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# Ligand-modulated metallophilicity: influence of $\pi$ -acceptor and $\sigma$ -donor strength†

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Metal–metal (M–M) closed-shell interaction, also known as metallophilicity, is frequently observed in  $d^{10}$  and  $d^8$  metal complexes featuring a close M–M distance. It has shown a significant impact on diverse chemical systems, influencing structural, catalytic, and photophysical properties. The strength of both M–M interactions and the resulting M–M distances is highly dependent on various types of coordinating ligands. Recent studies have revealed that metallophilicity is repulsive in nature due to strong M–M Pauli repulsion (Q. Wan, J. Yang, W.-P. To and C.-M. Che, Strong metal–metal Pauli repulsion leads to repulsive metallophilicity in closed-shell  $d^8$  and  $d^{10}$  organometallic complexes, *Proc. Natl. Acad. Sci. U. S. A.*, 2021, **118**, e2019265118). However, little is known about the role of ligands in M–M repulsions. Here, we elucidate how metal–ligand (M–L) coordination modulates M–M repulsion through two key mechanisms:  $\pi$ -backbonding and  $\sigma$ -donor interactions. By systematically evaluating ligands spanning a spectrum of  $\pi$ -accepting and  $\sigma$ -donating strengths, we uncover opposing ligand effects. Strong  $\pi$ -backbonding weakens M–M Pauli repulsion, enabling shorter intermetallic distances, whereas the  $\sigma$ -donating interaction increases the repulsion, lengthening M–M contacts. These computational insights establish a ligand-design framework for tuning metallophilicity in closed-shell metal complexes and advance the fundamental understanding of M–M interactions from the perspective of M–L coordination.

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## Introduction

Metallophilicity has been extensively studied in a wide range of  $d^8$  and  $d^{10}$  metal complexes, including Au(I), Au(III), Ag(I), Cu(I), Pt(0), Pd(0), Pt(II), Pd(II), Ir(I) and Rh(I).<sup>1–12</sup> These closed-shell metal complexes have important applications across various research fields, ranging from supramolecular self-assembly to advanced (opto)electronic materials and photocatalytic systems.<sup>2,7,8,13–16</sup> In supramolecular chemistry, M–M interactions are recognized as a key driving force guiding the self-assembly process.<sup>7,17,18</sup> In optoelectronics, incorporating M–M interactions into excited states would facilitate radiative decay and is particularly important in the design of phosphorescent materials.<sup>19–21</sup> In photocatalysis, M–M interactions facilitate hydrogen atom abstraction through an inner-sphere electron transfer mechanism.<sup>22–24</sup> A comprehensive understanding of M–M interactions is crucial, as it can provide valuable guidance for the molecular design. Despite its broad utility, the fundamental nature of metallophilicity has long been debated.<sup>3,5,6,25–35</sup> Recent breakthroughs show that metallophilicity is inherently repulsive, dominated by M–M Pauli repulsion

between metal centers rather than stabilizing interactions, as supported by computational and spectroscopic studies across Au(I), Ag(I), Cu(I), Rh(I), Pt(II), and Pd(II) systems.<sup>36–39</sup>

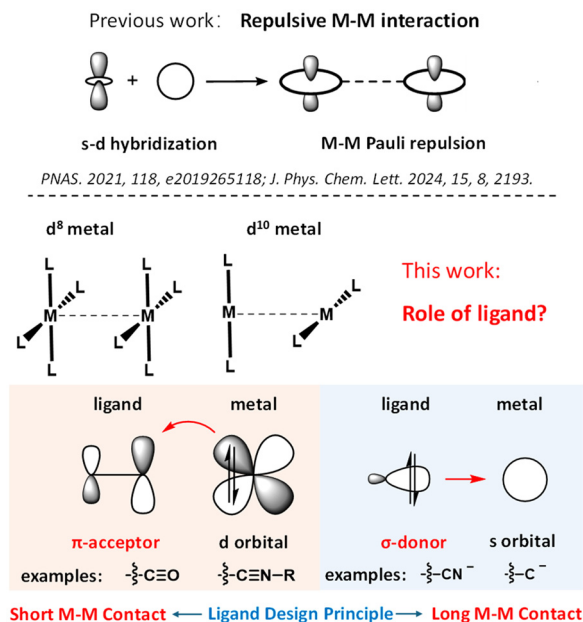
While the repulsive nature of metallophilicity has been revealed, the role of ligands in modulating this repulsion remains a critical unresolved question. Ligands are well known to exert dramatic control over M–M distances: for instance, Au(I) complexes exhibit Au–Au separations spanning from non-interacting ( $>3.8$  Å) to significantly shorter than the sum of van der Waals radii ( $\sim 2.5$  Å), depending on the ligand identity, as summarized in several comprehensive reviews.<sup>3,5</sup> Such variability highlights the dual capacity of ligands to either amplify or mitigate repulsive M–M interactions, yet the mechanistic origins of this modulation remain poorly defined. In this work, we address this gap by systematically dissecting how  $\sigma$ -donor and  $\pi$ -acceptor ligands tune M–M Pauli repulsion (Scheme 1).

Transition metal atoms can coordinate with a wide range of ligands, which are broadly categorized as  $\pi$ -acceptors,  $\pi$ -donors or  $\sigma$ -donors (Scheme 1).<sup>40</sup>  $\pi$ -Acceptor ligands, such as CO and nitriles ( $R-C\equiv N$ ), feature low-energy  $\pi^*$  orbitals that accept electron density from filled metal d-orbitals, enabling  $\pi$ -backbonding interactions. In contrast, strong  $\sigma$ -donor ligands like cyanide ( $CN^-$ ) and alkynides ( $R-C\equiv C^-$ ) coordinate through lone pairs, donating electron density to vacant metal orbitals (d/s/p). These contrasting bonding mechanisms exert

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**Scheme 1** Top: previous work, repulsive M-M interaction in closed-shell d<sup>8</sup> and d<sup>10</sup> transition metal complexes. Bottom: this work, the role of  $\pi$ -acceptor and  $\sigma$ -donor ligands in M-M interactions.

striking effects on metalphilicity: structural studies reveal that  $\pi$ -acceptors promote shortened M-M distances, while  $\sigma$ -donors favour elongation, as shown in Fig. 1 and 7.<sup>3,5</sup>

Here, we bridge this gap through computational investigation of ligand-mediated M-M interactions. By utilizing energy decomposition analysis (EDA) and natural orbital for chemical valence (NOCV) methods, we demonstrate that  $\pi$ -backbonding interactions delocalize electron density from metal centers to ligand  $\pi^*$  orbitals, attenuating M-M Pauli repulsion and enabling closer M-M contacts. Conversely, strong  $\sigma$ -donor ligands enhance metal ( $n+1$ )s-nd hybridization, intensifying Pauli repulsion and elongating M-M distances. These insights explain well the experimental structural trends shown in Fig. 1 and 7. Our findings not only resolve longstanding ambiguities in ligand-metalphilicity relationships but also establish ligand-design principles to tailor M-M interactions in functional materials for supramolecular assembly, catalysis, and optoelectronic applications.

## Results and discussion

### Role of $\pi$ -acceptor ligands

To investigate the role of  $\pi$ -acceptor ligands, we analysed a series of Rh(I) complexes with or without  $\pi$ -acceptor ligands (Fig. 1). Specifically, we compared the M-M distances in the following pairs using their corresponding X-ray crystal structures:

- Rh(terpy)Cl (**Rh-1**)<sup>41</sup> vs. [Rh(terpy)N≡CCH<sub>3</sub>]<sup>+</sup> (**Rh-2**)<sup>42</sup>
- Rh(CO)<sub>2</sub>ClNH<sub>2</sub>CH<sub>3</sub> (**Rh-3**)<sup>43</sup> vs. [Rh(CO)<sub>2</sub>(N≡CCH<sub>3</sub>)<sub>2</sub>]<sup>+</sup> (**Rh-4**).<sup>44</sup>

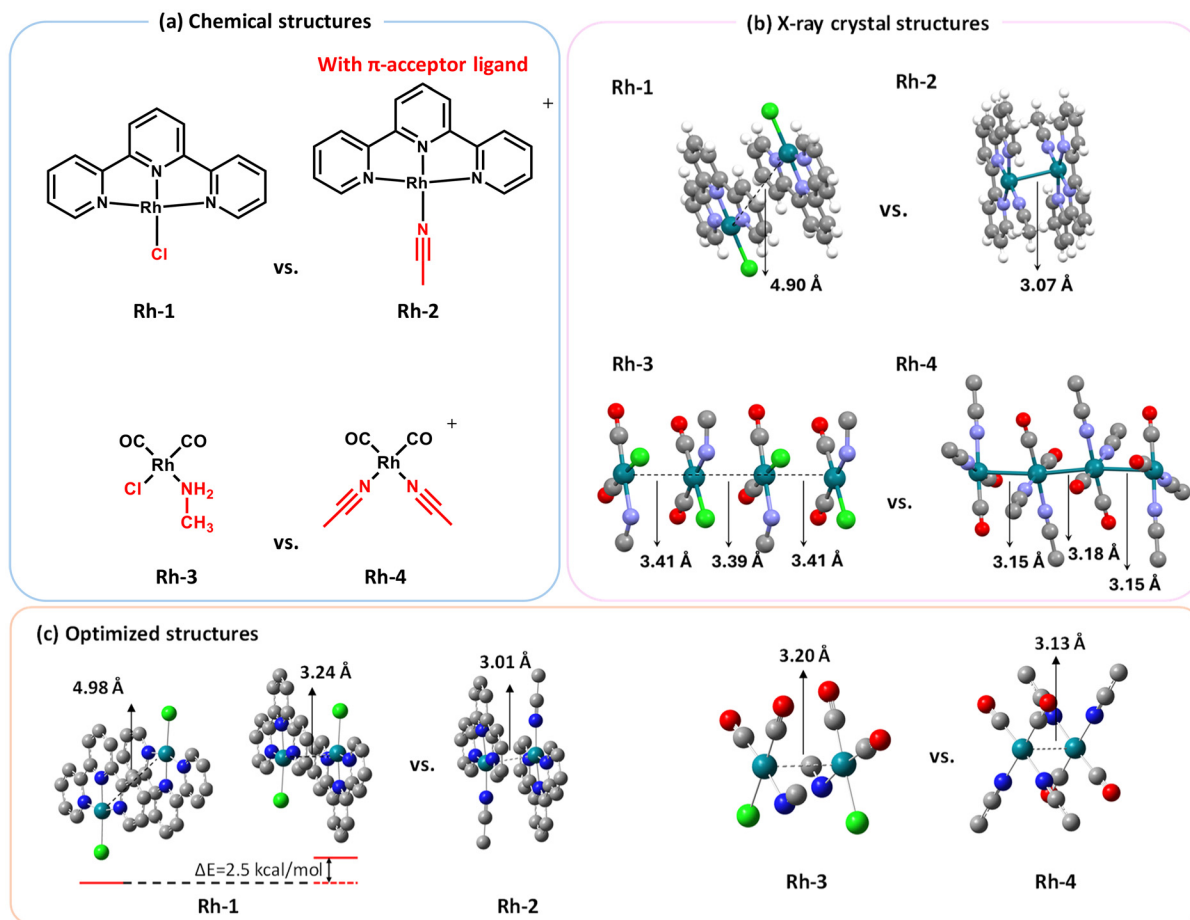
Here, **Rh-2** and **Rh-4** incorporate  $\pi$ -acceptor ligands (acetonitrile), while **Rh-1** and **Rh-3** lack such ligands. Notably, aside from Cl<sup>−</sup> and N≡CCH<sub>3</sub>, both **Rh-1** and **Rh-2** coordinate to the same terpyridine (terpy) ligand; aside from Cl<sup>−</sup>/NH<sub>2</sub>CH<sub>3</sub> and N≡CCH<sub>3</sub>, both **Rh-3** and **Rh-4** coordinate to the same CO ligand. These comparisons exclude the potential influence of other ligand effects such as London dispersion and so on. As shown in Fig. 1, key structural trends emerged in the structural comparison: replacing the Cl<sup>−</sup> ligand in **Rh-1** with the  $\pi$ -acceptor ligand CH<sub>3</sub>C≡N (**Rh-2**) shortened the Rh-Rh distance from 4.90 Å (**Rh-1**) to 3.07 Å (**Rh-2**). Substituting the Cl<sup>−</sup> and NH<sub>2</sub>CH<sub>3</sub> ligands in **Rh-3** with CH<sub>3</sub>C≡N ligands in **Rh-4** reduced the Rh-Rh distance from 3.39–3.41 Å (**Rh-3**) to 3.15–3.18 Å (**Rh-4**). These results demonstrate that coordination of  $\pi$ -acceptor ligand is closely related to the short M-M distances observed in the crystal structures of d<sup>8</sup> Rh(I) complexes.

The M-M distance trend observed in X-ray crystal structures was further validated through computational optimization of dimeric structures, thereby excluding contributions from the crystal packing effect. As shown in Fig. 1c, the optimized structures retain the key experimental trends (Fig. 1a and b): Rh complexes with  $\pi$ -acceptor ligands exhibit shorter Rh-Rh distances than analogues without such ligands (**Rh-1** vs. **Rh-2**; **Rh-3** vs. **Rh-4**). Notably, both optimized configurations of the [**Rh-1**]<sub>2</sub> dimer shown in Fig. 1c maintain longer Rh-Rh distances than those of [**Rh-2**]<sub>2</sub>. This consistency between solution-phase calculations and solid-state structural data indicates that the observed distance variations are primarily driven by ligand electronic properties rather than crystal packing or solvent effects.

For d<sup>8</sup>-d<sup>8</sup> closed-shell systems, the role of  $\pi$ -acceptor ligands in reducing M-M Pauli repulsion can be rationalized through an illustrative molecular orbital (MO) diagram shown in Scheme 2. First, we need to consider the ligand field splitting diagram of a metal with a d<sup>8</sup> electronic configuration.<sup>40</sup> As shown in Scheme 2, for a d<sup>8</sup> metal complex (e.g., Rh(I), Ir(I), Pt(II), or Pd(II)) with square planar coordination geometry, d<sub>x<sup>2</sup>-y<sup>2</sup></sub> is a strongly antibonding orbital with respect to M-L interactions and remains unoccupied, while d<sub>z<sup>2</sup></sub> is a weakly antibonding orbital and remains occupied.<sup>40</sup> The remaining three orbitals (d<sub>xz</sub>, d<sub>xy</sub> and d<sub>yz</sub>) are degenerate, non-bonding and occupied M-L orbitals.<sup>40</sup>

Next, we need to derive the formation of M-M bonding and antibonding orbitals from the occupied d orbitals of two d<sup>8</sup> metal atoms.<sup>40</sup> Defining the M-M bond axis as the z-axis, then the two d<sub>z<sup>2</sup></sub> orbitals can overlap in a  $\sigma$ -type manner to form a bonding  $\sigma$  orbital and an antibonding  $\sigma^*$  orbital (Scheme 2). Similarly, the d<sub>xz</sub> and d<sub>yz</sub> orbitals can overlap with another set of d<sub>xz</sub> and d<sub>yz</sub> orbitals in a  $\pi$ -fashion, forming two degenerate bonding  $\pi$  orbitals and two degenerate antibonding  $\pi^*$  orbitals. Since the  $\pi$  orbital overlap is smaller than the  $\sigma$  overlap, the energy splitting between the bonding and antibonding  $\pi$  orbitals is smaller than that between the bonding and antibonding  $\sigma$  orbitals. Furthermore, the d<sub>xy</sub> orbital can overlap with another d<sub>xy</sub> orbital in a  $\delta$ -type manner, resulting in





**Fig. 1** Chemical structures (a), X-ray crystal structures (b) and optimized structures (c) of  $d^8$  Rh(I) metal complexes with or without  $\pi$ -acceptor ligands studied in this work. Calculations were performed at the PBE0-GD3BJ/SDD(Rh)/6-31G\* (other elements) level of theory using the PCM solvent model for acetonitrile (ACN) in Gaussian 16.

bonding  $\delta$  orbitals and antibonding  $\delta^*$  orbitals. The  $\delta$  orbital overlap is even smaller than the  $\pi$  overlap, leading to an energy splitting between the bonding  $\delta$  and the antibonding  $\delta^*$  orbitals that is smaller than those for the  $\pi$  and  $\pi^*$  orbitals, as illustrated in Scheme 2.

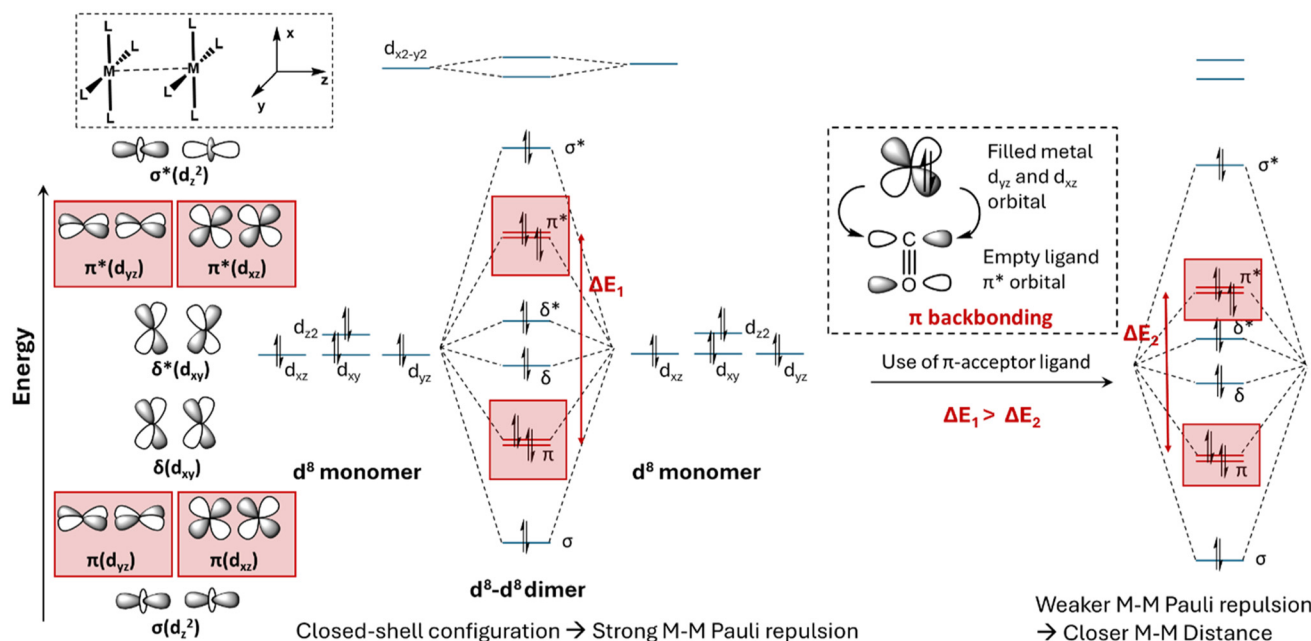
In Scheme 2, the M-M bonding orbitals—comprising one  $\sigma$ , two  $\pi$ , and one  $\delta$  orbital—along with the corresponding antibonding orbitals (one  $\sigma^*$ , two  $\pi^*$ , and one  $\delta^*$ ), are fully occupied. As demonstrated in our previous work on Rh(I), Pt(II) and Pd(II) systems, the closed-shell electronic configuration of two  $d^8$  metals with planar coordination geometry results in strong M-M Pauli repulsion, the magnitude of which is reflected in the energy splitting between the occupied M-M bonding and antibonding orbitals.<sup>36,37</sup>

The M-L coordination bond plays a critical role in determining both the energy levels and electron density distribution of the metal  $d$  orbitals.<sup>40</sup> This, in turn, modulates the extent of M-M orbital overlap, thereby influencing the magnitude of M-M Pauli repulsion and the resulting M-M distances observed in the crystal structures of these complexes. Among various types of M-L coordination bonds,  $\pi$ -backbonding is a

particularly important interaction in transition metal complexes. This interaction typically occurs when a ligand, such as carbon monoxide, accepts electron density from the metal through the overlap of a metal  $d$  orbital and a ligand  $\pi^*$  orbital. This process, termed “back-bonding,” stabilizes the metal’s  $d$  orbitals as the ligand acts as a  $\pi$ -acceptor. Consequently, M-M  $\pi$  and  $\pi^*$  orbitals are stabilized, as illustrated in Scheme 2. Since M-M Pauli repulsion is positively correlated with the degree of overlap between two occupied  $d$  orbitals,<sup>45,46</sup> it is expected that the presence of  $\pi$ -acceptor ligands weakens M-M Pauli repulsion arising from  $\pi$ -type interactions, leading to shorter M-M distances.

To computationally prove the speculation illustrated in Scheme 2, we performed Energy Decomposition Analysis (EDA),<sup>47–49</sup> Natural Orbitals for Chemical Valence (NOCV) analysis, molecular orbital (MO) diagram calculations, and high-level Domain-based Local Pair Natural Orbital Coupled Cluster with Singles, Doubles, and perturbative Triples (DLPNO-CCSD(T)) computations on Rh complexes, using the AMS (Amsterdam Modeling Suite) 2024.1<sup>50,51</sup> and ORCA 6.0 package.<sup>52,53</sup>





**Scheme 2** Schematic illustration of the role of the  $\pi$ -acceptor ligand in weakening the M–M Pauli repulsion for  $d^8$ – $d^8$  closed-shell metal complexes. Schematic illustration of the MO for  $d^{10}$ – $d^{10}$  closed-shell metal complexes has been provided in the ESI.†

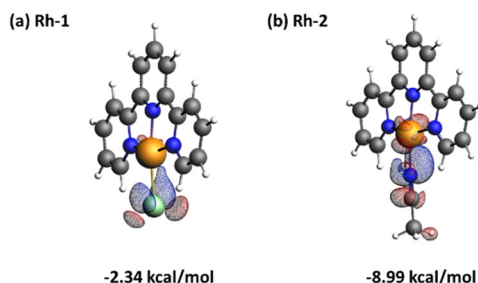
We first examined the formation of  $\pi$ -backbonding interactions between the metal and the ligand in complexes **Rh-1** and **Rh-2**. EDA-NOCV<sup>54</sup> analysis decomposes electron density changes into chemically understandable components using fragment molecular orbitals. Herein, the energetic quantification of M–L  $\pi$  interactions was calculated for the Rh–Cl bonds in the optimized **Rh-1** monomer and the Rh–N bonds in the optimized **Rh-2** monomer. The NOCV deformation density channels are depicted in Fig. 2 with the red and blue regions symbolizing the depletion and accumulation of electron density, respectively. One NOCV deformation density channel corresponding to Rh–Cl  $\pi$ -bonds shows electron transfer from the Rh atom to the Cl atom, with an orbital energy of  $-2.34 \text{ kcal mol}^{-1}$  of **Rh-1**. Similarly, one NOCV deformation density channel, representing Rh–N  $\pi$ -bonds, exhibits electron

transfer from the Rh atom to the N atom, with a stronger orbital energy of  $-8.99 \text{ kcal mol}^{-1}$  of **Rh-2**.

The difference in the energy of NOCV calculations between **Rh-1** and **Rh-2** indicates different  $\pi$ -accepting abilities between  $\text{Cl}^-$  and the  $\text{N}\equiv\text{CCH}_3$  ligand. For the  $\text{N}\equiv\text{CCH}_3$  ligand, the N atom forms a triple bond with the C atom, consisting of one  $\sigma$  bond and two  $\pi$  bonds. The two empty  $\pi^*$  orbitals in the  $\text{N}\equiv\text{CCH}_3$  ligand lie at relatively low energy, enabling efficient acceptance of d electron density from the Rh atom. The NOCV analysis suggests that stronger  $\pi$ -backbonding interaction occurs in **Rh-2**, leading to a reduction in the d orbital's electron density in **Rh-2** compared to **Rh-1**. This is expected to reduce the  $d\pi$ – $d\pi^*$  orbital splitting energy and weaken the M–M Pauli repulsion in the **Rh-2** dimer compared to the **Rh-1** dimer. Next, we conducted MO calculations on the **Rh-1** and **Rh-2** dimers to validate this hypothesis.

Based on this consideration, we examined and compared the  $d\pi$ – $d\pi^*$  orbital splitting energy in the optimized **Rh-1** and **Rh-2** dimers using MO calculations. As shown in Fig. 3, the orbital splitting energy between the M–M  $\pi(d_{xz}$ – $d_{xz})$  bonding orbital and the  $\pi^*(d_{xz}$ – $d_{xz})$  antibonding orbital was calculated to be  $0.408 \text{ eV}$  in **Rh-1** and  $0.393 \text{ eV}$  in **Rh-2**. The orbital splitting energy for M–M  $\pi$ -type interactions between the two Rh atoms is consistently smaller in **Rh-2** than that in **Rh-1**. In comparison to the **Rh-1** dimer, we expect that the smaller  $d\pi$ – $d\pi^*$  orbital energy splitting in the **Rh-2** dimer would lead to decreased M–M Pauli repulsion.<sup>36,37</sup> Finally, we calculated the M–M Pauli repulsion in the **Rh-1** and **Rh-2** dimers to prove this relationship.

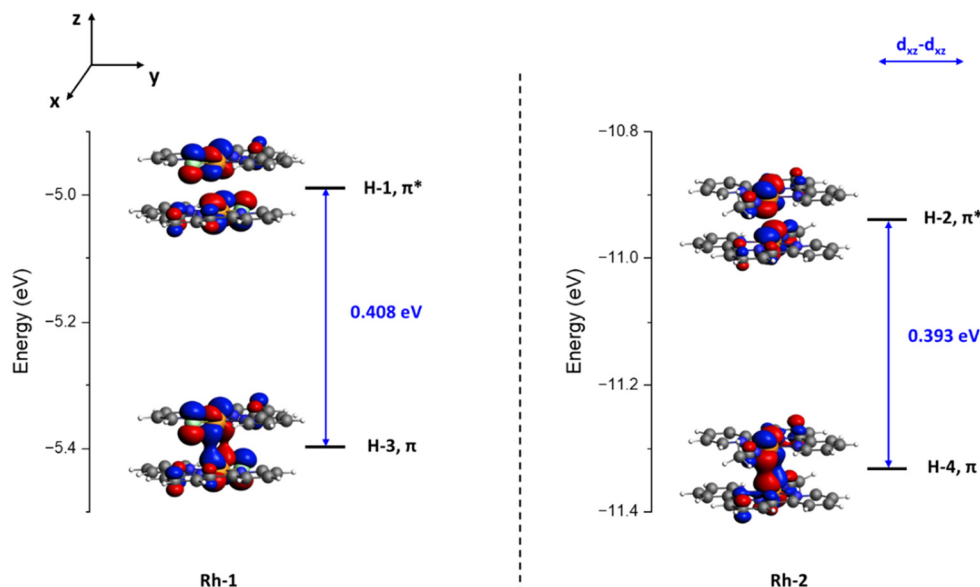
M–M Pauli repulsion for the **Rh-1** and **Rh-2** dimers has been calculated using the EDA method, and the results are



**Fig. 2** Calculated EDA-NOCV deformation density contributions to (a) the Rh → Cl  $\pi$ -type interaction in **Rh-1** and (b) the Rh → N  $\pi$ -type interaction in **Rh-2**, together with the calculated orbital interaction energies corresponding to each  $\pi$  interaction. Isovalue = 0.001. Charge transfer occurred from red to blue.







**Fig. 3** Calculated MO diagram for the **Rh-1** dimer (left) and the **Rh-2** dimer (right), highlighting the  $d\pi$  bonding and  $d\pi^*$  antibonding orbitals. Isovalue = 0.03.

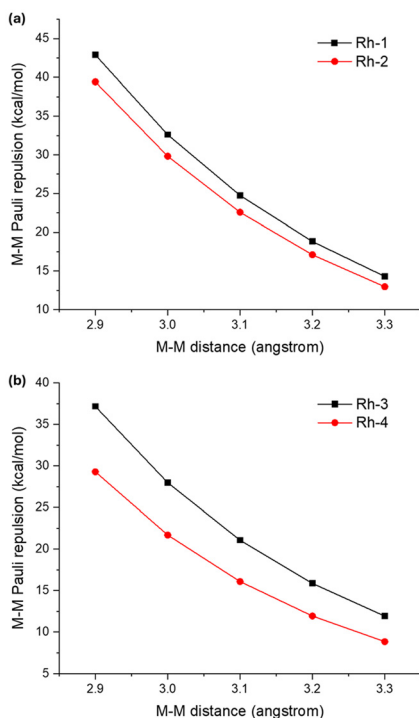
plotted in Fig. 4. In the calculations, the M–M distance in the **Rh(i)** dimers **1** and **2** was set from 3.3 Å to 2.9 Å, to compare the M–M Pauli repulsion at equal M–M distances. As shown in Fig. 4a, Rh–Rh Pauli repulsion increased with the decrease of

the Rh–Rh distance in both the **Rh-1** and **Rh-2** dimers, and the M–M Pauli repulsion in the **Rh-1** dimer is always stronger than that in the **Rh-2** dimer. For instance, at an Rh–Rh distance of 3.0 Å, the M–M Pauli repulsion was 32.60 kcal mol<sup>−1</sup> in the **Rh-1** dimer, compared to 29.82 kcal mol<sup>−1</sup> in the **Rh-2** dimer, representing a 2.78 kcal mol<sup>−1</sup> reduction. This calculated trend in M–M Pauli repulsion for **Rh-1** and **Rh-2** aligns with the trend observed in the M–M distances from their respective X-ray crystal structures (Fig. 1). The trend in Rh–Rh Pauli repulsion has been further validated through DLPNO-CCSD(T) calculations, which confirm that the **Rh-2** complex—featuring the  $\pi$ -accepting CH<sub>3</sub>CN ligand—exhibits weaker metal–metal Pauli repulsion than **Rh-1** (Fig. S1†).

The results of EDA, NOCV and MO calculations demonstrate a mechanistic pathway: the ligand's  $\pi$ -accepting ability directly influences the M–M  $d\pi$ – $d\pi^*$  orbital splitting energy, which in turn determines the magnitude of M–M Pauli repulsion. To summarize, coordination to  $\pi$ -acceptor ligands decreases the  $d\pi$ – $d\pi^*$ (M–M) orbital splitting and weakens M–M Pauli repulsion, facilitating the formation of close M–M contact in the **Rh-2** complex compared to **Rh-1**.

Similarly, NOCV calculations were performed on the **Rh-3** and **Rh-4** monomers, revealing total Rh → L  $\pi$ -bonding interactions of −17.42 and −18.99 kcal mol<sup>−1</sup>, respectively (Fig. S2†). At every M–M distance in Fig. 4b, stronger M–M Pauli repulsion is observed in the **Rh-3** dimer (with Cl<sup>−</sup>/NH<sub>2</sub>CH<sub>3</sub> ligands) than in the **Rh-4** dimer (with  $\pi$ -acceptor N≡CCH<sub>3</sub> ligands). The NOCV and EDA results suggest that the stronger  $\pi$ -backbonding interaction in **Rh-4** contributes to weaker M–M Pauli repulsion, and consequently, a shorter M–M distance in **Rh-4** (3.15–3.18 Å) than in **Rh-3** (3.39–3.41 Å) shown in Fig. 1.

In addition to the Rh(i)  $d^8$  systems, the relationship between the  $d^{10}$ – $d^{10}$  Pauli repulsion and the  $\pi$  acceptor ligand



**Fig. 4** Calculated M–M Pauli repulsion of (a) **Rh-1** and **Rh-2** dimers and (b) **Rh-3** and **Rh-4** dimers, by changing the M–M distance. For EDA calculations at non-equilibrium M–M distances, fragment geometries were frozen while varying metal–metal separation.



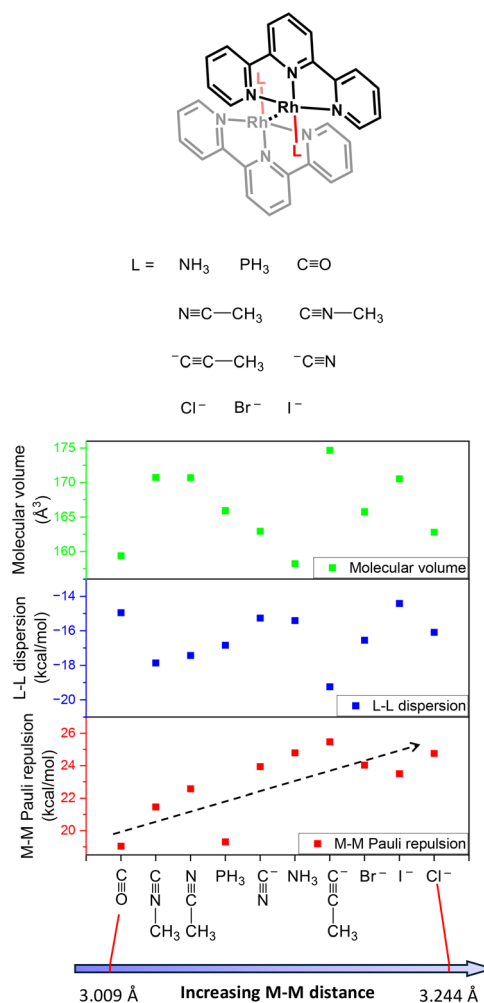
is demonstrated in the MO diagram of Scheme S1†. Different from the  $d^8$  closed-shell metal atom with four fully occupied metal d orbitals, the  $d^{10}$  metal atom has five fully occupied d orbitals.<sup>40</sup> These d orbitals, when interacting between two metal atoms, contribute to one  $\sigma$ -type bonding and antibonding orbital, two  $\pi$ -type bonding and antibonding orbitals, and two  $\delta$ -type bonding and antibonding orbitals. Similarly, the formation of  $\pi$ -backbonding interactions between the  $d^{10}$  metal atom and the ligand stabilizes the M–M  $\pi$ -type bonding and antibonding orbitals, leading to weakened M–M Pauli repulsion in  $d^{10}$ – $d^{10}$  metal complex systems (Scheme S1†).

To further validate the relationship between ligand  $\pi$ -accepting ability and M–M Pauli repulsion, we systematically evaluated diverse ligands ( $L = \text{NH}_3, \text{PH}_3, \text{CO}, \text{NCCH}_3, \text{CNCH}_3, \text{CCCH}_3, \text{CN}, \text{Cl}, \text{Br}, \text{I}$ ) in  $[\text{Rh}(\text{terpy})\text{L}]_2$  dimers (model system). Following structural optimization, we correlated M–M distances with three parameters: molecular volume (reflecting steric effects), ligand–ligand (L–L) dispersion interactions, and M–M Pauli repulsion energy. In Fig. 5, M–M distances show no clear correlation with the molecular volume or L–L dispersion but exhibit a positive correlation with both M–M Pauli repulsion energy and ligand  $\pi$ -accepting strength (Fig. 5, 6 and Table S1†). As shown in Fig. 5, increasing M–M distances in  $[\text{Rh}(\text{terpy})\text{L}]_2$  dimers ( $L = \text{CO} \rightarrow \text{Cl}^-$ ) correlate with progressively stronger M–M Pauli repulsion. Correspondingly, enhanced Pauli repulsion occurs concomitantly with weakened  $\text{Rh} \rightarrow \text{L} \pi$ -interactions, as evidenced by the EDA-NOCV analysis in Fig. 6. This absence of steric/L–L dispersion correlations, coupled with electronic dependence, provides evidence that ligand  $\pi$ -accepting capability primarily governs M–M distances through Pauli repulsion modulation.

Notably, for relatively large ligands, steric or dispersion effects could potentially dominate.<sup>33</sup> For instance, the dispersion interactions and softness of ligand fragments have previously been reported to govern intermolecular interactions in the  $\text{Rh}(\text{I})$  and  $\text{Au}(\text{I})$  closed-shell systems, while the metals play a relatively minor role.<sup>32–34,55,56</sup> The magnitude of ligand–ligand (L–L) dispersion interactions strongly depends on the ligand size.<sup>33,34</sup> Furthermore, recent studies demonstrate that ligand–metal interplay critically influences non-covalent interactions.<sup>55,57</sup> However, within the scope of this work—focused on investigating the electronic tuning effects *via* ligands—our analyses confirm that the observed M–M distance trends are linked to the  $\pi$ -accepting ligand modulation ability of M–M Pauli repulsion in smaller ligand systems.

### Role of $\sigma$ -donor ligands

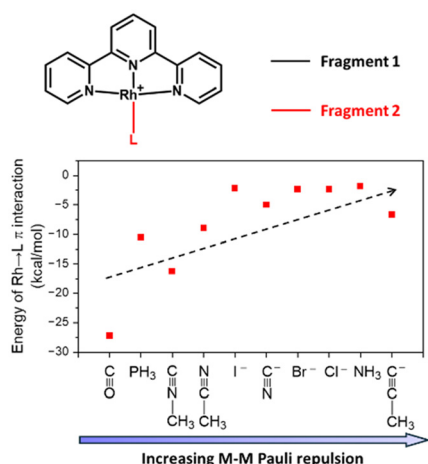
To explore the effect of the  $\sigma$ -donating ability of ligands on M–M interactions, several examples of  $\text{Pt}(\text{II})$  and  $\text{Au}(\text{I})$  complexes have been studied (Fig. 7). We compared the M–M distances in  $[\text{Au}(\text{SCN})_2]^-$  (**Au-1**) and  $[\text{Au}(\text{CN})_2]^-$  (**Au-2**), where  $\text{CN}^-$  is a stronger  $\sigma$ -donor. As shown in Fig. 7a, replacing the  $\text{SCN}^-$  ligand in **Au-1** with the  $\text{CN}^-$  ligand in **Au-2** resulted in a longer Au–Au distance of 3.64 Å, compared to 3.0064–3.070 Å in **Au-1** ( $\text{K}^+$  as the counteranion).<sup>55,58–60</sup> We also compared the M–M distances in  $[\text{Pt}(\text{terpy})\text{C}\equiv\text{CC}_6\text{H}_5]^+$  (**Pt-1**),  $\text{Pt}(\text{C}^-\text{N}^-\text{N})\text{C}\equiv\text{CC}_6\text{H}_5$  (**Pt-**



**Fig. 5** Model system  $[\text{Rh}(\text{terpy})\text{L}]_2$  and correlation analysis. Chemical structures of the  $[\text{Rh}(\text{terpy})\text{L}]_2$  dimer (top) with the corresponding plots of the molecular volume, L–L dispersion interaction, and M–M Pauli repulsion versus the M–M distance in the fully optimized geometries (bottom). All data points are tabulated in Table S1†.

**2**),  $[\text{Pt}(\text{C}^-\text{N}^-\text{N})\text{N}\equiv\text{CCH}_3]^+$  (**Pt-3**), and  $\text{Pt}(\text{N}^-\text{C}^-\text{N})\text{C}\equiv\text{CC}_6\text{H}_5$  (**Pt-4**), where **Pt-1** and **Pt-3** complexes have only one strong  $\sigma$ -donor ligand of  $\text{C}^-$ , while **Pt-2** and **Pt-4** complexes have two  $\text{C}^-$  ligands.<sup>61–64</sup> As shown in Fig. 7b, a close Pt–Pt distance of 3.37 Å was observed in the crystal structure of **Pt-1**. By changing the terpyridine ligand ( $\text{N}^-\text{N}^-\text{N}^-$ ) into the  $(\text{C}^-\text{N}^-\text{N})^-$  ligand and  $(\text{N}^-\text{C}^-\text{N})^-$  ligand, the distances became longer in **Pt-2** (6.43 Å) and **Pt-4** (4.89 Å). Substituting the  $\text{C}\equiv\text{CC}_6\text{H}_5$  ligand in **Pt-2** with the  $\text{N}\equiv\text{CCH}_3$  ligand in **Pt-3** resulted in a reduction of the Pt–Pt distance from 6.43 to 3.28 Å. Additionally, the Pt–Pt contact in **Pt-1** and **Pt-3** is nearly perpendicular to the planar Pt complex, revealing a face-to-face close packing form. In contrast, this angle deviates significantly from  $90^\circ$  in **Pt-2** and **Pt-4**. These structural findings suggest that in **Pt-2** and **Pt-4**, which feature two strong  $\sigma$ -donor  $\text{C}^-$  ligands, close Pt–Pt contacts are unfavourable. In contrast, **Pt-1** and **Pt-3** complexes, which contain only one  $\text{C}^-$  ligand, exhibit close Pt–Pt contacts.





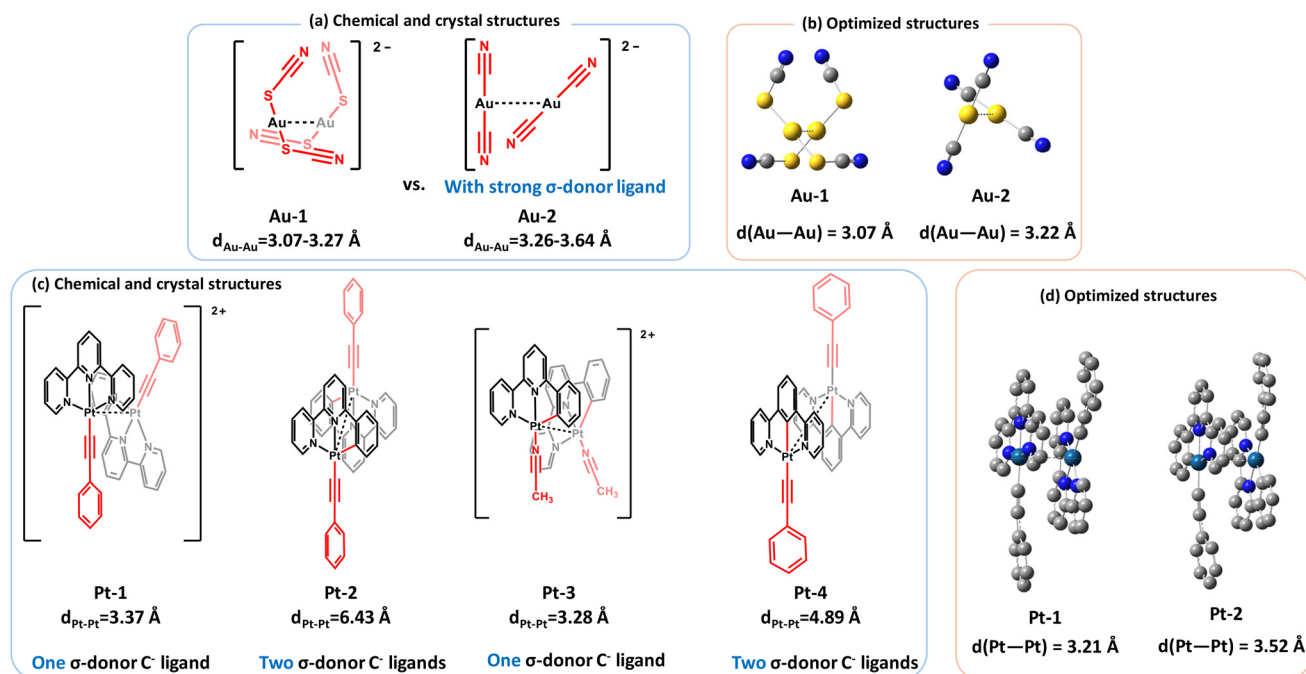
**Fig. 6** Calculated EDA-NOCV deformation orbital energies correspond to the  $\text{Rh} \rightarrow \text{L} \pi$  interaction in the model system of  $\text{Rh}(\text{terpy})\text{L}$ . All data points are tabulated in Table S1.†

The M–M distance trend observed in X-ray crystal structures was validated through computational optimization of the dimeric species. As shown in Fig. 7b and d, optimized structures replicate key experimental trends from Fig. 7a and c:  $[\text{Au-1}]_2$  exhibits shorter Au–Au distances than  $[\text{Au-2}]_2$ , and  $[\text{Pt-1}]_2$  shows reduced Pt–Pt distances compared to  $[\text{Pt-2}]_2$ . This con-

sistency between computational and crystallographic data confirms that the observed variations originate from intrinsic electronic properties rather than crystal packing effects.

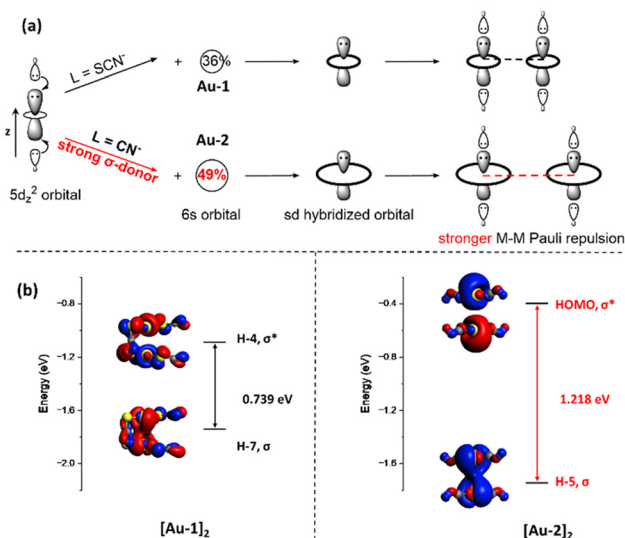
For the  $\sigma$ -bonding interaction between a metal and a ligand, the ligand donates electron density to the metal's empty d or s orbital through  $\sigma$ -type M–L interaction, which influences the composition, energy, and electron density distribution of metal orbitals. The role of strong  $\sigma$ -donor ligands in strengthening the M–M Pauli repulsion for  $d^8$ – $d^8$  and  $d^{10}$ – $d^{10}$  closed-shell systems is shown in Scheme S2† and Fig. 8a, respectively.

For the  $d^{10}$  metal complexes such as  $\text{Au}(\text{I})$ ,  $\text{Ag}(\text{I})$ , or  $\text{Cu}(\text{I})$ ,  $(n + 1)s$ -nd orbital hybridization plays a critical role in the formation of M–L  $\sigma$ -type coordination bonds ( $n = 3$  for  $\text{Cu}(\text{I})$ , 4 for  $\text{Ag}(\text{I})$  and 5 for  $\text{Au}(\text{I})$ ).<sup>36,37</sup> As shown in Fig. 8a, part of the empty  $(n + 1)s$  orbitals are used to accept electron density from  $\sigma$ -donor ligands. Defining the M–L bond direction as the z-axis, the  $(n + 1)s$ -nd orbital hybridization reduces the metal orbital's lobe size along the z-axis while increasing its size in the xy-plane.<sup>40</sup> When two metal atoms approach each other in the xy-plane, their sd-hybridized orbitals overlap, resulting in the formation of a bonding  $\sigma(\text{M–M})$  orbital and an antibonding  $\sigma(\text{M–M})^*$  orbital.<sup>36,37</sup> Coordination with a stronger  $\sigma$ -donor ligand enhances the extent of  $(n + 1)s$  orbital hybridization. This, in turn, increases the metal's electron density in the xy plane, resulting in greater orbital splitting and stronger M–M Pauli repulsion.



**Fig. 7** Chemical structures of (a)  $\text{Au}(\text{I})$  and (c)  $\text{Pt}(\text{II})$  metal complexes, together with the experimentally resolved M–M distance in their crystal structures. Note: for **Au-1** and **Au-2**, the Au–Au distance is different when changing the counteranion.<sup>5</sup> We compared **Au-1** and **Au-2** with the same counteranion here:  ${}^n\text{Bu}_4\text{N}[\text{Au-1}]$ :  $d(\text{Au–Au}) = 3.0700 \text{ \AA}$ ,  $\text{K}[\text{Au-1}]$ :  $d(\text{Au–Au}) = 3.0064/3.0430 \text{ \AA}$ .  $\text{K}[\text{Au-2}]$ :  $d(\text{Au–Au}) = 3.64 \text{ \AA}$ ,  ${}^n\text{Bu}_4\text{N}[\text{Au-2}]$ : no close Au–Au contact. Optimized structures of (b)  $\text{Au}(\text{I})$  and (d)  $\text{Pt}(\text{II})$  metal complexes. Calculations were performed at the PBE0-GD3BJ/SDD(Au, Pt)/6-31G\* (other elements) level of theory using the PCM solvent model for ACN in Gaussian 16.

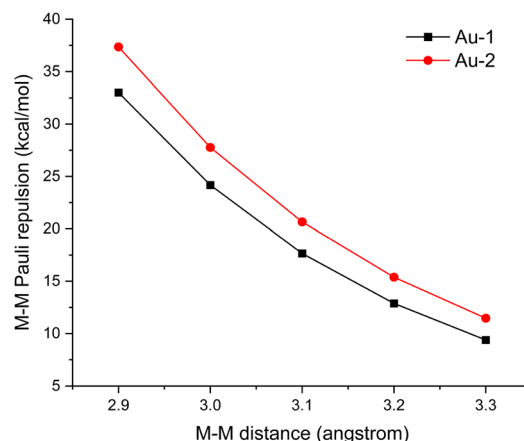




**Fig. 8** (a) Illustration scheme showing the role of the  $\sigma$ -donor ligand in strengthening the M-M Pauli repulsion for  $d^{10}$ - $d^{10}$  closed-shell metal complexes. Complexes **Au-1** and **Au-2** were used as the example. (b) Calculated MO diagram for the **Au-1** dimer (left) and the **Au-2** dimer (right) to highlight the  $d\sigma(M-M)$  bonding and  $d\sigma^*(M-M)$  antibonding orbitals. Isovalue = 0.03.

To validate the schematic diagram plotted in Fig. 8a, we conducted MO calculations on **Au-1** and **Au-2** dimers (Fig. 8b). Because of the stronger  $\sigma$ -donor ligand  $C\equiv N^-$  in **Au-2**, a greater contribution from the Au-6s orbital was observed in its 6s-5d hybridized orbital (49% in **Au-2** vs. 36% in **Au-1**). The increased Au-6s orbital contribution enlarges the metal orbital's lobe size in the  $xy$  plane, thereby increasing the orbital splitting energy in the **Au-2** dimer compared to the **Au-1** dimer. As shown in Fig. 8b, the orbital splitting energy between the M-M  $\sigma(6s-5d)$  bonding orbital and the M-M  $\sigma^*(6s-5d)$  antibonding orbital is calculated to be 0.739 eV in **[Au-1]<sub>2</sub>** and 1.218 eV in **[Au-2]<sub>2</sub>**. This larger orbital splitting energy in **[Au-2]<sub>2</sub>** indicates stronger Au-Au Pauli repulsion in **[Au-2]<sub>2</sub>** compared to that in **[Au-1]<sub>2</sub>**.

We next performed EDA calculations on the **Au-1** and **Au-2** dimers (Fig. 9). In the calculations, M-M distances in these Au (I) dimers were set from 3.3 Å to 2.9 Å to compare the M-M Pauli repulsion at equal M-M distances. Au-Au Pauli repulsion increased with the decrease of the Au-Au distance for both **Au-1** and **Au-2** dimers, and the M-M Pauli repulsion in the **Au-1** dimer is always stronger than that in the **Au-2** dimer. For example, at an Au-Au distance of 3.0 Å, the M-M Pauli repulsion is 27.7 kcal mol<sup>-1</sup> in the **Au-1** dimer, while it is 24.2 kcal mol<sup>-1</sup> in the **Au-2** dimer. Relaxation density (defined as the electron density difference between the metal atom after ligand coordination and its pre-coordination state) of Mulliken population on the metal's s-orbitals ( $\Delta s$ ) was examined to further characterize the ligand's  $\sigma$ -donating ability for complexes **Au-1** and **Au-2**. The calculated trend of M-M Pauli repulsion in **Au-1** and **Au-2** is consistent with the calculated relaxation density. Specifically, **Au-2** exhibits a larger 6s-orbital relaxation density ( $\Delta s = 0.9073$ ) with



**Fig. 9** Calculated M-M Pauli repulsion of **Au-1** and **Au-2** dimers by changing the M-M distance. For EDA calculations at non-equilibrium M-M distances, fragment geometries were frozen while varying metal-metal separation.

the  $CN^-$  ligand than **Au-1** exhibits with the  $SCN^-$  ligand ( $\Delta s = 0.8455$ ), indicating that the  $CN^-$  ligand donates more electron density to the Au-6s orbital. To summarize, coordination to strong  $\sigma$ -donor ligands increases the  $d\sigma(M-M)$ - $d\sigma^*(M-M)$  orbital splitting and strengthens the M-M Pauli repulsion, preventing the formation of close M-M contact in the **Au-2** complex compared to **Au-1**.

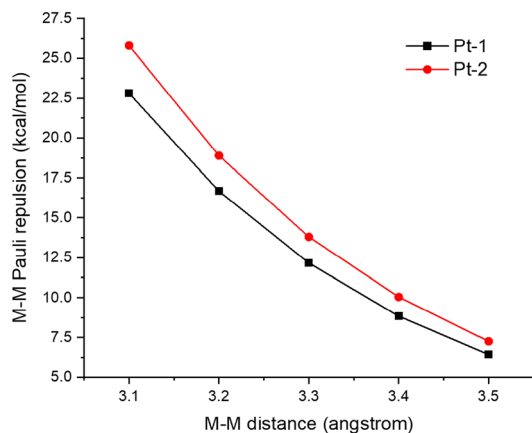
Let us now consider the role of  $\sigma$ -bonding interaction in  $d^8$ - $d^8$  systems. As previously noted, in square-planar  $d^8$  metal complexes, the  $nd_{x^2-y^2}$  orbital is unoccupied and strongly antibonding regarding the M-L coordination.<sup>37</sup> This orbital does not contribute to M-M Pauli repulsion. In contrast, the  $nd_{z^2}$  weakly antibonding orbital is fully occupied and significantly contributes to M-M Pauli repulsion through the M-M  $\sigma$ -type interaction (Scheme 2 and S2†).<sup>33</sup> Importantly, the  $nd_{z^2}$  orbital can accept the ligand's electron density by the hybridization with higher empty metal- $(n+1)s$  orbitals to weaken the M-L antibonding interaction ( $n = 4$  for Rh(I) and Pd(II),  $n = 5$  for Ir(I) and Pt(II)). The relative component of the hybridized  $(n+1)s$  orbitals is closely related to the  $\sigma$ -donating ability of ligands, similar to  $d^{10}$ - $d^{10}$  systems. We compared the Pt-6s orbital's population in **Pt-1** and **Pt-2**, where the **Pt-1** complex has one strong  $\sigma$ -donor  $C^-$  ligand and the **Pt-2** complex has two strong  $\sigma$ -donor  $C^-$  ligands.

As shown in Scheme S2,† more Pt-6s orbital's components show up in the Pt-5d<sub>z²</sub> orbital for **Pt-2** (23%) than in **Pt-1** (20%). In the calculated Mulliken population, **Pt-2** also has more Pt-6s orbital's relaxation density than **Pt-1** after ligand coordination ( $\Delta s = 0.1151$  in **Pt-2** vs. 0.0001 in **Pt-1**). The increased electron density in the Pt-6s orbital of the **Pt-2** complex compared to **Pt-1** is expected to amplify the Pt-Pt Pauli repulsion, which consequently elongates the Pt-Pt bond distance.

Next, we calculated and compared the M-M Pauli repulsion in **Pt-1** and **Pt-2** dimers (Fig. 10). M-M distances in the Pt(II) dimers were set from 3.5 Å to 3.1 Å for the comparison of M-M Pauli repulsion at equal M-M distances. As revealed in Fig. 10, the M-M Pauli repulsion in the **Pt-2** dimer is always stronger







**Fig. 10** Calculated M-M Pauli repulsion of Pt-1 and Pt-2 dimers by changing the M-M distance. For EDA calculations at non-equilibrium M-M distances, fragment geometries were frozen while varying metal-metal separation.

than that in the Pt-1 dimer, which is consistent with the electronic configuration calculations.

## Conclusions

In summary, we studied the relationship between ligands (broadly categorized as  $\pi$ -acceptors and/or  $\sigma$ -donors) and M-M Pauli repulsion in a series of  $d^8$  and  $d^{10}$  closed-shell metal complexes. With  $\pi$ -acceptor ligands such as nitriles and isocyanides, M-M Pauli repulsion is weakened by the transfer of electron density from the closed-shell metal atom to the  $\pi$ -acceptor ligand's empty orbital. This results in shorter M-M distances in their respective X-ray crystal structures. Conversely,  $\sigma$ -donor ligands, such as cyanide ( $\text{CN}^-$ ) and alkynides ( $\text{R-C}\equiv\text{C}^-$ ), strengthen M-M Pauli repulsion by transferring electron density from the ligand's lone pair of orbitals to the metal's empty ( $n+1$ ) s orbital, leading to longer M-M distances. This work provides valuable insights into the design of  $d^8$  and  $d^{10}$  metal complexes with controlled M-M distances by leveraging M-L coordination interactions. These findings also deepen the understanding of metallophilicity and offer a framework for tailoring M-M interactions in functional materials, with potential applications in catalysis, supramolecular assembly, and optoelectronics.

## Author contributions

Q. W. designed and supervised the project. Y. Z. carried out the calculations. Both authors contributed to the editing of the manuscript.

## Conflicts of interest

There are no conflicts to declare.

## Data availability

The data supporting this article have been uploaded as part of the ESI.†

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