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Thickness-dependent performance of antimony sulfide thin films as a photoanode for enhanced photoelectrochemical water splitting

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A two-step synthesis approach is employed for antimony sulfide thin films, which includes thermal evaporation followed by annealing in a sulfur atmosphere using chemical vapor deposition (CVD). The thickness of the films is systematically varied to study its impact on the material's properties. The orthorhombic crystal structure of each film is verified by Grazing Incidence X-ray Diffraction (GIXRD) analysis. Raman spectroscopy reveals thickness-dependent changes in the vibrational properties. Surface morphology and roughness are examined using atomic force microscopy (AFM) and field emission scanning electron microscopy (FESEM), with findings indicating that layer thickness significantly affects these surface characteristics. Energy-dispersive X-ray spectroscopy (EDS) and X-ray photoelectron spectroscopy (XPS) demonstrate that variations in film thickness influence the surface chemical composition and oxidation states. The Sb₂S₃ thin film with a thickness of 450 nm exhibited a band gap of 1.75 eV, indicating its potential for efficient light absorption. It also demonstrated a conductivity of 0.006 mA at an applied voltage of 1 V, reflecting its electrical transport properties. Furthermore, the film achieved a current density of 0.70 mA cm⁻², signifying enhanced charge transfer efficiency. These findings suggest that the 450 nm thick film offers an optimal balance of band gap, light absorption, and photocurrent density, making it the most suitable candidate for photoelectrochemical water-splitting applications.

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1. Introduction

Advancements in materials science are crucial for enhancing renewable energy technologies, particularly in the growth of resourceful, durable, and cost-effective materials for hydrogen energy production. A key area of focus is the improvement of photocatalysts and electrode materials to upsurge the effectiveness and scalability of water splitting. These resources must exhibit desirable properties such as nanostructures, low band gaps, and low resistivity to optimize their performance in photoelectrochemical (PEC) applications.¹ Significant research has been directed towards metal oxides, nitrides, and chalcogenides, each offering unique advantages and limitations. Meanwhile, metal oxides and nitrides possess high band gaps that limit their practical application.²⁻⁵ Chalcogenide compounds containing group 16 elements such as sulfur, selenium, and tellurium have emerged as promising alternatives due to their excellent optical properties, including high absorption

coefficients and low band gaps.^{6,7} Transition metal dichalcogenides, including materials like MoS₂, WS₂, MoSe₂, WSe₂, Sb₂S₃, and Sb₂Se₃ are particularly suited for water splitting because of their optimal band gaps (1–2 eV) and outstanding light absorption capabilities.^{8,9}

Among these antimony(III) sulfide (Sb₂S₃) films have attracted significant attention as promising alternative photocathodes for hydrogen generation in PEC cells. This interest stems from their distinct optoelectronic properties, including a narrow optical bandgap (E_g) of 1.7 eV and a high absorption coefficient (α) of 10⁴–10⁵ cm⁻¹ for photons with energy exceeding the bandgap.¹⁰ Additionally, Sb₂S₃ offers a high theoretical maximal photocurrent density of around 24.5 mA cm⁻² and an impressive theoretical solar-to-hydrogen efficiency of 28%, making it highly attractive for practical applications.¹¹ However, challenges such as insufficient absorption in thin films due to inadequate thickness, interference effects, and low material density, as well as issues related to inhomogeneity, internal scattering, and stress-induced defects in thicker films, can hinder absorption efficiency.¹²

Recent research has shown how important film thickness is in determining the performance of Ag₃SbS₃ thin films in PEC applications, with a 700 nm thick film demonstrating a photo-conversion efficiency of 0.68%, outperforming other materials.¹³ Additionally, research into the impact of photoanode

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thickness reveals that thinner photoanodes induce polarized states, increasing surface trapping and shortening the lifetime of photogenerated charges. Conversely, thicker photoanodes facilitate the removal of unstable surface layers, promoting oxygen evolution while preserving the integrity of the solid–solid interface.¹⁴

In this study, we explore how Sb_2S_3 thin film thickness influences morphology, elemental composition, and optical, electrical, and electrochemical properties in PEC water splitting. We aim to elucidate its role in charge transfer mechanisms, conductivity processes, and overall photoanode efficiency.

Our findings indicate that optimizing film thickness enhances charge transfer efficiency by stabilizing the interface and exposing active sites for oxygen evolution. This approach mitigates exciton recombination caused by deep-level defects in Sb_2S_3 , enhancing PEC performance. The optimized Sb_2S_3 photoelectrode achieves a photocurrent density of 0.7 mA cm^{-2} , surpassing the bare electrode (0.6 mA cm^{-2}).¹⁵ These findings provide a fundamental understanding of thickness-dependent charge transport mechanisms and offer a promising strategy for designing efficient Sb_2S_3 -based photoanodes for solar-driven water splitting. These results provide crucial insights into thickness-dependent charge transport and offer a promising strategy for designing efficient Sb_2S_3 -based photoanodes for solar-driven water splitting.

This work represents the first investigation into the role of Sb_2S_3 film thickness in PEC water-splitting applications. While Sb_2S_3 has been explored for various other applications, including photovoltaics,¹⁶ photosensitization of wide bandgap semiconductors,¹⁷ photocatalytic environmental remediation,¹⁸ battery storage,¹⁹ and even PEC solar cells.²⁰ Previously, it had been proposed to be an inactive water-splitting photocatalyst. In this work, we demonstrate how to optimize photocurrent density by varying film thickness. Despite facing challenges with photo-corrosion, we hypothesize that shielding Sb_2S_3 from degradation could improve its viability as a top cell absorber in devices, potentially through doping²¹ or heterojunction engineering.²² We identified three key areas for enhancement: morphology, stoichiometry, and optical properties. These improvements can be pursued by optimizing the film thickness, annealing in a sulfur atmosphere,²³ and applying doping techniques.²⁴

2. Experimental details

Antimony sulfide (Sb_2S_3) films were deposited by thermal evaporation using a fluorine-doped tin oxide-coated (FTO) glass substrate as the substrate and a molybdenum boat as the source holder. After ten minutes of ultrasonic cleaning in isopropyl alcohol, the substrates were dried with nitrogen purging before deposition. Cleaned substrate, and high-purity antimony sulfide (Sb_2S_3) powder (Thermo Fisher, 99.9%) was then placed into the thermal evaporation chamber. 5×10^{-6} mbar was the average pressure, which was maintained *via* liquid nitrogen, with a standoff distance of 16 cm between the source and the substrate. The substrate temperature was set to 150 °C during

Table 1 T1 to T4 film thicknesses as determined by a stylus profilometer

Sample name	Thickness (nm) \pm 10 nm
T1	350
T2	400
T3	450
T4	500

deposition. The deposition rate and film thickness were controlled and measured using a digital quartz crystal monitor. By setting the filament current to 85 A, the deposition rate was maintained between 3.5 and 4 \AA s^{-1} , with the final film thickness varying from approximately 350–500 nm. Following deposition, a custom chemical vapor deposition (CVD) apparatus was used to sulfurize (annealing in sulfur atmosphere) the films for 30 minutes at 400 °C under 900 mbar of vacuum. 100 mg of sulfur (Thermo Fisher, 99.9998% purity) was evaporated in a single-zone furnace with a nitrogen environment as defined in our previous work.²³ Following sulfurization, the films were cooled to room temperature under vacuum conditions. Film thicknesses were measured *via* a stylus profilometer (Bruker Dektak profilometer) and the experimental values are provided in Table 1.

3. Characterizations

Structural analysis of the antimony sulfide (Sb_2S_3) films was carried out using Grazing Incidence X-ray diffraction (GIXRD). Data was obtained with a Bruker AXS D8 Advance Diffractometer utilizing $\text{Cu K}\alpha$ radiation over an angular range of 10–60° with a scanning rate of $0.1^\circ \text{ min}^{-1}$ at room temperature. Raman spectroscopy was conducted in a backscattering configuration with a Nd: YVO_4 diode-pumped solid-state laser, employing a 532 nm excitation wavelength. The surface morphology was examined through scanning electron microscopy (SEM) (Zeiss EVO MA18) with an acceleration voltage of 10 kV, and the elemental composition was analyzed *via* energy-dispersive X-ray spectroscopy (EDS) (Oxford INCA X-act). X-ray photoelectron spectroscopy (XPS) measurements were performed with a SPECS (Germany) X-ray photoelectron spectroscope utilizing $\text{Al K}\alpha$ (1486.61 eV) radiation and applying a 13 kV anode voltage. The surface topography was investigated *via* atomic force microscopy (AFM) with the tapping mode Flex-Axiom AFM. Optical absorption spectra were obtained *via* a UV-visible spectrophotometer (Shimadzu UV-1900) over the range of 190–1100 nm. Photoluminescence measurements were carried out with a fluorescence spectrometer using a Nd: YVO_4 diode-pumped solid-state laser with an excitation wavelength of 420 nm.

The PEC performance of the photoelectrodes was measured using a CompactStat.h, IVIUM potentiostat connected to a standard three-electrodes PEC H-cell from Redox.me. The antimony sulfide thin film, platinum (Pt) wire, and Ag/AgCl electrode were used as working, counter, and reference



electrodes, respectively. A 0.1 M Na₂SO₄ aqueous solution (pH = 7) was used as the electrolyte. The potential relative to the reversible hydrogen electrode (RHE) was calculated *via* the appropriate conversion formula.

$$E_{\text{RHE}} = E_{\text{Ag}/\text{AgCl}} + E_{\text{Ag}/\text{AgCl}}^{\circ} + 0.059 \times \text{pH}$$

where $E_{\text{Ag}/\text{AgCl}}^{\circ} = 0.0197$ V vs. NHE at 28 °C.

The light source for PEC experiments was a 1-Sun LED Solar Simulator (Redoxme), providing a light intensity of 100 mW cm⁻². The films were exposed in a back side illumination geometry the exposed area was 0.5 cm². Photocurrent was recorded as a function of the voltage using linear sweep voltammetry (LSV) with a scan rate of 50 mV s⁻¹. Electrochemical impedance spectroscopy (EIS) was performed over a frequency range of 0.1 Hz to 100 kHz at 0.5 V sinusoidal voltage.

4. Results and discussion

4.1 Structural analysis

The Grazing Incidence X-ray diffraction (GIXRD) patterns of antimony sulfide (Sb₂S₃) are shown in Fig. 1 and are referenced to the standard JCPDS card (no. 42-1393). All samples, display an intense peak at 24.9° corresponding to the (1 3 0) orientation plane, revealing the orthorhombic stibnite structure of Sb₂S₃.

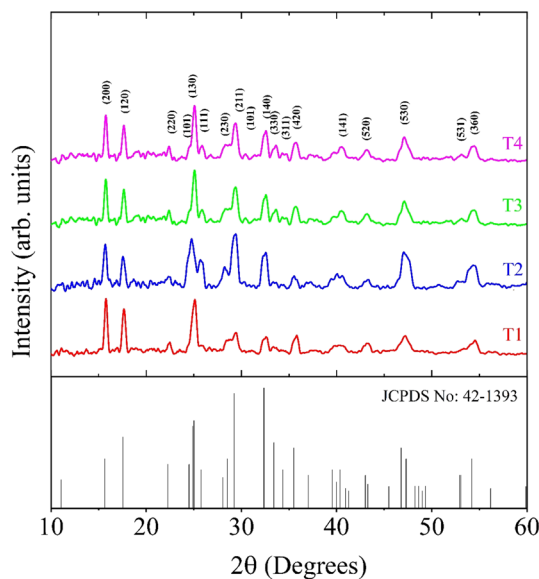


Fig. 1 GIXRD patterns of antimony sulfide films.

The enhanced intensity of the (1 3 0) peak with increasing film thickness (*i.e.*, longer deposition time) indicates preferential growth along the (1 3 0) crystallographic direction. This indicates that as the deposition time increases, the film growth mechanism favors alignment along the (1 3 0) plane, resulting in enhanced crystallographic orientation and increased peak intensity for this plane in the grazing incidence X-ray diffraction pattern.²⁵ Additionally, sample T2 showed a peak shift to a lower angle. This observed shift can be attributed to several factors. Compositional gradients, which involve variations in the film's composition, can alter the lattice parameters and affect the peak positions. Additionally, internal strain and variation in the lattice spacing resulted in a shift in the XRD peaks. Together, these factors reflect structural changes within the material, as indicated by the peak shift.²⁶

Table 2 lists the crystallographic parameters of the antimony sulfide films with various thicknesses. The data show a trend where increasing film thickness generally results in larger crystallite sizes and reduced microstrain. The observed decrease in microstrain suggests fewer lattice imperfections, as supported *via* the reduction of grazing incidence X-ray diffraction (GIXRD) peaks and enhance the crystallite size with increasing film thickness.²⁷ However, an anomaly is noted in sample T2, where the crystallite size unexpectedly decreases and the microstrain increases. This increase in microstrain for sample T2 may have contributed to a shift toward lower diffraction angles in the XRD patterns.²⁸ Furthermore, the lattice parameters vary with thickness, and for samples T3 and T4, these parameters closely match the expected values for antimony sulfide films.

Fig. 2 presents the Raman spectra of the antimony sulfide films, which show distinct modes at 109, 125, 146, 186, 238, 280, 300, and 495 cm⁻¹.^{29,30} The mode at 495 cm⁻¹ is assigned to the FTO substrate.³¹ Specifically, the modes at 109, 154, 186, and 280 cm⁻¹ correspond to the A_g modes, while the modes at 238 and 300 cm⁻¹ are related to the B_{1g}/B_{3g} modes. We can see the variation in the intensity with respect to the thickness. This variation in Raman intensity with variation in crystallite size occurs because larger crystallites have better scattering efficiency and fewer structural defects. In larger crystallites, more photons interact with the lattice, leading to stronger Raman signals. Additionally, reduced defects and less phonon confinement in larger crystals result in sharper and more intense Raman peaks.³² Here, we can observe that sample T2 has the highest intensity and shifts toward the lower side

Table 2 Crystal parameters of antimony sulfide films

Sample	Crystallite size <i>D</i> (nm)	Microstrain, ϵ ($\times 10^{-3}$)	Dislocation density, δ ($\times 10^{15}/\text{m}^2$)	Interplanar distance, d_{hkl} (Å)	Lattice parameters (Å)		
					<i>a</i> (11.22)	<i>b</i> (11.31)	<i>c</i> (3.83)
T1	18.0	8.9	3.0	3.58	11.3	11.47	3.87
T2	14.6	10.6	4.6	3.59	11.01	11.23	3.74
T3	19.4	8.2	2.6	3.64	11.18	11.38	3.78
T4	19.8	8.07	2.5	3.64	11.24	11.42	3.83



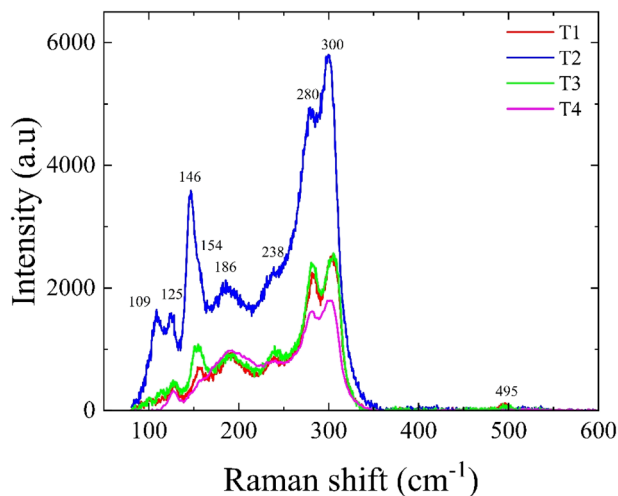


Fig. 2 Raman spectra of antimony sulfide films.

compared with the other samples. This may be due to the microstrain, which we observed *via* XRD. Microstrain generally affects Raman scattering by causing peak shifts due to distortions in the lattice structure. While it can alter the Raman spectra, it does not typically enhance the intensity directly. Instead, factors such as an optimal crystallite size and fewer defects often play a larger role in increasing the intensity.³³ XRD and Raman analyses showed that all the obtained samples were single-phase antimony sulfide.

4.2 Morphological and topographical analysis

Fig. 3 shows the morphology obtained through field emission scanning electron microscopy (FESEM). In samples T2 and T4

compact formation of grains is evident. In contrast, samples T1 and T3 display island-like structures. The noted changes in morphology with varying film thickness are likely attributable to alterations in grain size as the thickness increases due to extended deposition times, which allow grains to grow and coalesce. This variation in grain size could also be influenced by the change in adhesion properties of the molecules at different thicknesses, leading to the observed morphological differences.³⁴

Fig. 4 represents images of atomic force microscopy (AFM) that highlight the variations in surface morphology across the different samples. As the thickness of the samples changes, there is a corresponding variation in surface roughness, as indicated by the roughness parameters [average roughness (R_a) and root mean square roughness (R_q)] provided in Table 3. The notable differences in roughness are likely attributable to the morphological changes observed in Fig. 3. Specifically, sample T2 is rougher than samples T1, T3, and T4. The increased roughness at T2, as discussed in the morphological analysis, may be attributed to variations in deposition time, potentially influencing the adhesion properties of molecules at different thicknesses, thereby affecting film roughness.³⁵ It is further confirmed that the shifts toward lower angles observed in the XRD patterns, compared to other samples, could be attributed to the presence of the highest microstrain, as revealed by XRD analysis.

4.3 Compositional analysis

Energy-dispersive spectroscopy (EDS) was employed to conduct elemental analysis on the deposited antimony sulfide films. Table 4 provides the atomic percentage compositions of the films. Notably, the Sb to S (Sb:S) atomic ratio exhibited

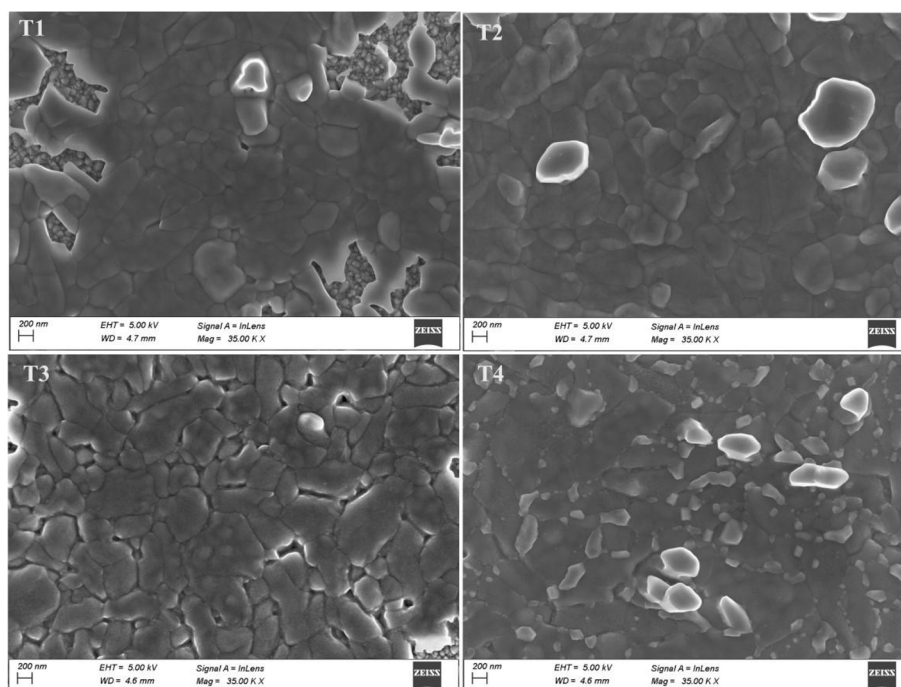


Fig. 3 FESEM images of antimony sulfide films.



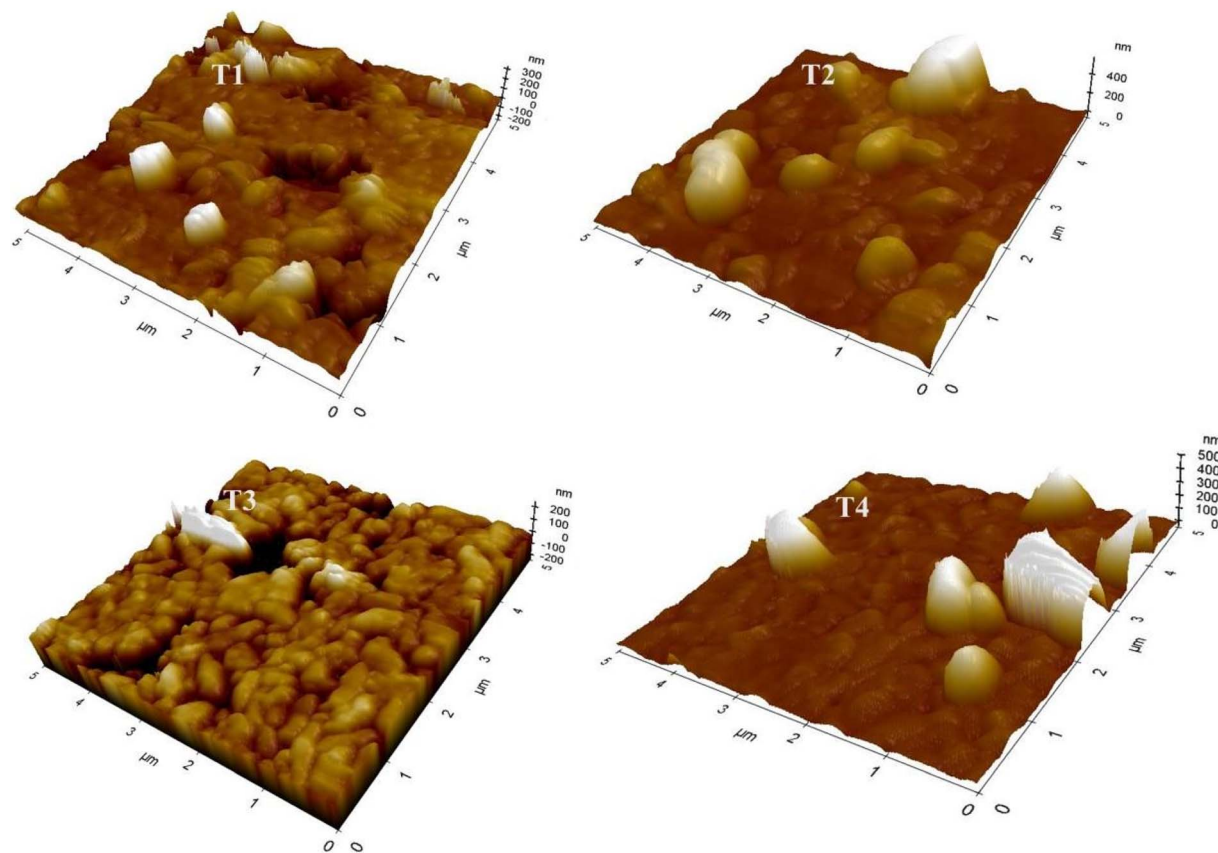


Fig. 4 AFM images of antimony sulfide films.

Table 3 Surface roughness parameters of antimony sulfide films

Sample name	R_a (nm) \pm 2 nm	R_q (nm) \pm 2 nm
T1	26.0	45.0
T2	72.6	103.4
T3	37.1	59.6
T4	22.4	33.6

variability corresponding to changes in film thickness. Specifically, sample T3 had an Sb:S ratio of 2:3.5, which closely approximates the expected stoichiometric ratio of 2:3.

X-ray photoelectron spectroscopy (XPS) was employed to analyze the chemical composition of the surface of the samples. The binding energy peaks were deconvoluted *via* a Gaussian-

Table 4 Atomic percentages of antimony sulfide films

Sample name	Atomic percentage (%)		
	Sb (40)	S (60)	S/Sb
T1	20	80	4
T2	33	67	2
T3	25	75	3
T4	34	66	1.9

Lorentzian function. Before data analysis, all spectra were referenced to the C 1s peak at 284.6 eV for calibration Fig. 5(a) displays the survey spectra of antimony sulfide films of various thicknesses. Fig. 5(b) presents the Sb 3d core spectra, illustrating a doublet of peaks at 528.9 eV ($3d_{5/2}$) and 538.2 eV ($3d_{3/2}$).¹⁵ The 9.4 eV splitting of the Sb 3d doublet indicates the presence of trivalent antimony (Sb^{3+}) in Sb_2S_3 .³⁶ The area ratio of the $3d_{5/2}$ peak to the $3d_{3/2}$ peak is approximately 1.5 (3:2), which is consistent with the expected ratio.³⁷ The two shoulder peaks observed at 529.2 eV and 530.9 eV depict the Sb $3d_{5/2}$ and Sb $3d_{3/2}$ of Sb_2O_3 , respectively, also showing a peak separation of 9.4 eV and an area ratio of 1.5. The peak at 532 eV reveals the presence of oxygen (O) in its 1s state.³⁸ These oxygen impurities are likely introduced during the sulfurization process due to residual oxygen in the tube furnace.³⁹ Fig. 5(c) presents the S 2p core spectra of the antimony sulfide films. In the sulfur spectrum, two overlapping peaks at 160.8 eV and 161.9 eV correspond to the $2p_{3/2}$ and $2p_{1/2}$ states, respectively. The area under the curve ratio of $2p_{3/2}$ to $2p_{1/2}$ is 2:1, with an energy separation of 1.1 eV, indicating the S^{2-} oxidation state.^{40,41} The S 2p peaks at 163 eV and 164 eV (S $2p_{3/2}$ and S $2p_{1/2}$) correspond to sulfur bound to oxygen, indicating the presence of sulfur oxides and sulfates.⁴² From both core spectra, it is evident that higher binding energy generally correlates with higher oxidation states of the material.⁴⁰ The XPS results demonstrate that these oxide states are present on the surface of the sample, likely because of



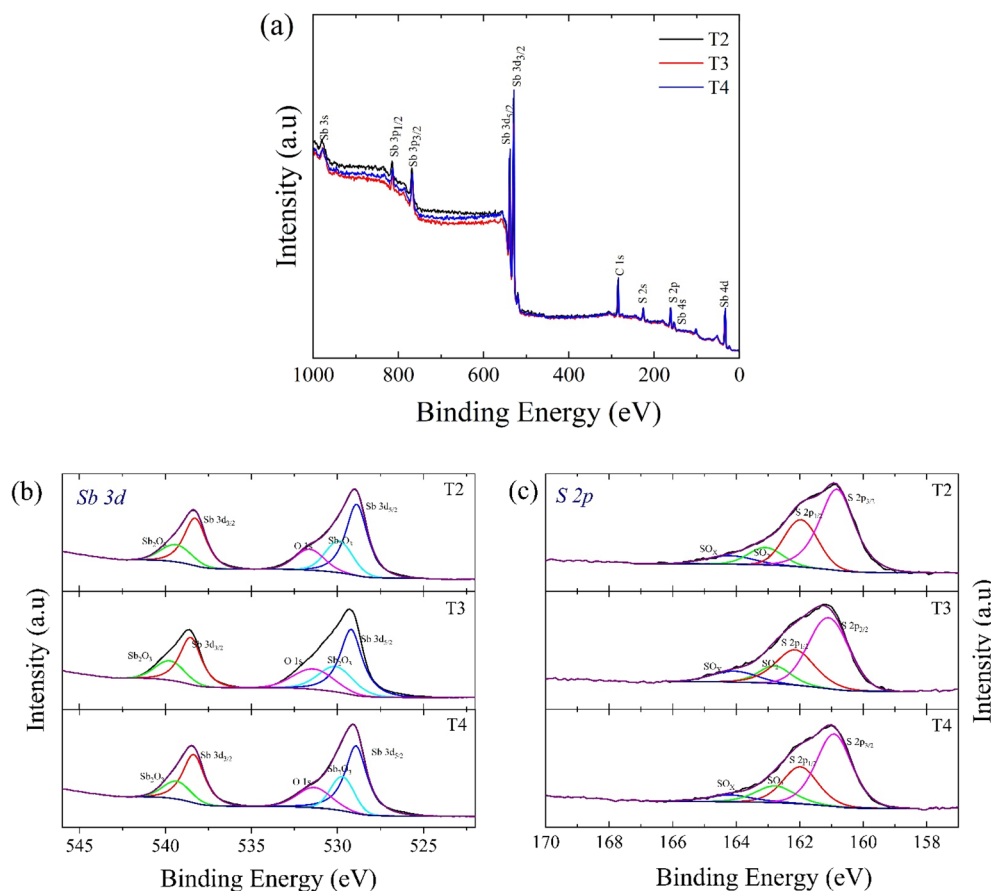


Fig. 5 (a) Survey spectra, (b) & (c) core spectra of Sb 3d and S 2p of the antimony sulfide film.

surface oxidation. This finding has been confirmed by the XRD pattern and Raman spectra showing no oxide peaks.⁴³

4.4 Optical analysis

Fig. 6 represents (a) absorbance spectra, (b) absorption coefficient spectra, (c) Tauc plots, and (d) photoluminescence spectra of antimony sulfide films. Table 5 provides the band gap values for each sample as deduced from Tauc plot fit. The analysis shows that both the absorbance and absorption spectra in Fig. 6(a) and (b) exhibit variations related to the thickness of the samples. Band gap values were obtained from the Tauc plots, as illustrated in Fig. 6. These are listed in Table 5 for each sample.

The observed variations in the absorbance and absorption coefficients are likely due to differences in the stoichiometric ratio, as indicated in Table 4. A higher atomic concentration of sulfur (T1: S/Sb = 4) leads to increased bandgap (T1 1.8 eV) because higher sulfur concentration strengthens the bonds within the material, which can shift the energy levels of the electronic states, resulting in a larger band gap.^{44,45} The band gap variations across the samples are attributed primarily to the incorporation of sulfur into the films and the presence of sulfur vacancies. These characteristics may cause additional electronic states to arise slightly above the valence band, thereby altering the absorption edge.⁴⁶ Additionally, differences in the grain size and stoichiometric ratio, as supported by the FESEM images

and EDS results in Table 4, contributed to the observed variations in the absorbance, absorption coefficient, and band gap. The stoichiometric ratio plays a crucial role in determining a material's band gap by introducing defects, altering the phase composition, and influencing atomic bonding. Deviations from ideal stoichiometry can create localized states, modify doping effects, and cause charge imbalances, leading to band gap narrowing or broadening, thus impacting the material's electronic properties.

The photoluminescence analysis was conducted to examine the optical properties of the samples. Fig. 6(d) presents the photoluminescence spectra of antimony sulfide films excited at a wavelength of 420 nm.⁴⁷ Upon excitation, emission peaks in the 460 nm range were observed, indicating blue luminescence.^{48,49} This blue emission is attributed to donor-acceptor ion transitions, commonly linked to defects or deep trap states.²⁹ A variation in the emission intensity with increasing film thickness was observed. These changes in emission intensity are likely due to a combination of physical thickness effects and stoichiometric variations that alter the material's electronic structure and optical behavior.⁵⁰

Increasing the film thickness causes variations in both the stoichiometric ratio and crystallite size, which subsequently affect the emission intensity. According to the properties of the donor and acceptor ions, where Sb^{3+} represents the donor ion



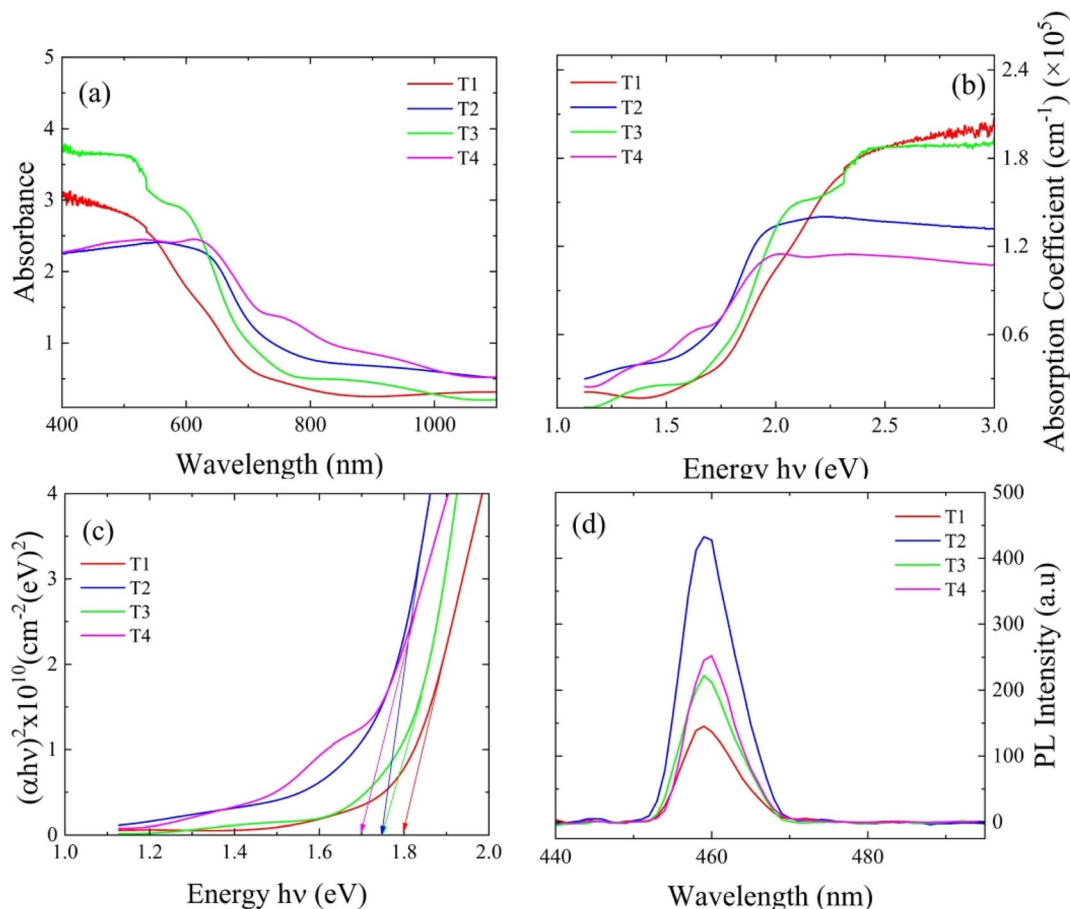


Fig. 6 (a) Absorbance spectra, (b) absorption coefficient spectra, (c) Tauc plots, and (d) photoluminescence spectra of antimony sulfide films.

Table 5 Band gap values of antimony sulfide films

Sample name	Band gap (eV)
T1	1.80
T2	1.75
T3	1.75
T4	1.70

and S^{2-} the acceptor ion, the films were classified into two types: antimony-rich (Sb-rich) and sulfur-rich (S-rich) films. Sb-rich films exhibit three electron trap states, termed donor defects, while S-rich films display two hole trap states, known as acceptor defects.^{51,52}

The primary defects identified include Sb interstitial (Sb_i), sulfur vacancy (V_s), antimony antisite (Sb_s), antimony vacancy (V_{Sb}), and sulfur antisite (S_{Sb}) defects. In Sb-rich Sb_2S_3 , the increase in V_s is attributed to sulfur deficiency. As a result, excess Sb atoms tend to occupy V_s sites rather than interstitial positions, due to the lower formation energy of Sb_s compared to Sb_i . Consequently, Sb-rich Sb_2S_3 exhibits a higher concentration of V_s and Sb_s defects.

In S-rich Sb_2S_3 , sulfur atoms infiltrate the lattice and fill V_s sites. Due to the higher formation energies of Sb_s and Sb_i defects

in a sulfur-rich environment, the formation of these defects is suppressed. Instead, the S-rich condition favors the formation of significant amounts of V_{Sb} , driven by the reduced formation energy. To maintain structural stability, some sulfur atoms may occupy V_{Sb} sites, leading to the formation of Sb_s antisite defects.⁵²

4.5 Electrical analysis

Fig. 7 presents the (a) linear and (b) semilogarithmic plots of the current-voltage (I - V) characteristics of the deposited antimony sulfide films. The data reveal that an increase in film thickness from T1 to T3 correlates with a significant increase in conductivity. Specifically, the T3 sample revealed a current of 0.006 mA at 1 V. This trend can be attributed to the corresponding decrease in film resistance by increasing the thickness.^{53,54} However, for thicknesses exceeding 450 nm, a decrease in conductivity was observed, as seen in sample T4. This reduction in conductivity is attributed to an increase in grain boundaries, as evidenced by SEM images. These structural imperfections lead to greater scattering and recombination of charge carriers, which in turn diminishes the overall conductivity.⁵⁵

4.6 Electrochemical analysis

Fig. 8(a) shows the cyclic voltammetry (CV) and Nyquist plots of the antimony sulfide films. In the CV curve shown in Fig. 8(a),



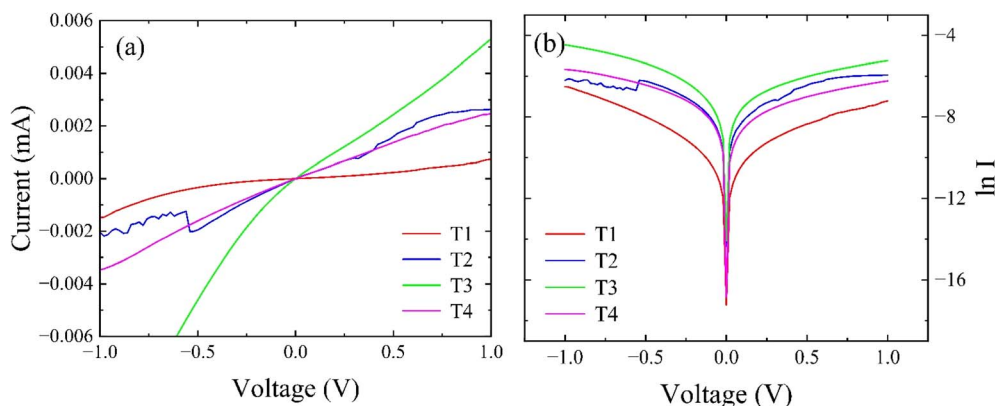


Fig. 7 (a) Linear and (b) semilog plots of the I - V characteristics of antimony sulfide films.

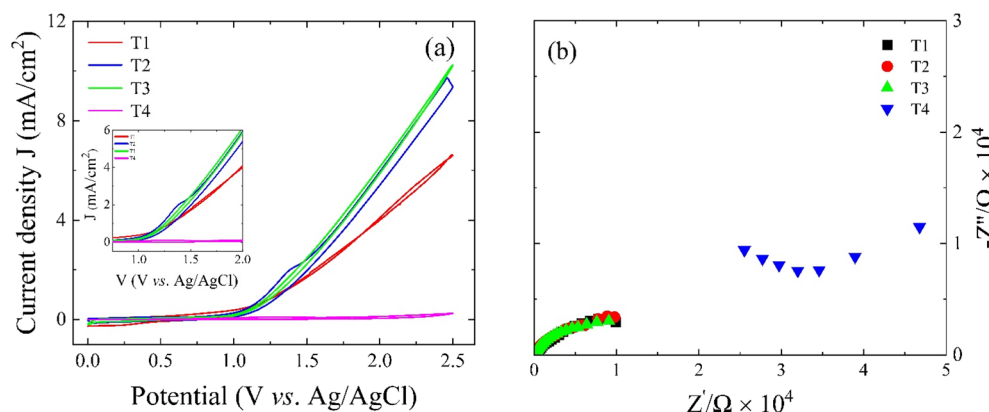


Fig. 8 (a) Cyclic voltammograms and (b) Nyquist plots of deposited antimony sulfide films.

an oxidation peak is observed at 1.3 V *versus* Ag/AgCl, with the peak intensity decreasing as the film thickness increases.⁵⁶ Initially, the current density increases with increasing film thickness (from T1 to T3) but for T3 it decreases. The oxidation peak in thinner films is attributed to the fact that a species undergoes oxidation during the electrochemical measurement. The variation in film current density concerning thickness is likely due to changes in conductivity, as depicted in Fig. 8. This variation in conductivity is influenced by defects and grain boundaries, which lead to increased scattering and recombination of charge carriers, ultimately affecting the complete conductivity of the films. The Nyquist plot in Fig. 8(b) representing the imaginary part of the impedance (Z'') as a function of the real part of the impedance (Z'), provides insights into the electrochemical properties at the interface electrode/electrolyte. This interface can be modeled with an equivalent circuit containing an $R_{SC}C_{SC}$ parallel circuit for the semiconductor part in series with another $R_H C_H$ parallel circuit corresponding to the Helmholtz layer formed inside the electrolyte near the electrode surface. Specifically, the internal resistance depicted in the plot represents the overall resistance encountered in the course of the charge transfer process. The observed impedance behavior specifically highlights the existence of charge transfer resistance, highlighting the difficulty of electron transfer at the

electrode–electrolyte interface.⁵⁷ As sample thickness increases, the charge transfer resistance decreases, likely because of the higher loading of antimony sulfide. Notably, sample T3 has a lower charge transfer resistance than the other samples, likely resulting from improved conductivity and better aggregation of antimony sulfide within the film.⁵⁸ The increase in charge transfer resistance observed in sample T4, as depicted in Fig. 8, may be due to the combined effects of low absorption coefficient, reduced surface roughness, and decreased electrical conductance at higher thicknesses.

The photocurrent density of antimony sulfide films varies with thickness due to changes in electrical conductivity, charge transport efficiency, and recombination processes. Fig. 9 illustrates the current density *versus* potential (J - V) characteristics of antimony sulfide films with varying thicknesses, measured under illuminated and dark conditions. Sample T1 exhibits a photocurrent density of 0.28 mA cm⁻² at 0.2 V *vs.* RHE. A progressive increase in current density is observed for samples T2 and T3, attributed to enhanced electrical conductivity and improved charge transport as the film thickness increases. This increase in thickness enhances charge carrier availability, facilitating efficient charge transfer to the electrode–electrolyte interface and resulting in a higher photocurrent density from T1 to T3. However, beyond an optimal thickness (T3), further



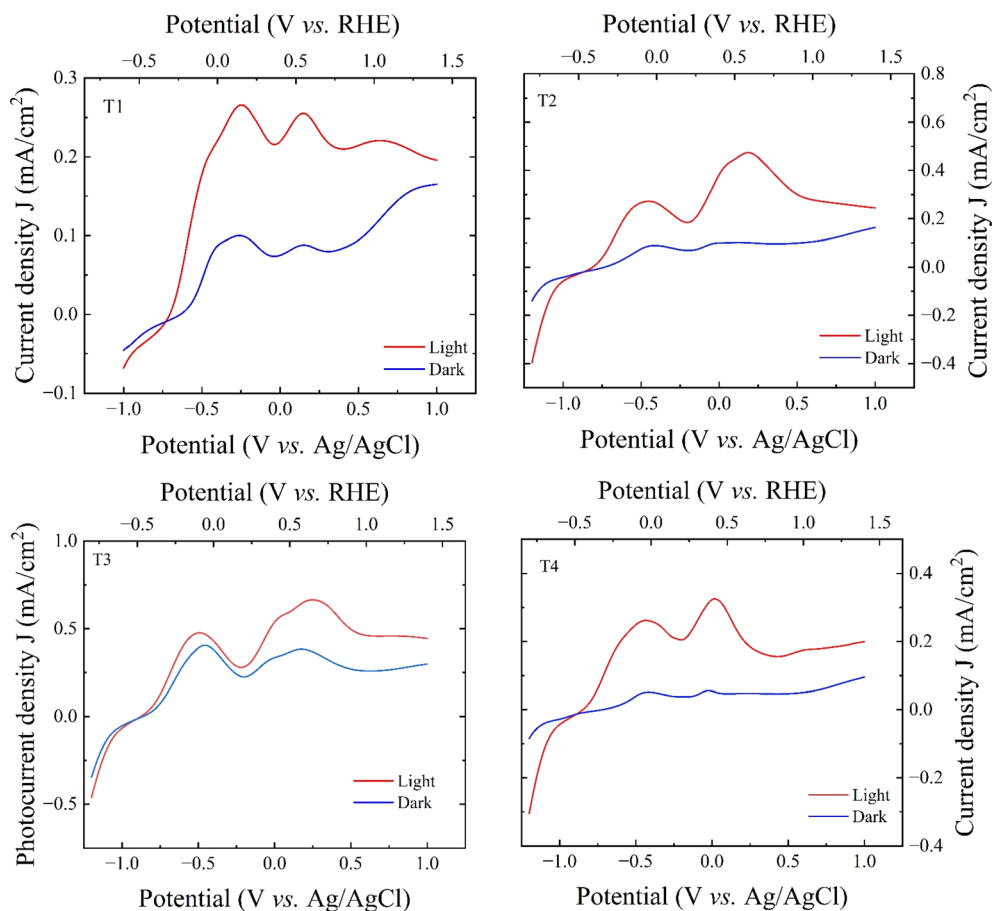


Fig. 9 LSV curves of antimony sulfide films.

film growth (T4) leads to excessive charge carrier recombination due to longer diffusion paths and increased internal resistance, ultimately reducing the photocurrent. This decline in current density in sample T4 is likely due to excessive film thickness, which promotes recombination and impedes conductivity, as evident from Fig. 7(a). Furthermore, the J - V curves exhibit two distinct humps, which may be associated with the oxidation of the material during PEC measurements. Among all samples, T3 achieves the highest photocurrent density of 0.70 mA cm^{-2} , which can be correlated with its superior electrical conductivity, as illustrated in Fig. 7. Enhanced conductivity generally improves charge carrier mobility, thereby accelerating charge transfer kinetics at the electrode–electrolyte interface. Consequently, this efficient charge transport contributes to an overall increase in photocurrent density, further supporting the optimization of film thickness for superior PEC performance.⁵⁹ However, during the experiment, substantial degradation of the samples was observed, likely due to electrode–electrolyte interactions leading to material corrosion, which adversely affects stability.^{60–62} Strategies such as elemental doping or the incorporation of heterojunctions may be effective in mitigating these stability issues while simultaneously improving charge transfer kinetics.

5. Conclusions

Thermally evaporated antimony sulfide thin films deposited on FTO substrates exhibited an orthorhombic stibnite structure across all samples, with crystallite size and microstrain varying with film thickness. Film thickness significantly affected vibrational properties, particularly the T_2 mode, where the 400 nm thick sample showed the highest intensity and peak shifts attributed to microstrain. Surface morphology analysis revealed that increasing thickness altered surface roughness, correlating with changes in grain size. Compositional analysis using EDS and XPS confirmed thickness-dependent stoichiometric variations and the presence of surface oxides. Among the investigated samples, the 450 nm thick film (T3) demonstrated optimal optoelectronic properties, with a band gap of 1.75 eV, high electrical conductance (0.006 mA at 1 V), and low charge transfer resistance. The obtained photocurrent density of 0.70 mA cm^{-2} under illumination highlighted its superior charge transport characteristics. However, surface oxidation remains a critical factor affecting long-term photocurrent stability during PEC operation. Addressing these oxidation-related challenges is essential for enhancing the durability of the film and the stability of the sample. The distinct performance of the 450 nm thick film positions it as a strong candidate for



photoelectrochemical water splitting, with further optimization expected to improve its effectiveness as a photoanode in PEC applications.

Data availability

The most data pertaining to this article is mentioned in the article itself. Any missing data can be provided on a reasonable request to the corresponding author.

Conflicts of interest

There are no conflicts of interest to declare.

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