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## Impact of boron desymmetrization on supramolecular polymerization of BODIPY dyes†

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Supramolecular polymers are often investigated for highly symmetric and planar molecules, such as typically explored BF<sub>2</sub>-substituted BODIPY dyes. However, it is surprising that the possibility of desymmetrizing the sp<sup>3</sup> hybridized boron centre of BODIPY dyes has remained unexplored in the context of supramolecular polymerization. Herein, we synthesized a new BODIPY derivative **2** with two different substituents at the boron (fluorine and phenyl), resulting in a system with two different  $\pi$ -surfaces, and analyzed its supramolecular polymerization in non-polar media. Notably, this symmetry reduction increases the complexity of the self-assembly by enabling the formation of an intermediate assembled state, which cannot be found in the symmetrical model BODIPY **1** with a BF<sub>2</sub> group. Different experimental and theoretical studies suggest that significant steric effects together with multiple potential intermolecular stacking modes of the BODIPY dyes lead to discrete nanoparticle intermediates that ultimately transform into more-ordered H-type supramolecular polymers at lower temperatures. Our results introduce a new design strategy for controlled supramolecular polymerization.

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## Introduction

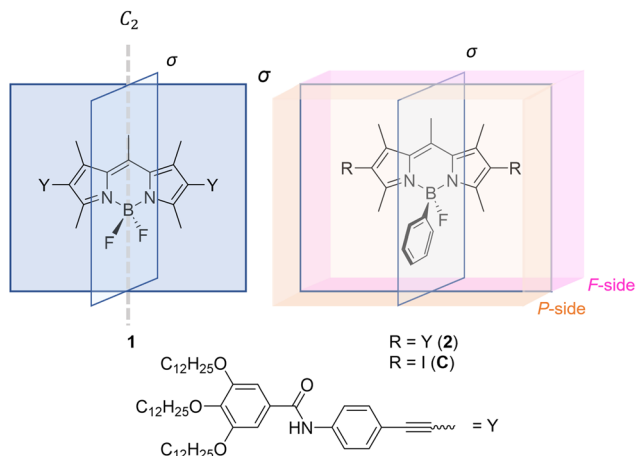
The controlled self-assembly of  $\pi$ -conjugated molecules has received considerable interest in recent years.<sup>1–4</sup> The rational design of chromophores has enabled the development of supramolecular functional materials with interesting properties like controlled morphologies,<sup>5–9</sup> favourable luminescence<sup>10,11</sup> or biomedical applications.<sup>12–14</sup> The main driving force towards controlled supramolecular polymers are non-covalent interactions, including  $\pi$ - $\pi$  interactions,<sup>15–17</sup> hydrogen bonds<sup>18,19</sup> or ionic interactions.<sup>20–22</sup> Due to the planar character of the  $\pi$ -bond, the design of new molecules that are suitable for self-assembly is usually limited to 1D or 2D functionalization of the chromophores.<sup>23–27</sup> However, the 2D character of chromophores can be overcome and expanded into the third dimension by rational molecular design, *e.g.* in helicenes.<sup>28,29</sup> Due to the repulsive steric interactions of their terminal rings, these molecules build screw-like structures that

go along with unique self-assembly features.<sup>30–32</sup> Furthermore, distorted  $\pi$ -surfaces are also observed for molecules such as bay-substituted perylene bisimide dyes, leading to planar chirality that give access to the emission of circularly polarized light.<sup>17,33</sup>

Recently, we reported that also the slight out-of-plane distortion of the  $\pi$ -surface of a bis(difluoroboron)-1,2-bis((1*H*-pyrrol-2-yl)methylene)hydrazine (BOPHY) derivative leads to increased complexity in the supramolecular polymerization when compared to a structurally related but planar boron dipyrromethene (BODIPY) derivative.<sup>34,35</sup> Although the oppositely oriented BF<sub>2</sub> groups of the BOPHY core lead to a slight planarity breaking, both sides of the  $\pi$ -surface are identical in terms of steric features. This is due to the symmetry resulting from the sp<sup>2</sup> hybridized carbon atoms that are included in the chromophores. To go beyond these systems, we aimed at designing self-assembling building blocks where the  $\pi$ -surface can be unsymmetrically functionalized orthogonal to the  $\pi$ -plane to yield a Janus-like  $\pi$ -surface. To address this goal, we exploited the planar BODIPY chromophore, as the sp<sup>3</sup> hybridized boron atom allows the selective substitution of one of the fluorine atoms, thus enabling desymmetrization in the third dimension.

Starting from our symmetrical model BODIPY **1**, we reduced the symmetry by removing the mirror plane aligned with the  $\pi$ -surface of the chromophore. This was achieved by replacement of a fluorine atom with a phenyl substituent at the boron atom to yield BODIPY derivative **2** (Scheme 1, for

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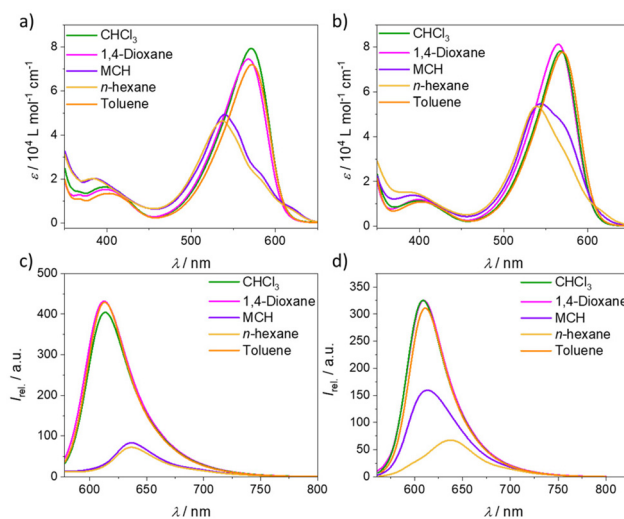
**Scheme 1** Molecular structures of the model BODIPY derivative **1** with  $C_2$ -symmetry, its unsymmetrical analogue **2** and reference compound **C**.

synthetic details see Fig. S1 in the ESI†). This molecular design enables a sterically more hindered surface on the phenyl-substituted side (P-surface) and a sterically less hindered side (F-surface, Scheme 1), which is expected to strongly influence the molecular stacking and self-assembly behaviour.

## Results and discussion

### Spectroscopic characterization and self-assembly of BODIPY derivative **2**

BODIPY dyes are well-known to self-assemble in solution into ordered nanostructures through various stacking modes.<sup>36–38</sup> The optical properties and self-assembly behaviour of BODIPY derivative **1** were already well discussed and investigated in previous reports.<sup>34,39</sup> **1** was found to form a single type of aggregate (one-dimensional H-type) with an antiparallel and slightly rotationally displaced packing of the chromophores that follows the cooperative mechanism. Diverse kinetic studies showed that the system is very robust and unaffected by various sample treatments (*e.g.* mechanical agitation, solvophobic quenching, *etc.*). Due to the high reproducibility and in-depth comprehension of this system, we selected it as a reference compound to understand the self-assembly behaviour of the structurally related BODIPY derivative **2**. To gain first insights into the properties of **2**, we performed solvent-dependent UV-vis and fluorescence measurements (Fig. 1). In solvents favouring the monomer state like  $CHCl_3$  and 1,4-dioxane, the absorption maximum, referring to the  $S_0 \rightarrow S_1$  transition, is located at  $\lambda = 568$  and  $564$  nm, respectively (Fig. 1,  $\epsilon = 8 \times 10^4$  L mol<sup>-1</sup> cm<sup>-1</sup>,  $c = 10$   $\mu$ M, 298 K). A band that corresponds to the  $S_0 \rightarrow S_2$  transition is also visible at lower wavelength ( $\lambda \approx 400$  nm). The emission maximum of **2** in these solvents is found at 609 nm, resulting in a Stokes shift of 41 nm in  $CHCl_3$  and 45 nm in 1,4-dioxane, which is similar to the spectroscopic properties of **1** in the same solvents.

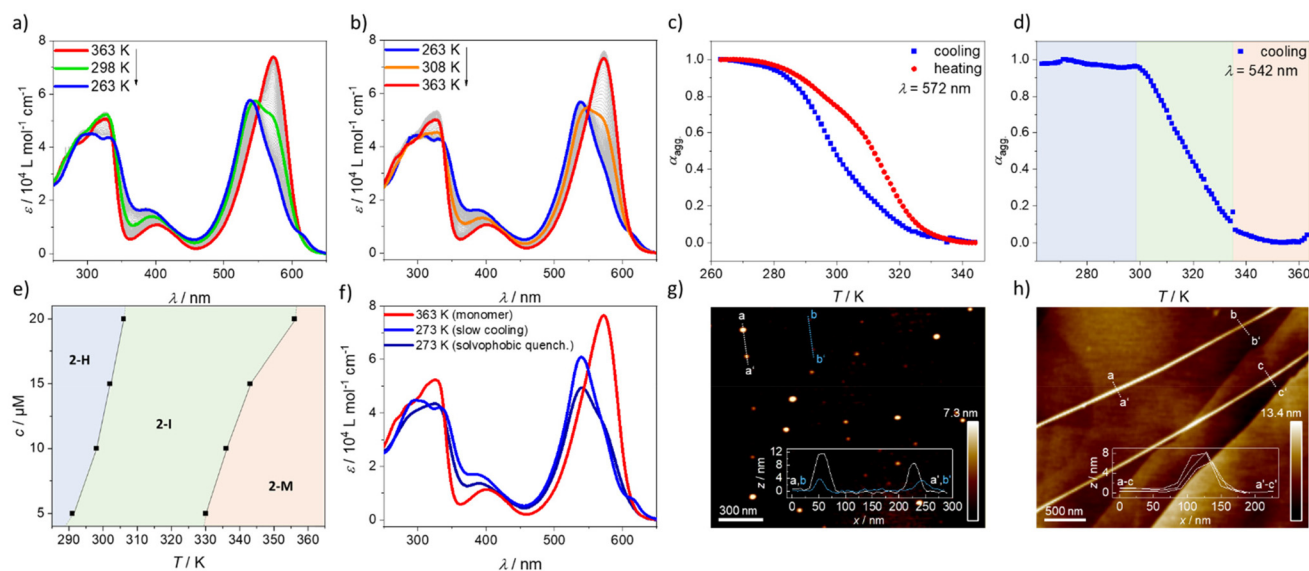


**Fig. 1** Solvent-dependent UV-vis and emission studies ( $\lambda_{exc.} = \lambda_{max. abs.}$ ) of **1** (a and c) and **2** (b and d) at  $c = 10$   $\mu$ M at room temperature.

However, in non-polar media under the same conditions ( $c = 10$   $\mu$ M, 298 K), the absorption maximum of **2** shifts to 540 nm in *n*-hexane and 545 nm in methylcyclohexane (MCH), accompanied by a reduction in the absorbance and the concomitant emergence of a red-shifted shoulder at 616 nm. These spectral features resemble those of **1** and agree with the formation of H-type aggregates with a rotationally displaced packing.<sup>34</sup> This is further supported by the concomitant quenching and red-shift of the emission maximum in both non-polar solvents (MCH:  $\lambda_{max.} = 613$  nm,  $\lambda_{stokes}$  shift = 68 nm, *n*-hexane:  $\lambda_{max.} = 638$  nm,  $\lambda_{stokes}$  shift = 98 nm) compared to the solvents favouring the monomer state (Fig. 1d). Interestingly, the spectral pattern of the H-type aggregates of **2** is similar to that of model compound **1** (Fig. 1a and c), indicating comparable intermolecular interactions and packing. However, the degree of aggregation for **2** in MCH seems to be less pronounced compared to **1** under similar conditions, as the spectrum of **2** shows a pronounced shoulder at 568 nm that is most probably related with the presence of residual monomers (Fig. 1b). A plausible reason for this difference could be the increased steric hindrance of **2** that is introduced by the phenyl group at the boron atom.

To gain further insights into the self-assembly process, variable temperature (VT)-UV-vis measurements were performed. Upon cooling a solution of **2** ( $c = 10$   $\mu$ M, 1 K min<sup>-1</sup>, MCH) from 363 K to 263 K, the absorption maximum shifts from 572 nm, corresponding to the monomeric species (**2-M**, red spectrum, Fig. 2a), to 538 nm, which can be associated with the H-type aggregate (**2-H**, blue spectrum, Fig. 2a). Furthermore, the red-shifted band at 616 nm that was previously observed in the solvent-dependent studies becomes apparent at lower temperatures (Fig. 2a). The H-type aggregate formation is independent of the cooling rate, indicating that the process is under thermodynamic control<sup>40–42</sup> (Fig. S8†). Compared to **1** at the same concentration, the onset of the





**Fig. 2** UV-vis cooling (a) and heating (b) experiments of **2** (MCH, 10  $\mu\text{M}$ , 1  $\text{K min}^{-1}$ ) with the corresponding cooling and heating curves, plotted at  $\lambda = 572$  nm (c). (d) Cooling curve of **2** plotted at  $\lambda = 542$  nm to demonstrate the two-step process (MCH, 10  $\mu\text{M}$ , 1  $\text{K min}^{-1}$ ). (e) Phase diagram extracted from cooling experiments of **2** at different concentrations. (f) UV-Vis experiments depicting the solvophobic quenching of **2** (10  $\mu\text{M}$ , MCH/ $\text{CHCl}_3$ ). (g) AFM image of **2-I** after spin-coating a 20  $\mu\text{M}$  solution of **2** (306 K, MCH) on HOPG and (h) AFM image of **2-H** after spin-coating a 20  $\mu\text{M}$  solution of **2** (273 K, MCH) on HOPG.

aggregation upon cooling starts at slightly lower temperatures. This effect results most probably from the increased steric hindrance that is introduced on the boron atom when a fluorine atom in **1** is replaced with a phenyl substituent in **2** (Fig. S9<sup>†</sup>). Upon heating aggregate **2-H** (1  $\text{K min}^{-1}$ , MCH, blue spectrum, Fig. 2b), the monomer is recovered at 363 K. However, by comparing the degree of aggregation ( $\alpha_{\text{agg}}$ ) of the cooling and heating processes at different temperatures (Fig. 2c), differences in the curve shape and a slight hysteresis are observed. A more detailed analysis of the cooling experiment shows that the absorption maximum of the monomer species of **2** shifts to shorter wavelengths at intermediate temperatures (298 K, green spectrum, Fig. 2a) without the appearance of a red-shifted shoulder at 616 nm. Furthermore, the bands at shorter wavelength (326 nm) corresponding to the absorption of the substituents at the 2- and 6-positions conjugated to the BODIPY core (Y in Scheme 1) remain unchanged. Only below a critical temperature (298 K, green spectrum, Fig. 2a) the red-shifted band appears rapidly (Fig. S10<sup>†</sup>) and the bands at short wavelength start to shift and split. These findings suggest the occurrence of an intermediate state (**2-I**,  $\lambda_{\text{max}} = 546$  nm) that transforms into the final H-type aggregate (**2-H**) after surpassing a concentration-dependent critical temperature (Fig. S10<sup>†</sup>). The shift from the monomeric species towards **2-I** occurs without a clear isobestic point, while the transformation from **2-I** to **2-H** exhibits an isobestic point at 542 nm. The two-step transformation from **2-M** to **2-H** via **2-I** can be visualized rather well when the degree of aggregation is plotted e.g. at the isobestic point between the intermediate and the aggregate (542 nm, Fig. 2d). At high temperatures, when mainly monomers are present, the degree of aggregation

remains nearly zero between 363 K and ca. 336 K. Below this temperature (336 K), the formation of **2-I** sets in, as indicated by the increase of  $\alpha_{\text{agg}}$  until a plateau is reached at around 298 K. Below this temperature, the proportion of **2-H** increases but no change in  $\alpha_{\text{agg}}$  is observed due to the characteristics of the isobestic point. This analysis was carried out for different concentrations to obtain a phase diagram (Fig. 2e and Fig. S11<sup>†</sup>) that reveals the general trends of the aggregation process. Increasing the concentration results in an increased onset temperature for the formation of **2-I** as well as for **2-H**.

In order to understand the disassembly of the aggregate, corresponding heating experiments were carried out (Fig. 2b). The studies reveal that the disassembly of **2-H** to the monomeric species occurs more gradually than in the cooling experiments. Nevertheless, an intermediate species **2-I\*** with an absorption maximum at  $\lambda_{\text{max}} = 546$  nm (similar to **2-I**) can be observed. However, in stark contrast to **2-I**, the red-shifted shoulder at 616 nm as well as the bands at around 300 nm resemble the spectral features of the aggregate, rather than those of the monomer. We assume that the cooperative effect in the transformation from **2-I** to **2-H** in the cooling process (Fig. S10a<sup>†</sup>) is caused by additional interactions of the side groups, most probably hydrogen bonds. This stabilization process does not allow to revert the intermediate **2-I** in the heating process but enables the formation of another intermediate species with interactions between the side groups. This results in a slightly more disturbed disassembly process, requiring higher temperatures.

Interestingly, the fact that reference compound **1** does not form any intermediate species (Fig. S12<sup>†</sup>) demonstrates that the replacement of the fluorine atom in **1** with a phenyl group



in **2** has a direct consequence on the observed self-assembly behaviour. We propose that repulsive steric interactions involving the phenyl group at the boron atom initially hinder close contacts of the BODIPY chromophores, resulting in species **2-I** with weak exciton coupling. Below a certain temperature, cooperative effects come into play and overcome the repulsive steric interactions, leading to the formation of the final aggregate **2-H**. To investigate the stability of **2-I**, time-dependent studies were performed. Unexpectedly, species **2-I** did not transform into aggregate **2-H** over several hours, demonstrating its thermodynamic stability (Fig. S13<sup>†</sup>). Therefore, we assume that **2-I** is a kinetically trapped off-pathway aggregate that forms *via* an anti-cooperative mechanism, analogous to previous literature reports.<sup>43</sup> Only below a certain critical temperature, the formation of **2-H** occurs *via* a cooperative pathway, indicated by the rapidly rising band at 616 nm (Fig. S10a<sup>†</sup>). This transformation could be fitted well to the nucleation–elongation model to obtain the thermodynamic parameters (Fig. S10a and Table S1<sup>†</sup>).<sup>44</sup> However, these data need to be taken with caution, as the process involves additional intermediate species that are not considered in the model. The anti-cooperative formation of **2-I** would explain the high stability of this species. Due to the favoured nucleation,<sup>2</sup> the monomer concentration is reduced, hindering the transformation to the thermodynamic aggregate *via* the monomer. Unfortunately, the curves describing the transformation from the monomer to **2-I** cannot be fitted to any suitable thermodynamic model, as this process cannot be isolated from the transformation to **2-H** at any wavelength (Fig. S10b<sup>†</sup>).

To prove the assumption of two competitive pathways, we investigated the morphology of **2** in the intermediate state and compared it to that of the final aggregate state **2-H**. To this end, we prepared solutions of **2-I** (MCH, 20  $\mu$ M, 310 K) and **2-H** (MCH, 20  $\mu$ M, 273 K), spin-coated them on highly ordered pyrolytic graphite (HOPG) and analysed the samples through atomic force microscopy (AFM). **2-I** formed small particles with a uniform size of *ca.* 4 to 12 nm (Fig. 2g and Fig. S20<sup>†</sup>). According to the size, this suggests the formation of discrete assemblies comprising a limited number of molecules. The aggregate **2-H**, however, shows elongated fibres (Fig. 2h and Fig. S21<sup>†</sup>) with a uniform width of 8 nm and several micrometres in length. These results corroborate the assumed pathway complexity, as the small particles of **2-I** suggest an anti-cooperative mechanism and the elongated structures of **2-H** are typical for aggregates that formed *via* a cooperative mechanism. Furthermore, the formation of well-structured, fibre-like H-type aggregates points out the drastic consequences of small structural changes in a molecular design, as they can lead to distinct pathways with different morphology. Despite **1** and **2** having similar spectral patterns for their respective aggregates, **1** was reported to form short worm-like structures,<sup>39</sup> while **2** self-assembles into much more extended fibres.

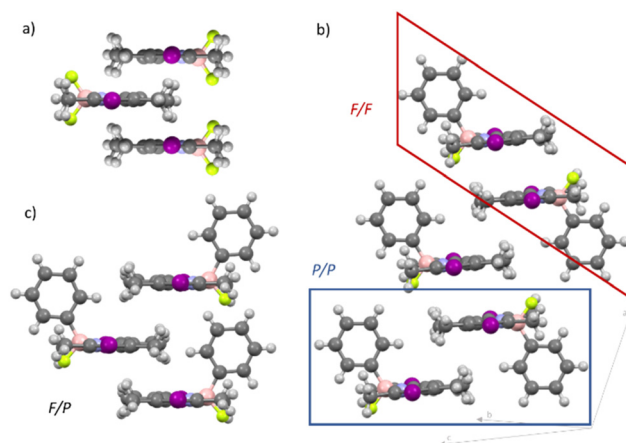
The increased steric hindrance of **2** compared to **1** prompted us to perform quenching experiments to rationalize whether other kinetic species may potentially be involved in the self-assembly process. Therefore, thermal (Fig. S14<sup>†</sup>) and

solvophobic quenching experiments (Fig. 2f) were performed. For the solvophobic quenching, a monomeric solution of **2** (CHCl<sub>3</sub>, *c* = 200  $\mu$ M, 273 K, 100  $\mu$ L) was prepared and rapidly diluted with cold MCH (273 K, 3900  $\mu$ L) to gain a 10  $\mu$ M solution that was subsequently tracked by UV-vis spectroscopy. These experiments demonstrate that the spectral features of the quenched solution resemble those of **2-H** that were observed in the cooling experiments ( $\lambda_{\text{max.}}$  = 540 nm). However, the formation of the typical red-shifted shoulder of **2-H** around 616 nm cannot be observed, indicating a similar but less ordered self-assembly compared to **2-H**. Nevertheless, this solution transforms over time into **2-H**, as evident by the emergence of the shoulder at 616 nm (Fig. S14a<sup>†</sup>). Similar observations were made for the thermally quenched solution (MCH, 10  $\mu$ M, 273 K, Fig. S14b and c<sup>†</sup>). In contrast to the findings for **2**, compound **1** does not show any significant differences in the spectral features upon thermal and solvophobic quenching (Fig. S15<sup>†</sup>), once again highlighting that the additional phenyl group of **2** plays an important role in the observed behaviour.

### Analysis of intermolecular interactions

Subsequently, we aimed at examining the packing of the molecules and the intermolecular interactions in the different self-assembled states. For this purpose, we rationalized potential aggregate formations, carried out different Fourier transform-infrared (FT-IR) and NMR experiments and performed density functional theory (DFT) calculations.

The decreased symmetry of the  $\pi$ -surface of **2** makes it challenging to predict the relative arrangement of the molecules upon aggregation. Commonly studied BODIPY dyes with a BF<sub>2</sub> group, such as **1**, are well known to stack in an antiparallel fashion both in solution and the solid state, where both sides of the  $\pi$ -surface have similar chemical environments due to the C<sub>2</sub>-symmetry of these molecules<sup>45,46</sup> (Fig. 3a). However,

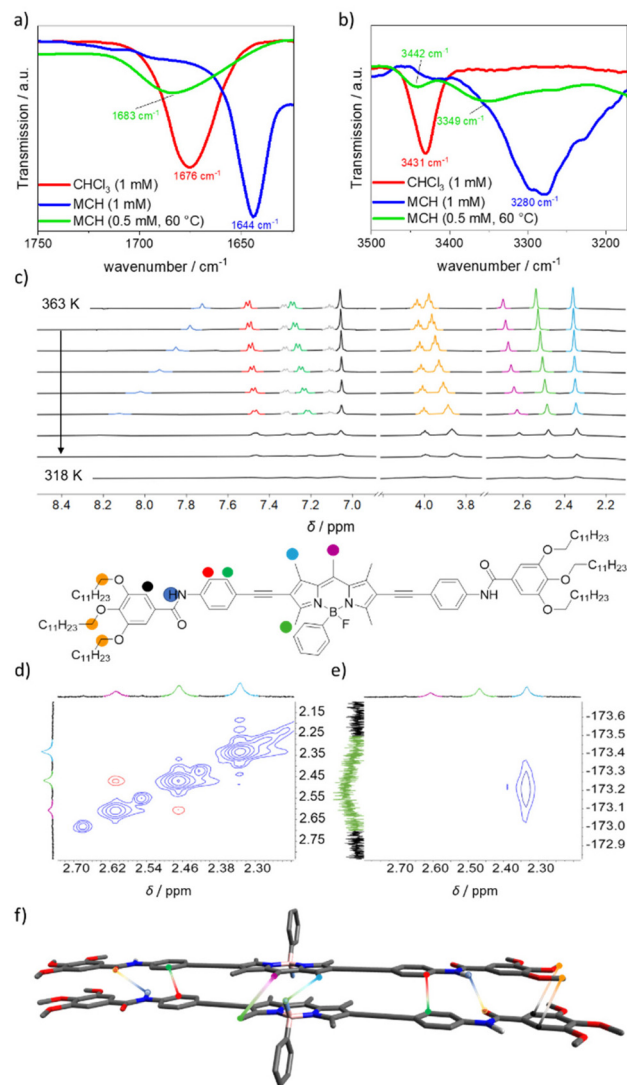


**Fig. 3** (a) Cartoon representation of the typical antiparallel stacking of a C<sub>2</sub>-symmetrical BODIPY derivative. (b) X-ray crystal structure of **C** that consists of F/F type (red frame) and P/P type (blue frame) dimers. (c) Cartoon representation of **C**, representing a potential F/P type stacking of the molecules.



this scenario changes when the  $C_2$ -symmetry is disrupted and two different  $\pi$ -surfaces emerge, as it is the case for **2** (the P- and the F-side, Scheme 1). On this basis, three different kinds of dimer stacks are conceivable (Fig. 3b and c): either both phenyl substituents (P/P) or both fluorine atoms (F/F) are on the same interacting side of the molecules, or they are on different sides (F/P). These three types of dimers may theoretically be able to combine in different ways to build larger aggregates, including potentially unordered structures. To represent the three different packing options, we crystallized the precursor molecule **C** (Fig. 3b, Fig. S22† and Scheme 1), which features a fluorine and a phenyl substituent on the boron, and iodine atoms instead of substituents Y at the 2- and 6-positions, from a dichloromethane solution. Intriguingly, in the crystal structure of **C**, two different types of dimers that could potentially form in a nucleation process for supramolecular polymerization can be observed: the sterically least favoured one (P/P, Fig. 3b, blue frame) along with the sterically most favoured (F/F, Fig. 3b, red frame). The last potential dimer packing that can occur (F/P) would lie in between the P/P and F/F in terms of steric effects. Although this packing is not found in the crystal structure, it likely exists in solution given that it is sterically more favoured than the P/P packing (Fig. 3c).

To examine the intermolecular interactions, we studied the different species **2-H** and **2-I** with different NMR and FT-IR techniques. For the investigation of the aggregate **2-H**, we first performed VT- $^1\text{H}$ -NMR experiments (MCH- $d_{14}$ , 1 mM, Fig. 4c). The formation of **2-H** is accompanied by a high-field shift of the methyl protons of the BODIPY dye at the *meso* (violet) and at the 3- and 5-positions (light green), indicating the formation of  $\pi$ - $\pi$  interactions. The methyl groups at the 1- and 7-positions (light blue) show only weak shifts, which can be explained by a merely weak participation in intermolecular interactions. Furthermore, a low-field shift of the amide protons (dark blue) is observed upon cooling. This can be associated with the formation of hydrogen bonds. At lower temperatures, the signals broaden and finally vanish due to the ongoing aggregation. To gain further information about the participation of hydrogen bonds, we performed FT-IR measurements in  $\text{CHCl}_3$  and MCH to compare the monomeric with the aggregated state (Fig. 4a and b). We found, that the carbonyl stretching shifts from  $\nu_{\text{C=O}} = 1676 \text{ cm}^{-1}$  in  $\text{CHCl}_3$  to  $\nu_{\text{C=O}} = 1644 \text{ cm}^{-1}$  in MCH, demonstrating a change in the chemical environment. Likewise, the N-H vibration shifts to lower wavenumbers when moving from  $\text{CHCl}_3$  ( $\nu_{\text{N-H}} = 3431 \text{ cm}^{-1}$ ) to MCH ( $\nu_{\text{N-H}} = 3280 \text{ cm}^{-1}$ ). These results are in good accordance with the formation of amide-amide N-H...O=C hydrogen bonds.<sup>40,47</sup> Interestingly, by comparing these results with reference BODIPY **1**, we found that the C=O and N-H stretching bands of **2-H** are shifted to lower wavenumbers compared to the aggregate of **1** ( $\Delta\nu_{\text{N-H}}(\text{MCH}) = 40 \text{ cm}^{-1}$ ,  $\Delta\nu_{\text{C=O}}(\text{MCH}) = 13 \text{ cm}^{-1}$ ), which can be traced back to stronger hydrogen bonds in **2-H** despite an increased steric hindrance in **2**. 2D-ROESY- and HOESY-NMR studies were carried out at elevated temperatures to obtain better resolution



**Fig. 4** FT-IR spectra of **2** in different solvents and at different concentrations showing the C=O (a) and N-H (b) stretching frequencies. (c) Temperature-dependent  $^1\text{H}$ -NMR experiments of **2** (MCH- $d_{14}$ , 1 mM,  $\Delta T = 5 \text{ K}$ ). (d)  $^1\text{H}$ - $^1\text{H}$ -ROESY- and (e)  $^1\text{H}$ - $^{19}\text{F}$ -HOESY-NMR studies of **2** at 333 K (1 mM, MCH- $d_{14}$ ) showing the slightly shifted and antiparallel BODIPY interactions, illustrated in (f).

of the signals in the aggregated state (Fig. 4d). In the  $^1\text{H}$ - $^1\text{H}$ -ROESY-NMR, a clear correlation was found between the *meso* methyl group and the methyl groups at the 3- and 5-positions from the BODIPY, indicating an antiparallel and slightly offset packing of the BODIPY dyes. This kind of interaction was already expected from the low-field shift of the corresponding proton signals from VT- $^1\text{H}$ -NMR experiments. The antiparallel packing is supported by correlations between the methyl groups at the 1- and 7-positions and the phenyl substituent (Fig. S17a†) at the boron position as well as with the fluorine atom, visible in the  $^1\text{H}$ - $^{19}\text{F}$ -HOESY spectrum (Fig. 4e). Also, the appearance of two signals for the alkoxy  $\text{OCH}_2$  methylene signal (orange) with a 2:1 ratio in the VT-NMR



experiments can be explained by ROESY-NMR. Here, a coupling of four out of six alkoxy methylene groups of the chains with the aromatic protons of the terminal phenyl ring (black) is observed, indicating a slight translational offset of the molecules (cartoon in Fig. 4f and Fig. S17b†).

For the investigation of **2-I**, the experimental conditions had to be optimized to identify the intermediate species. Hence, we carried out NMR and FT-IR experiments at a concentration of 0.5 mM, as it allowed us to correlate our results with UV-vis experiments. VT-UV-vis experiments revealed that **2-I** can be isolated at 333 K (Fig. S16†). Below this temperature, the formation of **2-H** was observed, similar to the previous experiments at lower concentrations. This also correlates with the results of VT-<sup>1</sup>H-NMR (Fig. S18†). Here, a signal broadening was noticed below 333 K, which is expected for the formation of extended aggregates. The VT-<sup>1</sup>H-NMR shows similar characteristics compared to the higher concentrated solution (Fig. 4). FT-IR experiments (MCH, 0.5 mM, 333 K) were conducted to gain more information about the hydrogen bonds (Fig. 4a). The C=O vibration ( $\nu_{\text{C=O}} = 1683 \text{ cm}^{-1}$ ) appears at nearly identical wavenumbers compared to the monomeric species in chloroform (1 mM,  $\nu_{\text{C=O}} = 1676 \text{ cm}^{-1}$ ), but shows a more pronounced broadening. The N-H vibration splits into two bands: a stretch that correspond to free N-H groups ( $\nu_{\text{N-H}} = 3442 \text{ cm}^{-1}$ ) and a band at lower wavenumbers ( $\nu_{\text{N-H}} = 3349 \text{ cm}^{-1}$ ) that indicates weak hydrogen bonding between the molecules. We therefore conclude that the formation of **2-I** is driven by  $\pi$ - $\pi$  interactions and rather diffuse and weak hydrogen bonds that do not include the carbonyl groups. The weak coordination of the amide groups goes in line with the lack of exciton coupling of the Y substituents that was observed in the UV-vis experiments (Fig. 2a). In 2D-ROESY experiments, however, we could find similar interactions in **2-I** and **2-H** that imply similar packings of the molecules (Fig. S19†). Nevertheless, the absence of strong hydrogen bonds in **2-I** indicates that the formation of **2-H** requires a rearrangement of the molecules, *e.g.* *via* the monomer.

To ultimately rationalize the packing and the preferred contact sides between the neighbouring molecules (F/F, P/P or F/P contacts) in **2-H**, we performed theoretical calculations. Previous NMR experiments do not allow to distinguish between the stacking possibilities due to the intrinsically similar interactions of the neighbouring molecules in the proposed packings. Hence, DFT (B3LYP/6-31G(d,p)) calculations were employed to optimize the monomer and compare the three proposed dimers. The dodecyloxy side chains at the peripheral phenyl rings were replaced with methoxy groups to reduce the computational costs. Subsequently, we performed calculations on the F/F, P/P and F/P dimers (Fig. S23†). From the optimized structure, it becomes evident that the P/P stacking is the energetically least favoured stacking ( $\Delta G = +29.851 \text{ kJ mol}^{-1}$ ). Due to the steric repulsion, only one amide group is involved in hydrogen bonds and the BODIPY chromophores only overlap slightly. The situation changes for the F/P and the F/F dimers. Both exhibit hydrogen bonds between each of the amide groups and both are well stabilized. The free

Gibbs energy favours the F/P assembly by  $11.673 \text{ kJ mol}^{-1}$  when compared to the F/F packing. Accordingly, we infer that the formation of F/P interactions in the supramolecular polymer **2-H** is the most probable scenario. However, given the similar stability of these two packings, we hypothesize that the competition between these two packing modes might be responsible for the observed pathway complexity.

## Conclusions

In conclusion, we have reported the supramolecular polymerization in nonpolar media of a new BODIPY derivative **2** with two unequal  $\pi$ -surfaces, stemming from the unsymmetrical functionalization of the boron atom with a fluorine and a phenyl substituent. To rationalize the impact of the symmetry breaking, the self-assembly of **2** was compared with that of a model  $C_2$ -symmetrical BODIPY compound **1** with a  $\text{BF}_2$  group. We found that **2** undergoes an H-type supramolecular polymerization (**2-H**) with similar spectral patterns to **1**. However, not only does the final morphology differ (short 1D fibers for **1** vs. long fibers for **2**), but also the pathway in which the aggregate is reached. In contrast to a single-step self-assembly of **1**, the introduction of an additional bulky (phenyl) substituent to the BODIPY boron atom in **2** increases the complexity of the self-assembly by enabling an intermediate species that most probably involves anticooperative effects. These findings can be associated with increased steric repulsion and the multiple possibilities in which unsymmetrical molecules can interact during nucleation. As a result, the system undergoes a monomer  $\rightarrow$  intermediate transformation upon cooling to intermediate temperatures (*ca.* 300–310 K), and a subsequent conversion to the final H-type aggregate (**2-H**) upon lowering the temperature to 263 K. We found that the packing of the intermediate state **2-I** (small nanoparticles) is similar to the final H-type fibers albeit with weaker and less ordered hydrogen bonds. Although three different dimer packings are feasible (F/F, P/P and F/P dimers), theoretical calculations suggest that only the F/F and F/P-type packings would allow for simultaneous H-bonding and aromatic interactions leading to the experimentally observed extended growth. Our findings introduce a new molecular design strategy for pathway-dependent controlled supramolecular polymerization.

## Author contributions

Tobias Benedikt Tischer: synthesis, investigation, methodology, writing – original draft. Zulema Fernández: theoretical calculations. Lorenz Borsdorf: AFM measurements. Constantin Daniliuc: X-ray diffraction analysis. Soichiro Ogi: conceptualization, writing – review and editing. Shigehiro Yamaguchi: writing – review and editing. Gustavo Fernández: conceptualization, writing – review and editing, supervision and funding acquisition.



## Data availability

The data supporting this article have been included as part of the ESI.† Deposition number 2386585† (for C) contains the supplementary crystallographic data for this paper.

## Conflicts of interest

There are no conflicts to declare.

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