Chemical Science



EDGE ARTICLE

View Article Online
View Journal | View Issue



Cite this: Chem. Sci., 2025, 16, 17165

dll publication charges for this article have been paid for by the Royal Society of Chemistry

Received 26th June 2025 Accepted 16th August 2025

DOI: 10.1039/d5sc04716a

rsc.li/chemical-science

Ultraviolet photodissociation of methanethiol (CH₃SH): revealing an S(¹D) atom elimination channel

Yucheng Wu,^{ab} Shunyang Zhou,^a Zijie Luo, ^{bac} Shuaikang Yang,^a Zhenxing Li, ^{bac} Yongxin Dong,^c Wei Hua,^a Quan Shuai,^a Dongxu Dai,^a Michael N. R. Ashfold, ^{bac} Kaijun Yuan ^{babf} and Xueming Yang ^{afg}

We report time-sliced velocity map imaging studies of the methyl (CH₃) and electronically excited sulfur (S(1D)) fragments formed following the photoexcitation of jet-cooled CH₃SH molecules in the $2^1A'' \leftarrow \tilde{X}^{1}A'$ absorption band (i.e. at wavelengths in the range $190 \le \lambda \le 210$ nm). Analyses of images of CH₃ fragments in their $v_2 = 0$, 1 and 2 vibrational levels confirm the perpendicular parent transition dipole moment and prompt bond fission and show that the ground state SH(X) partners are formed with an inverted vibrational population distribution, peaking at v = 2 at the shortest excitation wavelengths investigated. Most of the photolysis photon energy above that required to break the C-S bond is partitioned into product translational energy. Primary S(¹D) products are observed on excitation at $\lambda \leq 204$ nm and their relative yield is deduced to increase quite steeply with decreasing wavelength, but quantum yield estimates are beyond the scope of the present work. Image analysis reveals that the CH₄ partners are formed with a highly inverted vibrational population distribution, largely concentrated in the ν_4 bending mode. A possible formation mechanism for the $S(^1D)$ + CH_4 products is suggested, based on frustrated C-S bond extension on the initially populated 21A" potential energy surface (PES) and re-collision between the embryonic CH3 and SH moieties in the extended region of conical intersection between the 2¹A" and 1¹A" PESs en route to the target products. Cutting edge electronic structure calculations along with complementary ab initio molecular dynamics studies should help validate or overturn this envisaged mechanism.

Introduction

Metabolization of dimethylsulfoniopropionate (DMSP) produced by phytoplankton and other marine organisms in seawater is recognised as the major biogenic source of dimethyl sulfide (DMS) on Earth. Once emitted into the atmosphere, DMS is rapidly oxidized to become an important precursor of sulfated aerosols and cloud condensation nuclei. Microbial action can also demethylate DMSP to methanethiol (CH₃SH), but the roles of CH₃SH in the oceans and the atmosphere are only now starting to be explored at levels of detail similar to that

electronic absorption by CH₃SH lies at ultraviolet (UV) wavelengths close to the short wavelength end of the solar spectrum that penetrates through the ozone layer, so CH₃SH destruction in the troposphere is by reaction (notably oxidation by OH radicals⁹), not photochemistry.

hitherto focused on DMS.3-8 The long wavelength end of the

More widely, sulfur is one of the more abundant elements in the universe. The S/H ratio in the solar photosphere is \sim 1.3 \times 10⁻⁵ (ref. 10), though the deduced quantity of S-containing species in dense clouds in the interstellar medium (ISM) is currently far lower than this quoted cosmic abundance.11 One plausible explanation for this apparent depletion is that much of the sulfur is incorporated and processed within dust grains and icy mantles. The relative abundance and mobility of hydrogen in an ice matrix suggests that most of the sulfur released by sputtering, thermal- or photo-desorption from such surfaces will be in the form of H2S,12 but similar mechanisms could give rise to CH3SH in the ISM.13 CH3SH was first detected tentatively,14 then definitively,15 towards the prolific high-mass star-forming region Sagittarius B2 close to the Galactic center, and has subsequently been observed in a range of locations, e.g. towards the organic hot-core G327.3-06,16 the cold core B1,17 the

[&]quot;State Key Laboratory of Chemical Reaction Dynamics and Dalian Coherent Light Source, Dalian Institute of Chemical Physics, Chinese Academy of Sciences, 457 Zhongshan Road, Dalian, 116023, China. E-mail: kjyuan@dicp.ac.cn

^bUniversity of Chinese Academy of Sciences, Beijing 100049, China

^cMarine Engineering College, Dalian Maritime University, Liaoning, 116026, China ^dInstitute of Advanced Light Source Facilities, Shenzhen, Guangdong, 518100, China ^cSchool of Chemistry, University of Bristol, Bristol, BS8 1TS, UK. E-mail: mike. ashfold@bristol.ac.uk

^fHefei National Laboratory, Hefei 230088, China

^{*}Department of Chemistry and Center for Advanced Light Source Research, College of Science, Southern University of Science and Technology, Shenzhen 518055, China

dense, warm part of a high-mass star-forming region in Orion, ¹⁸ in the vicinity of the solar type protostar IRAS 16293-2422 (ref. 19) and the prestellar core L1544, ²⁰ in both the cold envelope and the hot gas around Cyg X-N12 (ref. 21) and in the hot gas surrounding G328.2551-0.5321. ²² The tentative identification of DMS in the atmosphere surrounding exoplanet K2-18 (ref. 23) has fuelled debate regarding the possibility of life on planets other than Earth, but other recent observational ²⁴ and laboratory²⁵ studies serve as a reminder that DMS may also arise *via* abiotic routes.

Chemical Science

Photodissociation can be an important contributor to CH₃SH destruction in many regions of the ISM. As noted above, the electronic absorption of CH₃SH is concentrated in the UV spectral region,26-29 comprising a broad region of continuous absorption that peaks at $\lambda \sim 235$ nm and extends to beyond 300 nm and more localised, structured Rydberg features at shorter wavelengths converging towards the first ionisation potential (IP = 9.2922 ± 0.0007 eV, corresponding to an excitation wavelength $\lambda = 133.4$ nm).³⁰ Photoexcitation within the first absorption band promotes an electron from the highest occupied molecular orbital (HOMO, a non-bonding 3p orbital centred on the S atom) to an excited orbital with mixed Rydberg (4s)/valence (σ^*) character.^{31,32} The resulting 1¹A" excited state is dissociative with respect to both the H₃CS-H and H₃C-SH bond extension coordinates (henceforth R_{S-H} and R_{C-S} , respectively), though the 1¹A" potential energy surface (PES) displays a modest barrier at short C-S separations. Prompt S-H bond rupture is thus the dominant fate when exciting CH₃SH at long wavelengths. The resulting $CH_3S(\tilde{X})$ fragments are formed with modest vibrational excitation in the ν_3 (C-S stretch) mode, but most of the photon energy in excess of that required to break the S-H bond $(D_0(H_3CS-H) = 30250 \pm 100 \text{ cm}^{-1}, \text{ ref. } 33)$ is partitioned into product translational motion along axes that are preferentially perpendicular to the transition dipole moment (and thus to the polarization vector ε of the photolysis laser radiation). The $CH_3(\nu = 0)$ fragment images recorded when exciting CH₃SH in the range $220 \le \lambda \le 233$ nm also display near limiting perpendicular recoil anisotropy and show that most of the partner SH(X) fragments are formed in their ground vibrational state.35 The relative yield of CH3 + SH products increases on tuning to shorter wavelengths, 36 particularly when exciting the more localised $2^1A'' \leftarrow \tilde{X}^1A'$ band that dominates the absorption spectrum at wavelengths in the range $190 \le \lambda \le 215$ nm, such that the two channels have similar quantum yields when photolyzing at $\lambda = 193$ nm.³⁷

The present work involves a detailed study of photofragmentation pathways following excitation of CH_3SH within the $2^1A'' \leftarrow \tilde{X}^1A'$ absorption band. The parent absorption spectrum at wavelengths around this transition is shown in Fig. 1, along with arrows showing the species probed at each wavelength investigated. This transition is attributed to a 4p/4s(a') \leftarrow HOMO(a'') electron promotion. The adiabatic $2^1A''$ state is bound with respect to extending R_{S-H} or R_{C-S} , but both dissociation coordinates can be accessed via efficient nonadiabatic coupling at a region of conical intersection (CI) between the $2^1A''$ and $1^1A''$ PESs in the near vertical Franck–Condon region. 31,32,38,39 Prior studies at $\lambda = 208$ nm (ref. 33) and

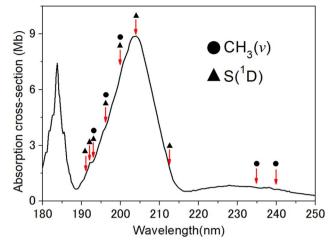


Fig. 1 Room temperature absorption spectrum of CH_3SH (after ref. 27 and 29), with the photolysis wavelengths investigated in the present work marked and labelled to show the species probed at each wavelength.

193 nm (ref. 34) revealed substantial vibrational excitation in the $CH_3S(\tilde{X})$ fragments from the former channel, most notably an obvious progression in the ν_3 (C–S) stretch mode. This points to some significant coupling between the C-S and S-H stretch motions during the dissociation process and the likely inadequacy of any one-dimensional representation of what is an inherently multi-dimensional problem. 38,39 CH3 fragments from the latter channel have been detected, generally with vibrational state specificity, following excitation at several wavelengths in the range $202 \le \lambda \le 210$ nm. ^{35,40-42} The SH(X) partners formed at these shorter excitation wavelengths populate a range of vibrational (ν) states, with distributions that peak at $\nu > 0$. The products from both channels again display preferential perpendicular recoil anisotropy, implying that non-adiabatic coupling via the 2¹A"/1¹A" CI is efficient and that dissociation occurs on a timescale comparable to, or faster than, the rotational period of the jet-cooled parent CH₃SH molecules.

Dissociation channels (1) and (2) are just two of many thermochemically allowed decay channels available to CH₃SH molecules when excited at $\lambda \sim 200$ nm. Focusing on spinallowed fragmentation processes yielded two products,

$$CH_3SH + hv \rightarrow CH_3 + SH$$
 $E_{th} = 25340 \pm 40 \text{ cm}^{-1}$ (1)

$$\rightarrow$$
 CH₃S + H $E_{th} = 30250 \pm 100 \text{ cm}^{-1}$, (2)

could be supplemented by one or more of the following fragmentations

$$CH_3SH + hv \rightarrow H_2CS + H_2$$
 $E_{th} = 11000 \pm 200 \text{ cm}^{-1} (3)$

$$\rightarrow$$
 S(¹D) + CH₄ $E_{\text{th}} = 27690 \pm 40 \text{ cm}^{-1}$ (4)

$$\rightarrow$$
 CH₂(\tilde{a}^1 A₁) + H₂S $E_{th} = 35260 \pm 40 \text{ cm}^{-1}$, (5)

and dissociations yielding products in higher excited states, e.g.

Edge Article Chemical Science

CH₃SH + hv
$$\rightarrow$$
 H₂CS(\tilde{A}) + H₂ $E_{th} = 27400 \pm 200 \text{ cm}^{-1}$ (6)
 \rightarrow S(1 S) + CH₄ $E_{th} = 40630 \pm 40 \text{ cm}^{-1}$, (7)

though we recognise that the energetic onsets for many of these later eliminations may be higher than these quoted thermochemical threshold (E_{th}) values as the process is likely to involve passage over an energy barrier. Inclusion of spin forbidden dissociations yielding, e.g., S(3P) atoms or triplet state CH2 radicals or excited triplet state H₂CS molecules (processes not currently included in this list) opens the range of possible dissociation processes yet further, and the onset of the lowest energy three-body dissociation channel (8) is only just above the energy of a $\lambda = 200$ nm photon.

$$CH_3SH + hv \rightarrow CH_3 + S(^3P) + H$$
 $E_{th} = 53730 \pm 40 \text{ cm}^{-1}$. (8)

The $D_0(H_3CS-H)$ value quoted for dissociation (2) is from ref. 33. All other threshold energies quoted here are derived using $\Delta_f H^0(0 \text{ K})$ values and uncertainties from ref. 43 and atomic term values from ref. 44, while the A-X term value for thioformaldehyde used in determining $E_{th}(6)$ is from Jacox.⁴⁵

H₂CS formation was predicted in early photolysis studies of CH₃SH at $\lambda = 185$ nm (ref. 46) and identified in later time-offlight mass spectrometry measurements following CH3SH photolysis at the ArF laser wavelength (193.3 nm).⁴⁷ S(³P) and S(¹D) atoms have been detected following 193.3 nm excitation of CH₃SH but, from measurements of the signal dependence upon photolysis laser intensity, these were deduced to arise from secondary photolysis of primary CH₃S and/or SH photoproducts.48 The present work extends high resolution imaging measurements of the CH₃(v) products down to the short wavelength end of the $2^1A'' \leftarrow \tilde{X}^1A'$ absorption band, *i.e.* $\lambda \ge 191$ nm (Fig. 1), thereby enabling new insights into the dissociation process (1), and provides first and definitive evidence for the operation of the rival dissociation channel (4) yielding $S(^{1}D)$ + CH₄ products.

Experimental 2.

The experiments employed a time-sliced velocity map ion imaging (TS-VMI) detection apparatus, details of which have been reported previously, 49-54 along with UV photolysis laser pulses in the range $190 \le \lambda \le 240$ nm and both UV and vacuum UV (VUV) probe laser pulses to detect, respectively, $CH_3(\nu)$ and S(¹D) photofragments. The CH₃SH sample was introduced into the source chamber as a pulsed supersonic beam (10% CH₃SH in Ar, Kylingas (99.5% purity)), where it was skimmed prior to entering (through a 2 mm hole in the first electrode) and propagating along the centre axis of the 23-plate ion optics assembly (IOA) within the spectrometer. The photolysis and probe laser beams intercepted the pulsed molecular beam at right angles, between the second and third plates of the IOA.

The requisite UV photolysis wavelengths were generated by frequency doubling or sum-frequency mixing using a table-top laser system. Light at $\lambda \ge 204$ nm (~ 1 mJ per pulse, pulse duration ~10 ns) was produced by frequency doubling the output of a dye laser (Sirah, PESC-G-24) pumped by the third harmonic (355 nm) output of a Nd:YAG laser (Continuum PL-9030). As such, the photolysis photon wavenumber in all cases was defined to sub-cm⁻¹ precision; the wavelengths used were chosen to be a whole number of nm. Light in the wavelength range $191 \le \lambda \le 199 \text{ nm}$ ($\sim 0.3 \text{ mJ}$ per pulse, pulse duration ~ 10 ns) was generated as the sum (i.e. $\omega_1 + \omega_2$) frequency output using a BBO crystal, with ω_1 set at the frequency corresponding to $\lambda_1 = 266$ nm (produced by frequency doubling the output of a 355 nm (the third harmonic output of a Continuum PL-9030 Nd:YAG laser) pumped dye laser (Sirah, PESC-G-24 operating at $\lambda \sim 532$ nm). The requisite ω_2 frequencies were produced using half of the second harmonic (532 nm) output from the same Continuum Nd:YAG laser to pump another dye laser (Sirah, PESC-G-24) generating photons with λ_2 wavelengths in the range 677–790 nm (\sim 8–10 mJ per pulse, duration of \sim 10 ns).

S(¹D) photoproducts were probed by Doppler scanning back and forth across the one photon absorption at $\lambda = 130.091$ nm, which populates the autoionizing 3p³(²D°)5s; ¹D° level of atomic sulfur. As previously reported, 50,55,56 these VUV photons were generated by four wave difference (i.e. $2\omega_3 - \omega_4$) frequency mixing of the frequency doubled output from one dye laser (at $\lambda_3 = 212.556$ nm) with the fundamental output of a second dye laser (at $\lambda_4 = 580.654$ nm) in a Kr/Ar gas mixture. The contrast between the structured and underlying continuum two-colour contributions to the S(1D) images was boosted by deliberately shifting the focus of the photolysis laser radiation 8 cm away from the interaction region, by translating the position of the 50 cm focal length lens used to focus this radiation. $CH_3(\tilde{X}, \nu)$ photoproducts were probed by two photon resonance enhanced multiphoton ionization (2 + 1 REMPI) using documented excitation wavelengths in the range $325 \le \lambda \le 333.5$ nm to sample population in the $\nu = 0$, $\nu_2 = 1$, $\nu_2 = 2$ and $\nu_1 = 1$ levels, ^{57,58} as described elsewhere.59 The requisite UV pulses (~2 mJ per pulse, duration of \sim 10 ns) were produced by frequency doubling the output of another (Sirah, PESC-G-24) dye laser operating at $\lambda \sim 650-667$ nm, which was pumped by the other half of the 532 nm output from the above Nd-YAG laser. The polarization vectors of the photolysis (ε_{phot}) and probe (ε_{probe}) laser radiation were both parallel to the front face of the detector, and the crossing angle between the photolysis and REMPI probe laser beam paths in the interaction region was \sim 7°, where the time delay between the respective pulses was in the range 10-20 ns.

The resulting S⁺/CH₃⁺ ions were accelerated through the rest of the IOA and detected with a dual microchannel plate assembly coupled with a P43 phosphor screen at the end of the 740 mm ion flight tube. The detector was time gated (15–20 ns) to select ions with m/z 32 (S⁺) or 15 (CH₃⁺) and to confirm that the signal was from the intended two-colour UV photolysis -REMPI/VUV probe scheme, and three images were recorded for each set of experimental conditions with: (i) both the photolysis and probe beams present in the interaction region; (ii) the photolysis beam present but the probe beam blocked, and (iii) the photolysis beam blocked and the probe beam present. For all two-colour images displayed in this article, the one-colour photoinduced background images recorded under conditions

(ii) and (iii) have already been subtracted from the image recorded under condition (i).

3. Results

3.1 CH₃ fragment imaging

Fig. 2 shows TS-VM images of the $CH_3(\nu=0)$ fragments obtained following photolysis of the jet-cooled CH_3SH sample at $\lambda=199,\ 196$ and 193 nm, at the high energy end of the $2^1A''\leftarrow \tilde{X}^1A'$ absorption band. Corresponding images of CH_3 products formed in the $\nu_2=1$ and $\nu_2=2$ levels (where ν_2 is the out-of-plane umbrella bending mode) are shown in Fig. S1 and S2 in the SI. In all cases, ε_{phot} is in the plane of the image, as shown by the double headed arrow included in panel (a) of each of these figures. Analysis of such images allows determination of (i) the total kinetic energy release, $P(E_T)$, distributions, given the

provisos of two body dissociation, momentum conservation and the partner fragment being SH, and (ii) the ($E_{\rm T}$ -dependent) recoil anisotropy. The latter is defined in terms of the anisotropy parameter, β , which is obtained by fitting the measured intensities to the expression $I(\theta) \propto [1 + \beta(P_2(\cos \theta))]$, where θ defines the angle of the recoil velocity vector relative to $\varepsilon_{\rm phot}$ and $P_2(\cos \theta)$ is the second Legendre polynomial. β takes limiting values of +2 and -1 in the case of axial recoil following one photon excitation via a transition dipole moment that lies, respectively, parallel and perpendicular to the breaking bond. The $P(E_{\rm T})$ and $\beta(E_{\rm T})$ plots (black and blue, respectively) are shown to the right of the corresponding images in Fig. 2, S1 and S2.

Each image shows a structured, anisotropic ring at a large radius and a secondary feature concentrated at the image centre. The former arises from the target process (1). The latter

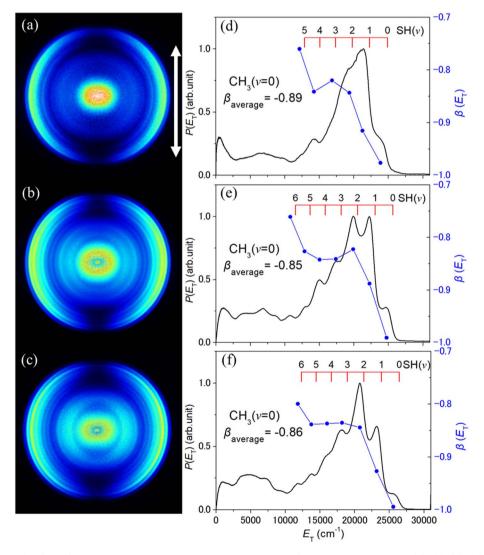


Fig. 2 TS-VM images of $CH_3(v=0)$ fragments formed by photolysis of jet-cooled CH_3SH molecules at $\lambda=(a)$ 199, (b) 196 and (c) 193 nm, with ϵ_{phot} aligned vertically in the plane of the image as indicated by the double headed arrow in (a). The corresponding $P(E_T)$ (black) and $\beta(E_T)$ (blue) distributions derived from each image are displayed in panels (d)–(f), with the relevant scales shown on, respectively, the left- and right-hand y-axes, along with the $\beta_{average}$ value determined over the indicated E_T range. The red combs above the respective $P(E_T)$ spectra indicate the maximum E_T value for the specified co-fragment formed with the probed $CH_3(v=0)$ species.

Edge Article Chemical Science

feature has been observed previously following excitation at other wavelengths in this region, 35,41,42 and is logically attributed to CH₃ fragments arising via unwanted probe-laser-induced photolysis of CH₃S($\tilde{\mathbf{X}}$) fragments formed via the rival primary fragmentation channel (2).

We consider this latter signal first. The thermochemical threshold energy for process (9)

$$CH_3S + h\nu \rightarrow CH_3 + S(^3P)$$
 $E_{th} = 24340 \pm 100 \text{ cm}^{-1}$ (9)

can be derived from the documented D₀(H₃CS-H) value³³ and the relevant $\Delta_f H^0(0 \text{ K})$ values given in ref. 43. The wavelength of the probe photons used for $CH_3(v = 0)$ detection ($\lambda =$ 333.46 nm, corresponding to a wavenumber of 29 988.6 cm⁻¹) falls within the CH₃S(Ã-X) absorption band. 60-62 The resulting CH₃S(Ã) fragments predissociate to CH₃ + S(³P) products (process (9)), most of which display preferential perpendicular recoil anisotropy. 63,64 CH₃(v = 0) fragments arising *via* one probe laser photon induced photolysis of primary $CH_3S(\tilde{X}, \nu = 0)$ fragments should thus appear with $E_{\rm T} \sim 5600~{\rm cm}^{-1}$. Several factors serve to blur such an analysis for the present experiments. The primary $CH_3S(X)$ fragments from $\lambda \sim 200$ nm photolysis of CH₃SH are formed with a range of kinetic energies, though momentum conservation ensures that most of any $E_{\rm T}$ released in dissociation process (2) is carried by the light H atom partner. More importantly, the $CH_3S(\tilde{X})$ fragments are distributed over a range of vibrational states, some with internal energies (E_{int}) approaching the limit allowed by energy conservation. 33,34 The relative absorption cross-sections of $CH_3S(\tilde{X})$ fragments in different ν levels at the probe wavelength of interest are not known, nor are many details of the dissociation dynamics of different $CH_3S(\tilde{A}, \nu)$ fragments. Conceivably, however, the probe-laser induced photodissociation of primary $CH_3S(\tilde{X})$ fragments in high E_{int} states could yield $CH_3 + S(^3P)$ products with $E_{\rm T} \sim 15\,000~{\rm cm}^{-1}$ or more, which would overlap with features attributable to $CH_3 + SH(X, higher \nu)$ products (primary process (1)) in the $P(E_T)$ spectra shown in Fig. 2, S1 and S2. The CH₃ fragments attributed to probe-laser induced photolysis of the ensemble of primary $CH_3S(X)$ products show some preferential perpendicular recoil anisotropy (Fig. S1).

We now focus on the feature at a larger radius in these $CH_3(\nu)$ images. As shown by the combs superposed above the $P(E_T)$ distributions derived from the $CH_3(\nu = 0)$, $CH_3(\nu_2 = 1)$ and $CH_3(v_2 = 2)$ images shown in Fig. 2, S1 and S2, the partner SH(X)fragments are formed in levels with ν up to \sim 6. For any set of $CH_3(\nu; \lambda)$ images, the onset at high E_T associated with the formation of SH(v = 0) co-fragments shows the expected shift to higher $E_{\rm T}$ as the photolysis wavelength is reduced. More strikingly, the most populated level in the SH(v) population distribution, P(v), clearly increases from v = 1 to v = 2 as the photolysis wavelength is decreased from 199 nm to 193 nm, but the various P(v) distributions appear to be relatively insensitive to the probed $CH_3(\nu)$ level. No signal attributable to $CH_3(\nu_1 = 1)$ fragments was observed when excited at the probe wavelength recommended58 for their detection. The SH product vibrational energy distributions can be put on a more quantitative footing by fitting the higher- E_T part of the $P(E_T)$ distributions using a set of suitably positioned Gaussian functions to represent the different SH(ν) levels. An illustrative decomposition of the data obtained at $\lambda=196$ nm is shown in Fig. S3. Of note, the signal attributable to the secondary photolysis of primary CH₃S fragments is likely to extend to $E_{\rm T}$ values attributed to primary CH₃ + SH(X, $\nu=5$ and, particularly, 6) products, so the present analysis almost certainly overestimates the relative yields of these SH(X, high ν) levels. The overlapping contributions from these two CH₃ radical sources may be responsible for the apparent reduction in the recoil anisotropy of the CH₃ + SH(X, higher ν) products.

Fig. 3(b) shows the SH P(v) distributions deduced from analysing the CH₃(v=0) images obtained at $\lambda=199$, 196 and 193 nm. These data are supplemented by the corresponding P(v) distributions reported previously following photolysis at $\lambda=204$ nm (ref. 41) and 210 nm.⁴² SH P(v) distributions determined when excited at longer wavelengths (in the $1^1A'' \leftarrow \tilde{X}^1A'$ absorption continuum) are shown in Fig. 3(a). The $\lambda=235$ and 240 nm data are derived from $P(E_T)$ spectra recorded as part of

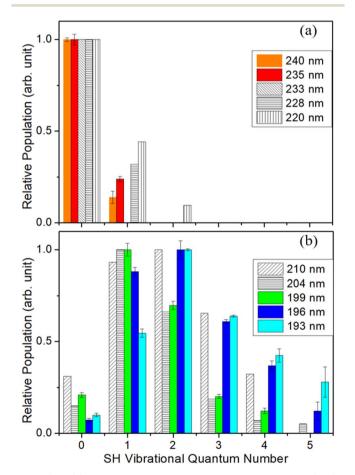


Fig. 3 SH P(v) distributions from photolysis of a jet-cooled CH₃SH sample deduced from analysing CH₃(v=0) images measured at $\lambda=$ (a) 240 and 235 nm (present work) and 233, 228 and 220 nm (from ref. 35) and (b) 210 nm (ref. 42), 204 nm (ref. 41), 199, 196 and 193 nm (present work), in each case normalized such that the most populated level is plotted with unit intensity. The error bars represent the standard error in the simulation, but do not allow for the likely overestimation of the relative population in the highest v levels (see Section 3.1).

the present study (shown in Fig. S4 in the SI), while the SH $P(\nu)$ distributions from CH₃SH photolysis at $\lambda=220$, 228 and 233 nm are from ref. 35. These long wavelength data are peripheral to the primary focus of the present paper but are included for completeness. All serve to reinforce previous conclusions that (i) the SH(X) fragments formed by photolysis within the $1^1A'' \leftarrow \tilde{X}^1A'$ continuum carry modest vibrational (and modest rotational) excitation, whereas those formed by photolysis in the $2^1A'' \leftarrow \tilde{X}^1A'$ band show inverted $P(\nu)$ distributions, (ii) the energy disposal in the SH(X) fragments formed at any given wavelength is relatively insensitive to any vibrational excitation in the probed CH₃(ν_2) fragment, and (iii) most of the photon energy in excess of that required to break the C–S bond is partitioned into product translational motion.

3.2 S(¹D) fragment imaging

Fig. 4(a) and (b) show TS-VM images of the $S(^1D)$ fragments formed when photolyzing jet-cooled CH_3SH molecules at $\lambda = 192$ and 191 nm. As noted in the Experimental section, the focus of the photolysis laser radiation was translated 8 cm from the interaction region to maximise the contrast between the structured signal of most interest from the underlying photolysis-laser-induced background signal. Even a cursory examination of these images shows that the $S(^1D)$ images recoil along axes

perpendicular to ε_{phot} . The left-hand panels in Fig. 5 show the corresponding S(1D) fragment images measured under the same experimental conditions when exciting at $\lambda = (a)$ 204, (b) 199, (c) 196 and (d) 193 nm. No fine structure was evident in the corresponding image measured at $\lambda = 212$ nm, as shown in Fig. S5. The panel to the right of each image shows the corresponding $P(E_T)$ distributions derived from (i) the two-colour image measured with both photolysis and probe lasers present (red trace), (ii) the one-colour probe laser only image (black trace), which peaks at low $E_{\rm T}$, and (iii) the difference (i.e. red minus black) distribution attributable to the real two-colour signal (blue trace). The $\beta(E_{\rm T})$ distribution derived from this real two-colour signal is also shown in each right-hand panel (blue points, joined by a dashed line) and the $\beta(E_T)$ distribution for the one colour probe only signal is also included in Fig. 4, 5 and S5 (black dots).

The two-colour images measured at $\lambda \leq 204$ nm contain (at least) three components. The structure attributable to the S(^1D) + CH_4 products formed *via* channel (4) is the feature of greatest novelty but we choose to interpret the overlapping unstructured contributions first. The probe-only signal (black traces in Fig. 4, 5 and S5) is attributable to the probe ($\lambda = 130.091$ nm) laser induced three body dissociation

$$CH_3SH + hv \rightarrow CH_3 + S(^1D) + H$$
 $E_{th} = 62970 \pm 40 \text{ cm}^{-1}, (10)$

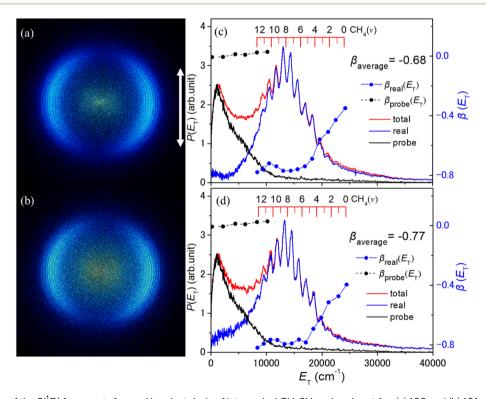


Fig. 4 TS-VM images of the $S^{(1)}$ D) fragments formed by photolysis of jet-cooled CH₃SH molecules at $\lambda = (a)$ 192 and (b) 191 nm, with ϵ_{phot} aligned vertically in the plane of the image as indicated by the double headed arrow in (a). Panels (c) and (d) show the $P(E_T)$ distributions derived from this two-colour image (red trace), from the 130.091 nm probe-laser only image (black trace) and the 'real' pump-probe two-colour $P(E_T)$ distribution obtained from the difference (blue trace), referenced to the left-hand y-axis scale. The black and blue dots show the $\beta(E_T)$ distributions derived from the corresponding traces, referenced to the right-hand y-axis scale. The $\beta_{average}$ value for the signal from photolysis laser induced dissociation determined over the indicated E_T range is also included in the inset. The red combs in panels (c) and (d) show the maximum E_T values for different $S(^1D) + CH_4(v_4)$ product channels.

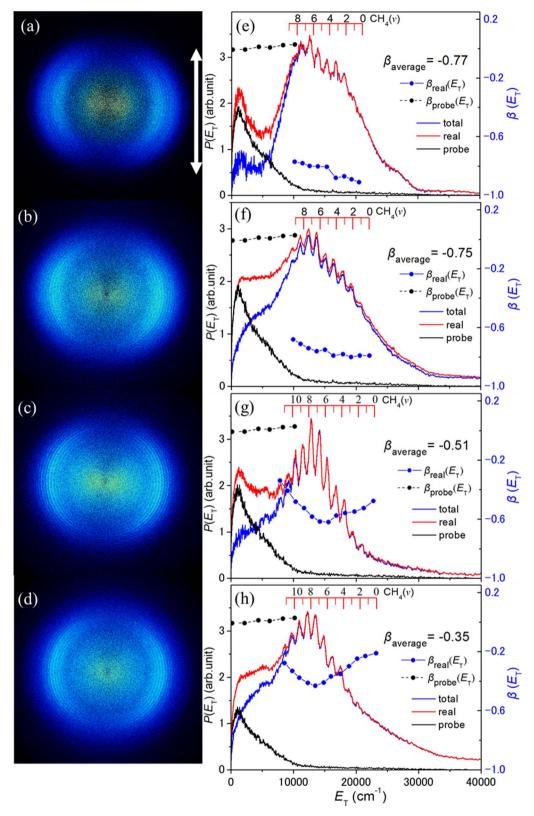


Fig. 5 TS-VM images of the S(1 D) fragments formed by photolysis of jet-cooled CH $_3$ SH molecules at λ = (a) 204, (b) 199, (c) 196 and (d) 193 nm, with ϵ_{phot} aligned vertically in the plane of the image as indicated by the double headed arrow in (a). Panels (e)–(h) show the $P(E_T)$ distributions derived from each two-colour image (red trace), from the 130.091 nm probe-laser only image (black trace) and the 'real' pump-probe two-colour $P(E_T)$ distribution obtained from the difference (blue trace), referenced to the left-hand y-axis scale. The black and blue dots show the $\beta(E_T)$ distributions derived from the corresponding traces, referenced to the right-hand y-axis scale. The $\beta_{average}$ value for the signal from photolysis laser induced dissociation determined over the indicated E_T range is also included in the inset. The red combs in panels (c) and (d) show the maximum E_T values for different S(1 D) + CH $_4(v_4)$ product channels.

products from which, by energy conservation arguments, they could appear with $E_{\rm T} \leq 13\,900~{\rm cm}^{-1}$. We note that the $E_{\rm T}$ values in the $P(E_{\rm T})$ traces displayed in Fig. 4, 5 and S5 are derived assuming two body dissociation and that the partner to the imaged S atom has m=16, and so is not strictly accurate for products arising via process (10). But the black trace in each plot sensibly declines to the baseline by this predicted upper limit $E_{\rm T}$ value and peaks at much lower translational energies, implying significant internal excitation in the CH₃ products formed at this VUV wavelength.

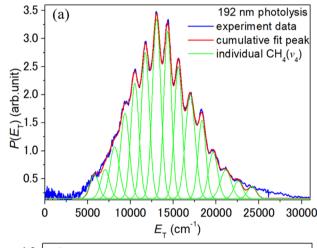
The second unstructured contribution lies under the structured features in the real two-colour signal (the blue traces in the displayed $P(E_T)$ spectra) and extends to E_T values much higher than that accessible via one photolysis photon induced formation of $S(^{1}D) + CH_{4}$ products. This contribution dominates the spectra recorded at longer wavelengths and is likely responsible for all the two-colour signal in the image measured at $\lambda = 212$ nm (Fig. S5). As Fig. S5 shows, the high $E_{\rm T}$ limit of this signal is sensibly consistent with that expected from resonance enhanced two-photolysis-photon induced dissociation of CH3SH via channel (10), i.e. $E_T \le 31\,370 \text{ cm}^{-1}$ at $\lambda = 212 \text{ nm}$. Inspection of the corresponding two-colour $P(E_T)$ spectra (blue traces) obtained at shorter wavelengths shows this high energy tail stretching out to progressively high $E_{\rm T}$ values consistent with the thermochemical prediction that two-photolysis-photon induced dissociation could yield products with $E_T \le 35\,070\,\mathrm{cm}^{-1}$ at $\lambda = 204\,\mathrm{nm}$ and $E_{\rm T} \le 37\,530~{\rm cm}^{-1}$ at $\lambda = 199~{\rm nm}$ (Fig. 5). But inspection of Fig. 4 also suggests that this two-photon contribution is comparatively minor in the $S(^{1}D)$ product images measured at λ = 191 or 192 nm, i.e. that process (4) – which we will now show to be responsible for the structure in these $P(E_T)$ spectra – contributes only weakly at $\lambda = 204$ nm but its relative importance increases steeply on tuning to shorter excitation wavelengths within the $2^1A'' \leftarrow \tilde{X}^1A'$ absorption band.

From here on, we focus on the structured signal attributed to $S(^1D) + CH_4$ product formation. CH_4 has four normal modes of vibration, which can be labelled by irreducible representations of the T_d point group, according to the symmetry of the associated normal coordinates. These are the symmetric (ν_1) and asymmetric (ν_3) stretching modes, with respective degeneracies of 1 and 3, and the ν_2 and ν_4 bending modes, with respective degeneracies of 2 and 3. The wavenumbers of the two stretch fundamentals are similar. The wavenumbers of the ν_2 and ν_4 fundamentals are also similar, and both are about half that of ν_1 and ν_3 . This leads to a well-defined polyad structure, with each polyad P_n defined by the integer n, where

$$n = 2(v_1 + v_3) + v_2 + v_4. (11)$$

The v_i in eqn (11) are the vibrational quantum numbers, and each (v_1, v_2, v_3, v_4) set defines a vibrational level, most of which – because of the degeneracies of the v_2 , v_3 and v_4 modes – split into multiple sub-levels. The levels within each polyad span a range of energies, the spread of which increases quite rapidly with $n.^{65,66}$ For future reference, we also note that the lowest energy levels within any P_n polyad are those associated with the nv_4 levels. 65,66

The observed peaks align fairly well with the $E_{T(max)}$ values predicted given the respective UV photon energies, $E_{th}(4)$ and the predicted wavenumber of the centre of gravity of each polyad with $n \le 9$,65 but the alignment between the ticks and the peaks degrades with increasing n. Better alignment is achieved by assuming a smaller peak spacing of $\sim 1330 \text{ cm}^{-1}$, shown by the combs superposed above the data in Fig. 4(c) and (d) and in the right hand panels in Fig. 5. This hints that the observed peaks are dominated by a sub-set of levels involving multiples of the ν_4 bending mode. Such an interpretation also fits better with the apparent near-constancy of the peak widths, the narrowness of which also implies that the CH4 fragments must be formed with only modest rotational excitation. As Fig. 6 shows, the vibrational state population distribution, $P(v_4)$, in the CH₄ fragments formed when photolyzing at $\lambda = 192$ nm (Fig. 4(a)) is highly inverted, peaking at $v_4 = 8$ and implying an average vibrational energy content in the CH4 fragments of \sim 10 000 cm⁻¹. The increased relative contribution to the S(1 D)



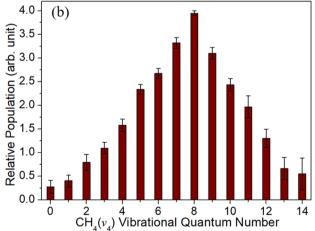


Fig. 6 (a) Illustrative decomposition of the $P(E_T)$ spectrum obtained from analysis of the TS-VM image of the $S(^1D)$ fragments formed by photolysis of jet-cooled CH $_3$ SH molecules at $\lambda=192$ nm using Gaussian functions to describe the relative populations in the partner CH $_4(v_4)$ levels. (b) Histogram illustrating the inverted CH $_4(v_4)$ vibrational population distribution so derived. The error bars represent the standard error in the simulation.

product yield from two-photolysis-photon induced dissociations in images recorded at longer wavelengths makes it hard to define any λ -dependence of the $P(\nu_4)$ distribution. The possible UV photofragmentation dynamics responsible for the deduced $S(^1D) + CH_4$ product state distribution is discussed along with the previously recognised dissociation processes (1) and (2) in the following section.

4. Discussion

The observation and characterization of the hitherto unreported S(1D) + CH₄ product channel (4) following photoexcitation of CH₃SH in the wavelength range $191 \le \lambda \le 204$ nm are the most striking findings from the present work. In what follows, we consider possible mechanisms for this process effectively an elimination of the central atom in an ABC framework to yield B + AC products. Reports of such photoinduced eliminations are becoming increasing commonplace,52 having recently been recognised in, for example: H2O (yielding H₂ + O(¹S) products),⁶⁷ H₂S (yielding H₂ molecules together with a sulfur atom, in the ¹D or ¹S excited state), ^{56,68} CO₂,⁶⁹ OCS,⁷⁰ CS₂,⁵¹ (yielding C(³P) atoms in each case), SO₂ (yielding S(1D) products)71 and, in a molecule of more similar size to that of present interest, CH_3NH_2 (yielding $NH(X^3\Sigma^-)$ + CH₄ products).^{72,73} These prior studies suggest various (molecule specific) mechanisms for such photo-induced eliminations, amongst which are a class of frustrated dissociations sometimes described as 'roaming'.74 In such cases, the topography of the excited state potential accessed by photoexcitation drives an initial bond extension, but the available energy is insufficient to enable complete bond fission. The partner moieties are thus drawn back towards one another and undergo what is effectively an intramolecular collision prior to reseparating as (potentially) an alternative product pair. The available data suggest that such a mechanism might account for the S(1D) + CH₄ fragment channel identified in the present work.

First, we reiterate that the operation of rival primary C-S and S-H bond fission processes following photoexcitation to the $2^{1}A''$ state at $\lambda \sim 200$ nm is well-established. The adiabatic potential of the $2^1A''$ state is bound in both R_{C-S} and R_{S-H} ; dissociation occurs on the 11A" PES, after non-adiabatic coupling in an extended region of CI between the 2¹A" and 1¹A" potentials.^{38,39} The present work reveals some vibrational excitation in the SH(X) fragments formed via channel (1) when excited in the 190 $\leq \lambda \leq$ 200 nm range, but also suggests that – as at longer excitation wavelengths - most of the photon energy in excess of that required to break the C-S bond is partitioned into $CH_3(\tilde{X}) + SH(X)$ product translational motion. The $CH_3S(\tilde{X})$ products arising via the rival channel (2) following excitation to the 21A" state and subsequent non-adiabatic coupling to the 1¹A" PES are distributed over many vibrational levels, with very obvious activity in the ν_3 (C-S) stretch mode, but, again, even the most internally excited H + $CH_3S(\tilde{X})$ products recoil with $E_{\rm T} \sim 1$ eV.^{33,34} Neither of these pathways obviously satisfy the picture of a frustrated dissociation and recollision.

Several other points merit note, however. The relative importance of the C-S bond fission process (1) increases markedly once exciting in the $2^1A'' \leftarrow \tilde{X}^1A'$ absorption band, suggesting that a larger fraction of the molecules photoexcited at these shorter wavelengths undergo initial C-S bond extension. This accords with the observed ν_5 (C–S) stretch progression in the weak emission spectrum from the parent CH₃SH molecules following excitation at $\lambda = 193.3$ nm.³⁶ Additionally, the $2^{1}A''$ PES is adiabatically bound, correlating to $CH_{3}(\tilde{X}) + SH(A)$ fragments when extending R_{C-S} (and to H + CH₃S(\tilde{A}) fragments when extending R_{S-H}). Thus, it is tempting to speculate that the observed S(1D) + CH₄ products might arise from a fraction of the photoexcited molecules that initially distort by extending R_{C-S} , avoid non-adiabatic coupling to the 11A" PES during this bond extension phase, reach the maximum R_{C-S} value consistent with energy conservation on the 2¹A" PES, then re-compress to enable a re-encounter between the CH3 and SH moieties and, potentially, an exothermic reaction on to $S(^{1}D) + CH_{4}$ products

Further mechanistic speculation at this stage is unwarranted, but it is very much hoped that the present data will inspire new high-level quantum chemical calculations of the 21A" and 11A" PESs and the non-adiabatic coupling between them (and, conceivably, to the \tilde{X}^1A' PES), together with ab initio molecular dynamics calculations designed to explore further the dynamics displayed by the products of dissociation processes (1), (2), (4) and, potentially, (3) and (5) - all of which could plausibly arise when exciting CH₃SH at $\lambda \sim 200$ nm. Future studies (experimental and/or theoretical) might also usefully search for processes yielding S(3P) atom products, estimate product quantum yields (branching fractions), and explore how these vary with the excitation wavelength. Extension of such contemporary photofragmentation studies to the oxygen analogue, methanol, the simplest and most abundant interstellar complex organic molecule observed in warm and cold environments, 75,76 should be an obvious priority, but searching for singlet products (e.g. $NH(a^{1}\Delta) + CH_{4}$) from photolysis of methylamine could also be very rewarding.

Author contributions

K. J. Y. and X. M. Y. supervised the research. K. J. Y. conceived the research. M. N. R. A. and K. J. Y. designed the experiments. Y. C. W., S. Y. Z., Z. J. L., S. K. Y., Z. X. L., Y. X. D., W. H. and Q. S. performed the experiments. Y. C. W., D. X. D., M. N. R. A., K. J. Y and X. M. Y. analysed the data. Y. C. W., K. J. Y. and M. N. R. A. wrote the manuscript. All authors discussed the results and commented on the manuscript.

Conflicts of interest

There are no conflicts to declare.

Data availability

The data supporting this study are available within the main text and the SI. See DOI: https://doi.org/10.1039/d5sc04716a.

Acknowledgements

The experimental work was supported by the National Natural Science Foundation of China (Grant No. 22241304, 22225303, 22403091, and 22173100), the Major Program of National Natural Science Foundation of China (No. 42494850 and 42494853), the National Natural Science Foundation of China (NSFC Center for Chemical Dynamics (Grant No. 22288201)), the Strategic Priority Research Program of the Chinese Academy of Sciences (Grant No. XDB0970000 and XDB0970200), the Innovation Program for Quantum Science and Technology (Grant No. 2021ZD0303304) and the Liaoning Revitalization Talents Program (Grant No. XLYC2402046). X. Yang also thanks the Shenzhen Science and Technology Program (Grant No. ZDSYS20200421111001787), and Z. Li thanks the Guangdong Science and Technology Program (Grant No. 2025A1515012671).

References

- 1 R. Jackson and A. Gabric, Microorganisms, 2022, 10, 1581.
- 2 R. P. Kiene, Mar. Chem., 2020, 54, 69-83.
- 3 C. L. Lee and P. Brimblecombe, *Earth-Sci. Rev.*, 2016, **160**, 1–18.
- 4 S. J. Lawson, C. S. Law, M. J. Harvey, T. G. Bell, C. F. Walker, W. J. de Bruyn and E. S. Saltzman, *Atmos. Chem. Phys.*, 2020, 20, 3061–3078.
- 5 G. A. Novak, D. B. Kilgour, C. M. Jernigan, M. P. Vermeuel and T. H. Bertram, *Atmos. Chem. Phys.*, 2022, 22, 6309–6325.
- 6 V. Gros, B. Bonsang, R. Sarda-Estève, A. Nikolopoulos, K. Metfies, M. Wietz and I. Peeken, *Biogeosciences*, 2023, 20, 851–867.
- 7 C. Wohl, G. L. Forster, P. M. Edwards, P. Suntharalingam and D. E. Oram, *Geophys. Res. Lett.*, 2025, **52**, e2025GL114929.
- 8 M. Rocco, E. Dunne, R. Salignat, A. Saint-Macary, M. Peltola, T. Barthelmeβ, G. Chamba, N. Barr, K. Safi, A. Marriner, S. Deppeler, C. Rose, J. Uitz, J. Harnwell, A. Engel, A. Colomb, A. Saiz-Lopez, M. J. Harvey, C. S. Law and K. Sellegri, J. Geophys. Res.: Atmos., 2025, 130, e2024JD041072.
- 9 G. S. Tyndall and A. R. Ravishankara, *Int. J. Chem. Kinet.*, 1991, 23, 483–527.
- 10 M. Asplund, N. Grevesse, A. J. Sauval and P. Scott, *Annu. Rev. Astron. Astrophys.*, 2009, 47, 481–522.
- 11 D. E. Anderson, E. A. Bergin, S. Maret and V. Wakelam, *Astrophys. J.*, 2013, 779, 141.
- 12 T. H. G. Vidal, J.-C. Loison, A. Y. Jaziri, M. Ruaud, P. Gratier and V. Wakelam, *Mon. Not. R. Astron. Soc.*, 2017, **469**, 435–447.
- 13 S. Narayanan, E. L. Piacentino, K. I. Öberg and M. Rajappan, *Astrophys. J.*, 2025, **986**, 10.
- 14 B. E. Turner, Astrophys. J., 1977, 213, L75.
- 15 R. A. Linke, M. A. Frerking and P. Thaddeus, *Astrophys. J.*, 1979, **234**, L139.
- 16 E. Gibb, A. Nummelin, W. M. Irvine, D. C. B. Whittet and P. Bergman, *Astrophys. J.*, 2000, 545, 309.

- 17 G. Cernicharo, N. Marcelino, E. Roueff, M. Gerin, A. Jiménez-Escobar and G. M. Muñoz, *Astrophys. J.*, 2012, **759**, L43.
- 18 L. Kolesniková, B. Tercero, J. Cernicharo, J. L. Alonso, A. M. Daly, B. P. Gordon and S. T. Shipman, *Astrophys. J.*, 2014, 784, L7.
- 19 L. Majumdar, P. Gratier, T. Vidal, V. Wakelam, J.-C. Loison, K. M. Hickson and E. Caux, Mon. Not. R. Astron. Soc., 2016, 458, 1859.
- 20 C. Vastel, D. Quénard, R. Le Gal, V. Wakelam, A. Andrianasolo, P. Caselli, T. Vidal, C. Ceccarelli, B. Lefloch and R. Bachiller, Mon. Not. R. Astron. Soc., 2018, 478, 5514.
- 21 M. el Akel, L. A. Kristensen, R. Le Gal, S. J. van der Walt, R. L. Pitts and F. Dulieu, *Astron. Astrophys.*, 2022, **659**, A100.
- 22 L. Bouscasse, T. Csengeri, A. Belloche, F. Wyrowski, S. Bontemps, R. Güsten and K. M. Menten, Astron. Astrophys., 2022, 662, A32.
- 23 N. Madhusudhan, S. Constatinou, M. Holmberg, S. Sarkar, A. A. A. Piette and J. I. Moses, *Astrophys. J. Lett.*, 2025, 983, L40.
- 24 N. Hänni, K. Altwegg, M. Combi, S. A. Fuselier, J. De Keyser, N. F. W. Ligterink, M. Rubin and S. F. Wampfler, *Astrophys. J.*, 2024, 976, 74.
- 25 N. W. Reed, R. L. Shearer, S. E. McGlynn, B. A. Wing, M. A. Tolbert and E. C. Browne, *Astrophys. J. Lett.*, 2024, 973, L38.
- 26 W. C. Price, J. P. Teegan and A. D. Walsh, *Proc. R. Soc. London, Ser. A*, 1950, 201, 600–609.
- 27 L. B. Clark and W. T. Simpson, *J. Chem. Phys.*, 1965, **43**, 3666–3672.
- 28 I. Tokue, A. Hiraya and K. Shobotake, *Chem. Phys.*, 1987, 116, 449–456.
- 29 G. L. Vaghjiani, J. Chem. Phys., 1993, 99, 5936-5943.
- 30 S. Choi, T. Y. Kang, K.-W. Choi, S. Han, D.-S. Ahn, S. J. Baek and S. K. Kim, *J. Phys. Chem. A*, 2008, **112**, 7191–7199.
- 31 J. E. Stevens, K. F. Freed, M. F. Arendt and R. L. Graham, *J. Chem. Phys.*, 1994, **101**, 4832–4841.
- 32 J. E. Stevens, H. W. Jang, L. J. Butler and J. C. Light, *J. Chem. Phys.*, 1995, **102**, 7059–7069.
- 33 S. H. S. Wilson, M. N. R. Ashfold and R. N. Dixon, *J. Chem. Phys.*, 1994, **101**, 7538–7547.
- 34 J. Segall, Y. Wen, R. Singer, M. Dulligan and C. Wittig, *J. Chem. Phys.*, 1993, **99**, 6600–6606.
- 35 G. A. Amaral, F. Ausfelder, J. G. Izquierdo, L. Rubio-Lago and L. Bañares, *J. Chem. Phys.*, 2007, **126**, 024301.
- 36 E. Jensen, J. S. Keller, G. C. G. Waschewsky, J. E. Stevens, R. L. Graham, K. F. Freed and L. J. Butler, *J. Chem. Phys.*, 1993, 98, 2882–2890.
- 37 S. B. Barone, A. A. Turnipseed, T. Gierczak and A. R. Ravishankara, J. Phys. Chem., 1994, 98, 11969–11977.
- 38 D. R. Yarkony, J. Chem. Phys., 1996, 104, 3639-3644.
- 39 D. R. Yarkony, J. Chem. Phys., 1996, 104, 7866-7881.
- 40 J. G. Izquierdo, G. A. Amaral, F. Ausfelder, F. J. Aoiz and L. Bañares, *ChemPhysChem*, 2006, 7, 1682–1686.
- 41 Z. C. Chen, Q. Shuai, A. T. J. B. Eppink, B. Jiang, D. X. Dai, X. M. Yang and D. H. Parker, *Phys. Chem. Chem. Phys.*, 2011, 13, 8531–8536.

42 D. V. Chicharro, S. M. Poullain, L. Rubio-Lago and

Edge Article

43 B. Ruscic and D. H. Bross, Active Thermochemical Tables (ATcT) values based on ver. 1.202 of the Thermochemical Network, 2024, available at https://atct.anl.gov/.

L. Bañares, J. Phys. Chem. A, 2019, 123, 8552-8561.

- 44 A. Kramida, Y. Ralchenko, J. Reader and NIST ASD Team, NIST Atomic Spectra Database (ver. 5.12), [Online] [2025, June 05], National Institute of Standards and Technology, Gaithersburg, MD, 2024, DOI: 10.18434/T4W30F. Available: https://physics.nist.gov/asd.
- 45 M. E. Jacox, *J. Phys. Chem. Ref. Data*, 1988, 17, 269–511, and references therein.
- 46 D. Kamra and J. M. White, *J. Photochem.*, 1975, 4, 361–373.
- 47 S. Nourbakhsh, K. Norwood, H.-M. Yin, C.-L. Liao and C. Y. Ng, *J. Chem. Phys.*, 1991, **95**, 946–954.
- 48 C.-W. Hsu, C.-L. Liao and C.-Y. Ng, *J. Chin. Chem. Soc.*, 1995, 42, 149–155.
- 49 Y. Chang, S. R. Yu, Q. M. Li, Y. Yu, H. L. Wang, S. Su, Z. C. Chen, L. Che, X. G. Wang, W. Q. Zhang, D. X. Dai, G. R. Wu, K. J. Yuan and X. M. Yang, *Rev. Sci. Instrum.*, 2018, 89, 063113.
- 50 J. M. Zhou, Y. R. Zhao, C. S. Hansen, J. Y. Yang, Y. Chang, Y. Yu, G. K. Cheng, Z. C. Chen, Z. G. He, S. R. Yu, H. B. Ding, W. Q. Zhang, G. R. Wu, D. X. Dai, C. M. Western, M. N. R. Ashfold, K. J. Yuan and X. M. Yang, Nat. Commun., 2020, 11, 1547.
- 51 Z. X. Li, M. Zhao, T. Xie, Z. J. Luo, Y. Chang, G. K. Cheng, J. Y. Yang, Z. C. Chen, W. Q. Zhang, G. R. Wu, X. G. Wang, K. J. Yuan and X. M. Yang, *J. Phys. Chem. Lett.*, 2021, 12, 844–849.
- 52 Z. X. Li, Y. L. Fu, Z. J. Luo, S. K. Yang, Y. C. Wu, H. Wu, G. R. Wu, W. Q. Zhang, B. N. Fu, K. J. Yuan, D. H. Zhang and X. M. Yang, *Science*, 2024, 383, 746–750.
- 53 S. K. Yang, Y. C. Wu, Z. J. Luo, Z. X. Li, W. Hua, Y. Chang, X. G. Wang, K. J. Yuan and X. M. Yang, *Chin. J. Chem. Phys.*, 2024, 37, 286–294.
- 54 Y. Chang, M. N. R. Ashfold, K. J. Yuan and X. M. Yang, *Natl. Sci. Rev.*, 2023, 10, nwad158.
- 55 Y. R. Zhao, Z. J. Luo, Y. Chang, Y. C. Wu, S. E. Zhang, Z. X. Li, H. B. Ding, G. R. Wu, J. S. Campbell, C. S. Hansen, S. W. Crane, C. M. Western, M. N. R. Ashfold, K. J. Yuan and X. M. Yang, *Nat. Commun.*, 2021, 12, 4459.
- 56 Y. R. Zhao, J. J. Chen, Z. J. Luo, Y. Chang, J. Y. Yang, W. Q. Zhang, G. R. Wu, S. W. Crane, C. S. Hansen, H. B. Ding, F. An, X. X. Hu, D. Q. Xie, M. N. R. Ashfold, K. J. Yuan and X. M. Yang, *Chem. Sci.*, 2023, 14, 2501–2517.
- 57 Z. C. Chen, A. T. J. B. Eppink, B. Jiang, G. C. Groenenboom, X. M. Yang and D. H. Parker, *Phys. Chem. Chem. Phys.*, 2011, 13, 2350–2355.

- 58 S. Mondal and K. P. Liu, *J. Phys. Chem. A*, 2019, **123**, 1514–1520.
- 59 Y. C. Wu, Z. J. Luo, S. K. Yang, S. J. McGoldrick, J. L. Mendham, Z. X. Li, S. Y. Zhou, Y. X. Dong, D. X. Dai, C. S. Hansen, M. N. R. Ashfold, K. J. Yuan and X. M. Yang, J. Chem. Phys., submitted.
- 60 A. B. Callear and D. R. Dickson, *Trans. Faraday Soc.*, 1970, 66, 1987–1995.
- 61 C. Anastasi, M. Broomfield, O. J. Nielsen and P. Pagsberg, *Chem. Phys. Lett.*, 1991, **182**, 643–648.
- 62 A. Kumar, P. K. Chowdhury, K. V. S. R. Rao and J. P. Mittal, Chem. Phys. Lett., 1992, 198, 406–412.
- 63 R. T. Bise, H. Choi, H. B. Pedersen, D. H. Mordaunt and D. M. Neumark, *J. Chem. Phys.*, 1999, **110**, 805–816.
- 64 C.-P. Liu, Y. Matsuda and Y.-P. Lee, J. Chem. Phys., 2003, 119, 12335–12341.
- 65 V. Boudon, M. Rey and M. Loëte, J. Quant. Spectrosc. Radiat. Transfer, 2006, **98**, 394–404.
- 66 K. Kefala, V. Boudon, S. N. Yurchenko and J. Tennyson, J. Quant. Spectrosc. Radiat. Transfer, 2024, 316, 108897.
- 67 Y. Chang, F. An, Z. C. Chen, Z. J. Luo, Y. Zhao, X. X. Hu, J. Y. Yang, W. Q. Zhang, G. R. Wu, D. Q. Xie, K. J. Yuan and X. M. Yang, *Nat. Commun.*, 2021, 12, 6303.
- 68 Y. R. Shao, J. J. Chen, Z. J. Luo, Z. X. Li, S. K. Yang, Y. Chang, F. An, Z. C. Chen, J. Y. Yang, G. R. Wu, W. Q. Zhang, X. X. Hu, D. Q. Xie, H. B. Ding, K. J. Yuan and X. M. Yang, *J. Phys. Chem. Lett.*, 2022, 13, 9786–9792.
- 69 L. Zhou, Y. C. Chang, Q.-Z. Yin and W. M. Jackson, *Science*, 2014, 346, 61–64.
- 70 W. T. Chen, L. Zhang, D. F. Yuan, Y. Chang, S. R. Yu, S. W. Wang, T. Wang, B. Jiang, K. J. Yuan, X. M. Yang and X. G. Wang, J. Phys. Chem. Lett., 2019, 10, 4783–4787.
- 71 Y. Chang, Y. L. Fu, Z. C. Chen, Z. J. Luo, Y. R. Zhao, Z. X. Li, W. Q. Zhang, G. R. Wu, B. N. Fu, D. H. Zhang, M. N. R. Ashfold, X. M. Yang and K. J. Yuan, *Chem. Sci.*, 2023, 14, 8255–8261.
- 72 J. O. Thomas, K. E. Lower and C. Murray, *J. Phys. Chem. Lett.*, 2012, 3, 1341–1345.
- 73 J. O. Thomas, K. E. Lower and C. Murray, *J. Phys. Chem. A*, 2014, **118**, 9844–9852.
- 74 A. G. Suits, Acc. Chem. Res., 2008, 41, 873-881.
- 75 E. Herbst and E. F. van Dishoeck, *Annu. Rev. Astron. Astrophys.*, 2009, **47**, 427–480.
- 76 V. Bariosco, L. Tinacci, S. Pantaleone, C. Ceccarelli,
 A. Rimola and P. Ugliengo, *Mon. Not. R. Astron. Soc.*, 2025,
 539, 82–94, and references therein.