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Efficiency boost in perovskite solar cells via TiO₂ nanodiscs embedded in the MoSe₂ electron transport layer revealed by optoelectronic simulations†

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To improve the performance of inverted perovskite solar cells (IPSCs), we introduce a novel approach to enhance the devices' efficiency notably using the Finite Element Method (FEM). Our novel strategy incorporates a cutting-edge metasurface-based reflector featuring titanium dioxide (TiO₂) nanodiscs within a MoSe₂ layer, employed as an electron transport layer (ETL). Demonstrating a substantial improvement in light reflection from the lower part of the structure, the TiO₂ nanodiscs as a metasurface-based reflector enhance electron transfer. Notably, the metasurface-based perfect reflector, incorporating TiO₂ nanodiscs, outperforms other TiO₂ nanocube variations with an impressive light reflectance of 97.95%. Exploring different materials for ETLs and hole transfer layers (HTLs), we identify molybdenum diselenide (MoSe₂) as a potent secondary absorbent material, featuring a smaller bandgap than the primary absorbent CH₃NH₃PbI₃ (MAPbI₃), thereby intensifying the electric field within the active layer and improving Power Conversion Efficiency (PCE). In the final evaluation, our inverted metasurface-based device structure (indium tin oxide (ITO)/cuprous oxide (Cu₂O)/MAPbI₃/TiO₂ nanodiscs and MoSe₂/aluminum (Al)/silicon dioxide (SiO₂)) significantly enhances the solar cell's electrical characteristics compared to the planar reference structure (ITO/copper(i) thiocyanate (CuSCN)/MAPbI₃/TiO₂/Al), with noteworthy increases in short circuit current density (J_{sc}), open circuit voltage (V_{oc}), and PCE values from 17.98 mA cm⁻² to 21.91 mA cm⁻², 1.03 V to 1.07 V, and 15.33% to 19.17%, respectively. This comprehensive investigation underscores the promising potential of our proposed inverted metasurface-based device structure for advancing solar cell technology.

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1. Introduction

The increasing demand for clean energy has become a critical issue due to the shortage of fossil fuel resources, the resulting air pollution, and the urgent need for immediate action to combat climate change.^{1–3} Therefore, the rapid development of new energy harvesting technologies to provide alternative clean, renewable, and sustainable energy from natural sources, such as solar—the most abundant energy resource—is crucial.^{4–7} In recent years, perovskite solar cells have achieved significant advancements, including higher PCE and lower manufacturing costs compared to the market-dominant silicon solar cells.^{8–11} Perovskite materials find application in diverse single-junction

and tandem solar cell architectures, both for indoor and outdoor environments.^{12–16} Among all different subgroups of perovskites, there has been extensive interest on metal halide perovskite solar cells. This is widely because of their remarkable optical and electrical properties such as noticeably high light absorption spectrum, simple fabrication processes, and substantial carrier mobility and diffusion lengths.^{17,18} The selection of MAPbI₃ is driven by the abundance of research literature and datasets, enabling a data-driven approach to model building.^{19–22} Therefore, MAPbI₃ was used as the perovskite active layer in this work due to the wide light absorption spectra (wavelengths region including 300 to 800 nm) and proper carrier transfer (high carrier mobility and permittivity).^{23–25} In order to enhance the solar cell's PCE, metasurface structures were developed to manipulate the incoming incident light.^{26–28} Furthermore, metasurface structures are easy to fabricate and have proven to increase the light absorption because of the internal light reflections that cause light interactions with other layers.^{29,30} Thus, deploying this technology in solar cell devices has raised attention and is currently used by

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many researchers.^{31–34} The integrated periodic nanostructured metasurfaces perform mainly through the Snell's law by trapping the incident light and inclining the total absorption of the solar cell.^{35,36} Therefore, opting the material, shape and size of the nanostructure have serious consequence in the system's efficiency.^{37–40} The TiO₂ nanodiscs can be grown on an Al substrate using a hydrothermal technique involving several steps.^{41,42} A research on a nanostructured cubic metasurface reflector consisting of three layers of TiO₂/Ag/SiO₂ proved to have a promising performance.⁴³ Therefore, in this paper, different shapes (nanodiscs and three various types of nanocubes) of TiO₂ were selected as the integrated metasurface nanostructure and the ETL on top of a planar Al/SiO₂ structure to reach the optimized structure. Then, the nanostructure's size and position's degrees were optimized for one of the nanocubic structures to reach the highest light reflection to increase the total light absorption of the solar cell. However, the light reflection was further improved by using TiO₂ nanodiscs instead. In order to increase the light reflection by the metasurface nanostructure, another optical study was done on discovering the material with suitable refractive index to surround the TiO₂ nanostructure. Thereby, ETL materials such as zinc oxide (ZnO), Al-doped zinc oxide (AZO), cerium oxide (CeO₂), and titanium dioxide–graphene (5%) (TiO₂–Graphen (5%) or TIGO-5) were studied. Plus, this is a well-recognized issue that TiO₂ only absorbs the visible light with the wavelength range of 300 to almost 400 nm.⁴⁴ With regard to this, it can be argued that an inverted solar cell would be more beneficial to increase the total light reflection inside the device since the TiO₂ nanostructure would be placed on a non-transparent contact like Al, which results in more light interaction in the ETL and consequently other layers. The IPSCs have become prominent in light of their advantageous characteristics like high stability and PCE.^{45,46} They consist of a reversed architecture, where the ETL and HTL are replaced by each other.⁴⁷ Hence, the HTL material was required to have a low light absorption rate in order for the incident light to reach the active layer and the ETL as well. For this reason, different materials including CuSCN, Spiro-OMeTAD, copper(i) iodide (CuI), and CuO₂ were studied as HTL. Additionally, another active layer, MoSe₂, was used to enhance the light absorption, which was located below the MAPbI₃ layer. MoSe₂ exhibits a compelling set of properties that make it a promising candidate for integration into perovskite solar cells.⁴⁸ Notably, its high electron mobility facilitates efficient charge transport, minimizing energy losses and boosting overall efficiency.^{49,50} Furthermore, MoSe₂ offers tunable electrical characteristics, allowing for precise energy level alignment within the perovskite layer, crucial for optimal charge transfer.⁵¹ Its strong light absorption across the visible and near-infrared spectrum maximizes the utilization of solar energy, while its robust chemical and thermal stability ensure long-term device performance and reliability under diverse operating conditions.⁵² Finally, the layered structure of MoSe₂ aligns well with the architecture of perovskite solar cells, potentially simplifying integration and enabling facile exfoliation for device fabrication.⁵³ Finally, the results were compared to reach the optimized device with the highest PCE,

resulted in the improvement in values of J_{sc} , V_{oc} , and PCE from 17.98 mA cm⁻² to 21.91 mA cm⁻², 1.03 V to 1.07 V, and 15.33% to 19.17%, respectively.

2. Theory

In this section, the electrical and optical models and their coupling were discussed. Additionally, the simulation details related to these models were also described.

2.1. Optoelectronic model

Two different sets of optical and electrical equations were analyzed in order to calculate various parameters, as shown in Fig. 1. The optical equations utilized Maxwell's equation (eqn (1)) to determine the electric field (E) with the input parameters including: refractive index (n) and extinction coefficient (k) of all layers.

$$\nabla \times (\nabla \times E) - k_0 \epsilon_r E = 0 \quad (1)$$

The calculation of the electric field resulted in the optical outputs such as reflection, absorption, and the optical generation rate for electrons and holes (G_n , G_p). The optical generation rate referred to the number of the generated electron-hole pairs at each point in the device due to the absorption of photons.

The dark current density and electron-hole distribution and recombination rate of electron and hole (R_n , R_p) was then calculated by solving the Poisson's and Continuity equations (eqn (2)–(4)), as well as the related equations to Shockley–Reed–Hall (SRH) recombination by applying the input parameters such as band gap (E_g), electron affinity (χ), dielectric constant (ϵ), electron mobility (μ_n), hole mobility (μ_p), density state of valence band (N_v), density state of conduction band (N_c), acceptor and donor doping (N_A, N_D), electron lifetime (τ_n), and hole lifetime (τ_p) of each layer.

$$\nabla \cdot (\epsilon_0 \cdot \epsilon_r \cdot \nabla \phi) = -\rho \quad (2)$$

$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla J_n + G_n - R_n \quad (3)$$

$$\frac{\partial p}{\partial t} = \frac{1}{q} \nabla J_p + G_p - R_p \quad (4)$$

The current density in the cell was solved by coupling the two sets of equations. Fig. 1 illustrates the flowchart of the optoelectronic model. A standard AM 1.5 with the light intensity of 100 mW cm⁻² perpendicular to the surface was used as the incoming light source. Different layers behaved differently after interacting with the incoming electric field according to their n and k . The dark current density, which has a direct relation with the amount of SRH recombination in the system, is calculated by solving the electric model without considering the incident electric field. To calculate the J_{sc} , both optical and electrical physics were solved, simultaneously. In this study, the



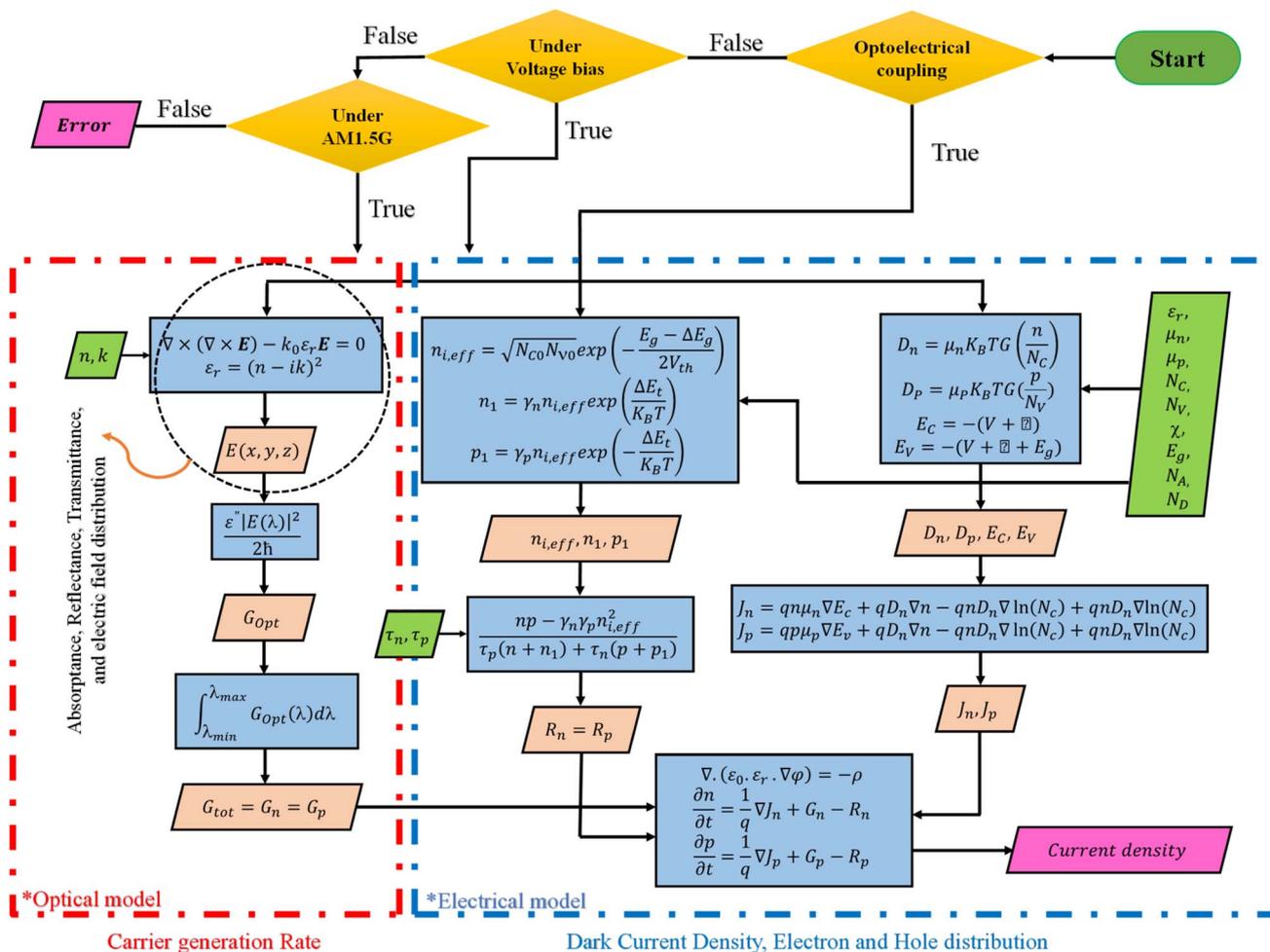


Fig. 1 Flowchart of the optoelectronic model to calculate solar cell parameters, where the red and blue sections are considered as separated optical and electrical models, respectively. In the optical model, the Maxwell's equation is solved to reach the carrier generation rate, whereas in the electrical model, the Poisson's and Continuity equations are solved to achieve the dark current density. The coupling of two models results in the photocurrent density of the solar cell.

wavelength in the range of 300 to 1200 nm was swept with 20 nm steps, and the current density spectrum was calculated in terms of wavelength. A sum of current densities in the

abovementioned wavelength range with 20 nm steps was added to the dark current density to calculate the total current density in terms of voltage. In addition to SRH recombination for the

Table 1 Electrical parameters for the used materials in the electrical model, where ϵ_r denotes the dielectric constant, N_C and N_V present the density states of conduction and valence bands, respectively, μ_n and μ_p show the electron and hole mobility, respectively, χ is the electron affinity, E_g is the band gap, N_A and N_D are the acceptor and donor doping, respectively, and τ_n and τ_p show the electron and hole lifetime, respectively

Parameters	ZnO	AZO	CeO ₂	TiGO-5	CuSCN	Spiro-OMeTAD	CuI	Cu ₂ O	MAPbI ₃	TiO ₂	MoSe ₂
ϵ_r	9	9	9	7.3	10	3	12.5	7.5	6.5	9	7.29
N_C (cm ⁻³)	2.2×10^{18}	2.2×10^{18}	1×10^{20}	1×10^{19}	1.79×10^{19}	2.2×10^{18}	2.8×10^{19}	2×10^{19}	1.6×10^{19}	1×10^{19}	9×10^{18}
N_V (cm ⁻³)	1.8×10^{19}	1.8×10^{19}	2×10^{21}	1×10^{19}	2.51×10^{19}	1.8×10^{19}	1×10^{19}	1×10^{19}	5.4×10^{19}	1×10^{19}	9×10^{18}
μ_n/μ_p (cm ² V ⁻¹ s ⁻¹)	100/25	10/5	100/25	12/12	25/25	$2 \times 10^{-4}/2 \times 10^{-4}$	100/43.9	8600/200	50/50	20/10	25/100
χ (eV)	4	4.4	3.5	4.24	1.7	2.2	2.1	3.4	3.93	4	4.32
E_g (eV)	3.3	3.35	4.6	3.1	3.6	3.0	3.1	2.2	1.55	3.2	1.4
N_A (cm ⁻³)	—	—	—	—	5×10^{18}	1×10^{18}	1×10^{18}	1×10^{18}	1×10^{13}	—	—
N_D (cm ⁻³)	1×10^{18}	1×10^{18}	1×10^{21}	5×10^{18}	—	—	—	—	—	5×10^{18}	3×10^{18}
τ_n/τ_p (ns)	0.03/0.03	0.03/0.03	5/2	2/2	5/5	0.1/0.1	3/3	3/3	8/8	5/2	2.4/2.4
Ref.	54	55	56 and 57	58–61	56	56 and 62	56 and 63	64 and 65	62 and 64	66	67



metal contact, surface recombination is also considered for the simulation. The surface recombination velocity is set at $1 \times 10^7 \text{ cm s}^{-1}$ for both electrons and holes.

2.2. Simulation details

The materials' complex refractive indices (n and k) are plotted in Fig. S3, S5† and 7g. Also, the electrical parameters of the solar cell's layers (ZnO, AZO, CeO₂, TiO₂-Graphen (5%) (TiGO-5), CuSCN, Spiro-OMeTAD, CuI, Cu₂O, MAPbI₃, TiO₂, MoSe₂),

including: ϵ , N_C , N_V , μ_n , μ_p , χ , E_g , N_A , N_D , τ_n and τ_p are provided in Table 1.

3. Results and discussion

3.1. Metasurface based reflector design

Fig. 2 shows the schematic of different types of structures in solar cells. The common inverted solar cell structures, as depicted in Fig. 2a, typically use bulk TiO₂ as ETL on the contact. To enhance the reflection from the TiO₂ surface, metasurface structures can be utilized. Unlike the natural behavior of materials, metasurface structures exhibit unique physical properties. Fig. 2b illustrates a reflector structure consisting of Al on SiO₂ (as anti-reflection coating) substrate. Also, Fig. 2c shows the structure when the bulk TiO₂ is applied on Al and SiO₂. Finally, the schematic of the proposed metasurface unit cell consisting of TiO₂ nanocubes on the Al and SiO₂ substrate is presented in Fig. 2d. TiO₂ is a preferred material for solar cell structures due to its high refractive index, smooth surface, and low heat loss, particularly in the visible region. The dimensions of the TiO₂ nanocubes, including length (L), width (W), and height (H), are represented by specific parameters. Additionally, the angle of the nanocubes relative to the horizontal axis is shown with the parameter α . Based on the results, it can be argued that the use of TiO₂ nanocubes on Al and SiO₂ substrate can improve the efficiency of photonic structures and solar cells

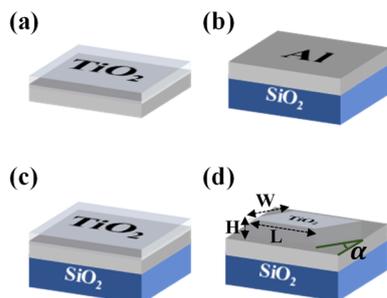


Fig. 2 (a) A common structure used as ETL in solar cells. (b) A reflector consisting of Al and SiO₂. (c) The common structure of ETL in solar cells on SiO₂ substrate. (d) TiO₂/Al/SiO₂ metasurface reflector structure.

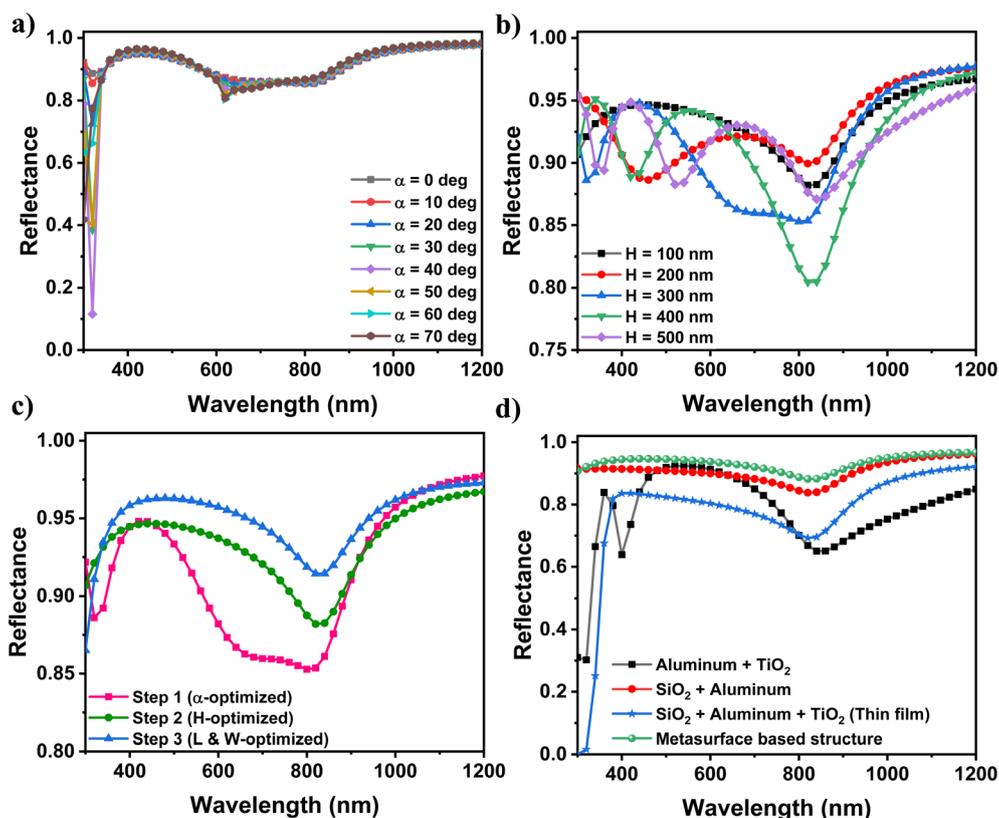


Fig. 3 The reflectance spectrum in terms of wavelength in the range of 300 to 1200 nm for (a) different values of α (nanocubic angle with horizontal axis), while considering $H = 300 \text{ nm}$, $L = 150 \text{ nm}$, and $W = 90 \text{ nm}$, (b) different values of H , while considering $\alpha = 0$, $L = 150 \text{ nm}$, and $W = 90 \text{ nm}$, (c) different optimizations including three steps, (d) comparing the optimized metasurface structure with other structures in Fig. 2.



by enhancing the reflection from the TiO₂ surface. The height of the Al and SiO₂ substrate is 150 nm and 300 nm, respectively, while the length and width of the unit cell are both equal to 250 nm. Opting the mentioned dimensions serves two purposes: first, the unit cell's dimensions must be smaller than the incident wavelength, and second, the structure must be easily manufactured, particularly through electron beam lithography. The main objective of designing the proposed metasurface structure is to develop a reflector that reflects the transmitted light through the structure, to increase the total absorption in the structure.

Each of the structures shown in Fig. 2 behaved differently to electromagnetic waves, and consequently, the amount of

Table 2 Reflectance values for (a) different values of α considering $H = 300$ nm, $L = 150$ nm and $W = 90$ nm, (b) different values of H considering $\alpha = 0^\circ$, $L = 150$ nm and $W = 90$ nm, (c) different values of L considering $\alpha = 0^\circ$, $H = 100$ nm and $L/W = 15/9$, and (d) different provided structures in Fig. 2

(a)		
$H/L/W$ (nm)	α°	Reflectance (%)
300/150/90	0	91.57
300/150/90	10	91.50
300/150/90	20	91.20
300/150/90	30	90.46
300/150/90	40	89.83
300/150/90	50	90.46
300/150/90	60	91.06
300/150/90	70	91.16
(b)		
$\alpha^\circ/L/W$ (nm)	H (nm)	Reflectance (%)
0/150/90	100	93.36
0/150/90	200	92.97
0/150/90	300	91.57
0/150/90	400	91.11
0/150/90	500	91.81
(c)		
α°/H (nm)	L (nm)	Reflectance (%)
0/100	90	91.82
0/100	110	92.29
0/100	130	92.82
0/100	150	93.36
0/100	170	93.90
0/100	190	94.47
0/100	210	94.20
(d)		
Structures	Reflectance (%)	
Fig. 1a	78.08	
Fig. 1b	90.64	
Fig. 1c	77.35	
Fig. 1d	94.47	

reflection from each structure varied. To increase the amount of reflectance from the reflector surface, various parameters, including H , L , W , and α were optimized. Therefore, the angle of the nanocube with the X -axis, represented by the parameter α , was swept from 0° to 70° with steps of 10° to. According to Fig. 3a and Table 2, a, and the incident power has the highest amount of reflection in the structure, reaching 91.57%, when the angle is 0° . The H , L and W parameters of the TiO₂ nanocubes are also effective in determining the amount of reflectance of the structure. By examining different heights of nanocubes from 100 to 500 nm with steps of 100 nm, different results were obtained. As the height was increased, the amount of reflectance from the surface decreased, and the highest amount of reflectance from the structure was at $H = 100$ nm according to Table 2, b. This could be explained by the frequency shift in the reflectance spectrum, where it reached the highest amount of reflection in the wavelength range of 300 to 1200 nm at this height. The amount of reflectance was increased to 93.36% by optimizing the height, which is more efficient compared to when the angle was optimized, 91.57%. Then, the L parameter was varied from 90 nm to 210 nm with steps of 20 nm while keeping the angle and height constant at their optimized values. The best result was obtained at $L = 190$ nm, which increased the amount of reflection up to 94.47% compared to the height optimization. The ratio of L/W was considered equal to 15/9. The optimal dimensions of the TiO₂ nanocubes were determined to be $\alpha = 0^\circ$, $H = 100$ nm, and $L = 190$ nm. By optimizing these parameters, the total reflectance of the structure in the wavelength range of 300 to 1200 nm reached a significant amount of 94.47%. The amount of reflection from the surface of the depicted structures in Fig. 2, is displayed in Fig. 3d. The metasurface structure exhibited a significant increase in reflectance compared to other structures, particularly the structure in Fig. 2a, which is commonly used in solar cells. According to Table 2, d, the amount of reflection in the wavelength range of 300 to 1200 nm in the provided metasurface structure in Fig. 2d was significantly increased from 78.08% to 94.47% compared to the given structure in Fig. 2a, which is commonly used in solar cells.

In order to enhance the light reflectance, other shapes of metasurface structures with TiO₂ was further investigated. The depicted nanocubic structure in Fig. 2a, proved to reflect 94.47% of the incident light power. According to Fig. 4e, the provided structures in Fig. 4b–d shared higher values of light reflectance equal to 95.83%, 96.52%, and 97.95%, respectively. It was apparent that the TiO₂ nanodiscs structure with the height of 100 nm and the diameter of 190 nm was selected as the most efficient structure for additional studies. Plus, the TiO₂ nanodiscs have proven transfer electrons more effectively. Fig. 4e displays the reflectance diagram of the reflector structure in terms of wavelength in the range of 300 to 1200 nm, when using different types of nanocubes and nanodiscs as depicted in Fig. 4a–d. After rigorous examination, it was discovered that the materials' refractive indices, which surrounds the nanostructures, had a noticeable impact on the amount of the reflectance from the metasurface-based reflector structure. As



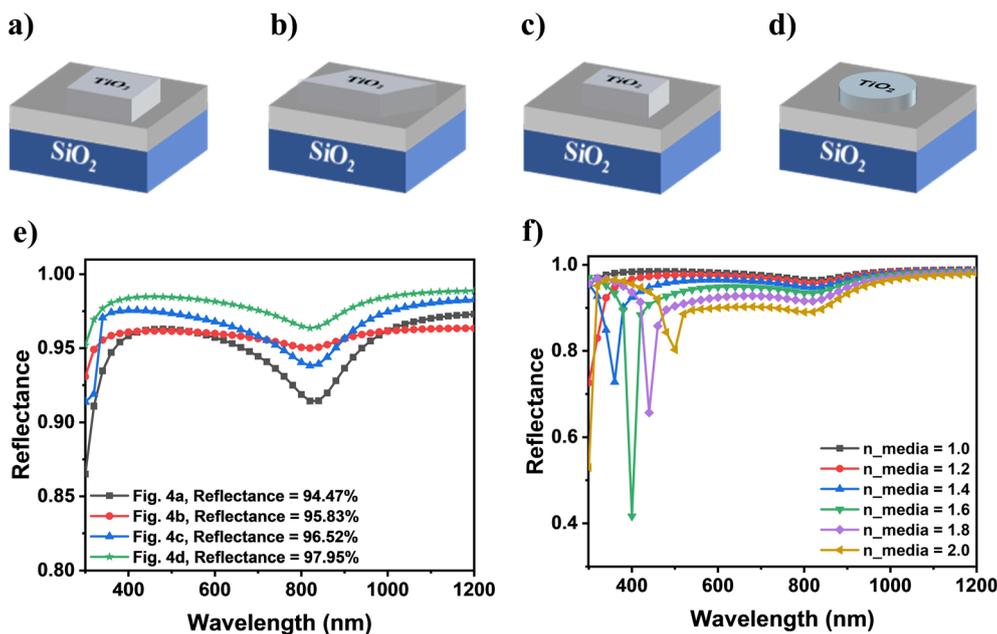


Fig. 4 The schematic of the metasurface reflector nanostructure with (a–c) three different types of nanocubes and (d) nanodiscs. (e) The reflectance spectrum of the depicted nanostructures in (a–d). (f) The reflectance spectrum for the metasurface reflector considering different refractive indices values (1 to 2 with steps of 0.2) for the media around the TiO₂ nanodiscs.

illustrated in Fig. 4f, by changing the refractive index of the media around the TiO₂ nanodiscs, the reflectance from the structure was changed. The investigated refractive index range was from 1 to 2 with steps of 0.2. According to the obtained results, the amount of reflection was higher in medias with lower refractive indices, which suggested that using materials with low refractive indices would result in more light absorption in the active layer of the solar cells. In conclusion, it can be claimed that photonic and metasurface nanostructures would improve the light absorption, ease fabrication, and improve efficiency in solar cells.

The purpose of designing the TiO₂/Al/SiO₂ reflector structure is to reflect light in order to reuse it in various applications such as solar cells and optical coatings. The combination of the mentioned materials allows for high reflectance at certain wavelengths. However, TiO₂ has better reflective properties than other materials due to its higher dielectric constant, which leads to a higher refractive index and results in absorbing the incoming light less. Fig. S1† shows the electric field distribution of the reflector structure for three various wavelengths of 360, 500, and 740 nm. It should be pointed out that a metasurface-based reflector can exhibit resonance, when a surface is located around one of the coupled resonance frequencies. Resonance in metasurface reflectors can be reached through various methods. Moreover, it is used to achieve total reflection or design tunable metasurface unit cells. Fig. S1† depicts the three different resonances that occurred in different wavelengths, which can cause the electric field to intensify around and above the nanodiscs, leading to an increase in the total absorption of the solar cell (refer to the ESI† for more details).

3.2. Solar cell design

It is worth noting that the refractive index of the upper layer of TiO₂ nanodiscs has an impact on the reflection spectra from this layer. To investigate this effect, the refractive index of this layer was swept from 1 to 4 with steps of 0.5, assuming the extinction coefficient to be equal to zero, which means no light absorption. By calculating the absorbance spectra from the surface of these layers, as provided in Fig. 5b, it was found that the refractive index range of 2 to 3 has the highest amount of light absorption in the structure. It is generally acknowledged that the enhancement in the reflectance spectra from the ETL can increase the light absorption in the active layer in an inverted solar cell. Plus, it prevents power loss by the metal contact. The finally proposed solar cell structure is presented in Fig. 5a. The structure is consisted of the following layers from top to bottom: ITO, HTL, active layer (MAPbI₃), ETL including TiO₂ nanodiscs and the other bulk material around these nanodiscs, Al, and SiO₂. ITO was used as a transparent contact for light to pass through. The band diagram of the proposed solar cell in the presence of the metasurface reflector is illustrated in Fig. 5c. In this structure, ZnO was used as the material around the TiO₂ nanodiscs, and CuSCN was used as the HTL. As plotted in Fig. 4f, the refractive index of the material around the TiO₂ nanodiscs did not have a significant effect on the reflectance spectra. Therefore, a material featuring a better carrier transfer should be used along with TiO₂, which can improve electron transfer, in addition to reflection improvement. According to Table 1, the MAPbI₃ layer has an electron affinity of 3.93 eV and a bandgap of 1.55 eV, which can result in efficient light absorption in the wavelength range of 300 to 800 nm. On the ETL side, TiO₂ and ZnO were used with the equal electron



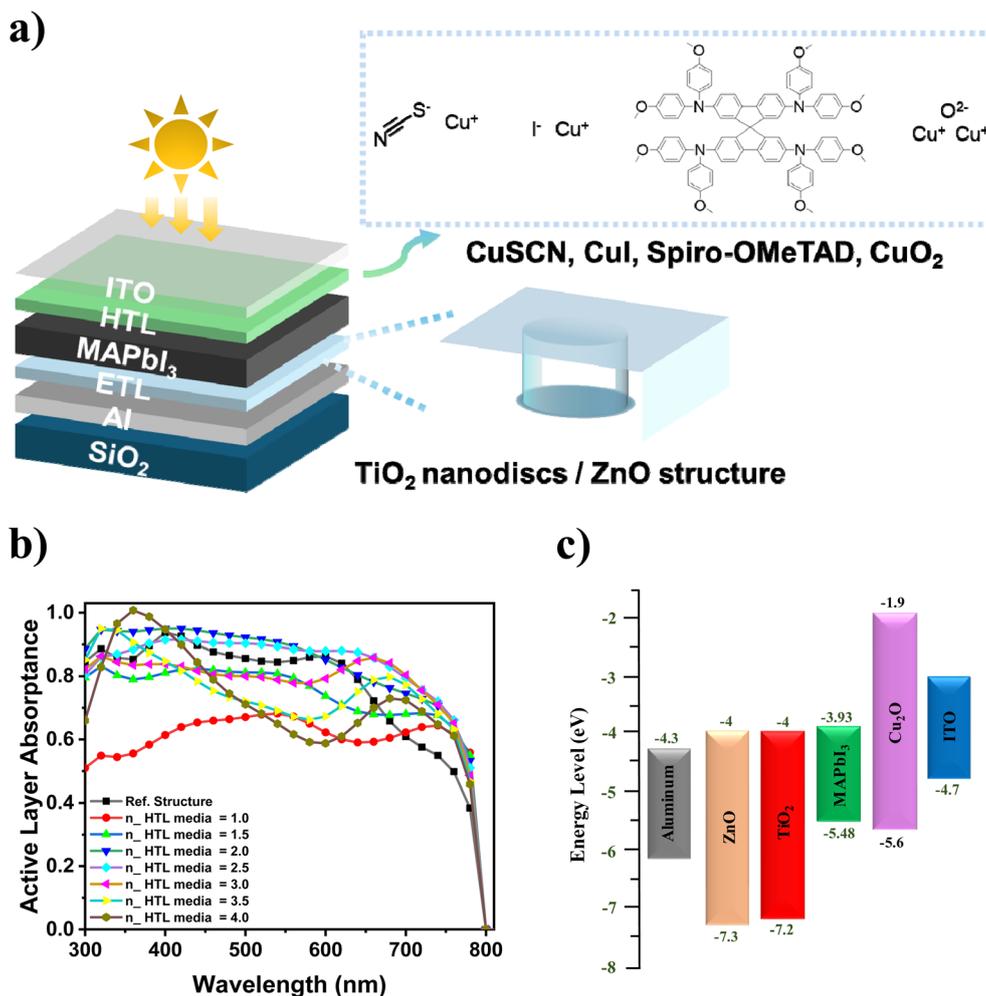


Fig. 5 (a) Schematic of the proposed solar cell with MAPbI₃ as the active layer. (b) The active layer (MAPbI₃) light absorbance spectrum with different HTL refractive indices between the range of 1 to 4 with steps of 0.5, while $k = 0$, and ZnO on TiO₂ nanodiscs as ETL. (c) Band diagram of the proposed structure, considering ZnO as the material around the TiO₂ nanodiscs, MAPbI₃ as active layer, and CuSCN as the HTL (all energy levels depicted in this diagram are based on the values listed in Table 1).

affinity value of -4 eV and the bandgap values of 3.2 eV and 3.3 eV, respectively. Due to the values of their conduction bands, these materials were able to transfer the electrons from the active layer into the Al contact with the work function of -4.3 eV. Furthermore, due to the values of their valence bands compared to MAPbI₃, they prevented the holes from entering the Al contact conveniently. On the HTL side, CuSCN was used with a bandgap of 3.4 eV, which does not absorb a noticeable portion of light in the visible region. Therefore, most of the incident light was able to reach the active layer. CuSCN with an electron affinity of 1.9 eV was considered as a suitable barrier to prevent the electrons from entering the transparent contact (ITO) and transferred the generated holes in the active layer into the ITO side with a valence band energy level difference of 0.18 eV with MAPbI₃.

It is obvious that applying different materials around TiO₂ nanodiscs can affect the amount of electron transfer. Although, according to the findings pertinent to Fig. 4f, different refractive indices did not have a significant effect on the amount of light reflection, the light transmission of these materials was

significant since they were not considered as absorbent materials. Therefore, various materials were considered for this purpose, including ZnO, AZO, CeO₂, and TiGO-5. The refractive indices and extinction coefficients of these materials in terms of wavelength in the range of 300 to 800 nm are shown in Figs. S2.†

As demonstrated in Fig. 6a, the absorbance spectra of the active layer (MAPbI₃) for different materials around the TiO₂ nanodiscs did not change significantly in the wavelength range of 300 to 800 nm and is almost the same for different materials, including ZnO, AZO, CeO₂, and TiGO-5. In order to generate electron-hole pairs, a photon with sufficient energy must be absorbed in the active layer. Therefore, the amount of carrier generation rate in different wavelengths, as depicted in Fig. 6b, is in accordance with the behavior of the absorbance diagram. Fig. 6c indicates the photocurrent density in terms of the voltage, where different values of the V_{oc} , can be explained by the different electrical characteristics of the materials. Bearing in mind that ZnO has a higher electron mobility compared to other materials, as stated in Table 1, as well as the 0.3 eV difference in its conduction band energy compared to the Al



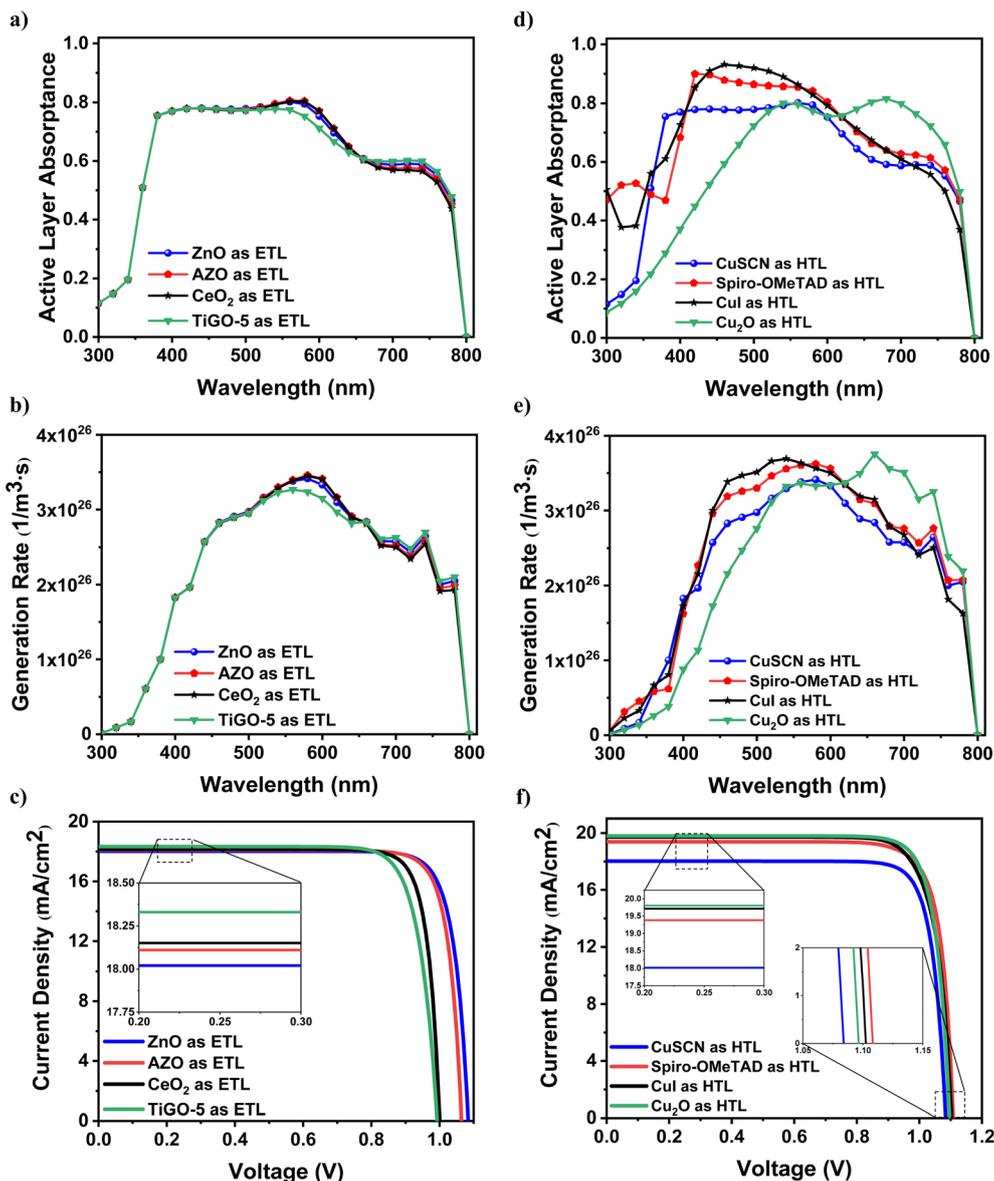


Fig. 6 The active layer (a) absorptance spectra, and (b) carrier generation rate in terms of wavelength, and (c) the photocurrent density in terms of voltage for the metasurface structure with different ETL materials around the TiO₂ nanodiscs. The active layer (d) absorptance spectra, and (e) carrier generation rate in terms of wavelength, and (f) the photocurrent density in terms of voltage for the metasurface structure with different HTL materials and ZnO as ETL.

work function, it resulted in a higher V_{oc} value. Likewise, according to Table 1, AZO shared almost the same V_{oc} value as ZnO due to their approximately equal electrical parameters. Despite ZnO and AZO, CeO₂ showed a relatively lower V_{oc} due to its lower conduction band than the Al work function. Similarly, TiGO-5 shared about the same V_{oc} value to CeO₂. This was because of the almost very close energy level of the TiGO-5 conduction band to the Al work function. The represented band diagram in Fig. S3† indicates that all the proposed materials as ETL can easily separate the electron from the active layer due to the lower energy level of their conduction band compared to the active layer's conduction band. Also, the holes were prevented to transfer into the metal contact in light of the same reason. It should be mentioned that the difference in V_{oc}

values is mainly contributed to the difference in the conduction band of these materials with the Al work function and also the electrical parameters of these materials according to Table 1.

Table 3 The values for photovoltaic parameter including J_{sc} , V_{oc} , PCE and FF for solar cells with TiO₂ nanodiscs embedded in various ETLs

ETL materials around the TiO ₂ nanodiscs	J_{sc} (mA cm ⁻²)	V_{oc} (V)	PCE (%)	FF
ZnO	18.02	1.08	16.40	0.84
AZO	18.11	1.06	16.25	0.84
CeO ₂	18.15	1.00	15.37	0.85
TiGO-5	18.33	0.99	14.83	0.82



Table 3 denotes the values of J_{sc} , V_{oc} , PCE, and FF of the solar cell, when the mentioned different materials around TiO_2 nanodiscs were investigated. According to Fig. 4f and S2(a),† the surface reflectance value is higher for the materials with lower refractive indices. Therefore, TiGO-5 represented the highest J_{sc} value of 18.33 mA cm^{-2} , owing to featuring the lowest refractive index among other ETL materials. However, as illustrated in Fig. 4f, the reflectance curves represented an insignificant difference for materials with different refractive indices. Also, as it was stated earlier, the structure with ZnO around TiO_2 nanodiscs presented the highest V_{oc} value of 1.08 V. Correspondingly, this structure showed the highest values of PCE and FF equal to 16.40% and 0.84, respectively.

According to Fig. 5b, the refractive indices of HTL materials appeared to have a striking impact on the amount of the active layer's light absorption. More particularly, materials with refractive indices between 2 to 3 showed the highest amount of light absorption by the active layer. Hence, it was concluded that the light interaction between the HTL and $MAPbI_3$ played an effective role in the amount of light reflection from the surface they shared and the proportion of the light that reached the active layer. In this regard, materials with different refractive indices were selected as HTL. These materials include CuSCN, Spiro-OMeTAD, CuI, and Cu_2O . Among these materials, the refractive indices of CuI and Cu_2O were in the target range. The refractive indices and extinction coefficients of these materials are shown in Fig. S4.† It was also observed that Cu_2O features a high extinction coefficient at lower wavelengths, which reduced the amount of active layer's light absorption at lower wavelengths.

The light reflection and absorption properties of a material can be described from its complex refractive index. The proportion of the transmitted light through the ITO and HTL layers and the reflected light from the layers below the active layer are considered as effective factors in the light absorption of the active layer. It was also distinguished earlier that employing materials with refractive indices in the range of 2 to 3 can increase the light reflection from the metasurface-based reflectors. Therefore, different HTL materials can cause different light absorption spectrum in the active layer. Fig. 6d shows the light absorbance curve of the active layer for different HTLs. Among the applied materials as HTL, CuI and Cu_2O , with refractive indices in the range of 2 to 3 in the visible light spectrum, the active layer experienced the highest light absorption in the wavelength range below 600 nm, and above 600 nm, respectively. According to Fig. 4b, the HTL medias with lower refractive indices resulted in lower light absorption spectrum in the active layer, and since the HTLs were assumed to absorb no light, this can be explained by the lower light reflection from the metasurface reflectors. Therefore, the HTL medias with lower refractive indices proved to effect on the light reflection from the metasurface reflectors. This is why the light absorption of the active layer was lower when CuSCN and Spiro-OMeTAD were employed as HTLs (Fig. 4d). Also, Fig. 6e shows the carrier generation rate of the active layer for different HTL materials. For the wavelengths below 600 nm, the light absorbance was the highest, when CuI was used as HTL, which is also evident in the carrier

generation rate diagram. For the wavelengths above 600 nm, the carrier generation rate was the highest when Cu_2O was used as HTL due to the significant light absorption of the active layer. In view of the fact that the carrier generation rate has a direct relation with J_{sc} , Fig. 6f, illustrates the highest J_{sc} values for the structures with Cu_2O and CuI as HTLs. Conversely, the structure with CuSCN as HTL represented the lowest J_{sc} value, which is due to its low refractive index and the noticeable light reflection reduction from the metasurface reflector. However, the V_{oc} values for different HTLs remained almost constant because the materials were all capable of effectively separating holes from the active layer and preventing the electrons from entering ITO contact. This can be explained through the energy band diagram of the HTLs, which are shown in Fig. S5.† The mentioned HTLs, also known as electron-blocking layers, including CuSCN, Spiro-OMeTAD, CuI, and Cu_2O , performed slightly different conduction bands equal to 2.03 eV, 1.73 eV, 1.83 eV, and 2.03 eV, respectively, which resulted in the moderately different V_{oc} . Additionally, the HTLs' electrical parameters, such as dielectric constant, electron and hole mobility, density of state for the conduction and valence band, electron and hole life-time and concentration of p-type doping, which are given in Table 1, were considered effective, too.

Table 4 presents the values of J_{sc} , V_{oc} , PCE, and FF for the proposed structure with different HTLs. As mentioned above, CuI and Cu_2O represented the highest values of almost the same J_{sc} equal to 19.71 mA cm^{-2} and 19.80 mA cm^{-2} , respectively. Although the V_{oc} values did not changed significantly, the Spiro-OMeTAD held the highest value equal to 1.11 V, due to its electrical characteristics. Nevertheless, the Cu_2O reported the highest value of PCE among other HTLs, equal to 18.15%, with the FF value of 0.83. Comparing the values in Tables 2 and 3, the PCE value was increased by 1.75%, when the HTL was changed from CuSCN to Cu_2O .

In order to enhance the light absorption, a second absorbent layer, $MoSe_2$, was introduced. $MoSe_2$ with a bandgap of 1.4 eV can absorb light in higher wavelengths compared to $MAPbI_3$. The $MoSe_2$ absorption range is from 300 to 920 nm, as provided in Fig. S6(b).† However, employing materials with higher refractive indices, around the TiO_2 nanodiscs did not have a critical impact on the light reflection from the surface of the metasurface reflector, as shown in Fig. 4f. Therefore, $MoSe_2$ can be used around the TiO_2 nanodiscs. The refractive indices and extinction coefficients of the other materials, including $MAPbI_3$, ITO, and TiO_2 , are depicted in Fig. S6.†

The planar reference structure, as shown in Fig. 7a, from top to bottom was composed of ITO as a transparent contact,

Table 4 The values for photovoltaic parameter including J_{sc} , V_{oc} , PCE and FF for solar cells with various HTLs

HTL materials	J_{sc} (mA cm^{-2})	V_{oc} (V)	PCE (%)	FF
CuSCN	18.02	1.08	16.40	0.84
Spiro-OMeTAD	19.38	1.11	17.74	0.82
CuI	19.71	1.10	17.73	0.80
Cu_2O	19.80	1.09	18.15	0.83



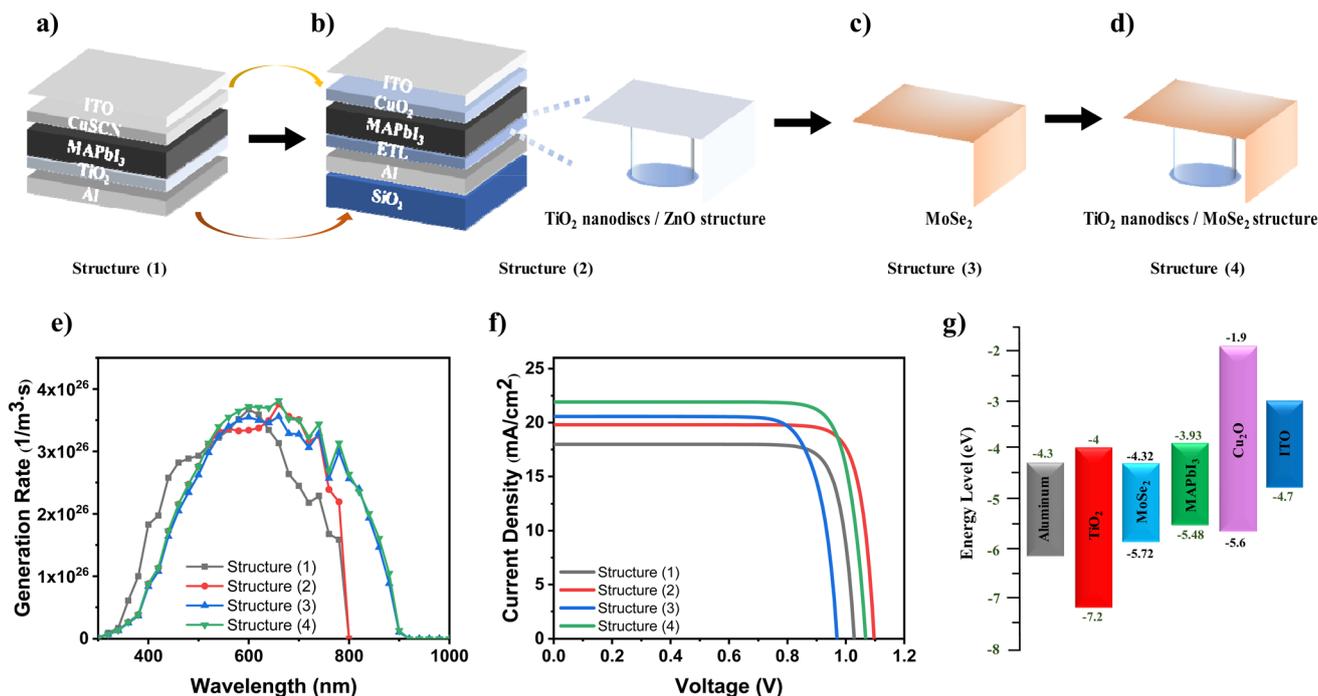


Fig. 7 The device architecture for (a) the planar reference structure (structure (1)), (b) the proposed metasurface-based structure with ZnO around the TiO₂ nanodiscs as ETL, and Cu₂O as HTL (structure (2)), (c) the planar structure with MoSe₂ as ETL and the second absorbent, and Cu₂O as HTL (structure (3)), and (d) the metasurface-based structure with MoSe₂ around TiO₂ nanodiscs as the second absorbent, and Cu₂O as HTL (structure (4)). The (e) carrier generation rate in terms of wavelength, and (f) current density in terms of voltage for the provided structures 1 to 4. (g) Band diagram of the structure 4 (all energy levels depicted in this diagram are based on the values listed in Table 1).

CuSCN as the HTL, MAPbI₃ as the active layer, TiO₂ as the ETL, and Al as the metal contact. Finally, four main structures were investigated, which can be explained as: (1) the planar reference structure (Fig. 7a), (2) the proposed metasurface-based structure with ZnO around the TiO₂ nanodiscs as ETL, and Cu₂O as HTL (Fig. 7b), (3) the planar structure with MoSe₂ as ETL and the second absorbent, and Cu₂O as HTL (Fig. 7c), and (4) the metasurface-based structure with MoSe₂ around TiO₂ nanodiscs as the second absorbent, and Cu₂O as HTL (Fig. 7d). The presence of TiO₂ nanodiscs can improve the efficiency of the system in two ways: (1) creating a metasurface reflector to increase the light reflection from the lower surface of the active layer to increase the J_{sc} value, and (2) improving the electron transfer through the nanodiscs, which increases V_{oc} . According to Fig. 7e, the carrier generation rate was higher in the structure (4) compared to the other structures owing to the presence of both the MoSe₂ as the second absorbent layer and the TiO₂ nanodiscs. The increase in the carrier generation rate, resulted in an increase in the J_{sc} value. According to Fig. 7f, which demonstrates the current density curves in terms of voltage, the J_{sc} value of the structure (4) is higher than the other structures due to the smaller bandgap of MoSe₂ compared to MAPbI₃, which absorbs light with wavelengths greater than 800 nm in the infrared region, as well as the higher light reflectance caused by the TiO₂ nanodiscs. Utilizing the TiO₂ nanodiscs along with ZnO in the structure (2) enhanced the carrier transfer rate, where both materials acted as hole blockers in the system.

Therefore, the V_{oc} value was increased slightly more than the other structures. On this account, structure (2) held a higher V_{oc} value compared to structures (3 and 4) in light of the lack of proper hole blocking in MoSe₂, which consequently increased the carrier recombination rate. Fig. 7g presents the energy band diagram related to the metasurface-based structure, where the TiO₂ nanodiscs were surrounded by the MoSe₂. In order to further discuss the lower V_{oc} of the structures (3 and 4), the electrical properties of MoSe₂ were studied. As indicated in Table 1, MoSe₂ has an electron affinity equal to 4.32 eV and a bandgap equal to 1.4 eV. As a consequence of the lower energy level of the conduction band in MoSe₂ compared to MAPbI₃, the generated electrons in MAPbI₃ layer were able to be transferred

Table 5 Comparison of the J_{sc} , V_{oc} , PCE and FF values for different structures of: (1) the planar reference structure, (2) the proposed metasurface-based structure with ZnO around the TiO₂ nanodiscs as ETL, and Cu₂O as HTL, (3) the planar structure with MoSe₂ as ETL and the second absorbent, and Cu₂O as HTL, and (4) the metasurface-based structure with MoSe₂ around TiO₂ nanodiscs as the second absorbent, and Cu₂O as HTL

Structure	J_{sc} (mA cm ⁻²)	V_{oc} (V)	PCE (%)	FF
1	17.98	1.03	15.33	0.84
2	19.80	1.09	18.15	0.83
3	20.56	0.97	15.75	0.78
4	21.91	1.07	19.17	0.82



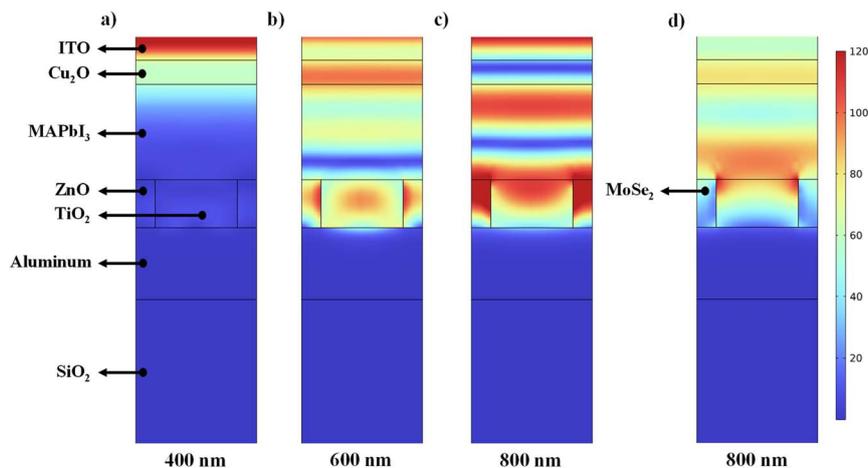


Fig. 8 Distribution of the electric field for the single junction metasurface reflector at different wavelengths of (a) 400 nm, (b) 600 nm, and (c) 800 nm. (d) Distribution of the electric field for the structure in the presence of MoSe₂ as the second absorbent layer at the wavelength of 800 nm.

into the metal contact, which resulted in lower V_{oc} values of the structures (3 and 4) compared to structure (2). However, the lower energy level of the valence band in MoSe₂ compared to MAPbI₃, was beneficial as it prevented the generated holes in MAPbI₃ from entering the metal contact and transferring the generated holes in MoSe₂ into the MAPbI₃.

Table 5 displays the values of the J_{sc} , V_{oc} , PCE, and FF for the structures (1 to 4). The J_{sc} value was enhanced from 19.80 mA cm⁻² to 21.91 mA cm⁻², when the MoSe₂ was applied into the structure (2) instead of the ZnO. However, as a result of the better electron transfer and hole blocking in the ZnO layer compared to MoSe₂, the value of the V_{oc} in structure (3) fell from 1.09 V to 1.07 V compared to structure (2). Accordingly, the total PCE value was increased from 18.15% to 19.17%, when MoSe₂ replaced the ZnO around the TiO₂ nanodiscs. Both structures (2 and 4) were improved, comparing their electrical characteristics including J_{sc} , V_{oc} , PCE, and FF to the planar structures (1 and 3) provided in Table 5, respectively. However, structure (4) was improved more significantly and reached the highest value of PCE equal to 19.17% as the champion structure among others. Also, the lack of optimal carrier separation in the system reduces the FF. By incorporating MoSe₂ into the structure, due to its narrow band diagram, some electrons generated in the MAPbI₃ layer enter the HTL. As illustrated in the band diagram of the structure (Fig. 7g), the electron entry into the HTL increases the series resistance, thereby reducing the FF. Additionally, some of these electrons recombine with the holes in the HTL, which leads to a decrease in the V_{oc} when using MoSe₂. However, considering the significant increase in J_{sc} compared to the decreases in V_{oc} and FF, it can be concluded that MoSe₂ is a better choice compared to ZnO.

Fig. 8 illustrates the electric field distribution in three different wavelengths of 400, 600, and 800 nm for the structure (2) and at the wavelength of 800 nm for the structure (4). At the wavelength of 400 nm, provided in Fig. 8a, only a small percentage of the electric field entered the active layer due to

light absorption in HTL (Cu₂O), where it absorbed approximately 30% of the input power, which is less than other materials. At the wavelength of 600 nm, according to Fig. 8b, the layers above the active layer (ITO and Cu₂O) showed no light absorption. In fact, Fig. S4(b) and S6(b)† indicate that most of the sunlight reached the active layer. About 75% of the incident light was absorbed in the active layer, together with the reflected light from the lower layers. Additionally, a strong electric field was formed around the TiO₂ nanodiscs due to their resonances. At the wavelength of 800 nm, according to Fig. 8c, the electric field reached the ETL. Due to the light reflection from the metasurface reflector and the resonance of the electric field between the nanodiscs, a strong field was formed in the ETL as well as the lower part of the active layer. Applying MoSe₂ as an absorbent material with the smaller bandgap than MAPbI₃ intensified the transmitted electric field through the active layer and increased the PCE. Fig. 8d shows the electric field distribution at the wavelength of 800 nm for structure (4). Moreover, a significant portion of the electric field around the TiO₂ nanodiscs that was depicted in Fig. 8c was absorbed by MoSe₂ in structure (4). According to Fig. 7e, the overall light absorbance and carrier generation of structure (4) in the presence of MoSe₂ at the wavelength of 800 nm seemed to have a significant portion, while for structure (2), this value is zero at this wavelength.

4. Conclusion

The efficiency of an inverted perovskite solar cell using a new perfect metasurface-based reflector with TiO₂ nanodiscs/MoSe₂ structure as ETL was enhanced through several optimizations. Initially, the perfect metasurface-based reflector structure (TiO₂/Al/SiO₂) was investigated and compared in different types of nanostructures, including: four different nanocubes and nanodiscs structures. It was explored that using TiO₂ as nanodiscs compared to TiO₂ as a bulk layer on an Al substrate



increased the reflectance from 78.08% to 97.95%. Subsequently, the light absorption in the active layer and the carrier transfer through the structure were improved. Additionally, different materials were investigated as the ETLs around the TiO₂ nanodiscs (ZnO, AZO, CeO₂, and TiGO-5), while other materials as an HTL above the active layer (CuSCN, Spiro-OMeTAD, CuI, and Cu₂O). The optimum structure consisting of ITO/Cu₂O/MAPbI₃/TiO₂ nanodiscs and ZnO/Al/SiO₂, obtained the values of the J_{sc} , V_{oc} , PCE, and FF equal to 19.80 mA cm⁻², 1.09 V, 18.15%, and 0.83, respectively. By applying MoSe₂ around TiO₂ nanodiscs with a smaller bandgap than the MAPbI₃, as the second absorbent in the infrared region, the electrical characteristics of the cell were improved compared to the planar reference invert solar cell structure (ITO/CuSCN/MAPbI₃/TiO₂/Al), where the values of the J_{sc} , V_{oc} , and PCE were increased from 17.98 mA cm⁻² to 21.91 mA cm⁻², 1.03 V to 1.07 V, and 15.33% to 19.17%, respectively.

Data availability

The data supporting this article have been included as part of the ESI.† Additional datasets generated and analyzed during the study, including optoelectronic simulations and characterization results, can be made available upon reasonable request by contacting the corresponding author.

Author contributions

Javad Maleki and Maryam Shahroostami: conceptualization, data curation, formal analysis, investigation, methodology, software, validation, visualization, writing – original draft, writing – review & editing. Siming Huang: investigation, validation, visualization, writing – review & editing. Mojtaba Abdi-Jalebi: conceptualization, data curation, formal analysis, funding acquisition, investigation, methodology, project administration, supervision, resources, validation, visualization, writing – review & editing.

Conflicts of interest

The authors declare no conflict of interest.

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