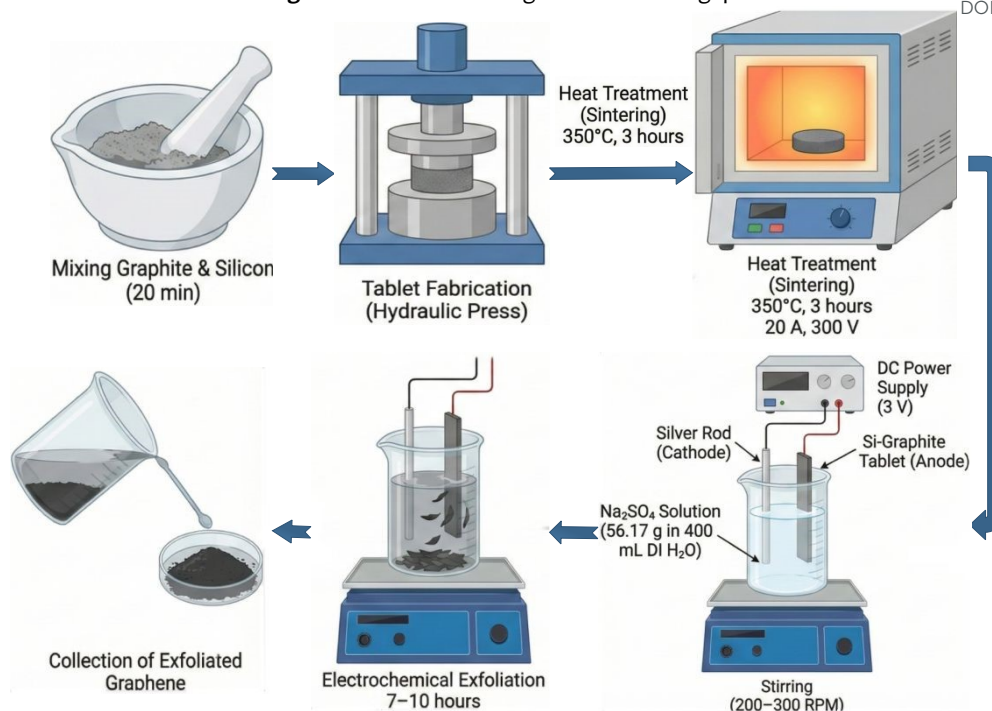
Graphene-silicon (Gr-Si) composites were produced by powder compaction, low-temperature sintering, and subsequent electrochemical exfoliation, which was further improved to provide materials with variable band gaps for semiconductor applications. By systematically varying the graphene-to-silicon precursor ratios (85:15, 80:20 and 75:25), the optical bandgap was successfully modulated from 1.25 eV to 1.56 eV, accompanied by a conductivity variation between 1.033 and 1.223 S m<sup>-1</sup>. Structural and chemical characterization using XRD, Raman spectroscopy, FTIR, UV-Vis spectroscopy, FESEM/EDX and thermal analysis revealed that silicon incorporation induces Si-O-C interfacial bonding and partial sp<sup>2</sup>→sp<sup>3</sup> rehybridization, leading to disruption of the π-electron network and progressive bandgap opening. Graphene-rich composites exhibited higher electrical conductivity, whereas silicon-rich compositions demonstrated enhanced thermal stability due to the formation of silica-like passivation layers. This scalable synthesis route establishes a direct structure-bandgap-conductivity correlation and highlights the potential of Gr-Si composites for optoelectronic, photovoltaic and next-generation semiconductor devices.

**Keywords:** Graphene, Silicon, Semiconductor, Bandgap, Composite, Electrochemical

## 1 INTRODUCTION

Graphene is a two-dimensional planar carbon nanostructure made up of sp<sup>2</sup>-hybridized carbon molecules that are one atom thick and form a hexagonal shape. Functional exogenous oxygen-bearing groups like hydroxyl, carbonyl, carboxylic and epoxy groups make up the edges of the nanoparticles. These groups increase the distance between the layers and make the atomic layer hydrophilic [1-5]. However, the zero-band gap in pure graphene limits its application in semiconductor devices, as shown in Figure 1. To make it easier to use graphene in electronics, scientists have tried different ways to create and change band gaps, such as chemical improvement, implanting heteroatoms and mixing it with other materials [6]. Nanostructured products have garnered considerable attention in scientific and engineering domains in recent years [7-9]. Graphene has amazing qualities. One layer of graphene has a Young's modulus of 1.0TPa [11], an intrinsic mobility of 2×10<sup>5</sup>cm<sup>2</sup>V<sup>-1</sup>s<sup>-1</sup> [10], and a thermal conductivity of (4.84±0.44) ×10<sup>3</sup> to (5.30±0.48) ×10<sup>3</sup>Wm<sup>-1</sup>K<sup>-1</sup> [12]. Graphene quantum dots exhibit unique optical, electrical, photoelectric and spin characteristics. At present, both top-down and bottom-up methods are used to make graphene quantum dots. Both of these methods take a lot of time, money and energy. To make GQDs, a single sheet of graphene is cut or broken up into particles with radii of less than 20 nm using the more effective downward method [13-16]. Graphene-based composites have been very interesting because graphene makes devices work better when added to them. Graphene composite materials facilitate numerous applications in photonics and optics [17], [18]. Hybrid graphene-silicon (Gr-Si) systems have become very popular among these methods. This is because silicon is a natural semiconductor and graphene makes the resulting composite stronger and better able to handle heat.



**Figure 1:** A schematic diagram of the bandgap.View Article Online  
DOI: 10.1039/D5NA01111F**Figure 2:** Schematic illustration of tunable bandgap exfoliated graphene-silicon (Gr-Si) fabrication process.

These graphene/silicon interfaces have now shown their usefulness in Schottky junction solar cells, broadband photodetectors, and tunable optoelectronic devices because they absorb more light, separate charges better and change how electronics work. Numerous studies [19-22] have demonstrated the synergistic interaction between the electronic characteristics of graphene and the semiconducting properties of silicon. Conversely, the material-level regulation of graphene-silicon hybrids remains underdeveloped. The majority of current research offers device-level demonstrations or methodologies for chemical doping. They don't work together to figure out how graphene/silicon composition, interfacial bonding Si-O-C domains, defect density, microstructural evolution, and bandgap modulation all work together. There are still not many scalable and reproducible ways to make Gr-Si composites with adjustable properties. The precise mechanisms that dictate the tunability of the bandgap in these hybrids are still not fully understood. In this paper, we show a controlled and scalable electrochemical exfoliation method for making graphene-silicon (Gr-Si) composites with adjustable band gaps. We mixed graphite and silicon powders in the right amounts (85:15, 80:20 and 75:25), pressed them together, heated them until they melted and then peeled them off to make very thin graphene sheets with silicon domains inside them. Figure 2 shows a diagram. For the first time, a direct link between composition, interfacial chemistry, defect structure and the change in the optical bandgap from 1.25 to 1.56 eV has been found by carefully changing the graphene-to-silicon ratio and fine-tuning the exfoliation parameters. To find the structure-property relationship that controls how the composites behave electronically and thermally, full characterizations are done using XRD, Raman spectroscopy, FTIR, UV-Vis spectroscopy, FESEM/EDX, TGA/DSC and four-point probe measurements. The initial comprehensive structure-bandgap-conductivity correlation for electrochemically exfoliated Gr-Si composites offers novel insights into the influence of Si incorporation on electronic states, microstructure, and transport properties.

So, compositional engineering gives Gr-Si hybrids new chances to be used in next-generation semiconductors, optoelectronics and photovoltaics.

Despite extensive reports on graphene-silicon heterostructures at the device level, most prior studies rely on chemical doping, CVD growth, or complex transfer processes, which limit scalability and reproducibility. Furthermore, the majority of existing works do not systematically correlate graphene-silicon composition with interfacial bonding, defect evolution, and bandgap modulation at the material level. In particular, the mechanisms governing bandgap tunability arising from Si-O-C interfacial domains and  $\pi$ -network disruption remain insufficiently understood.

In this study, we address these limitations by introducing a scalable electrochemical exfoliation strategy combined with controlled powder metallurgy processing to engineer graphene-silicon composites with tunable bandgaps. By varying silicon content within a percolation-preserving compositional window, we establish, for the first time, a comprehensive structure bandgap conductivity relationship in electrochemically exfoliated Gr-Si composites.



## 2. MATERIALS AND CHARACTERIZATION

Graphene was created by electro-exfoliating graphite with distilled water, sodium sulfate, graphite, silicon and a silver rod. Graphite and silicon were combined in a variety of molar ratios, including 85:15, 80:20 and 75:25. To create a homogenous mixture, a known amount of both powders was carefully mixed for roughly 20 minutes using a dry, clean mortar and pestle. A hydraulic jack was pressed from the bottom plate to compress the powder under pressure into a shape or size after 1 g of silicon-graphite combination was charged into a die to create the tablets. A high-temperature furnace was used to treat the resulting sintered tablet. For three hours, it was maintained at a steady temperature of 350°C while being exposed to an electrical input of 20 A and 300 V. A tougher and more resilient tablet resulted from the enhanced inter-particle bonding and grain boundary densification made possible by this heat exposure. All synthesis steps were repeated at least three times to ensure reproducibility and consistent trends were observed across batches. The electrochemical exfoliation voltage was 3 volts, electrolyte concentration was Na<sub>2</sub>SO<sub>4</sub>, 56.17 gm in 400 ml, stirring rate 200 to 300 rpm and the exfoliation duration was kept constant to minimize experimental variability. Uncertainty in the optical bandgap determination was estimated from linear-fitting errors in the Tauc plots, while electrical measurements were averaged across multiple probe placements. The graphite-silicon tablet served as the anode and a silver rod as the cathode. A DC power source supplied voltage with a steady current. With occasional stirring and 30-minute interval inspections, the exfoliation process lasted for seven to ten hours. Figure 3 illustrates that the exfoliated material eventually settled to the beaker's bottom and was gathered for additional processing.

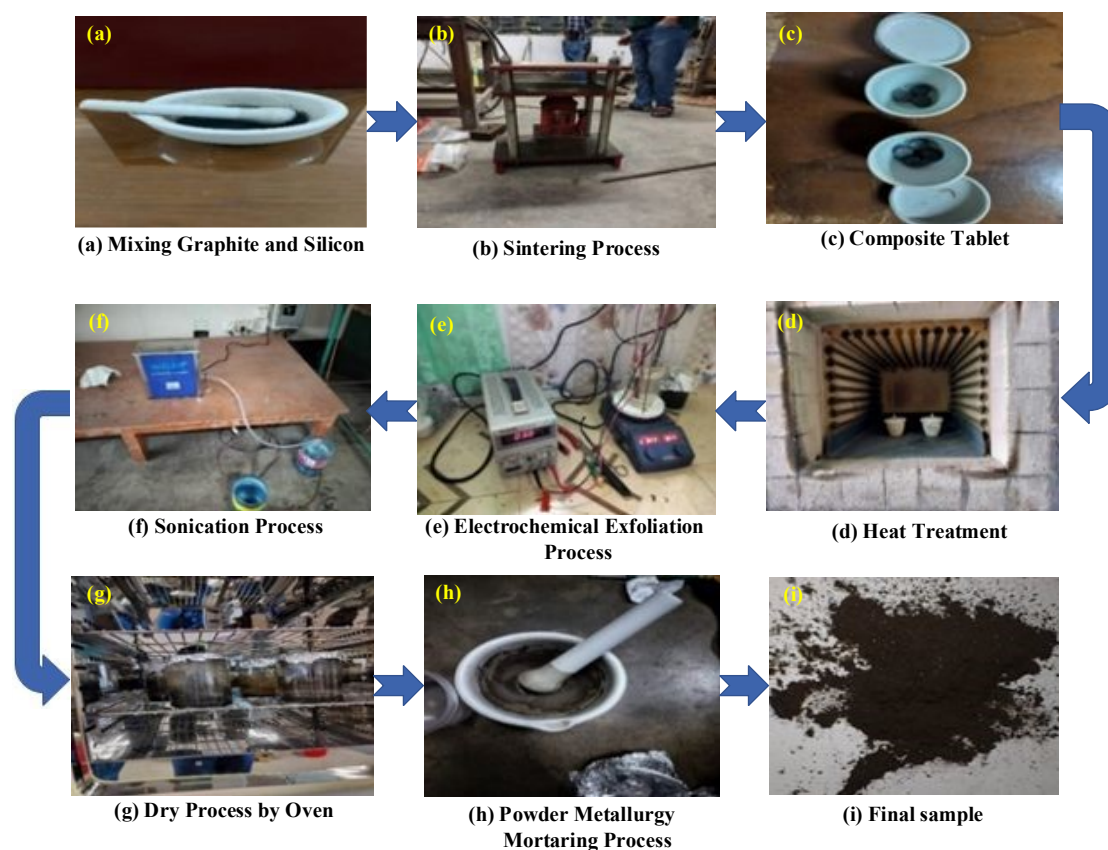


Figure 3. Process flow chart of exfoliated graphene-silicon composites.

## 3. RESULTS AND DISCUSSION

### 3.1 XRD Analysis

The crystal structure and phase purity of pure silicon and Si 15-25%-Graphene 75-85% mixtures were examined through XRD analysis. The scattered maxima of the specimens were identified by comparing them to the established reference Silicon Cubic (Diamond-type, Fd3m) JCPDS 00-026-1481 and Graphene hexagonal JCPDS 00-041-1487. The XRD patterns were collected in the 2θ range of 15°-65°, validating the successful synthesis of the Si-graphene composite and the crystalline characteristics of the prepared samples.

Table 1. Summary of XRD analysis from XRD peaks.

Peak Position (2θ)	Phase	Miller Indices	Intensity (a.u.)	Comments
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19.03°	Graphene	100	High (Gr 85% + Si 15%)	Presence of graphitic plane; multilayer graphene.
26.62°	Graphene	002	High (Gr 85% + Si 15%)	Characteristic graphene stacking reflection
28.07°	Silicon	111	Medium (Gr 85% + Si 15%)	Crystalline Si diffraction peak
29.03°	Silicon	200	Medium (Gr 85% + Si 15%)	Silicon cubic phase reflection
32.14°	Silicon	220	Medium (Gr 85% + Si 15%)	Stable Si orientation plane
29.12°	Silicon	111	High (Gr 80% + Si 20%)	Higher Si loading increases peak intensity
32.32°	Silicon	200	Medium (Gr 80% + Si 20%)	Silicon cubic phase reflection
33.92°	Silicon	311	Medium (Gr 80% + Si 20%)	Indicative of Si crystallinity
29.10°	Silicon	111	High (Gr 75% + Si 25%)	Dominant Si signal at high Si content
48.93°	Silicon	311	Lower (Gr 75% + Si 25%)	Reduced intensity due to graphene suppression

Table 1 shows the main diffraction peaks seen in Gr-Si composites, along with their Miller indices and relative intensities. For the Gr 85% + Si 15% sample, strong and sharp graphene reflections show up at 19.03° (100) and 26.62° (002). The broadened graphene (002) diffraction peak at 26.62° with reduced intensity compared to bulk graphite confirms partial exfoliation and reduced stacking order rather than pristine graphite. The absence of a diffraction peak near 10-12°, typically associated with graphene oxide, further excludes the formation of GO. Increased FWHM values and reduced crystallite sizes indicate few-layer graphene formation with induced microstrain resulting from electrochemical exfoliation and silicon incorporation. This shows that there are multilayer graphitic domains. Silicon reflections that match the 111, 200 and 220 planes show up at 28.07°, 29.03° and 32.14°, respectively. This means that the Si is well-crystallized and is embedded in the graphene matrix. As the Si content rises to 20-25 wt.%, the Si peaks get stronger, while the graphene peaks get weaker. This orderly change in peak intensity shows that the amount of Si is going up and the amount of long-range graphitic order is going down as the amount of silicon increases.

After 15% Si doping, graphene shows strong patterns at the diffraction peaks of 26.62° and 28.07°. The diffraction pattern of silicon has clear peaks that match the 111, 200, 220 and 311 planes. This shows that it is crystalline and has a pure phase. The 100 and 002 planes are visible in graphene. The XRD figure shows that the silicon peaks become stronger, while the characteristic graphene peaks get weaker as the amount of silicon increases. Previous studies have shown that the characteristic graphene diffraction peak is usually weaker in Si-graphene composites because silicon scatters light so much [23] [24]. We compared the peak positions and full width at half maximum (FWHM) of different phases and planes to standard values. This showed that the structure was well-ordered and crystalline. Emphasizing the effect of different amounts of silicon and graphene in the composites. 85 wt.% graphene and 15 wt.% silicon had the highest crystallinity, 85.39%, the largest crystallite size, 131.392 nm and the lowest dislocation density,  $5.792 \times 10^3 \text{ nm}^{-2}$  in the carbon phase. It also had the smallest distortion degree, 0.0688, which shows that it had better structural order, mechanical integrity and stability. Crystallographic parameters of the tested samples are listed in Table 2.

The d-spacing values are calculated using Bragg's law [25]:

$$i) d = \frac{\lambda}{2 \sin \theta} \quad (1)$$

Were,

$$\lambda = 1.5406 \text{ \AA} \text{ (Cu K radiation)} \text{ and } \theta = \frac{2\theta}{2}$$

The following formula is used to determine the degree of crystallinity (%):

$$ii) X_c = \frac{A_c}{A_c + A_a} \quad (2)$$

Were,

$A_c$  = Area of Crystalline

The following Scherrer's formula [26], [27] is used to compute the dimension of crystallite (nm):

$$iii) d = \frac{k\lambda}{2 \cos \theta} \quad (3)$$

Were,

$K = 0.9$  (constant of Scherrer)

$\lambda = 0.154006 \text{ nm}$  (X-ray source wavelength)

$\beta$  = Full width at half maximum in radians



$\theta$ =Top location in radians

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$$\text{iv) Displacement Density} = \frac{1}{D^2} \quad (4)$$

Were,  
D=Crystallite size (nm)

$$\text{v) Micro strain } \epsilon = \beta/4\tan\theta \quad (5)$$

Were,  
 $\beta$ =Full width at half maximum in radians  
 $\theta$ = Top location in radians

Table 2: Crystallographic parameters of the tested samples

Sample Name	% of Crystallinity	Average Crystalline Size (nm)	Dislocation Density ( $\text{nm}^{-2} \times 1000$ )	Average Macrostrain ( $\epsilon$ )
Gr 85%+Si 15%	85.39	132.392	5.792	0.0688
Gr 80%+Si 20%	66.11	115.736	7.465	0.0783
Gr 75%+Si 25%	73.84	126.978	6.202	0.0824

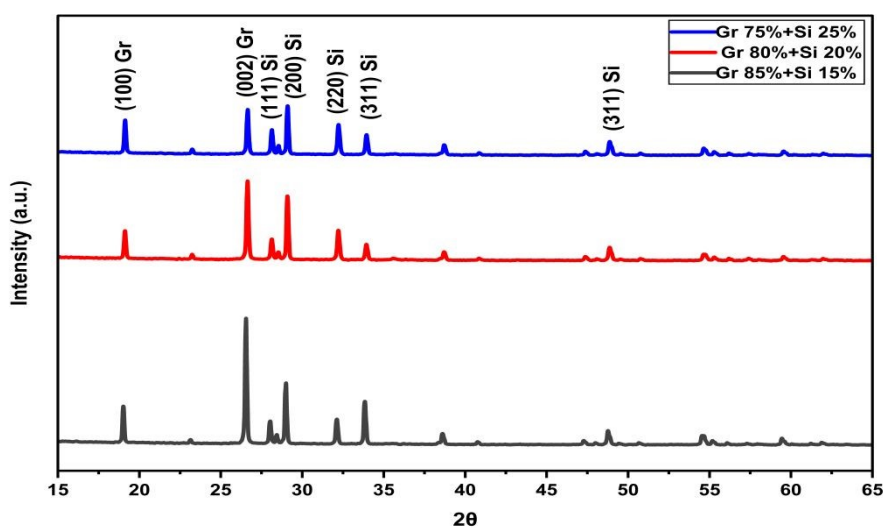


Figure 4: XRD pattern observed in Si 15-25wt.% and graphene 75-85 wt.% composites.

### 3.2 Raman Analysis

Figure 5 shows the intensity-normalised micro-Raman spectra of the exfoliated Gr-Si composites with varying compositions. All spectra exhibit a prominent and sharp peak at  $\sim 520 \text{ cm}^{-1}$ , which is assigned to the first-order optical phonon mode of crystalline silicon, as documented in previous studies [28]. The narrow linewidth and high intensity of this peak indicate that the silicon phase retains good crystallinity after composite formation, in excellent agreement with X-ray diffraction results and previous literature reports. Notably, the relative intensity of this Si peak increases systematically with increasing silicon content (from Si 15% to Si 25%), reflecting the compositional dependence of the Raman response. In addition to the silicon feature, all samples display the characteristic carbon-related bands, namely the D band ( $\sim 1350 \text{ cm}^{-1}$ ), G band ( $\sim 1580 \text{ cm}^{-1}$ ) and 2D band ( $\sim 2700 \text{ cm}^{-1}$ ), confirming the successful incorporation of graphene within the composite matrix. These peaks are caused by the second-order two-phonon scattering of graphene layers [29]. The G band arises from first-order  $E_{2g}$  phonon scattering, while the D band is defect-activated first-order scattering, and the 2D band is second-order two-phonon scattering. The 2D band is detected in all spectra; however, it appears broader and less intense than the G band. The estimated intensity ratio ( $I_G/I_{2D} \approx 0.59$ ) remains below unity for all samples, which is characteristic of few-layer graphene rather than monolayer graphene combined with the broadened 2D lineshape, is consistent with few-layer graphene rather than monolayer graphene. This observation is consistent with the expected outcome of the exfoliation process and indicates the formation of multilayer graphene with a certain degree of structural disorder. Furthermore, a gradual decrease in the intensity of the D, G, and 2D bands is observed with decreasing graphene content (from 85% to 75%), confirming the compositional tunability of the carbon phase within the composites. Subtle variations in the  $I_G/I_D$  ratio among the samples suggest minor differences in defect density and degree of disorder, likely arising from variations in graphene dispersion and its interaction with the silicon matrix. Overall, the Raman analysis confirms the coexistence of crystalline silicon and few-layer graphene, as well as their successful



integration, with composition-dependent structural characteristics that are expected to influence the functional performance of the Gr-Si composites. Article Online  
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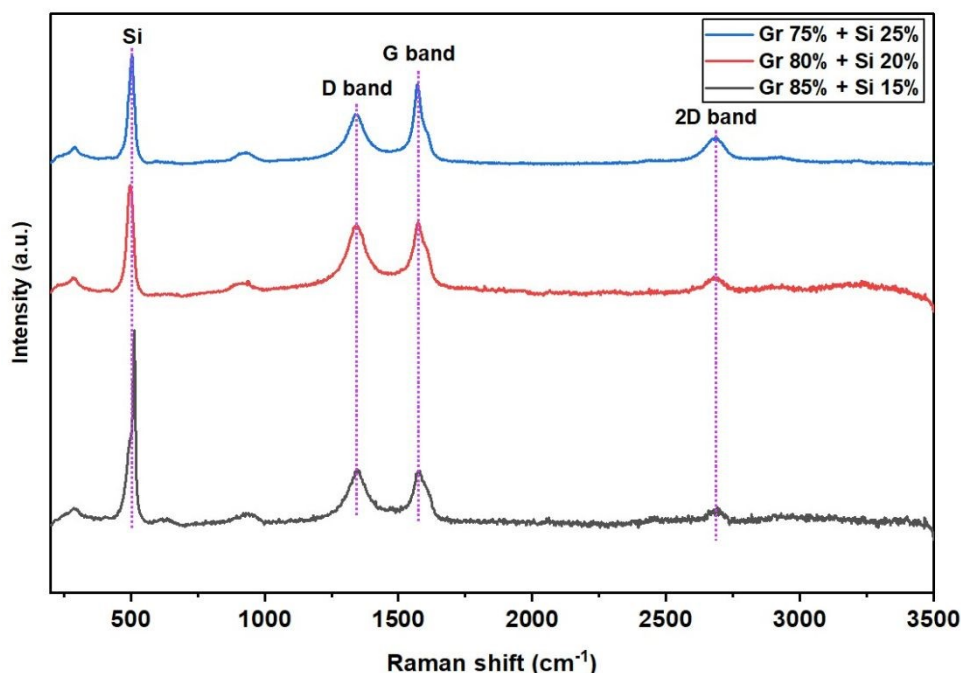


Figure 5: Micro Raman spectroscopy of Si-Graphene composites.

### 3.3 UV-Vis Analysis

Figure 6 shows the absorbance spectra of Gr-Si composites with different compositions that were measured between 300 and 1100 nm. All of the samples have a sharp absorption edge between 350 and 400 nm, which is typical of  $\pi$ - $\pi^*$  electronic transitions in graphene's  $sp^2$ -bonded carbon domains. The overall absorbance intensity increases steadily as the amount of graphene goes up. This is because the  $\pi$ -electron network gets bigger, which makes it better at collecting light. In Si-rich compositions, like Gr 75: Si 25, the absorption edge shifts slightly to the red. This is because of band tailing effects caused by defects at the Si-O-C interface and mid-gap states that are close to those defects [30]. On the other hand, the Gr 85%: Si 15% composite has a clearer absorption onset and a higher baseline absorbance across the visible-NIR range. This means that the carrier delocalization is better and the recombination loss is lower.

Using Tauc plots of  $(\alpha hv)^2$  versus photon energy ( $hv$ ), we figured out the optical bandgap ( $E_g$ ) of the graphene-silicon (Gr-Si) composites from the UV-Vis absorption data. We assumed that the electronic transition was allowed directly. We figured out the absorption coefficient  $\alpha$  from the measured absorbance and we extended the linear part of each  $(\alpha hv)^2$  curve to meet the energy axis at  $(\alpha hv)^2$ . This technique is frequently utilized for graphene-based semiconductors and thin-film hybrid materials [21],[23]. We used Tauc's relation [31] to figure out the samples' optical bandgap energy ( $E_g$ ) like this:

$$vi) (\alpha hv)^n = A(hv - E_g) \quad (6)$$

Were,

$$\alpha = \text{absorption coefficient} = 2.303 \frac{A}{t} \quad (A = \text{Absorbance, } t = \text{thickness}) [32]$$

$t$  = Thickness [33]

$hv$  = photon energy

$E_g$  = optical bandgap energy

$A$  = Constant

$n$  = Depending on the type of electronic transition

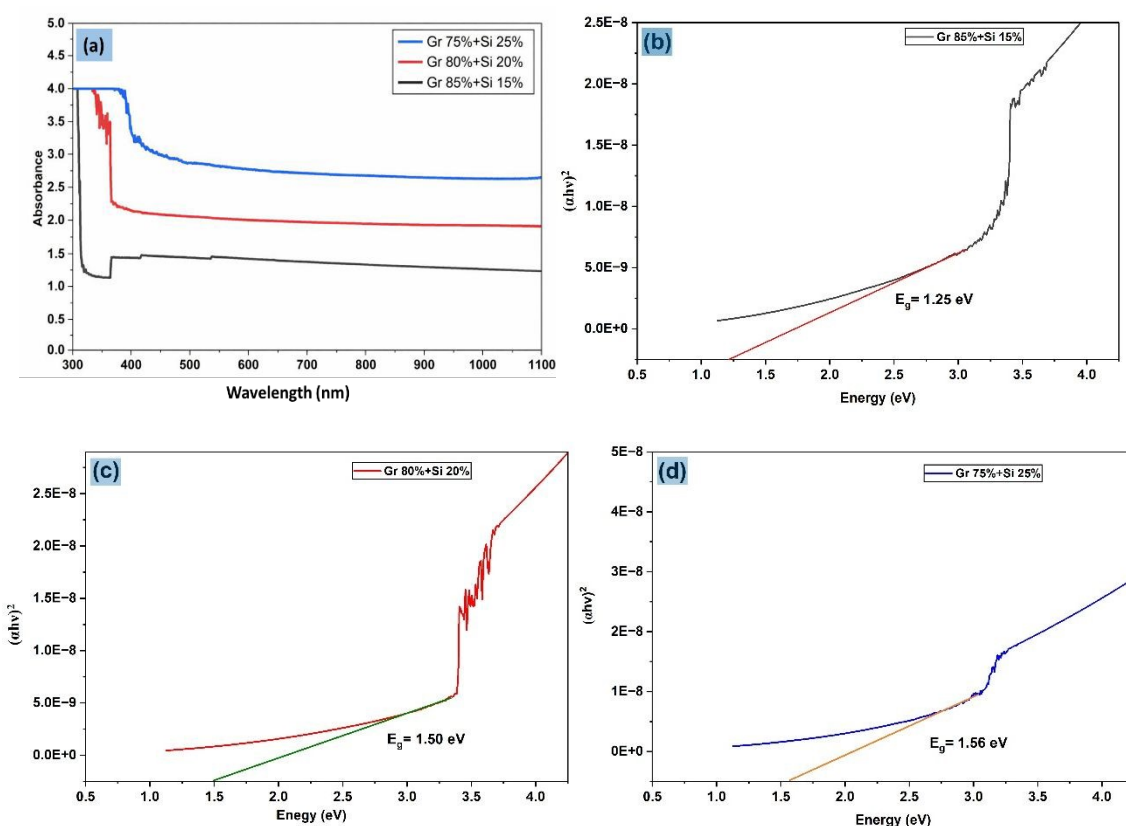
Figures 6 (a) to (d) show the Tauc plots for the Gr 85% + Si 15%, Gr 80% + Si 20% and Gr 75% + Si 25% composites. The optical bandgaps were extracted by linear extrapolation of  $(\alpha hv)^2$  versus photon energy, yielding values of 1.25 eV, 1.50 eV and 1.56 eV for increasing silicon content. This method is consistent with established graphene-silicon hybrid semiconductor



studies. So, as the silicon content goes from 15 wt.% to 25 wt.%, the optical bandgap goes from 1.25 eV to 1.56 eV. This shows that adding Si is a good way to change the electronic structure of Gr-Si hybrid systems.

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The gradual rise in  $E_g$  with Si loading is due to changes in the graphene  $\pi$ -electron network caused by silicon atoms and Si-O-C interfacial bonding. Adding Si creates localized electronic states and partial  $sp^2 \rightarrow sp^3$  rehybridization, which breaks up  $\pi$ -band delocalization and creates a finite bandgap. Similar mechanisms for widening the bandgap have been found in silicon-doped graphene, where Si 3p orbitals interact with graphene  $\pi$ -states to create energy levels that depend on the composition [34, 35]. These studies demonstrate that Si doping or Si-C hybridization can alter optical absorption edges, generate mid-gap states and enhance the effective bandgap via structural distortion and symmetry breaking.



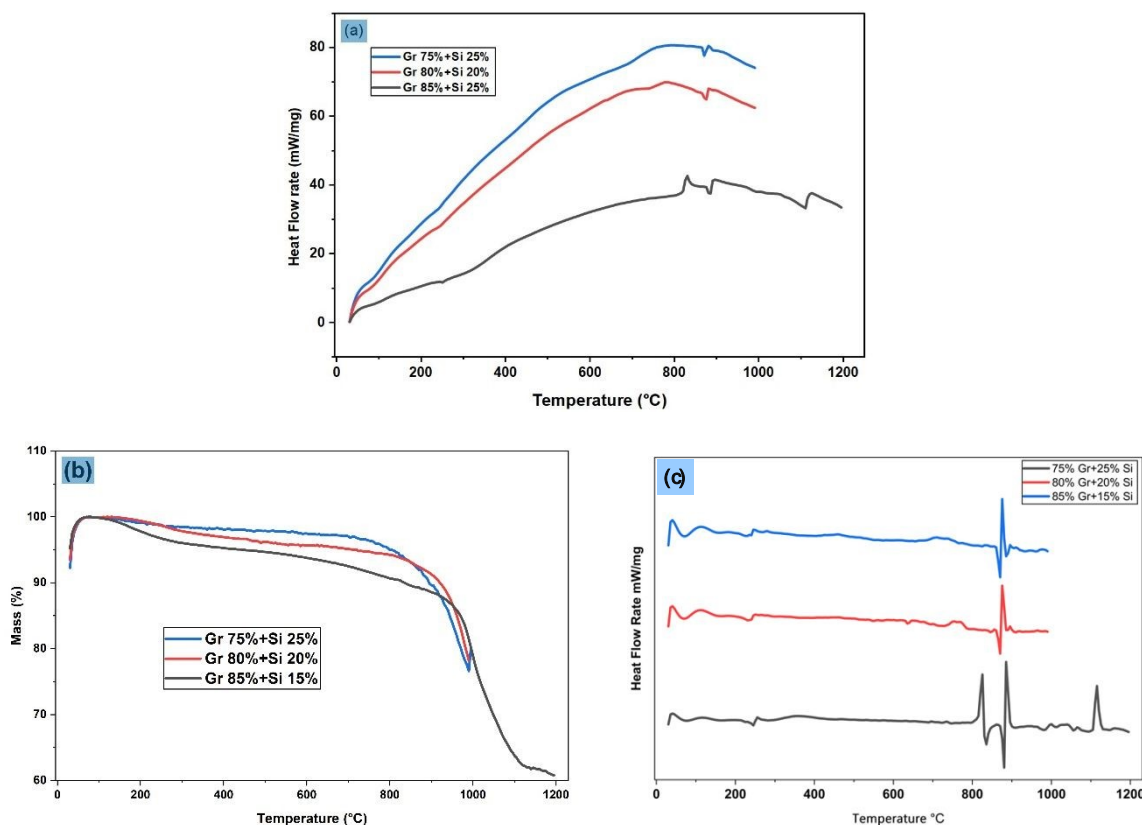
**Figure 6:**(a) UV-Vis absorbance spectra of graphene–silicon composites with different compositions in the 300–1100 nm range. The absorbance increases as the amount of graphene increases. (b)–(d) Tauc plots, which show  $(\alpha h\nu)^2$  as a function of photon energy, show that the bandgap changes between 1.25 and 1.56 eV depending on the composition. This is because of Si-O-C interfacial interactions and graphene  $\pi$ -state coupling.

### 3.4. Thermal Analysis (TGA-DSC)

The thermogravimetric profiles in Figure 7(b) show three areas where all the samples lost mass. Thermal stability was evaluated based on degradation onset temperature and residual mass behavior rather than total mass loss. Although the Si-rich composite exhibits higher total mass loss due to silicon oxidation, the delayed onset of major degradation and formation of protective Si-O-C networks indicate enhanced thermal resistance at elevated temperatures. However, the Gr-Si composites show a small increase in mass (1–3%) below 200 °C, especially in the Si-rich samples from this study, which they said was due to Si particles oxidizing and taking up oxygen at early stages [36]. This anomaly is due to the early oxidation of silicon particles and Si-O-C domains that are on the surface when there is oxygen in the air. A gradual decline means that the remaining organic groups are oxidizing, and the oxygen-containing groups C-O, C=O, and Si-O-C are breaking down. It is shown that adding Si improves thermal passivation by creating silica-like shells at the interface. Above 700 °C, the carbon lattice breaks down and Si turns into  $SiO_2$  in the air, which causes a sharp weight loss of 10–15 %. This shows that adding more silicon makes the material more thermally stable. Figure 7 (a) shows the differential scanning calorimetry (DSC) curves of graphene-silicon (Gr-Si) composites taken in air from room temperature to 1200 °C. The wide endothermic region below 300 °C is where moisture is released, and surface species break down. The successive exothermic peaks at higher temperatures (500–850 °C) are where defective carbon is oxidized and Si-O-C and  $SiO_2$  networks are formed. Figure 7 (c) shows derivative DSC curves (heat-flow rate versus temperature). There are several



exothermic maxima between 500 °C and 850 °C, which correspond to the oxidation of carbon domains in order and the restructuring of the Si-O-C network. The Si-rich composite Gr 75: Si 25) has the sharpest and most temperature-shifted exothermic peak, which shows that the protective SiO<sub>2</sub> passivation layer that forms at the interface makes the material more thermally stable and slows down the start of oxidation [37].



**Figure 7:** a) DSC, b) TGA and c) Derivative DSC thermograms of graphene-silicon composites showing initial mass gain due to silicon oxidation and enhanced thermal stability with higher Si content.

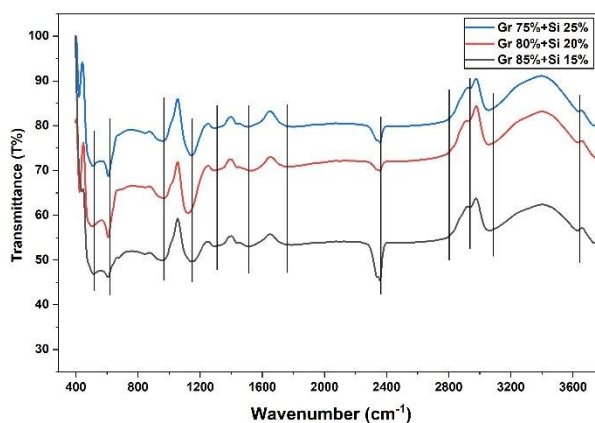
### 3.5. FTIR Analysis

Figure 8 shows the FTIR spectra of Gr-Si composites with different compositions (Gr 85% + Si 15%, Gr 80% + Si 20% and Gr 75% + Si 25%). It shows that the surface functionalization changes significantly as the silicon content goes up. Normalized FTIR peak intensities reveal a systematic increase in C=O and C-O vibrational modes with increasing silicon content, indicating enhanced oxidation and interfacial bonding. This trend supports bandgap widening via localized electronic states and partial sp<sup>2</sup>→sp<sup>3</sup> rehybridization. The peak at about 3066 cm<sup>-1</sup>, which corresponds to aromatic C-H stretching, is the same in all three samples. This shows that the sp<sup>2</sup> carbon domains in the graphene are still there [38]. The C=O stretching peak at about 1755 cm<sup>-1</sup> gets stronger as the amount of silicon increases. This means that more oxidation is happening and more carboxyl and ester groups are forming. As more silicon is added, the C-O stretching peak at 1180 cm<sup>-1</sup> gets stronger. This means that there are more oxygenated functional groups [39]. The peak around 2407 cm<sup>-1</sup>, which is due to O=C=O stretching, is stronger in the Gr 80% + Si 20% composite. This means that samples with moderate silicon content can hold more CO<sub>2</sub> [40]. The FTIR spectrum shows a peak at 667 cm<sup>-1</sup> that is usually linked to C-Br stretching or O-H bending in alcohols or phenols. It shows that there is either a carbon-bromine bond or a hydroxyl group, which gives us information about the material's chemical makeup [41]. The Si-O bond, which is usually seen around 512 cm<sup>-1</sup>, is not very clear in the FTIR spectra. So, these samples don't show any clear signs of Si-O bonding. Also, as the amount of silicon increases, the amount of oxidation, oxygen functionalities and Si-O bonding also increases. This makes the bandgap wider and the conductivity lower. The Gr 75% + Si 25% composite (S<sub>3</sub>) has the most oxidation and the biggest bandgap [42, 43]. The functional groups and chemical bonding of the tested samples are noted in Table 3.



**Table 3:** Compare the functional groups and chemical bonding of the tested samples.

Sample Name	Peak Position Wavenumber [ $\text{cm}^{-1}$ ]	Functional Groups	Class	Peak Details
S1	3633	O-H (Free) (Alcohol/Phenol)	Alcohol, Phenol	Strong, sharp
	3062	C-H (Aromatic)	Aromatic Hydrocarbon	Medium to weak
	2356	O=C=O	Carbonyl	Strong absorption
	2938	N-H (Amine salt)	Amine Salt	Strong, broad
	1770	C=O	Ester, Aldehyde	Strong, sharp
	1527	N-O	Nitro compound	Medium
	1299	C-O	Alcohol, Ether	Strong, sharp
	968	C=C	Alkene	Medium
	581	C-I	Halo compound	Strong, sharp
	609	C-Br	Halo compound	Strong
S2	3409	N-H stretching	Primary Amine	Medium
	2969	C-H stretching	Alkane	Medium
	2337	O=C=O	Carbon dioxide	Strong, sharp
	1677	C-H bending	Aromatic compound	Medium
	1369	O-H bending	Alcohol	Medium
	1068	C-N stretching	Amine	Strong
	647	C≡C stretching	Alkene	Medium
	497	C-Cl stretching	Alkyl chloride	Strong, sharp
S3	3629	O-H (Free) (Alcohol/Phenol)	Alcohol, Phenol	Strong, broad
	3054	C-H (Aromatic)	Aromatic Hydrocarbon	Medium to weak
	2356	O=C=O	Carbonyl	Strong absorption
	2938	N-H (Amine salt)	Amine Salt	Strong, broad
	1758	C=O	Ester, Aldehyde	Strong, sharp
	1515	N-O	Nitro compound	Medium
	1153	C-O	Alcohol, Ether	Strong, sharp
	964	C=C	Alkene	Medium
	513	C-I	Halo compound	Strong
	609	C-Br	Halo compound	Strong

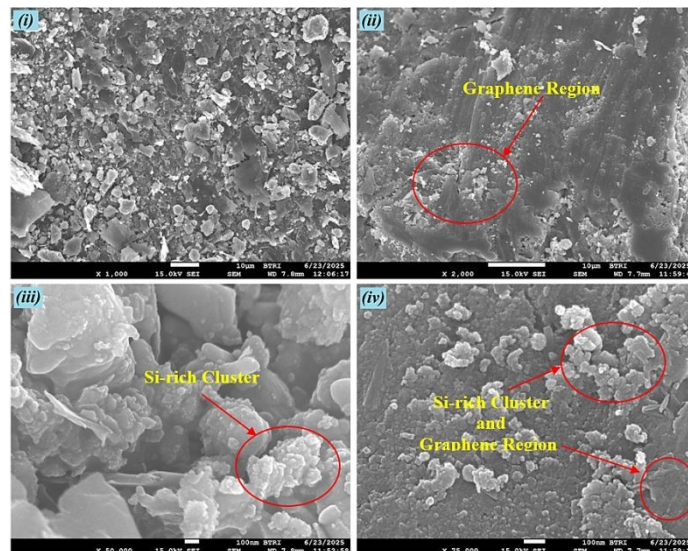
**Figure 8:** FTIR spectroscopy of silicon drop graphene samples.

### 3.6. FESEM Analysis

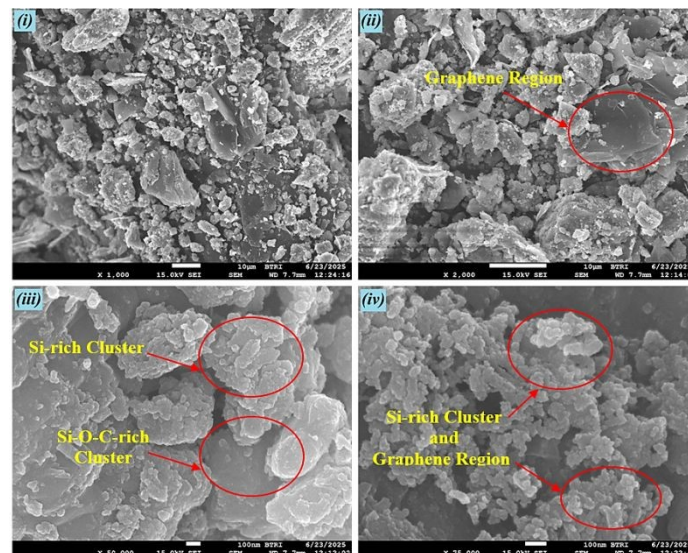
FESEM was employed to study the surface morphology and microstructural evolution of the Gr-Si composites. Figure 9 (a) depicts the microstructure of a 15% Si addition, characterized by smooth, wrinkled graphene sheets that create a



continuous conductive network, interspersed with a limited number of finely dispersed Si nanoparticles. Figure 9 (b) shows the enlarged graphene area where the lamellar carbon layers tightly overlap, which makes sure that there are few defect sites and a lot of electrical connectivity. This morphology will allow for metallic-like conduction with an optical band gap of about 1.25 eV. The homogeneous texture with low surface roughness also shows that the sheets stick together well and that Si does not fully agglomerate. This confirms that graphene keeps its natural two-dimensional continuity and high carrier mobility even when Si concentration is low [4]. This gets rough and grainy, with Si clusters evenly spread out throughout the graphene matrix. In the same way, the area marked in Figure 9 (b) had Si clusters inside the graphene flakes and the interfacial Si-O-C zones from partial oxidation and bonding between Si and C atoms were shown during thermal processing in Figure 9 (c). The morphology suggests that silicon atoms are partially integrated into the graphene framework, resulting in defect states at the interfaces and localized energy levels.

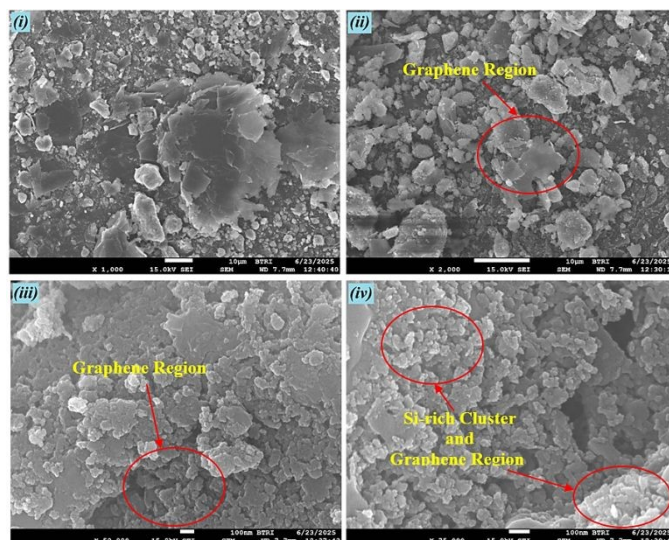


(a) FESEM images of graphene 85% + silicon 15% composites for sample one.



(b) FESEM images of graphene 80% + silicon 20% composites for sample two.

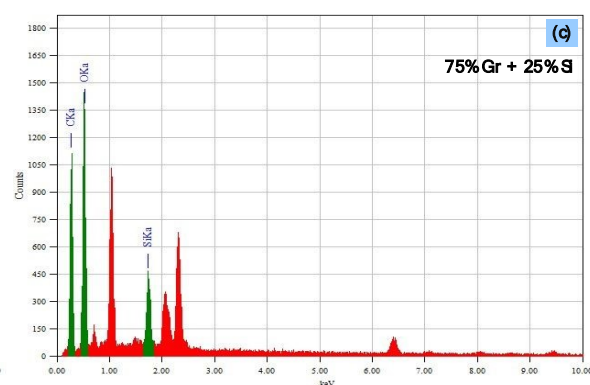
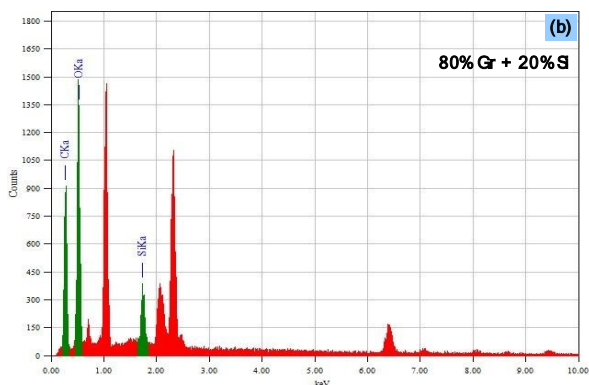
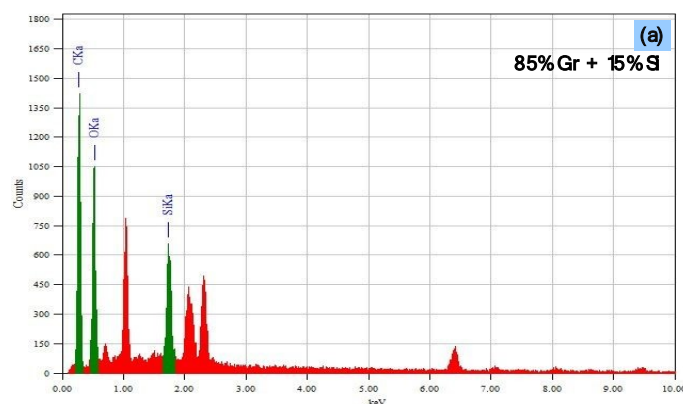




(c) FESEM images of graphene 75% + silicon 25% composites for sample three.

**Figure 9:** Field Emission Scanning Electron Microscopy images of graphene + silicon (a) Gr 85% + Si 15%, (b) Gr 80% + Si 20% and (c) Gr 75% + Si 25% composites for samples one, two and three.

The presence of conductive graphene and semiconducting Si-O-C domains elucidates the intermediate bandgap (1.50 eV) identified in UV-Vis/Tauc analysis. Similar hybrid interfaces have been documented to enhance charge separation and photocarrier longevity in graphene-Si composites [44], [45]. Figure 9(c) shows that 25% of the Si FESEM images are made up of big silicon granules and broken graphene sheets. This shows that Si is clumping together and graphene is less continuous. The Si-O-C-rich cluster zones show that oxidized silicon cores are surrounded by disordered carbon. This suggests that there is some phase separation and the formation of silica-like shells. The morphological transformation indicates a definitive shift from metallic graphene behavior to semiconductor-like performance, analogous to the findings in the Si-functionalized graphene systems reported by Xu et al. [46].



**Figure 10:** Energy-Dispersive X-ray spectra of graphene + silicon (a) Gr 85% + Si 15%, (b) Gr 80% + Si 20% and (c) Gr 75% + Si 25% composites.

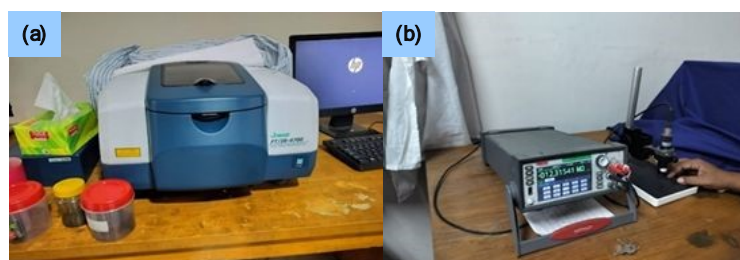


In Table 4, the elemental makeup of graphene-silicon composite samples with weight ratios of 85:15, 80:20, and 75:25, taken at magnifications of 1000 $\times$ , 2000 $\times$ , 5000 $\times$  and 7500 $\times$ . The table shows that C and O are the most common atomic elements in all of the samples, while Si is present in smaller amounts. The Gr 75% + Si 25% sample had the most carbon atoms (69.02 Atom%), which means it had a lot of graphene in it. The amount of silicon atoms went down as the amount of graphene went up.

**Table 4:** Elemental analysis of graphene-silicon (Gr-Si) composite samples at different weight ratios.

Sample	Element	(keV)	Mass (%)	Sigma	Atom (%)
Gr 85% + Si 15%	C K	0.277	52.31	0.30	60.60
	O K	0.525	42.13	0.51	36.64
	Si K*	1.739	5.56	0.12	2.76
	Total		100.00		100.00
Gr 80% + Si 20%	Element	(keV)	Mass%	Sigma	Atom%
	C K	0.277	50.53	0.27	58.50
	O K	0.525	45.44	0.50	39.50
	Si K*	1.739	4.03	0.10	2.00
Total		100.00		100.00	
Gr 75% + Si 25%	Element	(keV)	Mass%	Sigma	Atom%
	C K	0.277	61.98	0.27	69.02
	O K	0.525	35.80	0.48	29.93
	Si K*	1.739	2.22	0.07	1.06
Total		100.00		100.00	

Elemental identification based solely on FESEM contrast is insufficient. Therefore, EDX spectra were used to confirm compositional variations. It is noted that EDX is surface sensitive and semi-quantitative. Hence, measured silicon contents may differ from bulk precursor ratios due to exfoliation induced segregation, oxidation and spatial heterogeneity. Figures 10 (a) to (c) show the EDX spectra of Gr-Si composites with three different samples. The spectra show characteristic elemental peaks for C K $\alpha$  (0.28 keV), O K $\alpha$  (0.52 keV), and Si K $\alpha$  (1.74 keV). The spectra show that carbon, oxygen and silicon are all present in all of the samples, which proves that silicon was successfully hybridized into the graphene framework. The changes in relative intensity show changes in composition and shape that are in line with what FESEM shows. Figure 10 (a) shows that 85% Gr + Si 15% of the high carbon content (52 wt.%) and moderate oxygen level (42 wt.%) show that a continuous graphene matrix is dominant, with only a small amount of Si (5.6 wt.%). This composition matches the smooth, conductive shape seen in FESEM and explains why the optical bandgap is the lowest (1.25 eV). Figure 10 (b) shows an 80% Gr + 20% Si spectrum with 50.53 wt.% C, 45.44 wt.% O, and 4.03 wt.% Si. The Si percentage looks a little lower than in S1, but the higher oxygen content suggests that there is more interfacial oxidation and Si-O-C bonding. This agrees with the granular shape seen in FESEM for this sample. These oxidized domains create localized electronic states that make it harder for carriers to move around. This is consistent with the intermediate bandgap (1.50 eV) and the highest sheet resistance of the samples. The composition shows a transition zone where graphene sheets and Si-O-C regions are present at the same time, which makes the structure more disordered. Figure 10 (c) shows (Gr 75% + Si 25%), the relative Si peak intensity goes up and the O K $\alpha$  signal gets stronger. This suggests that Si oxidation and Si-O-C bond formation are both happening more quickly. The higher oxygen content (45 wt.%) shows that there is more interfacial hybridization, which is in line with the granular texture seen in FESEM. These Si-O-C linkages create localized electronic states that partially open the bandgap (1.50 eV) and keep the conductivity at a moderate level. The compositional trend (C > O > Si) and the changes in the Si and O peaks show that Si-O-C interfaces are forming over time, which control the pathways for charge transport. There have been reports of similar behavior in graphene-silicon nanohybrids, where controlled Si oxidation and interfacial bonding change how well the materials work for both light and electricity [47].



**Figure 11:** Four-point probe resistivity measurements of Graphene-Silicon (Gr-Si) composites; (a) Sheet Resistance ( $R_s$ ) and Conductivity ( $\sigma$ ) versus Silicon Content, (b) Graphical Correlation between Bandgap ( $E_g$ ) and Electrical Conductivity ( $\sigma$ ).



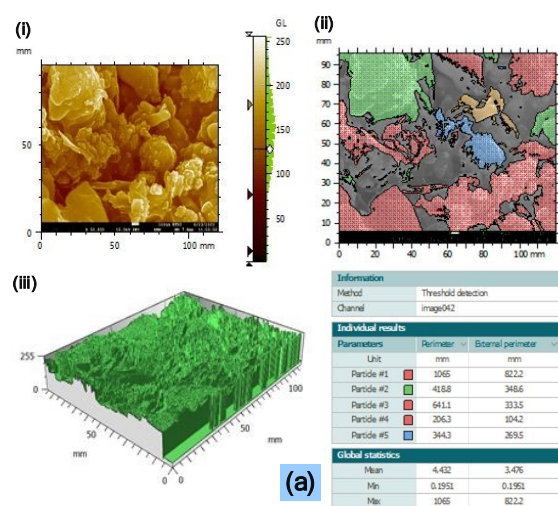
**Table 5:** Relationship between silicon content, bandgap modulation and electrical conductivity in graphene-silicon composites. Article Online  
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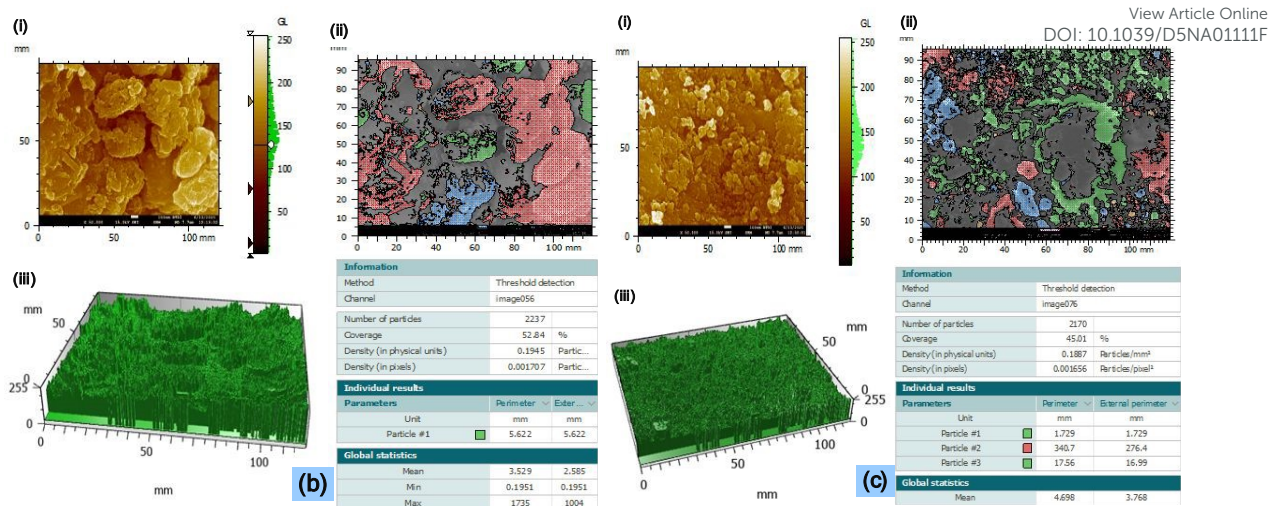
Nominal composition (Gr:Si)	Bandgap $E_g$ (eV)	Sheet resistance $R_s$ (k $\Omega$ )	Conductivity $\sigma$ (S m $^{-1}$ )
85:15	1.25	8.179	1.223
80:20	1.50	9.678	1.033
75:25	1.56	8.867	1.128

Table 5 shows a summary of electrical conductivity and how it relates to structure. Four-point probe measurements at room temperature indicate that conductivity diminishes as the optical bandgap expands from 1.25 eV (Gr 85% + Si 15%) to 1.56 eV (Gr 75% + Si 25%), which aligns with the gradual disruption of the  $\pi$ -conjugated graphene network due to Si incorporation. S1 (Gr 85%/Si 15%) has the lowest sheet resistance (8.179 k $\Omega$ ; 1.223 S m $^{-1}$ ) because FESEM shows that the graphene pathways are continuous and the Si is well-dispersed. S2 (Gr 80%/Si 20%) has the highest resistivity (9.678 k $\Omega$ ; 1.033 S m $^{-1}$ ) because the Si clusters are coarser, which stops percolation, and the Raman ID/IG is higher. S3 (Gr 75%/Si 25%) has the biggest bandgap, but it still regains some conductivity (8.867 k $\Omega$ ; 1.128 S m $^{-1}$ ) because graphene ribbons go around Si agglomerates and make new conductive networks. Similar to earlier research, adding Si to graphene broke up  $\pi$ -networks, raised the bandgap, and lowered conductivity [48]. These results, along with TGA/DSC evidence of silica-like passivation at high temperature, show that the quality of Si dispersion, not just the amount, controls the transport-bandgap trade-off in Gr-Si composites. This means that S1 can be used for contact/interconnect roles and S3 can be used for optoelectronic junctions where a wider gap with moderate conductivity is needed. Our findings for S1–S3 establish a correlation between the expanding band gap and diminishing conductivity in our Gr-Si composites [49].

### 3.7. SURFACE TOPOGRAPHY AND PARTICLE ANALYSIS

Figure 12 shows the analysis of the FESEM image application in Mountain 9, including surface topography, 3D views, and particles. The Gr85+Si15 sample had very small particles, about 0.22-0.30  $\mu\text{m}$ , and very large, with a perimeter of 900  $\mu\text{m}$  and a low particle density of about 0.15 particles/mm $^2$ . These big groups made sharp height changes in the 3D surface model, which made the topography uneven, with isolated high peaks being the most common. The surface became much more even as Si rose to 20 wt.%. The Gr80+Si20 sample had a particle density of about 0.19-0.21 particles/mm $^2$  and the range of particle sizes got smaller. The smallest features were still between 0.18 and 0.25  $\mu\text{m}$  and the largest particles had a perimeter of only about 600-650  $\mu\text{m}$ . The average size of the particles stayed around 3.0-3.3  $\mu\text{m}$ , which shows that the smaller particles were well-distributed. The corresponding 3D view showed a surface with a fine texture because there were a lot of moderate-height protrusions and no sudden changes in elevation, which is consistent with better dispersion. The particle density in the Gr75+Si25 sample was about 0.18-0.20 particles/mm $^2$ , which is similar to the 20% Si composition. However, the particle size distribution was wider again.





**Figure 12:** 3-D views, surface topography, and particle analysis of composite (a) Gr 85% + Si 15%, (b) Gr 80% + Si 20%, and (c) Gr75% + Si 25%.

The smallest sizes were between 0.20 and 0.28  $\mu\text{m}$ , and the biggest features had a perimeter of 720 to 780  $\mu\text{m}$ . The average size went up to about 4.4-4.8  $\mu\text{m}$ , which means that agglomeration happened again. The 3D topography showed wide, connected hills instead of separate peaks. The observation also showed that higher Si loading helps thicker and more connected surface domains grow. Overall, the best morphology was achieved at 20 wt.% Si, because height distribution, along with particle size and density, was able to create the most stable and uniform surface structure.

#### 4. CONCLUSIONS

This study introduces an innovative electrochemical exfoliation technique to manipulate the bandgap of graphene with regulated silicon integration, facilitating a distinct transition from semi-metallic to semiconducting properties. Graphene-silicon composites with composition dependent bandgaps ranging from 1.25 to 1.56 eV. Silicon acts as an interfacial modulator, inducing  $\pi$ -network disruption through Si-O-C bonding and localized electronic states. The quality of silicon dispersion, rather than absolute silicon content, governs the trade-off between bandgap opening and electrical conductivity. This is because graphene has a natural limit of zero bandgap. Through extensive structural, morphological, thermal, optical, and electrical analyses, we establish for the first time a complete structure-bandgap-conductivity correlation for electrochemically exfoliated Gr-Si composites, providing new insights into how compositional tuning influences electronic state modulation.

The cold sintering and controlled electrochemical exfoliation fabrication method used here is a scalable way to get around the current material-level limits of graphene-silicon hybrids by making it possible to engineer the bandgap in a precise and repeatable way. XRD and Raman studies show that mid-gap defect states appear and graphitic crystallinity is disrupted as silicon content rises. UV-Vis and Tauc analyses show that the bandgap widens depending on the composition, which is due to Si-O-C interfacial bonding and partial  $sp^2 \rightarrow sp^3$  rehybridization. The most important thing is that the quality of dispersion, not the amount of silicon, affects transport performance: Graphene-rich composites (Gr85:Si15) have high conductivity and low defect densities, which makes them good for high-frequency interconnects. Intermediate compositions (Gr80:Si20) have higher sheet resistance because Si-O-C clusters and silicon-rich materials (Gr75:Si25) partially regain conductivity through secondary graphene networks while also being more thermally stable.

Thermal analysis has shown that all of these composites are very stable at high temperatures. Si-rich samples, on the other hand, showed the formation of a protective  $\text{SiO}_2$  passivation shell above 700  $^\circ\text{C}$ , which shows that they can be used in high-temperature devices. FESEM and EDX characterization further illustrate composition-dependent morphological evolution, wherein well-dispersed silicon nanoparticles facilitate coherent electron pathways at low Si loading and provoke agglomeration at elevated Si content, thereby enhancing bandgap widening through localized electronic state formation.

In general and especially where tunable band gaps and controlled thermal stability are needed, these results show that Gr-Si composites are exciting materials for semiconductor, optoelectronic, and photovoltaic applications. So, the Gr75:Si25 composition is a good mix of semiconducting performance and moderate conductivity for optoelectronic junctions. Future research should concentrate on integrating density functional theory and molecular dynamics simulations to elucidate atomic-scale details regarding bandgap evolution and interfacial charge-transfer mechanisms.

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### Conflicts of Interest

The authors declare no conflicts of interest.

### REFERENCES

- [1] C. N. R. Rao, A. K. Sood, K. S. Subrahmanyam, and A. Govindaraj, "Graphene: the new two-dimensional nanomaterial.," *Angew. Chem. Int. Ed. Engl.*, vol. 48, no. 42, pp. 7752–7777, 2009, doi: 10.1002/anie.200901678.
- [2] J. Jinschek, E. Yucelen, H. Calderon, and B. Freitag, "Quantitative atomic 3-D imaging of single/double sheet graphene structure," *Carbon N. Y.*, vol. 49, pp. 556–562, Feb. 2011, doi: 10.1016/j.carbon.2010.09.058.
- [3] M. J. Allen, V. C. Tung, and R. B. Kaner, "Honeycomb carbon: A review of graphene," *Chem. Rev.*, vol. 110, no. 1, pp. 132–145, Jan. 2010, doi: 10.1021/cr900070d.
- [4] A. K. Geim and K. S. Novoselov, "The rise of graphene," *Nat. Mater.*, vol. 6, no. 3, pp. 183–191, 2007, doi: 10.1038/nmat1849.
- [5] A. K. Geim, "Graphene: status and prospects.," *Science*, vol. 324, no. 5934, pp. 1530–1534, Jun. 2009, doi: 10.1126/science.1158877.
- [6] Y. Zhang, Y.-W. Tan, H. L. Stormer, and P. Kim, "Experimental observation of the quantum Hall effect and Berry's phase in graphene," *Nature*, vol. 438, no. 7065, pp. 201–204, 2005, doi: 10.1038/nature04235.
- [7] E. Solati and D. Dorrani, "Investigation of the Structure and Properties of Nanoscale Grain-Size beta-Tantalum Thin Films," *Mol. Cryst. Liq. Cryst.*, vol. 571, pp. 67–76, Jan. 2013, doi: 10.1080/15421406.2012.741347.
- [8] Ş. Tjalu, M. Bramowicz, N. Naseri, S. Kulesza, and A. Ghaderi, "Surface Micromorphology and Fractal Geometry of Co/CP/X (X= Cu, Ti, SM and Ni) Nanoflake Electrocatalysts," *RSC Adv.*, vol. 6, Mar. 2016, doi: 10.1039/C6RA01791F.
- [9] Ş. Tjalu, M. Bramowicz, S. Kulesza, A. Shafiekhani, A. Ghaderi, and M. Ahmadi, "Gold Nanoparticles Embedded in Carbon Film: Micromorphology Analysis," *J. Ind. Eng. Chem.*, vol. 35, Dec. 2015, doi: 10.1016/j.jiec.2015.12.029.
- [10] K. Bolotin et al., "Ultrahigh Electron Mobility in Suspended Graphene," *Solid State Commun.*, vol. 146, pp. 351–355, Feb. 2008, doi: 10.1016/j.ssc.2008.02.024.
- [11] C. Lee, X. Wei, J. Kysar, and J. Hone, "Measurement of the Elastic Properties and Intrinsic Strength of Monolayer Graphene," *Science*, vol. 321, pp. 385–388, Jul. 2008, doi: 10.1126/science.1157996.
- [12] A. Balandin et al., "Superior Thermal Conductivity of Single-Layer Graphene," *Nano Lett.*, vol. 8, pp. 902–907, Apr. 2008, doi: 10.1021/nl0731872.
- [13] M. Bacon, S. Bradley, and T. Nann, "Graphene Quantum Dots," *Part. Part. Syst. Charact.*, vol. 31, Apr. 2014, doi: 10.1002/ppsc.201300252.
- [14] J. Shen, Y. Zhu, X. Yang, and C. Li, "ChemInform Abstract: Graphene Quantum Dots: Emergent Nanolights for Bioimaging, Sensors, Catalysis and Photovoltaic Devices.," *Chem. Commun. (Camb.)*, vol. 48, pp. 3686–3699, Mar. 2012, doi: 10.1039/c2cc00110a.
- [15] J. Zhang et al., "Strongly Green-Photoluminescent Graphene Quantum Dots for Bioimaging Applications," *Chem. Commun. (Camb.)*, vol. 47, pp. 6858–6860, Jun. 2011, doi: 10.1039/c1cc11122a.
- [16] C. Zhu, J. Zhai, and S. Dong, "Bifunctional fluorescent carbon nanodots: Green synthesis via soy milk and application as metal-free electrocatalysts for oxygen reduction," *Chem. Commun. (Camb.)*, vol. 48, pp. 9367–9369, Aug. 2012, doi: 10.1039/c2cc33844k.
- [17] Y. H. Lee, L. Polavarapu, N. Gao, P. Yuan, and Q.-H. Xu, "Enhanced Optical Properties of Graphene Oxide–Au Nanocrystal Composites," *Langmuir*, vol. 28, no. 1, pp. 321–326, Jan. 2012, doi: 10.1021/la204047a.
- [18] S. Bongu, P. Bisht, T. V. Thu, and A. Sandhu, "Multiple Nonlinear Optical Response of Gold Decorated-Reduced Graphene Oxide-Nanocomposite for Photonic Applications," *J. At. Mol. Condens. Nano Phys.*, vol. 2, pp. 207–214, Dec. 2015, doi: 10.26713/jamcnp.v2i3.346.
- [19] T. Feng et al., "Efficiency enhancement of graphene/silicon-pillar-array solar cells by HNO<sub>3</sub> and PEDOT-PSS," *Nanoscale*, vol. 4, no. 6, p. 2130, 2012, doi: 10.1039/c2nr12001a.
- [20] X. Miao et al., "High Efficiency Graphene Solar Cells by Chemical Doping," *Nano Lett.*, vol. 12, pp. 2745–2750, May 2012, doi: 10.1021/nl204414u.
- [21] X. An, F. Liu, Y. Jung, and S. Kar, "Tunable Graphene-Silicon Heterojunctions for Ultrasensitive Photodetection," *Nano Lett.*, vol. 13, Jan. 2013, doi: 10.1021/nl303682j.
- [22] K. S. Novoselov et al., "Electric Field Effect in Atomically Thin Carbon Films," *Science (80-.)*, vol. 306, no. 5696, pp. 666–669, 2004, doi: 10.1126/science.1102896.
- [23] C. Fu, C. Song, L. Liu, W. Zhao, and X. Xie, "High Reversible Silicon/Graphene Nanocomposite Anode for Lithium-Ion Batteries," *Int. J. Electrochem. Sci.*, vol. 11, no. 1, pp. 154–164, 2016, doi: https://doi.org/10.1016/S1452-3981(23)15833-0.
- [24] H. Huang, "The Effect of Commercialized Binders on Silicon Oxide Anode Material for HighCapacity Lithiumion Batteries," *Int. J. Electrochem. Sci.*, vol. 11, pp. 8697–8708, Oct. 2016, doi: 10.20964/2016.10.29.



- [25] L. Pevelen, D. Delphine, *Small Molecule X-ray Crystallography, Theory and Workflow* (Ed: J. C. Lindony), Academic Press 2010, pp. 2559-2576. Article Online DOI: 10.1039/D5NA01111F
- [26] S. Fatimah, R. Ragadhita, D. F. A. Husaeni, A. Bayu Dani Nandiyanto, *ASEAN J. Sci. Eng.* 2022, 2, 65.
- [27] D. J. Lim, N. A. Marks, M. R. Rowles, *Carbon* 2020, 162, 475.
- [28] K. Krishnamoorthy, M. Veerapandian, R. Mohan, and S.-J. Kim, "Investigation of Raman and photoluminescence studies of reduced graphene oxide sheets," *Appl. Phys. A*, vol. 106, pp. 501–506, Dec. 2011, doi: 10.1007/s00339-011-6720-6.
- [29] S. Dash, T. Dash, and T. Rout, "Preparation of graphene oxide by dry planetary ball milling technique under oxygen atmosphere," *IOP Conf. Ser. Mater. Sci. Eng.*, vol. 872, p. 12180, Jun. 2020, doi: 10.1088/1757-899X/872/1/012180.
- [30] M. Houmad, H. Zaari, A. Benyoussef, A. El Kenz, and H. Ez-Zahraouy, "Optical conductivity enhancement and band gap opening with silicon doped graphene," *Carbon N. Y.*, vol. 94, pp. 1021–1027, Nov. 2015, doi: 10.1016/j.carbon.2015.07.033.
- [31] P. R. Jubu, O. Adedokun, C. Mbakaan, A. Nathan-Abutu, E. Danladi, J. N. Tsaviv, P. I. Kyesmen, B. J. Akeredolu, A. T. Adepoju, F. Aungwa, Y. Yusof, O. S. Obaseki, Y. S. Elzawiei, M. Z. Pakhuruddin, *J. Mater. Sci.: Mater. Electron.* 2025, 36, 961.
- [32] T. Pandit, *Absorbance and Transmittance measurement of CsI thin films.* 2013.
- [33] R. Rai, B. K. Singh, *Absorbance and Transmittance Measurement of CsI Thin Films* 2013.
- [34] M. Shahrokhi and C. Leonard, "Tuning the band gap and optical spectra of silicon-doped graphene: Many-body effects and excitonic states," *J. Alloys Compd.*, vol. 693, pp. 1185–1196, 2017, doi: <https://doi.org/10.1016/j.jallcom.2016.10.101>.
- [35] S. Sahu and G. C. Rout, "Band gap opening in graphene: a short theoretical study," *Int. Nano Lett.*, vol. 7, no. 2, pp. 81–89, 2017, doi: 10.1007/s40089-017-0203-5.
- [36] G. Ge et al., "Manipulating Oxidation of Silicon with Fresh Surface Enabling Stable Battery Anode," *Nano Lett.*, vol. 21, no. 7, pp. 3127–3133, Apr. 2021, doi: 10.1021/acs.nanolett.1c00317.
- [37] C. Yin, X. Li, P. Chen, G. M. Kale, and B. Zhu, "Oxidation resistance and wettability of graphite/SiC composite," vol. 109, no. 7, pp. 629–637, 2018, doi: 10.3139/146.111644.
- [38] G. Surekha, K. V. Krishnaiah, N. Ravi, and R. Padma Suvarna, "FTIR, Raman and XRD analysis of graphene oxide films prepared by modified Hummers method," *Journal of Physics: Conference Series*, vol. 1495, p. 012012, Mar. 2020, doi: <https://doi.org/10.1088/1742-6596/1495/1/012012>.
- [39] Y. Liang, D. Wu, X. Feng, and K. Müllen, "Dispersion of graphene sheets in organic solvent supported by ionic interactions," *Advanced materials*, vol. 21, no. 17, p. 1679-1683, 2009, doi: 10.1002/adma.200803160.
- [40] S. Saxena, T. A. Tyson, S. Shukla, E. Negusse, H. Chen, and J. Bai, "Investigation of structural and electronic properties of graphene oxide," *Applied Physics Letters*, vol. 99, no. 1, p. 013104, Jul. 2011, doi: <https://doi.org/10.1063/1.3607305>.
- [41] N., Prabavathi & N. S. Nayaki, "Experimental Spectroscopic (FT-IR, FT-Raman, NMR) and DFT Studies of 7-methoxy-4-bromomethylcoumarin", *Journal of Environmental Nanotechnology*, vol. 3, no. 2, p. 108-121, 2014, doi: 10.13074/jent.2014.03.142072.
- [42] P. Shemella and S. K. Nayak, "Electronic structure and band-gap modulation of graphene via substrate surface chemistry," *Applied Physics Letters*, vol. 94, no. 3, Jan. 2009, doi: <https://doi.org/10.1063/1.3070238>.
- [43] C. V. Gomez, E. Robalino, D. Haro, T. Tene, P. Escudero, A. Haro, and J. Orbe, "Structural and Electronic Properties of Graphene Oxide for Different Degree of Oxidation," *Materials Today: Proceedings*, vol. 3, no. 3, pp. 796–802, Feb. 2016, doi: <https://doi.org/10.1016/j.matpr.2016.02.011>.
- [44] L. K. Pillari, K. Lessoway, and L. Bichler, "Carbon nanotube and graphene reinforced magnesium matrix composites: A state-of-the-art review," *J. Magnes. Alloy.*, vol. 11, pp. 1825–1905, Jun. 2023, doi: 10.1016/j.jma.2023.05.010.
- [45] J. Lee, C. Lee, J. Jun, D. Shin, and J. Jang, "A metal-oxide nanofiber-decorated three-dimensional graphene hybrid nanostructured flexible electrode for high-capacity electrochemical capacitors," *J. Mater. Chem. A*, vol. 2, Jul. 2014, doi: 10.1039/C4TA01695E.
- [46] M. Houmad, H. Zaari, A. Benyoussef, A. Kenz, and H. Ez-Zahraouy, "Optical Conductivity enhancement and band gap opening With Silicon doped graphene Optical Conductivity enhancement and band gap opening With Silicon doped graphene," *Carbon N. Y.*, vol. 94, Jul. 2015, doi: 10.1016/j.carbon.2015.07.033.
- [47] M. Houmad, H. Zaari, A. Benyoussef, A. El Kenz, and H. Ez-Zahraouy, "Optical conductivity enhancement and band gap opening with silicon doped graphene," *Carbon N. Y.*, vol. 94, pp. 1021–1027, 2015, doi: 10.1016/j.carbon.2015.07.033.
- [48] C. Chen et al., "An insitu iodine-doped graphene/silicon composite paper as a highly conductive and self-supporting electrode for lithium-ion batteries," *RSC Adv.*, vol. 7, no. 61, pp. 38639–38646, 2017, doi: 10.1039/C7RA06871A.
- [49] M. M. Ervasti, Z. Fan, A. Uppstu, A. V Krasheninnikov, and A. Harju, "Silicon and silicon-nitrogen impurities in graphene: Structure, energetics, and effects on electronic transport," *Phys. Rev. B*, vol. 92, no. 23, p. 235412, Dec. 2015, doi: 10.1103/PhysRevB.92.235412.



### Data availability

The data sets (Raw data) analyzed in this study are available from the corresponding author on reasonable request.

