

Cite this: *Dalton Trans.*, 2026, **55**, 2919

Two-dimensional metal–organic framework/graphene oxide composites as proton conductors: chemical grafting vs. physical blending

Chi-Lun Chuang,^a Chou-Hung Hsueh,^a Kuan-Chu Wu,^a Cheng-Hui Shen,^a Tzu-Chi Lin^a and Chung-Wei Kung^{*a,b,c}

A water-stable two-dimensional (2D) zirconium-based metal–organic framework (MOF), ZrBTB (BTB = 1,3,5-tri(4-carboxyphenyl)benzene), serves as a porous platform to integrate with graphene oxide (GO) in order to achieve ultrahigh proton conductivity (σ). Two synthetic methods, chemical grafting and physical blending, are used for preparing various nanocomposites composed of ZrBTB and GO. The porosity, morphology and proton-conducting characteristics of these composites and both pristine materials are investigated. The nanocomposite obtained by the grafting method with a GO loading of around 1 wt%, ZrBTB–0.01GO, possesses coordination bonds between GO and the hexa-zirconium nodes of ZrBTB; such chemical grafting reduces the number of accessible –OH/–OH₂ pairs on the MOF nodes. ZrBTB–0.01GO thus exhibits a worse proton-conducting performance, with a σ of $3.93 \times 10^{-3} \text{ S cm}^{-1}$ at 60 °C and 99% relative humidity (RH), compared to those of the pristine ZrBTB ($1.57 \times 10^{-2} \text{ S cm}^{-1}$) and pristine GO ($8.68 \times 10^{-3} \text{ S cm}^{-1}$), respectively. In contrast, considerably increased proton conductivities and decreased activation energies are achieved with the same GO loading by the physical blending method, with ZrBTB/0.01GO exhibiting a σ of $4.01 \times 10^{-2} \text{ S cm}^{-1}$ at 60 °C and 99% RH and an activation energy of 0.28 eV. At an optimal GO loading, the resulting 2D nanocomposite, ZrBTB/0.005GO, can achieve an ultrahigh σ of $1.03 \times 10^{-1} \text{ S cm}^{-1}$ at 60 °C under 99% RH with a low activation energy of 0.18 eV. The findings here suggest that compared to the commonly reported chemical grafting method, the physical blending method is more advantageous for preparing proton-conductive 2D MOF-based nanocomposites with more accessible proton-relaying functional groups and thus better performance.

Received 5th December 2025,
Accepted 12th January 2026

DOI: 10.1039/d5dt02918j

rsc.li/dalton

Introduction

Metal–organic frameworks (MOFs), as a class of highly ordered and designable porous materials, have undergone rapid development since the early 21st century.^{1,2} They have been extensively applied in various fields, including gas adsorption and storage,^{3,4} separation,^{5,6} catalysis,^{7,8} sensing,^{4,9,10} and energy storage.^{11,12} In addition to the aforementioned applications, benefiting from their high porosity, structural tunability to incorporate hydrophilic functional groups, and insulating nature for electrons, MOFs have also attracted great attention as promising proton-conductive materials.^{13–20} On the other hand, acidic water electrolyzers are crucial devices to produce green hydrogen, and proton exchange membrane fuel cells

(PEMFCs) are known as one of the most promising energy-conversion technologies for hydrogen utilisation.²¹ In both systems, the proton exchange membrane (PEM) capable of conducting protons in acidic aqueous environments plays a crucial role in determining the overall performance.²² Consequently, stable MOFs with superior proton-conducting characteristics are highly attractive materials for energy-related applications, including fuel cells and PEM electrolyzers.

It is well known that most MOFs exhibit poor chemical stability in aqueous media.²³ Recent developments in chemically robust tetravalent MOFs, including titanium-, zirconium-, cerium(IV)-, and hafnium-based frameworks, have opened new possibilities for utilising MOFs in humid or aqueous environments.^{23–26} Hence, these stable MOFs have been widely investigated as proton conductors, with most studies focusing on three-dimensional (3D) MOFs to date.^{15,27–29} Compared to 3D MOFs, two-dimensional (2D) MOFs with crystallographically well-defined 2D structures are supposed to offer shorter hopping distances for ionic species to move between stacked molecular sheets. However, studies reporting proton conduction in 2D tetravalent MOFs are still very rare.^{30–33} Recently,

^aDepartment of Chemical Engineering, National Cheng Kung University, Tainan City, 70101, Taiwan. E-mail: cwkung@mail.ncku.edu.tw

^bProgram on Key Materials, Academy of Innovative Semiconductor and Sustainable Manufacturing, National Cheng Kung University, Tainan City, 70101, Taiwan

^cCenter for Resilience and Intelligence on Sustainable Energy Research (RiSER), National Cheng Kung University, Tainan City, 70101, Taiwan



Wang and co-workers reported the first study on the proton-conducting characteristics of ZrBTB (BTB = 1,3,5-tri(4-carboxyphenyl)benzene), a six-connected 2D zirconium-based MOF, and its proton conductivity (σ) was found to be much higher than those of several pristine 3D tetravalent MOFs.³³ We thus reasoned that stable 2D MOFs, such as ZrBTB, should be promising candidates for designing high-performance proton conductors.

To further enhance the σ of the pristine MOFs, several strategies have been widely employed, such as the introduction of hydrophilic/proton-relaying functional groups and the incorporation of proton-conductive guest molecules into the MOFs.^{16,33–35} One of the effective approaches is the integration of a porous MOF with another proton-conductive material, such as graphene oxide (GO), to form nanocomposites with heterostructural interfaces.^{36–38} GO with a high degree of oxidation is electrically insulating, and it possesses abundant oxygen-containing functional groups to not only facilitate the adsorption of water molecules but also extend the proton-transport pathways.³⁹ MOF–GO composites have thus been reported as better proton conductors than both pristine materials. For example, ZIF-8@GO,³⁶ GO@UiO-66-NH₂,³⁷ and MIL-101/GO³⁸ were reported as excellent proton conductors with a Grotthuss-type mechanism.^{40,41} However, it should be noted that all the aforementioned studies investigating proton conduction in MOF–GO composites utilised 3D MOFs grown or grafted on GO, with chemical bonds formed between MOF crystals and GO layers. The resulting σ could usually be enhanced owing to the exfoliation of GO layers caused by the bulky MOF crystals present in between. In contrast, the physical blending of MOF crystals and GO was usually employed as a referencing method.^{36,38} However, it is worth mentioning that the chemical grafting or growth of MOF crystals on GO sheets should also reduce the accessible hydrophilic groups exposed on the external surface of both materials, which is not desirable for proton conduction. This effect may be more sig-

nificant when small MOF crystals or even 2D MOF nanosheets are used. The effects of synthetic approaches for MOF–GO composites on the resulting proton-conducting performance have barely been explored.

In this study, a chemically robust 2D MOF, ZrBTB,^{42–44} was employed as a platform for integration with GO. As the first study reporting composites with a 2D MOF and GO as proton conductors, two synthetic methods, including chemical grafting (ZrBTB–GO) and physical blending (ZrBTB/GO), were employed; see Fig. 1. With the use of physical blending, we aim to develop a new type of 2D MOF-based nanocomposite with an ultrahigh σ , featuring both shorter pathways for proton transport and a high density of protonic carriers. The results suggest that the optimised nanocomposite, ZrBTB/0.005GO, can achieve an ultrahigh σ of $1.03 \times 10^{-1} \text{ S cm}^{-1}$ at 60 °C under 99% relative humidity (RH) with a comparatively low activation energy (E_a) of 0.18 eV, which is significantly better than those of both the pristine materials and the chemically grafted composite.

Experimental

Chemicals

Zirconium(IV) chloride (ZrCl₄, 98%), potassium nitrate (KNO₃, 99%) and 1,3,5-tri(4-carboxyphenyl)benzene (H₃BTB, 97%) were obtained from Thermo Fisher Scientific. Benzoic acid (BA, 99.5%), potassium carbonate (K₂CO₃, 98%), sodium bromide (NaBr, 98%), sodium chloride (NaCl, 99%) and deuterium oxide (D₂O, 99.9 atom% D) were purchased from Sigma-Aldrich. *N,N*-Dimethylformamide (DMF, ≥99.8%), dimethyl sulfoxide (DMSO, 99%) and potassium chloride (KCl, 99.0%) were obtained from Duksan Pure Chemicals. Hydrochloric acid (HCl, 36.5–38.0%) was received from J.T. Baker. Potassium bromide (KBr, ≥99.0%) was purchased from Honeywell Fluka. Acetone (≥99%) and ethanol (95%) were

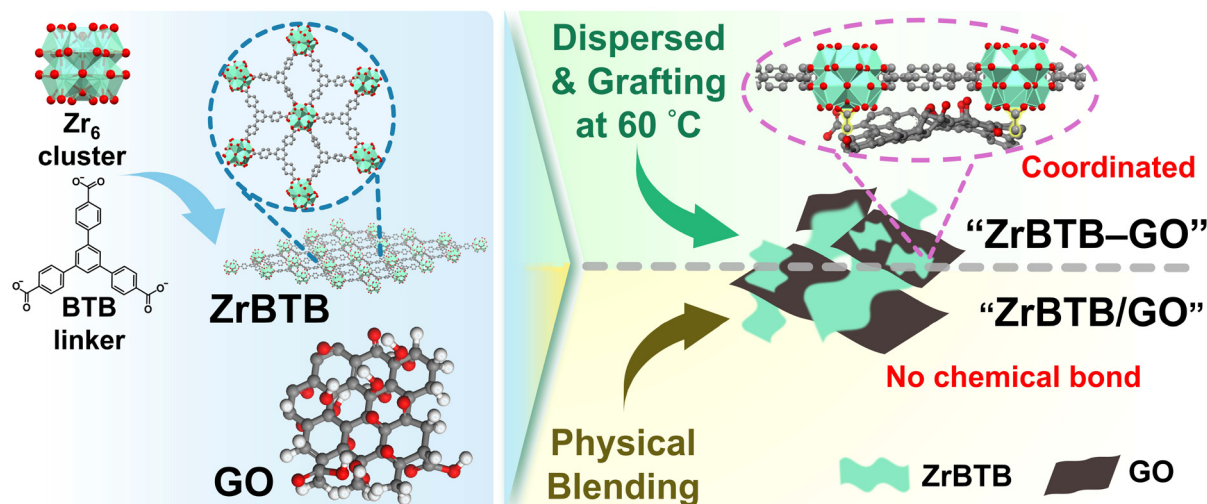


Fig. 1 Schematic diagram for the synthesis of MOF–GO nanocomposites by chemical grafting or physical blending.



obtained from Uni-Onward Co., Ltd, Taiwan. Graphene oxide (GO, model number: EW-GO1, 97%, oxygen content = 45%) powder was purchased from E-WAY Technology Co., Ltd, Taiwan. All the chemicals mentioned above were used as received without purification. Deionised water was used throughout the entire work.

Synthesis of materials

The synthesis of ZrBTB was performed according to the procedure in our recent work.^{45,46} First, ZrCl₄ (100 mg), H₃BTB (100 mg), and BA (6 g) were dissolved in a mixture of DMF (30 mL) and water (5 mL) using ultrasonication. The resulting homogeneous solution was then transferred into a DURAN® 100 mL glass laboratory bottle and sealed with a polytetrafluoroethylene (PTFE)-lined polybutylene terephthalate cap, and the bottle was then placed in an oven at 120 °C for 48 h. After the MOF growth, the mixture was cooled to room temperature (RT) and centrifuged at 5000 rpm to remove the supernatant. The centrifugation was performed three times, after adding 30 mL of fresh DMF every time, to effectively remove excess reagents. Afterwards, the powder was subjected to centrifugal washing with acetone (30 mL) three times to ensure complete solvent exchange, with immersion periods of 2 h, overnight and 2 h between each step. The powder was then activated in a vacuum oven at 60 °C overnight. Subsequently, to effectively remove the coordinated benzoate, 100 mg of the activated MOF powder was dispersed in a prepared solution containing DMSO (30 mL), water (0.4 mL) and 12 M HCl (0.8 mL) through ultrasonication, and the mixture was maintained at RT for 18 h. Then, the solid was washed with DMSO (30 mL) four times over 8 h, followed by solvent exchange with acetone three times. The “ZrBTB” powder was finally obtained after activation in a vacuum oven at 60 °C overnight.

For the synthesis of MOF/GO nanocomposites through physical blending, 10 mg of GO was first dispersed in 10 mL of ethanol by ultrasonication for 1 h. Thereafter, a certain amount of the GO-containing suspension was added into an agate mortar, followed by the evaporation of the solvent at RT. Subsequently, 100 mg of ZrBTB was added, and the mixture was ground for 3 min. The obtained materials prepared with 0.1, 0.5, 1.0 and 3.0 mL of the GO suspension were named “ZrBTB/0.001GO”, “ZrBTB/0.005GO”, “ZrBTB/0.01GO” and “ZrBTB/0.03GO”, respectively.

To perform the chemical grafting between ZrBTB and GO, the following procedure from a reported study was used.⁴⁷ GO (10 mg) was dispersed in ethanol (10 mL) by ultrasonication for 1 h, and ZrBTB (100 mg) powder was dispersed in ethanol (4 mL) to form another suspension. Then, the GO-containing suspension (1 mL) was added into the ZrBTB-based suspension in a 20 mL scintillation vial, and the mixture was ultrasonicated for 10 min. Thereafter, the vial was sealed with a PTFE-lined cap and heated in an oil bath at 60 °C for 24 h under stirring at 200 rpm. After the reaction, the obtained solid was washed with ethanol (10 mL) three times through centrifugation, and the solvent exchange was performed with acetone (10 mL) three times. The obtained material was then

dried in a vacuum oven at 60 °C overnight, and it was designated as “ZrBTB-0.01GO”.

Preparation of pellets for measuring conductivities

The procedure for preparing pellets was slightly modified from that in our recently published work.²⁹ First, to allow a fair comparison with MOF/GO nanocomposites prepared by physical blending, powders of ZrBTB, GO and ZrBTB-0.01GO were ground for 3 min in an agate mortar before pelletisation. Thereafter, approximately 15 mg of the sample powder was filled into a die set ($d = 7$ mm) equipped with a mini-pellet press (Specac), and the pelletising process was performed with a load of 0.3 tons for 3 seconds. This process was repeated three times to form a stable pellet. The thickness of each pellet was measured using a digital calliper (Mitutoyo). The resulting pellet was then sandwiched between two titanium plates and dried in a vacuum oven at 60 °C overnight before any measurement.

Measurements of proton and electrical conductivities

The entire procedure was slightly modified from that reported in our previous work.²⁹ For measuring its proton conductivity, the sandwiched pellet was first placed in the headspace of a hermetically sealed glass chamber loaded with water, with two wires connecting to both sides of the pellet. The chamber was then placed in an oven at 30 °C for 48 h to achieve 99% RH. Thereafter, electrochemical impedance spectroscopy (EIS) of the pellet was performed using a PGSTAT204 potentiostat/galvanostat (Autolab, Eco-Chemie, The Netherlands) equipped with an FRA2 module under a two-electrode mode. Each EIS measurement was performed at an applied voltage of 0 V, with an amplitude of 10 mV and a frequency range from 1 MHz to 1 Hz. The proton conductivity of the pellet at each temperature was measured every 5 °C from 30 °C to 60 °C. At each temperature, EIS measurement was performed 30 min after the chamber had stabilised at the targeted temperature. The measurement of each material was performed three times. The same process was employed to measure proton conductivities of the pellet in D₂O vapour, except that water was replaced by D₂O.

For measuring the proton conductivities under various levels of RH, the chamber was set at 30 °C, and the same process mentioned above was used except that water was replaced with a saturated aqueous solution of K₂CO₃, NaBr, NaCl, KCl, or KNO₃ loaded in the sealed glass chamber; it could thus achieve a RH of 45%, 58%, 75%, 84% or 94% in the headspace, respectively, measured using a humidity sensor (GM1361, Benetech).

For measurements of electrical conductivity, the sandwiched pellet was connected to the CHI6273E electrochemical workstation (CH Instruments Inc.) to perform current–voltage (I – V) measurements with a two-electrode setup.²⁹ The I – V curve of each pellet was recorded at RT in air.

Instrumentation

X-ray diffraction (XRD) patterns of the samples were recorded using a SmartLab (Rigaku) with CuK α radiation ($\lambda =$



0.15406 nm). Nitrogen adsorption–desorption isotherms were collected at 77 K using a 3FLEX (Micromeritics) after degassing the powder samples overnight at 110 °C with a VacPrep (Micromeritics). Scanning electron microscopy (SEM) images were collected using an SU-8230 (Hitachi) at an operating voltage of 10 kV. A JEM-1400 Flash (JEOL) was used to obtain transmission electron microscopy (TEM) images and energy-dispersive X-ray spectroscopy (EDS) data at an operating voltage of 120 kV. Fourier transform infrared (FTIR) spectra were recorded on a Nicolet 6700 (Thermo Fisher Scientific). For preparing FTIR samples, the material for testing was ground with KBr followed by a pelletising process to obtain a pellet. Thermogravimetric analysis (TGA) experiments of the samples were conducted on a thermal analyser (Setaram Labsys evo) under a nitrogen flow. For preparing TGA samples, the solid material was placed in a hermetically sealed chamber at 30 °C and 99% RH for 48 h, followed by the TGA measurement immediately. X-ray photoelectron spectroscopy (XPS) spectra were collected using a Theta Probe angle-resolved XPS spectrometer (Thermo Fisher Scientific). All spectra were corrected by referencing the strongest C 1s peak originating from C–C bonds to 284.8 eV before analysing the data. During the synthesis of ZrBTB and its composites, a Heraeus Megafuge 16 centrifuge with a TX-400 rotor (Thermo Fisher Scientific) was utilised for centrifugation during all washing and solvent-exchange steps.

Results and discussion

Materials characterisation

ZrBTB was synthesised by a solvothermal method using BA as the modulator during crystal growth, followed by treatment with hydrochloric acid to replace the capping benzoate with terminal –OH/–OH₂ ligands on its hexa-zirconium nodes.⁴⁵ Nanocomposites composed of ZrBTB and the commercially available GO were then prepared by two synthetic methods. The crystallinities of all materials were first verified through powder XRD. As shown in Fig. 2(a) and Fig. S1, the XRD patterns of ZrBTB and all composites are consistent with the simulated pattern of ZrBTB, with the main diffraction peaks located at 5.2°, 9.0° and 10.4°, which correspond to the planes of 100, 110 and 200, respectively.^{42,45,46,48} This finding suggests that the crystalline structure of ZrBTB was successfully formed during MOF growth, and its crystallinity can be well preserved after integration with GO. In addition, the strong diffraction peak in the experimental XRD pattern of GO located at 11.1°, attributed to the interlayer distance between stacked graphene sheets,⁴⁹ is not observable in the patterns of all composites. Even in the pattern of ZrBTB/0.03GO with a relatively high GO loading of almost 3 wt%, this diffraction peak from stacked GO is still not observed (see Fig. S1). This absence indicates that in each nanocomposite, the stacked GO layers become mostly separated into small grains with poor long-range order; it thus confirms the successful formation of uniform composites with ZrBTB sheets and less stacked GO layers.

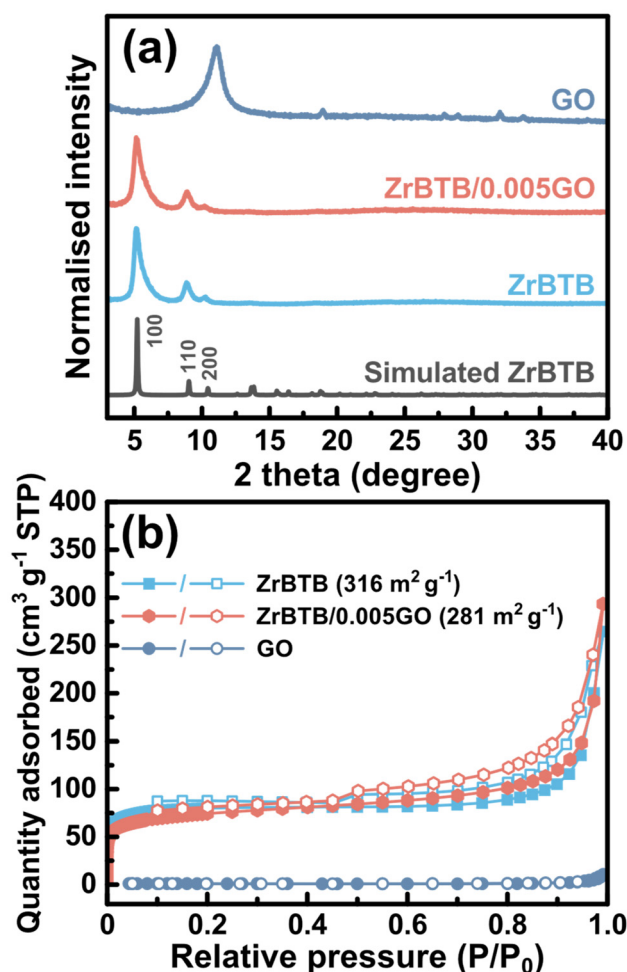


Fig. 2 (a) Powder XRD patterns and (b) N₂ adsorption–desorption isotherms of ZrBTB, GO and ZrBTB/0.005GO. The simulated pattern of ZrBTB is also shown in (a) and BET specific surface areas are provided in (b). Solid and empty scatters in (b) indicate data for adsorption and desorption, respectively.

Nitrogen gas (N₂) adsorption–desorption measurements were performed at 77 K to examine the porosity of each material. The obtained isotherms and specific surface areas calculated based on the Brunauer–Emmett–Teller (BET) method are shown in Fig. 2(b) and Fig. S2(a). All isotherms of ZrBTB and the MOF-based nanocomposites exhibit a steep gas uptake at low relative pressures and a hysteresis loop in the high-pressure region, corresponding to the microporosity of the stacked 2D sheets of ZrBTB and the mesoporous cavity between large aggregates of MOF sheets, respectively.⁵⁰ Furthermore, the BET surface area of ZrBTB was estimated as 316 m² g⁻¹. All these characteristics are consistent with those reported for ZrBTB.^{48,50} In contrast, as revealed in Fig. 2(b), GO shows negligible nitrogen uptake at 77 K. After integrating ZrBTB with GO, the BET surface area of each composite becomes slightly smaller than that of the pristine ZrBTB, but all composites still possess BET surface areas larger than 230 m² g⁻¹. It indicates that the major porosity of ZrBTB is still



preserved in every nanocomposite. Pore size distributions of the MOF-based materials were also extracted from their isotherms by using the model of density functional theory (DFT) for carbon-slit pores. As shown in Fig. S2(b), ZrBTB and all nanocomposites possess two kinds of main pore sizes ranging around 0.7–0.8 nm and 1.2–1.3 nm, which originate from the interlayer space between stacked molecular sheets of ZrBTB and the aperture on the 2D sheets of ZrBTB, respectively.^{45,48,50}

SEM and TEM images of ZrBTB, GO and three representative nanocomposites were collected in order to probe their morphologies, and the results are shown in Fig. 3. For ZrBTB, as revealed in Fig. 3(a), it is composed of loosely stacked 2D sheets with a flower-like morphology, and a similar morphology can also be observed in its TEM image (Fig. 3(f)). This morphology is consistent with that reported for ZrBTB in previous studies.^{46,51} On the other hand, GO exhibits a morphology of aggregates in its SEM image (Fig. 3(b)), but exfoliated GO sheets can be found in its TEM image (Fig. 3(g)). For all the three nanocomposites, including ZrBTB–0.01GO, ZrBTB/0.01GO and ZrBTB/0.005GO, a similar flower-like morphology of the 2D MOF can be observed in their SEM images. However, as shown in Fig. 3(c–e), tiny particles randomly attaching onto the flower-like MOF sheets can be found. Such morphologies become more obvious in the TEM images of these nanocomposites (Fig. 3(h–j)), revealing that these materials are composed of aggregates of loosely packed tiny 2D sheets. The findings here indicate that the incorporation of a small amount of GO into ZrBTB, either by chemical grafting or physical blending, can result in a significant morphological change from both pristine materials, but the 2D characteristic morphology can still be observed in all nanocomposites. The representative nanocomposite, ZrBTB/0.005GO, was then subjected to EDS measurements, and the results are shown in Fig. S3. The elemental mapping data in Fig. S3(d) indicate that all elements, including zirconium, carbon and oxygen, are uni-

formly located on those loosely packed 2D sheets in the nanocomposite.

To probe the interaction between ZrBTB and GO in the composites synthesised by various methods, FTIR spectra of all materials were collected with the use of their pellets ground with KBr. The obtained data are shown in Fig. S4, and the normalised spectra are plotted in Fig. 4. The characteristic peaks of ZrBTB located at 1416 and 1544 cm^{-1} are associated with the symmetric and asymmetric stretching vibrations of coordinated bidentate carboxylate groups, respectively.⁵² Other key peaks at 1589 and 1605 cm^{-1} , though not clearly defined

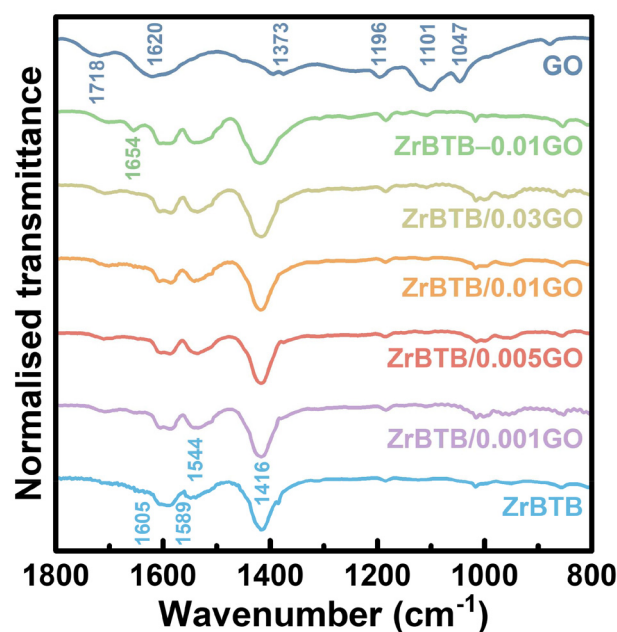


Fig. 4 FTIR spectra of ZrBTB, GO and all nanocomposites. Characteristic peaks are listed.

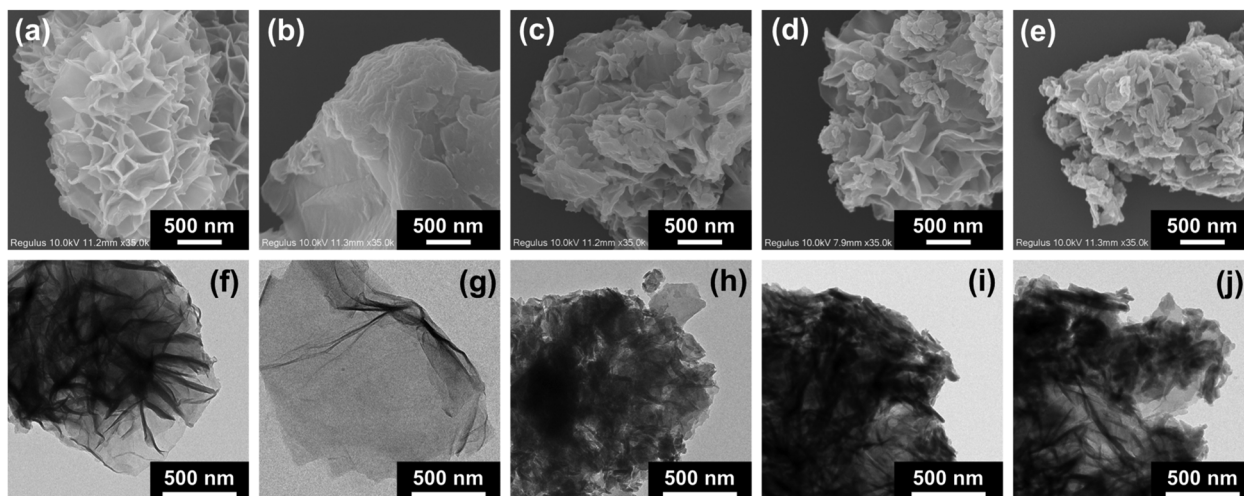


Fig. 3 SEM and TEM images of (a and f) ZrBTB, (b and g) GO, (c and h) ZrBTB–0.01GO, (d and i) ZrBTB/0.01GO and (e and j) ZrBTB/0.005GO.



in the literature, could be attributed to the C=C stretching vibrations of the aromatic ring in the BTB linker,^{45,48} or the C=O vibrations of coordinated unidentate carboxylate groups.⁵³ These characteristic peaks can be observed in all MOF-based materials. This result indicates that the coordination between linkers and hexa-zirconium nodes in the 2D MOF is well preserved in all composites. The FTIR spectrum of the pristine GO also displays several characteristic peaks, including the C=O stretching vibration of the carboxyl group at 1718 cm⁻¹, the O-H bending vibration of adsorbed water at 1620 cm⁻¹, the O-H bending vibration at 1373 cm⁻¹, the C-O stretching vibration of the epoxy group at 1196 cm⁻¹, and the C-O stretching vibration of alkoxy groups at 1101 and 1047 cm⁻¹, respectively; these features agree well with those of GO with a high degree of oxidation reported in previous studies.⁵⁴⁻⁵⁶ After integrating GO into ZrBTB, a small peak of C=O stretching from the free carboxylic group of GO, located at approximately 1700 cm⁻¹, can be observed in the FTIR spectra of all nanocomposites. It is worth noting that a new peak at 1654 cm⁻¹ can be clearly observed, exclusively in the FTIR spectrum of ZrBTB-0.01GO, which should be associated with the C=O/C-O vibrational features from the direct coordination of the carboxylate group in GO to the hexa-zirconium node of ZrBTB to form C-O-Zr bonds.³⁸ The finding here clearly indicates that, although ZrBTB-0.01GO and ZrBTB/0.01GO show almost the same crystallinity, porosity and morphology as discussed previously, they do possess quite different interactions between GO and the MOF. After the chemical grafting between co-dispersed GO and ZrBTB sheets at 60 °C, significant coordination bonds between GO and the MOF are present in the resulting ZrBTB-0.01GO.

XPS measurements were performed to further characterise ZrBTB, GO, ZrBTB-0.01GO and ZrBTB/0.01GO. From the survey spectra shown in Fig. S5, signals of the corresponding elements in each material can be detected. High-resolution XPS spectra were then recorded in the C 1s, O 1s and Zr 3d regions, and the obtained results are shown in Fig. S6, S7 and S8, respectively. Peaks in the region of C 1s were deconvoluted and identified according to a previous study.⁵⁷ On the other hand, the characteristic peaks of Zr-MOFs and GO in the region of O 1s were identified by following previous studies.⁵⁷⁻⁶⁰ As shown in Fig. S6 and S7, for both C 1s and O 1s regions, peaks of the corresponding functional groups in ZrBTB and GO can be found in their corresponding spectra, respectively. Peaks from both pristine materials are also in general observable in the spectra of both composites in Fig. S6 and S7. In addition, an obvious difference can be observed in the XPS spectra collected in the region of Zr 3d. As shown in Fig. S8, the physically blended composite ZrBTB/0.01GO shows almost the same XPS spectrum as that of the pristine ZrBTB, while the spectrum of ZrBTB-0.01GO reveals a positive shift of 0.4 eV in its Zr 3d peaks. Such a peak shift should be attributed to the change in the coordination environments of the hexa-zirconium nodes of the MOF, rendering the zirconium atoms more electron-deficient.^{61,62} Agreeing with the FTIR observation, the XPS result here clearly indicates the coordi-

nation of GO to the hexa-zirconium nodes of ZrBTB in the composite synthesised by chemical grafting.

Proton conductivity of ZrBTB, GO and nanocomposites

Since proton conductors usually require an electrically insulating feature, the electrical conductivities of ZrBTB, GO, and three representative nanocomposites were first quantified by measuring *I-V* curves of their pellets. Detailed protocols for such a two-probe *I-V* method can be found in our previous work.^{29,50} *I-V* data are plotted in Fig. S9, and the calculated values of electrical conductivity are listed in Table S1. ZrBTB and GO exhibit electrical conductivities of 3.46×10^{-12} S cm⁻¹ and 7.93×10^{-9} S cm⁻¹, respectively, suggesting that both pristine materials are electrical insulators. For all nanocomposites tested here, their electrical conductivities are in the range of 10^{-11} S cm⁻¹, which are around two orders of magnitude lower compared to that of GO.

The proton-conducting performances of all materials were then evaluated by performing EIS measurements. The σ of each material was measured at various temperatures (*T*) using

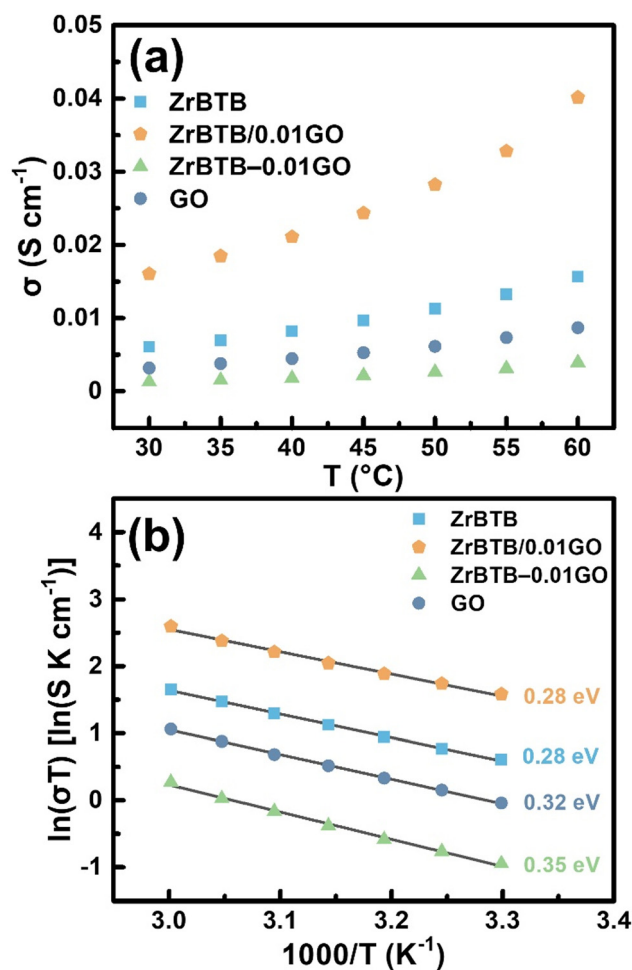


Fig. 5 (a) Proton conductivity at each temperature and 99% RH, and (b) Arrhenius plot obtained from the proton conductivities of ZrBTB, GO, ZrBTB-0.01GO and ZrBTB/0.01GO. Activation energies are listed in (b).



EIS spectra and calculated according to the method described in Section S1 of the SI. The resulting Nyquist plots of ZrBTB, GO, ZrBTB-0.01GO and ZrBTB/0.01GO are shown in Fig. S10, and the obtained σ values are plotted in Fig. 5(a). At 60 °C and 99% RH, the proton conductivities of the pristine ZrBTB and GO are $1.57 \times 10^{-2} \text{ S cm}^{-1}$ and $8.68 \times 10^{-3} \text{ S cm}^{-1}$, respectively, which are consistent with the values reported in previous studies.^{33,39} After the integration of both materials by the grafting method, ZrBTB-0.01GO exhibits a σ of $3.93 \times 10^{-3} \text{ S cm}^{-1}$ at 60 °C and 99% RH, which is inferior to both pristine materials. In contrast, ZrBTB/0.01GO, the composite prepared by physical blending, achieves a high σ of $4.01 \times 10^{-2} \text{ S cm}^{-1}$, which is about 2.5-fold higher than that of the pristine ZrBTB and 4.6-fold higher than that of the pristine GO. The E_a for proton conduction in each material was subsequently calculated from the linear fitting of Arrhenius plots; see the detailed protocols in Section S1 of the SI. The results for the above materials are shown in Fig. 5(b). The pristine materials, ZrBTB and GO, exhibit E_a values of 0.28 eV and 0.32 eV, respectively, which are consistent with previously reported values.^{33,39} For the composites, ZrBTB-0.01GO exhibits an E_a of 0.35 eV, which is higher than those of both pristine materials. In contrast, ZrBTB/0.01GO shows an E_a of 0.28 eV, the same as that of the pristine ZrBTB. The findings here clearly indicate that with around 1 wt% of GO in the 2D MOF-GO nanocomposite, the physically blended material can outperform both pristine materials for proton con-

duction, while the chemically grafted composite shows the worst proton-conducting performance compared to others.

Since the E_a values of ZrBTB, GO and the two nanocomposites are all below 0.4 eV, such results suggest that the proton transport within these materials predominantly proceeds *via* the Grotthuss mechanism. However, ZrBTB-0.01GO, in terms of both σ and E_a , exhibits the worst performance compared to both pristine materials, which is unexpected compared to the findings of several previously reported chemically grafted MOF-GO composites with 3D MOFs.³⁶⁻³⁸ One plausible reason is that the pristine ZrBTB already exhibits a higher proton conductivity than GO. Therefore, the prevention of GO from stacking through the coordination-driven grafting of ZrBTB sheets is insufficient to further enhance the proton conductivity. Furthermore, the coordination of ZrBTB and GO through carboxylate groups, as evidenced by the FTIR spectra, should largely reduce the numbers of accessible proton-relaying groups in both materials and thus disrupt the hydrogen-bonding network essential for proton transport. Such effects thus result in a decrease in σ and an increase in E_a after chemically grafting the 2D MOF and GO together. The results here suggest that for synthesising proton-conductive nanocomposites with 2D MOFs and GO, the physical blending method, which does not generate coordination bonds between both materials, should be more advantageous compared to the chemical grafting method that is typically used for 3D MOFs.

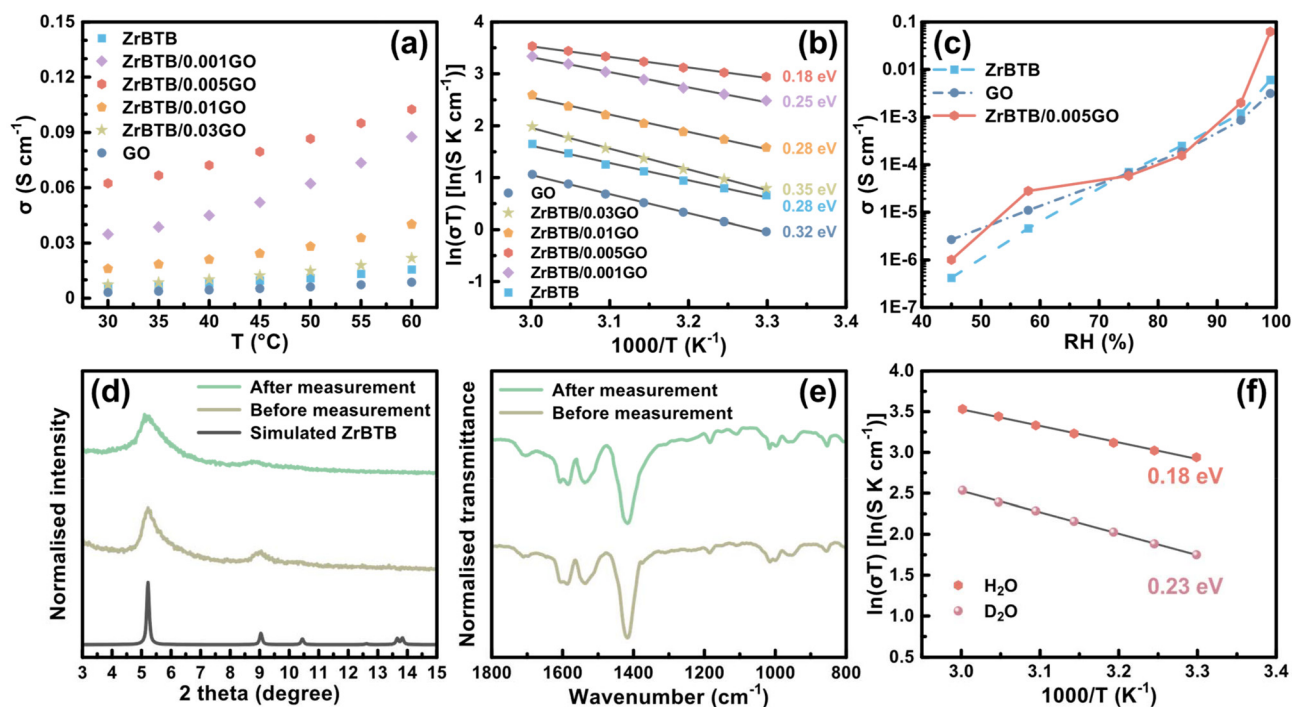


Fig. 6 (a) Proton conductivities of ZrBTB, GO and various physically blended nanocomposites, measured at different temperatures and 99% RH, and (b) the corresponding Arrhenius plots from (a). (c) Proton conductivities of ZrBTB, GO and ZrBTB/0.005GO, measured under various RH levels at 30 °C. (d) Grazing-incidence XRD patterns and (e) FTIR spectra of ZrBTB/0.005GO pellets, measured before and after the proton conductivity measurements under 99% RH at various temperatures from 30 °C to 60 °C. (f) Arrhenius plots of ZrBTB/0.005GO measured under the atmospheres of H₂O and D₂O vapours, respectively. Activation energies are listed in (b) and (f), and the simulated XRD pattern of ZrBTB is provided in (d) for comparison.



To clarify the effect of GO loading on the proton-conducting performance of the resulting material, a series of composites with varying GO loadings were prepared by physical blending for measuring their proton conductivities. Nyquist plots are shown in Fig. S11, and the results are summarised in Fig. 6(a and b). It can be found that all physically blended composites tested here exhibit higher values of σ compared to both pristine materials, and ZrBTB/0.005GO shows the optimal performance among all materials. An ultrahigh σ of $1.03 \times 10^{-1} \text{ S cm}^{-1}$ and a low E_a of 0.18 eV can be achieved at 60 °C and 99% RH, outperforming many of those highly proton-conductive MOF-based and GO-based materials reported previously (Table S2).

Since RH strongly influences the proton conduction in MOF-based materials, varying RH levels is an effective approach to investigate the underlying conducting mechanism. Proton conductivities of ZrBTB/0.005GO and both pristine materials were measured at 30 °C under 94, 84, 75, 58 and 45% RH, respectively, and the obtained Nyquist plots are shown in Fig. S12. As illustrated in Fig. 6(c), the σ values of both ZrBTB and GO increase gradually from around $10^{-6} \text{ S cm}^{-1}$ to $10^{-3} \text{ S cm}^{-1}$ when the RH rises from 45% to 99%. ZrBTB/0.005GO exhibits a similar trend in σ to both pristine materials in the region below 84% RH. However, its σ increases sharply and becomes outperforming compared to both pristine materials, only when the RH increases from 84% to 99%. This observation indicates that ZrBTB/0.005GO requires sufficient water uptake to achieve higher σ compared to both pristine materials, exhibiting ultrahigh conductivity under a RH of 94% or that approaching saturation. In addition, ZrBTB/0.005GO can preserve its crystallinity and coordinating characteristics after both the pelletisation and EIS measurements under 99% RH, as evidenced by both the grazing-incidence X-ray diffraction (GIXRD) patterns and FTIR spectra shown in Fig. 6(d) and (e).

To probe the kinetic isotope effect of water vapour and verify that protons are the species being transported within these materials, σ values at various temperatures were measured in the presence of D_2O vapour for ZrBTB/0.005GO and both pristine materials. The resulting Nyquist plots are shown in Fig. S13, and the corresponding Arrhenius plots are shown in Fig. 6(f) and Fig. S14. The σ values of ZrBTB and GO in the presence of D_2O vapour are reduced by factors of 1.6 and 1.4, respectively, compared to those measured under H_2O vapour over the entire temperature range. The results here are consistent with the theoretical expectation that σ values should be inversely proportional to the square root of the atomic mass of protonic carriers.²⁹ Notably, the difference in σ observed for ZrBTB/0.005GO under H_2O and D_2O is around 3-fold, substantially greater than those for both pristine materials. Such an enhanced sensitivity to the isotopic substitution indicates the presence of augmented proton-hopping transport strongly relying on the hydrogen-bonding network and hydrophilic functional groups,^{33,63} as observed in other proton-conductive materials as well.^{33,64,65}

Water uptake in ZrBTB, GO and nanocomposites

The results discussed previously indicate that 2D ZrBTB–GO nanocomposites prepared by chemical grafting and physical blending can achieve quite different proton-conducting characteristics, and the physically blended composite can exhibit significantly higher σ than both pristine materials when the RH in the testing environment is close to saturation. We thus hypothesised that such differences in proton-conducting behaviours should be attributed to the affinity between the materials and water molecules coming from the external environment. To test this hypothesis, ZrBTB, GO and three representative nanocomposites were exposed to an environment of 99% RH at 30 °C for 2 days, immediately followed by the TGA measurement of each sample in order to quantify their water uptake from the humid environment.^{29,34} The obtained TGA data are shown in Fig. 7(a). The weight loss of each sample before the plateau at 110 °C, a temperature slightly higher than the boiling point of water, indicates the

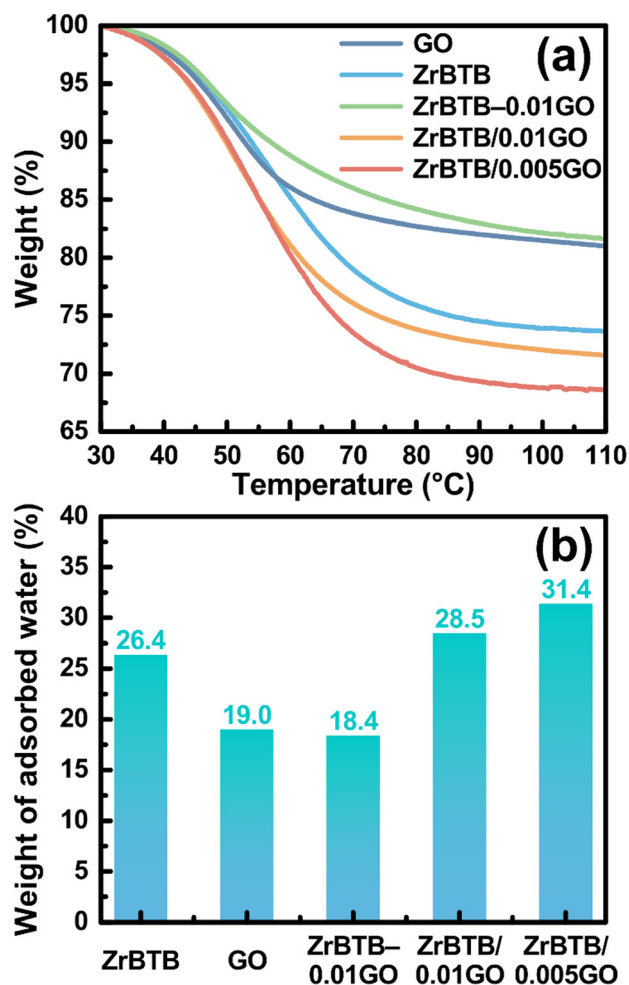


Fig. 7 (a) TGA curves of various materials, after 2 days of exposure to the environment under 99% RH at 30 °C. Ramping rate: 1° min^{-1} . (b) Mass fractions of adsorbed water in each material prior to the TGA measurement, extracted from the data shown in (a).



amount of physically adsorbed water in the material.²⁹ The resulting water uptakes are plotted in Fig. 7(b). It can be seen that the trend in TGA results is exactly consistent with that in the proton-conducting performance, with ZrBTB/0.005GO exhibiting the highest σ value and ZrBTB-0.01GO showing the lowest σ value. The E_a values, derived from the Arrhenius plots, also follow the same trend, with ZrBTB/0.005GO exhibiting the lowest E_a and ZrBTB-0.01GO exhibiting the highest E_a . The findings here indicate that for these materials, the amount of water in the hydrogen-bonding network governs the Grotthuss-type proton transport. With a small amount of GO, *i.e.*, around 1 wt% or 0.5 wt% here, to form composites with ZrBTB, it creates more voids between the flower-like aggregates of ZrBTB sheets and transforms the morphology of the material into loosely packed tiny 2D sheets (see Fig. 3); such morphologies should provide more locations to capture and accumulate water molecules. However, for ZrBTB-0.01GO, its chemical coordination between GO and the MOF reduces the amounts of accessible hydrophilic functional groups in both materials, leading to reduced water uptake. With the physically blended nanocomposites, more voids between aggregates of MOF sheets are present compared to the pristine ZrBTB, with hydrophilic groups from both ZrBTB and GO well preserved. They can thus adsorb more water molecules from the humid air under a high RH, which is the cause for their better proton-conducting performance compared to those of both the pristine ZrBTB and pristine GO.

Conclusions

ZrBTB and its nanocomposites with GO were successfully synthesised by performing either chemical grafting or physical blending, enabling a systematic investigation of their proton-conducting behaviours. Nanocomposites synthesised by both methods possess a similar morphology, with aggregates of loosely packed small 2D sheets. However, the coordination between GO and ZrBTB through carboxylate groups can be found in the chemically grafted nanocomposite, as evidenced by FTIR data. All MOF-based materials can preserve their crystallinity and coordination characteristics after the pelletising process and EIS measurements under 99% RH. The chemically grafted nanocomposite shows worse proton-conducting performance compared to the pristine materials. On the other hand, the physically blended composites can outperform both pristine materials as proton conductors, with ZrBTB/0.005GO achieving an ultrahigh σ of $1.03 \times 10^{-1} \text{ S cm}^{-1}$ at 60 °C under 99% RH and a relatively low E_a of 0.18 eV. This superior performance is attributed to the cooperative effect of hydrophilic functional groups from GO and the hydroxyl groups on the nodes of ZrBTB, which enhances the water uptake and promotes the formation of a continuous hydrogen-bonding network. As the first study reporting proton-conducting behaviours in composites composed of 2D MOFs and GO, the findings here suggest that for 2D MOF-based materials, physical blending is a simpler and more effective strategy to synthesise

nanocomposites than chemical grafting, aiming to develop high-performance proton-conductive materials.

Conflicts of interest

There are no conflicts to declare.

Data availability

All data supporting this article have been included in the manuscript or included as part of the supplementary information (SI). Supplementary information: additional experimental data. See DOI: <https://doi.org/10.1039/d5dt02918j>.

Acknowledgements

This work is financially supported by the National Science and Technology Council (NSTC) of Taiwan, under the following grants: 112-2223-E-006-003-MY3 and 113-2923-E-006-006-MY2. This study is also in part funded by the Ministry of Education (MOE) of Taiwan, under the Yushan Young Fellow Program (MOE-112-YSFEE-0005-001-P2) and the Higher Education Sprout Project under National Cheng Kung University (NCKU). We also thank the Core Facility Center of NCKU for conducting both TEM and TGA measurements.

References

- H. Furukawa, K. E. Cordova, M. O'Keeffe and O. M. Yaghi, *Science*, 2013, **341**, 1230444.
- S. Kitagawa, R. Kitaura and S. I. Noro, *Angew. Chem., Int. Ed.*, 2004, **43**, 2334–2375.
- L. J. Murray, M. Dincă and J. R. Long, *Chem. Soc. Rev.*, 2009, **38**, 1294–1314.
- J. L. Obeso, V. B. López-Cervantes, C. V. Flores, A. Martínez, Y. A. Amador-Sánchez, N. S. Portillo-Velez, H. A. Lara-García, C. Leyva, D. Solis-Ibarra and R. A. Peralta, *Dalton Trans.*, 2024, **53**, 4790–4796.
- J.-R. Li, J. Sculley and H.-C. Zhou, *Chem. Rev.*, 2012, **112**, 869–932.
- T. Chen and D. Zhao, *Coord. Chem. Rev.*, 2023, **491**, 215259.
- Y. Liu, A. J. Howarth, J. T. Hupp and O. K. Farha, *Angew. Chem., Int. Ed.*, 2015, **54**, 9001–9005.
- S. H. Park, H. M. Kim, M. L. Díaz-Ramírez, S. Lee and N. C. Jeong, *Chem. Commun.*, 2024, **60**, 14577–14580.
- L. E. Kreno, K. Leong, O. K. Farha, M. Allendorf, R. P. Van Duyne and J. T. Hupp, *Chem. Rev.*, 2012, **112**, 1105–1125.
- M. Ko, L. Mendecki, A. M. Eagleton, C. G. Durbin, R. M. Stolz, Z. Meng and K. A. Mirica, *J. Am. Chem. Soc.*, 2020, **142**, 11717–11733.
- D. Sheberla, J. C. Bachman, J. S. Elias, C.-J. Sun, Y. Shao-Horn and M. Dincă, *Nat. Mater.*, 2016, **16**, 220–224.



- 12 G. Lee, G. Park and S. S. Park, *J. Am. Chem. Soc.*, 2024, **146**, 29767–29772.
- 13 S. Bureekaew, S. Horike, M. Higuchi, M. Mizuno, T. Kawamura, D. Tanaka, N. Yanai and S. Kitagawa, *Nat. Mater.*, 2009, **8**, 831–836.
- 14 T. Yamada, K. Otsubo, R. Makiura and H. Kitagawa, *Chem. Soc. Rev.*, 2013, **42**, 6655–6669.
- 15 P. Ramaswamy, N. E. Wong and G. K. Shimizu, *Chem. Soc. Rev.*, 2014, **43**, 5913–5932.
- 16 W. J. Phang, H. Jo, W. R. Lee, J. H. Song, K. Yoo, B. Kim and C. S. Hong, *Angew. Chem., Int. Ed.*, 2015, **54**, 5142–5146.
- 17 K. I. Otake and H. Kitagawa, *Small*, 2021, **17**, 2006189.
- 18 K.-P. Chen, Y. Ma, H.-X. Ren, C.-X. Zhang and Q.-L. Wang, *Dalton Trans.*, 2024, **53**, 8716–8721.
- 19 K.-X. Zhao, G.-Q. Zhang, X.-R. Wu, H.-B. Luo, Z.-X. Han, Y. Liu and X.-M. Ren, *ACS Appl. Electron. Mater.*, 2025, **7**, 3164–3175.
- 20 S. Biswas, D. K. Thapa and L. Mandal, *Dalton Trans.*, 2025, **54**, 1750–1769.
- 21 K. Jiao, J. Xuan, Q. Du, Z. Bao, B. Xie, B. Wang, Y. Zhao, L. Fan, H. Wang and Z. Hou, *Nature*, 2021, **595**, 361–369.
- 22 Y. Wang, D. F. R. Diaz, K. S. Chen, Z. Wang and X. C. Adroher, *Mater. Today*, 2020, **32**, 178–203.
- 23 A. J. Howarth, Y. Liu, P. Li, Z. Li, T. C. Wang, J. T. Hupp and O. K. Farha, *Nat. Rev. Mater.*, 2016, **1**, 1–15.
- 24 S. Yuan, J.-S. Qin, C. T. Lollar and H.-C. Zhou, *ACS Cent. Sci.*, 2018, **4**, 440–450.
- 25 M.-D. Tsai, K.-C. Wu and C.-W. Kung, *Chem. Commun.*, 2024, **60**, 8360–8374.
- 26 C.-H. Wu, K.-C. Wu, C.-H. Shen and C.-W. Kung, *Coord. Chem. Rev.*, 2025, **538**, 216704.
- 27 X. Chen and G. Li, *Inorg. Chem. Front.*, 2020, **7**, 3765–3784.
- 28 G. K. Shimizu, J. M. Taylor and S. Kim, *Science*, 2013, **341**, 354–355.
- 29 W. H. Ho, S.-C. Li, Y.-C. Wang, T.-E. Chang, Y.-T. Chiang, Y.-P. Li and C.-W. Kung, *ACS Appl. Mater. Interfaces*, 2021, **13**, 55358–55366.
- 30 M. Szufla, A. Choroś, W. Nitek and D. Matoga, *Chem. – Eur. J.*, 2022, **28**, e202200835.
- 31 M. Szufla, J. A. R. Navarro, K. Góra-Marek and D. Matoga, *ACS Appl. Mater. Interfaces*, 2023, **15**, 28184–28192.
- 32 M. K. Sarango-Ramírez, M. Donoshita, Y. Yoshida, D. W. Lim and H. Kitagawa, *Angew. Chem., Int. Ed.*, 2023, **62**, e202301284.
- 33 Y. Zheng, Y. Liu, H. Li, Z. Yang, W. Wu, J. Zhang, J. Wu and J. Wang, *Adv. Funct. Mater.*, 2025, **35**, 2500151.
- 34 F. Yang, H. Huang, X. Wang, F. Li, Y. Gong, C. Zhong and J.-R. Li, *Cryst. Growth Des.*, 2015, **15**, 5827–5833.
- 35 S.-L. Zheng, C.-M. Wu, L.-H. Chung, H.-Q. Zhou, J. Hu, Z. Liu, Y. Wu, L. Yu and J. He, *ACS Energy Lett.*, 2023, **8**, 3095–3101.
- 36 L. Yang, B. Tang and P. Wu, *J. Mater. Chem. A*, 2015, **3**, 15838–15842.
- 37 Z. Rao, K. Feng, B. Tang and P. Wu, *J. Membr. Sci.*, 2017, **533**, 160–170.
- 38 P. M. Unnikrishnan, G. Premanand and S. K. Das, *Inorg. Chem.*, 2025, **64**, 3506–3517.
- 39 M. R. Karim, K. Hatakeyama, T. Matsui, H. Takehira, T. Taniguchi, M. Koinuma, Y. Matsumoto, T. Akutagawa, T. Nakamura, S.-I. Noro, T. Yamada, H. Kitagawa and S. Hayami, *J. Am. Chem. Soc.*, 2013, **135**, 8097–8100.
- 40 L. Glasser, *Chem. Rev.*, 1975, **75**, 21–65.
- 41 K. D. Kreuer, A. Rabenau and W. Weppner, *Angew. Chem., Int. Ed. Engl.*, 1982, **21**, 208–209.
- 42 I. S. Kim, J. Borycz, A. E. Platero-Prats, S. Tussupbayev, T. C. Wang, O. K. Farha, J. T. Hupp, L. Gagliardi, K. W. Chapman, C. J. Cramer and A. B. F. Martinson, *Chem. Mater.*, 2015, **27**, 4772–4778.
- 43 Z. Hu, E. M. Mahdi, Y. Peng, Y. Qian, B. Zhang, N. Yan, D. Yuan, J.-C. Tan and D. Zhao, *J. Mater. Chem. A*, 2017, **5**, 8954–8963.
- 44 L. Feng, Y. Qiu, Q.-H. Guo, Z. Chen, J. S. W. Seale, K. He, H. Wu, Y. Feng, O. K. Farha, R. D. Astumian and J. F. Stoddart, *Science*, 2021, **374**, 1215–1221.
- 45 Y.-L. Chen, C.-H. Shen, C.-W. Huang and C.-W. Kung, *Mol. Syst. Des. Eng.*, 2023, **8**, 330–340.
- 46 T.-C. Lin, K.-C. Wu, J.-W. Chang, Y.-L. Chen, M.-D. Tsai and C.-W. Kung, *Dalton Trans.*, 2024, **53**, 11426–11435.
- 47 K.-C. Wu, M.-D. Tsai, C.-H. Wu, T.-H. Yang, Y.-L. Chen and C.-W. Kung, *APL Mater.*, 2024, **12**, 031114.
- 48 J.-W. Chang, T.-C. Lin, Y.-L. Chen, P.-C. Han, S.-C. Yang, M.-D. Tsai, K. C. W. Wu and C.-W. Kung, *CrystEngComm*, 2024, **26**, 2852–2861.
- 49 L. Stobinski, B. Lesiak, A. Malolepszy, M. Mazurkiewicz, B. Mierzwa, J. Zemek, P. Jiricek and I. Bieloshapka, *J. Electron Spectrosc. Relat. Phenom.*, 2014, **195**, 145–154.
- 50 M.-D. Tsai, Y.-L. Chen, J.-W. Chang, S.-C. Yang and C.-W. Kung, *ACS Appl. Energy Mater.*, 2023, **6**, 11268–11277.
- 51 R. Shimoni, Z. Shi, S. Binyamin, Y. Yang, I. Liberman, R. Ifraemov, S. Mukhopadhyay, L. Zhang and I. Hod, *Angew. Chem., Int. Ed.*, 2022, **61**, e202206085.
- 52 Y. Li and R. T. Yang, *Langmuir*, 2007, **23**, 12937–12944.
- 53 S. Ullah, M. A. Bustam, M. A. Assiri, A. G. Al-Sehemi, G. Gonfa, A. Mukhtar, F. A. Abdul Kareem, M. Ayoub, S. Saqib and N. B. Mellon, *Microporous Mesoporous Mater.*, 2020, **294**, 109844.
- 54 S. Thakur and N. Karak, *Carbon*, 2012, **50**, 5331–5339.
- 55 G. Srinivas, J. W. Burrell, J. Ford and T. Yildirim, *J. Mater. Chem.*, 2011, **21**, 11323–11329.
- 56 V. Brusko, A. Khannanov, A. Rakhmatullin and A. M. Dimiev, *Carbon*, 2024, **229**, 119507.
- 57 R. Al-Gaashani, A. Najjar, Y. Zakaria, S. Mansour and M. A. Atieh, *Ceram. Int.*, 2019, **45**, 14439–14448.
- 58 L. Liu, Z. Qiao, X. Cui, C. Pang, H. Liang, P. Xie, X. Luo, Z. Huang, Y. Zhang and Z. Zhao, *ACS Appl. Mater. Interfaces*, 2019, **11**, 23039–23049.
- 59 C.-H. Shen, Y. Zhao, H. N. Nam, L. Zhu, Q. M. Phung, V. Austen, M. Kim, D. Jiang, X. Wei, T. Yokoshima, C.-W. Kung and Y. Yamauchi, *Chem. Sci.*, 2025, **16**, 7026–7038.



- 60 T. Shen, X. Feng, L. Fan, L. Li, Y. Zhang, B. Chai, U. Khan and X. Wang, *Cryst. Growth Des.*, 2025, **25**, 5312–5320.
- 61 J. Liu, Z. Li, X. Zhang, K.-I. Otake, L. Zhang, A. W. Peters, M. J. Young, N. M. Bedford, S. P. Letourneau, D. J. Mandia, J. W. Elam, O. K. Farha and J. T. Hupp, *ACS Catal.*, 2019, **9**, 3198–3207.
- 62 A. R. Heiba, M. O. Abdel-Salam, T. Yoon and E. El Sawy, *Nanoscale*, 2025, **17**, 459–473.
- 63 T. E. DeCoursey and V. V. Cherny, *J. Gen. Physiol.*, 1997, **109**, 415–434.
- 64 D. W. Kang, K. S. Lim, K. J. Lee, J. H. Lee, W. R. Lee, J. H. Song, K. H. Yeom, J. Y. Kim and C. S. Hong, *Angew. Chem., Int. Ed.*, 2016, **55**, 16123–16126.
- 65 F. Yang, G. Xu, Y. Dou, B. Wang, H. Zhang, H. Wu, W. Zhou, J.-R. Li and B. Chen, *Nat. Energy*, 2017, **2**, 877–883.

