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Methacrylate-based copolymers as tunable hosts for triplet–triplet annihilation upconversion†

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The ability to convert light to higher energies through triplet–triplet annihilation upconversion (TTA-UC) is attractive for a range of applications including solar energy harvesting, bioimaging and anti-counterfeiting. Practical applications require integration of the TTA-UC chromophores within a suitable host, which leads to a compromise between the high upconversion efficiencies achievable in liquids and the durability of solids. Herein, we present a series of methacrylate copolymers as TTA-UC hosts, in which the glass transition temperature (T_g), and hence upconversion efficiency can be tuned by varying the co-monomer ratios (*n*-hexyl methacrylate (HMA) and 2,2,2-trifluoroethyl methacrylate (TFEMA)). Using the model sensitiser/emitter pair of palladium(II) octaethylporphyrin (PdOEP) and diphenylanthracene (DPA), the upconversion quantum yield was found to increase with decreasing glass transition temperature, reaching a maximum of $1.6 \pm 0.2\%$ in air at room temperature. Kinetic analysis of the upconversion and phosphorescence decays reveal that increased PdOEP aggregation in the glassy polymers leads to a competitive non-radiative relaxation pathway that quenches the triplet state. Notably, the threshold intensity is highly sensitive to the glass transition temperature, ranging from 1250 mW cm^{-2} for PHMA₉₀TFEMA₁₀ ($T_g = -9.4 \text{ }^\circ\text{C}$) to $\sim 200 \text{ mW cm}^{-2}$ for more 'glassy' hosts, e.g. PHMA₃₃TFEMA₆₇ ($T_g = 20.1 \text{ }^\circ\text{C}$), suggesting the TTA-UC mechanism switches from diffusion-based collisions to triplet exciton migration at localised sensitiser–emitter pairs.

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Introduction

Two low-energy photons can be converted to one high-energy photon using sensitised triplet–triplet annihilation upconversion (TTA-UC).^{1,2} This process has wide-reaching applications including solar energy harvesting,^{3–5} photo-catalysis,⁶ sensors,^{7–9} bioimaging,^{10–14} and anti-counterfeiting.¹⁵ However, the uptake for some of these applications has been limited by the practicality of incorporating suitable TTA-UC systems in mass-produced devices.^{16,17} The complexity of the TTA-UC mechanism places particular demands on both the luminophores and the host materials in which it is intended to occur.

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Fig. 1 illustrates the TTA-UC mechanism, which uses a pair of luminophores: a sensitiser and an emitter. For efficient TTA-UC, the sensitiser should strongly absorb the incident photons to populate the first singlet excited state (S_1), before undergoing intersystem crossing (ISC) at a high rate to populate the first

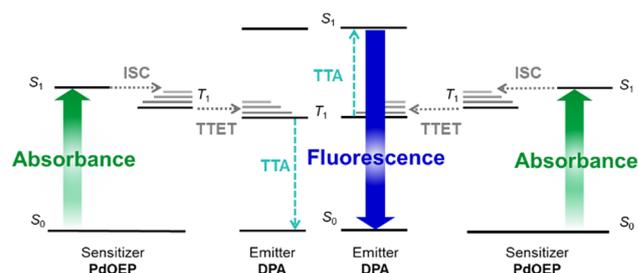


Fig. 1 Energy level diagram illustrating the TTA-UC mechanism: absorption of an incident photon (green arrow) by a sensitiser, such as palladium(II) octaethylporphyrin (PdOEP), resulting in excitation from the ground state (S_0) to the first singlet excited state (S_1); intersystem crossing (ISC, grey dashed arrow), to the first triplet excited state (T_1); triplet–triplet energy transfer (TTET, grey dashed arrow) from the sensitiser to an emitter such as 9,10-diphenylanthracene (DPA); triplet–triplet annihilation (TTA, turquoise dashed arrows) between two T_1 excited emitters, such that one populates the S_1 state and the other returns to S_0 ; fluorescence (blue arrow) of the S_1 emitter to return to the ground state.

triplet excited state (T_1). On collision between a triplet-excited sensitizer and a ground-state emitter, a good energy match between T_1 states should facilitate effective triplet–triplet energy transfer (TTET). Two triplet-excited emitters may then collide and undergo triplet–triplet annihilation (TTA), such that one emitter populates the higher-energy S_1 state, while the second relaxes to the ground-state.¹⁸ Finally, this singlet-excited emitter should fluoresce with a high photoluminescence quantum yield (Φ_{PL}), emitting a photon at the desired upconverted energy. These multimolecular processes require a high degree of luminophore mobility for the necessary collisions to occur. High rates of chromophore diffusion have been found to be the main factor in enhancement of the TTET efficiency, providing the emitter concentration is not too large ($>10^{-1}$ M).¹⁹

As a result, efficient TTA-UC has been demonstrated in liquid solutions, with upconversion quantum yields (Φ_{UC}) exceeding 35% (as a two-photon process Φ_{UC} has a maximum cap of 50%).^{20,21} However, potential issues with leakage and solvent evaporation make liquid systems impractical for device integration, so the development of solid-state hosts is desirable. These hosts should facilitate triplet exciton diffusion, allow homogeneous distribution of luminophores without aggregation, and possess high optical transparency, whilst retaining a robust yet flexible mechanical scaffold.²² Furthermore, for systems operating in air, ingress of molecular oxygen is detrimental due to quenching of triplet excited states.²³ Therefore, unless the system is to be fabricated and encapsulated in air-free conditions,^{24,25} an ideal host should also act as an effective barrier to oxygen.^{26,27}

Due to their tuneable properties, organic polymers have emerged as promising TTA-UC hosts.²² The host performance strongly depends on the glass transition temperature, T_g . Typically the TTA-UC efficiency is maximised using a polymer host whose T_g is below the required operating temperature.²⁸ Chromophore mobility has been shown to have a crucial effect on UC activity, previously shown by lowering the operating temperature of a low T_g ethyleneoxide-epichlorohydrin polymer host.²⁸ Above room temperature, the UC emission was clearly visible, whereas no UC was observed below 280 K which was comparable to the T_g of the material. It was concluded that for low chromophore concentrations where triplet exciton migration is not possible, fast mobility is essential for successful TTA-UC and the polymer host should be in its rubbery state. Cross-linked elastomers such as polyurethanes and alkyl acrylates are in their rubbery state at usual operating temperatures and provide excellent mechanical stability;^{28,29} however, they are difficult to recycle or reuse, which poses ethical questions around their sustainable use.³⁰ Meanwhile, uncrosslinked polymers such as siloxanes operating well above their T_g behave increasingly as liquids, presenting the same issues of leakage and instability.³¹ To achieve a compromise between stability, efficiency, and processability, the ideal host would therefore be an uncrosslinked polymer whose T_g is close to, but still below, the intended working temperature.

Methacrylate polymers have been previously investigated as TTA-UC hosts.^{32–34} While upconversion was observed in poly(methyl methacrylate) (PMMA) glasses (a $T_g = 92$ °C), a

high chromophore concentration (0.005% w/w PdOEP and 25% w/w DPA) was required.³⁴ The strategy of increasing emitter concentration can lead to aggregation and subsequent quenching of emission.³³ Polyacrylate elastomers of decreasing T_g were investigated by Monguzzi *et al.*²⁹ at low chromophore concentrations (0.1 mM PdOEP, 10 mM DPA), where it was concluded that the diffusion length of excited chromophores is extended as the rigidity of host decreases. The host with the lowest T_g of -62 °C achieved the highest TTA-UC efficiency of 21%, highlighting the importance of the host state on overall performance.

Previous studies have focused on the development of a specific host system whose properties are tuned to suit specific operation conditions.^{19,35,36} Here, we take an alternative approach, in which we design a series of structurally-related polymer hosts, whose properties can be tuned to the device requirements. This is achieved through copolymerisation of a low T_g methacrylate monomer with a significantly higher T_g comonomer such that, by varying the ratio, a spectrum of intermediate T_g values may be targeted. The resultant copolymer hosts were doped with the benchmark TTA-UC sensitizer–emitter pair of palladium(II) octaethylporphyrin (PdOEP) and 9,10-diphenylanthracene (DPA) to assess the correlation between the thermal properties of the host, and the TTA-UC characteristics. Using a combination of steady-state and time-resolved spectroscopic analysis, we propose that the TTA-UC mechanism transitions from diffusion-based collisions to triplet exciton migration at localised sensitizer as the glass transition temperature of the host increases.

Results and discussion

Design of the copolymer host system

The choice of host system was determined by several requirements: (i) the polymer class should satisfy the outlined general host demands for TTA-UC; (ii) there should be a wide selection of functional monomers available; (iii) the identified monomers should be compatible with copolymerisation. The methacrylate homopolymer family has previously been demonstrated to be a suitable host for TTA-UC, with Φ_{UC} from 0.08–3.5%,^{37–39} primarily using alkyl-functional monomers due to their low glass transition temperatures (T_g).²⁹ Moreover, there is a wide variety of available methacrylate monomers suited to copolymerisation, which affords tunability in the T_g . We opted to use *n*-hexyl methacrylate (HMA), and 2,2,2-trifluoroethyl methacrylate (TFEMA) as monomers (Fig. 2), since their respective homopolymers show significant difference in their T_g s (-5 °C and 74 °C, respectively).⁴⁰

A series of poly(*n*-hexyl methacrylate-*co*-2,2,2-trifluoroethyl methacrylate) copolymers was synthesised using reversible-addition fragmentation transfer (RAFT) polymerisation, using 2-cyanopropan-2-yl dodecyl trithiocarbonate as a chain transfer agent (CTA)⁴¹ and varying ratios of TFEMA (*m*) and HMA (*n*) to tune the T_g , as outlined in Fig. 3 and Table 1. RAFT polymerisation was chosen to give control over the molecular weight (M_w) of the synthesised polymers, and as such minimise any contribution from M_w variation to the photophysical properties.



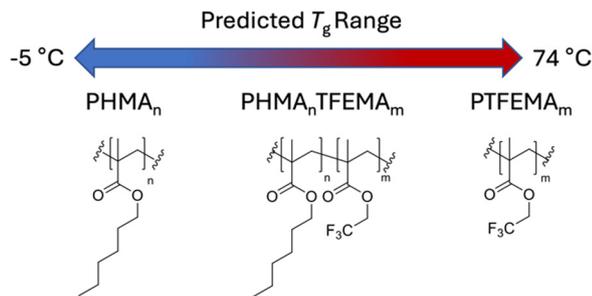


Fig. 2 Chemical structures of the homopolymers, PHMA and PTFEMA, and the resultant copolymer poly(*n*-hexyl methacrylate-*co*-2,2,2-trifluoroethyl methacrylate) (PHMA_{*n*}TFEMA_{*m*}). The monomer ratios, *n* and *m*, were varied to access a range of glass transition temperatures, *T*_g.

All synthesised polymers were in the range 7400–16 400 g mol⁻¹ (Table 1). A key drawback of the RAFT method is that the CTA can give strong colouration to the final material, ranging from pink to yellow depending on the exact CTA used.⁴² In our case, the CTA resulted in a strong yellow colouration (see Fig. S4, ESI[†]), which would lead to parasitic absorption that would be detrimental to the TTA-UC efficiency. To overcome this, the CTA end group was removed post-polymerisation *via* reduction with azobisisobutyronitrile (AIBN) and tributyltin hydride to leave a single hydrogen as the chain end (Fig. 3, step 2). This mechanism has been reported to have a quantitative yield and to be particularly effective for methacrylic polymers.⁴³ The resulting end-reduced polymers showed no coloration and retained good optical clarity (Fig. S4, ESI[†]). The final copolymers are denoted as PHMA_{*n*}TFEMA_{*m*}, where *n* and *m* are the molar percentages for HMA and TFEMA, respectively. All polymers, before and after end-reduction were fully characterised by ¹H and ¹³C nuclear magnetic resonance (NMR) spectroscopy (Tables S1, S2 and Fig. S8–S39, ESI[†]) and size-exclusion chromatography (SEC, Fig. S40, ESI[†]). Full details of the synthetic methodology and characterisation data are available in the ESI[†] (Sections 1–3).

Thermal and flow properties

The ratios of HMA and TFEMA used to fabricate the copolymers were identified using the using the Fox equation (eqn (1)),⁴⁴

which predicts the expected *T*_g of the copolymer based on the weight fractions *W*_{*i*} of each comonomer and the glass transition temperatures of their respective homopolymers (*T*_{g,*i*}):

$$\frac{1}{T_g} = \frac{W_1}{T_{g,1}} + \frac{W_2}{T_{g,2}} \quad (1)$$

It should be noted that the similar molecular weights of HMA (170.25 g mol⁻¹) and TFEMA (168.11 g mol⁻¹) mean the weight fractions are approximately equal to the molar fractions of each comonomer, and will therefore be referred to as “fraction” for simplicity. The Fox equation predicts a decrease in *T*_g with the increasing fraction of HMA, notably dropping to *T*_g ~ 22 °C with the inclusion of 60% HMA (Table 1), suggesting this is the minimum weight percentage required for a high-diffusion TTA-UC host. Qualitative inspection of the synthesised polymers supports this predicted transition point, with polymers containing 67–100% HMA existing as highly viscous liquids, while <50% resulted in glassy solids. Notably, PHMA₃₃TFEMA₆₇ exhibited a variable form, from a highly viscous liquid to a glassy solid as the laboratory temperature varied (between 18–24 °C).

Differential scanning calorimetry (DSC) was performed to assess the validity of the predicted to experiment *T*_g values. As shown in Table 1, the measured *T*_g values are all notably lower than those predicted, which can be rationalised based on the relatively low *M*_w of the polymers, which results in a higher free volume due to the chain ends and hence lower *T*_g than those calculated for infinite chains. While these lower *T*_g values will result in a shift towards more liquid-like behaviour for all copolymers, the compositional region in which the behaviour of the synthesised polymers shifts from glassy to viscous liquid remains around 40–50% HMA. Crucially, the trend is as predicted: the *T*_g decreases as HMA content is increased, such that the PHMA_{*n*}TFEMA_{*m*} series spans a range of properties from glassy solids to viscous liquids at room temperature.

Steady-state optical properties

Polymers were doped with PdOEP (0.3 mM) and DPA (30 mM) as the TTA-UC sensitizer–emitter pair and cast as thick films (~200 μm) on glass substrates. This chromophore pair has been extensively studied for green-blue upconversion, providing a useful benchmark to assess the host performance.^{35,45,46} The sensitizer-to-emitter ratio was selected based on previous studies using samples of similar thickness to facilitate comparison.^{26,29} Fig. 4a presents the transmittance spectra for both undoped and doped PHMA₅₀TFEMA₅₀ copolymer films (see Fig. S42–S44, ESI[†] for corresponding data for all other samples). The undoped copolymers exhibit good optical transmittance in the visible range (85–95%). When doped with the TTA-UC luminophore pair, the polymer films retain high transmittance above 600 nm, but absorb strongly from 490–560 nm (PdOEP, Q band) and 320–420 nm (DPA S₀ → S₁ and PdOEP Soret band) as expected. The individual absorption and emission spectra of DPA and PdOEP in tetrahydrofuran solutions (5 μM) are provided in Fig. S45, ESI[†] for reference.

Photoluminescence quantum yields (Φ_{PL}) were measured for polymer films doped with 30 mM of DPA (DPA-only, Table S4, ESI[†]) and for the TTA-UC doped films (PdOEP:DPA, Table 2) to

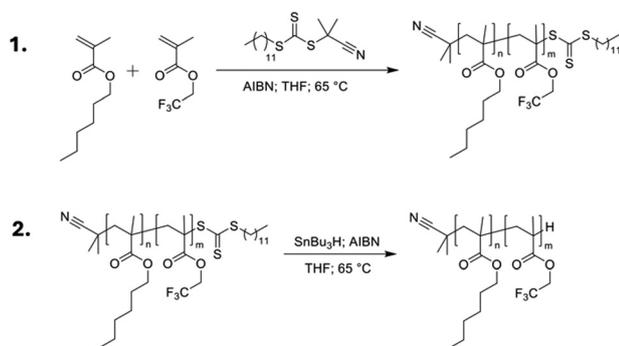


Fig. 3 Synthesis of methacrylate copolymers. Step 1: polymerisation of *n*-hexyl methacrylate and 2,2,2-trifluoroethyl methacrylate *via* RAFT polymerisation. Step 2: Radical reduction of trithiocarbonate using AIBN and tributyl tin hydride to remove coloration from the final polymers.



Table 1 Thermal properties and molecular weights of methacrylate (co)polymers

Polymer	Form	$T_{g(\text{Calc})}^a$ (°C)	T_g^b (°C)	M_n^c (g mol ⁻¹)	M_w^d (g mol ⁻¹)	D^e
PTFEMA ₁₀₀	Glassy solid	74	43.3	7400	10 300	1.40
PHMA ₃₃ TFEMA ₆₇	Glassy solid ^f	43	20.1	10 200	13 600	1.33
PHMA ₅₀ TFEMA ₅₀	Rubbery solid	29	16.6	10 700	14 300	1.34
PHMA ₆₀ TFEMA ₄₀	Highly viscous liquid	22	10.4	11 500	15 000	1.31
PHMA ₆₇ TFEMA ₃₃	Highly viscous liquid	17	1.0	12 100	15 600	1.29
PHMA ₈₀ TFEMA ₂₀	Viscous liquid	8	-8.49	14 300	18 200	1.27
PHMA ₉₀ TFEMA ₁₀	Viscous liquid	1	-9.42	15 700	19 500	1.24
PHMA ₁₀₀	Viscous liquid	-5	-10.26	16 400	20 500	1.25

^a Calculated glass transition temperature (from eqn (1)). ^b Measured glass transition temperature (from DSC). ^c Number-average molecular weight (from SEC). ^d Weight-average molecular weight (from SEC). ^e Dispersity, calculated as M_w/M_n . ^f Form varied with laboratory temperature due to overlap with T_g .

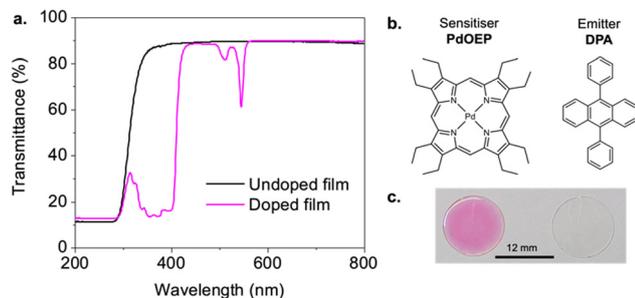


Fig. 4 Optical properties of methacrylate copolymer hosts. (a) Transmittance spectra of undoped (black) and doped (pink) PHMA₅₀TFEMA₅₀ copolymer film. (b) Chemical structures of sensitizer PdOEP and emitter DPA. (c) Photographs of a lumophore-doped film (left), and undoped film (right).

understand the impact of the host. In DPA-only samples, upon direct excitation at 375 nm, the Φ_{PL} values exceeded 85% for all hosts (Fig. 5a), with the PHMA₆₇TFEMA₃₃ sample reaching 96%, comparable to that of DPA in solvents.⁴⁷ In contrast, the PdOEP:DPA-doped polymer films exhibited a significant decrease in Φ_{PL} , ranging from 33 to 50% due to parasitic absorption from the Soret band of PdOEP. As shown in Fig. 5a, the Φ_{PL} increases with increasing HMA content, from 49.7% in PHMA₆₇TFEMA₃₃ to 52.8% in PHMA₁₀₀ - the lowest T_g matrix.

Upconversion characterisation

To quantify the upconversion activity in the different polymer hosts, the upconversion quantum yield was measured (see Section 2.6.1, ESI†). The Φ_{UC} can be expressed as:³⁵

$$\Phi_{\text{UC}} = \frac{1}{2} f \Phi_{\text{ISC}} \Phi_{\text{TTET}} \Phi_{\text{TTA}} \Phi_{\text{PL}} \quad (2)$$

where Φ_{ISC} , Φ_{TTET} , Φ_{TTA} and Φ_{PL} are the quantum yields for ISC, TTET, TTA, and emitter photoluminescence, respectively. The factor f represents the probability that a singlet excited state is produced per TTA step.⁴⁸ Note that we will use the conventional description for reporting Φ_{UC} , which is capped at 50% due to the factor of $\frac{1}{2}$ in eqn (2).

Fig. 5b illustrates how Φ_{UC} varies as a function of HMA content in the polymer. The increase in Φ_{UC} with HMA content can be attributed to the increase in chromophore mobility with decreasing T_g , with the highest Φ_{UC} ($1.6 \pm 0.2\%$) obtained for PHMA₁₀₀, the most liquid-like host. A transition can also be seen from a steep increase at lower HMA content to a shallower slope above $\sim 67\%$ HMA. This corresponds to the transition in host properties, from glassy solids with T_g s above room temperature, to viscous liquids with sub-ambient T_g values. For the more liquid-like hosts, the diffusivity will be higher, and hence less limiting to the upconversion efficiency. Notably, PHMA₆₇TFEMA₃₃ presents promising results, as it combines a Φ_{UC} of $1.3 \pm 0.3\%$, comparable to those in lower- T_g hosts, with solid-like behaviour. This is preferable for applications at and above

Table 2 Key photophysical properties of PdOEP : DPA (0.3 mM : 30 mM) doped PHMA_nTFEMA_m hosts

Polymer	Φ_{PL}^a (%)	Φ_{UC}^b (%)	UC lifetimes (collection at 440 nm) ^c				Phosphorescence lifetimes (collection at 660 nm) ^d			
			τ_1 (ms)/ f_1 (%)	τ_2 (ms)/ f_2 (%)	$\langle \tau_{\text{UC}} \rangle$ (ms)	χ^2	τ_1 (ms)/ f_1 (%)	τ_2 (ms)/ f_2 (%)	$\langle \tau_{\text{phos}} \rangle^a$ (ms)	χ^2
PHMA ₁₀₀	53 ± 2	1.6 ± 0.2	1.1/29	3.1/71	2.52	1.085	0.36/22	2.3/78	1.88	1.453
PHMA ₉₀ TFEMA ₁₀	42 ± 3	1.4 ± 0.1	0.2/38	0.59/62	0.44	1.381	0.51/17	1.4/83	1.24	1.362
PHMA ₈₀ TFEMA ₂₀	46 ± 4	1.2 ± 0.2	0.18/42	0.48/58	0.35	1.365	0.29/11	1.4/89	1.27	1.413
PHMA ₆₇ TFEMA ₃₃	50 ± 4	1.3 ± 0.2	0.18/35	0.61/65	0.46	1.522	0.15/22	1.5/78	1.21	1.451
PHMA ₆₀ TFEMA ₄₀	43 ± 4	1.2 ± 0.2	0.15/41	0.43/59	0.32	1.388	*	*	*	*
PHMA ₅₀ TFEMA ₅₀	42 ± 3	0.7 ± 0.2	0.16/39	0.45/61	0.34	1.247	*	*	*	*
PHMA ₃₃ TFEMA ₆₇	34 ± 6	0.5 ± 0.1	0.11/57	0.27/43	0.18	1.358	0.09/29	0.71/71	0.53	1.400
PTFEMA ₁₀₀	34 ± 4	—	0.027/72	0.074/28	0.04	1.134	*	*	*	*

^a Photoluminescence quantum yield of the emitter DPA ($\lambda_{\text{ex}} = 375$ nm, $\lambda_{\text{em}} = 380$ –530 nm). ^b Upconversion quantum yield at excitation power intensity of 1 W cm^{-2} (532 nm). ^c Lifetime fitting data for UC decay curves (average lifetime ($\langle \tau_{\text{UC}} \rangle$), individual lifetimes (τ_i) and fractional contributions (f_i) and goodness-of-fit (χ^2)). $\lambda_{\text{ex}} = 532$ nm, $\lambda_{\text{em}} = 440$ nm. ^d Lifetime fitting data for phosphorescence decay curves (average lifetime ($\langle \tau_{\text{phos}} \rangle$), individual lifetimes (τ_i) and fractional contributions (f_i) and goodness-of-fit (χ^2)). $\lambda_{\text{ex}} = 532$ nm, $\lambda_{\text{em}} = 660$ nm. *There is no detectable signal.



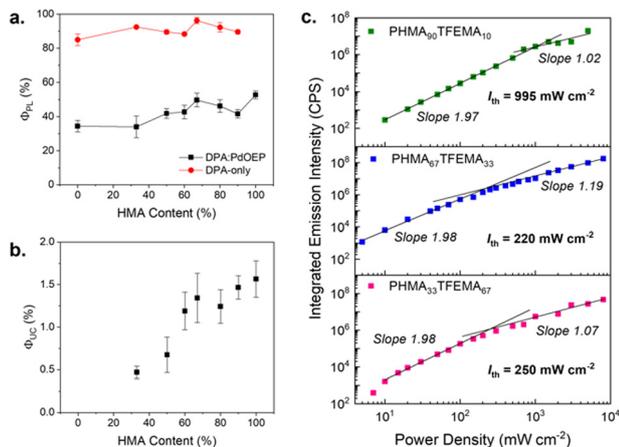


Fig. 5 Upconversion and photoluminescence of methacrylate copolymers. (a) Photoluminescence quantum yield of the emitter DPA with sensitizer present (red circles) and without sensitizer (black squares) with increasing hexyl content of the copolymer, $\lambda_{\text{ex}} = 375$ nm, $\lambda_{\text{em}} = 380$ –530 nm. (b) Upconversion quantum yield at excitation intensity of 1 W cm^{-2} , $\lambda_{\text{ex}} = 532$ nm, $\lambda_{\text{em}} = 380$ –500 nm. (c) Threshold intensity was determined from a logarithmic plot of integrated UC emission intensity at varying laser power (532 nm) for PHMA₉₀TFEMA₁₀ (green), PHMA₆₇TFEMA₃₃ (blue), and PHMA₃₃TFEMA₆₇ (pink).

room temperature, compared to more liquid films like those of PHMA₁₀₀. While these efficiencies are below that of the highest performing acrylate hosts of 21%,²⁹ the purpose of this study is to investigate the tunability of host properties *via* copolymerisation to suit a variety of applications at varying operating temperatures. Such studies in methacrylate polymers are usually conducted in melt-processed glassy matrixes with high chromophore concentrations *via* triplet exciton migration,^{32–34} which rely on close chromophore proximity, and are avoided in the low T_g host and mechanism presented here.

Threshold intensity

Fig. 5c shows the threshold intensity (I_{th}) for three PdOEP:DPA-doped copolymers: PHMA₉₀TFEMA₁₀, PHMA₆₇TFEMA₃₃ and PHMA₃₃TFEMA₆₇. These hosts were selected to compare systems with low, ambient and high T_g values. I_{th} , also known as the power density threshold, is the excitation at which the Φ_{TTA} is 50% of the maximum value.^{48,49} Experimentally, I_{th} is determined by locating the point at which the TTA-UC emission intensity shifts from a quadratic (where the slope of the emission intensity relative to power is 2) to a first-order dependence on incident light intensity (where the slope is 1).⁵⁰ Systems with lower thresholds are generally favoured and will convert more efficiently under low-power sources such as natural sunlight (100 mW cm^{-2}).^{48,51} It should be noted that, in our systems, at laser power densities $> 500 \text{ mW cm}^{-2}$, photobleaching occurred before the full emission spectrum could be measured which led to a plateau at high intensities, which is not an uncommon issue with this measurement.⁵² Photodegradation at high power was quantified for each film (Fig. S46, ESI[†]), where copolymers show a higher resistance to degradation than the mono-polymers. All samples showed fast photodegradation at very high laser power

(4 W cm^{-2}) as demonstrated by the decreased UC emission intensity and the appearance of clear spot on the film at the point of laser incidence.

For the more liquid-like PHMA₉₀TFEMA₁₀ ($T_g = -9.4$ °C) the transition from quadratic to first-order dependence transition is clearly discernible and corresponds to $I_{\text{th}} = 995 \text{ mW cm}^{-2}$. In contrast, for the more ‘glassy’ hosts, PHMA₃₃TFEMA₆₇ ($T_g = 20.1$ °C) and PHMA₆₇TFEMA₃₃ ($T_g = 1.0$ °C), this transition is more subtle and leads to significantly lower I_{th} values, around 220–250 mW cm^{-2} . This large decrease in the I_{th} value indicates that changing the T_g of the host can completely alter the TTA-UC mechanism. Low threshold intensities have been reported for crystalline TTA-UC materials, which often exhibit large triplet exciton diffusion lengths of the acceptor molecules due to localised chromophore aggregation.^{53–56} It seems plausible that increased chromophore aggregation will occur in the higher T_g hosts (see following section), which would support a localised exciton diffusion mechanism and the observed decrease in I_{th} . Note that although the steady-state photoluminescence spectra indicate a reduction in the first vibronic peak due to reabsorption (Fig. S47, ESI[†]), lifetime measurements suggest DPA aggregation also makes a contribution (Fig. S48 and Table S5, ESI[†]). Compared to the monoexponential decay of DPA in dilute THF solution, DPA-only doped methacrylate copolymer films show a biexponential decay, where the short lifetime, τ_1 (~ 5 ns) corresponds to quenched species due to aggregation, and τ_2 (~ 10 ns) exhibits a longer lifetime associated with the reabsorption effect. Importantly, the contribution of τ_1 is significantly higher for DPA-only PTFEMA₁₀₀, supporting the argument that stronger aggregation is observed in the high T_g hosts. In contrast, in ‘soft’ hosts, such as PHMA₉₀TFEMA₁₀, Brownian diffusion of the acceptor molecules is expected to determine and limit the maximum achievable triplet exciton diffusion length of the acceptor,⁵⁷ leading to the large I_{th} . This change in mechanism would be expected to affect multiple parameters, such as the rates of ISC, TTET, and triplet lifetimes.⁴⁹

Time-resolved emission measurements

Lifetime measurements were performed to gain better understanding of the variation in upconversion behaviour across the series of polymers. The polymers vary significantly in T_g and, consequently, in viscosity of the host matrix, which is expected to affect the kinetics of excited state decay of both luminophores (DPA and PdOEP). Fig. 6a shows the upconversion decay traces for the series of PdOEP:DPA-doped polymer films (collected at 440 nm). The UC lifetimes (τ_{UC}) values were estimated by tail-fitting a biexponential decay function. Fig. 6b and Table 2 summarise the average UC lifetimes ($\langle \tau_{\text{UC}} \rangle$), which ranged from 0.04 ms in PTFEMA₁₀₀ to 2.51 ms in PHMA₁₀₀. In agreement with the Φ_{UC} values, a significant increase in the $\langle \tau_{\text{UC}} \rangle$ is observed as the HMA content of the polymer films increased. This suggests that for the viscous liquid film (PHMA₁₀₀), the higher Φ_{UC} value is accompanied by a longer $\langle \tau_{\text{UC}} \rangle$ indicating reduced competitive quenching mechanisms in this matrix.

The average phosphorescence lifetimes in air ($\langle \tau_{\text{phos}} \rangle$, collection at 660 nm, Fig. S48, ESI[†]) were obtained for the same



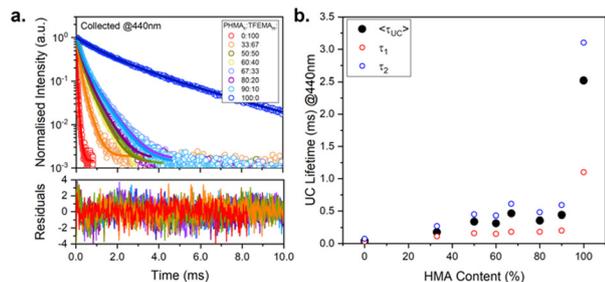


Fig. 6 Upconversion kinetic studies on PdOEP:DPA (0.3 mM : 30 mM) doped in different PHMA_nTFEMA_m host matrices. (a) Upconversion decay traces (open symbols) and biexponential decay fits (solid lines) and corresponding residuals. $\lambda_{\text{ex}} = 532$ nm, $\lambda_{\text{em}} = 440$ nm. (b) Corresponding upconversion average lifetimes (filled black circles) and individual lifetime components (τ_1 and τ_2) obtained from decay fits.

samples from biexponential tail fits, comprising a short lifetime component, $\tau_1 \sim 0.1\text{--}0.5$ ms, and a longer-lived component, $\tau_2 \sim 1.4$ ms, as summarised in Table 2. We assign τ_2 to the natural phosphorescence lifetime, since in previous studies of PdOEP-only in polymer hosts, a characteristic τ_{phos} of 1.4 ms was observed in the absence of oxygen.⁵⁸ The shorter lifetime observed in these PdOEP:DPA-doped polymer samples, τ_1 , is assigned to the quenched triplet state, which arises from competing relaxation pathways such as aggregation-caused quenching, oxygen quenching and TTET. We observed that polymer hosts with particularly low T_g , e.g. PHMA₉₀TFEMA₁₀, exhibit a high fractional contribution of τ_2 , even in air. Meanwhile, those with higher T_g – e.g. PHMA₃₃TFEMA₆₇ – show greater fractional contribution from τ_1 , leading to a significantly shorter $\langle\tau_{\text{phos}}\rangle$. This same trend is also observed in the analogous PdOEP-only doped films (Fig. S49 and Table S5, ESI[†]).

To further understand the triplet relaxation mechanisms, phosphorescence decay curves (Fig. 7 and Table S6, ESI[†]) were measured for representative low (PHMA₁₀₀) and high (PTFEMA₁₀₀) T_g homopolymer hosts doped with PdOEP-only, in air and N₂ atmospheres. In both cases, τ_2 became more

prevalent with deaeration, further supporting the assignment of this component to the natural phosphorescence lifetime of the unquenched species (around 1.2–1.5 ms). Its presence even under ambient conditions suggests that the second coverslip used to aid sample handling prevented further oxygen ingress during the measurement, with any local oxygen already present inside the matrix being rapidly photoconsumed, especially in PHMA₁₀₀ where oxygen is expected to diffuse more readily. On the other hand, for PTFEMA₁₀₀, even in the absence of oxygen, the decay trace still exhibited a significant contribution from τ_1 , suggesting that aggregation-caused quenching is a significant relaxation pathway. We note that as DPA is not present in these samples, the contribution from TTET is also excluded. The phosphorescence lifetime analysis thus supports the presence of increased PdOEP aggregation in the high T_g polymer hosts. Aggregation is also evident in the DPA fluorescence lifetimes (Fig. S48 and Table S5, ESI[†]), where PdOEP:DPA-doped films exhibit shorter average lifetimes than DPA-only doped films, suggesting additional non-radiative pathways for these samples. This can result in longer triplet exciton diffusion lengths for the acceptors, consistent with the lower threshold intensity observed in these samples.

Conclusions

We have demonstrated the formation of tunable TTA-UC host materials through RAFT copolymerisation of two structurally-similar methacrylate monomers, HMA and TFEMA. Through variation of the co-monomer ratios, a series of PHMA_nTFEMA_m copolymers were synthesised, whose glass transition temperatures spanned a range from below to above room temperature. For the model TTA-UC pair (PdOEP:DPA), the upconversion quantum yields increased with decreasing glass transition temperature, across the range 0.5–1.6% in air at room temperature. This is accompanied by a five-fold decrease in the threshold intensity, which we propose is due to increased luminophore aggregation in the higher T_g hosts, leading to a switch in mechanism from diffusion-based collisions to triplet exciton migration at localised sensitiser–emitter pairs. Notably, this work demonstrates the versatility of copolymerisation to optimise the physical form to meet a required operating temperature or application, without dramatically changing the upconversion efficiency. Further tuning of the chromophore concentrations and limiting of oxygen diffusion and subsequent quenching could be employed in future work to improve efficiency of these hosts. The insight gained through this study will inform the design of next generation TTA-UC hosts, with the aim of increasing the upconversion efficiency in real-world conditions.

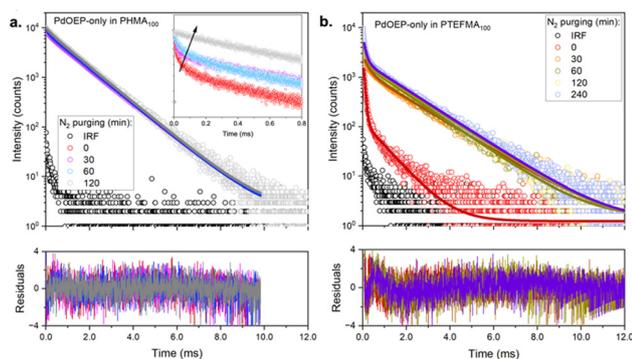


Fig. 7 Phosphorescence decay traces of PdOEP-only (0.3 mM) doped in: (a) PHMA₁₀₀ and (b) PTFEMA₁₀₀ under varying durations of nitrogen (N₂) purging. Inset: Zoom in on the first millisecond regimes to highlight the increase in the shorter lifetime after increased N₂ purging. Measurements were performed using a 532 nm excitation source and collected at 660 nm.

Data availability

Data for this article are available at the University of Cambridge Apollo repository at <https://www.repository.cam.ac.uk/home>.



Author contributions

M. J. B.: conceptualisation, methodology, investigation, data curation, formal analysis, writing – original draft, writing – reviewing & editing. A. R. C.: investigation, data curation, formal analysis, writing – original draft, reviewing & editing. L. G. F.: investigation, data curation, formal analysis, writing – original draft, writing – reviewing & editing. G. H. B. M.: data curation, formal analysis, writing – reviewing & editing. N. W. F.: investigation, data curation. R. D.: resources. J. J. G. K.: investigation and data curation. B. L. C.: investigation and data curation. R. C. E.: conceptualisation, project administration, resources, funding acquisition, supervision, writing – reviewing & editing.

Conflicts of interest

There are no conflicts to declare.

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Notes and references

- C. A. Parker, C. G. Hatchard and T. A. Joyce, *Nature*, 1965, **205**, 1282–1284.
- C. A. Parker and C. G. Hatchard, *Proc. R. Soc. London, Ser. A*, 1962, **269**, 574–584.
- M. J. Y. Tayebjee, D. R. McCamey and T. W. Schmidt, *J. Phys. Chem. Lett.*, 2015, **6**, 2367–2378.
- J. Pedrini and A. Monguzzi, *J. Photonics Energy*, 2017, **8**, 1.
- D. Beery, T. W. Schmidt and K. Hanson, *ACS Appl. Mater. Interfaces*, 2021, **13**, 32601–32605.
- T. Yu, Y. Liu, Y. Zeng, J. Chen, G. Yang and Y. Li, *Chem. – Eur. J.*, 2019, **25**, 16270–16276.
- Y. Sun, L. Wei, S. Zhu, P. Jin, C. He, Q. He, C. Yang and W. Wu, *Sens. Actuators, B*, 2023, **387**, 133764.
- M. Xu, X. Zou, Q. Su, W. Yuan, C. Cao, Q. Wang, X. Zhu, W. Feng and F. Li, *Nat. Commun.*, 2018, **9**, 2698.
- A. A. Mendonsa, C. C. Soeldner, N. E. Mudd, S. C. Saccomano and K. J. Cash, *ACS Sens.*, 2023, **8**, 3043–3050.
- Q. Dou, L. Jiang, D. Kai, C. Owh and X. J. Loh, *Drug Discovery Today*, 2017, **22**, 1400–1411.
- W. Lin, J. Li, H. Feng, F. Qi and L. Huang, *J. Anal. Test.*, 2023, **7**, 327–344.
- S. H. C. Askes, W. Pomp, S. L. Hopkins, A. Kros, S. Wu, T. Schmidt and S. Bonnet, *Small*, 2016, **12**, 5579–5590.
- O. S. Kwon, H. S. Song, J. Conde, H. Kim, N. Artzi and J.-H. Kim, *ACS Nano*, 2016, **10**, 1512–1521.
- B. Zhang, K. D. Richards, B. E. Jones, A. R. Collins, R. Sanders, S. R. Needham, P. Qian, A. Mahadevegowda, C. Ducati, S. W. Botchway and R. C. Evans, *Angew. Chem., Int. Ed.*, 2023, **62**, e202308602.
- W. Yin, T. Yu, J. Chen, R. Hu, G. Yang, Y. Zeng and Y. Li, *ACS Appl. Mater. Interfaces*, 2021, **13**, 57481–57488.
- B. Joarder, N. Yanai and N. Kimizuka, *J. Phys. Chem. Lett.*, 2018, **9**, 4613–4624.
- A. K. Williams, B. J. Davis, E. R. Crater, J. R. Lott, Y. C. Simon and J. D. Azoulay, *J. Mater. Chem. C*, 2018, **6**, 3876–3881.
- L. G. Franca, D. G. Bossanyi, J. Clark and P. L. dos Santos, *ACS Appl. Opt. Mater.*, 2024, DOI: [10.1021/acsaom.4c00041](https://doi.org/10.1021/acsaom.4c00041).
- A. Monguzzi, R. Tubino and F. Meinardi, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2008, **77**, 155122.
- Y. Zhou, F. N. Castellano, T. W. Schmidt and K. Hanson, *ACS Energy Lett.*, 2020, **5**, 2322–2326.
- N. Yanai, K. Suzuki, T. Ogawa, Y. Sasaki, N. Harada and N. Kimizuka, *J. Phys. Chem. A*, 2019, **123**, 10197–10203.
- M. J. Bennison, A. R. Collins, B. Zhang and R. C. Evans, *Macromolecules*, 2021, **54**, 5287–5303.
- V. Gray, K. Moth-Poulsen, B. Albinsson and M. Abrahamsson, *Coord. Chem. Rev.*, 2018, **362**, 54–71.
- S. Balushev, K. Katta, Y. Avlasevich and K. Landfester, *Mater. Horiz.*, 2016, **3**, 478–486.
- D. Dzebo, K. Moth-Poulsen and B. Albinsson, *Photochem. Photobiol. Sci.*, 2017, **16**, 1327–1334.
- A. R. Collins, B. Zhang, M. J. Bennison and R. C. Evans, *J. Mater. Chem. C*, 2024, **12**, 6310–6318.
- A. J. Svagan, D. Busko, Y. Avlasevich, G. Glasser, S. Balushev and K. Landfester, *ACS Nano*, 2014, **8**, 8198–8207.
- T. N. Singh-Rachford, J. Lott, C. Weder and F. N. Castellano, *J. Am. Chem. Soc.*, 2009, **131**, 12007–12014.
- A. Monguzzi, M. Mauri, A. Bianchi, M. K. Dibbanti, R. Simonutti and F. Meinardi, *J. Phys. Chem. C*, 2016, **120**, 2609–2614.
- D. J. Fortman, J. P. Brutman, G. X. De Hoe, R. L. Snyder, W. R. Dichtel and M. A. Hillmyer, *ACS Sustainable Chem. Eng.*, 2018, **6**, 11145–11159.
- J. Fang, C. Zhou, Y. Chen, L. Fang, W. Wang, C. Zhu, Y. Ni and C. Lu, *ACS Appl. Mater. Interfaces*, 2020, **12**, 717–726.
- S. H. Lee, M. A. Ayer, R. Vadrucci, C. Weder and Y. C. Simon, *Polym. Chem.*, 2014, **5**, 6898–6904.
- S. Raišys, S. Juršėnas and K. Kazlauskas, *Sol. RRL*, 2022, **6**, 2100873.
- S. H. Lee, J. R. Lott, Y. C. Simon and C. Weder, *J. Mater. Chem. C*, 2013, **1**, 5142.
- R. R. Islangulov, J. Lott, C. Weder and F. N. Castellano, *J. Am. Chem. Soc.*, 2007, **129**, 12652–12653.
- S. H. Lee, D. C. Thévenaz, C. Weder and Y. C. Simon, *J. Polym. Sci., Part A: Polym. Chem.*, 2015, **53**, 1629–1639.
- A. Monguzzi, A. Oertel, D. Braga, A. Riedinger, D. K. Kim, P. N. Knüsel, A. Bianchi, M. Mauri, R. Simonutti, D. J. Norris and F. Meinardi, *ACS Appl. Mater. Interfaces*, 2017, **9**, 40180–40186.



- 38 L. Jiang, L. Zeng, M. Zhang, L. Huang and D. Pang, *Adv. Opt. Mater.*, 2024, **12**, 2301879.
- 39 D. Yildiz, C. Baumann, A. Mikosch, A. J. C. Kuehne, A. Herrmann and R. Göstl, *Angew. Chem., Int. Ed.*, 2019, **58**, 12919–12923.
- 40 J. Brandrup, E. H. Immergut, E. A. Grulke, A. Abe and D. R. Bloch, *Polymer Handbook*, John Wiley & Sons, 4th edn, 2005.
- 41 B. A. Abel and C. L. McCormick, *Macromolecules*, 2016, **49**, 465–474.
- 42 S. Perrier, P. Takolpuckdee and C. A. Mars, *Macromolecules*, 2005, **38**, 2033–2036.
- 43 Y. K. Chong, G. Moad, E. Rizzardo and S. H. Thang, *Macromolecules*, 2007, **40**, 4446–4455.
- 44 T. G. Fox, *Bull. Am. Phys. Soc.*, 1956, **1**, 123.
- 45 T. N. Singh-Rachford and F. N. Castellano, *Coord. Chem. Rev.*, 2010, **254**, 2560–2573.
- 46 V. Gray, D. Dzebo, M. Abrahamsson, B. Albinsson and K. Moth-Poulsen, *Phys. Chem. Chem. Phys.*, 2014, **16**, 10345–10352.
- 47 L. Hou, A. Olesund, S. Thurakkal, X. Zhang and B. Albinsson, *Adv. Funct. Mater.*, 2021, **31**, 2106198.
- 48 A. Monguzzi, R. Tubino, S. Hoseinkhani, M. Campione and F. Meinardi, *Phys. Chem. Chem. Phys.*, 2012, **14**, 4322.
- 49 N. A. Durandin, J. Isokuortti, A. Efimov, E. Vuorimaa-Laukkanen, N. V. Tkachenko and T. Laaksonen, *J. Phys. Chem. C*, 2019, **123**, 22865–22872.
- 50 P. Duan, D. Asthana, T. Nakashima, T. Kawai, N. Yanai and N. Kimizuka, *Faraday Discuss.*, 2017, **196**, 305–316.
- 51 T. Ogawa, N. Yanai, A. Monguzzi and N. Kimizuka, *Sci. Rep.*, 2015, **5**, 10882.
- 52 F. Edhborg, A. Olesund and B. Albinsson, *Photochem. Photobiol. Sci.*, 2022, **21**, 1143–1158.
- 53 P. B. Merkel and J. P. Dinnocenzo, *J. Lumin.*, 2009, **129**, 303–306.
- 54 P. Duan, N. Yanai, H. Nagatomi and N. Kimizuka, *J. Am. Chem. Soc.*, 2015, **137**, 1887–1894.
- 55 P. Duan, N. Yanai and N. Kimizuka, *J. Am. Chem. Soc.*, 2013, **135**, 19056–19059.
- 56 K. Narushima, S. Hirata and M. Vacha, *Nanoscale*, 2017, **9**, 10653–10661.
- 57 A. Monguzzi, J. Mezyk, F. Scotognella, R. Tubino and F. Meinardi, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2008, **78**, 195112.
- 58 P. Douglas and K. Eaton, *Sens. Actuators, B*, 2002, **82**, 200–208.

