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# A DFT study for hydrogen storage application on pristine magnesium dicarbide ( $\text{MgC}_2$ ) monolayer

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The hydrogen storage potential of pure  $\text{MgC}_2$  was systematically investigated using density functional theory (DFT) calculations. The phonon dispersion and *ab initio* molecular dynamics (AIMD) simulations confirmed the dynamic and structural stability of  $\text{MgC}_2$ , reinforcing its suitability as a promising hydrogen storage material. The electronic structure analysis revealed that pure  $\text{MgC}_2$  exhibits semiconducting behavior with a band gap of 0.25 eV, and transforms into a metallic state upon hydrogen adsorption. Hydrogen molecules were adsorbed onto the  $\text{MgC}_2$  surface via physisorption, with an average adsorption energy of 0.286 eV, indicating moderate binding strength suitable for reversible hydrogen storage. Hirshfeld charge analysis demonstrated that  $\text{MgC}_2$  transfers 0.041 e, 0.139 e, and 0.259 e to 1, 4, and 8 hydrogen molecules, respectively, highlighting charge redistribution upon adsorption. The calculated hydrogen storage capacity of 2.05% suggests a feasible adsorption mechanism. Additionally, AIMD simulations at 400 K confirmed that hydrogen adsorption does not induce significant distortions in the  $\text{MgC}_2$  framework, further validating its thermal and mechanical stability. These findings underscore the potential of  $\text{MgC}_2$  as an efficient hydrogen storage material for sustainable energy applications, offering a promising pathway for the development of next-generation clean energy technologies.

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## 1. Introduction

One of the main and most efficient ways to address the enormous rise in energy consumption and environmental protection is the development of clean, sustainable energy sources like wind, solar, and hydrogen.<sup>1–4</sup> However, since hydrogen is not readily available in nature, it must be stored for use in transportation, as well as other mobile and stationary applications.<sup>5,6</sup> There are two types of hydrogen storage technologies: physical-based and material-based. The first category consists of currently in-use technologies that require specialized, big-volume containers to store hydrogen as compressed gas, cold/cryocompressed gas, and liquid.<sup>7</sup> The second method involves chemical reactions within materials at the atomic or molecular level. It is still quite difficult to find a substance that can safely, effectively, and affordably store hydrogen for real-world applications.<sup>8,9</sup> Recently, advancements in hydrogen storage technology have made it possible to store a substantial amount of

hydrogen fuel safely, which has tremendously benefited the development of cars that run on hydrogen. Over the years, many efforts have been made to produce pure hydrogen ( $\text{H}_2$ ) without harming the environment. According to the standards of the U.S. Department of Energy (DOE), a material used for hydrogen storage must have a storage gravimetric density of more than 5.5 wt%.<sup>10</sup>

In recent years, there has been enormous interest in nano-structured materials to store hydrogen by an adsorption mechanism.<sup>11,12</sup> These factors have led to the proposal of many materials, including  $\text{MgH}_2$ ,  $\text{Mg}(\text{BH}_4)_2$ , and  $\text{LiBH}_4$ , for the solid-state storage of hydrogen. Despite having a high gravimetric density of hydrogen, these materials exhibit significant dissociation temperatures and thermodynamic stability.<sup>13,14</sup> The hydrogen storage capacity and recovery efficiency of carbon-based nanomaterials, such as graphene, are significantly influenced by lithium- and sodium-decorated inorganic graphenylene.

Although surface adsorption of  $\text{H}_2$  is inefficient for hydrogen release, it is appropriate for improving the quantity of storage. Additionally, to keep the reservoir pressure high once  $\text{H}_2$  has been retrieved, an absorbent gas, such as  $\text{N}_2$ ,  $\text{CH}_4$ , or  $\text{CO}_2$ , is required. This leads to competitive adsorption between  $\text{H}_2$  and the cushion gas, significantly affecting  $\text{H}_2$  recovery and purification in the recovered fluids.<sup>15</sup> Consequently, none of them is a good option. On the other hand, magnesium hydride ( $\text{MgH}_2$ ) has attracted significant attention due to its high volumetric

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capacity (110 g L<sup>-1</sup>) and gravimetric density (7.6 wt%), non-toxicity, good reversibility, and vast natural reserves. Furthermore, when supplying mixed gases free of harmful components, materials like MgH<sub>2</sub> show enormous potential for hydrogen filtration and selective storage.<sup>16,17</sup> Qin *et al.* conducted first-principles calculations to investigate the structural, mechanical, electronic, dynamic, thermodynamic, optical, and hydrogen storage properties of XMoH<sub>3</sub> (X = Na, K, Rb, Cs). Their findings confirm the thermodynamic, mechanical, and dynamic stability of these perovskites based on formation energy, elastic constants, and phonon spectra.<sup>18</sup>

Recent investigations have explored the potential of newly proposed 2D MXene materials, including Sc<sub>2</sub>C, Cr<sub>2</sub>C, Ti<sub>2</sub>C, Ti<sub>2</sub>N, and V<sub>2</sub>C for hydrogen storage applications. Liu *et al.* investigated the use of V<sub>2</sub>C (vanadium carbide) and Ti<sub>3</sub>C<sub>2</sub> (titanium carbide) MXenes in combination to improve the capacity of MgH<sub>2</sub> to absorb and desorb hydrogen. By incorporating 2V<sub>2</sub>C/Ti<sub>3</sub>C<sub>2</sub>, the activation energy for the hydrogen release process of magnesium hydride (MgH<sub>2</sub>) was successfully reduced by 36%, resulting in a reversible capacity of 6.3%.<sup>19</sup>

The SiB<sub>2</sub> monolayer, recognized for its lightweight nature, high carrier mobility, and exceptional stability, holds significant potential as an electrode material. Its structure is similar to that of MgB<sub>2</sub>, which has been extensively studied.<sup>20,21</sup> The MgB<sub>2</sub> has been investigated for applications such as hydrogen storage and alkali metal ion batteries. The hydrogenation of bulk MgB<sub>2</sub> requires high pressure (90–100 MPa) and temperature (390–400 °C) and is time-intensive. Li *et al.* studied nanoscale MgB<sub>2</sub> hydrogenation *via* mechanical milling, though the reaction pathway and intermediate states remain unclear.<sup>22</sup> Theoretical studies suggest that Li decoration can enhance hydrogen storage capacity to 7.26 wt%,<sup>23</sup> while Na and Li addition improves H<sub>2</sub> storage efficiency in C<sub>2</sub>O by 13.08 wt%.<sup>24</sup> In our previous study, we employed DFT to investigate optimized structures, thermodynamic properties, and H<sub>2</sub> storage capacities of MgB<sub>2</sub>. The results highlight a subtle charge transfer from Li/Na/K to the MgB<sub>2</sub> monolayer, enhancing its electro-positive nature and improving electrostatic interactions with H<sub>2</sub> molecules. The maximum H<sub>2</sub> adsorption occurs with nine H<sub>2</sub> molecules for Li and eight H<sub>2</sub> molecules for Na and K. The corresponding adsorption energies range from −0.24 to −0.21 eV for Li, −0.22 to −0.20 eV for Na, and −0.25 to −0.20 eV for K.<sup>25</sup> Previously, we have investigated various 2D materials and transition metal chalcogenides (TMDs) for alkali and alkaline earth metal-ion batteries, focusing on their structural, adsorption, electronic properties, open-circuit voltage (OCV), theoretical capacity, and thermodynamic stability.<sup>26–31</sup>

From the above discussion, it is clear that a thorough analysis of pristine MgC<sub>2</sub>'s capacity to store hydrogen, taking van der Waals correction parameters into account, has not yet been carried out. This gap inspired us to conduct a DFT calculation from the ground up to examine the hydrogen-adsorbed MgC<sub>2</sub> system's electrical characteristics, binding/adsorption energies, and desorption temperatures.

## 2. Computational detail

The Amsterdam density functional (ADF-BAND) software was used in the DFT framework, and Slater Type Orbitals (STO) was primarily used to create molecular orbitals for every structure.<sup>32</sup> Based on the framework offered by the Perdew–Burke–Ernzerhof (PBE) functional, the generalized gradient approximation (GGA) was used throughout the study to compute the exchange–correlation energies required for structure optimization. Given that van der Waals forces are the main interaction forces between H<sub>2</sub> and the substrate material, we addressed the shortcomings of DFT by integrating the Grimme semi-empirical correction technique (DFT-D3) throughout the calculations in handling dispersive interactions. Additionally, our method used double-zeta polarized (DZP) basis sets and excellent numerical quality, together with settings that permitted non-frozen cores.<sup>33</sup> For fully relaxed geometries, we impose strict requirements, with an energy optimization threshold of 10<sup>−5</sup> eV. We maintained 10<sup>−3</sup> Å as the step convergence criterion and 0.02 eV Å<sup>−1</sup> as the gradient convergence criterion. Hirshfeld charge analysis was used to probe the arrangement of electrical ions between the absorption setups and optimum substances.

The adsorption energy of hydrogen molecules on pure MgC<sub>2</sub> is calculated by following the formula.

$$E_{\text{ads}} = \frac{E_{\text{MgC}_2} + nE_{\text{H}_2} - E_{(\text{MgC}_2 - n\text{H}_2)}}{n} \quad (1)$$

where  $E_{\text{MgC}_2}$  denotes the total energy of MgC<sub>2</sub>,  $E_{\text{MgC}_2} + nE_{\text{H}_2}$  denotes the energy of hydrogen adsorbed MgC<sub>2</sub>, and  $n$  denotes the amount of hydrogen molecule,  $E_{\text{H}_2}$  is the energy of a single H<sub>2</sub> molecule.<sup>34</sup>

To determine hydrogen adsorption on the MgC<sub>2</sub> nanosheet, we calculated the binding energy ( $E_{\text{BE}}$ ) of the H<sub>2</sub> molecule using the following relation<sup>35</sup>

$$E_{\text{BE}} = E_{\text{MgC}_2} + nE_{\text{H}_2} - E_{n\text{H}_2 - \text{MgC}_2} \quad (2)$$

where,  $E_{\text{MgC}_2}$ ,  $nE_{\text{H}_2}$ ,  $n$ , and  $E_{n\text{H}_2 - \text{MgC}_2}$  are the total energies of pristine MgC<sub>2</sub> nanosheet, an isolated H<sub>2</sub> molecule,  $n$  is the number of H<sub>2</sub> molecules and MgC<sub>2</sub> nanosheet with adsorbed H<sub>2</sub> molecules, respectively.

To determine the gravimetric storage capacity, the following formula was used:

$$\text{wt}\% = \frac{n \times M_{\text{H}_2}}{n \times M_{\text{H}_2} + M_{\text{MgC}_2}} \times 100\% \quad (3)$$

where  $M_{\text{H}_2}$  is the molecular mass of the H<sub>2</sub> molecule,  $M_{\text{MgC}_2}$  molecular mass of the host MgC<sub>2</sub>,  $n$  is the number of adsorbed H<sub>2</sub> molecules.

The desorption temperature ( $T_{\text{d}}$ ) of hydrogen-adsorbed systems is crucial for practical applications. We used the Van't Hoff equation to analyze the desorption temperature.<sup>35,36</sup>

$$T_{\text{D}} = \frac{E_{\text{ads}}}{k_{\text{b}}} \times \left( \frac{\Delta S}{R} - \ln P \right)^{-1} \quad (4)$$



where  $k_b$  is Boltzmann constant ( $1.380 \times 10^{-23} \text{ J K}^{-1}$ ),  $\Delta S$  represents the difference for hydrogen entropy switching from gaseous to the liquid state ( $75.44 \text{ J mol}^{-1} \text{ K}^{-1}$ ),  $R$  is the universal gas constant ( $8.314 \text{ J mol}^{-1} \text{ K}^{-1}$ ), and equilibrium pressure is 1 atm.

The desorption energy ( $E_{\text{des}}$ ) is calculated as follows:

$$E_{\text{des}} = E_{\text{H}_2} + E_{\text{MgC}_2 + (n-1)\text{H}_2} - E_{\text{MgC}_2 + n\text{H}_2} \quad (5)$$

where  $E_{\text{des}}$  and  $E_{\text{MgC}_2 + (n-1)\text{H}_2}$  are gradual desorption energy and net energy of  $\text{MgC}_2$  with  $(n-1)$   $\text{H}_2$  adsorption, respectively.

### 3. Results and discussion

#### 3.1 Structural and electronic properties of $\text{MgC}_2$ monolayer

The relaxed crystal structure of the  $\text{MgC}_2$  monolayer within the  $P6mm$  space group is shown in Fig. 1(a). We have previously computationally designed  $\text{MgB}_2$  and analyzed its structural, electronic, and adsorption properties. This material was investigated for hydrogen storage applications after being decorated with alkali metals.<sup>25</sup> Similarly, the  $\text{MgC}_2$  structure was designed computationally by replacing the B atoms in  $\text{MgB}_2$  with C atoms.  $\text{MgB}_2$  consists of two atomic layers and exhibits a geometry similar to that of  $\text{BeB}_2$ ,  $\text{HfB}_2$ ,  $\text{TiB}_2$ ,  $\text{FeB}_2$ , and  $\text{MgB}_2$  monolayers.<sup>36–39</sup>

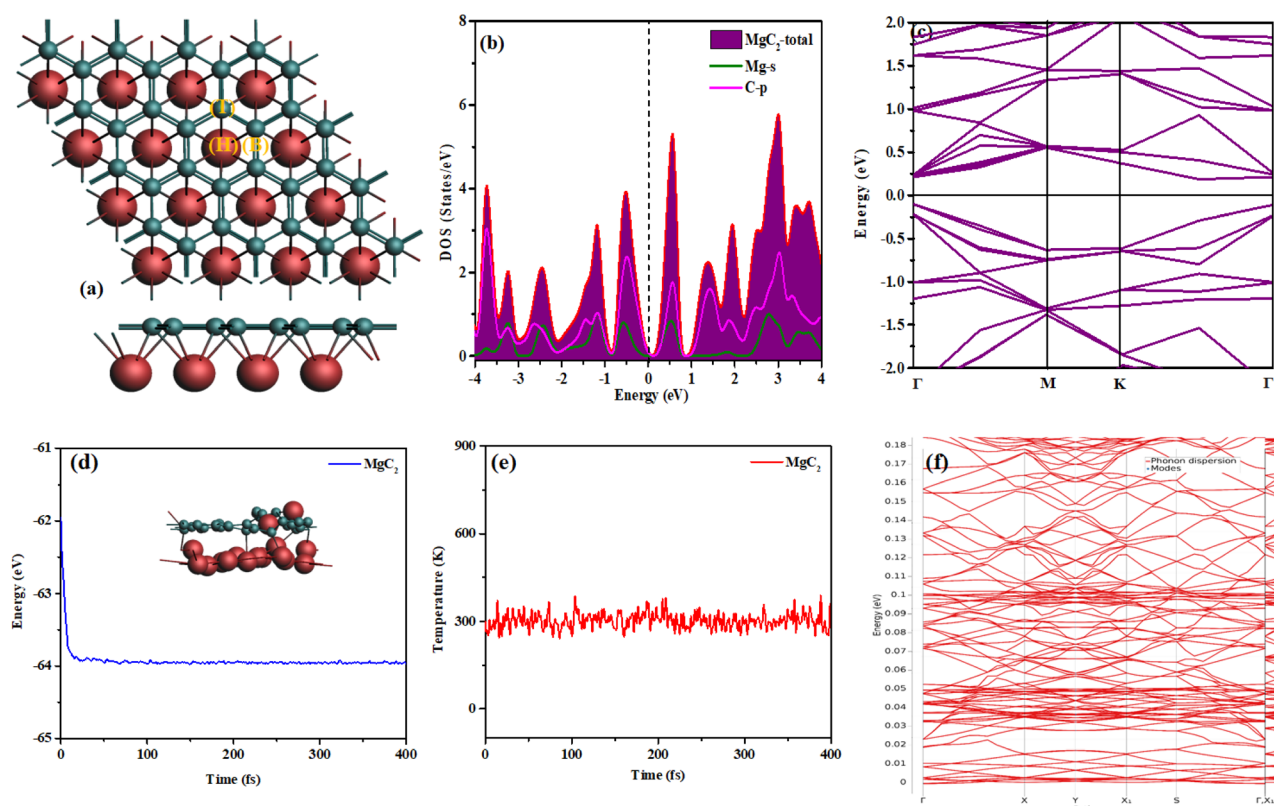
The upper atomic layer, composed of C atoms, follows a honeycomb arrangement similar to graphene, and the lower

**Table 1** Lattice parameters, bond length, and height of  $\text{MgC}_2$  monolayer

	Lattice parameters (Å)	Bond length(Å)	Height (Å)
<b><math>\text{MgC}_2</math>(this work)</b>	<b>3.181</b>	<b>2.14</b>	<b>1.46</b>
$\text{MgB}_2$	3.017	2.42	1.68
$\text{BeB}_2$	3.032	1.84	0.44
$\text{HfB}_2$	3.165	2.32	1.41
$\text{TiB}_2$	3.083	1.73	1.19

atomic layer consists of Mg atoms positioned at the center of the C–C honeycomb. The optimized lattice parameters of  $\text{MgC}_2$  are  $a = b = 3.181 \text{ Å}$ , which are slightly larger than those of the  $\text{BeB}_2$  monolayer ( $a = b = 3.032 \text{ Å}$ ) and  $\text{MgB}_2$  ( $a = b = 3.017 \text{ Å}$ ). Additionally, the monolayer thickness of  $\text{MgC}_2$  is  $1.460 \text{ Å}$ , significantly greater than that of  $\text{BeB}_2$  ( $0.443 \text{ Å}$ ) and less than  $\text{MgB}_2$  ( $1.683 \text{ Å}$ ). A comparison of the lattice parameters is provided in Table 1.

Fig. 1(b and c) illustrates the band structure and density of states (DOS) of  $\text{MgC}_2$ . The DOS results indicate that the material exhibits semiconducting behavior with a bandgap of  $0.30 \text{ eV}$ . The band structure indicates the direct bandgap of  $0.3 \text{ eV}$  at  $\Gamma$  point which is consistent with the DOS. The total DOS is mainly governed by the p-orbitals of carbon atoms, with only a minor contribution from the s-orbitals of magnesium atoms. Fig. 1(b)



**Fig. 1** Optimized structure of (a)  $\text{MgC}_2$ , (b) density of states (DOS), and (c) band structure (d) energy–time (e) temperature–time (f) phonon of  $\text{MgC}_2$ . Green-colored atoms represent boron, while maroon-colored atoms represent magnesium.



Table 2 Binding energy of MgC<sub>2</sub> at studied sites

Adsorption sites	TOP (T)	Bridge (B)	Hollow (H)
Binding energy (eV)	0.23	0.25	0.25

reveals that the s-orbital of Mg and the p-orbital of C overlap in the valence band.

MD calculations were conducted at 300 K for the pure MgC<sub>2</sub> monolayer. As observed in Fig. 1d and e, the energy graph peaks at −62 eV, while the temperature graph reaches a maximum value of 460 K. The continuous fluctuations in both graphs indicate the structural stability of MgC<sub>2</sub> at 300 K. The phonon calculations for the pure MgC<sub>2</sub> monolayer reveal no imaginary frequencies throughout the Brillouin zone, indicating its dynamic stability as shown in Fig. 1f. The absence of imaginary peaks in the phonon dispersion spectrum confirms that the structure remains stable under small perturbations and does not undergo spontaneous distortions or soft modes. This result validates the robustness of the material and supports its potential for practical applications.

### 3.2 Adsorption of hydrogen molecule on pristine MgC<sub>2</sub>

This section examines the interaction of H<sub>2</sub> molecules with MgC<sub>2</sub>, where a single H<sub>2</sub> molecule was placed at various distinct locations, as shown in Fig. 1(a). The hydrogen adsorption sites include (i) the hollow site of the hexagonal boron ring (H), (ii) the top of a boron atom (T), and (iii) the bridge site between two boron atoms (B). The structures were sequentially optimized by varying the adsorption sites. The binding energy was calculated using the formula in eqn (2), and the corresponding values for different sites are presented in Table 2.<sup>35</sup>

The binding energy estimates for these adsorption sites are presented in Table 2, and all values fall within the DOE criteria range (0.2–0.8 eV per H<sub>2</sub>), making them suitable for hydrogen adsorption. The binding energy values for the H and B sites were found to be similar or higher. However, hydrogen initially placed at the B site migrated to the H site, indicating that the H site is more stable than the B site. Consequently, additional hydrogen molecules were sequentially adsorbed onto the H site, and their binding energies were calculated, as shown in Table 3.<sup>25,40,41</sup> The 5.73 Å is the ideal upward distance for a deposited

H<sub>2</sub> molecule. In reference, the mean bond length of a free H<sub>2</sub> molecule is approximately 0.74 Å. An H<sub>2</sub> molecule adsorbed on the structure and caused a little increase in bond length to 0.75 Å.<sup>42</sup> The retention of H<sub>2</sub> molecular properties is indicated by the change in bond length that is seen throughout the procedure of adsorption and optimization. To determine the optimal adsorption site, adsorption energies were analyzed. Following structural optimization, the hexagonal hollow site emerged as the most favorable location for H<sub>2</sub> adsorption, with an average binding energy of 0.24 eV.<sup>10,43</sup>

### 3.3 Adsorption and electronic properties of H<sub>2</sub> molecule adsorbed on MgC<sub>2</sub>

Based on the binding energy, it is concluded that the most favorable atom decoration occurs at the H site. Fig. 2 illustrates the optimized structures of hydrogenated MgC<sub>2</sub>. Hydrogen molecules were sequentially adsorbed onto the H site, up to a total of 8H<sub>2</sub>, and their binding energy, adsorption energy, weight percentage, desorption energy, and desorption temperature were calculated, as detailed in Table 3.

Subsequently, each structure was optimized, and a vertical distance of 3.02 Å was determined to be optimal for the adsorbed H<sub>2</sub> molecule. It is discovered that the free H<sub>2</sub> molecule's bond length is 0.74 Å. The H–H bond length elongates to 0.78 Å when one H<sub>2</sub> molecule is added to the system. The H–H bond length in the 2H<sub>2</sub>–8H<sub>2</sub> molecule adsorption system is determined to be 0.78–0.80 Å. The H–H bond length determined throughout the adsorption and optimization procedures demonstrates that the molecule's structure remains intact.<sup>42</sup>

As shown in Fig. 3, the total and partial DOS were calculated for the Mg-s, C-p, H-s, and s orbitals of all metal elements to evaluate the system's electrochemical properties. Fig. 3(a) reveals that absorption of a single molecule of hydrogen transforms the behavior of MgC<sub>2</sub> from semiconducting to metallic, resulting in a zero bandgap. The p orbital of the C atom significantly contributes to the total DOS, while the s orbital of hydrogen exhibits overlapping peaks at −4.5 eV in the valence band, and hybridized peaks between 6–8 eV in the conduction band. Similarly, Fig. 3(b) demonstrates that after the adsorption of four hydrogen molecules, MgC<sub>2</sub> retains its metallic nature with a zero bandgap. The p orbital of the C atom remains the dominant contributor to total DOS, while the s orbital of hydrogen displays hybridized peaks within the ranges

Table 3 Binding, adsorption energy, weight percentage, and desorption temperature values

	Binding energy (eV)	Adsorption energy (eV per H <sub>2</sub> )	Weight percent (%)	Desorption energy (eV)	Desorption temperature (K)
MgC <sub>2</sub> –1H <sub>2</sub>	0.25	0.290	0.26	0.29	370.52
MgC <sub>2</sub> –2H <sub>2</sub>	0.40	0.242	0.52	0.20	309.19
MgC <sub>2</sub> –3H <sub>2</sub>	0.62	0.234	0.77	0.21	298.97
MgC <sub>2</sub> –4H <sub>2</sub>	0.82	0.226	1.03	0.20	288.75
MgC <sub>2</sub> –5H <sub>2</sub>	1.03	0.223	1.29	0.21	284.92
MgC <sub>2</sub> –6H <sub>2</sub>	1.24	0.220	1.54	0.20	281.08
MgC <sub>2</sub> –7H <sub>2</sub>	1.44	0.218	1.80	0.20	278.53
MgC <sub>2</sub> –8H <sub>2</sub>	1.62	0.640	2.05	0.17	817.71





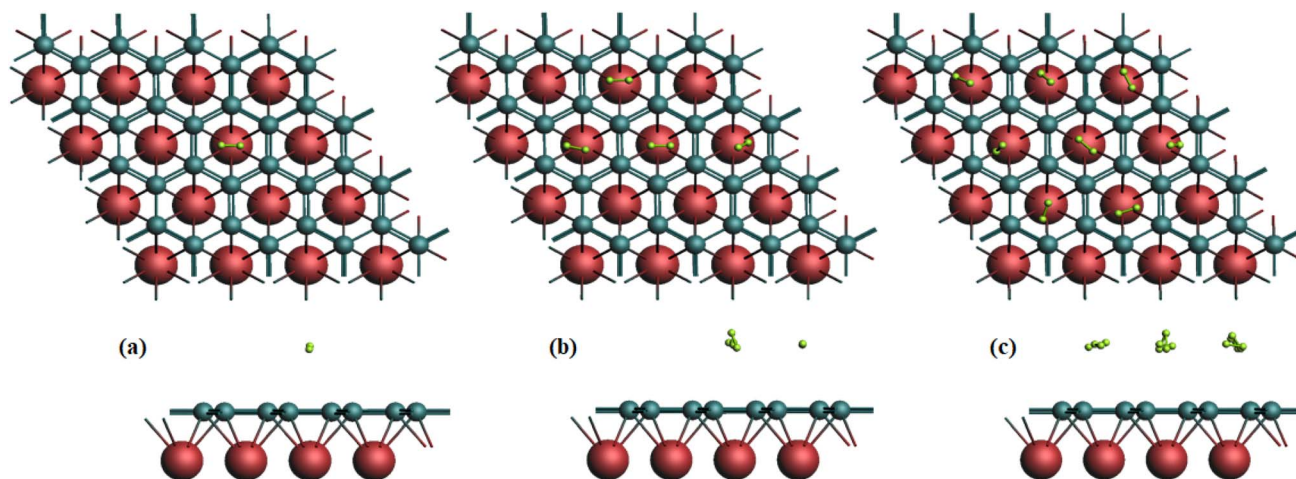


Fig. 2 Optimized structures of (a) 1H<sub>2</sub>- (b) 4H<sub>2</sub>- (c) 8H<sub>2</sub>-adsorbed MgC<sub>2</sub>.

of  $-4$  to  $-6$  eV in the valence band and  $6$  to  $8$  eV in the conduction band.

In Fig. 3(c), the total DOS of MgC<sub>2</sub> after the adsorption of eight hydrogen molecules shows a reduction in the bandgap. The p orbital of the C atom remains the primary contributor to the total DOS, while the s orbital of hydrogen exhibits overlapping peaks at  $7.1$  eV in the conduction band and hybridized peaks within the ranges of  $-4$  to  $-6$  eV in the valence band and  $6$  to  $8$  eV in the conduction band. In summary, the adsorption between MgC<sub>2</sub> and hydrogen is primarily driven by the polarizing effect of the localized voltage generated by charged particles and their spatial recombination.

After the adsorption of hydrogen molecules onto the MgC<sub>2</sub> monolayer, charge analysis calculations were conducted. The results reveal that 1, 4, and 8 hydrogen molecules withdraw charges of  $0.041$  e,  $0.139$  e, and  $0.259$  e from the surface, respectively. For MgC<sub>2</sub>, the typical hydrogen desorption temperature ranged from  $370$  K to  $496$  K for 1H<sub>2</sub>. As illustrated in Fig. 4, these values were derived from adsorption energies calculated within a pressure range of  $1$ – $10$  atm.

Furthermore, three critical temperatures related to hydrogen desorption were determined: the maximum temperature at

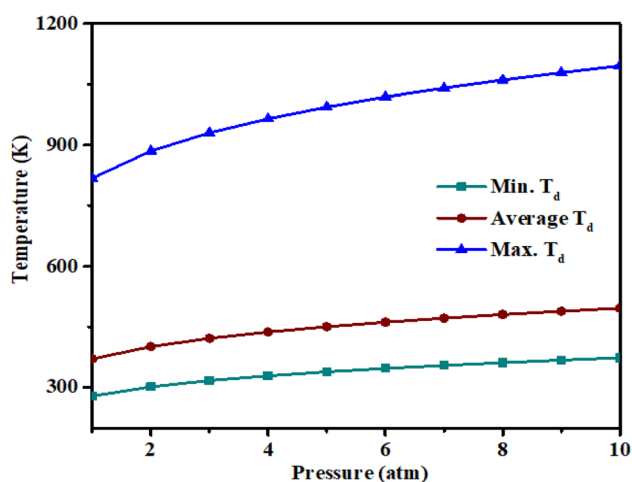


Fig. 4 Minimum, average, and maximum values for desorption temperature after hydrogen adsorption on MgC<sub>2</sub>.

which desorption takes place (max-TD), the minimum temperature necessary to start hydrogen release (min-TD), and the average desorption temperature (Av-TD). In our study, the

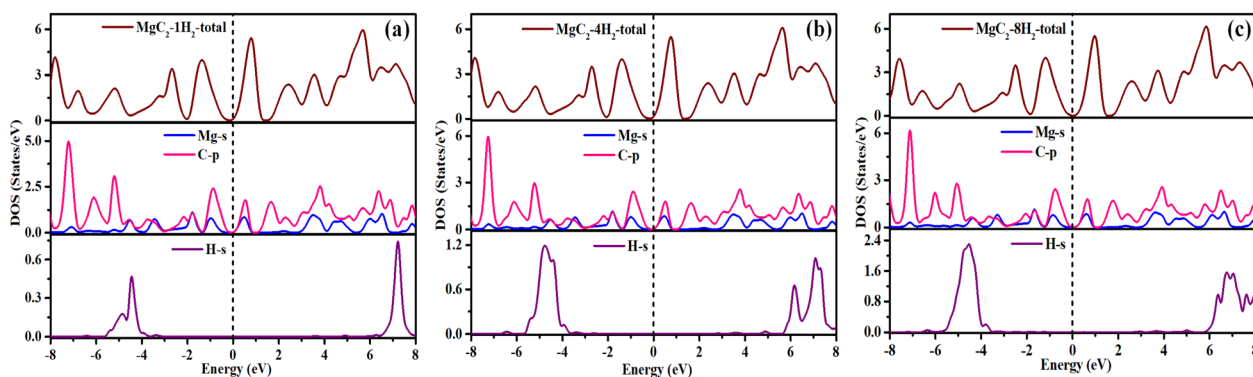


Fig. 3 DOS of (a) 1H<sub>2</sub>, (b) 4H<sub>2</sub>, (c) 8H<sub>2</sub> adsorbed MgC<sub>2</sub>.

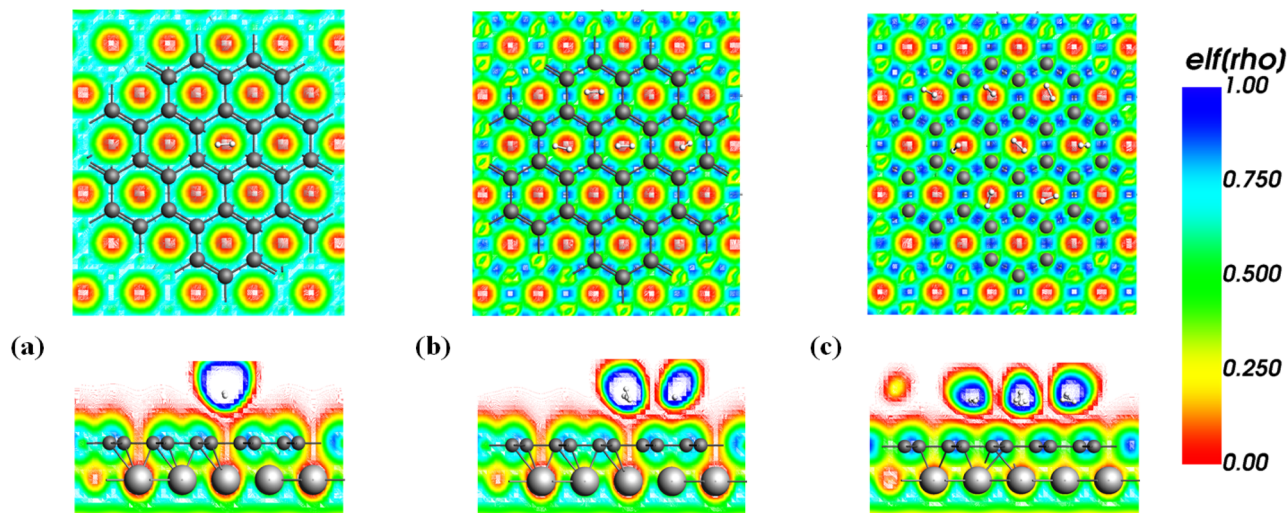


Fig. 5 (a)  $\text{MgC}_2\text{-1H}_2$ , (b)  $\text{MgC}_2\text{-4H}_2$ , (c)  $\text{MgC}_2\text{-8H}_2$ .

average desorption temperature of  $\text{MgC}_2$  monolayers was calculated to range from 370 K to 496 K, indicating their suitability for fuel cell applications.<sup>44</sup>

### 3.4 Electron localization function (ELF)

The charge carrier density mechanism serves as a gauge of charge confinement and scattering and is connected to the kinetic energy density.<sup>45,46</sup> Its values range from 0.5 to 1.0, indicating areas where electrons are present and engaged in bonding or nonbonding interactions, including lone pairs, covalent bonds, or atomic shells. On the other hand, values between 0.0 and 0.5 show that there are small intermolecular interactions, and electrons are dispersed throughout the atomic shells.<sup>47</sup> The ELF map highlights regions of increased electron concentration and shows how charge carrier density changes on surfaces and adsorbents during complex creation.

As seen in Fig. 5(a) a value near 0.5 is indicated by green, indicating the presence of covalent bonds, and a large charge density is indicated by the increased red coloration in the  $\text{MgC}_2$  layer after deposition of  $1\text{H}_2$ .<sup>48</sup> Following the uptake of four hydrogen molecules, Fig. 5(b) shows a striking color shift with an increase in green and a decrease in red and yellow, suggesting covalent electron sharing between the hydrogen

molecules and the  $\text{MgC}_2$  layer. The bright red hue represents a decrease in surface electrons, while the increasing green color and the presence of a light blue shade in Fig. 5(c) indicate charge transfer from the  $\text{MgC}_2$  layer to the hydrogen molecules, with up to eight hydrogen molecules adsorbed.

### 3.5 *Ab initio* molecular dynamics (AIMD)

AIMD calculations were performed to assess the configurational stability of the system after hydrogen adsorption on  $\text{MgC}_2$  at elevated temperatures such as 400 K. Simulations were conducted with a time step of 0.25 femtoseconds for a total duration of approximately 4.5 picoseconds at a temperature of 400 K. Fig. 6 presents the temperature and energy fluctuation graphs at 400 K.

As can be seen from Fig. 6(a1), the system attained higher energy values of  $-1706.5$  eV when one hydrogen molecule was adsorbed on  $\text{MgC}_2$ . Fig. 6(a2) shows that the temperature peaks peaked at 5000 K. Based on Fig. 6(b1), it can be concluded that the system achieved higher energy values of  $-1855.9$  eV when the eight hydrogen molecules were adsorbed on  $\text{MgC}_2$ . Fig. 6(b2) shows the temperature peaks around 4500 K. The steady fluctuations of the peaks indicate the stability of the  $\text{MgC}_2$  sheet after hydrogen adsorption at 300 K.

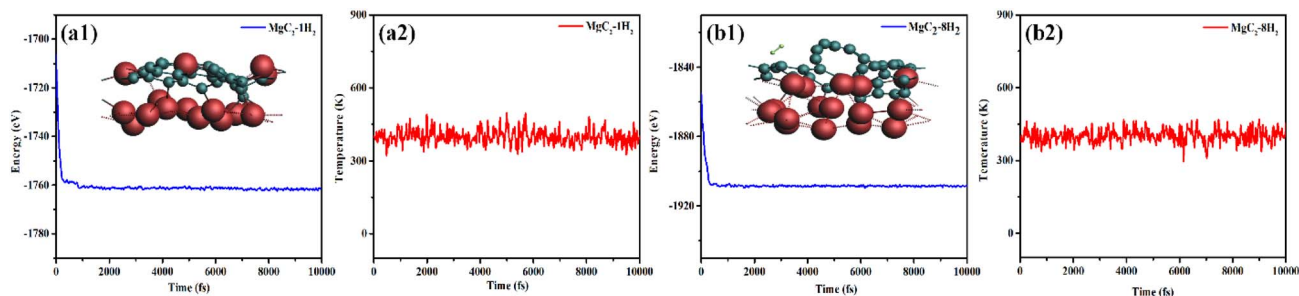


Fig. 6 Energy and temperature graph of (a1 and a2)  $\text{MgC}_2\text{-1H}_2$  and (b1 and b2)  $\text{MgC}_2\text{-8H}_2$ .



## 4. Conclusion

In this study, the adsorption behavior of hydrogen on pure  $\text{MgC}_2$  was systematically analyzed using density functional theory (DFT) and *Ab Initio* Molecular Dynamics (AIMD) simulations. The structural stability of  $\text{MgC}_2$  was thoroughly evaluated using phonon dispersion calculations and AIMD simulations, confirming its robustness as a potential hydrogen storage material.  $\text{MgC}_2$  demonstrates a gravimetric hydrogen capacity of 2.05 wt%, capable of adsorbing up to eight hydrogen molecules. The hydrogen changes the behavior of the pure  $\text{MgC}_2$  to metallic. The average adsorption energy for  $\text{MgC}_2$  was calculated to be 0.286 eV. Hirshfeld charge analysis revealed charge transfers of 0.041, 0.139, and 0.259 e to 1, 4, and 8 hydrogen molecules to the pure  $\text{MgB}_2$  structures respectively. AIMD simulations confirmed that the pure  $\text{MgC}_2$  structure remains stable without any distortions upon hydrogen adsorption at 400 K. The comprehensive analysis of structural, electronic, and adsorption properties, along with binding energy calculations, desorption temperature and energy assessments, charge density analysis, and electron localization function (ELF) mapping, all indicate that  $\text{MgC}_2$  is a highly promising candidate for hydrogen storage applications. The strong interaction between  $\text{MgC}_2$  and hydrogen molecules, coupled with favorable adsorption and desorption characteristics, suggests its potential for practical hydrogen storage solutions. These findings provide a theoretical foundation for further experimental validation and the development of  $\text{MgC}_2$ -based hydrogen storage materials.

## Data availability

Data is available on request from the authors.

## Conflicts of interest

There are no conflicts to declare.

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