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# Photosalient effect and reversible photochromic photoluminescence driven by cascade [2 + 2] cycloaddition reaction and water adsorption in a 0D hybrid metal halide†

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The ionic structure of hybrid metal halides (HMHs) is conducive to the integration of the photoactivity of organic components and mechanical properties of inorganic components. Therefore, HMHs can function as a suitable platform to advance photomechanical performance due to their great structural designability and wide compositional diversity. Herein, a notable photosalient effect was observed in 0D ionic (LH)<sub>2</sub>InCl<sub>5</sub> (**1**, L = 4-styrylpyridine). Upon irradiation with 450 nm light, **1** undergoes a rare single-crystal-to-single-crystal transformation to (L<sub>2</sub>H<sub>2</sub>)[In(H<sub>2</sub>O)Cl<sub>5</sub>] (**1a**), along with distinct photochromic photoluminescence and various photomechanical responses, driven by the cascade photoinduced [2 + 2] cycloaddition of LH<sup>+</sup> and water sorption by coordinatively unsaturated In<sup>3+</sup>. A detailed mechanistic investigation unraveled the photosalient effect by initiation of the cycloaddition reaction and enhancement by water sorption, which resulted in a phase transformation that is more energetically and kinetically favorable through decreasing Gibbs free energy (1.51 eV) and a kinetical energy barrier (0.29 eV). Additionally, the partially reversible transformation between LH<sup>+</sup> monomers and dimers endows **1a** with excellent photoswitch performance due to the variation in the  $\pi$  electron-conjugated degree for LH<sup>+</sup>. At last, by integrating powdered **1** with polymethyl methacrylate (PMMA), the **1**–PMMA hybrid matrix membrane exhibited rapid and controllable responsive behaviors, including bending and curling, through controlling the direction of the concentration gradient. This work demonstrates that ionic HMHs might function as a promising platform for advanced photomechanical and photochromic materials.

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## Introduction

Actuating materials are thought to hold great promise in the fields of soft robotics, artificial muscles, and biomimetics because they are capable of a kinetic energy output to implement complex and rapid macroscopic movements under external stimuli including light, heat, force, and chemical

reagents.<sup>1–8</sup> Among the different stimuli, light is advantageous because it imparts a high degree of spatial and temporal control without any physical contact, as well as precise operation over properties including intensity, wavelength, and polarization.<sup>9,10</sup>

Currently, the motions of photoactuators are usually triggered through photothermal (molecular sorption/desorption and phase transition) or photochemical (photoisomerization and photodimerization) processes.<sup>10</sup> Among these strategies, [2 + 2] photocycloaddition has been widely employed to translate nanoscopic structural perturbation into macroscopic motions such as jumping, curling, bending, twisting, and disintegrating.

To trigger [2 + 2] cycloaddition in crystals, parallel olefin groups should satisfy Schmidt's topochemical criteria to undergo structural transformation.<sup>11–14</sup> Therefore, such photo-actuator materials are mostly based on molecular crystals (MCs)<sup>15–20</sup> and coordination polymers (CPs)<sup>21–28</sup> because of the adequate understanding of their structural design and tunability. Although a great deal of effort has been recently devoted to the development of new photocycloaddition crystals to unravel the structure–activity relationships,<sup>29–37</sup> few effective strategies have been proposed to improve photomechanical

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performance due to the limited understanding of a cascade of processes including strain generation, accumulation, transmission, and release during the photomechanical response.

Different from MCs and CPs mainly assembled through non-covalent and coordinated interaction, respectively, ionic crystals (ICs) are assembled through coulombic interaction with more dense packing. Such a stacking mode usually endows them with a larger Young's modulus, implying that a more intense photomechanical response might occur than that of MCs and CPs under same variation in cell volume. Among various ICs, hybrid metal halides (HMHs) are a class of ionic semiconductors derived from halide perovskites that has recently attracted enormous attention.<sup>38–45</sup>

Because there are great advantages when structural and compositional tunability is enabled for HMHs, various strategies<sup>46–54</sup> have been proposed to tune the structural features of HMHs, including dimensionality, layer number, and polyhedron connection modes that will affect the light-to-work energy conversion and transmission processes of interior strain in the photomechanical response. Therefore, HMHs might function as a satisfactory model to reveal the structure–activity relationships of the photomechanical response.

With this in mind, few works have successfully achieved [2 + 2] cycloaddition in HMHs. For instance, hybridizing 4-pyridineacrylic acid with BiCl<sub>3</sub> results in photochemical reactions, including photosolvent motions and photochromism after ultraviolet (UV) irradiation.<sup>55</sup> The [2 + 2] photocycloaddition of 4-vinyl pyridine derivatives induces photochromism and photoinduced-cracking in Zn- and Pb-based HMHs.<sup>56</sup> These two examples demonstrate the feasibility of photocycloaddition in HMHs. However, the photomechanical properties of HMHs have not been systematically explored.

In this work, 4-styrylpyridine (**L**) was successfully hybridized with InCl<sub>3</sub> to form a zero-dimensional (0D) (LH)<sub>2</sub>InCl<sub>5</sub> (denoted as **1**) with coordinatively unsaturated In<sup>3+</sup> sites. The ideally parallel stacking of LH<sup>+</sup> dimers enables a fast and thorough photocycloaddition reaction, followed by two-step H<sub>2</sub>O sorption on In<sup>3+</sup> sites with a trigonal bipyramidal geometry. Furthermore, the significant lattice reorganization caused by photocycloaddition and H<sub>2</sub>O sorption results in notable photomechanical motions and photochromic photoluminescence (PL) of **1**.

## Results and discussion

Compound **1** was prepared through a cooling crystallization method, and the synthesis details are found in the ESI.† To determine the crystal phase of the resulting sample, single crystal X-ray diffraction (SCXRD) was conducted, and the crystal structure of **1** is shown in Fig. 1a. It was observed that **1** crystallizes in triclinic space group *P* $\bar{1}$  with *Z* = 2 (Table S1†) and consists of discrete [InCl<sub>5</sub>]<sup>2−</sup> polyhedra and LH<sup>+</sup> dimers with anti-parallel alignment. Different from other In-based HMHs adopting an octahedral coordination,<sup>57–60</sup> a rare [InCl<sub>5</sub>]<sup>2−</sup> polyhedron with trigonal bipyramidal geometry is formed in **1**, which is caused by the charge balance and the strong interactions between the LH<sup>+</sup> dimer and the [InCl<sub>5</sub>]<sup>2−</sup> polyhedron.

As shown in Fig. 1a, the short distance between the phenyl ring and the [InCl<sub>5</sub>]<sup>2−</sup> polyhedron hinders the coordination of free Cl<sup>−</sup> or H<sub>2</sub>O with coordinatively unsaturated In<sup>3+</sup>. The strong intermolecular interactions were confirmed from the Hirshfeld surface analysis (Fig. S1†) that quantitatively determined the predominant contributions from the N–H⋯Cl (42.1%) and H⋯H (30.7%) bonding interactions. [InCl<sub>5</sub>]<sup>2−</sup> polyhedra are interconnected through coulombic and hydrogen bonding interactions with cationic LH<sup>+</sup> (Fig. S2†). LH<sup>+</sup> dimers exhibit two stacking modes in **1**, as denoted by the cyan and orange olefins. Each stacking mode shows two types of disorder, with an occupancy of approximately 69% (solid cyan) and approximately 74% (solid orange) for the predominant alignment, respectively.

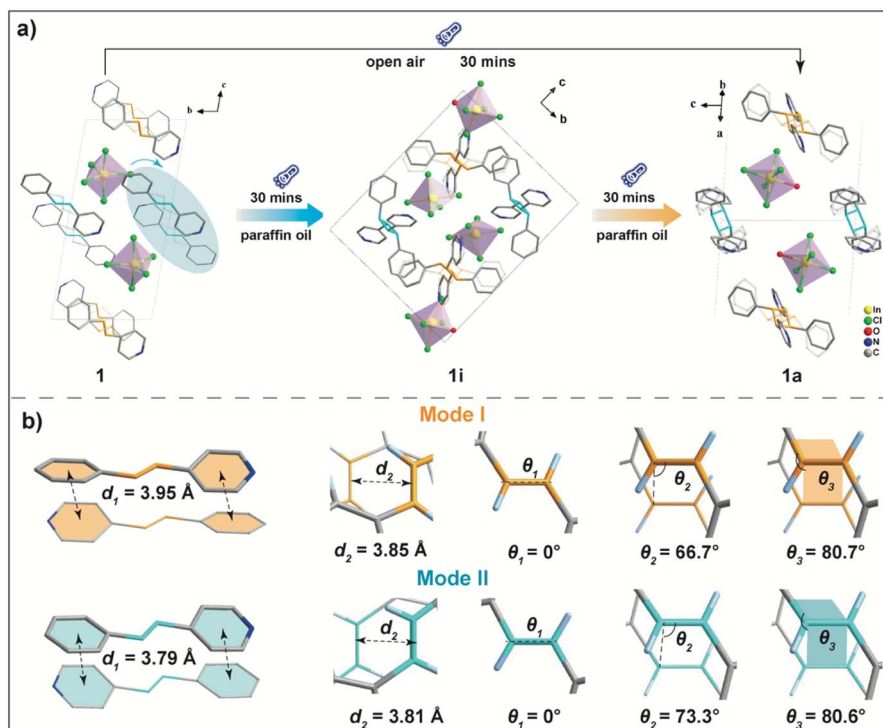
Upon irradiation of 450 nm light for 30 min in open air, a single-crystal-to-single-crystal transformation occurred from **1** to (L<sub>2</sub>H<sub>2</sub>)[In(H<sub>2</sub>O)Cl<sub>5</sub>] (**1a**), with the crystal data listed in Table S1.† Compared to **1**, the LH<sup>+</sup> photodimers in **1a** revolve away the [InCl<sub>5</sub>]<sup>2−</sup> polyhedron upon photocycloaddition, which forces the rotation of the [InCl<sub>5</sub>]<sup>2−</sup> polyhedron and creates sufficient free space for In<sup>3+</sup> to adsorb H<sub>2</sub>O. Once single crystals of **1** were placed in paraffin oil to slow the sorption of H<sub>2</sub>O, some **1** single crystals transferred to an intermediate phase of (L<sub>2</sub>H<sub>2</sub>)<sub>2</sub>[In(H<sub>2</sub>O)Cl<sub>5</sub>][InCl<sub>5</sub>] (**1i**), with complete photodimerization of LH<sup>+</sup>. Interestingly, only half of the [InCl<sub>5</sub>]<sup>2−</sup> polyhedra in **1i** were coordinated with H<sub>2</sub>O molecules in paraffin oil. A further increase in the reaction time resulted in the phase transformation from **1i** to **1a**.

To understand the fast and thorough photocycloaddition, the alignments of the LH<sup>+</sup> dimers in **1** are shown in Fig. 1b. Because of the similar structural change for each disordered part, only the high occupancy of disordered olefin is discussed. All LH<sup>+</sup> dimers in **1** align in an antiparallel head-to-tail manner, with distances of 3.95 Å (mode I) and 3.79 Å (mode II) between the pyridyl and phenyl rings. Profiting from their  $\pi$ – $\pi$  stacking interaction, the olefin bonds are perfectly aligned with an angle of  $\theta_1 = 0.0^\circ$ . Although LH<sup>+</sup> dimers slip along the lateral axis ( $\theta_2 \neq 90^\circ$ ) and slightly incline ( $\theta_3 \neq 90^\circ$ ), the distances of the olefin bonds (3.85 Å for mode I and 3.81 Å for mode II) are well satisfied with Schmidt's criteria.

Despite the crystal system and space group of **1** remaining the same after phase transformation, significant variations were observed for the crystal structure due to the thorough photocycloaddition reaction of the LH<sup>+</sup> dimers. As shown in Fig. 2a, the average breadth of the LH<sup>+</sup> dimers in mode I increased from 4.07 Å (**1**) to 6.57 Å (**1i**) and 6.2661 Å (**1a**), along with a reduction in the length from 11.07 Å (**1**) to 9.85 Å (**1i**) and 10.10 Å (**1a**). A similar structural change was also observed for LH<sup>+</sup> dimers in mode II. In comparison with the centrosymmetric LH<sup>+</sup> dimers in **1** and **1a**, non-centrosymmetric LH<sup>+</sup> photodimers were formed in **1i**. This difference is due to the presence of [InCl<sub>5</sub>]<sup>2−</sup> and [In(H<sub>2</sub>O)Cl<sub>5</sub>]<sup>2−</sup> pairs that interact with terminal pyridinium with a different degree of force (Fig. S3†).

The significant geometric variation of LH<sup>+</sup> dimers also drives the rotation of all [InCl<sub>5</sub>]<sup>2−</sup> polyhedra, which creates sufficient free volume for the following H<sub>2</sub>O sorption. As shown in Fig. 2b, compared to the dense stacking of **1** and **1a** (free volume





**Fig. 1** (a) Crystal structures and phase transformation of **1**, **1i**, and **1a**. Light source: 450 nm, 260 mW cm<sup>-2</sup>. (b) Alignments of LH<sup>+</sup> dimers in **1**.  $d_1$  is defined by the mean distance between the centroid of the phenyl and pyridyl rings.  $d_2$  is defined by the distance between the parallel C=C bonds.  $\theta_1$  denotes the angle formed between two adjacent C=C bonds in the top view.  $\theta_2$  is defined by the angle involving the C=C bond and one of the carbon atoms from the adjacent C=C bond. The orange and cyan planes for  $\theta_3$  are defined by the two C=C bonds and their adjoining single bond.

percentage = 0%), free volume with a percentage of 5.1% is created near the [InCl<sub>5</sub>]<sup>2-</sup> and [In(H<sub>2</sub>O)Cl<sub>5</sub>]<sup>2-</sup> polyhedra in **1i**. The adsorption of H<sub>2</sub>O molecules generates octahedral [In(H<sub>2</sub>O)Cl<sub>5</sub>]<sup>2-</sup> pairs with a distance of 5.72 Å and an intermolecular hydrogen bonding length of 3.25 Å. Upon phase transformation from **1i** to **1a**, all [InCl<sub>5</sub>]<sup>2-</sup> pairs transform to [In(H<sub>2</sub>O)Cl<sub>5</sub>]<sup>2-</sup> pairs, concurrent with a decrease in the distance of the [InCl<sub>5</sub>]<sup>2-</sup> pairs from 7.45 Å to 5.75 Å (Fig. 2a) and the reconstruction of an intermolecular hydrogen bonding network (Fig. S4†).

To further understand the process of phase transformation, the lattice parameters of a single crystal of **1** were monitored under different irradiation times. As shown in Fig. 2c, upon the photocycloaddition reaction of LH<sup>+</sup> dimers, the *a*-axis length gradually decreased, along with an increase in the *b*-axis length within 40 s, while the *c*-axis length remained nearly constant and the cell volume (*V*) slightly increased by approximately 1.1%. Interestingly, increasing the irradiation time to 50 s resulted in a sharp increase in *V* to 4.3%, along with the first sudden change in *a*- and *c*-axis lengths. Additionally, a second sudden change was observed for all lattice parameters once the irradiation time was extended to 90 s.

Such an evolutionary tendency suggests a two-step phase transformation during the irradiation. Within the first 40 s, the slowly varied lattice parameters indicate that the photodimerization reaction is ongoing. Once the lattice stress sufficiently accumulates to drive the rotation of [InCl<sub>5</sub>]<sup>2-</sup>, H<sub>2</sub>O

sorption will occur, accompanied by the phase transformation from **1** to **1i**. When extending the irradiation time to 90 s, the H<sub>2</sub>O sorption by the other half of [InCl<sub>5</sub>]<sup>2-</sup> induces the phase transformation from **1i** to **1a**, resulting in a total increase in cell volume to 8.0%.

Density functional theory (DFT) calculations were also conducted to understand the cycloaddition process. Optimization of **1** was performed, using their X-ray structures as the initial structures. The C<sub>15</sub>–C<sub>40</sub> bond length was fixed at specific values during the optimization process to examine the impact of this constraint upon the reaction pathway. As shown in Fig. 2d, this process is highly endothermic by 1.31 eV, with an electronic energy barrier of 2.09 eV. Once the distance of C<sub>15</sub>–C<sub>40</sub> decreased to 2.0 Å, a maximum value was obtained. A further decrease in the distance to 1.6 Å will result in a local minimum value of energy, and was consistent with the experimental results of the bond length of C–C in cyclobutane, at approximately 1.57 Å (obtained at 200 K).

Fig. 2e further reveals the H<sub>2</sub>O sorption effect on the free energy variation of the phase transformation. Once H<sub>2</sub>O sorption is involved, the energy variation changes from endothermic (0.57 eV) to exothermic (–0.94 eV), with a reduction in the energy barrier of 0.29 eV. This highlights the critical role of H<sub>2</sub>O sorption in enhancing the cycloaddition reaction by synergistically modulating the thermodynamic equilibria and lowering the kinetic energy barrier.

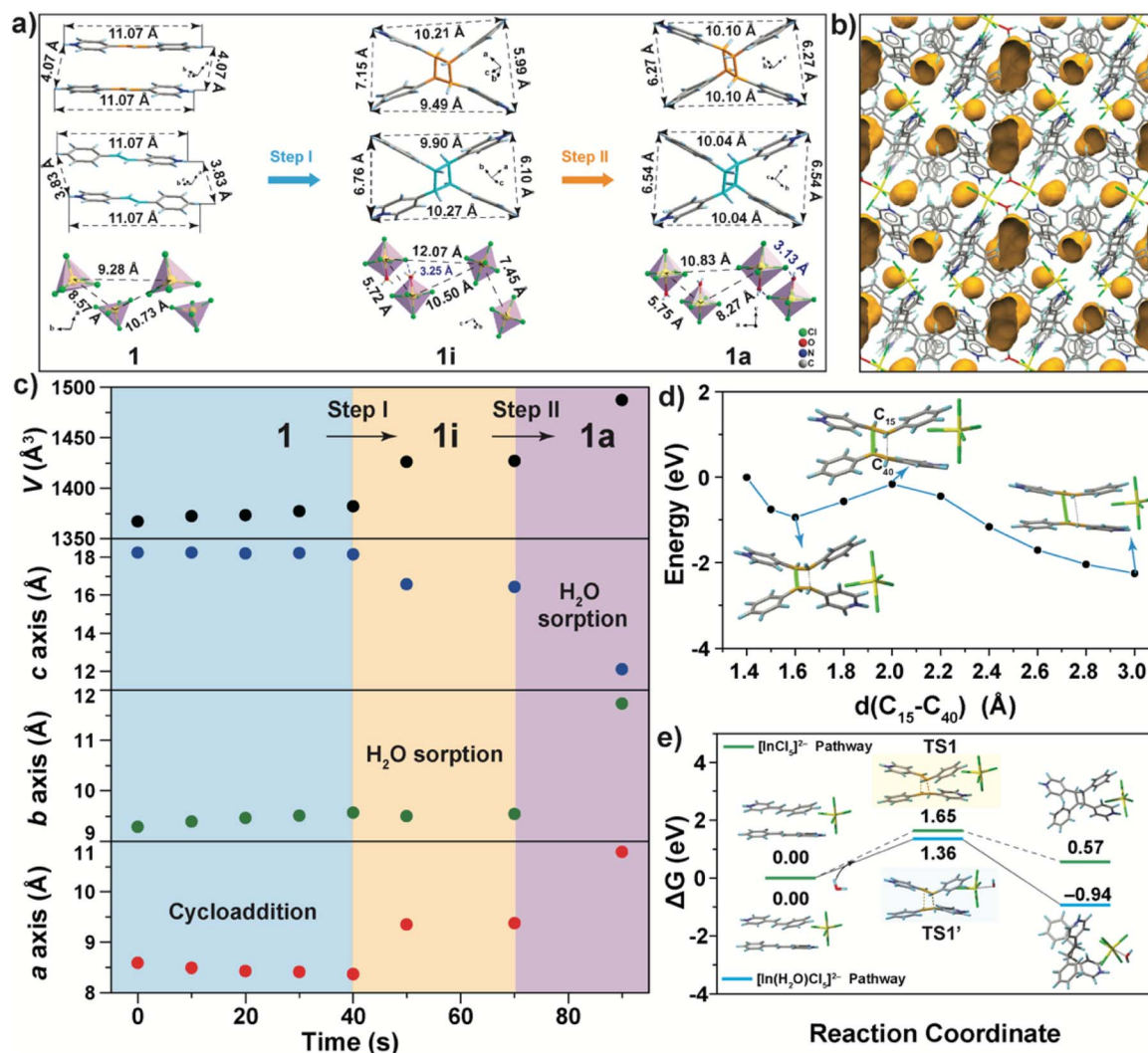


Fig. 2 (a) Geometrical variation of LH<sup>+</sup> dimers and the [InCl<sub>5</sub>]<sup>2-</sup> polyhedron in **1**, **1i**, and **1a**. (b) Schematic illustration of the free volume in **1i** (probe radius: 1.0 Å; approx. grid spacing: 0.2 Å). (c) Lattice parameters of **1** at different irradiation times (light source: 450 nm, 260 mW cm<sup>-2</sup>). (d) Energy profile of the [2 + 2] cycloaddition reaction with constrained C-C bond length. The C<sub>15</sub>-C<sub>40</sub> bond length was fixed during the optimization process to examine the impact of this constraint on the reaction pathway. (e) Computational study of LH<sup>+</sup> dimers, the [InCl<sub>5</sub>]<sup>2-</sup> polyhedron, LH<sup>+</sup> photodimers, and the [In(H<sub>2</sub>O)Cl<sub>5</sub>]<sup>2-</sup> octahedron involved in [2 + 2] cycloaddition reactions.

Due to the instability of **1i**, only the phase purities of powdered **1** and **1a** were determined through powder X-ray diffraction (PXRD). As shown in Fig. 3a, the diffraction patterns of **1** and **1a** match the simulated results, indicating the high phase purity of the as-prepared samples. To confirm the H<sub>2</sub>O sorption, the thermogravimetric analysis (TGA) of powdered **1** and **1a** is shown in Fig. 3b. The decomposition temperature of **1** was approximately 193 °C, while **1a** exhibited the first weight loss platform starting at approximately 150 °C with a weight loss ratio of 2.60%, corresponding well with the theoretical weight ratio (2.67%) of H<sub>2</sub>O in **1a**.

The presence of coordinated H<sub>2</sub>O was also confirmed by energy-dispersive spectroscopy (EDS) mapping. The EDS mapping (Fig. S5†) of **1a** exhibits a morphology-dependent distribution for all elements including O, while this was insignificant for the O element in **1**. Fourier transform infrared (FT-

IR) spectroscopy (Fig. 3c) shows a notable band at 3430 cm<sup>-1</sup> caused by the O-H stretching vibration for **1a**, while this band is absent for **1**. Additionally, the intensity of the =C-H bending vibration peak (964 cm<sup>-1</sup>) of **1** dramatically decreased upon phase transformation, which verified the cascade [2 + 2] cycloaddition and H<sub>2</sub>O sorption reactions.

To evaluate the rate of photocycloaddition, the samples were digested in dimethyl sulfoxide-*d*<sub>6</sub> (DMSO-*d*<sub>6</sub>) to conduct proton (<sup>1</sup>H) nuclear magnetic resonance (<sup>1</sup>H NMR) spectroscopy at various time intervals. As shown in Fig. 3d and S6,† the signals of 8.85 (H<sub>a</sub>) and 8.21 ppm (H<sub>b</sub>) for **1** resulted from the pyridyl ring, while the peaks of 7.76 (H<sub>d</sub>), 7.55 (H<sub>e</sub>), and 7.48 (H<sub>f</sub>) ppm were derived from the phenyl ring, respectively. The peaks at 8.01 (H<sub>c</sub>) and 7.44 ppm (H<sub>g</sub>) for **1** were assigned to olefinic protons. Upon 450 nm irradiation, the intensity of these peaks gradually decreased, and a new set of peaks appeared, including



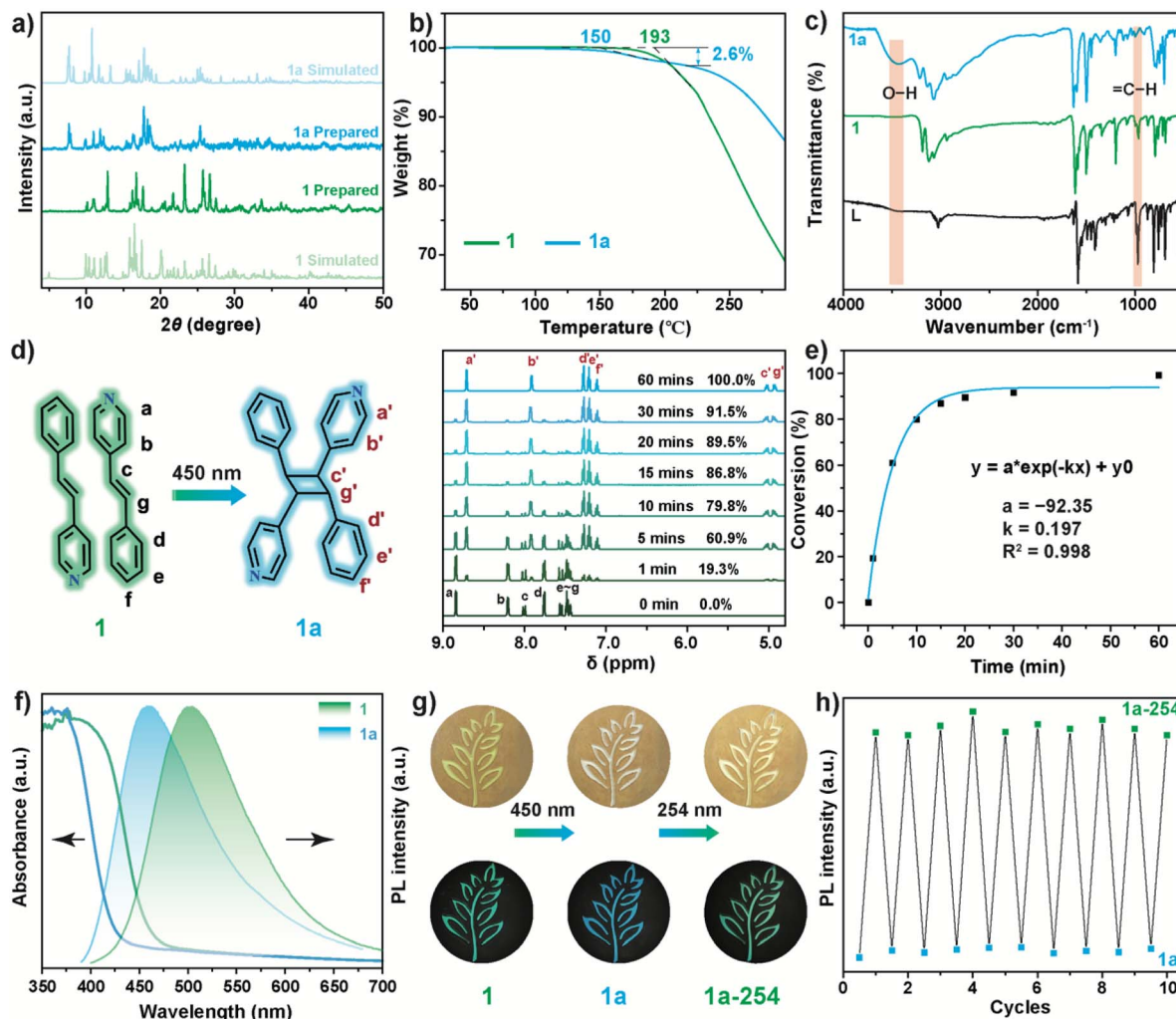


Fig. 3 (a) PXRD patterns, (b) TGA curves, and (c) FT-IR spectroscopy of **1** and **1a**. (d)  $^1\text{H}$  NMR spectra (solvent:  $\text{DMSO}-d_6$ ) of **1** with different irradiation times. (e) Conversion ratios for the photodimerization of **1**. (f) UV-Vis diffuse reflectance and PL spectra of **1** and **1a**. (g) Digital images of **1**, **1a**, and **1a-254** ( $6.1 \text{ mW cm}^{-2}$ ). (h) Reversibility test between **1a** and **1a-254**; the PL intensity was monitored at 502 nm, and the irradiation time was 10 min for 254 and 450 nm.

two peaks at 4.93 ppm and 5.02 ppm, which were derived from cyclobutane. By integrating the peak area of  $\text{H}_a$ , the photocycloaddition proportion was determined, and is shown in Fig. 3e. The photodimerization of  $\text{LH}^+$  was nearly complete within 60 min. The fitting results indicate the first-order reaction kinetics of  $\text{LH}^+$  photocycloaddition with a rate constant of  $0.197 \text{ min}^{-1}$ , which is higher than that in CPs.<sup>61</sup>

UV-Vis diffuse reflectance and photoluminescence (PL) spectra were also obtained to determine the influence of [2 + 2] photocycloaddition on the optical properties of the samples. The absorption onset at approximately 470 nm exhibited by **1** (Fig. 3f) was consistent with the absorption of  $\text{LH}^+\cdot\text{Cl}$  (Fig. S7†). Upon illumination with 450 nm light, the absorption onset of **1a** underwent a blueshift to approximately 430 nm, caused by the disruption of the  $\pi$  electron conjugation of  $\text{LH}^+$ . Similarly, the photodimerization reaction of  $\text{LH}^+$  also greatly influenced the PL of **1**. As shown in Fig. 3f, the broadband emission peaked at 502 nm and then blueshifted to 460 nm due to the

photocycloaddition reaction. This photochromic PL was also observed for  $\text{LH}\cdot\text{Cl}$  (Fig. S8†), in which the [2 + 2] photocycloaddition reaction was confirmed by the  $^1\text{H}$  NMR spectra (Fig. S9†).

These results indicate the photochromic origin resulting from the [2 + 2] cycloaddition reaction of  $\text{LH}^+$ , rather than the  $\text{H}_2\text{O}$  sorption by  $[\text{InCl}_5]^{2-}$  polyhedra. Consistent with the UV-Vis and PL spectra, **1** exhibited a distinct photochromic behavior from yellow-green to white under room light (Fig. 3g), along with the emission color transformation from green to blue under UV light. Upon irradiating **1a** with 254 nm (**1a-254**) for 10 min, the color of **1a-254** reverted to a faint yellow-green, concurrent with the emissive color returning to green. This color evolution indicates the partial reversibility between  $\text{LH}^+$  photodimers and  $\text{LH}^+$  monomers, which was confirmed from the  $^1\text{H}$  NMR.

As shown in Fig. S10,† a dedimerization ratio of approximately 13.2% (8 h) was obtained. The PL spectrum of **1a-254**

(Fig. S11†) also indicates the successful dedimerization of  $\text{LH}^+$  dimers. Increasing the irradiation time for 254 nm resulted in the falling of the main peak at 460 nm and the rising of the shoulder peak at 502 nm. By monitoring the PL intensity at 502 nm (Fig. 3h), excellent reversibility between **1a** and **1a**-254 was observed, indicating the great potential of **1a** in photo-switching. However, due to the irreversible  $\text{H}_2\text{O}$  sorption, **1a** cannot revert back to **1**, as evidenced from the PXRD patterns (Fig. S12†).

Because of the significant changes in lattice parameters caused by the cascade photocycloaddition and  $\text{H}_2\text{O}$  sorption, single crystals of **1** exhibited distinct macroscopic photomechanical motions. As shown in Fig. 4a and Video S1,† once

irradiated with 450 nm light, the crystal ( $1.22 \times 0.30 \times 0.15 \text{ mm}^3$ ) jumps away at 0.5 s, followed by the motions of cracking, splitting, and hopping. Upon irradiation with 405 nm blue light, the crystal ( $1.08 \times 0.72 \times 0.15 \text{ mm}^3$ ) shows a bursting motion at 3.0 s and thoroughly disintegrates at 25 s (Fig. 4b and Video S2†). Nevertheless, irradiating crystal **1** ( $0.75 \times 0.35 \times 0.14 \text{ mm}^3$ ) with 365 nm light resulted in gradual photochromic PL with insignificant photomechanical motions (Fig. 4c and Video S3†).

These excitation wavelength-dependent photoresponsive behaviors were attributed to the discrepancies in absorption coefficients and the power density of excited light that induces varying degrees of interior stress. To unravel the contribution of

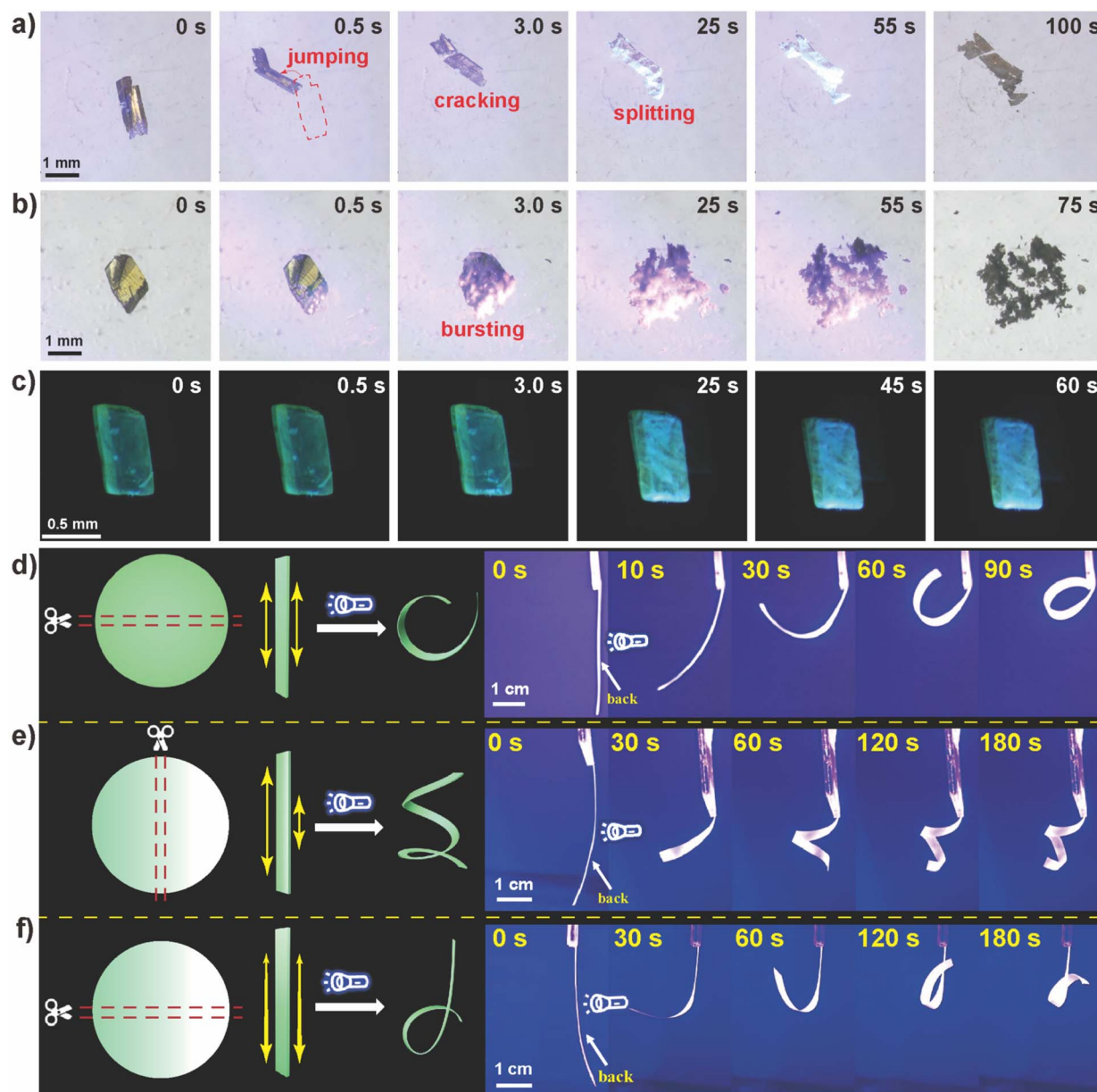


Fig. 4 Digital images of crystal **1** under (a) 450 nm ( $260 \text{ mW cm}^{-2}$ ), (b) 405 nm ( $260 \text{ mW cm}^{-2}$ ), and (c) 365 nm ( $80 \text{ mW cm}^{-2}$ ) irradiation with different times. Images of (d) 1-PMMA with uniform distribution, (e) 1-PMMA with a perpendicular concentration gradient, and (f) 1-PMMA with a parallel concentration gradient upon irradiation with 450 nm light.

photocycloaddition and H<sub>2</sub>O sorption, crystal **1** was placed in a Schlenk tube evacuated for 2 h at 80 °C, and then filled with dry nitrogen to exclude the sorption of H<sub>2</sub>O. As shown in Fig. S13,† upon irradiating crystal **1** with 450 nm light, photo-mechanical behaviors such as jumping and cracking continue to be observed (Video S4†), suggesting that the photomechanical effect of **1** was initiated by the cycloaddition reaction of **LH**<sup>+</sup>, rather than H<sub>2</sub>O sorption. However, the photomechanical strength seems much weaker than that of **1** in open air, indicating the critical role of H<sub>2</sub>O sorption in enhancing the photosalient effect.

Although the increase in cell volume was relatively smaller than that of **L**-based CPs (Table S2†), the photomechanical behavior of **1** was more intense in our case. Such change was ascribed to the ionic structure of HMMs with dense stacking that exhibits relatively higher brittleness. Although an understanding of the mechanism during the photoresponsive process can be acquired through a well-defined crystal structure obtained from X-ray diffraction, it remains a great challenge to combine the motions of discrete microscale crystals into one organized macroscale system.

In an effort to produce a more predictable mechanical behavior,<sup>62</sup> powdered **1** was mixed with polymethyl methacrylate (PMMA) to produce a hybrid matrix membrane (**1**-PMMA). The PXRD pattern of **1**-PMMA indicated that the crystal structure of **1** is preserved in the composite membrane (Fig. S14†). The pure PMMA film did not exhibit any photomechanical behavior (Fig. S15†), indicating no photoinduced PMMA response. Upon perpendicularly exposing the backside of **1**-PMMA (5.0 × 40 mm<sup>2</sup>) to the light source (Fig. 4d and Video S5†), **1**-PMMA promptly bends away from the light source. The longer the irradiation time, the more severely the film bends. This was also observed at 405 and 365 nm (Fig. S16†).

To understand the kinematics and mechanism of photoinduced deformation, we quantified the bending angles ( $\varphi$ ) as a function of exposure time ( $t$ ) by recording a series of snapshot images. As shown in Fig. S17,† when exposing **1**-PMMA for 10 s, the bending angle reached 50°, and a bending angle of 305° was achieved within 90 s. However, small bending angles of 182° and 112° were observed for 405 and 365 nm, respectively, within 90 s. The correlation between the bending angle and irradiation time can be described according to the following equation:<sup>63</sup>

$$\varphi = \kappa(1 - \exp(-t/\tau)) \quad (1)$$

where  $\varphi$  denotes the bending angle,  $\kappa$  denotes the response time constant,  $t$  denotes the irradiation time, and  $\tau$  denotes the actuation response time coefficient. The  $\tau$  values were calculated to be 69, 26, and 71 for 450, 405, and 365 nm, respectively, indicating that the photomechanical response of **1**-PMMA is more sensitive to 405 nm. However, **1**-PMMA exhibited the largest  $\kappa$  value (417), suggesting the most thorough photocycloaddition reaction under the irradiation of 450 nm. As shown in Fig. 3f, **1** and **1a** exhibit notable absorption below approximately 400 nm, implying that 365 and 405 nm light would be hindered from penetrating deeply into the film. As

a result, a more thorough photocycloaddition reaction was achieved for 450 nm.

Interestingly, when the front side of **1**-PMMA was irradiated, the film exhibited a phototropic bending phenomenon (Fig. S18† and Video S6†) that was different from the photo-responsive behavior elicited by irradiation from the backside. To investigate this, scanning electron microscopy (SEM) and EDS mapping were conducted to reveal the powder distribution in the composite membrane. As shown in Fig. S19a,† **1** is mainly distributed in the backside of the membrane due to the gravity effect, while being relatively sparse in the front side.

This was further confirmed by the EDS mapping, in which the signals of In, Cl, and N were significantly stronger on the backside. Therefore, once exposed to 450 nm illumination, **1**-PMMA generated a more significant expansion on the backside (Fig. S19b†). Consequently, **1**-PMMA always exhibited the bending direction away from the backside, irrespective of the irradiation direction. The  $\kappa$  and  $\tau$  value was calculated to be 115 and 69, respectively, once the backside of **1**-PMMA was irradiated (Fig. S20†). Although the irradiation direction would not alter the sensitivity of **1**-PMMA toward light wavelength, the photocycloaddition ratio of **LH**<sup>+</sup> on the backside was smaller due to the limitation of light penetration, and thereby, a smaller bending rate was observed for front side irradiation.

By leveraging the sedimentation of **1** during the preparation of the film, we attempted to tilt the mold at a certain angle (Fig. S21†) to prepare **1**-PMMA with a concentration gradient. Interestingly, the film with a perpendicular concentration gradient (Fig. 4e and Video S7†) shows a curling responsive behavior due to the different degrees of volume expansion in both of the long sides. Once the concentration gradient was parallel to the long side of the belt film, no significant response was observed in the upper side of **1**-PMMA due to the lower concentration of **1** (Fig. 4f and Video S8†), while notable bending was observed for the lower section with a higher **1** concentration.

The reversible photochromic PL of **1a** encouraged us to explore its potential in anti-counterfeiting. As shown in Fig. 5a, a pine tree pattern composed of five groovy parts was filled with **1** powder. Upon 450 nm irradiation, the irradiated region exhibited a distinct photochromic PL due to the dimerization of **LH**<sup>+</sup>. Further irradiating at 254 nm caused a change in color from blue to green, while no change was observed for the blocked part. Thus, the pine tree pattern shows various emissive colors for different regions, which can be further erased or recovered upon 450/254 nm irradiation, greatly increasing the complexity and difficulty of decryption. Similarly, different emissive colors can be adopted for information storage.

As shown in Fig. 5b, the blue- and green-emissive dots are defined to be “0” and “1”. Based on binary ASCII code, the 4 × 8 matrix pattern gives a valid code of “Stpy”, which can be successfully erased upon 450 nm irradiation. Furthermore, given the rapid photomechanical response of **1**, several pieces of **1**-PMMA film were tailored to petal-like shapes to mimic the blooming process of a flower. As shown in Fig. 5 and Video S9,† once the handcrafted item was irradiated with 450 nm light, the **1**-PMMA films responded immediately and gradually bent away









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