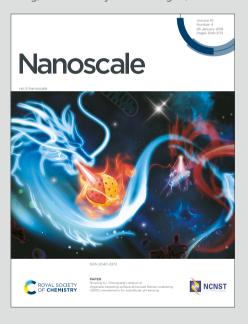




Nanoscale

Accepted Manuscript

This article can be cited before page numbers have been issued, to do this please use: Y. Cai, Z. Gu, S. Yang, C. B. Murray and C. Kagan, *Nanoscale*, 2025, DOI: 10.1039/D5NR03413B.



This is an Accepted Manuscript, which has been through the Royal Society of Chemistry peer review process and has been accepted for publication.

Accepted Manuscripts are published online shortly after acceptance, before technical editing, formatting and proof reading. Using this free service, authors can make their results available to the community, in citable form, before we publish the edited article. We will replace this Accepted Manuscript with the edited and formatted Advance Article as soon as it is available.

You can find more information about Accepted Manuscripts in the <u>Information for Authors</u>.

Please note that technical editing may introduce minor changes to the text and/or graphics, which may alter content. The journal's standard <u>Terms & Conditions</u> and the <u>Ethical guidelines</u> still apply. In no event shall the Royal Society of Chemistry be held responsible for any errors or omissions in this Accepted Manuscript or any consequences arising from the use of any information it contains.



View Article Online DOI: 10.1039/D5NR03413B

COMMUNICATION

Shape Effects on the 2D Self-Assembly of Lithographically Fabricated Nanoparticles

Yi-Yu Cai, ‡ab Zeyu Gu, ‡a Shengsong Yang, c Christopher B. Murray, cd and Cherie R. Kagan acd*

Received 00th January 20xx, Accepted 00th January 20xx

DOI: 10.1039/x0xx00000x

We fabricate and polymer ligate square, pentagonal, and hexagonal nanoplates and investigate the effect of nanoparticle (NP) shape on the self-organization of their dispersions at the liquid-air interface into two-dimensional assemblies. Hexagonal NPs assemble into interlocking hexagonal lattices, exhibiting the highest translational order with a characteristic correlation length ($\xi 0.5$) extending over nine NP building blocks. Square NP assemblies are disordered, characterized by $\xi 0.5$ averaging around three NP units, consistent with the propensity of NPs to slide and form coexisting $\Lambda 1$ and $\Lambda 0$ lattices. Pentagonal NP assemblies show six-fold coordination, but do not fully tessellate the plane, yielding anisotropic assemblies with intermediate order and $\xi 0.5$ of about seven NP units. In binary NP assemblies, mixed pentagonal and hexagonal NP assemblies with shared six-fold coordination show greater ordering in comparison to assemblies of square and hexagonal NPs.

Introduction

The shape of nanoparticles (NPs), analogous to atomic valency, directs the order and motif of self-assembled structures. 1-3 Geometrically, NP shape significantly impacts thermodynamically favorable arrangement for dense packing.4 Polygons with fewer than seven vertices, interacting primarily through steric hindrance, show a preference for edge-to-edge alignment.5 due to shape-induced directional entropic forces for generating more free volume.⁶⁻⁸ Studies on the formation of two-dimensional structures or ordered lattices of polygons have employed both computational simulations and experimental methods at various length scales. 9-12 Regular hexagonal-shaped particles typically form a close-packed hexagonal lattice across nano-,13, 14 micro-,15 and millimeter scales,16 mirroring their intrinsic six-fold symmetry. The interlocking geometry of their superlattice results in a rigid structure that enhances translational order. ^12 Square assemblies display more intrinsic disorder because they can slide along columns and rows, allowing for greater flexibility. The degree of corner rounding affects the packing density, as sharp corners result in a perfect square lattice. 9 while rounded corners lead to Λ_1 and Λ_0 lattice structures. ^17 Experimental research has demonstrated the assembly of square particles into two-dimensional planes at both the nano ^14 and micrometer scales, ^18-20 showcasing diverse lattice structures. Unlike squares and hexagons, the regular pentagon cannot fully tessellate a plane.

The study of pentagonal NP assemblies spans multiple scales.¹⁰ At the millimeter scale, pentagonal styrene particles subjected to hard-core interactions and weak hydrodynamic interactions on air tables exhibit the ice-ray packing,²¹ the densest known local configuration characterized by alternating stripes of particles in two distinct orientations.²² Lithographically-defined microscale epoxy-based polymer particles assembled in aqueous environments form hexagonal rotator crystal structures, dominated by "hard" short-range repulsive in-plane interactions. 19, 23, 24 Translational and rotational jamming prevent the formation of the ice-ray packing. At the nanometer scale, shape customization has largely originated from advancements in chemical synthesis techniques, where achieving high uniformity in size and shape is a critical factor for the self-assembly of superlattices.^{25, 26} Colloidal penta-twinned elongated nanoprisms with pentagonal cross-sections self-assemble from dispersions into multilayered structures that adopt configurations ranging from the dense ice-ray packing to the less dense Dürer packing.²² Chemical routes to synthesize nanoscale pentagonal plates are rare,²⁷ thus, to our knowledge, there are few studies on the twodimensional assembly of nanoscale pentagonal particles. Studies on the assembly of pentagonal-shaped molecules have also been conducted. The interactions that govern their assembly significantly diverge from those in particle assemblies, thus falling outside the scope of this paper.²⁸

In this study, we use nanoimprint lithography and thermal evaporation to fabricate Au nanoplates with precise control over size and shape, including square, pentagonal, and

^{a.} Department of Electrical and Systems Engineering, University of Pennsylvania, Philadelphia, PA 19104, USA

b Macao Institute of Materials Science and Engineering (MIMSE), Faculty of Innovation Engineering, Macau University of Science and Technology, Taipa, Macao, 999078 China

^c Department of Chemistry, University of Pennsylvania, Philadelphia, PA 19104,

d. Department of Materials Science and Engineering, University of Pennsylvania, Philadelphia, PA 19104, USA, E-mail: kagan@seas.upenn.edu

[‡] These authors contributed equally.

Supplementary Information available: [details of any supplementary information available should be included here]. See DOI: 10.1039/x0xx00000x

Journal Name

hexagonal geometries with identical side lengths, as a model system for studying two-dimensional self-assembly. These nanoplates are assembled into 2D extended monolayers via the liquid-air interface assembly technique.²⁹ Our quantitative analysis reveals that the translational order of the Au nanoplates follows the trend: hexagonal > pentagonal > square shape, attributed to the shape symmetry, interlocking capability based on coordination number, and tiling ability. Furthermore, binary mixtures of pentagonal and hexagonal NPs, which exhibit a common six-fold coordination, yield greater translational

order than their combinations with squares, demonstrating the

Results and discussion

This article is licensed under a Creative Commons Attribution 3.0 Unported Licence.

Open Access Article. Published on 19 September 2025. Downloaded on 9/21/2025 11:32:02 PM.

influence of shape on self-assembly.

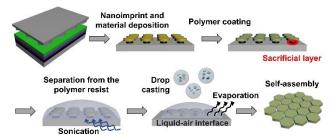


Figure. 1 Schematic of the preparation and self-assembly of polygonal NP dispersions. Square, pentagonal, and hexagonal Au NPs are fabricated via nanoimprint lithography and thermal evaporation of thin films. After liftoff of the resist, NPs remain bound to the wafer. These NPs are subsequently coated with a thiolated polystyrene polymer layer, facilitating their harvesting through sonication and stable dispersion in liftoff solvents. The self-assembly process is carried out at the liquid-air interface, with NPs dispersed in a hydrophobic phase (toluene, tetrahydrofuran, chloroform) atop a hydrophilic ethylene glycol layer. As the ethylene glycol evaporates, assembled NP layers deposit on the underlying wafers.

The square, pentagonal, and hexagonal Au NPs are produced via a combination of nanoimprint lithography and thermal evaporation techniques (Figure 1), as detailed in our prior work.30 Briefly, e-beam lithography is employed to create imprint templates with the desired polygonal shapes. Here, templates are designed to produce NPs with uniform side lengths of approximately 220 nm. Wafers are coated with resist trilayers, and the shapes are transferred into the top thermal Nanonex resist layer and middle polydimethylglutarimide (PMGI) undercut layer through imprinting and developing. Subsequently, a layer of Au or other materials is physically evaporated onto the patterned resist layers, and upon lift-off, yields the NPs anchored by the bottom sacrificial Durimide layer. To ensure the dispersion of the NPs in solvents, the Au NPs are coated with a thiolated-polystyrene polymer (numberaverage molecular weight $M_n = 50k$) in a tetrahydrofuran (THF) solution. The coated NPs are then harvested from the wafer by immersion and sonication in N-methyl-2-pyrrolidone (NMP). Through a series of centrifugation and redispersion steps in organic solvents, we remove excess polymer ligands and concentrate approximately 108 NPs into less than 5 µL of a

mixed organic solvent of either chloroform and toluene, or THE and toluene, with a volume ratio opings, or the spective of the specific of the

The Au NP dispersions are then dropped onto an ethylene glycol subphase, forming distinct hydrophobic organic phases. As the organic solvent evaporates, NP films emerge on the surface of the ethylene glycol. Wafers are submerged in the ethylene glycol subphase, and the ethylene glycol is evaporated through vacuum drying, depositing the NP films on the wafers. NP assemblies are subsequently imaged in the SEM, and the images are analyzed to quantify their structure, including the shapes, positions, orientations, and long-range order of the NPs. The analysis techniques applied to all images include Fourier transforms (FFTs), 2D nearest-neighbor positional distributions, 2D autocorrelation functions, and full pair position probability distributions, allowing for comprehensive characterization of the assembly structures. 12 Detailed orientational maps and lattice-type classifications are provided in Figure S2.

Figure 2 shows the distinct structures for square, pentagonal, and hexagonal Au NP assemblies. Hexagonal Au NP assemblies form highly uniform (Figure 2C(I)), close-packed hexagonal lattice structures as depicted in Figure 2C(I) inset.14 Uniform particle alignment is evidenced by consistent coloration across the image in Figure S2C, where colors represent particle orientation angles, as well as by the sharp peaks in FFTs (Figure 2C(II)). 2D nearest-neighbor positional distributions (Figure 2C(III)) quantify the likelihood of locating a neighboring NP at a specific distance. 12 They reveal a nearly isotropic spread in distance with a major-to-minor axis FWHM ratio of 1.09 ± 0.07, the closest to unity among all shapes studied. Square Au NP assemblies exhibit a coexistence of Λ_0 and Λ_1 lattices (Figure 2A, S2A). In the Λ_0 -lattice, square NPs align in rows with their faces offset by half a side length, shown in blue in the inset of Figure 2A(I). In contrast, the Λ_1 -lattice, colored pink in the same inset, is characterized by squares with rounded corners that touch, forming a rhombic symmetry. 18 The lattice-type detection and orientational analysis used to distinguish Λ_0 from Λ_1 are illustrated in Figure S2A. The 2D FFTs of the images for square Au NP assemblies show discernible peaks (Figure 2A(II)), though less pronounced compared to those for other NP shapes, reflecting the disorder from the coexistence of multiple lattices (Figure 2A(II)). The 2D nearest-neighbor positional distributions (Figure 2A(III)) of square NPs show the highest anisotropy with a major-to-minor axis FWHM ratio of 2.39 ± 0.29, indicating strong directional differences in the translational correlations.

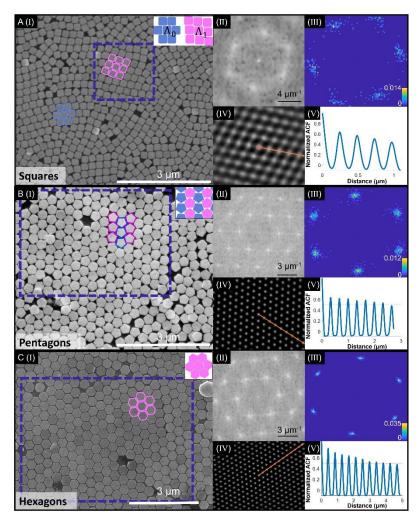


Figure. 2 Structural analysis of (A) square, (B) pentagonal, and (C) hexagonal Au NP assemblies. Each set includes: (I) SEM images, regions with colored outlines and insets showing representative lattices, (II) FFTs of SEM images, (III) 2D probability distributions of nearest-neighbor positional patterns, (IV) 2D autocorrelation functions, and (V) linecut profiles from the 2D autocorrelation in the marked regions.

Pentagonal Au NP assemblies exhibit ice-ray packing with antiparallel arrangements indicated by alternating purple and blue colors in Figure 2B and Figure S2B.²² Other areas form a hexagonal rotator crystal phase,²³ maintaining positional order but lacking consistent orientational order. FFTs of pentagonal Au NP assemblies reveal a hexagonal pattern (Figure 2B(II)), consistent with distinct peaks in 2D nearest-neighbor positional distributions (Figure 2B(III)). Quantitative analysis of 2D nearest-neighbor positional distributions shows intermediate level of anisotropy (major-to-minor axis FWHM ratio of 1.80 \pm 0.12) between hexagonal and square assemblies. However, the expected rectangular symmetry of antiparallelpatterned ice-ray packing is absent in the FFTs, likely due to local variations and imperfections in the antiparallel structure observed in the image.34

To quantitatively compare the long-range positional order across different NP shapes, we utilized 2D autocorrelation functions derived from SEM images (Figure 2A-C ($\rm IV$)). These functions measure the self-similarity of the image as the displacement vector changes. The $\xi_{0.5}$ value of translational

order is defined at points where the oscillation in the line profiles drops to less than 50% of their initial value (Figure 2A-C (V)). To determine the $\xi_{0.5}$ value for each polygonal NP assembly, we start with a well-ordered region in each SEM image to measure the $\xi_{0.5}.$ As the area expands, the $\xi_{0.5}$ value typically reaches a maximum before decreasing due to the inclusion of less ordered regions. This approach enables us to identify the maximum analysis area for each shape, capturing the best-case scenario for order without being constrained by a fixed area across different shapes, as exemplified by the selected areas shown in Figure 2A-C (m V). Using this method, we observe that hexagonal NPs exhibit a maximal $\xi_{0.5}$ value of 9.2 \pm 0.8 repeating units (Table S1), in contrast to 6.8 ± 0.4 for pentagonal Au NPs. Square Au NPs show a lower maximal $\xi_{0.5}$ value of 3.2 ± 0.4 , likely due to the presence of multiple lattices and their propensity to slide.

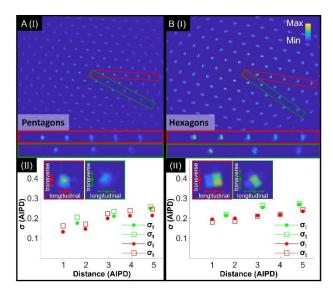


Figure. 3 Analysis of the 2D pair position distribution function for (A) pentagonal and (B) hexagonal Au NP assemblies. Each set includes: (I) 2D pair position probability distributions with corresponding linecuts in the $[11\overline{2}0]$ (red) and $[1\overline{1}00]$ (green) directions. (II) Insets: definitions of the width of the position probability distribution for the first peak away from the center. Peak widths, in units of average interparticle distance (AIPD), for the $[11\overline{2}0]$ (red) and $[1\overline{1}00]$ (green) directions in both longitudinal and transverse orientations are plotted against interparticle distance.

The translational order in pentagonal and hexagonal Au NP assemblies is further analyzed through the full pair position probability distribution.35 which assesses the likelihood of locating an NP at a specific location, given that another NP is

positioned at the origin (Figure 3). A broader peak in this distribution indicates increased disordeP,ार्निविधिताराष्ट्रिकारिकार्याका accumulation of lattice disorder in both pentagonal and hexagonal Au NP assemblies across various lattice directions, and in longitudinal (σ_l) and transverse (σ_t) orientations. However, the key distinction is that well-ordered hexagonal Au NP assemblies exhibit isotropic translational order, with σ_l roughly equal to σ_t . In contrast, pentagonal Au NP assemblies consistently show anisotropic translational order, with $\sigma_{\!\scriptscriptstyle I}$ smaller than σ_t , revealing a dominance of shearing over compressive or tensile disorder. The anisotropy in pentagonal Au NP assemblies is attributed to their inability to tessellate a plane without creating voids, thereby facilitating the sliding of rows or columns. While this allows for more flexibility than for hexagonal NP assembly, pentagonal NPs still achieve a higher degree of order than square NPs as their six-fold coordination partially restricts NP translation. While lateral capillary forces from contact-line distortions may contribute to particle interactions at short range,36 the systematic shape-dependent assembly trends observed are consistent with geometric constraints and entropic forces dominating the assembly process for these sub-micron particles.4,29

We mix combinations of the polygonal Au NPs to form binary mixtures, specifically pentagonal-hexagonal and squarehexagonal combinations (Figure 4). Both binary mixtures exhibit a degree of phase separation, where particles of the same shape tend to cluster together, as visible in the SEM images (Figure 4A,B(I)). The pentagonal-hexagonal mixture has sharper and more distinct peaks in the FFTs (Figure 4A,B(II)), indicating better long-range order. The 2D nearest-neighbor position distributions reveal a clearer hexagonal pattern for the pentagonal-hexagonal mixture (Figure 4B(III)). In contrast, the

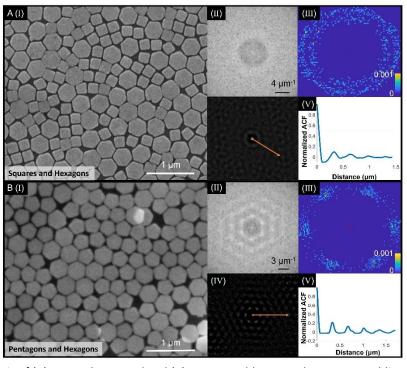


Figure. 4 Structural analysis of (A) square-hexagonal and (B) pentagonal-hexagonal Au NP assemblies. Each set includes: (I) SEM images and their (II) FFTs, (III) 2D probability distributions of nearest-neighbor positional patterns, (IV) 2D autocorrelation functions, and (V) linecut profiles from the 2D autocorrelation.

Journal Name

COMMUNICATION

square-hexagon mixture shows a dual-ring structure without distinct peaks (Figure 4A(III)), reflecting different nearestneighbor distances for different shapes and a less uniform local structure. The autocorrelation analysis (Figure 4A,B(IV)) and the line profiles for both mixtures (Figure 4A,B(V)) show peak heights below 0.5, indicating limited coherence. The pentagonal-hexagonal mixture displays more pronounced and regularly spaced peaks compared to the square-hexagon mixture, indicating higher translational order and longer-range correlations. The split peaks visible in the pentagonalhexagonal autocorrelation likely result from the coexistence of multiple characteristic distances between particles of the same and different shapes in the binary system. This improved ordering in the pentagonal-hexagonal system benefits from their geometric compatibility, both having six nearest neighbors. The lower order in the square-hexagonal mixture¹¹ is likely due to the incompatibility between the different lattice structures as squares can form both $\Lambda_{0^{\text{-}}}$ and $\Lambda_{1}\text{-lattices}$ and hexagons pack in a hexagonal lattice.

Conclusions

This article is licensed under a Creative Commons Attribution 3.0 Unported Licence

Den Access Article. Published on 19 September 2025. Downloaded on 9/21/2025 11:32:02 PM.

In this study, we explore how polygonal shape influences the 2D self-assembly of lithographically fabricated nanoplates, with a focus on square, pentagonal, and hexagonal NPs. Analysis of translational order reveals a clear hierarchy: hexagonal NPs form the most ordered assemblies due to their interlocking hexagonal lattices, square NPs exhibit the most disorder with coexisting Λ_1 and Λ_0 lattices, while pentagonal NPs show intermediate, anisotropic order. This behavior stems from the geometries of each shape and their ability to tile a plane. Similar shape-driven hierarchies of translational order have been directly observed for colloidal nanocrystals at the tens-ofnanometers scale, where van der Waals and ligand-mediated interactions dominate. In this regime, ligand length/coverage sets the interparticle gap and effectively rounds particle corners, thereby modulating directional entropic forces and, in turn, translational coherence.4, 12 Beyond the assembly of a single shape, we show that binary mixtures of pentagonal and hexagonal NPs achieve higher translational order through shared coordination numbers, in contrast to the comparatively disordered square-hexagon mixtures. To our knowledge, this provides nanoscale experimental evidence that coordination compatibility can mitigate demixing and enhance order in a binary polygonal system (pentagon-hexagon); by contrast, prior simulations of other symmetric polygon mixtures (e.g., hexagon-square/triangle) predict demixing unless shape- or edge-specific patches are introduced.¹¹

Our approach to NP fabrication and assembly is also applied to create 2D extended assemblies of binary mixtures of hexagonal Au and Ag NPs, as well as hexagonal Ge NPs, suggesting its broad applicability (Figure S3). This research opens avenues for further exploration into the self-assembly of NPs with anisotropic, complex shapes (such as convex and concave polygons), varied sizes, and compositions with entropic and enthalpic patchy interactions. 11, 37-39 By integrating NPs of various types, we plan to design functional materials that leverage the combined properties of the constituent particles for applications in optics and opto-electronics. For the sample, close-packed monolayers of hexagonal nanoplates provide uniform nearest-neighbor gaps and edge sites that yield reproducible near-field hot spots and enable quantitative SERS.40, 41 Our previous work showed that symmetry and interparticle gaps in polygonal NP metamolecules tune the Fano lineshapes and polarization response³⁰ and correlated singledomain spectroscopy on binary nanocrystal superlattices further demonstrates that lattice symmetry and composition systematically control near-field coupling and resonance energy.⁴² For electronic/optoelectronic devices, edge-to-edge assembly of cubic-shaped quantum dots with short interparticle spacings enhance coupling and charge transport in quantumdot assemblies, compared to those composed of spherical quantum dots. 43, 44

Conflicts of interest

There are no conflicts to declare.

Data availability

The data supporting this article have been included as part of the Supplementary Information. The data that support the findings of this study are available from the corresponding author upon reasonable request.

Acknowledgements

The authors gratefully thank the Office of Naval Research Multidisciplinary University Research Initiative Award ONR N00014-18-1-2497 for primary support of the work, including NP nanofabrication, assembly, and characterization. Y.C. acknowledges the Science and Technology Development Fund (FDCT), Macao SAR (Grant No. 0024/2024/RIB1, 0136/2024/RIA2) and the National Natural Science Foundation of China (Grant No. 12404421) for support during the writing of this manuscript. The nanofabrication was conducted partially at the Singh Center for Nanotechnology, which was supported by the NSF National Nanotechnology Coordinated Infrastructure Program under grant NNCI-2025608. This research also utilized the electron beam lithography infrastructure at the Center for Functional Nanomaterials (CFN), a U.S. Department of Energy Office of Science User Facility located at Brookhaven National Laboratory, under Contract No. DE-SC0012704.

References

- J. Gong, R. S. Newman, M. Engel, M. Zhao, F. Bian, S. C. Glotzer and Z. Tang, Nat. Commun., 2017, 8, 14038.
- S. C. Glotzer and M. J. Solomon, Nat. Mater., 2007, 6, 557-2.
- G. van Anders, N. K. Ahmed, R. Smith, M. Engel and S. C. Glotzer, ACS Nano, 2014, 8, 931-940.
- M. A. Boles, M. Engel and D. V. Talapin, Chem. Rev., 2016, **116**, 11220-11289.

Vanoscale Accepted Manuscrii

COMMUNICATION Journal Name

- 5. J. A. Anderson, J. Antonaglia, J. A. Millan, M. Engel and S. C. Glotzer, Phys. Rev. X, 2017, 7, 021001.
- 6. G. van Anders, D. Klotsa, N. K. Ahmed, M. Engel and S. C. Glotzer, Proc. Natl. Acad. Sci. U. S. A., 2014, 111, E4812-E4821.
- 7. E. S. Harper, G. van Anders and S. C. Glotzer, Proc. Natl. Acad. Sci. U. S. A., 2019, 116, 16703-16710.
- 8. T. Vo and S. C. Glotzer, Proc. Natl. Acad. Sci. U. S. A., 2022, 119, e2116414119.
- 9. C. Avendaño and F. A. Escobedo, Soft Matter, 2012, 8,
- T. Schilling, S. Pronk, B. Mulder and D. Frenkel, Phys. Rev. 10. E, 2005, 71, 036138.
- 11. J. A. Millan, D. Ortiz, G. van Anders and S. C. Glotzer, ACS Nano, 2014, 8, 2918-2928.
- 12. J. C. Ondry, L. B. Frechette, P. L. Geissler and A. P. Alivisatos, Nano Lett., 2022, 22, 389-395.
- 13. F. Qin, X. Cui, Q. Ruan, Y. Lai, J. Wang, H. Ma and H.-Q. Lin, Nanoscale, 2016, 8, 17645-17657.
- X. Ye, J. E. Collins, Y. Kang, J. Chen, D. T. N. Chen, A. G. 14. Yodh and C. B. Murray, Proc. Natl. Acad. Sci. U. S. A., 2010, 107, 22430-22435.
- 15. T. D. Clark, J. Tien, D. C. Duffy, K. E. Paul and G. M. Whitesides, J. Am. Chem. Soc., 2001, 123, 7677-7682.
- 16. N. Bowden, I. S. Choi, B. A. Grzybowski and G. M. Whitesides, J. Am. Chem. Soc., 1999, 121, 5373-5391.
- 17. Y. Jiao, F. H. Stillinger and S. Torquato, Phys. Rev. Lett., 2008, **100**, 245504.
- 18. J.-M. Meijer, V. Meester, F. Hagemans, H. N. W. Lekkerkerker, A. P. Philipse and A. V. Petukhov, Langmuir, 2019, 35, 4946-4955.
- 19. K. Zhao, R. Bruinsma and T. G. Mason, Proc. Natl. Acad. Sci. U. S. A., 2011, 108, 2684-2687.
- 20. L. Rossi, V. Soni, D. J. Ashton, D. J. Pine, A. P. Philipse, P. M. Chaikin, M. Dijkstra, S. Sacanna and W. T. M. Irvine, Proc. Natl. Acad. Sci. U. S. A., 2015, 112, 5286-5290.
- 21. Y. L. Duparcmeur, A. Gervois and J. P. Troadec, J. Phys.: Condens. Matter, 1995, 7, 3421.
- 22. J. Marcone, W. Chaâbani, C. Goldmann, M. Impéror-Clerc, D. Constantin and C. Hamon, Nano Lett., 2023, 23, 1337-1342.
- 23. K. Zhao and T. G. Mason, Phys. Rev. Lett., 2009, 103, 208302.
- 24. P.-Y. Wang and T. G. Mason, *Nature*, 2018, **561**, 94-99.
- 25. C. B. Murray, C. R. Kagan and M. G. Bawendi, Annu. Rev. Mater. Sci., 2000, 30, 545-610.
- C. B. Murray, C. R. Kagan and M. G. Bawendi, Science, 26. 1995, 270, 1335-1338.
- 27. J. J. Velázquez-Salazar, L. Bazán-Díaz, Q. Zhang, R. Mendoza-Cruz, L. Montaño-Priede, G. Guisbiers, N. Large, S. Link and M. José-Yacamán, ACS Nano, 2019, 13, 10113-10128.
- 28. C. Yuan, N. Xue, X. Zhang, Y. Zhang, N. Li, Q. Xue, T. Wu, S. Hou and Y. Wang, Chem. Commun., 2019, 55, 5427-5430.
- 29. A. Dong, J. Chen, P. M. Vora, J. M. Kikkawa and C. B. Murray, Nature, 2010, 466, 474-477.
- 30. Y.-Y. Cai, A. Fallah, S. Yang, Y. C. Choi, J. Xu, A. Stein, J. M. Kikkawa, C. B. Murray, N. Engheta and C. R. Kagan, Adv. Mater., 2023, 35, 2301323.
- 31. J. C. Conrad and M. L. Robertson, JACS Au, 2023, 3, 333-343.

- J. W. Yu, H. Yun, W. B. Lee and Y. Kim, Adv. Sci., 2024, 11 32. 2406720. DOI: 10.1039/D5NR03413B
- 33. M. A. Boles and D. V. Talapin, J. Am. Chem. Soc., 2015, **137**, 4494-4502.
- 34. S. Pichler, M. I. Bodnarchuk, M. V. Kovalenko, M. Yarema, G. Springholz, D. V. Talapin and W. Heiss, ACS Nano, 2011, **5**, 1703-1712.
- 35. B. H. Savitzky, R. Hovden, K. Whitham, J. Yang, F. Wise, T. Hanrath and L. F. Kourkoutis, Nano Lett., 2016, 16, 5714-
- 36. P. A. Kralchevsky, N. D. Denkov and K. D. Danov, Langmuir, 2001, 17, 7694-7705.
- N. Pakalidou, J. Mu, A. J. Masters and C. Avendaño, Mol. 37. Syst. Des. Eng., 2020, 5, 376-384.
- 38. J. A. Millan, D. Ortiz and S. C. Glotzer, Soft Matter, 2015, **11**, 1386-1396.
- 39. E. S. Harper, R. L. Marson, J. A. Anderson, G. van Anders and S. C. Glotzer, Soft Matter, 2015, 11, 7250-7256.
- 40. L. Xing, Y. Xiahou, X. Zhang, W. Du, P. Zhang and H. Xia, ACS Appl. Mater. Interfaces, 2022, 14, 13480-13489.
- 41. S. Simitha, D. Mohan, S. Regi, J. Jacob and V. Ipe Thomas, Mater. Today: Proc., 2023, 93, 66-72.
- 42. X. Ye, J. Chen, B. T. Diroll and C. B. Murray, Nano Lett., 2013, 13, 1291-1297.
- 43. W.-k. Koh, S. R. Saudari, A. T. Fafarman, C. R. Kagan and C. B. Murray, Nano Lett., 2011, 11, 4764-4767.
- 44. P. Liljeroth, K. Overgaag, A. Urbieta, B. Grandidier, S. G. Hickey and D. Vanmaekelbergh, Phys. Rev. Lett., 2006, 97, 096803.

Open Access Article. Published on 19 September 2025. Downloaded on 9/21/2025 11:32:02 PM.

BY This article is licensed under a Creative Commons Attribution 3.0 Unported Licence.

The data supporting this article have been included as part of the Supplementary Information. The data that support the findings of this study are available from the corresponding author upon reasonable request.