

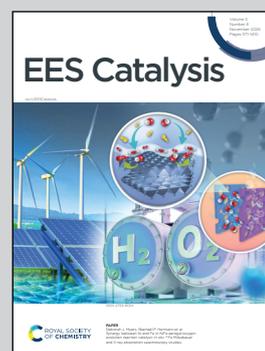
Showcasing research from the group of Dr Jörg Meyer, Theoretical Chemistry, Leiden Institute of Chemistry, Leiden University, The Netherlands.

Vibrational excitation in plasma catalysis: how important are dynamical effects?

Vibrationally excited  $N_2$  molecules have been suggested to boost the efficiency of heterogeneously catalysed ammonia synthesis, where  $N_2$  dissociation on a metal surface is the rate-determining step. To quantify the enhancement of the dissociation rate coefficient relative to  $N_2$  molecules in the vibrational ground state, we have performed molecular dynamics simulations in the quasi-classical framework on accurate high-dimensional potential energy surface obtained from density functional theory calculations. The enhancement is significantly more pronounced than found previously based on the so-called Fridman-Macheret  $\alpha$  model, which has important implications for both thermal and plasma-enabled catalysis.

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## Vibrational excitation in plasma catalysis: how important are dynamical effects?†

Floris van den Bosch,  Nick Gerrits  and Jörg Meyer \*

Plasma catalysis offers a promising alternative to current ammonia production processes, due to the combination of high selectivity of heterogeneous catalysis and efficient activation of nitrogen in the plasma. However, the theoretical understanding of how various plasma processes contribute to efficiency improvements remains limited. The pioneering work of Metha *et al.* (*Nat. Catal.*, 2018, **1**, 269) extended the standard formulation of transition state theory by making it vibrational state-specific through the use of the Fridman–Macheret  $\alpha$  model. The resulting microkinetic model accounted for vibrational contributions under the non-equilibrium conditions of a plasma reactor. In this work, we critically examine the prototypical chemical process of activated  $N_2$  reactivity on ruthenium through explicit rate coefficient calculations using state-of-the-art molecular dynamics, based on a potential energy surface previously validated against molecular beam experiments. Our findings reveal that vibrational activation is significantly more effective in promoting surface reactivity than predicted by the Fridman–Macheret  $\alpha$  model, which fails to capture the full complexity of state-specific contributions. Furthermore, our calculations indicate that vibrational activation is also the primary driver of highly activated thermal catalytic reactions. These results provide a valuable benchmark to guide the development of future state-specific microkinetic models for heterogeneous and plasma catalysis.

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### Broader context

Pre-activating reactants by a plasma holds great potential to boost heterogeneously catalysed processes, especially in scenarios where dissociative chemisorption (DC) reactions are rate-limiting. However, the fundamental atomic-scale understanding and quantification of this enhanced performance are still lacking. The pioneering work on plasma-enabled ammonia synthesis for DC of  $N_2$  introduced the Fridman–Macheret (FM)  $\alpha$  model, which is widely used to extend standard micro-kinetic modelling to account for vibrational excitations in plasma catalysis. While this model is computationally inexpensive, it relies on assumptions based on gas-phase chemistry. In this article, we scrutinise the FM- $\alpha$  model for  $N_2$  dissociation on Ru(0001) using explicit molecular dynamics simulations on an accurate potential energy surface. We find that vibrational excitation plays a far more dominant role in  $N_2$  dissociation than the FM- $\alpha$  model predicts, and the reasons behind this discrepancy are analysed. This is a crucial step towards developing accurate effective models to replace the FM- $\alpha$  model.

## 1 Introduction

The synthesis of ammonia is a pivotal chemical reaction for the world's food supply. To date, it has been dominantly driven by the Haber–Bosch process,<sup>1</sup> with the dissociative chemisorption of  $N_2$  on a catalyst's metal surface being the rate controlling step.<sup>2,3</sup> This elementary reaction step has been the subject of many studies – both experimental<sup>4–6</sup> and theoretical<sup>7–13</sup>—and

ruthenium has been found to be the optimal single element catalyst for the Haber–Bosch process.<sup>14–17</sup> Despite optimisations, this process's power consumption remains high even under optimal reaction conditions. Current efforts to improve the energy efficiency of the process involve applying plasma-enhanced catalysis. This approach aims to overcome the kinetic limitations of traditional thermal (heterogeneous) catalysis, potentially significantly reducing the operational temperature and pressure.<sup>18</sup> With reactants being excited in the plasma, numerous (synergistic) effects can increase reactivity, such as rovibrational and electronic excitation, ionisation and dissociation, modification of the catalyst surface through etching and charging, and the presence of electric and magnetic fields.<sup>19–21</sup> Moreover, the combination of a non-equilibrium plasma with a

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heterogeneous catalyst has been demonstrated to be able to surpass the sum of its parts.<sup>22,23</sup> However, plasma catalysis is also considerably more complex than heterogeneous catalysis. The aforementioned plasma-induced effects might all play a role in plasma catalysis, or some not at all, depending on the operating conditions, the catalyst, and the reactor design.<sup>20</sup> This makes it challenging to investigate and disentangle these effects through experiments alone. Therefore, simulations are necessary to elucidate the role of individual plasma-induced effects and thus enhance our understanding of the entire process.

Vibrational excitation of N<sub>2</sub> is known to increase the dissociative chemisorption (DC) rate.<sup>24–27</sup> The ground-breaking work of Mehta *et al.*<sup>28</sup> for plasma-enhanced ammonia production has made the DC of N<sub>2</sub> a prototypical reaction and an important showcase for the importance of vibrational excitations in plasma catalysis.<sup>20</sup> This notion has been reached based on simulations for the turnover frequency (TOF) for ammonia synthesis under realistic catalytic reaction conditions. The underlying microkinetic models (MKMs) rely on the Fridman-Macheret (FM)  $\alpha$  model<sup>29</sup> to quantify the effect of vibrational excitation on the reaction rate coefficient of the DC of N<sub>2</sub>.<sup>28,30</sup> The FM  $\alpha$  model is rooted in gas-phase reactions, where it models the vibrational contribution in diatom–atom reactions by computing the relative ratio of the forward and backward barrier heights. This model has become the workhorse approach for modelling vibrational excitation in plasma catalysis.<sup>31,32</sup>

Recently, questions have arisen regarding the relative importance of vibrational excitation compared to other plasma effects.<sup>21,33</sup> For instance, the concentration of vibrationally excited molecules is typically much lower compared to the concentration of radicals in dielectric barrier discharges, which are the plasmas most commonly used in plasma catalysis.<sup>21</sup> It is important to note that the FM  $\alpha$  parameter cannot be directly measured experimentally. Instead, it requires models fitted to experimental data or theoretical calculations to validate its value. While a simple gas-phase reaction may be reasonably well described by a model like TST and the FM  $\alpha$ , molecule–metal surface reactions typically involve significantly more complex potential energy surfaces (PESs) and necessitate high-dimensional models. The applicability of the FM  $\alpha$  model to these complex reactions has already been questioned by Kedalo *et al.*<sup>34</sup> for N<sub>2</sub> + Ru(0001) based on comparison with MD simulations. While their study should be considered with caution (see Section 3.3 and Section S2.2 in the ESI†), they found that the effect of vibrational excitation exceeds the FM  $\alpha$  model's predictions but is not as significant as the effect of translational energy. Recently, some of us have suggested that the FM  $\alpha$  approach performs poorly, both qualitatively and quantitatively, for DC in general. This is based on an extensive analysis of theoretical molecular dynamics (MD) studies available in the literature.<sup>35</sup> For instance, the effect of the vibrational excitation of polyatomic molecules on the reactivity depends on the specific vibrational mode(s) being excited.<sup>35–40</sup> For methane, excitation of the vibrational stretch modes is more effective than exciting the bend modes.<sup>36–38</sup> Similarly, exciting overtones generally leads to complex distributions of near-degenerate vibrational states that again yield considerably different reaction probabilities.<sup>39,40</sup>

Furthermore, the FM  $\alpha$  model seems to significantly underestimate the effect of vibrational excitation for a variety of DC reactions.<sup>35</sup> In late barrier systems, the bobsleigh effect can make molecules with high incidence energy deviate from the minimum energy path (MEP), and consequently they experience much higher barrier heights, lowering reactivity<sup>40–42</sup> (also often referred to as Polanyi's rules<sup>43</sup>). Additionally, sterical hindrance in the MEP combined with dynamical effects can have a major influence on the effect of rotational and vibrational excitation of reactants.<sup>44</sup> Unfortunately, the FM  $\alpha$  model is not able to capture these effects, since it only considers the forward and backward reaction barrier heights of the ground state PES and the absolute vibrational energy.

MD approaches are able to capture these complex effects of vibrational excitation, offering an intriguing alternative to the FM  $\alpha$  model.<sup>45</sup> In this study, we explore how and why the FM  $\alpha$  model and some simple extensions deviate considerably from full-scale MD simulations. For the latter, we build on our previous work that established an accurate MD model for N<sub>2</sub> dissociation on a Ru(0001) surface, meticulously validated against the gold standard in gas-surface dynamics provided by molecule-beam experiments.<sup>11,12,46</sup> This model employs quasi-classical trajectory (QCT) calculations using a machine-learned high-dimensional potential energy surface (PES) based on DFT calculations, including all relevant degrees of freedom of both the metal surface and the molecule. We utilise this model to compute the effect of vibrational excitation on the reaction probabilities and uncover a nontrivial relationship between vibrational excitation and reactivity. Crucially, we show that vibrational excitation plays a significantly more important role in plasma catalysis than previously anticipated, influencing the predicted TOF of ammonia in plasma catalysis. If the rate-controlling reaction step is highly activated, our findings suggest that the same can hold for conventional thermal catalysis. This could have serious implications for the modelling of heterogeneous catalysis; to the best of our knowledge, vibrational excitation is consistently neglected in MKMs for thermal catalysis. Finally, we analyse the reaction dynamics to elucidate the disparity in reactivity between conventional TST methods and our dynamics-based approach, concluding that dynamical effects play an important role in plasma catalysis.

## 2 Methods

### 2.1 Microkinetic model

Our study builds on the MKM originally developed by Mehta *et al.*<sup>28</sup> for modelling plasma-enabled catalysis. In this MKM, each vibrational level is treated as a distinct reactant, with its own specific reaction rate and a concentration directly proportional to the vibrational distribution. We consider three distributions for the vibrational states of the reactant N<sub>2</sub> molecules:

- (1) A 'ground-state-only' distribution ( $\nu = 0$ ), as typically used to model thermal catalysis.
- (2) A Boltzmann distribution to account for the population of vibrational states in thermal equilibrium at a given vibrational temperature  $T_{\text{vib}}$  that is the same as the gas temperature  $T_{\text{gas}}$ .



(3) A Treanor distribution<sup>47</sup>

$$f(\nu; T_{\text{vib}}, T_{\text{gas}}) = B \exp \left[ -\frac{\hbar\omega\nu}{T_{\text{vib}}} + \frac{x_c \hbar\omega\nu^2}{T_{\text{gas}}} \right], \quad (1)$$

to describe the non-equilibrium population of vibrational states, which is commonly used in plasma catalysis.<sup>28,48,49</sup> The Treanor probability density  $f(\nu; T_{\text{vib}}, T_{\text{gas}})$  is a function of the vibrational state  $\nu$  at given temperatures  $T_{\text{vib}}$  and  $T_{\text{gas}}$ , which usually are significantly different from each other.  $B$  is a normalization constant,  $\omega$  is the vibrational frequency, and  $x_c$  is the anharmonicity coefficient.

Throughout this work, we use this MKM to compute relative TOFs of ammonia synthesis, which are independent of pressure. This highlights the fact that different reaction rate coefficients for the  $\text{N}_2$  dissociation step cause the same relative differences for the TOFs for both the industrial high-pressure (100 atm for the Haber–Bosch process) and the typical low-pressure conditions in plasmas (1 atm). We assume a vibrational temperature of  $T_{\text{vib}} = 3000$  K for the plasma, like in the work of Mehta *et al.*<sup>28</sup> As mentioned in Section 1, other plasma-induced effects (*e.g.*, the presence of radicals) can also play a role, which we neglect here, in order to focus on the effect of vibrational excitation.

## 2.2 Reaction rate coefficients

**2.2.1 Transition state theory.** Usually, reaction rate coefficients in an MKM are obtained based on transition state theory (TST) in the form of an Arrhenius equation:

$$k = A \exp \left[ -\frac{E_a}{k_B T_{\text{gas}}} \right]. \quad (2)$$

In this equation,  $k_B$  is the Boltzmann constant,  $E_a$  is the activation barrier, including the effects of zero-point energy, and  $A$  is a frequency factor related to the entropy barrier and pressure. All activation barriers  $E_a$  in the work of Mehta *et al.*<sup>28</sup> are determined by scaling relationships (SR), *i.e.*, using fitted trends between adsorption energy and reaction barriers across many metal surface species.<sup>50</sup> Focussing on  $\text{N}_2$  dissociation in this work, we opt to use a more accurate activation barrier of  $E_a^{\text{TS}} = 1.83$  eV on the basis of a DFT-based transition state (TS) search<sup>11,51</sup> for this rate-limiting step. This is a substantially higher barrier than the  $E_a^{\text{SR}} = 1.48$  eV barrier from the SR. For all other elementary reaction steps, we keep the SR-based values as determined by Mehta *et al.*<sup>28</sup> for terraces on the Ru(0001) surface unchanged.

**2.2.2 Fridman–Macheret  $\alpha$  models.** To make the rate coefficient for  $\text{N}_2$  dissociation vibrational state specific, Mehta *et al.*<sup>28</sup> have suggested to use the Fridman–Macheret (FM)  $\alpha$  model, which has originally been developed for gas phase reactions. According to this model, reaction rate coefficients are obtained from a simple extension of the Arrhenius equation:

$$k_\nu = A \exp \left[ -\frac{E_a - \alpha E_\nu}{k_B T_{\text{gas}}} H(E_a - \alpha E_\nu) \right]. \quad (3)$$

Here,  $H$  is the Heaviside step function, and most importantly, the FM  $\alpha$  quantifies how effectively the vibrational

energy  $E_\nu$  (excluding zero-point energy) reduces the effective barrier height. The FM  $\alpha$  model aims to enforce the Polanyi rules<sup>43</sup>—which state that, in a late-barrier system, vibrational excitation has a greater impact on reactivity—by linking the ratio between the forward and backward reaction barriers to the “lateness” of the barrier and, consequently, to  $\alpha$  as follows:

$$\alpha = \frac{E_a^{\text{forward}}}{E_a^{\text{forward}} + E_a^{\text{backward}}}. \quad (4)$$

Here, the forward and backward reactions are the DC and desorption of  $\text{N}_2$ , respectively, and can be obtained with quantum chemistry calculations, such as DFT, either directly through a transition state (TS) search, or indirectly using the SR. Using scaling relationship activation energies results in an FM  $\alpha^{\text{SR}}$  of 0.37 and using the DFT activation energy for the forward dissociation reaction (keeping the adsorption energy equal) a slightly larger  $\alpha^{\text{DFT}} = 0.39$  is obtained. We refer to this rate calculation method as the ‘TST + FM’ level of theory, appending ‘@SR’ or ‘@TS’ when using the activation energy of  $\text{N}_2$  dissociation from scaling relationships or a DFT transition state search, respectively.

In the field of gas-surface dynamics, vibrational efficacies (VEs)

$$\eta_\nu(P) = \frac{E_{\text{inc}}^0(P) - E_{\text{inc}}^\nu(P)}{E_{\text{vib}}^\nu - E_{\text{vib}}^0} \quad (5)$$

can be considered the equivalent of  $\alpha$  in the FM  $\alpha$  model, *i.e.*, they quantify how much more efficiently vibrational energy  $E_{\text{vib}}^\nu$  increases the reaction probability  $P$  relative to the same amount of translational incidence energy  $E_{\text{inc}}^\nu$  (taken to be perpendicular to the surface) for a molecule in vibrational state  $\nu$ . While VEs have been measured for the dissociative chemisorption of a number of molecules,<sup>52,53</sup> no accurate data are available for  $\text{N}_2$ , since a suitable experimental technique to prepare molecular beams in specific vibrational states is yet to be found. To assess the accuracy of the FM  $\alpha$ , we have obtained a mean  $\bar{\eta}$  (detailed in Section 2.4 and Section S4 in the ESI†) using  $\text{N}_2$  dissociation probability curves  $P_\nu(E_{\text{inc}})$  as a function of incidence energy  $E_{\text{inc}}$  per vibrational state  $\nu$  calculated through MD simulations. This  $\bar{\eta}$  we use as a substitute to the FM  $\alpha$  in eqn (3) in order to compute rate coefficients. We refer to this method as the ‘TST +  $\bar{\eta}$ @TS’ level of theory in the following.

**2.2.3 Kinetic gas theory.** The last method for obtaining vibrational-state-specific rate coefficients forgoes the Arrhenius equation entirely by making use of the expressions derived from kinetic gas theory.<sup>54</sup> Instead of approximating the reactivity with a single barrier height, this requires reaction probability curves  $P_\nu(E_{\text{inc}})$  from MD to define rate coefficients:

$$k_\nu = A \langle P_\nu \rangle (T_{\text{gas}}), \quad (6)$$

where  $A$  is a frequency constant as in eqn (2) and (3), and  $\langle P_\nu \rangle (T_{\text{gas}})$  is the ensemble-averaged reaction probability at temperature  $T_{\text{gas}}$ , which is given by:

$$\langle P_\nu \rangle (T_{\text{gas}}) = \int_0^\infty f_{\text{inc}}(E_{\text{inc}}; T_{\text{gas}}) P_\nu(E_{\text{inc}}) dE_{\text{inc}}. \quad (7)$$



In this integral,  $f_{\text{inc}}(E_{\text{inc}}; T_{\text{gas}})$  is the probability density of the 1D Maxwell–Boltzmann distribution for the incidence energy  $E_{\text{inc}}$  according to the gas temperature  $T_{\text{gas}}$  given by:

$$f_{\text{inc}}(E_{\text{inc}}; T_{\text{gas}}) = \frac{1}{\sqrt{\pi k_{\text{B}} T_{\text{gas}} E_{\text{inc}}}} \exp\left[-\frac{E_{\text{inc}}}{k_{\text{B}} T_{\text{gas}}}\right]. \quad (8)$$

Here,  $E_{\text{inc}} = \frac{1}{2} m v_{\text{inc}}^2$ , where  $m$  is the molecular mass of an  $\text{N}_2$  molecule and  $v_{\text{inc}}$  is its center of mass's velocity component perpendicular to the surface. We calculate the  $\text{N}_2$  dissociation probability  $P_{\nu}(E_{\text{inc}})$  using MD simulations, as detailed in Section 2.3.3. Rate coefficients obtained with this method are referred to as the 'MD' level of theory in the following.

## 2.3 Molecular dynamics simulations

**2.3.1 High-dimensional neural network potential.** All MD simulations were performed using the high-dimensional neural network potential (HDNNP) constructed by Shakouri *et al.*<sup>11</sup> according to the Behler–Parinello approach.<sup>55</sup> Briefly, this potential is based on 25 000 single-point RPBE-DFT<sup>56</sup> calculations sampling the interaction of  $\text{N}_2$  molecules with the Ru(0001) surface, which is modelled by a slab of 7 layers using a  $3 \times 3$  supercell, where the bottom layer is kept fixed throughout all simulations. Of these 25 000 single points, 5000 were for a relaxed surface at 0 K surface temperature, and the remaining 20 000 were generated taking lattice expansion and surface atom displacements into account. In this work, the potential is evaluated in the LAMMPS code<sup>57</sup> using the ML-HDNNP package's interface to the n2p2 library.<sup>58,59</sup>

**2.3.2 Initial conditions.** The initial height of  $\text{N}_2$  above the surface is 5.5 Å. Because normal energy scaling has been observed for  $\text{N}_2$  dissociation probabilities on Ru(0001),<sup>60</sup> we only consider perpendicular incidence with a center of mass velocity  $v_{\text{inc}} = \sqrt{\frac{2E_{\text{inc}}}{m}}$ . The other initial conditions of the  $\text{N}_2$  molecule (*i.e.*, orientation, phase of the vibration, and the position relative to the unit cell) are sampled quasi-randomly according to a low-discrepancy sequence as defined in ref. 61. The same set of initial conditions ( $\theta$ ,  $\phi$ ,  $x$ ,  $y$ , and vibrational phase within its period) is used for all incidence energies and vibrational states. The surface temperature for this simulation was 575 K and we expanded the lattice constant of the slab according to experimental thermal expansion factors.<sup>62</sup> A Nosé–Hoover canonical ensemble (NVT) simulation using LAMMPS for a total of 150 ps with a 0.5 fs time step and a 50 fs damping time is used to extract 100 different snapshots of surface configurations with displaced atoms.

**2.3.3 Quasi-classical trajectory method.** The initial conditions are sampled separately by MD trajectories at different incidence energies according to the QCT approach.<sup>63</sup> In the QCT method, the  $\text{N}_2$  molecule is initialised with vibrational energy equal to the eigenenergies of the Hamiltonian of the free diatom. The vibrational states and corresponding energies are computed from the diatom potential *via* the Fourier grid Hamiltonian method.<sup>64</sup> Further details of the QCT method for  $\text{N}_2$  on Ru(0001) are provided in ref. 12. The dissociation probabilities were calculated for

incidence energies ranging from 0.25 to 10.0 eV in steps of 0.25 eV and for vibrational quantum numbers  $\nu \in \{0..10\}$ . We only simulate the rotational ground state  $J = 0$ , because previous work has shown that rotational excitation has limited to no effect on the reaction probability of  $\text{N}_2$  on Ru(0001).<sup>9</sup> The number of trajectories used for computing the reaction probabilities was  $10^4$  ( $10^5$ ) for total energies (*i.e.* incidence and vibrational energy, including the zero-point energy) larger (smaller) than 2.25 eV. All MD trajectories are calculated with the LAMMPS code using the standard NVE integrator with a time step of 0.33 fs.

**2.3.4 Reaction probabilities.** The reaction outcome of each MD trajectory is determined by the following criteria: the  $\text{N}_2$  molecule is considered to have reacted (dissociatively) if its bond length is larger than 2.65 Å. The molecule is considered to be scattered if it is further than 5 Å away from the surface with the centre-of-mass velocity vector pointing away from the surface. If neither occurs within the maximum simulation time of 36.3 ps, the molecule is considered to be trapped on the surface. Since most (99.9%) reactive trajectories occurred within the first 3.3 ps, we neglect the possibility that trapped molecules might react (dissociatively) after 36.3 ps and count them as non-reactive. We calculate  $\text{N}_2$  dissociation probability  $P_{\nu}(E_{\text{inc}})$  using an estimator:

$$\hat{p}_{\nu}(E_{\text{inc}}) = \frac{N_{\text{diss}}}{N_{\text{diss}} + N_{\text{scat}} + N_{\text{trap}}} = \frac{N_{\text{diss}}}{N_{\text{total}}}. \quad (9)$$

Here, the total number of trajectories is the sum of dissociated, scattered and trapped MD trajectories, respectively, at each  $E_{\text{inc}}$  and vibrational quantum number  $\nu$  (only explicitly denoted on the left-hand side of eqn (9) for brevity).

We have compared different approaches to numerically evaluate the integral in eqn (7) using the results from our MD calculations (see Section S2 in the ESI† for details). The most robust method uses the carefully constructed fitting function

$$S_{\nu}(E_{\text{inc}}) = \gamma_{\nu} \exp[-\beta_{\nu} E_{\text{inc}}^{-2\nu}]. \quad (10)$$

to obtain a continuous representation of  $\hat{p}_{\nu}(E_{\text{inc}})$ . Here,  $\alpha_{\nu}$ ,  $\beta_{\nu}$ , and  $\gamma_{\nu}$  are the optimal fitting parameters per vibrational state  $\nu$ , which we have obtained from a least-squares fit as tabulated in Section S2.1 of the ESI.†

The HDNNP was designed with an emphasis on configurations where the total initial energy of the  $\text{N}_2$  molecule (both translational and vibrational components) and the Ru surface does not exceed 15 eV. Consequently, the chance increases that the HDNNP needs to extrapolate towards parts of the PES that are less-well covered by the original DFT data set for MD trajectories at higher initial energies. We account for these uncertainties in  $\hat{p}_{\nu}(E_{\text{inc}})$  and propagate them in the form of 95% confidence bounds. Using the exact Clopper–Pearson method,<sup>65</sup> we compute the 95% confidence intervals for the estimators. For the lower bound, we include the number of non-conclusive trajectories terminated because of extrapolation  $N_{\text{extr}}$  in the total as if they are non-reactive ( $N_{\text{total}}^{\dagger} = N_{\text{diss}} + N_{\text{scat}} + N_{\text{trap}} + N_{\text{extr}}$ ), such that the lower bound is from the 95% confidence interval of a binom( $N_{\text{diss}}$ ,  $N_{\text{total}}^{\dagger}$ ) distribution. Likewise, for the upper bound, we also include the inconclusive extrapolative trajectories in the reactive trajectories



$N_{\text{diss}}^{\dagger} = N_{\text{diss}} + N_{\text{extr}}$ , such that the upper bound is from the 95% confidence interval of a binom( $N_{\text{diss}}^{\dagger}$ ,  $N_{\text{total}}^{\dagger}$ ) distribution. This allows both the uncertainty in the estimator due to finite sampling and the outcome of the trajectories due to the extrapolation of the HDNNP to be accounted for (see Fig. S1 in the ESI†). These errors are dominated by the former (latter) at low (high) incidence energies. Uncertainties in the resulting reaction rate coefficients  $k_{\nu}$  and TOFs from the MKM calculations are accounted for by distinct fits of eqn (10) to the lower bound and upper bound of the estimators, giving us effective 95% confidence intervals ( $p_{-}$  and  $p_{+}$ ) for a given fit parameter  $p$ . Additionally, we account for the uncertainty introduced by the least-squares procedure  $\sigma_{\pm}$  at the confidence bounds  $p_{\pm}$  by expanding the individual bounds as follows:

$$p'_{\pm} = \hat{p} + \text{sgn}(p_{\pm} - \hat{p}) \sqrt{(p_{\pm} - \hat{p})^2 + z\sigma_{\pm}} \quad (11)$$

Here,  $z = 1.96$  corresponds to 95% confidence intervals of the least-squares uncertainties and we assume that the fit uncertainty is independent of the input data uncertainty allowing us to sum the variances associated with the confidence intervals. This results in 3 sets of parameters for eqn (10) per vibrational state: a lower bound, an estimator and an upper bound (see Section S2.1 in the ESI†), which are used to compute the error bars for results at the 'MD' level of theory. Although the least-squares uncertainties are substantial for high vibrational quantum numbers, they ultimately hardly affect the uncertainties of the derived  $k_{\nu}$  and TOFs.

#### 2.4 Vibrational efficacy calculations

As described in Section 2.2.2, we have used a mean VE  $\bar{\eta}$  to substitute the FM  $\alpha$  in eqn (3) for vibrational state-specific rate coefficients. However,  $\eta_{\nu}(P)$  as defined in eqn (5) depends on the vibrational state-specific reaction probability and is not trivially reduced to a single scalar  $\bar{\eta}$ . In this section, we will discuss methods for computing  $\bar{\eta}$  as done in the literature and compare to our approach. One method is to pick a fixed reaction probability and take the mean over the vibrational quantum number such that

$$\bar{\eta}(P) = \frac{1}{N+1} \sum_{\nu=0}^N \eta_{\nu}(P). \quad (12)$$

This method considers only a single point of the reaction probability curves, which needs to be chosen carefully. In previous work for  $\text{N}_2$  dissociation on Ru(0001),  $P$  has been chosen between 0.1 and 0.5, resulting in  $\bar{\eta} \approx 1.6$  for  $\nu \leq 3$ .<sup>10,12</sup>

Here, we extend the original approach by aligning and incorporating the reaction probability curves for different vibrational states over a broader range of incidence energies. To achieve this, we express the numerator in eqn (5) as a sum of shifts between the successive vibrational states,  $E_{\text{inc}}^0 - E_{\text{inc}}^{\nu} = \Delta E_{\text{inc}}^{\nu} = \sum_{i=1}^{\nu} \delta E_{\text{inc}}^i$ . The  $\delta E_{\text{inc}}^i$  values are obtained as optimal shifts by minimising the total square difference integrals

$$D(\delta E_{\text{inc}}^{\nu}) = \int_{0.25 \text{ eV}}^{10 \text{ eV}} [P_{\nu}(E - \delta E_{\text{inc}}^{\nu}) - P_{\nu-1}(E)]^2 dE. \quad (13)$$

separately for each vibrational state  $\nu$ . The  $\nu$ -specific vibrational efficacies are then given by:

$$\eta_{\nu} = \frac{\Delta E_{\text{inc}}^{\nu}}{\Delta E_{\text{vib}}^{\nu}} = \frac{\sum_{i=1}^{\nu} \delta E_{\text{inc}}^i}{E_{\text{vib}}^{\nu} - E_{\text{vib}}^0}. \quad (14)$$

Finally, the mean value  $\bar{\eta}$  is obtained as the slope of a linear fit through the points ( $\Delta E_{\text{vib}}^{\nu}$  and  $\Delta E_{\text{inc}}^{\nu}$ ). For the details of this procedure and estimation of its uncertainty, see Section S4 in the ESI.†

## 3 Results: ammonia synthesis on Ru(0001)

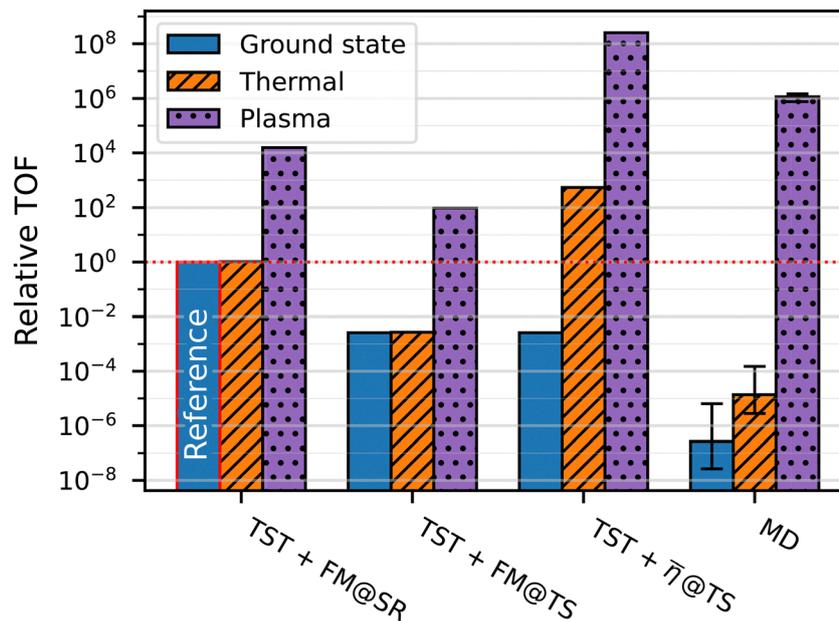
### 3.1 Turnover frequency for ammonia production

In Fig. 1, we present the relative TOFs of ammonia synthesis with the DC rate of  $\text{N}_2$  computed at the aforementioned separate levels of theory: TST + FM@SR, TST + FM@TS, TST +  $\bar{\eta}$ @TS, and MD defined in Section 2.2. Additionally, for each level of theory, TOFs are calculated for all three vibrational distributions (see Section 2.1), *i.e.*, a ground state ( $T_{\text{vib}} = 0$  K), thermal ( $T_{\text{vib}} = 673$  K) and plasma distribution ( $T_{\text{vib}} = 3000$  K). The TOFs presented in Fig. 1 are given relative to the reference ammonia TOF reported as 'plasma-off' results by Mehta *et al.*,<sup>28</sup> which, in our notation, corresponds to TST + FM@SR with a ground-state vibrational distribution. To ensure an adequate comparison, the vibrational temperatures and the gas temperature ( $T_{\text{gas}} = 673$  K) have been chosen identical to this previous work.

At this level of theory, Mehta *et al.*<sup>28</sup> demonstrated that the plasma vibrational distribution produces a significantly higher TOF—four orders of magnitude greater—compared to the ground-state-only result. Here, we show that the TOF increases by only 3% when a thermal distribution of  $\text{N}_2$  vibrational states is considered, as opposed to only the ground state. In other words, TST + FM@SR predicts that vibrationally excited species contribute minimally to ammonia production under thermal conditions—consistent with common expectations. As mentioned above, the first refinement of the reaction rate (*i.e.*, TST + FM@TS) involves computing the barrier height directly instead of relying on scaling relationships. In general, the TOFs for all vibrational distributions have decreased roughly by the same amount (almost 3 orders of magnitude), which is directly caused by the 0.35 eV increase in activation energy. Although this increase of  $E_{\text{a}}^{\text{forward}}$  also increases the FM  $\alpha$  ( $|E_{\text{a}}^{\text{forward}}| \gg E_{\text{a}}^{\text{backward}}$  in eqn (4)), the vibrational enhancement of the reactivity is negligible in the ground state and for thermal conditions, and only slightly larger under plasma conditions compared to TST + FM@SR.

At the next higher level of theory, TST +  $\bar{\eta}$ @TS, we observe a qualitative shift, with the vibrational ground state and thermal distributions now differing by more than 5 orders of magnitude. This suggests that, when the FM  $\alpha = 0.39$  is replaced by the much higher  $\bar{\eta} = 1.8$ , vibrationally excited species dominate the TOF in the MKM not only under plasma conditions but also under thermal conditions. Given the increased contributions of





**Fig. 1** Relative turnover rates for ammonia production on Ru(0001) under industrial reaction conditions for different levels of theory using the 'plasma-off' result from the study by Mehta *et al.*<sup>28</sup> as the reference. For each level of theory, we report the TOF for the vibrational ground state ( $T_{\text{vib}} = 0$  K), a thermal gas ( $T_{\text{vib}} = 673$  K) where the vibrational states are Boltzmann distributed, and a plasma ( $T_{\text{vib}} = 3000$  K) where the vibrational states are Treanor distributed.  $T_{\text{gas}}$  is 673 K in all three cases. The levels of theory are (i) SR activation energies with the FM  $\alpha$  model as described in ref. 28, (ii) DFT-calculated activation energy<sup>11,51</sup> with the FM  $\alpha$  model, (iii) DFT[calculated] activation energies with the mean VE  $\bar{\eta}$  from QCT and (iv) explicit reaction probabilities computed with MD using the QCT method.

vibrationally excited  $\text{N}_2$ , one might expect a larger relative difference in TOFs between thermal and plasma distributions. However, the effect remains similar, with the TOF increasing by 5 orders of magnitude compared to 4 orders of magnitude previously. In Section 3.2, we will show that, at higher vibrational levels, the reaction becomes effectively 'barrierless', leading to a saturation of the reaction rate.

Finally, at our highest level of theory—rate coefficients from MD—the difference between the ground state and thermal distributions is larger than in the FM  $\alpha$ -based results but remains considerably smaller than the  $\bar{\eta}$ -based results. Conversely, the gap between the plasma distribution and the ground state or the thermal distribution is immense, spanning 10 orders of magnitude—a difference unlikely to be observed in experiments. However, we argue that the MD simulations capture a realistic trend, as discussed in Section 3.2. Instead, the likely overestimation of the population of highly excited vibrational states by the Treanor distribution, as also reported in ref. 30, may account for this discrepancy. This underscores how the partitioning of energy can have a substantial impact on the predicted reaction rates in plasma catalysis. Finally, the ground state TOF at the MD level of theory is significantly lower than at the TST-based levels, indicating a higher effective activation energy for  $\text{N}_2$  dissociation. In Section 3.4, we will show that this effective increase arises from dynamical effects.

### 3.2 Analysis of $\text{N}_2$ dissociation probability

To better understand the overall ammonia TOFs computed *via* the MKMs, we analyse the DC of  $\text{N}_2$  on Ru(0001) in detail, as

this elementary step is known to have a pronounced effect for the conditions considered in the previous section ( $T_{\text{gas}} = 683$  K). Since the prefactor  $A$  is identical in eqn (3) and (6), relative differences between the reaction rate coefficients  $k_\nu$  are best compared by the thermally averaged reaction probabilities. In Fig. 2, we present  $\langle P_\nu \rangle$  as a function of the vibrational quantum number  $\nu$  for the four levels of theory previously introduced and discussed. This figure reasserts that the change in activation energy between the SR and the explicit TS only induces a shift in the  $\text{N}_2$  reaction probability. On a logarithmic scale, the exponent in eqn (3) simplifies to a linear dependence on  $E_\nu$ —consequently on  $\nu$  when anharmonicity is negligible. Therefore, the nearly identical slopes for TST + FM@SR and TST + FM@TS explain the negligible differences in the effect of vibrational excitation on the ammonia TOFs between these two levels of theory in Fig. 1.

When the FM  $\alpha$  is substituted with  $\bar{\eta}$ , the slope is notably increased, reaching  $\langle P_\nu \rangle = 1$  already for vibrational states  $\nu \geq 4$ . In other words, according to the FM +  $\bar{\eta}$ @TS level of theory,  $\text{N}_2$  DC on Ru(0001) is effectively barrierless and thus is guaranteed to occur starting from the 4<sup>th</sup> vibrationally excited state. For both TST + FM@SR and TST + FM@TS, the first vibrationally excited state  $\nu = 1$  is only one order of magnitude more reactive than  $\nu = 0$ , much less than the decrease in population between the two states for a thermal Boltzmann distribution. On the other hand, for TST +  $\bar{\eta}$ ,  $\nu = 1$  is four orders of magnitude more reactive, more than enough to allow this state dominate reactivity over the vibrational ground state despite its lower population. The individual contributions of each vibrational state to



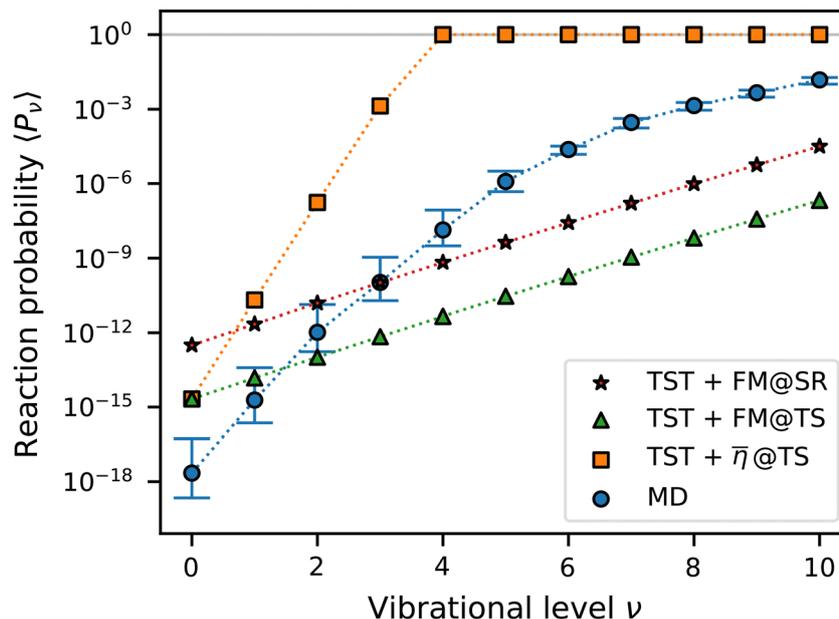


Fig. 2 Thermally averaged reaction probabilities per vibrational level at different levels of theory: (i) activation energies from the SR<sup>28</sup> with the FM  $\alpha$  model<sup>29</sup> (red stars), (ii) activation energy for N<sub>2</sub> dissociation from a DFT TS search<sup>11,51</sup> with the FM  $\alpha$  model (green triangles), (iii) the same activation energies with the mean VE from MD (orange squares), and (iv) reaction probability explicitly computed using MD.

the total reaction probability is further discussed in Section S3 of the ESI.† All of this explains why the increase in TOF from a thermal to a plasma distribution in Fig. 1 for TST +  $\bar{\eta}$ @TS is smaller than expected *a priori*: the reaction rate saturates at high vibrational quantum numbers, which is precisely where the plasma distribution differs most significantly from the thermal distribution in terms of vibrational state populations. Nevertheless, the substantial impact of the plasma on the ammonia TOF is maintained, as the total number of vibrationally excited species remains higher in the plasma distribution compared to the thermal one.

Finally, MD-based reaction probabilities are the only ones that strongly deviate from a linear dependence on the vibrational level. While they initially exhibit a slope similar to  $\bar{\eta}$ , this slope gradually decreases with increasing vibrational quantum number, appearing to asymptotically approach  $\langle P_\nu \rangle = 1$ . This change in slope explains the smaller difference between the TOFs for the ground state and thermal distribution, as only the first few vibrational quantum numbers contribute significantly to the N<sub>2</sub> dissociation rate (see Section S3 in the ESI†). The fact that the TOFs in the vibrational ground state at the TST + FM@TS and TST +  $\bar{\eta}$ @TS levels of theory are more than three orders of magnitude higher than those from MD is a direct consequence of the corresponding decrease in reaction probability at  $\nu = 0$  by a similar amount.

### 3.3 Vibrational efficacy

As outlined in Section 2.4, the mean vibrational efficacy, denoted by  $\bar{\eta}$ , is determined by aligning reaction probability curves for subsequent vibrational levels,  $\nu$ , and averaging over the resulting  $\eta_\nu$ . Notably,  $\nu$ -specific vibrational efficiencies,  $\eta_\nu$ , exhibit minimal variations with respect to  $\nu$  (as detailed in

Section S4 in the ESI†). This behaviour is similar to the FM  $\alpha$  parameter, which remains constant for each vibrational level. However, the mean vibrational efficacy calculated from our MD simulations,  $\bar{\eta} = 1.805$ , is significantly higher than the FM  $\alpha = 0.387$ .  $\bar{\eta} > 1$  implies that vibrational energy is more efficient than the same amount of translational energy in promoting the dissociation of N<sub>2</sub> on Ru(0001). In other words, vibrational excitation through plasma excitation is more efficient than thermal heating of the gas. This phenomenon cannot be captured by the FM  $\alpha$  model, which, by definition (see eqn (4)), restricts  $\alpha$  to the range of 0 to 1.

Going beyond the alignment of entire reaction curves, our molecular dynamics simulations enable us to delve deeper into the dependence of  $\eta_\nu$  on the incidence energy according to eqn (5). This analysis allows us to examine the vibrational efficiencies in greater detail. Fig. 3a illustrates that the  $\nu$ -specific vibrational efficacies are indeed significantly smaller at low incidence energies, exhibiting an increase with each vibrational level. Notably,  $\eta_\nu(E_{\text{inc}})$  rapidly increases until approximately  $E_{\text{inc}} \approx 3$  eV, where they attain a steady value of  $\eta_\nu \approx 2$  for each vibrational quantum number. This observation suggests that the thermally averaged state-specific vibrational efficacies

$$\langle \eta_\nu \rangle (T_{\text{gas}}) = \int_0^\infty f_{\text{inc}}(E_{\text{inc}}; T_{\text{gas}}) \eta_\nu(E_{\text{inc}}) dE_{\text{inc}} \quad (15)$$

strongly depend on the vibrational state. Fig. 3b reveals that the average efficacy  $\langle \eta_\nu \rangle$  varies from  $\langle \eta_{\nu=1} \rangle = 0.14$  to  $\langle \eta_{\nu=10} \rangle = 1.0$  for a thermal gas at 1000 K, exhibiting an almost linear increase with the vibrational quantum number. Notably, all these values are significantly lower than the mean efficacy  $\bar{\eta}$ , because the thermal averages are primarily influenced by lower incidence energy contributions (see the inset of Fig. 3a).



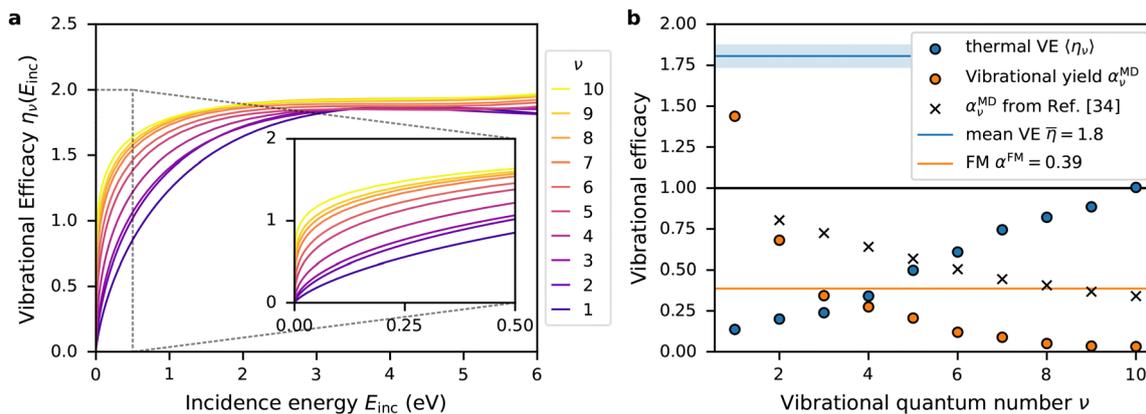


Fig. 3 (a) VE as a function of incidence energy per vibrational state, with the inset zooming into  $0 \text{ eV} \leq E_{\text{inc}} \leq 0.5 \text{ eV}$ . (b) Comparison of the vibrational state independent FM  $\alpha^{\text{FM}}$  (eqn (4), orange line) and mean VE  $\bar{\eta}$  (eqn (14), blue line), as well as the state-specific VE  $\langle \eta_{\nu} \rangle$  ( $T_{\text{gas}} = 1000 \text{ K}$ ) (eqn (15), blue circles) and MD-derived  $\alpha_{\nu}^{\text{MD}}$  ( $T_{\text{gas}} = 1000 \text{ K}$ ) (eqn (16), orange circles). Results for  $\alpha_{\nu}^{\text{MD}}$  ( $T_{\text{gas}} = 1000 \text{ K}$ ) as obtained by Kedalo *et al.*<sup>34</sup> are also included (black crosses).

Previously, Kedalo *et al.*<sup>34</sup> proposed an approach to extract FM  $\alpha$  values from the reaction probability curves generated through MD simulations. This involves solving for  $\alpha$  in eqn (3), which yields

$$\alpha_{\nu}^{\text{MD}}(T_{\text{gas}}) = \frac{\log\langle P_{\nu} \rangle - \log\langle P_{\nu-1} \rangle}{E_{\text{vib}}^{\nu} - E_{\text{vib}}^0} k_{\text{B}} T_{\text{gas}}. \quad (16)$$

We have included the results from ref. 34 in Fig. 3(b) (black crosses) for  $\alpha_{\nu}^{\text{MD}}$  ( $T_{\text{gas}} = 1000 \text{ K}$ ) for  $\text{N}_2$  dissociation on Ru(0001). In addition, we have also applied eqn (16) to the thermally averaged reaction probability obtained from our own MD simulations according to eqn (7) for the same temperature (orange circles in Fig. 3b). The results for  $\alpha_{\nu}^{\text{MD}}$  ( $T_{\text{gas}} = 1000 \text{ K}$ ) from the two different sets of calculations differ quite notably from each other. Most importantly, as  $\nu$  increases, our results do not approach  $\alpha^{\text{FM}}$  but rather converge to zero. We think that systematic shortcomings in the methodology of Kedalo *et al.*<sup>34</sup> are causing these differences: firstly, their barrier height for  $\text{N}_2$  dissociation on Ru(0001) computed with the PBE density functional (1.9 eV) is larger than our value underlying our HDNNP computed with the RPBE density functional (1.84 eV<sup>11</sup>). This is in large disagreement with the SBH17 database,<sup>51</sup> where the PBE and RPBE barrier heights have been computed as 1.49 eV and 1.95 eV, respectively. The values from the SBH17 follow the common trend that RPBE yields larger barrier heights for molecule–metal surface reactions than PBE. The small deviation of the aforementioned RPBE values (0.1 eV) is caused by slight differences in the DFT setup used in both studies, whereas the much larger difference between the PBE values (0.4 eV) is clearly beyond this error margin. Secondly, the root mean square error (RMSE) for the forces of the neural network fit for the potential energy surface reported by Kedalo *et al.*<sup>34</sup> (0.2–0.4 eV Å<sup>-1</sup>) is an order of magnitude larger than what is usual (0.02–0.03 eV Å<sup>-1</sup>).<sup>40</sup> Thirdly, in their MD simulations, Kedalo *et al.*<sup>34</sup> shift the Maxwell–Boltzmann distribution for the velocities to considerably higher energies, effectively assuming

that 80% of the total (translational) velocity of  $\text{N}_2$  molecules is distributed into the component perpendicular to the surface (see Fig. S4 in the ESI†). Furthermore, their MD simulation yields higher  $\text{N}_2$  dissociation probabilities for the flat Ru(0001) surface than that for the stepped Ru(113) surface—opposite to what has been obtained in other computational work.<sup>7,66,67</sup>

In light of these concerns regarding the DFT calculations, machine-learned potential and MD simulations of Kedalo *et al.*,<sup>34</sup> we now focus our comparison of  $\alpha_{\nu}^{\text{MD}}$  with  $\langle \eta_{\nu} \rangle$  on the results obtained from our own MD simulations (orange circles and blue squares in Fig. 3(b)). Starting from  $\alpha_1^{\text{MD}} \approx 1.4$ , there is a monotonous decrease with vibrational quantum number  $\nu$ , which still leaves  $\alpha_2^{\text{MD}}$  significantly larger than the FM  $\alpha$  given by eqn (4). The fundamentally opposite trends for  $\langle \eta_{\nu} \rangle$  and  $\alpha_{\nu}^{\text{MD}}$  with respect to  $\nu$  arise from the different order of operations. For  $\langle \eta_{\nu} \rangle$ , one first calculates the difference between the reaction probability curves and then integrates this difference over the velocity distribution. The definition of  $\langle \eta_{\nu} \rangle$  in eqn (15) consistently generalises the standard definition of vibrational efficacies used in gas-surface dynamics, completely analogous to  $\langle P_{\nu} \rangle$  in eqn (7). However, we do not consider  $\langle \eta_{\nu} \rangle$  as a suitable substitute to introduce  $\nu$ -dependence to the FM  $\alpha$ , as this would effectively double-count statistical effects in the TST and MD descriptions. On the other hand, for  $\alpha_{\nu}^{\text{MD}}$ , one first integrates  $P_{\nu}(E_{\text{inc}})$  over the Maxwell–Boltzmann distribution and then calculates the difference between the thermally averaged reaction probabilities  $\langle P_{\nu} \rangle$ . Consequently, vibrational-state-related effects are not blended with a combination of (different) statistical treatments at the TST and MD levels. By construction, if the results for  $\alpha_{\nu}^{\text{MD}}$  were to be substituted for the constant  $\bar{\eta}$  in eqn (3), the resulting TOFs would closely align with the MD reference results shown in Fig. 1. Clearly,  $\alpha_1^{\text{MD}}(1000 \text{ K}) > 1$  is the key enhancement that cannot be achieved with the FM  $\alpha$  model.

Altogether, we present compelling evidence that capturing the impact of vibrational excitation on the reaction rate coefficient using the FM  $\alpha$  approach, which currently stands as the



state-of-the-art in plasma catalysis modelling, presents substantial challenges. According to the Polanyi rules,<sup>43</sup> this effect can be estimated based on the potential energy surface. Early (late) barriers typically exhibit low (high) vibrational enhancement. Therefore, we anticipate similar implications for the efficiency of plasma excitation for other highly activated late barrier reactions, such as CH<sub>4</sub>,<sup>40,68,69</sup> HCl,<sup>44,70</sup> and CH<sub>3</sub>OH,<sup>71,72</sup> where high VEs have also been observed.

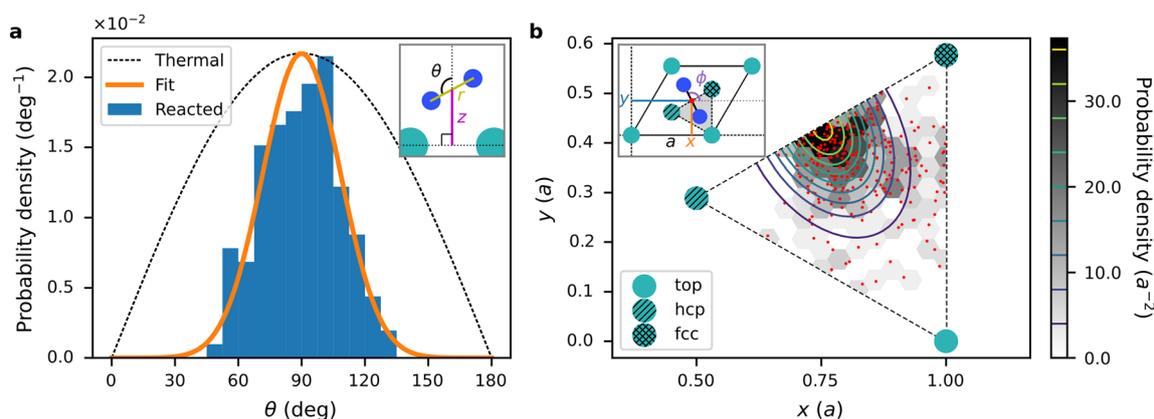
### 3.4 Dynamical effects

We continue our analysis by focusing on the TOFs for the vibrational ground state as presented in Fig. 2. The results based on reaction probabilities obtained from our MD simulations are significantly lower than those based on TST-based methods. This difference is likely caused by dynamical effects that are not accounted in any of the latter approaches. The conventional formulation of TST, as employed here, primarily considers the minimum energy path (MEP) for the dissociation of N<sub>2</sub>, with approximate vibrational corrections. This MEP is defined by a single reaction coordinate, which, in the simplest case for DC, encompasses only the bond length and the molecule–surface distance. In contrast, our MD simulations include a much higher number of degrees of freedom: 6 for the N<sub>2</sub> molecule and 162 resulting from the 54 moving atoms in the slab model for the Ru(0001) surface. Their dynamical interplay can facilitate dissociation along paths that differ substantially from the MEP.

To analyse these dynamical effects, we extracted statistics from MD trajectories of dissociated N<sub>2</sub> molecules in the vibrational ground state at a high incidence energy ( $E_{\text{inc}} = 3.25$  eV) above the barrier along the MEP on our PES ( $E_{\text{a}}^{\text{TS}} = 1.83$  eV). Even at such a high incidence energy the dissociation probability obtained from the MD simulations is less than 100% (see Section S1, ESI†). In other words, not all molecules are efficiently steered towards the MEP and can thus encounter barriers higher than  $E_{\text{a}}^{\text{TS}}$ . In our statistical analysis, we focus on

the coordinates depicted in the insets of Fig. 4 that are traditionally not included in the reaction coordinate: the angular orientation of the molecular axis of N<sub>2</sub> by the polar angle  $\theta$  and its lateral position ( $x$  and  $y$ ) relative to the surface unit cell, all based on the N<sub>2</sub> centre of mass as the origin. As described in Section 2.3.2, our trajectory ensemble is initially uniformly distributed over both orientations and lateral positions. According to TST, all reactive trajectories will cross the same minimal energy barrier along the MEP, which implies that there should be no preferential initial conditions. In Fig. 4, we show the distributions for the initial orientation  $\theta$  and  $xy$  position of the subset of dissociated N<sub>2</sub> molecules, which have been rescaled such that their maxima coincide with those of the distributions for all trajectories. A sharper distribution for reacted trajectories indicates a degree of inaccessibility of the MEP for initial conditions further from the MEP, leading to a lower reaction probability compared to TST. By computing the ratio between the integrals over these (rescaled) distributions and their original counterparts for all trajectories, we can estimate the reduction in reactivity in particular degrees of freedom due to dynamical effects.

As shown in Fig. 4a and b, respectively, both the  $\theta$  and  $xy$  distributions for the reacted molecules deviate considerably from the uniform distribution of all molecules. The very narrow distributions in Fig. 4(b) illustrate that dissociation dominantly occurs for molecules that are initially placed over the bridge site, in agreement with previous findings.<sup>10,11</sup> This significant narrowing of the distribution for  $xy$  due to surface corrugation leads to a substantial reduction of reactivity by more than a factor of 100. Fig. 4a shows that the distribution for  $\theta$  narrows similarly for the reacted molecules, but the reactivity reduction in this case is only a factor of 2–3 $\times$ . In the ESI,† we show that the effect of the azimuthal orientation of the molecular axis on the reactivity is even smaller (see Fig. S7, ESI†). There we also demonstrate that  $xy$ -related reactivity reduction strongly depends



**Fig. 4** Distributions for initial (a) orientation  $\theta$  and (b)  $xy$  position of N<sub>2</sub> molecules above the Ru(0001) surface for  $E_{\text{inc}} = 3.25$  eV ( $> E_{\text{a}}^{\text{TS}} = 1.83$  eV) and  $\nu = 0$ . The coordinates are depicted in the insets. (a) The initial distribution for all trajectories (both reacted and non-reacted, black dashed line, eqn (S6) in the ESI†) is compared to the distribution of the subset comprised by the reacted trajectories (blue bars). To guide the eye, the latter has been fitted to the initial distribution multiplied by a Gaussian centred at  $\theta = 90^\circ$  (orange line, eqn (S7) in the ESI†). (b) The distribution for the initial  $xy$ -position of reacted trajectories in the irreducible wedge of the unit cell, as depicted in the inset.  $a$  denotes the surface lattice constant. Individual reacted trajectories are in red, the histogram is in greyscale, and the fitted bimodal Gaussian distribution is the contour lines. For better comparison and to quantify dynamical steering effects, the reacted distributions have been rescaled such that their maxima coincide with those of the distributions for all trajectories.



on the incidence energy and ranges from a  $50\times$  at low incidence energy to a  $10\times$  reduction at high incidence energy.

Altogether, assuming that the aforementioned angular and translational degrees of freedom are independent, the total reduction in reactivity due to dynamical effects is a factor of 20–250 $\times$  compared to TST, depending on the incidence energy (see Section S5 in the ESI†). This accounts for more than half of the  $10^3$ – $10^4\times$  reduction in the reaction probability shown in Fig. 2. Recrossing events in our MD simulations, *i.e.*, molecules which do not dissociate despite crossing the barrier, could be responsible for a further reduction of the reaction probability compared to TST. Our recrossing analysis reveals that this is not the case (see Section S6 in the ESI†). Consequently, the bobsleigh effect is very likely responsible for the remaining reduction in reactivity due to dynamical effects, since it is closely related to the movement along the remaining molecular degree of freedom, the *z*-coordinate.<sup>40–42</sup> In principle, MEP deviations and barrier recrossing could be accounted for in TST by including entropic contributions and transmission factors in the reaction rate coefficients for the vibrational ground. However, accounting for these effects requires additional computational efforts, which is why we have not further investigated this here.

## 4 Discussion: implications and recommendations for thermal catalysis and plasma catalysis

For thermal catalysis, vibrational excitation of reactants is generally not considered to significantly affect the results. However, in this work, we predict that highly activated thermal catalysis (*e.g.*, ammonia synthesis) is indeed significantly affected by vibrationally excited states, even if their population is comparatively low. Dynamical effects, such as the bobsleigh effect, considerably limit the effectiveness of translational energy, whereas vibrational energy is not affected as much. Therefore, for a chemical reaction where the TS is characterised by a large barrier height and a very elongated bond, we suspect that simulating only the vibrational ground state is insufficient. Interestingly, dynamical effects can also influence other aspects of the reaction. For instance, moving surface atoms can alter reactivity locally.<sup>73,74</sup>

For plasma catalysis, we found that the effects of vibrational excitation are much more pronounced than previously anticipated based on the basis of the FM  $\alpha$  model. Some of us have recently observed that, in general, the vibrational efficacies for DC reactions scale with the absolute barrier height.<sup>35</sup> In comparison to other metal surfaces, Ru(0001) has a relatively low barrier height for the DC of N<sub>2</sub>. Other metals thus likely exhibit even higher vibrational efficacies than described by the FM  $\alpha$  model. Consequently, the predicted optimal catalyst and operating conditions might change considerably, depending on the level of theory employed to account for the vibrational enhancement of DC reactions.

Similarly, it is often unclear which kind of plasma-induced effects are dominant.<sup>20</sup> Using different plasma catalytic models

that try to disentangle the various effects might lead to qualitatively different conclusions. For example, the Langmuir–Rideal reaction mechanism (also, incorrectly, known as Eley–Rideal<sup>75</sup>), where a reaction takes place directly between a gas phase reactant (in plasma catalysis typically a radical) and a surface adsorbate, is often studied along vibrationally excited reactants. For the plasma-catalytic non-oxidative coupling of methane, it has been predicted that the importance of radicals and vibrational excitation of reactants depends on the binding strength of the catalyst surface.<sup>30,76,77</sup> Similarly, for CO<sub>2</sub> hydrogenation, radicals were predicted to be much more important.<sup>31</sup> In a kinetic investigation of ammonia synthesis on Fe, adsorption of radicals, Langmuir–Rideal reactions, and vibrational excitation were all found to be important.<sup>73</sup> Bayer *et al.*,<sup>33</sup> on the other hand, concluded on the basis of a combination of simulations and experiments that the loss of vibrationally excited N<sub>2</sub> on Fe catalysts was due to vibrational relaxation (*i.e.*, non-reactive energy loss), and not DC. Furthermore, for dry reforming of methane on Ni, it was predicted that the Langmuir–Rideal reaction mechanism dominates, compared to vibrational excitation.<sup>32</sup> Generally, it has been concluded that Langmuir–Rideal reactions and adsorption of radicals are much more important than vibrational excitation of reactants. However, in all cases, the FM  $\alpha$  approach was employed. Based on our work, it can be expected that vibrational excitation plays a significantly more important role than has previously been accounted for by catalytic models. Furthermore, it is likely that the competing Langmuir–Rideal reaction mechanism occurs less often than typically modelled, especially when dynamical effects are considered.<sup>78</sup> This could have positive implications for the practical application of plasma catalysis. As discussed above, it is often believed that radicals are the driving force behind the enhanced chemical reactivity. However, designing a catalyst and a reactor in such a way that radicals are not ‘quenched’, which reduces the availability of radicals for desired reaction steps, is not facile. But if vibrational excitation is much more important than thus far thought, practical plasma catalysis might be more easily achieved. In short, the way dynamical effects are modelled in plasma catalysis can alter predictions considerably.

In this work, we have also observed that scaling relationships can lead to a considerable underestimation of the barrier height, necessitating explicit TS search calculations instead. Similarly, it can be expected that, for the majority of catalytically important reactions, standard DFT calculations will underestimate the barrier height, even when the TS is directly computed instead of employing the SR.<sup>79</sup> Self-interaction errors are likely the main culprit and are most prominent when the charge transfer at the TS is large.<sup>80,81</sup> Future improvements in electronic structure calculations are crucial for obtaining more reliable PESs, including barrier heights for MKMs. Fortunately, for N<sub>2</sub> + Ru(0001), the employed PES yields very good agreement with currently available experimental data from well-defined molecular beam experiments.<sup>11,46</sup>

We also acknowledge that the employed QCT method may have certain limitations, such as neglecting nuclear quantum effects. Although tunnelling is likely insignificant for a molecule



as heavy as N<sub>2</sub>, zero-point energy violation does occur in our simulations (see Fig. S5 in the ESI†), leading to a slight overestimation of the reaction rate of lower vibrational states. In the future, by substituting the QCT approach with ring polymer MD, nuclear quantum effects can be simulated approximately, while keeping the simulations tractable.<sup>82</sup> For the DC of methane on a moving Pt(111) surface, this approach has been demonstrated to accurately reproduce results from experiments, even for low incidence energies and high vibrational temperatures.<sup>83</sup> However, to the best of our knowledge, ring polymer MD simulations cannot yet be performed for specific vibrational states, necessitating methodological advancements to simulate vibrational state populations found in plasma catalysis.

Despite all methodological challenges, dynamical effects remain a crucial important aspect in determining reaction rates of molecules on metal surfaces. For instance, in the case of heterogeneous catalytic cracking of ammonia on Fe surfaces, which is the reverse of the Haber–Bosch process investigated in this work, dynamical effects play a significant role not only in the diffusion coefficient of ammonia but also in reaction steps involving surface adsorbates.<sup>84</sup> Therefore, we conclude that dynamical effects cannot be ignored in determining the reaction rates of molecules on metal surfaces. In the context of plasma catalysis, it appears that dynamical effects are even more important than in thermal catalysis, which unfortunately cannot be fully captured by TST models (at present). Accurate MD calculations for gas-surface dynamics are challenging to perform, as evidenced by the discrepancies between the present work and the work by Kedalo *et al.*<sup>34</sup> (see Section 3.3 and Section S2.2 in the ESI†). In the meantime, if available in the literature, vibrational state-resolved reaction probabilities measured or calculated for molecular beams can be utilised to guide catalytic modelling. Unfortunately, accurate data are scarce, particularly for higher vibrational states and very low reaction probabilities.

Accurate determination of the vibrational enhancement of DC reactions, as exemplified in this work, requires the computation of numerous reaction probability curves, each requiring many trajectories. Therefore, the remaining discrepancy between the most advanced TST model and the MD simulations for N<sub>2</sub> dissociation on Ru(0001) prompts further research to develop computationally efficient alternatives to account for the impact of vibrational excitation and dynamical effects on reaction rate coefficients. Although the mean VE does improve the FM  $\alpha$  model by reproducing at least the initial trend of the MD results, there is room for improvement as the mean VE does overestimate the effect of vibrational excitations. The sudden vector projection model<sup>85</sup> is an alternative to performing full dimensional MD simulations, coming at a much lower computational cost. Unfortunately, this model still struggles with the same assumption as the FM  $\alpha$  model that vibrational energy cannot be more effective than translational energy in promoting reactivity, and that dynamical effects such as the bobsleigh effect are not explicitly accounted for. Nevertheless, some of us have recently shown that a combination of the absolute barrier height, equilibrium and transition state bond lengths, and the sudden vector projection model manages to reproduce VEs of various molecule–metal

surface reactions reasonably well at a similar computational cost as the FM  $\alpha$  model.<sup>35</sup> Furthermore, it was observed that the curvature features of the MEP could significantly enhance the prediction of VEs with simplified static models.

## 5 Conclusions

In this study, we explored the impact of vibrational excitation on reactants in heterogeneous and plasma catalysis, and how this effect influences predictions obtained with microkinetic models. Our work marks a significant milestone in scrutinising the influence of vibrational effects on reaction rate coefficients for strongly activated dissociative chemisorption, which can be strongly driven by dynamical effects. Specifically, we computed vibrational-state-specific reaction rate coefficients for the dissociative chemisorption of highly vibrationally excited N<sub>2</sub> on Ru(0001), which is the rate-limiting step in ammonia synthesis. These reaction rate coefficients were obtained using molecular dynamics simulations at a level that is the state of the art in the field of gas-surface dynamics. They were subsequently employed in a MKM for ammonia synthesis on Ru(0001). The MD simulations revealed orders of magnitude lower reaction rate coefficients for N<sub>2</sub> and concomitant TOFs for ammonia compared to the reaction rate coefficients from transition-state-theory, which are the current workhorse for MKMs. Going beyond scaling relationships in the practical implementation of TST brings the results closer to the MD-based reference results, which strongly encourages investment of the additional computational effort to calculate transition states and their corresponding energy barriers, as this is still much easier than performing MD simulations.

Even under typical thermal catalytic conditions, the small fraction of vibrationally excited N<sub>2</sub> molecules dominate the rate coefficient for N<sub>2</sub> dissociation due to the very high vibrational efficacy. Consequently, typical MKMs for thermal catalysis, which only consider the vibrational ground state, prove insufficient for a highly activated chemical process like ammonia synthesis. Most importantly, we predict that the reactivity gap between thermal catalysis and plasma catalysis due to vibrational excitation is substantially larger than previously anticipated. MKMs used in plasma catalysis do commonly account for vibrational excitation with the FM  $\alpha$  approach, which unfortunately leads to a large underestimation of the reactivity of vibrationally excited N<sub>2</sub>. Employing vibrational efficacies determined from MD simulations yields a much larger impact of vibrational excitation on the reaction rate and improves the agreement between the TST model and the MD approach. Complex features from the PES, such as corrugation and anisotropy, along with dynamical effects like the bobsleigh effect, are responsible for the remaining discrepancy between the two approaches. It remains to be seen if a simple and computationally efficient modification of the Arrhenius equation can account for these effects. Our results emphasise the significance of how energy is partitioned and modelled in a chemical system, as it can significantly alter predictions. Since the energy distributions in plasma catalysis



can be highly out of equilibrium, it is probable that dynamical effects play a substantial role not only in the DC of vibrationally excited reactants but also in other plasma-induced effects (e.g., Langmuir–Rideal reactions). For the time being, we conclude that the gold standard for determining the reaction rate coefficients for DC reactions in heterogeneous catalysis and plasma catalysis involves performing dynamical simulations.

## Author contributions

FvdB: data curation, formal analysis, investigation, methodology, software, validation, visualisation, and writing – original draft preparation. NG: project administration, supervision, and writing – review and editing. JM: conceptualization, funding acquisition, resources, supervision, and writing – review and editing.

## Conflicts of interest

There are no conflicts to declare.

## Data availability

MD input files and trajectory outcomes including uncertainty analysis and scripts needed to reproduce all results from this article and the ESI<sup>†</sup> are available at zenodo.org (<https://doi.org/10.5281/zenodo.15006139>). Additional data can be obtained from the authors upon reasonable request.

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