



Cite this: *Phys. Chem. Chem. Phys.*, 2026, **28**, 3550

Lipid bilayer membranes with asymmetrically distributed LPC and DAG

Chang Liu,^{id}^{ab} Zhongjie Han,^{ab} Rui Ma^{id}^{*c} and Chen Song^{id}^{*ab}

The complex chemical and biophysical characteristics of biomembranes are influenced by the asymmetric distribution of specific lipids. *In vitro*, the introduction of lysophosphatidylcholine (LPC) into one leaflet of lipid bilayers is frequently utilized to regulate membrane protein activity. *In vivo*, the conversion of phosphatidylinositol 4,5-bisphosphate (PIP₂) to diacylglycerol (DAG) in one leaflet can also modulate membrane protein activities. However, the effects of such variations in lipid composition, lipid quantity, and particularly the asymmetry of lipid distribution on the properties and morphology of biomembranes remain to be fully elucidated. Through molecular dynamics simulations, we demonstrate that the asymmetric distribution of LPC and the asymmetric conversion of PIP₂ induce asymmetric alterations in membrane structures and lipid dynamics. Such alterations can generate an imbalance in the lateral pressure distribution between the two leaflets, potentially leading to membrane curvature. The extent of membrane curvature is also influenced by the length and degree of unsaturation of the lipid acyl tail chain. Our findings underscore the critical role of lipid asymmetry in shaping biomembrane structure and dynamics, providing new insights into the regulation of membrane proteins and cellular functions mediated by these specific lipids.

Received 25th August 2025,
 Accepted 12th January 2026

DOI: 10.1039/d5cp03245h

rsc.li/pccp

1 Introduction

Biological membranes are semipermeable barriers that compartmentalize cellular and organelle contents. They play a crucial role in facilitating various biological processes, such as regulating the transport of chemical substances, establishing and maintaining transmembrane solute gradients, and modulating intercellular recognition and adhesion.^{1,2} Biological membranes are primarily composed of lipids, membrane proteins, and carbohydrates. The composition of lipids in biological membranes can vary significantly and is vital to the characteristics of the membrane and its biological functions.^{3,4} The simplest geometries of biological membranes are planar and spherical. However, membranes can undergo morphological changes across a wide range of length scales, from nanometers to micrometers, and over diverse time scales.^{5–7} In addition to planar and spherical morphologies, biological membranes can also adopt other long-lived shapes, as observed in the endoplasmic reticulum (ER), the Golgi complex (GC), and the mitochondrial inner membrane (MIM).^{8–10} In contrast to

these persistent membrane architectures that define organelle and cellular structures, transient membrane deformations play essential roles in dynamic processes, including endocytosis, exocytosis, immune responses, cell motility, and division.^{11–13}

Membrane properties are influenced by lipid composition through various physical and chemical characteristics, including head group size, charge, acyl chain length, and degree of saturation.¹⁴ Among these characteristics, the relative size of the lipid head group and the hydrophobic tail significantly influences the shape of the membrane. For instance, phosphatidylcholine (PC), the most abundant zwitterionic phospholipid in eukaryotic cell membranes and subcellular organelles,¹⁵ possesses a cylindrical shape due to its two acyl chains and large choline head group, which promotes the formation of a planar bilayer. In contrast, lysophosphatidylcholine (LPC), a derivative of PC consisting of a single acyl chain, exhibits a conical molecular geometry that promotes positive membrane curvature and facilitates the formation of normal hexagonal or micellar phases.^{16,17} Conversely, diacylglycerol (DAG), which lacks a phosphate head group and thus presents a smaller polar region, adopts an inverted conical shape that promotes negative membrane curvature and favors the formation of inverse hexagonal phases.

These lipid characteristics determine membrane morphology and also regulate the functions of membrane proteins. For instance, LPC, when asymmetrically intercalated into one leaflet of a bilayer membrane, is known to facilitate the gating of

^a Center for Quantitative Biology, Academy for Advanced Interdisciplinary Studies, Peking University, Beijing 100871, China. E-mail: c.song@pku.edu.cn

^b Peking-Tsinghua Center for Life Sciences, Academy for Advanced Interdisciplinary Studies, Peking University, Beijing 100871, China

^c College of Physical Science and Technology, Xiamen University, Xiamen 361005, China. E-mail: ruima@xmu.edu.cn

mechanosensitive ion channels, such as the bacterial MscL and MscS channels,^{18–22} as well as the eukaryotic TRAAK and TREK-1 channels.²³ Similarly, phosphatidylinositol 4,5-bisphosphate (PIP₂), a crucial anchoring lipid in the inner leaflet, undergoes hydrolysis by phospholipase C to produce DAG and inositol-1,4,5-triphosphate (IP3), thereby regulating numerous signaling pathways.^{24–27} This conversion can activate multiple ion channels, including TRPC3, TRPC5, and TRPC6.²⁸ Therefore, alterations in lipid composition can influence not only membrane properties but also the activation of membrane proteins.

Molecular dynamics (MD) simulations are a valuable tool for examining biological membranes,^{29–31} as they provide a comprehensive atomistic understanding of biological systems when combined with experimental observations.³⁰ However, a detailed atomic-level investigation into the influence of asymmetrically distributed lipid molecules, particularly specific lipids extensively used in experiments such as LPC and DAG, on membrane properties has not been thoroughly explored. Therefore, in this study, we employed all-atom MD simulations to investigate how variations in lipid composition influence membrane characteristics.

We specifically investigate the influence of lipid headgroup and acyl chain structures on membrane properties by symmetrically incorporating 25% LPC of varying lengths into POPC lipid bilayers, or by symmetrically converting 25% PIP₂ with different saturation levels into DAG (Fig. 1). For the asymmetric

systems, the lipid compositional variations were confined to the lower leaflet. Furthermore, we conducted these lipid modifications exclusively in one leaflet of the bilayer to replicate experimental conditions for activating MS channels, allowing us to examine how these asymmetric lipid compositions affect membrane properties. Our results demonstrate that the asymmetrical insertion of LPC thinned the membrane, asymmetrically altered the order of the lipid acyl tail chains in the two leaflets, increased membrane fluidity, and induced curvature toward the LPC-inserted leaflet. Conversely, the asymmetrical conversion of PIP₂ to DAG thickened the membrane, maintained lipid order with only minor changes, increased fluidity, and caused the membrane to curve away from the DAG leaflet.

2 Results

2.1 Impact on the overall structure of lipid bilayers

Area per lipid (APL) and membrane thickness are commonly used metrics for assessing the structural properties of lipid bilayers during MD simulations. In conventional bilayer systems, membrane expansion is typically characterized by a single area-per-lipid (APL) value, obtained by dividing the total membrane area by the number of lipids. This measure implicitly assumes a uniform molecular architecture in which each lipid contributes one headgroup and two acyl chains. However, this assumption breaks down in our heterogeneous membranes containing LPC and DAG. LPC possesses only a single

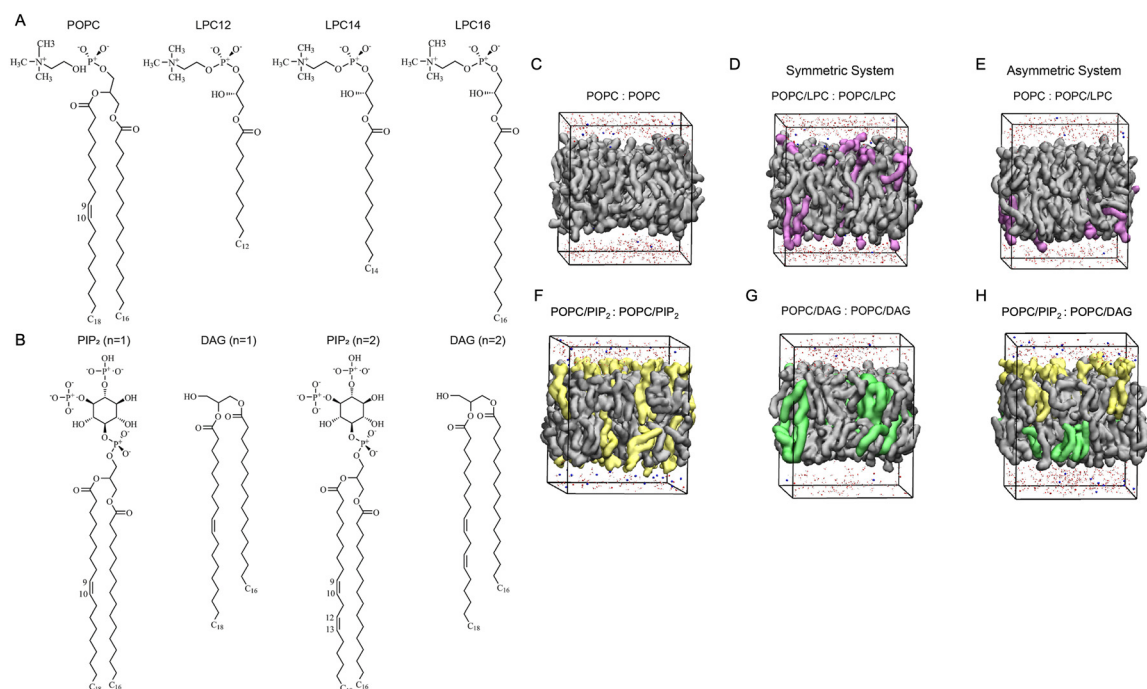


Fig. 1 Lipid molecules used in our molecular dynamics (MD) simulations and simulation systems of various lipid compositions. (A) Illustration of POPC and LPC lipids with varying acyl chain lengths. (B) Representation of PIP₂ or DAG with different degrees of unsaturation in their acyl chains, where n denotes the degree of unsaturation. (C) Simulation composed of pure POPC lipids. (D) and (E) Simulations composed of POPC and asymmetrical or symmetrical LPC. (F) and (G) Simulations composed of POPC and symmetrical PIP₂ or DAG. (H) Simulation composed of POPC, asymmetrical PIP₂, and DAG. POPC lipids are shown with grey surfaces, LPC lipids with pinkish-purple surfaces, PIP₂ lipids with yellow surfaces, and DAG lipids with green surfaces. For clarity, only part of the ions and water molecules are shown.

acyl chain due to its truncated tail, whereas DAG lacks a headgroup entirely. These structural deviations disrupt the canonical correspondence between headgroup count and tail-chain count, making a single APL metric insufficient for describing the lateral organization of such bilayers. To overcome this limitation, we separately calculated the area per lipid headgroup (APLH) and the area per lipid tail (APLT) (Fig. 2 and 3) and also explicitly calculated the total membrane surface area for each system, with the corresponding values summarized in Table S1. APLH is defined as the projected membrane area in the XY plane divided by the total number of lipid headgroups, while APLT is defined as the projected XY-plane area divided by the total number of acyl tails. This decomposition enables us to quantify membrane expansion in a manner that reflects the distinct structural contributions of headgroups and tails. Because DAG lacks phosphate-containing headgroups, it was excluded from the APLH calculation. Together, these metrics provide a more accurate description of how headless DAG lipids and single-tailed LPC lipids differentially modulate headgroup spacing and tail-chain packing within the bilayer.

In symmetrical bilayer systems, the insertion of LPC resulted in an approximately 5.4% decrease in APLH and a 5.1% increase in APLT compared to pure POPC lipid bilayers. Conversely, the conversion of PIP₂ to DAG led to an approximate 25% increase in APLH and a 5.9% decrease in APLT compared to the POPC/PIP₂ lipid bilayers (Fig. 2 and 3). In asymmetrical bilayer systems, the insertion of LPC or the conversion of PIP₂ to DAG in the lower leaflet resulted in varied effects on the upper and lower leaflets of the membrane. When LPC was inserted in the lower leaflet, the APLH of this leaflet decreased by approximately 12.2%. However, it caused an approximate 9.6% increase of APLH in the upper leaflet compared to pure POPC lipid bilayers (Fig. 2A). The APLT increased by around 10% in the upper leaflet and decreased by 2.5% in the lower leaflet (Fig. 3A). When PIP₂ was converted to DAG in the lower leaflet, the APLH of this leaflet increased by approximately 30%. However, this conversion led to an approximately 2.5% decrease of APLH in the upper leaflet (Fig. 2B). Since DAG lipids

share the same tail acyl chain as PIP₂ lipids, the conversion of PIP₂ to DAG had a minor effect on the average lipid tail in both symmetrical and asymmetrical bilayer systems (Fig. 3B). The asymmetric bilayer system exhibited a more pronounced impact on the upper leaflet compared to the corresponding symmetric bilayer systems, in the absence of lipid composition changes. The length of the LPC tail chain did not significantly affect the area per phosphorus headgroup and acyl tail chain, while the unsaturation of the PIP₂ or DAG tail chain had a minor effect (Fig. 2 and 3).

Variations in bilayer thickness were observed following alterations in lipid composition. As the insertion length of the LPC acyl chain tail increased, the reduction in thickness was less pronounced compared to that of pure POPC. Similarly, an increase in the unsaturation of PIP₂ converted to DAG resulted in a smaller increase in bilayer thickness. LPC insertion resulted in a decrease in bilayer thickness, while PIP₂ conversion led to an increase. These changes in bilayer thickness were more pronounced in symmetric bilayer systems compared to asymmetric bilayer systems (Fig. 4).

We note that when POPC is mixed with LPC at a 3 : 1 ratio, the difference in acyl-chain length between the two species may, in principle, promote partial demixing or domain formation through inter-leaflet coupling, as reported at higher LPC concentrations.³² However, within the timescale of our simulations, clustering of a few LPC molecules was only occasionally observed, and LPC lipids remained highly mobile and largely dispersed throughout the bilayer (Fig. S1A and B). Both visual inspection of the final configurations and the calculated LPC-LPC radial distribution functions (Fig. S1) showed no evidence of large-scale segregation, indicating that the POPC:LPC mixture remained essentially homogeneous over the course of the simulations.

2.2 Impact on the order parameter and fluidity of lipid molecules

The deuterium order parameter ($|S_{CD}|$) of lipid tails is a critical indicator of the orientation and ordering of the carbon chains.

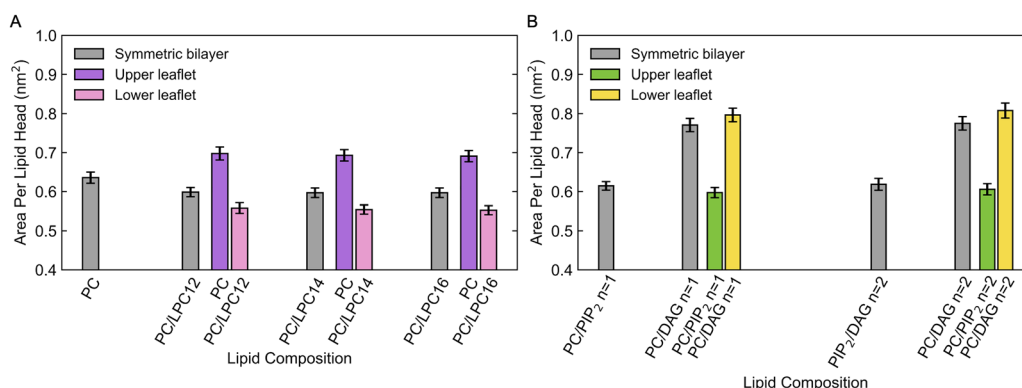


Fig. 2 Area per lipid head in lipid bilayer systems with LPC insertion (A) and PIP₂ conversion (B). Grey bars denote symmetric bilayers with identical lipid compositions in the upper and lower leaflets. For asymmetric bilayers, purple and green bars correspond to the upper leaflet, while pink and yellow bars correspond to the lower leaflet. Lipid compositions are indicated below the bars. Error bars represent the standard deviation obtained from five independent MD trajectories. The parameter n denotes the degree of unsaturation.

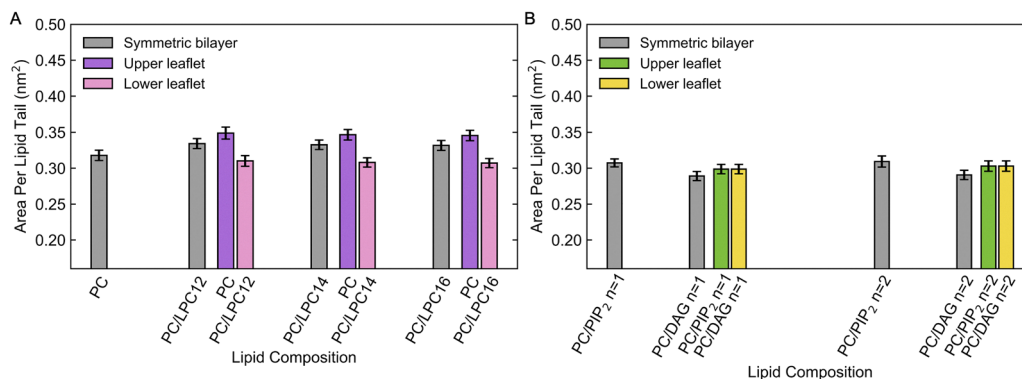


Fig. 3 Area per lipid tail in lipid bilayer systems with LPC insertion (A) and PIP₂ conversion (B). Grey bars denote symmetric bilayers with identical lipid compositions in the upper and lower leaflets. For asymmetric bilayers, purple and green bars correspond to the upper leaflet, while pink and yellow bars correspond to the lower leaflet. Lipid compositions are indicated below the bars. Error bars represent the standard deviation obtained from five independent MD trajectories. The parameter n denotes the degree of unsaturation.

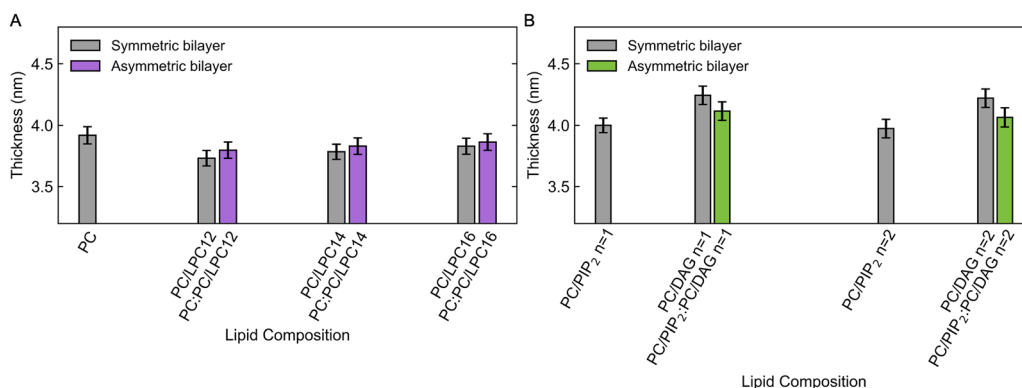


Fig. 4 Lipid bilayer thickness in systems with LPC insertion (A) and PIP₂ conversion (B). Grey bars denote symmetric bilayers with identical lipid compositions in the upper and lower leaflets, while purple and green bars denote asymmetric bilayers. Lipid compositions are indicated below the bars; for asymmetric bilayers, the compositions of upper and lower leaflets are separated by a colon. Error bars represent the standard deviation obtained from five independent MD trajectories. The parameter n denotes the degree of unsaturation.

The order parameters range from 1 to -0.5 , where a value of 1 indicates a fully ordered arrangement along the membrane normal, while lower values indicate a progressively disordered distribution that is perpendicular to the membrane normal.³³ To evaluate how lipid composition influences the structure of lipid bilayers, we analyzed $|S_{CD}|$ of the acyl chains (eqn (1) and Fig. 5, 6).

In symmetric bilayer systems, LPC insertion resulted in an increase in membrane disorder (Fig. 5A and B), while the conversion of PIP₂ to DAG led to a slight increase in lipid ordering (Fig. 6A and B). In asymmetric bilayer systems, LPC insertion exhibited distinct effects on the two lipid leaflets. In the lower leaflet where LPCs were inserted, the lipid acyl chains demonstrated a slight increase in ordering, while the upper layer exhibited a greater degree of disorder (Fig. 5C and D). These effects were less pronounced in bilayer systems involving PIP₂-DAG conversion (Fig. 6C and D). These observed changes were consistent with the results of the APLT analysis. As illustrated in Fig. 3, larger APLT values indicated an increased available space for the lipid acyl chains, which can accommodate

a more disordered tail. The order parameters of the POPC acyl chain were not significantly influenced by the length or degree of unsaturation of the lipid acyl chain (Fig. 5 and 6).

The lateral diffusion coefficient of lipids is commonly used to quantify the mobility of lipid molecules within the membrane plane. This parameter is crucial to elucidate the dynamic properties of biological membranes, which play a key role in the regulation of various biological processes, such as lipid phase transitions and clustering of lipids and proteins. In our atomistic MD simulations, we determined the lateral diffusion coefficients of lipid molecules by analyzing the mean square displacement of lipid phosphorus atoms (Fig. S2). The diffusion coefficient was calculated by fitting a straight line to the data from 100 to 150 ns, as described in eqn (2), for both symmetric simulation systems and asymmetrical LPC insertion systems. For the PIP₂-DAG conversion systems, a straight line was fitted to the data from 150 ns to 200 ns. The correction for periodic boundary conditions when calculating lipid diffusion coefficients^{34,35} was not applied, under the assumption that relative diffusion coefficients are largely insensitive to this

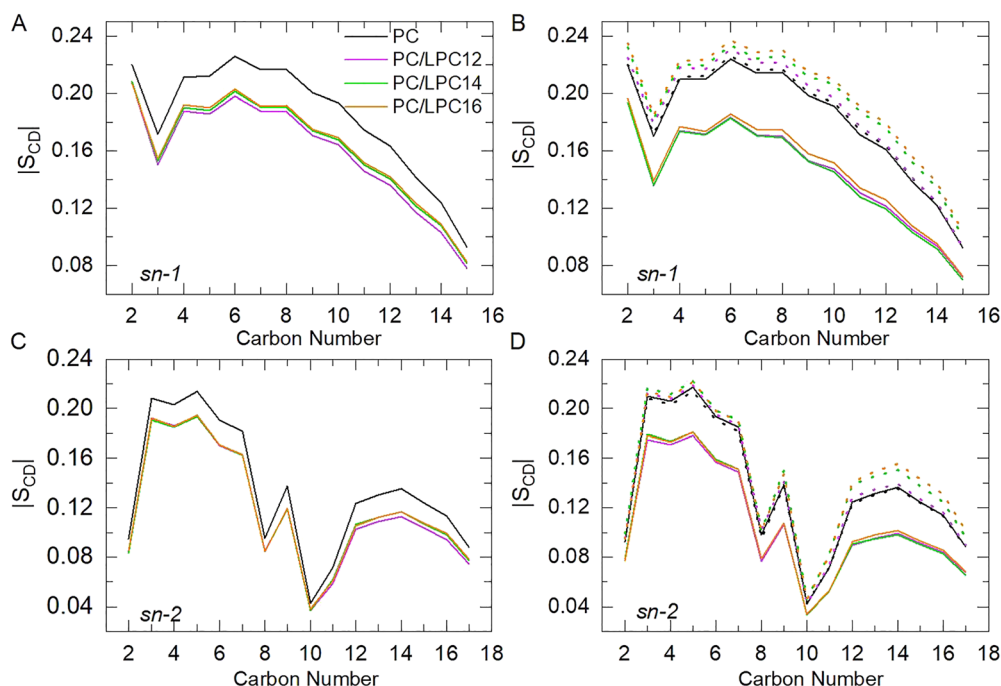


Fig. 5 Acyl chain order parameters of POPC lipids in LPC insertion bilayer systems. The order parameters of the two acyl chains of POPC, designated as sn-1 and sn-2, were analyzed in both symmetric (A) and (B) and asymmetric (C) and (D) systems. The solid line represents the order parameters for the upper leaflet, while the dotted plot denotes the order parameters for the lower leaflet (C) and (D). n denotes the degree of unsaturation.

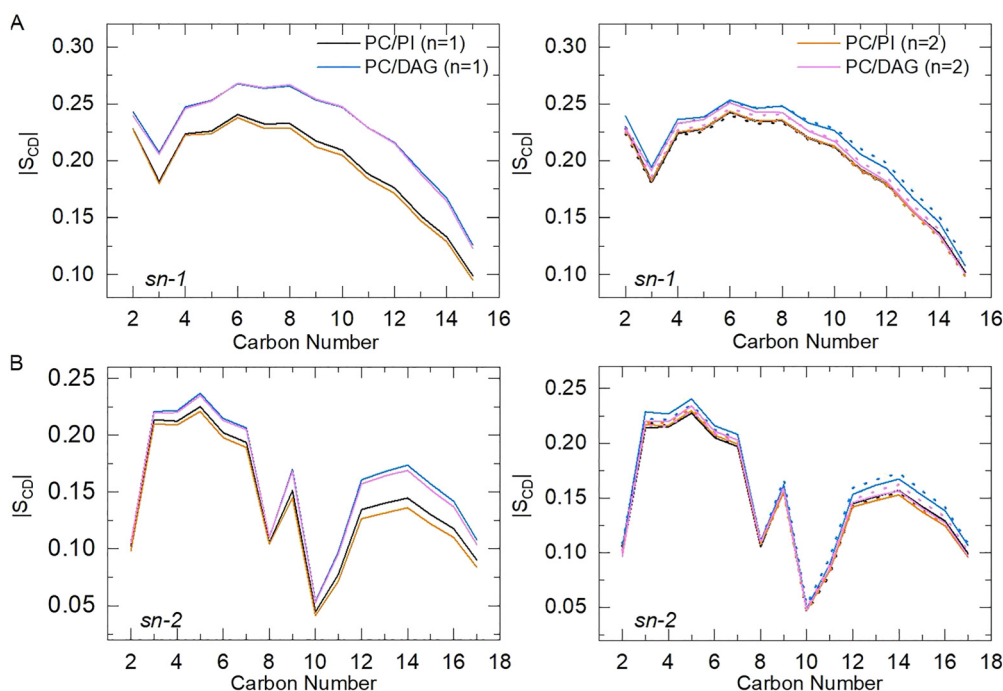


Fig. 6 Acyl chain order parameters of POPC lipids in PIP₂ conversion bilayer systems. The order parameters of the two acyl chains of POPC, designated as sn-1 and sn-2, were analyzed in symmetric (A) and (B) and asymmetric (C) and (D) bilayer systems. The solid line represents the order parameters for the upper leaflet, while the dotted plot denotes the order parameters for the lower leaflet (C) and (D).

correction. In addition to POPC diffusion, we also computed the lateral diffusion coefficients of LPC, PIP₂, and DAG molecules in both symmetric and asymmetric bilayers (Fig. S3).

Across all lipid types, diffusion is consistently slower in asymmetric bilayers than in their symmetric counterparts. Furthermore, LPC insertion increased the fluidity of the POPC lipid

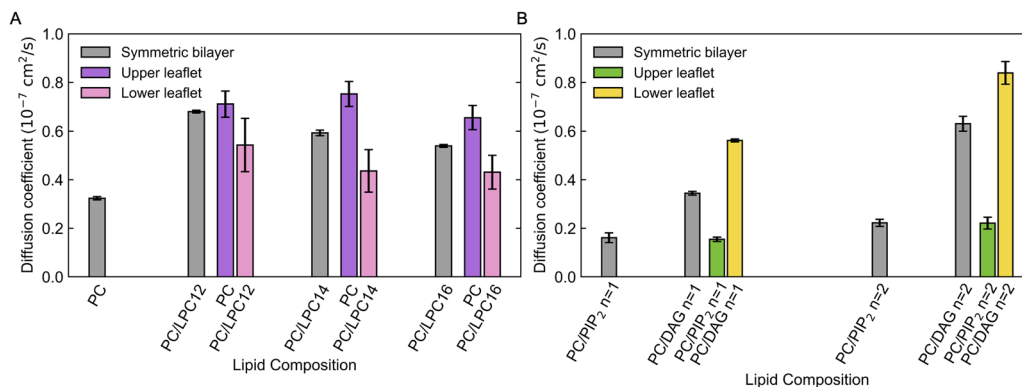


Fig. 7 Lateral diffusion coefficients of POPC lipid molecules in symmetric or asymmetric LPC insertion bilayer systems (A) and lateral diffusion coefficients of lipid molecules in symmetric or asymmetric PIP₂ conversion bilayer systems (B). In each panel, sets of histograms illustrate the results for three distinct systems: the symmetric system on the left, the upper leaflet system in the middle, and the lower leaflet-altered system on the right. In the symmetric system, only one histogram is displayed due to the equal composition in both leaflets. *n* denotes the degree of unsaturation.

bilayer in both symmetric and asymmetric bilayer systems. The diffusion coefficient of the LPC-inserted leaflet decreased with increasing length of the LPC acyl chain in both symmetric and asymmetric bilayer systems. However, no significant differences were observed in the upper leaflets of the asymmetric bilayer systems (Fig. 7A). The conversion of PIP₂ to DAG in the symmetric and asymmetric bilayer systems increased the fluidity of the POPC lipid bilayers. In both types of systems, the diffusion coefficient of the leaflet undergoing PIP₂ conversion increased with the unsaturation of the acyl chain. However, compared to the leaflet where PIP₂ is converted to DAG, the fluidity of the upper leaflet in the lipid bilayers of asymmetric conversion did not show significant changes (Fig. 7B).

2.3 Asymmetrical alteration induced asymmetric pressure profiles across lipid bilayers

To further elucidate the effect of changes in lipid composition on membrane properties, we calculated the lateral pressure profiles (LPPs) along the membrane normal. LPPs depict the spatial distribution of lateral stresses within a lipid bilayer, encompassing repulsive pressure components from lipid–lipid interactions and cohesive hydrophobic tension that promotes segregation of lipid chains from the aqueous environment.

These components exhibit an uneven distribution across the bilayer.³⁶ The pressures across the monolayer must be balanced for a lipid monolayer to remain flat. If the lateral pressure of the chain and the head group region is unbalanced, the shape of the monolayer will change. In a symmetric bilayer where the composition and quantity of lipids are identical in both monolayers, they will exhibit the same LPPs and counteract each other. Therefore, we focus on the analysis of LPPs in asymmetric simulation systems (Fig. 8 and 9). In addition, we computed the density profiles of the phosphorus and nitrogen atoms of the lipids to identify the spatial regions corresponding to different parts of the phospholipid molecules.

As can be seen in Fig. 8 and 9, a positive lateral pressure (I) was observed in the headgroup region, primarily due to steric, hydrational, and electrostatic effects. These interactions were overall repulsive, but might incorporate attractive contributions, such as hydrogen bonding interactions. In the vicinity of the glycerol backbone region (II), located just above/below the lipid headgroups, an attractive force arose because of the unfavorable contact of the hydrocarbon chains with water, known as the hydrophobic effect. Within the hydrocarbon interior of the membrane (III), attractive van der Waals interactions between the chains were counteracted by repulsive

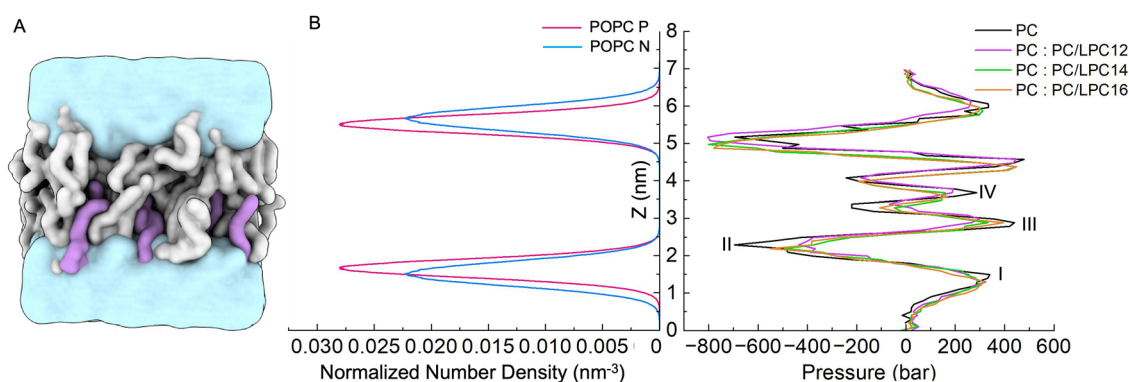


Fig. 8 Simulation system of PC: PC/LPC lipid bilayer (A) and lateral pressure profiles together with the density distributions of P and N atoms of POPC in asymmetric LPC insertion systems with varying acyl chain lengths (B).

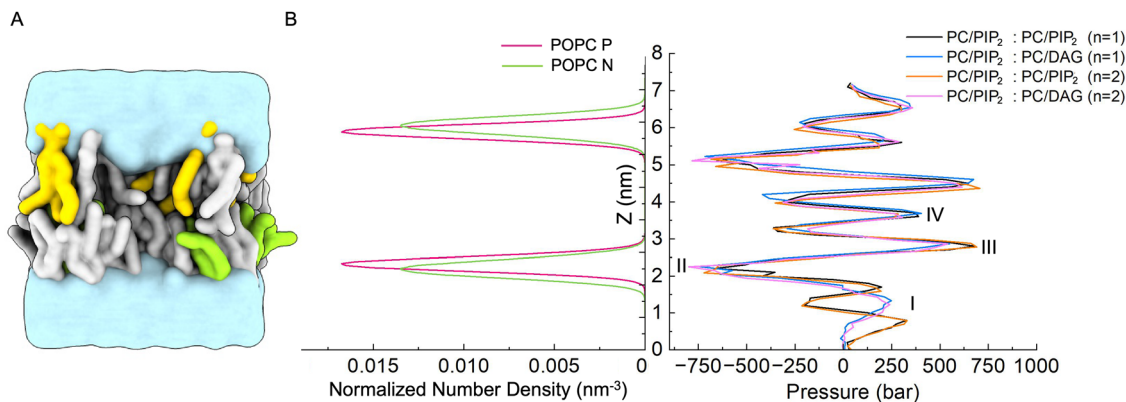


Fig. 9 Simulation system of PC/PIP₂: PC/DAG lipid bilayer (A) and lateral pressure profiles together with the density distributions of the P and N atoms of POPC in asymmetric PIP₂ conversion to DAG systems with varying degrees of acyl chain unsaturation (B). *n* denotes the degree of unsaturation.

interactions resulting from the thermal motions of the chains, resulting in a net positive lateral pressure that tends to expand the membrane. Additionally, the lateral pressure profile exhibited a positive peak at the bilayer center (IV), which was related to the geometries of specific lipid types.^{37,38}

In systems with LPC insertion, compared to pure POPC membranes, the attraction in region (II) and the repulsion in region (III) were reduced (Fig. 8). In simulation systems where PIP₂ is converted into DAG (Fig. 9), the repulsion between the headgroups of POPC and PIP₂ generated two peaks in region (I) (Fig. 9, black and orange lines). DAG lacks a headgroup, thus exhibiting only one peak resulting from the repulsion between the headgroups of POPC lipids in region (I) (Fig. 9, blue and pink lines). Furthermore, this conversion also reduced the repulsion in the lipid tail double-bond region (III), as evidenced by the weaker peak intensities of the blue and pink curves compared to the black and orange ones in Fig. 9.

To analyze the differential stress between the leaflets, we integrated the lateral pressure $\pi(z)$ along the membrane normal *z* direction from the bulk water to the bilayer center (Fig. 8B and 9B). The differential stress of the two leaflets was then obtained and denoted as $\Delta\sigma$ ($\langle\sigma_{\text{lower}}\rangle - \langle\sigma_{\text{upper}}\rangle$). The mean of $\Delta\sigma$, derived from the analysis of five independently repeated simulation trajectories, was approximately 0 mN m⁻¹ in pure POPC and POPC: POPC/PIP₂ (*n* = 1, 2) lipid bilayers (Table 1). This finding suggests that membranes composed of these lipid constituents tend to form planar structures. Interestingly, the lipid composition asymmetry of the POPC and PIP₂ in our simulation system did not lead to a significant differential stress between the two leaflets.

In contrast, the insertion of LPC and the conversion of PIP₂ into DAG induced differential stress in the two membrane leaflets in opposite directions. For LPC insertion, the $\Delta\sigma$ fluctuated within the range of -9.97 mN m⁻¹ to -12.83 mN m⁻¹, whereas the conversion of PIP₂ into DAG resulted in the $\Delta\sigma$ fluctuating within the range of 5.56 mN m⁻¹ to 7.32 mN m⁻¹ (Table 1). This phenomenon can be attributed to LPC having a relatively larger lipid head group and a single-chain acyl chain (Fig. 1A), which leads to increased crowding on the inserted side. Conversely, the conversion of PIP₂, which has a larger head group, into DAG, which has an almost negligible head group (Fig. 1B), results in a less crowded lipid head group region in the lower leaflet. Moreover, as the length of the LPC acyl chain inserted into the lipid increased and the unsaturation of the acyl chain converted from PIP₂ to DAG increased, the differential stress became more pronounced (Table 1). These results indicate that both the acyl chain length and the degree of unsaturation of the lipid tails significantly affect the LPPs of the lipid bilayers.

2.4 Asymmetric pressure profiles induced membrane curvature

The distribution of lateral pressures and tensions across lipid monolayers and bilayers is essential for determining their spontaneous curvature. The insertion of LPC and the conversion of PIP₂ to DAG introduce asymmetrical stress within the bilayer, potentially leading to curvature. This tendency for the membrane to bend can be characterized by the spontaneous curvature c_0 .

We estimated the spontaneous curvature induced by lipid changes *via* eqn (5) based on the integral of the first moment of

Table 1 Differential stress between two leaflets and corresponding membrane curvature. The error bars represent the standard deviation of $\Delta\sigma$ or c_0 obtained from five independent MD trajectories

Upper leaflet	POPC	POPC	POPC	POPC	POPC/PIP ₂	POPC	POPC/PIP ₂	POPC
Lower leaflet	POPC	POPC/LPC12	POPC/LPC14	POPC/LPC16	POPC/PIP ₂ (<i>n</i> = 1)	POPC/DAG (<i>n</i> = 1)	POPC/PIP ₂ (<i>n</i> = 2)	POPC/DAG (<i>n</i> = 2)
$\Delta\sigma$ (mN m ⁻¹)	0.46 ± 0.60	-9.97 ± 0.76	-11.25 ± 1.08	-12.83 ± 0.71	0.84 ± 0.90	5.56 ± 1.01	-0.10 ± 1.2	7.32 ± 1.74
c_0^+ (nm ⁻¹)	0.02 ± 0.01	0.04 ± 0.01	0.03 ± 0.01	0.07 ± 0.02	0.00 ± 0.01	-0.02 ± 0.01	0.00 ± 0.01	-0.03 ± 0.02
c_0^- (nm ⁻¹)	0.00 ± 0.01	0.44 ± 0.02	0.49 ± 0.02	0.53 ± 0.01	0.00 ± 0.01	-0.20 ± 0.02	0.00 ± 0.01	-0.33 ± 0.01
C_0 (nm ⁻¹)	0.01 ± 0.01	-0.20 ± 0.01	-0.23 ± 0.01	-0.23 ± 0.01	0.00 ± 0.01	0.09 ± 0.02	0.00 ± 0.01	0.15 ± 0.02

LPPs obtained from MD simulation analyzes. In addition to the effective curvature of the entire bilayer, the leaflet-level spontaneous curvatures were also calculated to examine the individual contributions from each monolayer. The direction of membrane curvature was defined from the perspective of an observer located below the lower leaflet of the bilayer, looking towards the upper leaflet. A concave curvature of the membrane (bending away from the lower leaflet) was defined as positive, while a convex curvature (bending toward the lower leaflet) was defined as negative. As shown in Table 1, the spontaneous curvature of POPC and POPC: POPC/PIP₂ was approximately 0 nm⁻¹. The curvature trends under different lipid compositions are illustrated in Fig. S4. The insertion of LPC and conversion of PIP₂ to DAG resulted in membrane curvature in opposite directions. Specifically, LPC insertion into the lower leaflet induced negative membrane curvature, ranging from -0.20 nm⁻¹ to -0.23 nm⁻¹, while the conversion of PIP₂ to DAG in the lower leaflet induced positive membrane curvature, ranging from 0.09 nm⁻¹ to 0.15 nm⁻¹. Furthermore, as the LPC tail chain lengthened and the unsaturation of DAG increased, the absolute value of the spontaneous curvature induced by these changes became more pronounced.

3 Discussion

The composition of the lipid bilayer, in which membrane proteins are embedded, has been demonstrated to play a crucial role in regulating protein function. The lipid bilayer's composition influences membrane protein function not only through specific biochemical interactions and lipid binding, such as the electrostatic attraction between polybasic motifs in proteins and negatively charged lipid headgroups (*e.g.*, PIP₂),³⁹ but also through the physical properties of membranes themselves.^{40–42} Our results indicate that changes in lipid composition, caused by LPC insertion or PIP₂-DAG conversion, can significantly impact membrane properties, including APL, thickness, fluidity, lateral pressure, and potentially curvature. These alterations in membrane properties can affect membrane morphology and regulate the activities of membrane-embedded or associated proteins, such as mechanosensitive ion channels.

The LPC insertion and PIP₂-DAG conversion exhibit distinct effects on APLH, indicating that the packing within the lipid headgroup region is influenced by both the structure of the lipid fatty acyl chain and the organization of the lipid headgroup itself. The hydrophobic thickness of the lipid bilayer is anticipated to closely match that of any protein embedded within it. This alignment is crucial due to the significant energetic penalty associated with exposing either fatty acyl chains or hydrophobic amino acids to an aqueous environment. Any disparity in hydrophobic thickness between the lipid bilayer and the protein is likely to result in distortion of either the lipid bilayer, the protein, or both, in order to minimize this mismatch. In cases of extreme hydrophobic mismatch, a membrane protein may be excluded from the lipid bilayer, as observed in simple model transmembrane α -helices.⁴³

Moreover, extreme mismatch could also induce the formation of non-bilayer phases by the lipids, particularly at low molar ratios of lipid to protein.^{43–45}

The insertion of LPC or the conversion of PIP₂ to DAG influences membrane thickness. Understanding the impact of changes in lipid composition on membrane thickness could yield valuable insights for experimental studies. For instance, alterations in membrane thickness can regulate the gating of mechanosensitive channels such as MscL,¹⁸ MscS,⁴⁶ and TREK-2,⁴⁷ all of which are known to be sensitive to membrane thickness. Additionally, MscS and PIEZO channels are affected by the degrees of lipid unsaturation,^{46,48,49} likely due to changes in membrane thickness or fluidity. Remarkably, similar lipid remodeling occurs physiologically during T-cell activation, where DAG generation coupled with PUFA enrichment creates a combinatorial lipid signature that potentially coordinates PIEZO-mediated calcium signaling in immune responses.⁵⁰ This highlights the critical role of membrane lipid composition in modulating immune cell function and mechanotransduction pathways.

The transition from the liquid crystalline phase to the gel and fluid phases induces a pronounced shift in the physical properties of a lipid bilayer, which can significantly influence the activities of membrane proteins, such as the clustering of these proteins. The symmetric or asymmetric insertion of LPC into the membrane lipid layer has distinct effects on the order parameters of the two leaflets of the lipid bilayers. Nonetheless, both insertion modes of LPCs contribute to an increase in membrane fluidity. The conversion of PIP₂ to DAG did not significantly affect the membrane order parameter but had a substantial impact on membrane fluidity. Alterations in these lipid components manifest effects on membrane order parameters and dynamics, offering a method for studying the conformation of membrane proteins, for example, Ca²⁺, Na⁺, and K⁺-ATPases.^{45,51,52}

Lipid lateral pressure profiles play a crucial role in the conformational equilibrium of proteins, constituting a vital aspect of protein function.⁵³ Mechanosensitive channels, known for their responsiveness to membrane forces and curvature,⁵⁴ can be activated by membrane stretch, lateral pressure, and local membrane curvature, as well as by the incorporation of amphipathic lipids.⁵⁵ This activation could be closely linked to the intricate relationship between mechanosensitive channels and membrane lipids. Amphipathic compounds, such as chlorpromazine, local anesthetics, and LPC could induce channel activation through insertion^{18,19,56} or the conversion of PIP₂ into DAG in one monolayer of the lipid bilayers.²⁸ Whether this activation was triggered by local membrane curvature or changes in the transbilayer pressure profile without local curvature remained unknown.^{57–59}

Membrane tension and bending not only drive gating transitions in mechanosensitive channels but also impact the organization of membrane proteins.⁶⁰ Certain proteins exhibit a propensity to partition into membrane regions with specific curvature, while others demonstrate flexibility, adapting to a range of curvature states.^{61,62} The lipid structures of LPC and

DAG are not inclined to form planar bilayer structures; instead, they induce membrane curvature. This property can be leveraged in *in vitro* experiments to regulate the localization of specific proteins. For instance, the spatial distribution of PIEZO1 channels, influenced by membrane curvature, avoids localization at highly curved membrane protrusions such as filopodia and instead is enriched in membrane invaginations.^{63–65}

In summary, this study examines the effects of symmetric and asymmetric LPC insertion, as well as the conversion of PIP₂ to DAG, on the structure, dynamic properties, and morphology of membranes. The findings contribute to a more comprehensive understanding of the roles of LPC and DAG in regulating membrane structures and the functions of membrane proteins. This work not only enriches our insight into related physiological processes but also serves as a valuable reference for future investigations.

4 Methods

4.1 Simulation systems

The molecular dynamics simulation systems consisted of symmetric or asymmetric mixtures of POPC/LPC, POPC/PIP₂, or POPC/DAG lipid monolayers, as well as pure POPC lipids. The LPC lipids insertion system consisted of 144 lipids and approximately 6417 water molecules. The PIP₂ conversion system was composed of 128 lipids and approximately 5297 water molecules (Table 2). Na⁺ and Cl⁻ ions were added to the system to achieve charge neutrality and a physiological ion concentration of 0.15 M.

4.2 Molecular dynamics simulations

All MD simulations were performed using GROMACS 2018.6.⁶⁶ The simulations utilized the CHARMM36 force field and the CHARMM implementation of the TIP3 water model.^{67–69} The temperature was maintained at a constant temperature of 300 K using the Nose–Hoover thermostat.^{70,71} The pressure was maintained at 1 bar using the Parinello–Rahman semiisotropic barostat^{72,73} with a time constant of 1 ps and a compressibility of $4.5 \times 10^{-5} \text{ bar}^{-1}$. A cutoff of 1.2 nm was used for the Lennard-Jones interaction, with a smooth shift function from

1.0 to 1.2 nm. The particle mesh Ewald (PME) method was used for long-range electrostatics.⁷⁴ After energy minimization, we ran 200-ps NVT and 500-ps NPT simulations to preliminarily equilibrate the systems, and then we carried out 500-ns NPT production simulations for each lipid simulation system (Table 2).

4.3 Analysis of MD trajectories

If not otherwise stated, the last 400 ns trajectories were used for the analysis. The order parameters of the lipid tails were analyzed using the following definition:

$$|S_{\text{CD}}| = \frac{\langle 3 \cos^2 \theta_z - 1 \rangle}{2} \quad (1)$$

where θ_z is the angle between the C–H bond and the membrane normal (z direction).

The lateral diffusion coefficient of the lipids was calculated with the following equation:

$$\langle r^2(t) \rangle = \langle (r(t) - r(0))^2 \rangle = 4Dt \quad (2)$$

where $\langle r^2(t) \rangle$ is the two-dimensional mean squared displacement (MSD) of phosphorus atoms, D is the diffusion coefficient, and t is the simulation time. A time interval of 150–200 ns or 100–150 ns was used to perform a linear fit and obtain the value of D in PIP₂ conversion systems and other systems. The error estimate was performed by calculating the difference in the diffusion coefficients obtained from the two halves of the fit interval, as implemented in the GROMACS-msd tool.

The lateral pressure profiles (LPPs) were calculated using the GROMACS-LS tool.⁷⁵ The calculation of LPPs was based on the last 400 ns of the simulation data (Table 2), with the coordinates and velocities saved every 20 ps. An electrostatic cutoff of 2.0 nm was used for the LPP calculation.⁷⁵ The profiles were calculated with about 80 slabs, corresponding to an approximate slab width of 1 Å. The outputs of GROMACS-LS were the components of the stress tensor, σ , as a function of the z coordinate. LPP was calculated using the formula

$$\pi(z) = P_{\text{L}}(z) - P_{\text{N}} \quad (3)$$

Table 2 List of the simulation systems

	Upper leaflet: lower leaflet	Lipid amounts	Duration (ns)	Replicates
Symmetric bilayer composition	POPC/LPC12: POPC/LPC12	48/16: 48/16	500	5
	POPC/LPC14: POPC/LPC14	48/16: 48/16		
	POPC/LPC16: POPC/LPC16	48/16: 48/16		
	POPC/PIP ₂ : POPC/PIP ₂ ($n = 1$)	48/16: 48/16		
	POPC/DAG: POPC/DAG ($n = 1$)	48/16: 48/16		
	POPC/PIP ₂ : POPC/PIP ₂ ($n = 2$)	48/16: 48/16		
	POPC/DAG: POPC/DAG ($n = 2$)	48/16: 48/16		
Asymmetric bilayer composition	POPC: POPC/LPC12	64: 48/16		
	POPC: POPC/LPC14	64: 48/16		
	POPC: POPC/LPC16	64: 48/16		
	POPC/PIP ₂ : POPC/DAG ($n = 1$)	48/16: 48/16		
	POPC/PIP ₂ : POPC/DAG ($n = 2$)	48/16: 48/16		

where $P_L(z)$ is the lateral component of the pressure tensor ($P_L(z) = 0.5[P_{XX}(z) + P_{YY}(z)]$), and P_N is the normal component ($P_N = P_{ZZ}$).

The spontaneous curvature estimation was calculated using continuum mechanics theory. When two lipid monolayers were combined to form a bilayer, we used the midplane between the two monolayers to represent the shape of the bilayer, whose position was denoted as $z = 0$ along the normal direction of the bilayer. Assuming that the thickness of the two monolayers was d^+ and d^- , we calculated the spontaneous curvature of each monolayer via the integral^{76,77}

$$\int_0^{d^+} z\pi(z)dz = \kappa^+ c_0^+ \quad \int_0^{-d^-} z\pi(z)dz = \kappa^- c_0^- \quad (1)$$

where κ^\pm and c_0^\pm denoted the bending rigidity and spontaneous curvature of the two monolayers, respectively, and $\pi(z)$ was the LLP defined in (3). In practice, we performed integration from the midplane of the bilayer to bulk water. The spontaneous curvature of the bilayer then reads as

$$\kappa c_0 = \int_{-d^-}^{d^+} z\pi(z)dz = \kappa^+ c_0^+ - \kappa^- c_0^- \quad (4)$$

For both the bilayer and the monolayer, the first moment of the LLP $\pi(z)$ gave the product of the bending rigidity and the spontaneous curvature. In order to obtain the spontaneous curvature, the values for the bending rigidity κ and κ^\pm were needed which in principle could be obtained by a buckling simulation.⁷⁸ Here, for simplicity, we assumed $\kappa^\pm = 10k_B T$ and $\kappa = \kappa^+ + \kappa^- = 20k_B T$, which were typical values for the plasma membrane.⁷⁹ We neglected the effect that the bilayer bending rigidity κ could be different for the two monolayers due to compositional asymmetry. This may affect the quantitative accuracy of the calculated spontaneous curvature, but not the sign of it (bending direction of the membranes).

Conflicts of interest

The authors declare no conflicts of interest.

Data availability

The simulation data supporting this article are available from the corresponding author upon reasonable request.

Supplementary information (SI) is available. See DOI: <https://doi.org/10.1039/d5cp03245h>.

Acknowledgements

This work was supported by the Innovative Research Groups of the National Natural Science Foundation of China (T2321001 to C.S.) and the Postdoctoral Fellowship of the Peking-Tsinghua Center for Life Sciences (Z.H.). C.S. was supported in part by the Frontier Innovation Fund of Peking University Chengdu Academy for Advanced Interdisciplinary Biotechnologies. Part of the

computation was performed on the computing platform of the Center for Life Sciences at Peking University.

References

- H. Watson, Biological membranes, *Essays Biochem.*, 2015, **59**, 43–69.
- A. Mangiarotti and R. Dimova, Biomolecular condensates in contact with membranes, *Annu. Rev. Biophys.*, 2024, **53**, 319–341.
- T. Harayama and H. Riezman, Understanding the diversity of membrane lipid composition, *Nat. Rev. Mol. Cell Biol.*, 2018, **19**, 281–296.
- O. Ali and A. Szabó, Review of eukaryote cellular membrane lipid composition, with special attention to the fatty acids, *Int. J. Mol. Sci.*, 2023, **24**, 15693.
- J. Zimmerberg and M. M. Kozlov, How proteins produce cellular membrane curvature, *Nat. Rev. Mol. Cell Biol.*, 2006, **7**, 9–19.
- M. M. Kozlov, F. Campelo, N. Liska, L. V. Chernomordik, S. J. Marrink and H. T. McMahon, Mechanisms shaping cell membranes, *Curr. Opin. Cell Biol.*, 2014, **29**, 53–60.
- P. Bassereau, R. Jin, T. Baumgart, M. Deserno, R. Dimova, V. A. Frolov, P. V. Bashkirov, H. Grubmüller, R. Jahn, H. J. Risselada and L. Johannes, The 2018 biomembrane curvature and remodeling roadmap, *J. Phys. D: Appl. Phys.*, 2018, **51**, 343001.
- J. Nixon-Abell, C. J. Obara, A. V. Weigel, D. Li, W. R. Legant, C. S. Xu, H. A. Pasolli, K. Harvey, H. F. Hess, E. Betzig and C. Blackstone, Increased spatiotemporal resolution reveals highly dynamic dense tubular matrices in the peripheral ER, *Science*, 2016, **354**, aaf3928.
- M. S. Ladinsky, D. N. Mastronarde, J. R. McIntosh, K. E. Howell and L. A. Staehelin, Golgi structure in three dimensions: functional insights from the normal rat kidney cell, *J. Cell Biol.*, 1999, **144**, 1135–1149.
- T. G. Frey and C. A. Mannella, The internal structure of mitochondria, *Trends Biochem. Sci.*, 2000, **25**, 319–324.
- R. Sougrat, A. Bartesaghi, J. D. Lifson, A. E. Bennett, J. W. Bess, D. J. Zabransky and S. Subramaniam, Electron tomography of the contact between T cells and SIV/HIV-1: implications for viral entry, *PLoS Pathog.*, 2007, **3**, e63.
- H. T. McMahon and J. L. Gallop, Membrane curvature and mechanisms of dynamic cell membrane remodelling, *Nature*, 2005, **438**, 590–596.
- M. Kaksonen and A. Roux, Mechanisms of clathrin-mediated endocytosis, *Nat. Rev. Mol. Cell Biol.*, 2018, **19**, 313–326.
- D. Casares, P. V. Escribá and C. A. Rosselló, Membrane lipid composition: effect on membrane and organelle structure, function and compartmentalization and therapeutic avenues, *Int. J. Mol. Sci.*, 2019, **20**, 2167.
- T. A. Lagace and N. D. Ridgway, The role of phospholipids in the biological activity and structure of the endoplasmic

- reticulum, *Biochim. Biophys. Acta, Mol. Cell Res.*, 2013, **1833**, 2499–2510.
- 16 A. R. Battle, P. Ridone, N. Bavi, Y. Nakayama, Y. A. Nikolaev and B. Martinac, Lipid-protein interactions: Lessons learned from stress, *Biochim. Biophys. Acta, Biomembr.*, 2015, **1848**, 1744–1756.
 - 17 M. A. Zhukovsky, A. Filograna, A. Luini, D. Corda and C. Valente, Phosphatidic acid in membrane rearrangements, *FEBS Lett.*, 2019, **593**, 2428–2451.
 - 18 E. Perozo, A. Kloda, D. M. Cortes and B. Martinac, Physical principles underlying the transduction of bilayer deformation forces during mechanosensitive channel gating, *Nat. Struct. Biol.*, 2002, **9**, 696–703.
 - 19 E. Perozo, D. M. Cortes, P. Sompornpisut, A. Kloda and B. Martinac, Open channel structure of MscL and the gating mechanism of mechanosensitive channels, *Nature*, 2002, **418**, 942–948.
 - 20 B. Corry and B. Martinac, Bacterial mechanosensitive channels: experiment and theory, *Biochim. Biophys. Acta, Biomembr.*, 2008, **1778**, 1859–1870.
 - 21 T. Nomura, C. G. Cranfield, E. Deplazes, D. M. Owen, A. Macmillan, A. R. Battle, M. Constantine, M. Sokabe and B. Martinac, Differential effects of lipids and lyso-lipids on the mechanosensitivity of the mechanosensitive channels MscL and MscS, *Proc. Natl. Acad. Sci. U. S. A.*, 2012, **109**, 8770–8775.
 - 22 R. Strutt, J. W. Hindley, J. Gregg, P. J. Booth, J. D. Harling, R. V. Law, M. S. Friddin and O. Ces, Activating mechanosensitive channels embedded in droplet interface bilayers using membrane asymmetry, *Chem. Sci.*, 2021, **12**, 2138–2145.
 - 23 F. Maingret, A. J. Patel, F. Lesage, M. Lazdunski and E. Honoré, Lysophospholipids open the two-pore domain mechano-gated K⁺ channels TREK-1 and TRAAK, *J. Biol. Chem.*, 2000, **275**, 10128–10133.
 - 24 M. J. Berridge, Inositol trisphosphate and calcium signaling, *Nature*, 1993, **361**, 315–325.
 - 25 B.-C. Suh and B. Hille, PIP₂ is a necessary cofactor for ion channel function: how and why?, *Annu. Rev. Biophys.*, 2008, **37**, 175–195.
 - 26 L. Cocco, M. Y. Follo, L. Manzoli and P.-G. Suh, Phosphoinositide-specific phospholipase C in health and disease, *J. Lipid Res.*, 2015, **56**, 1853–1860.
 - 27 J. H. C. Luong, I. H. Cheong, X. F. Zhang, Z. Kozlakidis and H. Wang, Structure and Roles of Phospholipase C (PLC), Phosphatidylinositol 4, 5-bisphosphate (PIP₂), and Inositol 1, 4, 5-trisphosphate (IP₃) in Metabolism and Disease: A Systematic Review, *Innov. Digit. Health Diagn. Biomark.*, 2025, **5**, 1–13.
 - 28 R. Gutorov, B. Katz, E. Rhodes-Mordov, R. Zaguri, T. Brandwine-Shemmer and B. Minke, The role of membrane lipids in light-activation of drosophila TRP channels, *Biomolecules*, 2022, **12**, 382.
 - 29 J. Loschwitz, O. O. Olubiyi, J. S. Hub, B. Strodel and C. S. Poojari, Computer simulations of protein-membrane systems, *Prog. Mol. Biol. Transl. Sci.*, 2020, **170**, 273–403.
 - 30 M. L. Berkowitz, Detailed molecular dynamics simulations of model biological membranes containing cholesterol, *Biochim. Biophys. Acta, Biomembr.*, 2009, **1788**, 86–96.
 - 31 C. Liu, Q. Zhong, K. Kang, R. Ma and C. Song, Asymmetrical calcium ions induced stress and remodeling in lipid bilayer membranes, *Phys. Chem. Chem. Phys.*, 2025, **27**, 740–753.
 - 32 G. Enkavi, M. Javanainen, W. Kulig, T. Róg and I. Vattulainen, Multiscale simulations of biological membranes: the challenge to understand biological phenomena in a living substance, *Chem. Rev.*, 2019, **119**, 5607–5774.
 - 33 T. J. Piggot, J. R. Allison, R. B. Sessions and J. W. Essex, On the calculation of acyl chain order parameters from lipid simulations, *J. Chem. Theory Comput.*, 2017, **13**, 5683–5696.
 - 34 B. A. Camley, M. G. Lerner, R. W. Pastor and F. L. Brown, Strong influence of periodic boundary conditions on lateral diffusion in lipid bilayer membranes, *J. Chem. Phys.*, 2015, **143**, 243113.
 - 35 R. M. Venable, H. I. Ingólfsson, M. G. Lerner, B. S. Perrin Jr, B. A. Camley, S. J. Marrink, F. L. Brown and R. W. Pastor, Lipid and Peptide Diffusion in Bilayers: The Saffman-Delbrück Model and Periodic Boundary Conditions, *J. Phys. Chem. B*, 2017, **121**, 3443–3457.
 - 36 R. S. Cantor, Lipid composition and the lateral pressure profile in bilayers, *Biophys. J.*, 1999, **76**, 2625–2639.
 - 37 K. J. Boyd, N. N. Alder and E. R. May, Molecular dynamics analysis of cardiolipin and monolysocardiolipin on bilayer properties, *Biophys. J.*, 2018, **114**, 2116–2127.
 - 38 C.-P. Chng, Y. Sadosky, K. J. Hsia and C. Huang, Curvature-regulated lipid membrane softening of nano-vesicles, *Extreme Mech. Lett.*, 2021, **43**, 101174.
 - 39 W. D. Heo, T. Inoue, W. S. Park, M. L. Kim, B. O. Park, T. J. Wandless and T. Meyer, PI (3, 4, 5) P₃ and PI (4, 5) P₂ lipids target proteins with polybasic clusters to the plasma membrane, *Science*, 2006, **314**, 1458–1461.
 - 40 J. A. Lundbæk, S. A. Collingwood, H. I. Ingólfsson, R. Kapoor and O. S. Andersen, Lipid bilayer regulation of membrane protein function: gramicidin channels as molecular force probes, *J. R. Soc., Interface*, 2010, **7**, 373–395.
 - 41 M. Ashrafuzzaman, R. E. Koeppe and O. S. Andersen, Intrinsic lipid curvature and bilayer elasticity as regulators of channel function: a comparative single-molecule study, *Int. J. Mol. Sci.*, 2024, **25**, 2758.
 - 42 J. Bigay and B. Antonny, Curvature, lipid packing, and electrostatics of membrane organelles: defining cellular territories in determining specificity, *Dev. Cell*, 2012, **23**, 886–895.
 - 43 R. J. Webb, J. M. East, R. P. Sharma and A. G. Lee, Hydrophobic mismatch and the incorporation of peptides into lipid bilayers: a possible mechanism for retention in the Golgi, *Biochemistry*, 1998, **37**, 673–679.
 - 44 J. C. Holthuis and A. K. Menon, Lipid landscapes and pipelines in membrane homeostasis, *Nature*, 2014, **510**, 48–57.
 - 45 A. Lee, Lipid-protein interactions in biological membranes: a structural perspective, *Biochim. Biophys. Acta, Biomembr.*, 2003, **1612**, 1–40.
 - 46 F. Xue, C. D. Cox, N. Bavi, P. R. Rohde, Y. Nakayama and B. Martinac, Membrane stiffness is one of the key determinants of E. coli MscS channel mechanosensitivity, *Biochim. Biophys. Acta, Biomembr.*, 2020, **1862**, 183203.

- 47 P. Aryal, V. Jarerattanachai, M. V. Clausen, M. Schewe, C. McClenaghan, L. Argent, L. J. Conrad, Y. Y. Dong, A. C. Pike, E. P. Carpenter and T. Baukowitz, Bilayer-mediated structural transitions control mechanosensitivity of the TREK-2 K2P channel, *Structure*, 2017, **25**, 708–718.
- 48 P. Ridone, S. L. Grage, A. Patkunarajah, A. R. Battle, A. S. Ulrich and B. Martinac, “Force-from-lipids” gating of mechanosensitive channels modulated by PUFAs, *J. Mech. Behav. Biomed. Mater.*, 2018, **79**, 158–167.
- 49 L. O. Romero, A. E. Massey, A. D. Mata-Daboyn, F. J. Sierra-Valdez, S. C. Chauhan, J. F. Cordero-Morales and V. Vásquez, Dietary fatty acids fine-tune Piezo1 mechanical response, *Nat. Commun.*, 2019, **10**, 1200.
- 50 F. Sasaki, M. Hayashi, Y. Mouri, S. Nakamura, T. Adachi and T. Nakashima, Mechanotransduction via the Piezo1-Akt pathway underlies Sost suppression in osteocytes, *Biochem. Biophys. Res. Commun.*, 2020, **521**, 806–813.
- 51 A. Starling, J. East and A. Lee, Effects of gel phase phospholipid on the Ca²⁺ + -ATPase, *Biochemistry*, 1995, **34**, 3084–3091.
- 52 F. Cornelius, Modulation of Na, K-ATPase and Na-ATPase activity by phospholipids and cholesterol. I. Steady-state kinetics, *Biochemistry*, 2001, **40**, 8842–8851.
- 53 R. S. Cantor, Lateral pressures in cell membranes: a mechanism for modulation of protein function, *J. Phys. Chem. B*, 1997, **101**, 1723–1725.
- 54 J. Gullingsrud and K. Schulten, Lipid bilayer pressure profiles and mechanosensitive channel gating, *Biophys. J.*, 2004, **86**, 3496–3509.
- 55 N. Bavi, C. D. Cox, Y. A. Nikolaev and B. Martinac, Molecular insights into the force-from-lipids gating of mechanosensitive channels, *Curr. Opin. Physiol.*, 2023, 100706.
- 56 C. Kung; Y. Saimi and B. Martinac, *Current Topics in Membranes and Transport*, Elsevier, 1990, vol. 36, pp. 145–153.
- 57 H. Hong and J. U. Bowie, Dramatic destabilization of transmembrane helix interactions by features of natural membrane environments, *J. Am. Chem. Soc.*, 2011, **133**, 11389–11398.
- 58 J. T. Marinko, H. Huang, W. D. Penn, J. A. Capra, J. P. Schleich and C. R. Sanders, Folding and misfolding of human membrane proteins in health and disease: from single molecules to cellular proteostasis, *Chem. Rev.*, 2019, **119**, 5537–5606.
- 59 N. Bavi, C. D. Cox, Y. A. Nikolaev and B. Martinac, Molecular insights into the force-from-lipids gating of mechanosensitive channels, *Curr. Opin. Physiol.*, 2023, **36**, 100706.
- 60 I. Levental and E. Lyman, Regulation of membrane protein structure and function by their lipid nano-environment, *Nat. Rev. Mol. Cell Biol.*, 2023, **24**, 107–122.
- 61 A. Tian and T. Baumgart, Sorting of lipids and proteins in membrane curvature gradients, *Biophys. J.*, 2009, **96**, 2676–2688.
- 62 F.-C. Tsai, M. Simunovic, B. Sorre, A. Bertin, J. Manzi, A. Callan-Jones and P. Bassereau, Comparing physical mechanisms for membrane curvature-driven sorting of BAR-domain proteins, *Soft Matter*, 2021, **17**, 4254–4265.
- 63 C. A. Haselwandter and R. MacKinnon, Piezos membrane footprint and its contribution to mechanosensitivity, *eLife*, 2018, **7**, e41968.
- 64 Y.-C. Lin, Y. R. Guo, A. Miyagi, J. Levring, R. MacKinnon and S. Scheuring, Force-induced conformational changes in PIEZO1, *Nature*, 2019, **573**, 230–234.
- 65 G. Vaisey, P. Banerjee, A. J. North, C. A. Haselwandter and R. MacKinnon, Piezo1 as a force-through-membrane sensor in red blood cells, *eLife*, 2022, **11**, e82621.
- 66 M. J. Abraham, T. Murtola, R. Schulz, S. Páll, J. C. Smith, B. Hess and E. Lindahl, GROMACS: High performance molecular simulations through multi-level parallelism from laptops to supercomputers, *SoftwareX*, 2015, **1**, 19–25.
- 67 J. B. Klauda, R. M. Venable, J. A. Freites, J. W. O’Connor, D. J. Tobias, C. Mondragon-Ramirez, I. Vorobyov, A. D. MacKerell Jr and R. W. Pastor, Update of the CHARMM all-atom additive force field for lipids: validation on six lipid types, *J. Phys. Chem. B*, 2010, **114**, 7830–7843.
- 68 A. D. MacKerell Jr, D. Bashford, M. L. Bellott, R. L. Dunbrack Jr, J. D. Evanseck, M. J. Field, S. Fischer, J. Gao, H. Guo, S. Ha and D. Joseph-McCarthy, All-atom empirical potential for molecular modeling and dynamics studies of proteins, *J. Phys. Chem. B*, 1998, **102**, 3586–3616.
- 69 S. R. Durell, B. R. Brooks and A. Ben-Naim, Solvent-induced forces between two hydrophilic groups, *J. Phys. Chem.*, 1994, **98**, 2198–2202.
- 70 S. Nosé, A molecular dynamics method for simulations in the canonical ensemble, *Mol. Phys.*, 1984, **52**, 255–268.
- 71 W. G. Hoover, Canonical dynamics: Equilibrium phase-space distributions, *Phys. Rev. A:At., Mol., Opt. Phys.*, 1985, **31**, 1695.
- 72 M. Parrinello and A. Rahman, Polymorphic transitions in single crystals: A new molecular dynamics method, *J. Appl. Phys.*, 1981, **52**, 7182–7190.
- 73 S. Nosé and M. Klein, Constant pressure molecular dynamics for molecular systems, *Mol. Phys.*, 1983, **50**, 1055–1076.
- 74 T. Darden, D. York and L. Pedersen, Particle mesh Ewald: An N log(N) method for Ewald sums in large systems, *J. Chem. Phys.*, 1993, **98**, 10089–10092.
- 75 J. M. Vanegas, A. Torres-Sánchez and M. Arroyo, Importance of force decomposition for local stress calculations in biomembrane molecular simulations, *J. Chem. Theory Comput.*, 2014, **10**, 691–702.
- 76 B. Róycki and R. Lipowsky, Spontaneous curvature of bilayer membranes from molecular simulations: Asymmetric lipid densities and asymmetric adsorption, *J. Chem. Phys.*, 2015, **142**, 02B601.
- 77 A. Hossein and M. Deserno, Spontaneous curvature, differential stress, and bending modulus of asymmetric lipid membranes, *Biophys. J.*, 2020, **118**, 624–642.
- 78 M. Hu, P. Diggins and M. Deserno, Determining the bending modulus of a lipid membrane by simulating buckling, *J. Chem. Phys.*, 2013, **138**, 214110.
- 79 M. Deserno, Fluid lipid membranes: From differential geometry to curvature stresses, *Chem. Phys. Lipids*, 2015, **185**, 11–45.