



Interface-induced fast Li⁺ transport in mixed ionic–electronic conductors

 Chenchen He,^{id}^a Zhuoxun Zhang^{id}^a and Tao Li^{id}^{*ab}

Cite this: DOI: 10.1039/d6cc00741d

 Received 4th February 2026,
Accepted 12th May 2026

DOI: 10.1039/d6cc00741d

rsc.li/chemcomm

Interfacial instability between lithium metal and solid-state electrolytes limits the performance of all-solid-state lithium metal batteries (ASSLBs), leading to parasitic reactions, non-uniform Li⁺ flux, and dendrite growth. Here, we develop a composite interlayer composed of the anti-perovskite Li₂OHCl_{0.75}Br_{0.25} (AP) and carbon nanotubes (CNTs) to enhance both interfacial stability and ionic transport. The AP–CNT interlayer exhibits enhanced Li⁺ conductivity arising from interfacial electron transfer from AP to CNTs, which generates a built-in electric field that facilitates Li⁺ migration. Lithium symmetric cells incorporating this interlayer achieve a high critical current density of 2.4 mA cm⁻² at 55 °C. This design integrates chemical robustness with coupled ion–electron transport, offering a generalizable strategy for safe, dendrite-free, and high-performance ASSLBs.

All-solid-state lithium metal batteries have emerged as promising candidates for next-generation electric vehicle energy storage. Their promise lies in solid-state electrolytes (SSEs), which offer high mechanical strength, single-ion (Li⁺) conduction, and intrinsic non-flammability, enabling enhanced safety and the prospect of ultra-high energy density cycling.^{1,2} However, their practical implementation is severely hindered by unstable interfaces between SSEs and lithium metal anodes.³ Owing to the strong reducing nature of lithium metal, spontaneous interfacial reactions occur upon contact with SSEs, leading to the formation of a chemically and structurally heterogeneous interphase. In sulfide-based SSEs, such reduction produces ionically resistive decomposition products (*e.g.*, Li₂S-, LiF-, and Li₃P-rich phases),^{4,5} which disrupt continuous Li⁺ transport, induce non-uniform ion flux and local overpotential accumulation, and ultimately promote dendrite formation and cell failure.⁶ To address these challenges, artificial interlayers—widely explored in metal electrodes and electrocatalysts^{7–9}—have been introduced to chemically decouple the Li/SSE interface, preserve Li⁺

transport continuity, and homogenize current distribution, thereby enhancing interfacial stability in ASSLBs.^{10,11}

Pure ionic conductors are widely employed as artificial interlayers owing to their favourable chemical stability against Li metal; however, their low room-temperature Li⁺ conductivity and high interfacial energy often result in severe interfacial polarization and inhomogeneous ion flux.^{12–14} Anti-perovskite (AP) solid electrolytes exemplify this limitation. Shen *et al.* demonstrated that an AP interlayer at the Li/Li₃InCl₆ interface significantly improved chemical compatibility and current distribution, increasing the critical current density from 2.4 to 4.2 mA cm⁻² and enabling stable Li plating/stripping for over 2400 h.¹⁵ Nevertheless, the intrinsically low ionic conductivity of AP (~10⁻⁵ S cm⁻¹ at room temperature) restricts Li⁺ transport, typically requiring elevated temperatures (>80 °C),¹⁶ which compromises the practical applicability and safety of SSBs.

In this context, mixed ionic–electronic conductors (MIECs) formed by incorporating conductive carbon into ionic conductors offer a viable strategy to lower nucleation overpotential, alleviate interfacial polarization, and reduce lithium deposition behaviour.¹⁷ However, most reported MIEC systems remain limited by insufficient ionic conductivity, resulting in aggravated polarization, and diminished control over lithium deposition during prolonged cycling.^{18,19} Therefore, preserving the intrinsic chemical stability of AP while enhancing Li⁺ transport through rational composite design remains a key challenge for achieving uniform lithium deposition and long-term interfacial stability.²⁰

Here, we design and construct a composite interlayer composed of Li₂OHCl_{0.75}Br_{0.25} and carbon nanotubes to regulate coupled ionic and electronic transport at the lithium metal/sulfide solid electrolyte interface, thereby enhancing interfacial stability and electrochemical kinetics in all-solid-state lithium metal batteries. By systematically tuning the AP/CNT mass ratio (1–50 wt% CNT), we identify a non-linear dependence of ionic conductivity on CNT content, with a maximum at 5 wt% CNT, which preserves chemical stability while markedly enhancing Li⁺ transport. Unlike previously

^a Department of Chemistry, Virginia Tech, Blacksburg, VA, 24061, USA.

E-mail: tli25@vt.edu

^b X-ray Science Division, Argonne National Laboratory, Lemont, Illinois 60439, USA


reported MIEC interlayers that rely on enhanced electronic conductivity to reduce polarization,^{17,21} our strategy exploits the work-function mismatch between AP and CNT to induce interfacial electron transfer until Fermi-level equilibration is achieved, resulting in charge redistribution and the establishment of a built-in electric field at the interface. This field converts boundary-limited transport in pristine AP into faster interfacial Li⁺ migration (Fig. 1), representing a fundamentally different ion–electron coupling mechanism. As a result, the ionic conductivity of the composite interlayer is enhanced more than tenfold relative to pristine AP, leading to substantially improved interfacial kinetics. Benefiting from these synergistic effects, Li symmetric cells incorporating the composite interlayer achieve a high critical current density of 2.4 mA cm⁻² at 55 °C. This work demonstrates a general design strategy for stabilizing lithium metal interfaces through synergistic ion–electron regulation.

The Li₂OHCl_{0.75}Br_{0.25} SSE was synthesized following the procedure reported in the supplementary information (SI). An AP-*x*% CNT composite (*x* = 1, 3, 5, 10, 50) was prepared by ball milling AP and CNTs at a specific weight ratio. High-resolution X-ray diffraction (XRD) patterns of pure AP and its mixture with CNTs are shown in Fig. S1 (SI). No new diffraction peaks were observed, except for an additional peak at 7.88° corresponding to CNTs, suggesting that no structural change occurs upon mixing AP with CNTs. The electronic conductivity of AP–CNT composites was measured at 55 °C using the chronoamperometry (CA) method. The electronic conductivity increases monotonically with the CNT content. Once the CNT content exceeds 3 wt%, the composite transitions into an electronic conductor (Fig. S2, SI). Specifically, AP–5% CNT shows an electronic conductivity of 5.4 × 10⁻² S cm⁻¹ (Fig. S3, SI), which is more than eight orders of magnitude higher than that of pristine AP (1.27 × 10⁻¹⁰ S cm⁻¹). Owing to their distinct electronic conductivity behaviours, the ionic conductivity of AP-*x*% CNT composites with low CNT content (*x* = 0, 1, 3) was evaluated using EIS, while samples with higher CNT content (*x* = 5, 10, 50) were characterized by the chronoamperometry method (Fig. 2a and b). Unlike electronic conductivity, the ionic

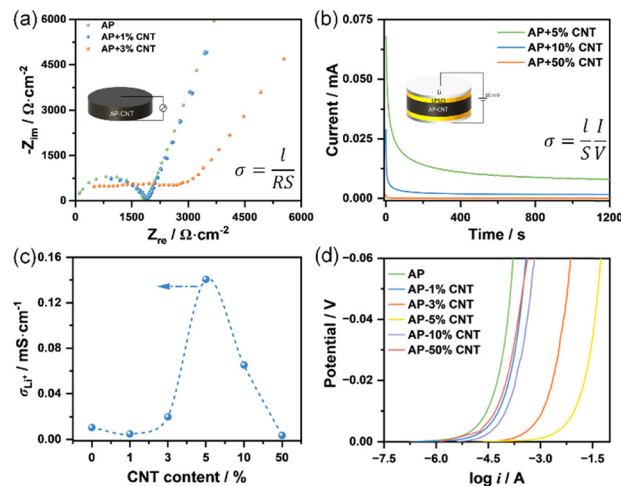


Fig. 2 (a) EIS profiles of SS/AP-*x*% CNT/SS (*x* = 0, 1, 3) cells at 55 °C. (b) CA measurement of Li/LPSCI/AP-*x*% CNT/LPSCI/Li (*x* = 5, 10, 50) cells at 55 °C. The applied voltage is 10 mV. (c) Calculated ionic conductivity of AP-*x*% CNT from (a) and (b). (d) Tafel curves of symmetric Li/AP-*x*% CNT/LPSCI/AP-*x*% CNT/Li (*x* = 1, 3, 5, 10 and 50) cells at 55 °C.

conductivity displays a non-linear evolution, reaching a maximum of 1.41 × 10⁻⁴ S cm⁻¹ at 5 wt% CNT, which is much higher than that of pristine AP (1.05 × 10⁻⁵ S cm⁻¹), as shown in Fig. 2c. This behaviour is noteworthy, as the ionic conductivity typically decreases with increasing CNT content. These results are further supported by Tafel measurements. Fig. 2d shows the Tafel curves of symmetric cells incorporating AP–CNT composite interlayers obtained at 55 °C. Consistent with the ionic conductivity results, the AP–5% CNT interlayer exhibits the highest exchange current density of 0.02 mA cm⁻². The elevated exchange current density indicates accelerated interfacial kinetics and reduced energy barriers for Li⁺ transfer, demonstrating that the composite interlayer facilitates efficient ion transport at the Li interface rather than merely improving electronic conductivity.

To gain deeper insight into the mechanism underlying the enhanced ionic conductivity of AP–5% CNT, Raman spectroscopy was carried out on AP–CNT composites by varying CNT contents. As shown in Fig. 3a and Fig. S4, the samples containing CNTs exhibit a prominent D band at 1350 cm⁻¹ and a G band at 1575 cm⁻¹, whereas pristine AP shows no significant Raman features. Quantitative analysis of the I_D/I_G ratio is presented in Fig. 3b. Pure CNTs display the lowest I_D/I_G ratio, consistent with their high graphitic order and minimal defect density. Upon mixing with AP, the I_D/I_G ratio increases and reaches a maximum at 5 wt% CNT content, indicating the introduction of additional structural disorder within the CNTs. This enhanced defect density likely originates from the perturbation of the sp² carbon network caused by the electronic interaction between CNTs and AP.

To validate this interfacial electronic interaction, the surface potential of the AP–CNT composites was further examined using Kelvin probe force microscopy (KPFM). Variations in the contact potential difference (CPD) directly reflect changes

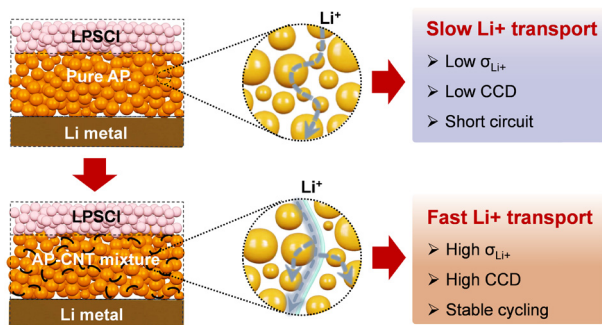


Fig. 1 Schematic illustration of Li⁺ transport facilitated by the built-in electric field within the AP–CNT composite interlayer. While Li⁺ transport in pure AP occurs along particle boundaries, the AP–CNT interface generates an internal electric field that accelerates Li⁺ migration, creating highly efficient transport pathways.



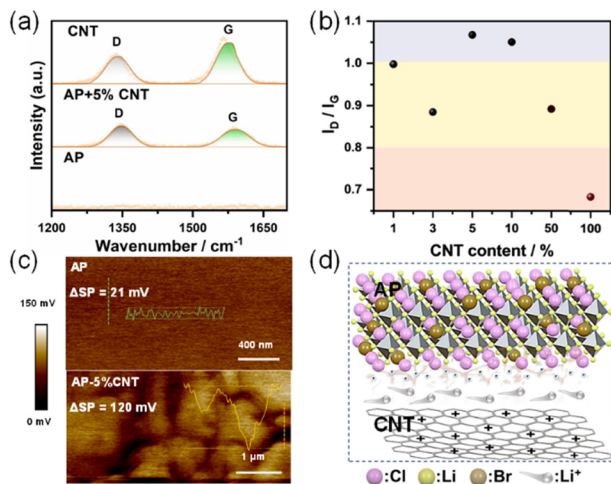


Fig. 3 (a) Raman spectra of AP, CNT and AP-5% CNT. (b) I_D/I_G ratios of AP-CNT composites with varying CNT contents (0, 1, 3, 5, 10, 50, and 100). (c) KPFM images and contact potential differences of AP and AP-5% CNT. (d) Scheme of the built-in electric field.

in surface potential, with the average CPD values extracted from line profile measurements. As shown in Fig. 3c and Fig. S5, pristine AP exhibits an average CPD of 21 mV, whereas AP-1%, AP-5%, and AP-10% CNT display a substantially higher CPD of 63, 120 and 100 mV, respectively. The non-monotonic variation of CPD with CNT content, with a maximum at 5 wt% CNT, suggests that the surface potential change is not solely determined by the intrinsic work-function difference between AP and CNTs, but instead reflects interfacial charge redistribution. As illustrated in Fig. S6, the work-function mismatch drives electron transfer from CNTs to AP until Fermi-level equilibration is achieved, leading to electron accumulation on the AP surface and the formation of positively charged vacancies on the CNT surface. This charge redistribution induces band bending in the AP phase and establishes a built-in electric field at the interface, as shown in Fig. 3d.^{22,23} Such a field provides an additional electrostatic driving force that alters the local potential energy landscape of Li^+ . Therefore, it effectively lowers the migration energy barrier and facilitates interfacial ion transport, accounting for the enhanced ionic conductivity observed in the AP-5% CNT system.^{22,24}

The impact of the AP-5% CNT interlayer on Li plating/stripping was investigated in Li/LPSC/Li symmetric cells, in which the LPSC electrolyte was sandwiched between two AP or AP-5% CNT interlayers. The symmetric cells were charged and discharged at stepwise-increased current densities for 1 h. As shown in Fig. 4a and b, the cell with the AP interlayer short-circuited immediately at a current density of 0.1 mA cm^{-2} , whereas the AP-5% CNT interlayer enabled a much higher critical current density of 2.4 mA cm^{-2} , demonstrating significantly enhanced interfacial stability. This value is also higher than that without an interlayer (Fig. S7, SI). The effect of the interlayer was further evaluated in Li/LPSC/AIO@NMC811 full cells with AP or AP-5% CNT interlayer at 55°C . For comparison, Li/LPSC/AIO@NMC811 full cells without an interlayer

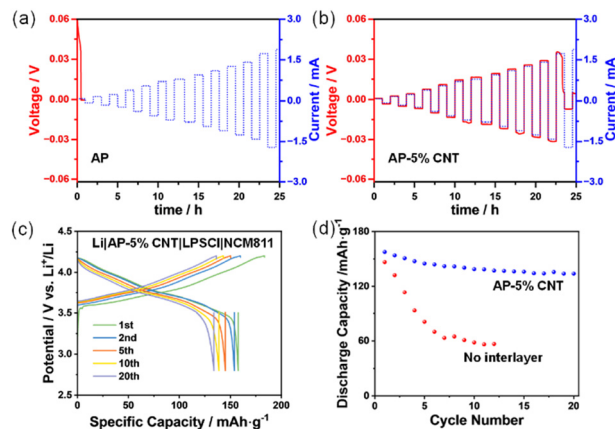


Fig. 4 (a) Galvanostatic cycling of (a) Li/AP/LPSC/AP/Li cells and (b) Li/AP-5% CNT/LPSC/AP-5% CNT/Li cells at stepwise-increased current densities at 55°C . (c) Charge-discharge curves of the Li/AP-5% CNT/LPSC/AIO@NMC811 cell at 55°C . (d) Cycling performance of solid-state batteries with/without the AP-CNT interlayer at a rate of 0.1C.

were also assembled. The capacity of the Li/LPSC/LSO@NMC811 full cell without an interlayer rapidly decayed to 58.3 mAh g^{-1} after only 10 cycles (Fig. S8, SI). This degradation arises from severe side reactions between the Li metal and the LPSC electrolyte as well as the pore formation during Li metal stripping at low pressures due to the limited creep rate of the Li metal, which leads to increased polarization and suppressed capacity utilization.^{5,25-28} After adding an AP interlayer, the battery short-circuits even during the first charge cycle (Fig. S9, SI). The failure is attributed to the low ionic conductivity of AP, which readily promotes dendrite growth. Furthermore, the poor wettability (Fig. S10, SI) between the AP and Li metal creates localized current hotspots, further accelerating non-uniform Li deposition and leading to rapid short-circuiting.²⁹ In contrast, the Li/LPSC/LSO@NMC811 full cell with AP-5% CNT exhibits superior performance, delivering a specific discharge capacity of 157.7 mAh g^{-1} in the first cycle and maintaining a discharge capacity of 134.0 mAh g^{-1} after 20 cycles. Fig. 4d highlights the advantages of the AP-CNT interlayer by comparing the cycling performance of solid-state batteries with and without the interlayer. This can be attributed to the multifunctional nature of the AP-CNT interlayer. On the one hand, the built-in electric field generated by AP-5% CNT facilitates Li^+ migration across the interface. This interfacial effect effectively reduces ion transport resistance and ensures a more uniform Li^+ flux during cycling. Furthermore, the AP-CNT layer stabilizes lithium metal, mitigates detrimental side reactions between the lithium metal anode and the sulfide electrolyte, and withstands interfacial stress during repeated plating/stripping. Overall, these synergistic effects lead to slower capacity decay and higher discharge capacity retention in cells with the AP-5% CNT interlayer.

In summary, we have demonstrated an AP-CNT composite interlayer that simultaneously enhances interfacial stability and Li^+ transport in all-solid-state lithium-metal batteries. By incorporating a small fraction of CNTs into the anti-perovskite



$\text{Li}_2\text{OHCl}_{0.75}\text{Br}_{0.25}$, interfacial electron transfer generates a built-in electric field, thereby facilitating Li^+ migration and promoting a more homogeneous interfacial ion flux. As a result, the AP-5% CNT interlayer exhibits an ionic conductivity higher than that of pristine AP and enables stable lithium plating/stripping with a high critical current density of 2.4 mA cm^{-2} at 55°C . When applied in full cells, this multifunctional interlayer markedly improves cycling stability and capacity retention by mitigating parasitic reactions, suppressing dendrite growth, and accommodating interfacial stress. Beyond this specific system, the concept of built-in electric fields to accelerate Li^+ transport could be extended to anodes and other solid-state electrolyte systems, such as oxide- and polymer-based electrolytes,^{30,31} providing a general and effective strategy for designing high-performance, dendrite-free solid-state lithium-metal batteries.

Conflicts of interest

There are no conflicts to declare.

Data availability

The data supporting this article have been included as part of the supplementary information (SI). Supplementary information: XRD, chronoamperometric curves, Raman spectra, galvanostatic cycling curves, charge-discharge curves and further experimental details. See DOI: <https://doi.org/10.1039/d6cc00741d>.

Acknowledgements

We gratefully acknowledge the financial support from the Foundation for Transportation Technologies Office of the Department of Energy of the United States. The authors are grateful for the startup from VT. This research was performed on APS beam time award (DOI: <https://doi.org/10.46936/APS-190182/60014244>) from the Advanced Photon Source and Center for Nanoscale Materials, a U.S. Department of Energy (DOE) Office of Science user facility operated for the DOE Office of Science by Argonne National Laboratory under Contract No. DE-AC02-06CH11357.

References

- 1 E. P. Alsac, D. L. Nelson, S. G. Yoon, K. A. Cavallaro, C. Wang, S. E. Sandoval, U. D. Eze, W. J. Jeong and M. T. McDowell, *Chem. Rev.*, 2025, **125**, 2009–2119.
- 2 S. E. Sandoval, C. G. Haslam, B. S. Vishnugopi, D. W. Liao, J. S. Yoon, S. H. Park, Y. Wang, D. Mitlin, K. B. Hatzell, D. J. Siegel, P. P. Mukherjee, N. P. Dasgupta, J. Sakamoto and M. T. McDowell, *Nat. Mater.*, 2025, **24**, 673–681.
- 3 W. D. Richards, L. J. Miara, Y. Wang, J. C. Kim and G. Ceder, *Chem. Mater.*, 2015, **28**, 266–273.
- 4 T. Swamy, X. Chen and Y.-M. Chiang, *Chem. Mater.*, 2019, **31**, 707–713.
- 5 D. H. S. Tan, E. A. Wu, H. Nguyen, Z. Chen, M. A. T. Marple, J.-M. Doux, X. Wang, H. Yang, A. Banerjee and Y. S. Meng, *ACS Energy Lett.*, 2019, **4**, 2418–2427.
- 6 M. Zhou, Y. Liao, L. Li, R. Xiong, G. Shen, Y. Chen, T. Huang, M. Li, H. Zhou and Y. Zhang, *J. Energy Chem.*, 2023, **85**, 181–190.
- 7 H. Chen, C. Meng, Z. Jiao, A. Yuan and H. Zhou, *Energy Fuels*, 2025, **39**, 4069–4078.
- 8 S. Qiang, J. Chen, S. Huang, H. Xu, X. Zhuo, H. Zhou, A. Yuan, H. Zhou and Y. Qiao, *J. Colloid Interface Sci.*, 2026, **701**, 138722.
- 9 Z. Zhao, J. Chen, H. Zhou, B. Hu, C. Meng, A. Yuan and Y. Qiao, *Metals Adv.*, 2026, **40**, 62–70.
- 10 Y. Su, L. Ye, W. Fitzhugh, Y. Wang, E. Gil-González, I. Kim and X. Li, *Energy Environ. Sci.*, 2020, **13**, 908–916.
- 11 P. Ren, X. Wang, B. Huang, Z. Liu and R. Liu, *J. Energy Storage*, 2024, **82**, 110200.
- 12 X. Fan, X. Ji, F. Han, J. Yue, J. Chen, L. Chen, T. Deng, J. Jiang and C. Wang, *Sci. Adv.*, 2018, **4**, eaau9245.
- 13 J. Su, M. Pasta, Z. Ning, X. Gao, P. G. Bruce and C. R. M. Grovenor, *Energy Environ. Sci.*, 2022, **15**, 3805–3814.
- 14 R. Guo, K. Zhang, W. Zhao, Z. Hu, S. Li, Y. Zhong, R. Yang, X. Wang, J. Wang, C. Wu and Y. Bai, *Energy Mater. Adv.*, 2023, **4**, 0022.
- 15 L. Shen, Z. Wang, S. Xu, H. M. Law, Y. Zhou and F. Ciucci, *Nat. Commun.*, 2025, **16**, 3687.
- 16 Z. D. Hood, H. Wang, A. Samuthira Pandian, J. K. Keum and C. Liang, *J. Am. Chem. Soc.*, 2016, **138**, 1768–1771.
- 17 Z. Wang, J. Xia, X. Ji, Y. Liu, J. Zhang, X. He, W. Zhang, H. Wan and C. Wang, *Nat. Energy*, 2024, **9**, 251–262.
- 18 S.-Q. Li, Z. Wang, Z. Sun, X. Zheng, F.-H. Du, J.-N. Yang, Y.-W. Zhang, Y. Sun, H. Zong, M.-L. Guo, X.-Y. Wu, Y. Xu, G. Wang, C. Wu, Q. Zhang and K.-X. Wang, *ACS Energy Lett.*, 2025, **10**, 5972–5981.
- 19 R. Xu, F. Han, X. Ji, X. Fan, J. Tu and C. Wang, *Nano Energy*, 2018, **53**, 958–966.
- 20 Y. Zhao and L. L. Daemen, *J. Am. Chem. Soc.*, 2012, **134**, 15042–15047.
- 21 C. Yang, L. Zhang, B. Liu, S. Xu, T. Hamann, D. McOwen, J. Dai, W. Luo, Y. Gong, E. D. Wachsman and L. Hu, *Proc. Natl. Acad. Sci. U. S. A.*, 2018, **115**, 3770–3775.
- 22 L. Du, Y. Zhang, Y. Xiao, D. Yuan, M. Yao and Y. Zhang, *Energy Environ. Sci.*, 2025, **18**, 2949–2961.
- 23 G. Xiao, K. Yang, Y. Qiu, P. Shi, G. Zhong, X. An, Y. Ma, L. Chen, S. Guo, J. Mi, Z. Han, T. Hou, H. Yan, Y. Tian, X. Zhang, Y. Cao, M. Liu, Z. Zhou and Y. B. He, *Adv. Mater.*, 2025, e2415411, DOI: [10.1002/adma.202415411](https://doi.org/10.1002/adma.202415411).
- 24 M. Yang, K. Yang, Y. Wu, Z. Wang, T. Ma, D. Wu, L. Yang, J. Xu, P. Lu, J. Peng, Z. Jiang, X. Zhu, Q. Gao, F. Xu, L. Chen, H. Li and F. Wu, *ACS Nano*, 2024, **18**, 16842–16852.
- 25 S. Wenzel, S. J. Sedlmaier, C. Dietrich, W. G. Zeier and J. Janek, *Solid State Ionics*, 2018, **318**, 102–112.
- 26 J. M. Doux, H. Nguyen, D. H. S. Tan, A. Banerjee, X. Wang, E. A. Wu, C. Jo, H. Yang and Y. S. Meng, *Adv. Energy Mater.*, 2019, **10**, 1903253.
- 27 C. Wang, Y. Liu, W. J. Jeong, T. Chen, M. Lu, D. L. Nelson, E. P. Alsac, S. G. Yoon, K. A. Cavallaro, S. Das, D. Majumdar, R. Gopalaswamy, S. Xia and M. T. McDowell, *Nat. Mater.*, 2025, **24**, 907–916.
- 28 L. Di, Z. Huang, L. Gao, Y. Zuo, J. Zhu, M. Sun, S. Zhao, J. Zheng, S. Han and R. Zou, *Sci. Adv.*, 2025, **11**, eadw9590.
- 29 G. Cen, H. Yu, R. Xiao, L. Ben, R. Qiao, J. Zhu, X. Zhang, G. Liu, K. Jiang, X. Yao, H. Zhang and X. Huang, *Nat. Sustainability*, 2025, **8**, 1360–1370.
- 30 M. Liu, W. Xu, S. Liu, B. Liu, Y. Gao and B. Wang, *Adv. Sci.*, 2024, **11**, e2402915.
- 31 J. Kondo, Y. Iga, K. Otani, R. Noba, H. Hikosaka, Y. Kondo, Y. Katayama and Y. Yamada, *J. Phys. Chem. C*, 2025, **129**, 15231–15238.

