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Multipolar Origin and Polarization-Controlled High-Q Quasi-BIC Fano Resonances in Dielectric Metasurfaces for Sensing Applications†

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We designed an multifunctional all-dielectric metasurface, employing cuboid structures patterned with bow-tie-shaped nanoholes, exhibiting multiple Fano resonances induced by quasi-bound states in the continuum (quasi-BICs) through structural asymmetry. Among them, several resonant modes demonstrated high quality factors in the range of 10^3–10^4, along with near-unity modulation depth and strong spectral contrast. The optical responses were analyzed utilizing the finite-difference time-domain (FDTD) method, with Fano profiles fitted to theoretical models and the BIC-governed modes validated via the squared inverse ratio law. Furthermore, multipolar decomposition and electromagnetic spatial field profile revealed the origins of the resonance, while LC circuit modeling provided additional physical insight into the Fano profiles. The proposed metasurface also exhibited strong polarization dependence, indicating its potential for active optical switching. Finally, refractive index sensing performance, including the potential of detecting *Vibrio cholerae* in proper environment, reached a sensitivity of 342 nm/RIU and a figure of merit of 217.14 RIU^{-1}. Advancing the control of high-Q quasi-BIC Fano resonances, this study highlights Fano resonators' potential for refractive index sensing and active switching.

1 Introduction

Metasurfaces with engineered strong light–matter interactions enable advanced and multifunctional control over optical wavefronts, including amplitude, phase, polarization, and dispersion [1–5]. Their compactness, high design flexibility, and ability to integrate with existing platforms make them ideal candidates for next-generation technologies in flat optics, compact imaging, on-chip photonic circuits, and broadband optical communication [6–8]. By harnessing sharp spectral features arising from resonant effects, such as Fano resonances and bound states in the continuum (BIC), these planar photonic structures provide exceptional capability for light manipulation at subwavelength scales [9–12]. This ability depends critically on the material's inherent optical characteristics and the precise engineering of resonant nanostructures. Metasurfaces are broadly categorized into plasmonic and all-dielectric types, each distinguished by their resonance mechanisms and electromagnetic response [13]. Plasmonic metasurfaces, typically composed of noble metals or graphene, rely on localized surface plasmon resonances to achieve strong near-field confinement in the visible to near-infrared range [14,15]. However, their practical utility is constrained by inherent ohmic losses, which limit Q factors and degrade performance in applications requiring high spectral selectivity [16]. Notably, all-dielectric metasurfaces constructed from high refractive index materials, such as silicon, germanium, gallium phosphide, support low-loss Mie resonances and enable the excitation of high-Q modes, including BICs, through careful control of geometric and symmetry parameters. These structures exhibit minimal absorption losses, enhanced field confinement, and are compatible with CMOS fabrication processes, making them particularly attractive for scalable and low-loss photonic integration [17–19].

Among the various multipolar excitations in nanophotonics, the toroidal dipole (TD), first introduced by Zeldovich in 1957, has emerged as a distinctive mode arising from closed-loop displacement currents, representing a unique excitation in all-dielectric metasurfaces with weak radiation and behavior markedly different from conventional electric (ED) and magnetic dipole (MD) responses [20]. Although TD modes were previously overlooked due to masking by dominant multipoles, recent studies have demonstrated that their excitation can be selectively enhanced through structural asymmetry, near-field coupling,

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and engineered high-index dielectric arrays[21,22]. These configurations support a high Q factor with suppressed radiative loss, making them promising for applications in sensing and nonlinear optics. However, the relatively weak near-field confinement of TD modes highlights further optimization in geometry and material design to maximize their photonic functionality.

Ultra-narrow optical resonances with exceptionally high Q factors and minimal losses play a vital role in modern nanophotonics, enabling enhanced light–matter interactions essential for applications such as high-sensitivity sensors, nonlinear optics, slow-light devices, narrowband absorbers, and spectral filters[4,23–27]. A common route to achieving such resonances is through Fano resonance, resulting from interference between radiative and non-radiative modes, producing sharp asymmetric line shapes[28,29]. This effect becomes particularly pronounced when combined with BICs, arising from non-radiating localized states embedded within the continuous radiation spectrum[30]. Although ideal BICs possess infinite Q factors and zero linewidth, practical implementation is achieved through quasi-BICs, where slight structural asymmetry induces weak coupling to free-space radiation, resulting in finite but extraordinarily high Q factors[31–34]. Such mechanisms have been widely utilized to excite high-Q TD Fano resonances in all-dielectric metasurfaces[35]. Pang *et al.* introduced square defects into Si tetramer arrays to realize polarization-independent quasi-BICs[36]. Sun *et al.* employed a nanorod–ring unit with broken symmetry, enabling tunable TD resonances by varying the nanorod length[35]. Li *et al.* achieved BIC-governed high-Q resonances in a Si-based hollow metasurface by incorporating asymmetric air holes[22]. Additionally, Lv *et al.* investigated quadrupolar quasi-BICs in GaP nanodisk arrays activated via lateral displacement[13]. However, beyond geometric asymmetry, this effect can also be realized through material asymmetry. Yu *et al.* demonstrated dual-band, polarization-insensitive TD modes using permittivity-asymmetric Si–InAs cuboid tetramers[37]. Although GaP exhibits a high refractive index comparable to Si, comprehensive investigations of GaP-based metasurfaces, especially regarding BIC-governed Fano resonances, remain insufficient. Moreover, hollow-structured GaP metasurfaces tailored for such resonance phenomena have yet to be thoroughly explored.

In this paper, we present an all-dielectric metasurface exhibiting multiple Fano resonances in the near-infrared (NIR) regime, predominantly governed by MD and TD excitations. The incorporation of structural asymmetry—achieved by modifying both the dimension and orientation of a single triangular nanohole of bow-tie-shaped nanoholes—introduced quasi-BIC-supported Fano resonances, resulting in additional pronounced Fano resonances. These resonances exhibited both exceptionally high quality factors and modulation depths. The Fano profiles for both symmetry and asymmetry configurations were fitted using the theoretical model, while the BIC mode was validated by calculating the squared inverse ratio law. Additionally, the Fano resonances were interpreted using an equivalent LC circuit model to elucidate the underlying resonance mechanism. To further analyze the origin of resonance, the contributions of various multipolar moments were investigated through spatial distributions of the electric and magnetic fields. The impact of the polarization angle on optical performance was investigated, demonstrating potential for being utilized as an optical switch. Moreover, the impact of the structural parameters was also investigated. Finally, the sensing capability of the structure was assessed by changing the ambient refractive index to simulate the presence of Vibrio cholerae. This all-dielectric metasurface, with lower complexity in design, offers promising prospects for advanced applications in RI sensing and active photonic switching.

2 Structure Design and Methodology

Figure 1(a) depicts our proposed all-dielectric metasurface, comprising a two-dimensional periodic array of GaP nanocuboids with square cross-section and bow-tie-shaped nanoholes, deposited on a SiO_2 substrate along the x and y axes with a periodicity, P = 650 nm. The thickness of the substrate was considered to be $t_{SiO_2} = 200$ nm. Moreover, Figs. 1(b) and (c) present the cross-sectional views along the xy- and yz-planes, respectively. The side of the square-shaped cross-section of GaP nanocuboids was w = 550 nm, while the thickness was $t_{GaP} = 150$ nm. The bow-tie-shaped nanoholes were comprised of two equilateral triangles with side lengths s_1, and s_2, separated from center to center with a distance of 275 nm. To facilitate the BIC, an asymmetry parameter δ was introduced, defined as $\delta = s_2 - s_1$. The complex refractive index of GaP and SiO_2 were adopted from Bond *et al.*[38] and Palik *et al.*[39], respectively [see ESI for details].
We employed the finite-difference time-domain (FDTD) approach in Ansys Lumerical to quantitatively investigate the optical properties of our proposed structure. The simulation configuration is illustrated in Fig. 1(d). Since our proposed metasurface exhibited periodicity along the x and y axes, we employed periodic boundary conditions to reduce computational space and simulation time while applying 12-layer perfectly matched layer (PML) boundary conditions along both directions of the z-axis to eliminate parasitic reflections from the structure. An additional mesh with dimension 4 nm × 4 nm × 4 nm was incorporated to simulate precisely. The simulation was conducted by considering the structure immersed in a liquid medium, such as water, with a refractive index of 1.33 at the temperature of 300 K. A CW-normalized plane wave source was incident from the top of the structure with a wavelength ranging from 900 nm to 1600 nm. We employed a frequency-domain field and power monitor beneath the structure to quantify transmittance $T(\lambda)$, which is the ratio of transmitted power, $P_T(\lambda)$, to the incident power, $P_I(\lambda)$, expressed as follows[4],

$$T(\lambda) = \frac{P_T(\lambda)}{P_I(\lambda)}. \quad (1)$$

3 Results and Discussion

3.1 Dimensional Asymmetry

The asymmetry in dimensions of the bow-tie-shaped nanoholes allowed our proposed structure to exhibit resonances governed by BIC. Figure 2(a) illustrates the transmittance spectra of the

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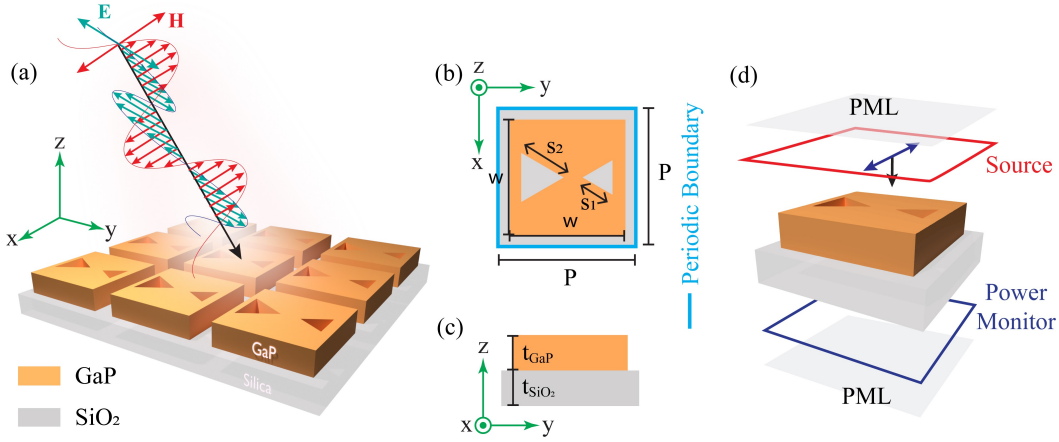

Fig. 1 (a) 3D schematic illustration of the proposed all-dielectric bow-tie etched metasurface, (b) x-y, and (c) x-z plane view of the structure. The optimized structural parameters: P = 650 nm, w = 550 nm, t_{SiO_2} = 200 nm, t_{GaP} = 150 nm, and $\delta = s_2 - s_1$. (d) Optical simulation configuration of the proposed structure.

structure for $\delta = 0$ nm ($s_1 = s_2 = 130$ nm), and $\delta = 100$ nm ($s_1 = 130$ nm, $s_2 = 230$ nm). When $\delta = 0$ nm, the structure maintained a distinct in-plane symmetry, resulting in two sharp Fano profiles at $\lambda = 1044$ nm (Mode I), and 1217 nm (Mode II). The structure also supported symmetry-protected BICs with theoretically infinite Q factors but remained undetectable in the transmittance spectrum. Moreover, we analyzed the finite size of the structure (20×20 array) for this configuration, resulting in broader but qualitatively consistent resonance behavior [See ESI for details]. However, the in-plane symmetry was disrupted when we considered $\delta = 100$ nm, allowing the non-radiative bound states to couple with the free space continuum and access the radiation channel[22]. Consequently, the radiation channel transformed into a resonant mode, allowing energy leakage and outward radiation, which led to the transformation of symmetry-protected BICs into quasi-BICs[13]. Therefore, we observed three additional Fano profile resonances at $\lambda = 911$ nm (Mode III), 1159 nm (Mode IV), and 1278 nm (Mode V) governed by quasi-BICs. Meanwhile, the Fano profile (Mode I and II) of this structure experienced a blue shift at $\lambda = 1036$ nm, and 1195 nm, respectively.

The performance of the metasurface in optical applications can be evaluated based on the modulation depth, spectral contrast ratio, and Q-factor of the Fano resonance. The modulation depth is defined by[37],

$$Modulation\ depth = \frac{T_{peak} - T_{dip}}{T_{peak}} \times 100\% \quad (2)$$

where T_{peak}, and T_{dip} are the maximum and minimum values of the transmittance spectrum, respectively. Moreover, we calculated the spectral contrast ratio by utilizing these two parameters,

$$Spectral\ contrast\ ratio = \frac{T_{peak} - T_{dip}}{T_{peak} + T_{dip}} \times 100\%. \quad (3)$$

To observe the transmittance spectra qualitatively, the Fano model can be expressed as,

$$T_{Fano} = \left| a_1 + ia_2 + \frac{b}{\omega - \omega_{s_\circ} + i\gamma} \right|^2, \quad (4)$$

where a_1, a_2, and b are the real number constants, $\omega_{s_\circ}$, and ω are the resonance and angular frequency, and γ is the overall damping loss. The radiative Q-factor, Q_{rad}, can be defined as $Q_{rad} = \omega_{s_\circ/2}\gamma$. We employed Eq. 4 to fit the calculated modes I and II for $\delta = 0$ nm and III to V for $\delta = 100$ nm. Figures 2 (b) and (c) present the calculated and fitted Fano profiles (Mode I and II) for $\delta = 0$ nm and (Mode III to V) for 100 nm, demonstrating that the fitted curve closely aligns with the calculated data. Table 1 summarizes the extracted Q_{rad} values from these fittings, modulation depth, and spectral contrast ratio for modes I, II ($\delta = 0$ nm) and I-V ($\delta = 100$ nm). A simplified LC circuit based on the stable-input impedance mechanism can be employed to replicate the Fano resonance profile. Introducing a capacitor or an inductor in parallel with the LC circuit enables the coupling between a broadband bright mode and a narrowband dark mode, respectively, thus giving rise to an asymmetric Fano-like resonance[40]. Figure 2(d) depicts the LC circuit representations for Modes I and II ($\delta = 0$ nm), where a complementary capacitor and inductor were incorporated for Modes I and II, respectively. The corresponding stable-input impedances of these configurations are as follows.

$$Z_{Mode\ I} = -\frac{j}{\omega C_c} \frac{(\omega - \omega_{s_\circ})(\omega + \omega_{s_\circ})}{(\omega - \omega_{s_\circ}\sqrt{(C_s/C_c)+1})(\omega + \omega_{s_\circ}\sqrt{(C_s/C_c)+1})} \quad (5)$$

$$Z_{Mode\ II} = \frac{j\omega L_s L_c}{(L_s + L_c)} \frac{(\omega - \omega_{s_\circ})(\omega + \omega_{s_\circ})}{(\omega - \omega_{s_\circ}/\sqrt{(L_c/L_s)+1})(\omega + \omega_{s_\circ}/\sqrt{(L_c/L_s)+1})} \quad (6)$$

We can define transmittance as the ratio of incident power (P_{input}) to transmitted power (P_{output}). Figure 2(e) depicts both the LC circuit fitted and calculated Modes I and II ($\delta = 0$ nm), demonstrating a good match between the fitted data and the calculated result [see ESI for details].

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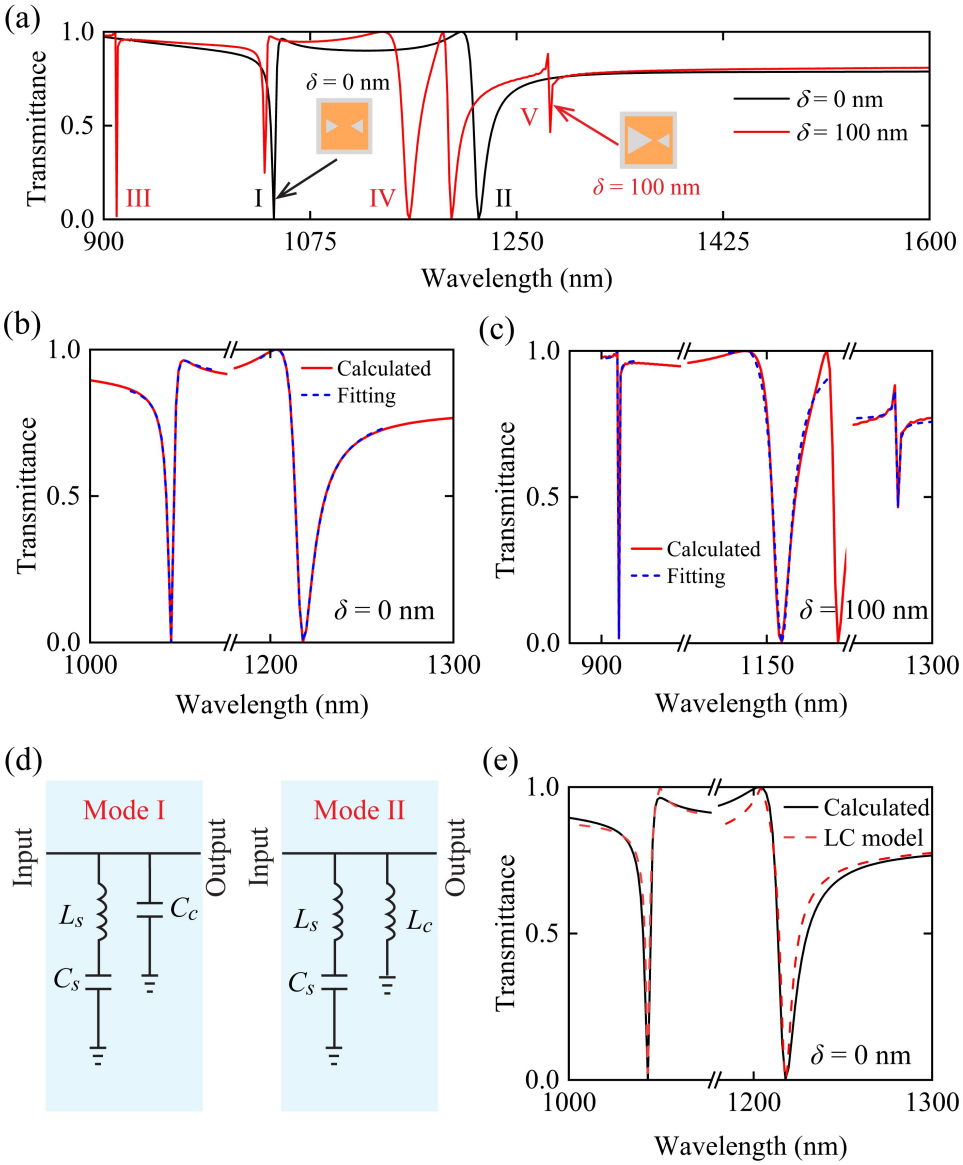

Fig. 2 (a) Transmittance spectra of our proposed structure for $\delta = 0$ nm, and 100 nm. Plot of FDTD calculated and fitted graph for (b) $\delta = 0$ nm near $\lambda = 1044$ nm, and 1217 nm (c) $\delta = 100$ nm near $\lambda = 911$ nm, 1159 nm, and 1278 nm under TM-polarized incident light. (d) Schematics of the LC circuit models for Mode I and II ($\delta = 0$ nm). (e) FDTD calculated and fitted curve with LC circuit for Mode I and II for $\delta = 0$ nm.

Table 1 Extracted values of quality factor, modulation depth, and spectral contrast ratio for modes I, II ($\delta = 0$ nm) and III–V ($\delta = 100$ nm).

δ	Mode	Quality Factor (Q_{rad})	Modulation Depth (%)	Spectral Contrast Ratio (%)
0 nm	I	2.90×10^2	99.93	99.87
	II	1.10×10^2	99.78	99.57
100 nm	I	6.14×10^2	74.51	59.39
	II	1.54×10^2	99.95	99.89
	III	3.17×10^3	98.50	97.04
	IV	8.96×10^1	99.50	99.00
	V	6.38×10^4	47.15	30.85

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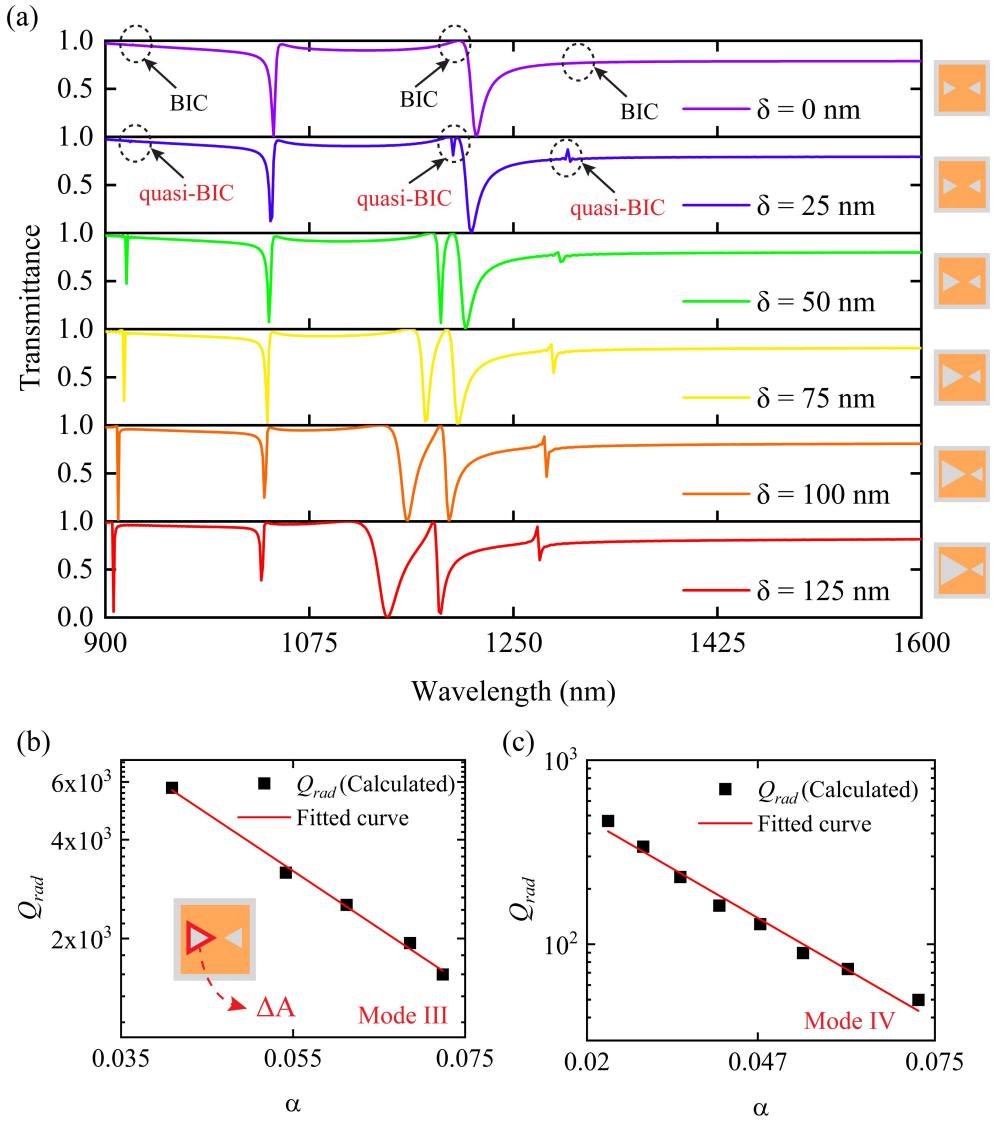

Fig. 3 Transmittance spectra for different asymmetry parameters, $\delta = 0$ nm to 125 nm under TM-polarized incident light (b) and (c) relationship between the Q-factor, Q_{rad}, and the degree of asymmetry α for Mode III and IV, respectively. α was defined by the ratio of ΔA and A. The solid red line represents the fitted data, indicating an inverse squared relationship between the Q_{rad} and α. Additionally, the inset of (b) highlights the magnitude of ΔA, denoted by the marked red region.

Figure 3(a) provides further insight into the formation of symmetry-breaking quasi-BIC-supported Fano resonance profiles by illustrating the transmittance spectra for $\delta = 0$ nm to 125 nm. At $\delta = 0$ nm, the transmittance spectrum exhibited three symmetry-protected BICs with a theoretically infinite Q-factor. When $\delta \neq 0$ nm, the in-plane symmetry was broken, leading to the appearance of three distinct Fano profiles (Mode III to V) supported by symmetry-breaking quasi-BIC. At $\delta = 25$ nm, we found these resonances at $\lambda = 921$ nm, 1198 nm, and 1299 nm, while at $\delta = 125$ nm, these resonances experienced a blue shift to $\lambda = 911$ nm, 1159 nm, and 1278 nm respectively. Moreover, Modes I and II exhibited a blue shift, similar to the discussion in the previous section. These blue shifts resulted from the decrease in the effective refractive index of the metasurface with the increase of the asymmetry degree α. Here, α represent normalized perturbation, defined as the ratio of ΔA to A, where ΔA and A represented the area of the broken part (i.e., reduced area due to symmetry breaking ($\delta \neq 0$ nm) of bow-tie shaped nanohole) and GaP based dielectric part, respectively. Here, we adopted the normalized perturbation since it quantifies the fractional modification of the GaP dielectric region rather than a single structural displacement. Moreover, the Q_{rad} experienced a decrease with the increase of α, resulting from the wider radiation channel, which gave rise to more energy loss through radiation. Figures 3(b) and (c) illustrate the extracted Q_{rad} as a function of α for Modes III and IV, which maintained the square inverse ratio law $Q_{rad} \propto (\alpha)^{-2}$. This trend satisfied the fundamental criterion for symmetry-protected BIC, indicating that the Fano resonance was governed by the BIC. This linear relation resulted from the incorporated α, which represented multiple geometric contributions simultaneously, rather than the

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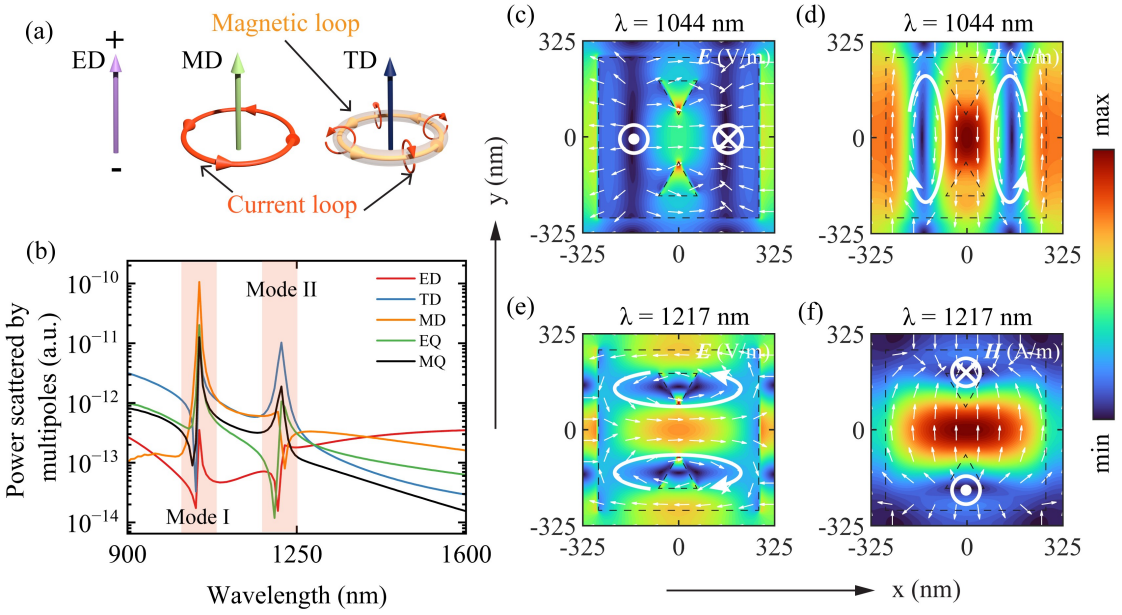

Fig. 4 (a) 3D schematic representation of ED, MD, and TD modes. (b) Scattering power contributions from individual multipole moments—including ED, TD, MD, EQ, and MQ—for the proposed structure at $\delta = 0$ nm, where the red-shaded region highlights resonant modes I and II. Spatial distribution of electric and magnetic fields in the x-y plane (c-d) at $\lambda = 1044$ nm and (e-f) at $\lambda = 1217$ nm. White arrows indicate the orientation of the corresponding field vectors.

actual symmetry-breaking perturbation.

The underlying physical mechanism responsible for the emergence of Fano resonances can be elucidated through multipole decomposition. This method resolves the total response into distinct contributions, including the electric dipole (ED), magnetic dipole (MD), electric quadrupole (EQ), magnetic quadrupole (MQ), and toroidal dipole (TD), moments of these are mathematically defined as[13],

$$P = \frac{1}{i\omega}\int J\, d^3r, \tag{7}$$

$$M = \frac{1}{2c}\int [\mathbf{r}\times\mathbf{J}]\, d^3r, \tag{8}$$

$$Q_{ab} = \frac{1}{2i\omega}\int \left[r_a J_b + r_b J_a - \frac{2}{3}(\mathbf{r}\cdot\mathbf{J})\right] d^3r, \tag{9}$$

$$M_{ab} = \frac{1}{3c}\int [(\mathbf{r}\cdot\mathbf{J})_a r_b + r_a[\mathbf{r}\times\mathbf{J}]_b r_a]\, d^3r, \tag{10}$$

$$T = \frac{1}{10c}\int \left[(\mathbf{r}\cdot\mathbf{J})\mathbf{r} - 2r^2\mathbf{J}\right] d^3r, \tag{11}$$

where $\mathbf{J}$, $\mathbf{r}$, and c represent the current density, displacement vector, and speed of light, respectively, while a and b denote the Cartesian tensor components.

Figure 4(a) presents the 3D schematics of the ED, MD, and TD modes. The corresponding scattering power contributions of each multipole (ED, MD, TD, EQ, MQ) for the case of $\delta = 0$ *nm* are shown in Fig. 4(b). As observed, at a distinct resonance wavelength, a specific multipolar component dominated the scattering response. Specifically, Mode I exhibited a dominant MD contribution, while Mode II was primarily influenced by the TD response. Consequently, the electromagnetic field distribution was primarily governed by the strongest dipolar excitation. Figures 4(c) and (d) show the spatial distributions of the electric and magnetic fields at $\lambda = 1044$ nm, while Figs. 4(e) and (f) correspond to $\lambda = 1217$ nm.

At $\lambda = 1044$ nm, the magnetic field distribution in the xy-plane featured two oppositely circulating loops, which drove an anticlockwise vortex of the electric field in the xz-plane. This configuration reflects the dominant influence of MD excitation at this resonance, oriented in the negative y direction. In contrast, at $\lambda = 1217$ nm, the electric field assumed a similar looped structure in the xy-plane, while the magnetic field exhibited a clockwise vortex in the yz-plane, revealing the dominance of TD excitation at this resonance in the positive x direction. For $\delta = 100$ nm, a pronounced TD excitation was identified at $\lambda = 911$ nm, whereas MD excitation was found to dominate at $\lambda = 1159$ nm and $\lambda = 1278$ nm [see ESI for details].

3.2 Rotational Asymmetry

The rotation of an individual triangular nanohole introduces in-plane symmetry breaking within our proposed structure. To examine this effect, we investigated the rotational dependence of a single nanohole in our structure with $\delta = 0$ nm, rotating from 0° to 360°. As shown in Fig. 5(a), Mode I remained at $\lambda = 1044$ nm throughout the entire rotation range, indicating its robustness to rotational asymmetry. In contrast, Mode II redshifted from $\lambda = 1217$ nm to $\lambda = 1221$ nm as the rotation increased from 0° to 180°. Upon further rotation to 360°, Mode II gradually blueshifted, ultimately returning to its original resonance wavelength at $\lambda = 1217$ nm. Meanwhile, an additional resonance was observed near Mode II with the rotation of this single nanohole. At a rotation angle of 90°, this resonance emerged at $\lambda = 1207$ nm with a modulation depth of 46.46 %. As the nanohole rotated further to 180°, the resonance exhibited a blue shift to $\lambda = 1201$ nm,

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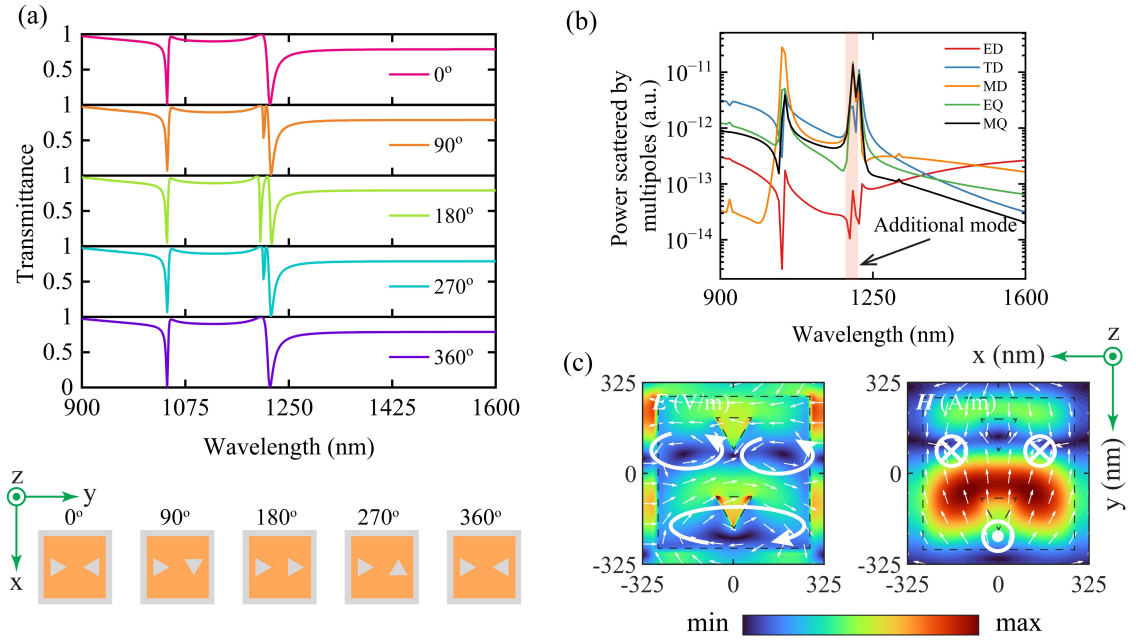

Fig. 5 (a) Transmittance spectra of the proposed structure, incorporating the rotational asymmetry (one triangular nanohole rotated from 0° to 360° as in the figure). (b) Scattering power contributions from individual multipole moments at the rotation of 180°. (c) Electric and magnetic field distribution profile at $\lambda = 1201$ nm at 180°.

accompanied by a significantly enhanced modulation depth of 93.53 %. The corresponding Q_{rad} was 9.86×10^2 at both 90° and 270°, but decreased to 4.40×10^2 at 180°. The observed blue shift and Q_{rad} reduction with increasing asymmetry suggest that this resonance originated from a quasi-BIC. To further elucidate the origin of the additional Fano resonance, the scattering power contributions from individual multipole moments were analyzed for a rotational angle of 180°. As shown in Fig. 5(b), the resonance at $\lambda = 1201$nm is predominantly governed by the MD, EQ, and MQ modes. This interpretation is further supported by the electric and magnetic field distributions in Fig. 5(c), where two anti-clockwise and one clockwise loops in the electric field confirm the coexistence of MD and EQ characteristics. Additionally, the localized out-of-plane magnetic field excitation indicates a strong presence of MD and MQ contributions.

3.3 Impact of Light Polarization

Owing to the asymmetric configuration of the proposed structure, we analyzed its optical response under varying polarization angles, φ, of the incident light. The analysis was carried out for φ ranging from 0° to 90°, considering structural asymmetries of $\delta = 0$ nm and 100 nm. The corresponding transmittance spectra are shown in Fig. 6(a) and (b). For $\delta = 0$ nm, Mode I maintained its resonance wavelength and modulation depth, though its Q-factor decreased to 1.38×10^2. In contrast, Mode II exhibited a gradual reduction in modulation depth with increasing φ, and fully vanished at $\varphi = 90°$. Interestingly, a new Fano profile emerged at $\lambda = 1175$ nm as φ increased, characterized by a Q-factor of 4.34×10^2 and a modulation depth of 99.96 % at $\varphi = 90°$. Figure 6(c) illustrates the transmittance at $\lambda = 1175$ nm, and 1217 nm for $\varphi = 0°$ to 360°, presenting the obvious on-off state of the corresponding resonance modes. For $\delta = 100$ nm, all Fano profiles except Mode I exhibited a gradual reduction of modulation depth with increasing φ, and totally disappeared at $\varphi = 90°$. However, four new Fano profiles started to appear with the increase of φ. At $\varphi = 90°$, new four Fano profiles were at $\lambda = 936$ nm, 1012 nm, 1030 nm, and 1148 nm with modulation depths of 73.06 %, 26.64 %, 80.07 %, and 98.21 %, respectively. The Q-factors of these four Fano profiles were 4.70×10^3, 2.95×10^3, 6.64×10^2, and 1.73×10^3, respectively. The transmittance at these particular wavelengths are shown in Fig. 6(d) for $\varphi = 0°$ to 360° to clearly present the on-off state. The aforementioned characteristics of both structures with $\delta = 0$ nm and 100 nm allow them to be utilized as an active optical switch.

3.4 Impact of Structural Parameters

The impact of the key structural parameters, namely, the thickness of the GaP cuboids, t_{GaP}, and the distance between nanoholes, $d_{nanoholes}$, on the optical performance of the structure was investigated. The corresponding optical performance is discussed below.

3.4.1 Impact of thickness of GaP cuboids

To examine the impact of the t_{GaP}, we varied the t_{GaP} from 50 nm to 250 nm for $\delta = 0$ nm and 100 nm, while keeping other structural parameters constant. Figures 7(a) and (b) illustrate the corresponding transmittance spectra for $\delta = 0$ nm and 100 nm, respectively. For $\delta = 0$ nm, the Mode I and II were located at $\lambda = 902$ nm and $\lambda = 997$ nm at $t_{GaP} = 50$ nm. However, as t_{GaP} increased, both Fano profiles exhibited a significant red shift, reaching $\lambda = 1230$ nm and $\lambda = 1332$ nm at $t_{GaP} = 250$ nm, primarily due to the increased effective refractive index of the structure. More-

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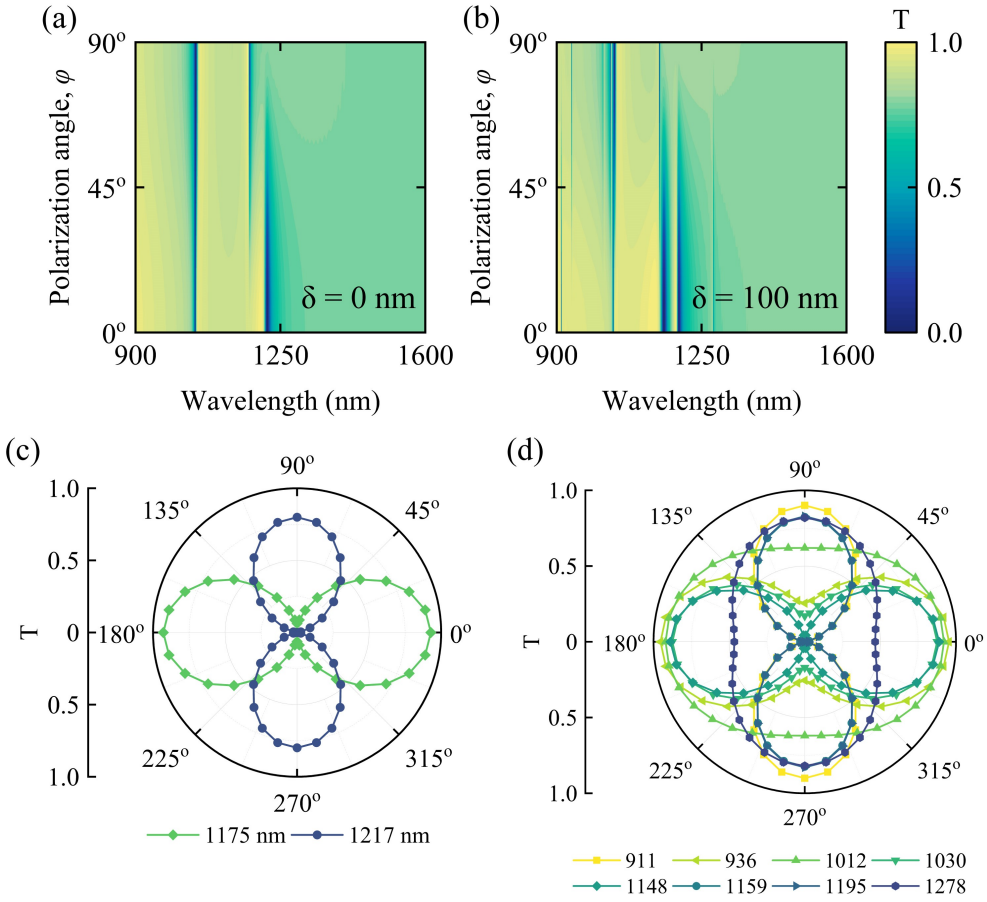

Fig. 6 Transmittance spectra for (a) $\delta = 0$ *nm*, and (b) $\delta = 100$ *nm* at polarization angle, $\varphi = 0°$ to 90°. Polar plot of transmittance spectra for $\varphi = 0°$ to 360° at (c) λ = 1175 nm and 1217 nm for δ = 0 nm and at (d) λ = 911 nm, 936 nm, 1012 nm, 1030 nm, 1148 nm, 1159 nm, 1195 nm, and 1278 nm for δ = 100 nm.

over, some additional resonant modes were observed for further increase of $t_{GaP} > 175$ nm, suggesting the emergence of additional resonant pathways due to increased optical thickness. However, for $\delta = 100$ nm, these additional modes appeared earlier due to introducing asymmetry in nanoholes, and a single mode (Mode III) was found for our proposed structure with t_{GaP} = 150 nm. Despite the asymmetry ($\delta = 100$ nm), the overall mode evolution with increasing t_{GaP} followed a similar trend as the symmetric case ($\delta = 0$ nm). Specifically, at t_{GaP} = 50 nm, the Mode I and II were at $\lambda \sim 900$ nm and $\lambda = 990$ nm, redshifted to $\lambda = 1214$ nm and 1302 nm at t_{GaP} = 250 nm. Meanwhile, Mode III emerged around t_{GaP} of 145 nm at $\lambda = 906$ nm, and gradually shifted to $\lambda = 998$ nm at t_{Gap} = 250 nm. In addition, Modes IV and V, initially located at $\lambda = 980$ nm and 1036 nm for t_{GaP} = 50 nm, were redshifted to $\lambda = 1251$ nm and 1405 nm, respectively, as t_{GaP} increased to 250 nm.

3.4.2 Impact of the distance between nanoholes

The distance between the triangular nanoholes, $d_{nanoholes}$, in the bow tie-shaped structure was varied to investigate the impact on optical responses for both symmetric ($\delta = 0$ nm) and asymmetric (100 nm) cases. The range of $d_{nanoholes}$ was considered from 0 nm to 200 nm, and the resulting transmittance spectra are presented in Figs. 7(c) and (d) for symmetric ($\delta = 0$ nm) and asymmetric (100 nm) cases, respectively. In a symmetric structure, the Mode I remained almost at the same resonance wavelength with the increase of $d_{nanoholes}$. However, the Mode II experienced both redshift and blueshift, starting at $\lambda = 1143$ nm for $d_{nanoholes}$ = 0 nm, and ending at $\lambda = 1212$ nm for $d_{nanoholes}$ = 200 nm. For the asymmetric case, a similar dual-shift behavior was observed for Mode IV, which shifted from $\lambda = 1135$ nm to $\lambda = 1152$ nm across the same range of $d_{nanoholes}$, while both Modes I and II remained at nearly the same resonance wavelengths. Notably, Mode III and V exhibited contradictory trends in activation: Mode III appeared at larger $d_{nanoholes}$, and disappeared at smaller $d_{nanoholes}$, while Mode V displayed inverse behavior.

3.5 Sensing of *Vibrio Cholerae*

Owing to the high quality factor, Q_{rad} exhibited by our proposed structure, we further explored its potential for refractive index (RI) sensing. Specifically, the structure was employed to detect *Vibrio cholerae*, a bacterium with an approximate refractive index of 1.365[41]. Nevertheless, practical detection of *Vibrio cholerae* would depend on implementing proper biochemical surface functionalization and microfluidic handling of the analyte. The corresponding transmittance spectra for a single bacterium for both $\delta = 0$ nm and 100 *nm* are illustrated in Fig. 8(a), where the associ-

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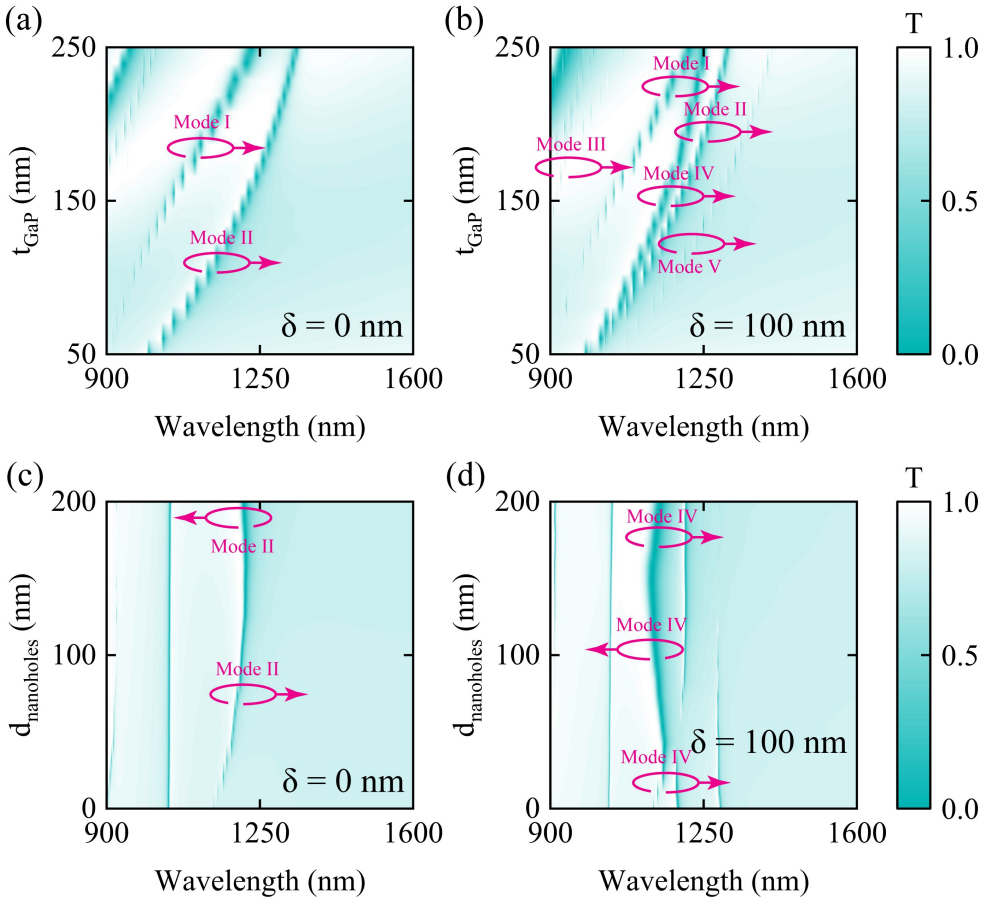

Fig. 7 Transmittance spectra for varying t_{GaP} from 50 nm to 250 nm at (a) $\delta = 0$ nm and (b) $\delta = 100$ nm and for varying $d_{nanoholes}$ from 0 nm to 200 nm at (c) $\delta = 0$ nm and (d) $\delta = 100$ nm.

ated resonance wavelength shifts are also marked. For both cases, the most prominent resonance shift was associated with Mode I, which was 10.75 nm and 10.58 nm for $\delta = 0$ nm and 100 nm, respectively; however, the modulation depth was lower (74.51 %) for $\delta = 100$ nm. Interestingly, Mode III at $\delta = 100$ nm offered a strong trade-off, exhibiting both a high Q_{rad} and significant modulation depth, alongside a resonance shift of 7.13 nm. Figures 8(c) and (d) depict the transmittance spectra for an increasing number of *V. cholerae* bacteria (0 to 10) for $\delta = 0$ nm, and 100 nm, respectively. In both cases, the resonance wavelength exhibited a redshift, attributed to the increase in the effective refractive index introduced by additional bacteria cells. The corresponding performance parameters, including sensitivity (S) and figure of merit (FOM) are shown in Figs. 9(a) and (b), respectively, which were calculated by,

$$S = \frac{\Delta\lambda}{\Delta n}, \tag{12}$$

$$FOM = \frac{S}{FWHM}. \tag{13}$$

Here, $\Delta\lambda$, Δn, and FWHM represent the resonance wavelength shift, the change in refractive index, and the full width at half maximum of the resonance, respectively. As previously discussed, Mode I exhibited the most prominent resonance wavelength shifts for both $\delta = 0$ nm and $\delta = 100$ nm, which was directly reflected in its sensitivity. For single bacterium detection, Mode I at $\delta = 0$ nm yielded the highest sensitivity among all modes, with a value of 306.85 nm/RIU. A comparable sensitivity of 302.28 nm/RIU was observed for Mode I at $\delta = 100$ nm. In both symmetric and asymmetric cases, Mode I exhibited a consistent increase in sensitivity with the number of bacteria, reaching approximately 342 nm/RIU when ten bacteria were present.

Since the FOM is strongly related to the quality factor, Mode III ($\delta = 100$ nm) achieved significantly higher FOM values than the other modes, maintaining an average around 200 RIU^{-1}, peaking at 217.14 RIU^{-1} for three bacteria. Interestingly, despite its superior sensitivity, Mode I at $\delta = 100$ nm showed relatively lower FOM values, fluctuating around 124 RIU^{-1}. Consequently, Mode III ($\delta = 100$ nm) outperformed in terms of FOM, indicating its potential advantage for applications requiring high spectral resolution and detection precision.

4 Comparative Analysis

Table 2 summarizes the key performance metrics of previously reported all-dielectric metastructures exhibiting Fano resonances, highlighting the compelling performance balance achieved by the proposed bowtie-shaped nanohole-based cuboid metasurface. Zhang *et al.* employed a lucky knot structure using Si and BaF_2 materials in the mid-infrared region, yielding a relatively low Q-

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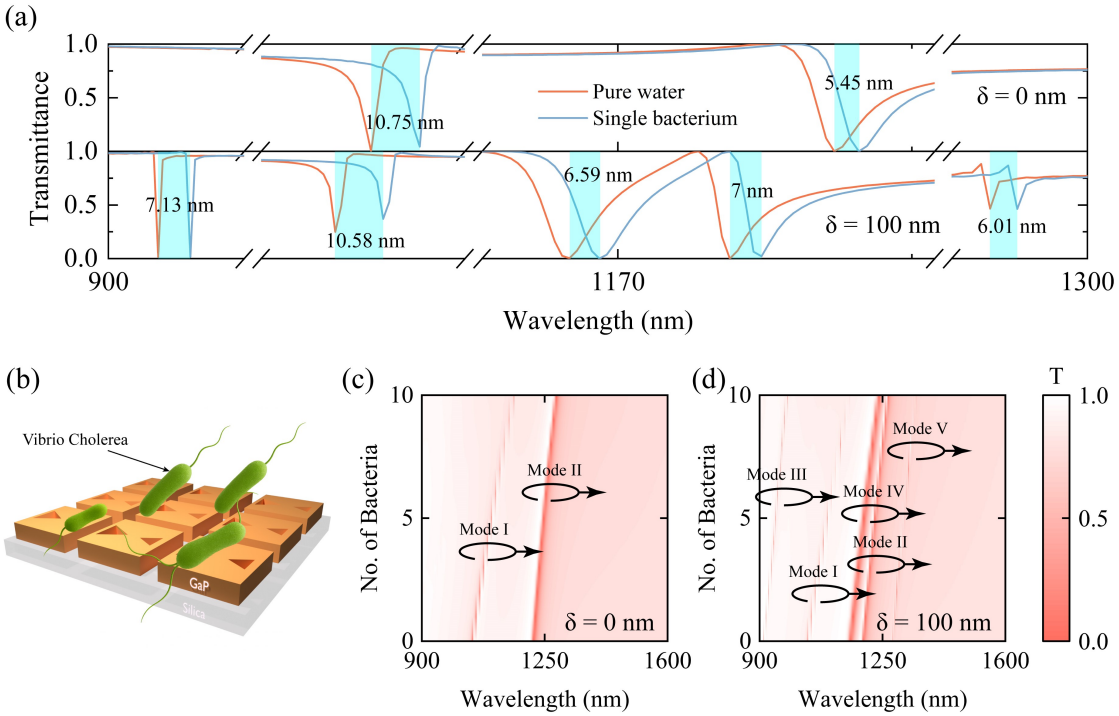

Fig. 8 (a) Transmittance spectra for pure water and a single *Vibrio cholerae* bacterium for $\delta = 0$ nm and 100 nm. Resonance shifts due to the presence of the bacterium are indicated by blue shading. (b) Schematic illustration of *V. Cholerae* sensing by the proposed structure. (c) and (d) Transmittance spectra for 0–10 V. cholerae bacteria at $\delta = 0$ nm and 100 nm, respectively.

factor of 5.20×10^2 and a FOM of only 32.7 RIU^{-1}, although the authors reported a high sensitivity (S) of 986 nm/RIU[42]. Li *et al.* designed a structure with a nanocube with two square air holes operating in the 1100–1600 nm range, which achieved an ultra-high Q-factor of 1.54×10^5 and nearly 100 % modulation depth by exploiting TD and MD resonances[22]. However, the derived FOM of 389 RIU^{-1}, and the sensitivity 287.5 nm/RIU remained moderate. In another work, cuboid tetramer clusters reported by Yu *et al.* demonstrated strong multipolar interactions (TD, MQ) and yielded a FOM of 910 RIU^{-1}, with a Q-factor of 1.13×10^4 and S = 182 nm/RIU, but again operated over a narrower spectral window (1080–1200 nm)[37]. Lv *et al.* pushed the sensing performance further with a four-pillar with GaP, showing an extremely high S of 488.99 nm/RIU and an exceptional FOM of 2.51×10^5 RIU^{-1}, facilitated by a combination of EQ, MQ, and MD resonances[13]. Another structure comprised of GaP-based hollow nanocuboids was proposed by Algorri *et al.*, reaching a Q-factor of around 10^5 along with the sensitivity of 200 nm/RIU[43]. Pang *et al.* introduced a nanoblock tetramer cluster incorporating square defects, achieving a Q-factor of 1.66×10^4 and nearly 100 % modulation depth. Their structure supported TD and MQ excitations, delivering an excellent sensitivity of 256 nm/RIU and a high FOM of 2519.7 RIU^{-1} [36]. Nousios *et al.* utilized Si on sapphire substrate to make a square slot etched structure with a silicon bridge. By changing the width of this bridge, they achieved resonances with Q-factors of up to 10^7, governed by TD, MQ, and EQ[44]. Algorr *et al.* proposed a similar structure using Si and SiO_2, achieving a similar Q-factor, sensitivity S = 58.9 nm/RIU, and FOM = 11.7×10^4 RIU^{-1} [45]. Li *et al.* proposed an I-shaped bar and Φ-shaped disk configuration with structural asymmetry based on distance variation[22]. The design achieved a Q-factor of 1.77×10^4 and high sensitivity S = 784.8 nm/RIU through TD and MD excitation.

In contrast, our work presents a bowtie-shaped nanohole-based cuboid comprised of GaP, designed for the near-infrared region, a range more compatible with on-chip integration and bio-sensing. The incorporated geometric asymmetry enables the excitation of both TD and MD modes, peaking Q-factor of 6.38×10^4, near-unity modulation depth, and an appreciable sensitivity of 342 nm/RIU. The corresponding FOM of 217.14 RIU^{-1} surpassed most of recent works while maintaining a compact and fabrication-friendly configuration.

5 Suggested Fabrication Techniques

Although our study was conducted through numerical simulation, previous studies have corroborated the experimental feasibility of fabricating the proposed structure. The fabrication sequence, from the initial SiO_2 substrate to the completed device, is schematically presented in Fig. 10. The GaP cuboid layer can be deposited through radio-frequency (RF) sputtering[46], a well-established technique for achieving uniform and high-quality semiconductor thin films. To define the bow-tie-shaped nanohole patterns, electron beam lithography (EBL) can be employed. This process starts with the spin-coating of an EBL-sensitive resist layer on top of the GaP surface. After patterning the resist using high-resolution EBL writing, the design can be transferred into the GaP layer via inductively coupled plasma (ICP) etching[22]. Subsequent removal of the residual resist will complete the fabrication, yield-

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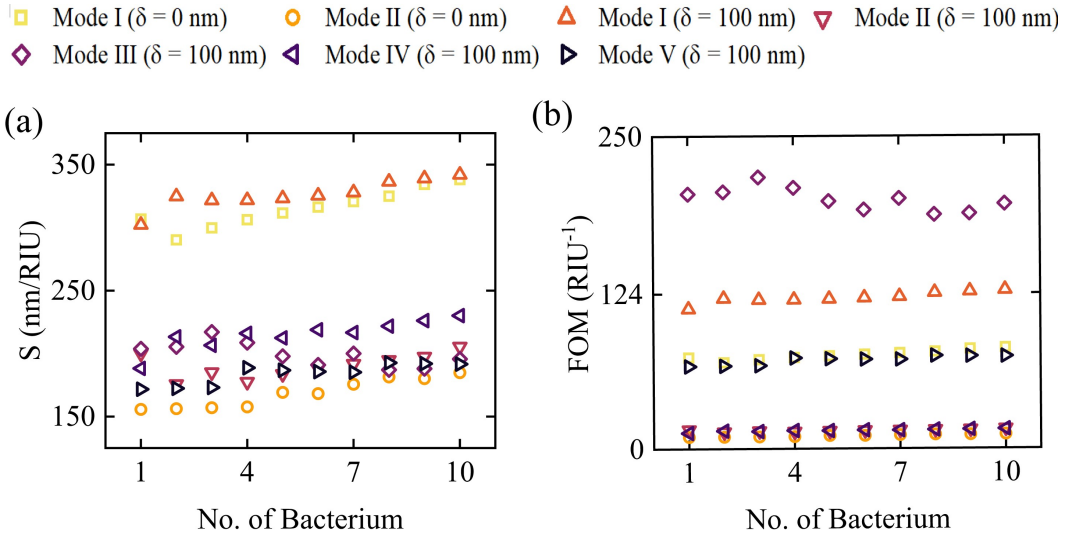

Fig. 9 (a) Sensitivity (S) and (b) figure of merit (FOM) corresponding to the resonance shifts for 1–10 bacteria under both asymmetry parameters ($\delta = 0$ nm and 100 nm.)

ing the final device structure with the designed nanophotonic features. Since practical fabrication inevitably introduces nanoscale deviations from the nominal geometry, we conducted a study on the fabrication tolerance by considering the effect on Q_{rad} by dimensional perturbation of ±10 nm, ±5 nm, and 0 nm. [See ESI for details]

6 Conclusions

We proposed a multifunctional all-dielectric metasurface, harnessing precisely engineered structural asymmetry to produce multiple high-Q quasi-BIC Fano resonances. The physical origins of these modes were clarified through multipolar decomposition, revealing dominant MD and TD contributions. To further validate the resonance features, LC circuit modeling and theoretical Fano profile fitting were employed, while the BIC nature was confirmed through the square inverse law. The proposed structure offered a remarkable combination of spectral sensitivity, modulation depth, and design simplicity with Q-factor approaching 10^4, pronounced polarization selectivity, and notable refractive index sensing performance, where sensitivity of 342 nm/RIU and FOM of 217.14 RIU^{-1} were achieved for *Vibrio Cholerae*. Furthermore, its CMOS compatibility and tunable optical response significantly broadened its applicability, reflecting it as a powerful candidate for next-generation integrated photonic technologies, including high-performance biosensing, active optical modulation, nonlinear light–matter interactions, and beyond.

Author contributions

Soikot Sarkar: Conceptualization, Methodology, Visualization, Software, Investigation, Writing – original draft, Writing – review & editing. **Ahmed Zubair**: Conceptualization, Methodology, Visualization, Resources, Supervision, Writing – original draft, Writing – review & editing.

Conflicts of interest

The authors declare no conflicts of interest.

Data availability

The simulation files for this article are available on GitHub: `https://github.com/ahmedzubair003/Quasi_BIC_Fano_Dielectric_Metasurfaces`

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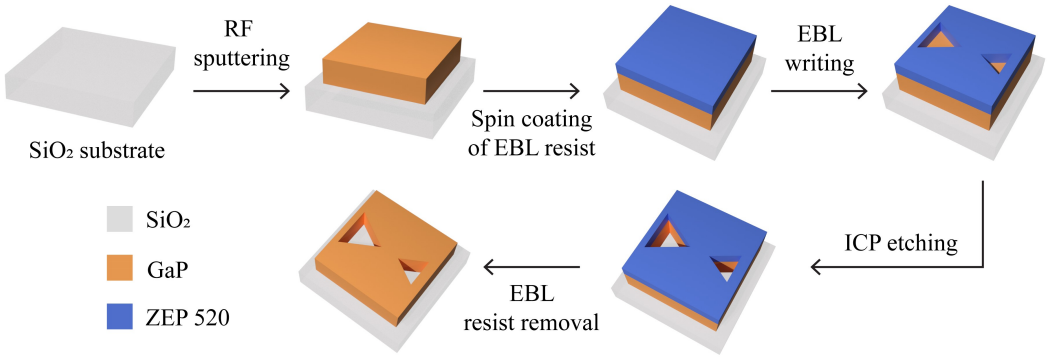

Fig. 10 Suggested fabrication techniques of our proposed metasurface structure with possible steps from SiO_2 substrate to final structure.

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Table 2 Comparison of all-dielectric metastructures exhibiting Fano resonances, highlighting key performance metrics. The last row corresponds to the proposed structure.

Structure Type	Material	Operating Wavelength (nm)	Asymmetry Parameter	Quality Factor	Modulation Depth	Excitation Mode	S (nm/RIU)	FOM (RIU^{-1})	Ref
Lucky knot	Si, BaF_2	6500–8000	–	5.20×10^2	–	–	986	32.7	42
Nanocube with two square air holes	Si, SiO_2	1100–1600	Air hole dimension	1.54×10^5	nearly 100%	TD, MD	287.5	389	22
Cuboid tetramer clusters with symmetric structural parameters	–	1080–1200	Permittivity	1.13×10^4	nearly 100%	TD, MQ	182	910	37
Four cylindrical pillars	GaP, MgF_2	1030–1180	Placement of two pillars	1.47×10^4	–	EQ, MQ, MD	488.99	2.51×10^5	13
hollow nanocuboid	GaP, SiO_2	-	hole position	$\sim 10^5$	-	ED	200	-	43
Nanoblock tetramer clusters	Si, SiO_2	1000–1350	Square defects	1.66×10^4	nearly 100%	TD, MQ	256	2519.7	36
square slots etched	Si, Sappier	1300-2000	width of silicon bridge	$\sim 10^7$	-	TD, MQ, EQ	-	-	44
square slot rings etched	Si, SiO_2	1400 - 1700	width of silicon bridge	$\sim 10^7$	-	-	58.9	11.7×10^4	45
I-shaped bar and ϕ-shaped disk	Si	1100–1300	Distance between bar and neighboring disks	1.77×10^4	–	TD, MD	784.8	–	22
Bowtie-shaped nanohole-based cuboid	GaP, SiO_2	900–1600	Dimension of triangular nanohole	6.38×10^4	nearly 100%	TD, MD	342	217.14	This Work

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Data Availability Statement

Data for this article are available at GitHub at https://github.com/ahmedzubair003/Quasi_BIC_Fano_Dielectric_Metasurfaces.git