


 Cite this: *EES Sol.*, 2025, 1, 1061

Stabilizing perovskite solar cells at 85 °C via additive engineering and MXene interlayers

Fanny Baumann, Jessica Padilla-Pantoja, Jose Manuel Caicedo-Roque, Masoud Karimipour, Naji Vahedigharehchopogh, Jose Santiso, Belén Ballesteros, Ramses A. Miranda Gamboa, Zhenchuan Tian, Sonia R. Raga * and Monica Lira-Cantú *

The commercialization of Perovskite Solar Cells (PSCs) is currently a reality. The detection of early instabilities, especially at high temperatures is vital for the successful widespread implementation of the technology. Commercial devices must sustain temperatures as high as 85 °C in order to surpass standardized tests and certification for the final PV product. However, *in situ* and *operando* stability analysis and detailed structural and electronic properties of full devices are still rarely found in the literature. In this work, we carried out *in situ* operational stability testing at 85 °C, complemented by *in situ* X-ray diffraction, impedance spectroscopy, photoluminescence, current–voltage measurements and electron microscopy. Our results demonstrate a large lattice expansion in the halide perovskite which provokes a clear voltage drop in the PSC. While no changes in lattice constants were observed over time at 85 °C, we noticed a reversible formation of an amorphous “carbon rich” surface shell material surrounding the perovskite grains. This material is linked to a decrease in shunt resistance, and the increase of ionic conductivity. The latter triggered the gradual photovoltaic performance loss observed in our PSC at high temperature. Additionally, we demonstrate the possibility to delay this PSC degradation by employing stability-enhancing methods such as additive engineering and the application of functionalized 2D Ti₃C₂ MXene interlayers to the PSC. Our work showcases the value of complementing stability tests with advanced characterization, significantly showcasing the value of *operando* structural studies.

 Received 28th June 2025
 Accepted 19th September 2025

DOI: 10.1039/d5el00104h

rsc.li/EESolar

Broader context

Amid today's urgent push for sustainable energy, halide perovskite solar cells (PSC) have emerged as a promising renewable technology. Despite significant progress, PSC devices still struggle to exceed a one-year lifespan, falling short of industrial standards and limiting commercialization. Rapid detection of early degradation through accelerated testing over just a few hours can save time and resources, advancing PSC development more efficiently. The variety of dynamic processes and complex degradation pathways in PSC under operation conditions can only be understood through *in situ* characterization. This study investigates the evolution of HP crystal lattice under applied bias, illumination and high temperature with a customized *in situ* XRD setup. We show that targeted design modifications in the PSC can prevent light- and heat-induced perovskite amorphisation into a phase that creates new shunt paths and increases carrier losses. This work highlights the need for strategies that reinforces the crystal structure by targeting grain boundaries and mitigates thermal expansion mismatches through interfacial engineering.

Introduction

Before perovskite solar cells (PSC) can satisfy industrial needs with warranty for acceptable performance,^{1,2} it is fundamental to mitigate photovoltaic loss under high temperature and variable light-irradiance conditions experienced under real operation.^{3,4} Photovoltaic modules heat up to an average of ~45–50 °C and up to ~65 °C under operation.^{5,6} At extreme conditions,

such as space stations or satellites, both radiation and temperature levels are more intense, with temperatures up to 85 °C measured.⁷ 65 °C and 85 °C are the temperature thresholds defined in the ISOS consensus protocols for PSCs stability,⁸ and the IEC 61215 standard.^{9,10}

Usually, high temperatures are reached at high irradiation, yet empirical calculations suggest that the internal heating of the perovskite layer might be higher than the measured module temperature.¹⁰ Operational testing at illumination and elevated temperature might be more predictive of outdoor stability.^{11,12} Both light and temperature provide excess energy into the PSC system, which can destabilize the halide perovskite (HP)

Catalan Institute of Nanoscience and Nanotechnology (ICN2), CSIC and the Barcelona Institute of Science and Technology (BIST), Building ICN2, Campus UAB, Bellaterra, E-08193, Barcelona, Spain. E-mail: sonia.ruizraga@icn2.cat; monica.lira@icn2.cat



structure leading to material and device degradation. The extra energy in the PSC system may produce damaging structural effects in the HP,¹³ decrease ionic migration activation energies, and induce sequential degradation processes from intrinsic changes of the HP.^{14,15} However, most studies on temperature-induced degradation only report local material changes in isolated single crystals or films, eluding the interactions of the selective layers, applied bias or photogeneration in the degradation, while the effects on device performance have been observed in separate experiments with post-mortem investigations eluding reversible phenomena. Here, we join local and general effects, to understand the initial degradation pathways of full PSC devices. In addition, the interplay between varying light intensity and temperature is not frequently reported for full devices but it is critical to get an accurate picture of mechanisms triggering degradation, particularly relevant for outdoor applications.^{16–18}

Successful strategies to mitigate PSC degradation at higher temperatures and illumination include compositional tuning, interlayers to avoid delamination, interfacial engineering, and molecular additives in the bulk.^{4,19–22} Additives in the HP solution can induce changes in the film crystal orientation and relaxation of internal strain in the crystal grains.^{20–23} Moreover, additives that incorporate in the bulk HP can lead to formation of 2D/3D structures. Recently, the use of additives that reside only at grain boundaries and surfaces is becoming more frequent.^{21,24,25} The phosphonate additive 3-phosphonopropionic acid (H3PP) had proven to passivate shallow defects and provoke a small compression in the perovskite lattice, with the consequence of improved stability at high illumination.^{21,23} Similarly, interlayers including MXene-functionalized with the H3PP additive had proven to greatly improve stability of PSCs under outdoor testing.²⁶ However, the effect of these additions had not been analyzed at higher temperatures.

In this work, we investigated full PSCs devices under constant light irradiation of $\sim 100 \text{ mW cm}^{-2}$ (1 sun), in N_2 and at $85 \text{ }^\circ\text{C}$ by operational tracking of maximum power point (MPPT) and with *in situ* X-ray diffraction (XRD) while under bias near MPP. We compared state-of-the-art triple cation-based PSCs (REF) to modified devices (MOD) with additive engineering in the HP precursor and functionalized-MXene interfacial layer between the HP and spiro-OMeTAD. The HP and the MXene were both modified by adding the organic molecule 3-phosphonopropionic acid (H3PP).

MPP operational stability testing was complemented with characterization including *in situ* X-ray diffraction (XRD) analyses of full devices during operational conditions, electrochemical impedance spectroscopy (EIS) at varying light intensities, photoluminescence (PL), current density–voltage (J – V), and transmission electron microscopy (TEM). This combination of advanced methods allowed us to identify simultaneous changes in lattice constants and electrical output at high temperature and illumination while correlating structural and performance changes with the observed electrochemical properties of the devices. We show how reversible stages of degradation in full devices at high temperatures induce structural transformations that facilitate inter-grain amorphization

without formation of PbI_2 or lattice parameter evolution over time. This contrasts with previous works where PbI_2 is reportedly the end product of irreversible thermal degradation of perovskite crystals. Our results also show a photocarrier loss mechanism causing the MPP decay that is attributed to electrical shunt paths through the amorphous material originated at grain boundaries.

This work is, to our knowledge, the first study where structural dynamics by XRD and device performance is monitored jointly in full PSC devices under thermal conditions of $85 \text{ }^\circ\text{C}$ and illumination stress.

Results and discussion

Operational stability at $85 \text{ }^\circ\text{C}$ with detailed characterization before and after (B & A)

Devices investigated in this work were produced with a standard recipe and with stability enhancing modifications for comparison. We produced control n–i–p solar cells with precursor composition $\text{Cs}_{0.08}\text{MA}_{0.12}\text{FA}_{0.8} \text{Pb}(\text{I}_{0.88}\text{Br}_{0.12})_3$ (FA: formamidinium, MA: methylammonium) and device stack FTO/ TiO_2 -c/ TiO_2 -mp/HP/spiro-OMeTAD/Au (hereafter REF). Modified devices (hereafter MOD) incorporate the H3PP additive in the perovskite precursor and an interlayer of delaminated 2D MXene Ti_3C_2 functionalized with H3PP between the perovskite and the HTM, having a final architecture: FTO/ TiO_2 -c/ TiO_2 -mp/H3PP:HP/H3PP:MXene/spiroOMeTAD/Au (Fig. 1a), with the corresponding champion photovoltaic performance shown in Fig. 1b. Operational testing under 100 mW cm^{-2} spectral-matching LED illumination and at $85 \text{ }^\circ\text{C}$ under maximum power point (MPP) operational stability test (ISOS-L2) with N_2 flow was endured for more than 20 h. Fig. 1c shows the set up employed for stability analysis, and Fig. 1d the home-made setup employed for *in situ* characterization.

The power conversion efficiency (PCE) evolution at $85 \text{ }^\circ\text{C}$ under illumination in Fig. 2b and the efficiencies before and after operational testing (Fig. 2a and c, respectively) show that the modification with H3PP and MXene interlayer successfully improved stability in MOD PSCs, in accordance with previous observations at lower temperatures.^{21,26} At these severer conditions, all PSCs of both varieties lose a significant part of their initial function, with MOD PSCs showing a more stabilized performance once at high temperature.

Fig. 2d and e shows the corresponding J – V curves before and after MPPT testing, together with the V_{MPP} – J_{MPP} evolution (displayed as gradient-colored dots depending on time), hereafter MPP wandering plot. The MPP wandering plot lets us observe a different slope in the V_{MPP} – J_{MPP} changes over time; once the REF sample was operated at higher temperature the MPP voltage dropped progressively from 0.8 V to 0.7 V, before suddenly falling further to 0.6 V in the after J – V -scan (Fig. 2d), while MOD devices showed a 0.05 V initial improvement in V_{MPP} during the first hours of operation, followed by small accompanied voltage and current losses. V_{MPP} in REF decayed 0.1 V more (from 0.8 V to 0.6 V) compared to in MOD, where V_{MPP} even improved 0.05 V at the start of operation, followed by small voltage and current losses. We suggest this strategic way of



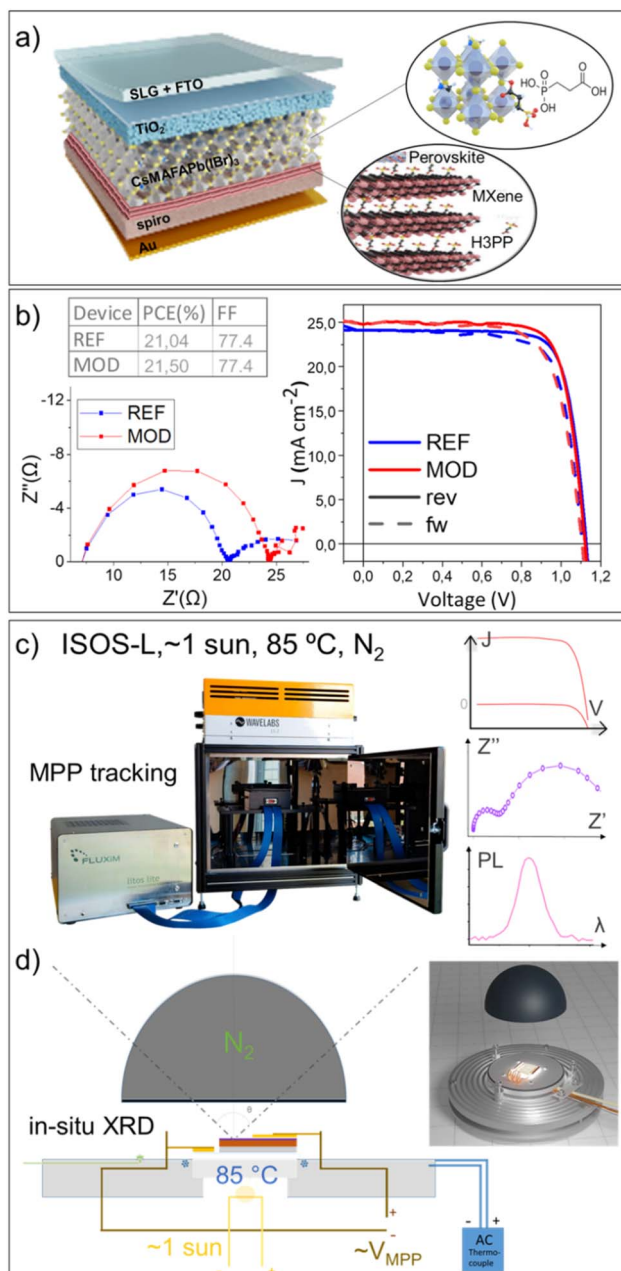


Fig. 1 PSCs analyzed by *in situ* characterization at 85 °C and constant irradiation of $\sim 100 \text{ mW cm}^{-2}$. (a) Schematic representation of the modified n-i-p PSC devices (MOD) studied. (b) Initial EIS and J - V data of the champion devices before experiments (c) LITOS LITE™ commercial setup employed for the MPPT stability analysis of PSCs at 85 °C and $\sim 100 \text{ mW cm}^{-2}$ illumination and (d) *in situ* XRD at $\sim V_{\text{MPP}}$ at 85 °C and $\sim 10^5 \text{ mW cm}^{-2}$ illumination.²³

presenting the MPP decay as it contains more information than only the PCE ($V_{\text{MPP}} \times J_{\text{MPP}}/P_{\text{in}}$).

Ionic and electronic changes

In order to reveal how the electronic and ionic mechanisms in our devices differed with the added modifications, we carried out electrochemical impedance spectroscopy (EIS),

photoluminescence (PL) emission and time-resolved PL decay (TRPL) on all PSCs before and after the operational stress tests. EIS was measured at the open circuit voltage (V_{OC}) under a logarithmic array of light intensities from 0.001 to 103.5 mW cm^{-2} . Fig. 3a displays the EIS data fitted to a standard Matryoshka circuit, as introduced by Bisquert *et al.*, (see Fig. S1).^{27,28} In high efficiency PSCs, the resistances extracted from the high- and the low-frequency (LF) Nyquist arcs (R_{HF} and R_{LF} , respectively) are attributed to a unique physical process in the perovskite, where the LF contribution originates from the phase-delayed electrical response by perovskite mobile ions. While parameters obtained from six different cells of each variation were very similar in all fresh devices (represented as grey symbols in Fig. 3), a parameter distribution was found depending on device-to-device variation and the time that EIS was measured after the MPP test. The fitted EIS values PSC measured 10 min and 90 min after stress from each variation (MOD, REF) are represented by red circles and blue squares respectively, with cells measured in-between laying within the shaded red and blue areas (longer times shown in Fig. S2).

EIS analyses show three main changes in MOD and REF devices measured after the high temperature operational test, compared to before.

The smallest (parallel) resistance has a predominant contribution to the resulting resistance value from EIS. Recombination resistance (R_{rec}) shows at high light intensities, while shunt resistance dominates in dark and low illumination. In Fig. 3b, the resistances are represented as $R_{\text{SUM}} = R_{\text{HF}} + R_{\text{LF}}$.²⁹ First, a change in the recombination mechanisms in the perovskite film is deduced by reduced slope of R_{SUM} vs. V_{OC} (eqn (S2)) at the highest light intensities. The increased apparent-ideality-factor (n_{ap}) and the reduced intersect (R_0) observed, suggest higher recombination rates due to increased number of defect trap states after stress, resulting in lowered V_{OC} at high light intensities.³⁰ Second, at lower light intensities, a large decrease in R_{SUM} occurred in parallel to a V_{OC} drop at the low light intensity range. We recently identified this change as a shunt-like mechanism causing carrier loss within the perovskite film, named as perovskite shunt resistance (R_{psh}).²⁹ In contrast with traditional shunt resistance of a photovoltaic device, R_{psh} does not appear to have a constant value as would be found for an ohmic contact between the selective layers. Instead, R_{psh} depends on the electric field without simple exponential behavior, suggesting it is a material property of the perovskite film. The decrease in R_{psh} is significantly more severe for the REF devices, extending towards higher light intensities and competing with the R_{rec} signal even at the highest illumination. Third, a reduction of between one and three orders of magnitude of the low frequency time constant (τ_{LF}) (Fig. 3c), $\tau_{\text{LF}} = C_{\text{LF}} \cdot R_{\text{LF}}$ (eqn (S1)), indicates increased ion conductivity, in-turn dependent on the ionic density and mobility.³¹⁻³⁵ After operation at 85 °C, both the REF and the MOD low frequency phase signal contribution (Fig. S3) shifted right towards higher frequencies, leading to lower time constants (τ_{LF}) in the fittings (Fig. 3c). This increase in ion-related contributions was delayed in MOD devices, particularly at lower light intensities.



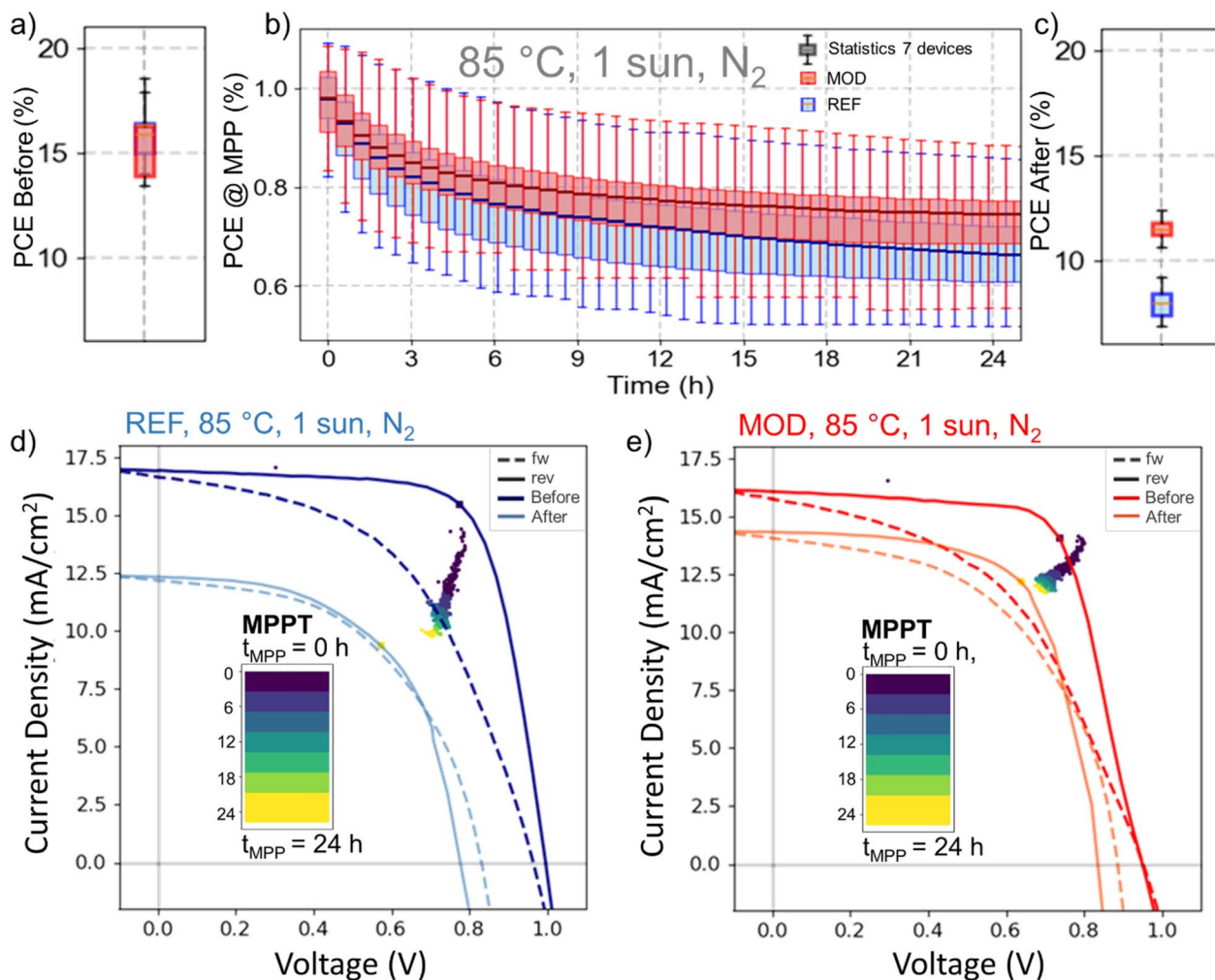


Fig. 2 Operational stability test of PSCs analyzed at 85 °C and 100 mW cm⁻² (~1 sun) LED illumination using LITOS LITE™ for seven REF PSCs (blue) and seven MOD PSCs (red). (a) Power conversion efficiency (PCE) before operational testing for REF and MOD devices extracted from the *J*-*V* reverse scan measured under AM1.5G solar simulator. (b) LITOS LITE™ MPP tracking at 85 °C, normalized to initiation of the test (boxplots marking the statistics include 7 PSC). (c) Power conversion efficiency (PCE) after operational testing for the same devices as in (a). (d) and (e) MPP wandering plot showing the change in *V*_{MPP} and *J*_{MPP} for two PSC of initial equal performance (in measurement shown in (a)), plotted on the before and after *J*-*V* curve taken immediately before (cold cells) and after (warm cells) MPP test, (d) without modification (REF) and (e) with MXene:H3PP interface and HP:H3PP bulk modification (MOD).

Ion mobility or density simulations

When comparing the experimental EIS data through Nyquist and Bode-plots (Fig. S3) with simple drift diffusion simulated data (Fig. S4),^{17,32,36} we found that the plot shifts observed could be better replicated by increasing the cation mobility, rather than increasing the cation density (Fig. S4).

We propose that the τ_{LF} changes could be related to the creation of high mobility channels for ions. It is worth noting that fresh devices before operation showed substantially higher τ_{LF} for MOD compared to REF PSCs, τ_{LF} extraction for MOD devices was particularly uncertain since the LF arc extended beyond the longest modulation wavelength measured. This indicates an initial effect of the H3PP and MXene modifications, possibly by hindering ionic channel pathways or bonding to loosen mobile ionic states (such as iodine vacancies) at the surfaces.³⁷⁻⁴² The remarkably low τ_{LF} , observed in REF PSCs

measured immediately after stress, were jointly observed with effects of shunting seen on dark *J*-*V* curves (Fig. S2). Depending on how long time has passed after the operational test when the EIS was measured, both REF and MOD parameters showed a strong return towards initial values (reversibility). Still, MOD values showed a smaller initial change after stress, especially in these “recovering” regions (Fig. 3c and S2). The R_{psh} , R_{rec} , ion mobility increase, and V_{OC} drop appear to be reversible phenomena with some irreversible component, and the employed modifications interceded with these reversible processes.

In summary, the functionalized MXene and the presence of the bulk additive mitigated recombination and shunt losses after the high temperature stress test, resulting in higher V_{OC} retention than REF devices at high and low light intensities, respectively. The modification in PSCs led to an initially



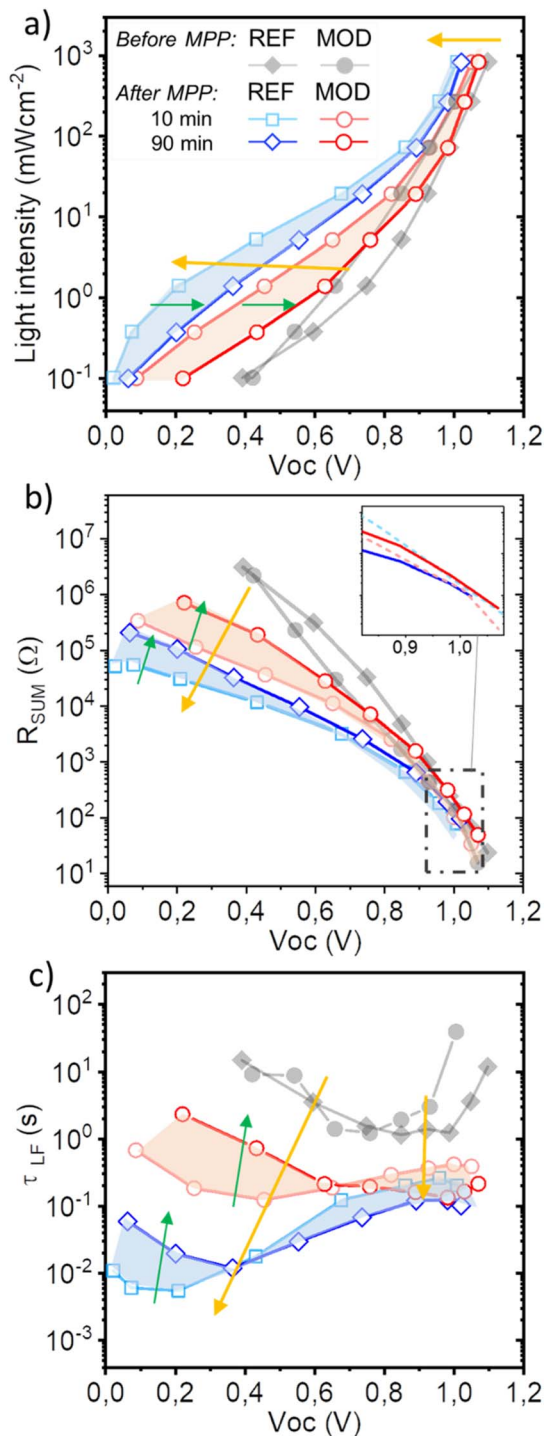


Fig. 3 EIS at a logarithmic array of light intensities, before and after MPP operation at 85 °C for ~24 h. Values obtained from samples measured 10 min and 90 min after operation from each variation (MOD/REF) with shadowed approximate range of additional operated solar cells for visual guide. Yellow arrows mark the degradation trend. (a) Light intensity/ V_{OC} relationship. (b) Total resistance $R_{SUM} = R_{HF} + R_{LF}$. (c) The low frequency time-constant, τ_{LF} . Dark green arrows mark the recovery phenomena of τ_{LF} depending on the time of EIS measurement after stress.

increased low frequency time constant (τ_{LF}), related to a lowered ionic conductivity on fresh devices, and minimized the decrease of τ_{LF} after temperature stress, in comparison with REF PSCs. Below, we reveal that abovementioned performance losses through R_{rec} , R_{psh} and τ_{LF} in REF can be attributed to the formation of a decomposition-related material located at the perovskite grain boundaries, which appearance is reduced in the presence of H3PP and MXene (in MOD devices).

Small bandgap changes and loss of generation

To relate the observed V_{OC} changes to possible changes in the perovskite bandgap, photoluminescence emission spectra (PL) was measured from the incident light side of the n-i-p solar cells before and after stress. No new peaks or protruding shoulders could be observed on PL emission spectra for either of the samples after stress, confirming that the perovskite did not undergo significant halide segregation or phase separation into any photo-emissive material.⁴³ However, a ~6 nm blue shift in PL (fitted by a Gaussian function), Fig. 4a and b, and a 5-fold drop in emission intensity was observed after the stress test. The decrease in PL intensity was coherent for both MOD and REF PSCs, with the blue shift being slightly more pronounced in the REF samples. As will be shown in XRD studies below, the thermal lattice expansion is similar in both MOD and REF, and HP lattices are not in an expanded state at the end of tests once thermal stress is removed. The PL blue shift after stress cannot be attributed to thermal expansion. While deep defects or structural changes would cause larger shifts, smaller PL shifts have been reported to stem from interaction with electron states close to the valence- and conduction-band, such as electron-phonon interactions, conversion of shallow defects to non-radiative interaction, and cation migration.⁴⁴⁻⁴⁶ Additional support that the blue shift is defect-related was seen by TRPL (Fig. 4c). We assign the reduced signal at longer lifetimes to trap-mediated paths in the radiative recombination process, apparently higher in REF than in MOD after stress. This is probably due to a loss of photoactive material, defective photo-generating pathways or shallow defect reactions.⁴⁷

In situ structural and performance evolution at 85 °C

By employing a custom-made X-ray diffraction (XRD) holder (Fig. 1b) in a common use XRD parallel beam setup under similar conditions and timeframe as the MPPT test, operational testing could be coupled with the structural evolution in full PSC devices. XRD scans were acquired continuously over a timeframe similar to the MPP study while monitoring the device electronic output by J - V curves and current output at approximate V_{MPP} . Structural effects were induced independently by the heating stage (low (27 °C) and high (85 °C)), and from illumination (dark and at 100 mW cm⁻²), all at constant N₂ flow enclosed by a graphite dome. Illumination was shone from the solar cell glass side by a white LED (Fig. S1a) calibrated to produce a similar J_{SC} as at 1 sun AM1.5G conditions.

Fig. 5 shows gathered structural and performance data for two PSC, one REF and one MOD PSC, through stages of the experiment including varied light and heat, and the main and



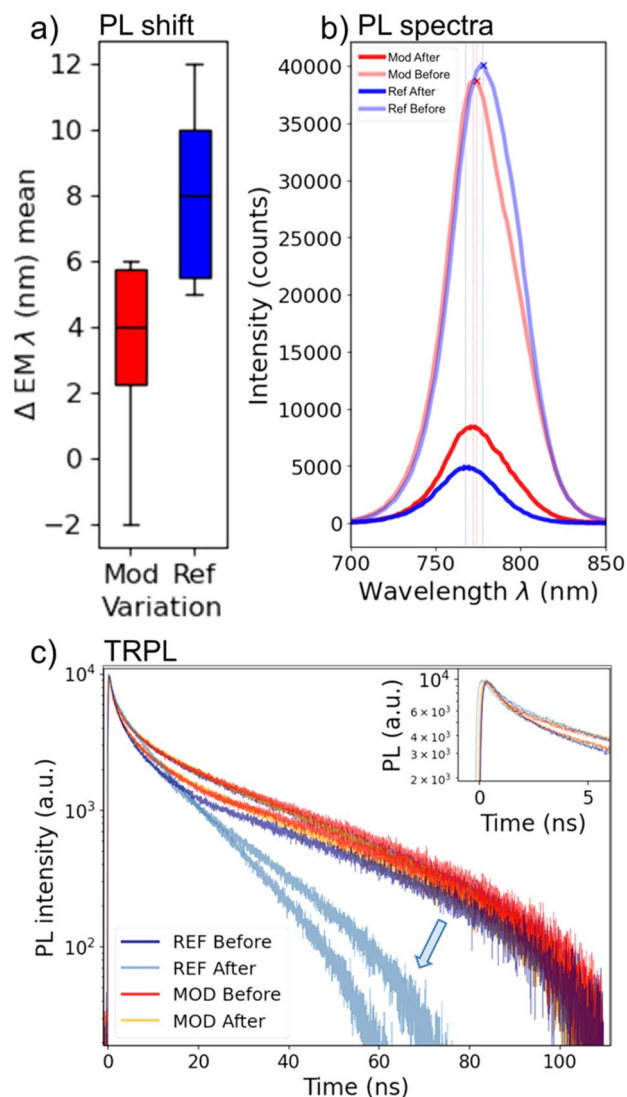


Fig. 4 PL before and after operational testing at 85 °C. (a) Statistical distribution of shift (Δ) in the Gaussian position fitting of PL spectra emission wavelength (EM λ). A mean of spectra positions acquired before and after TRPL, compared before and after operational stability. (b) An example of PL spectra to visualize the PL intensity drop after operation. (c) TRPL of two representative REF/MOD PSC taken before and after operational stability at 85 °C, showing how the REF PSCs lose luminescence intensity at longer timescales.

longest stress block at 85 °C and 105 mW cm⁻² illumination for over 10 h. The experiment can be summarized in five stages as follows:

- (1) Dark at room temperature (RT, 27 °C) (Fig. 5a-stage 1).
- (2) Subsequent light and temperature addition from RT to 85 °C, and repeatedly turning the light on and off (Fig. 5a-stage 2).
- (3) Holding the temperature at 85 °C under illumination overnight for over 10 hours (Fig. 5a-stage 3).
- (4) Cool-down and light on and off (Fig. 5a-stage 4).
- (5) Switching the illumination on/off at RT and comparison of illumination and conductive heating (Fig. 5a-stage 5).

We observed that the integrity of all crystalline HP is maintained throughout the experiment with only different degree of

isotropic thermal expansion independent on time, while the average intensity of all XRD peaks from the REF sample showed a small gradual decrease at 85 °C and illumination (Fig. 5c), indicating that a small part of the perovskite crystal is gradually lost. On the contrary, the MOD sample, where grains are surrounded by the H3PP molecule and the MXene, did not showcase this decrease (Fig. 5c and S7b). This HP peak intensity decrease was the only crystalline change observed in the HP in agreement with the gradual performance (PCE) decrease observed for the REF sample (Fig. 5d). Our results suggest that while part of the perovskite material is converted into a non-crystalline component, the crystalline integrity of the remaining HP is maintained.

Upon applying the illumination, a discrete 2θ left shift was seen in all perovskite diffraction peaks (Fig. S5), indicating an increase of the lattice parameter (Fig. 5b), while peaks from the substrate (fluoride doped tin oxide (FTO)), remained unaffected (Fig. S7). When the set temperature was elevated to 85 °C, a second larger shift was noted, consistent with the known effect of thermal lattice expansion.⁴⁸ The lattice parameter at 85 °C and illumination remained steady along a ~ 14 h period, without additional changes in PbI₂ content or the creation of new crystalline features (Fig. S7a). After fitting all perovskite peaks with a Pseudo Voigt function and applying Nelson-Riley regression to obtain the lattice parameter (a) at the time of each measurement, no deviation from cubic structure was detected in any of the conditions, including no major changes in microstructure (Fig. S7c). Surprisingly, apart from an initially smaller lattice parameter in the MOD devices,²³ the magnitude of the changes observed in the a -value for REF and MOD were the same, within the resolution of the experiment. Both REF and MOD HP lattice parameters transitioned between seemingly discrete values (a jump up slightly below 0.01 Å from the illumination and ~ 0.02 Å when heating from 27 to 85 °C) (Fig. 5a, b, and S6). Once stabilized at higher temperatures, the gradual loss in peak intensity seen on the XRD (at around 6–16 h in Fig. 5c, stage 3) can be linked to the gradual decrease in the PCE output observed for REF (Fig. 5d stage 3, to below 10%) under continued exposure to 85 °C heating and ~ 100 mW cm⁻² LED illumination. The PCE of the MOD device shows a flat curve upon prolonged exposure to the harsh conditions (Fig. 5d), indicating that the performance of MOD devices is better maintained at higher temperature than for the REF devices, in agreement with the response observed under the stability analysis shown in Fig. 2, under the ISOS-L3 protocol.

The J - V data measured during the *in situ* XRD (Fig. 5d-h) provided additional understanding of performance consequences of the crystalline changes, and accentuated the importance of the MXene:H3PP interlayer in MOD solar cells.⁴⁹ J - V curves taken before and after stress in XRD (Fig. 5e and f, J - V curve A and E), correspond well to J - V curves observed before and after stability tests under MPP tracking, however, the changes observed on J - V curves for the REF sample during stress were much larger than expected. We observed an initial sharp drop in V_{OC} (Fig. 5e, J - V curve B) when raising the temperature up to 85 °C, this response is also observed in the calculated power conversion efficiency of the REF sample



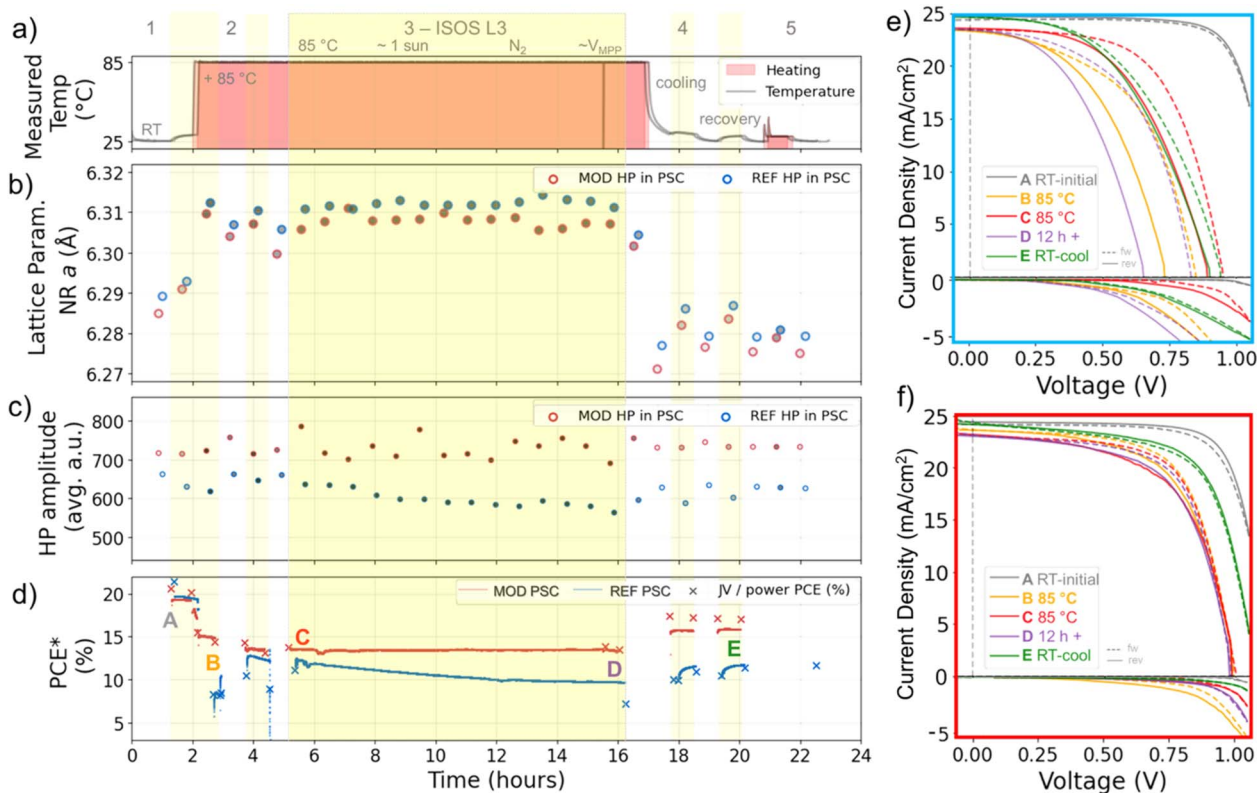


Fig. 5 Simultaneous structural and operational changes during *in situ* XRD of 2 devices, reference PSC (REF) and PSC with bulk HP:H3PP additive modification and MXene:H3PP interface modification (MOD). The output is shown only for the connected pixel, while the area of XRD is outside of the biased area (see Fig. 1b). (a) Temperature setting and measured temperature, top marked 5 key stages during the experiment. (b) Lattice parameter from y -axis intercept in Nelson–Riley fitting (Fig. S4) of XRD scans at 45 min intervals, (c) average amplitude of collective HP peaks on each XRD scan. (d) PCE calculated from output current density at approximated V_{MPP} and $J-V$ curves, the yellow shade encompassing (a)–(d) indicates the periods where light is on. (e) and (f) Selection of $J-V$ curves (marked A–E in (d)) taken at 40 mV s^{-1} to show diode progression for the MOD device (red) and for REF device (blue), and $J-V$ curves in dark conditions, closest to the points in the approximate PCE plot of $J-V$ s (e) MOD (f) REF. fw-forward scan, rev-reverse scan.

(Fig. 5d, $J-V$ curve B). V_{OC} -loss is a common trend upon operation of PV at higher temperatures, but the detected drop of above 0.2 V (Fig. 5e, $J-V$ curve B) for REF solar cells is double that reported in previous works.⁵ During the lengthening of interatomic distances taking place upon thermal lattice expansion, as observed in Fig. 5b, the tightly bound HP/HTM interface could lose contact or band-alignment, causing the voltage drop observed in Fig. 5d–f.^{48,50} Strain compensation by the mesoporous TiO_2 layer in the identical device structure has been observed as a lack of residual strain transfer from the FTO layer to the HP after glass cool down post-fabrication in our previous work.²³ It is possible that the MXene provides a similar function, leading to a lack of strain response in the MOD devices as both interfaces at the ETL and HTL sides of the perovskite may compensate the interfacial strain produced during heating. At the heat-up, during stage 2, the PCE of both REF and MOD samples showcased a drop of 50% and 25%, respectively, a change in the initial efficiency from 20% to 10% for the REF sample and from 20% to 15% for the MOD sample. However, the fast drop is followed by a fast recovery of the efficiency in the REF sample, from 10% to 12% (Fig. 5d and e $J-V$ curve C). However, the sudden drop and fast recovery of the efficiency detected for the REF sample is not observed when the

use of the MXene:H3PP interlayer is employed, an indication of the beneficial effect of the MXene interface which reduces this effect in half or completely prevents the sharp drop, possibly preventing future device degradation (Fig. 5d and f $J-V$ curves A, B, C).⁵¹

Both REF and MOD solar cells transition to inverted hysteresis at 85 °C (Fig. 5e and f, $J-V$ curves B–D).⁵² In addition to MOD showing a smaller drop of V_{OC} upon raising the temperature, we also observe a smaller hysteresis. While the REF sample undergoes a large increase in the hysteresis index, not stemming from reduced fill factor, but rather from a 100–150 mV gradually lower V_{OC} in the reverse scan direction (at speed 40 mV s^{-1} , Fig. 5e, $J-V$ curves C and D), the absolute hysteresis index at 85 °C remains relatively small in the MOD device (Fig. 5f, $J-V$ curves C and D). Hysteresis is known to originate as a combination of factors related to electrode poling or ionic accumulation at the interface with the selective contact, modulating the internal electric field and altering the carrier transport and extraction.^{53–55} Despite that the exact mechanisms of hysteresis remain elusive, we attribute the increase in hysteresis in REF to the temperature – increased ionic motion inside the perovskite layer, in accordance with the lowered τ_{LF} seen in EIS. The increased mobility would result in larger ionic



build-up, widening the depletion layer near the contacts. Thermal expansion (Fig. 5b) is likely accompanied by a significantly wider perovskite band gap.^{49,56} As a consequence, both interfacial band bending and wider gap will affect the band alignment between the perovskite and the selective contacts, probably originating the $J-V$ changes and inducing inverted hysteresis. The retained small hysteresis found in MOD is likely due to the hindered ion migration through the layer, also shown by larger τ_{LF} values from EIS, and prevention of the formation of an interfacial depletion layer by the MXene interlayer. Additionally, REF PSC dark $J-V$ shows lower shunt resistance (Fig. 5e) than MOD (Fig. 5f), in line with our results from EIS. Upon cooling, the REF $J-V$ reverse scan recovers to a point of similar V_{OC} as the forward scan, with similar inversed hysteresis as MOD but with both lower current and voltage output (Fig. 5e and f, $J-V$ curves D and E).

To summarize, *in situ* XRD revealed two mechanisms of PSC loss at 85 °C; the first related to interfacial strain response during thermal expansion, the second to loss of crystalline HP.

At 85 °C, thermal lattice expansion leads to interfacial strain that produces a (partially reversible) voltage drop, and the use of MXene:H3PP interlayer greatly prevented this effect. However, the main cause of the gradual performance decay at 85 °C, was revealed to be the gradual loss of crystallinity in the bulk HP.

Upon cooling, partial peak recovery in REF (Fig. 5c) suggests recrystallization, implying the degraded phase remains in the film, but is undetectable by our XRD measurements. The HP structural evolution directly impacts device PCE, which also shows partial reversibility. In the presence of the MXene and

H3pp additive (MOD device), HP XRD peaks did not decline, and the earlier degradation of the PSC, observed in the REF device, was prevented.

Our results from EIS, PL and *in situ* XRD are in agreement with separate observations of modified ionic movement, the appearance of inversed hysteresis, apparent shunting dependent on the electric field, and increased recombination and interface delamination above temperature thresholds.^{57–67} This is the first time these phenomena are reported jointly on the same timeline, revealing their direct connection.

Morphology investigations

In order to reveal more about the morphological consequences of the crystal decomposition observed in the REF solar cells during XRD analysis, we employed *ex situ* complimentary transmission electron microscopy (TEM) with energy dispersive X-ray spectroscopy (EDX), and cross section scanning electron microscopy (CS-SEM). Solar cells from the same batch as in Fig. 5 were split in two in preparation for CS-SEM (Fig. S8) and investigated with and without treatment at $\sim 100 \text{ mW cm}^{-2}$ and 85 °C under high N_2 flow for 20 hours. After testing, HP samples were obtained by scratching the surface of the PSC, and the resulting powder was analyzed by HRTEM (Fig. 6a–f).

HRTEM of un-treated MOD and REF HP samples (Fig. S9) looked similar with clear distinguishable HP grains, as expected. However, the REF sample treated at 85 °C showed a new material enclosing the HP crystallite grains (Fig. 6a–c and S10). This core-shell nanostructure encompasses a well-crystalline

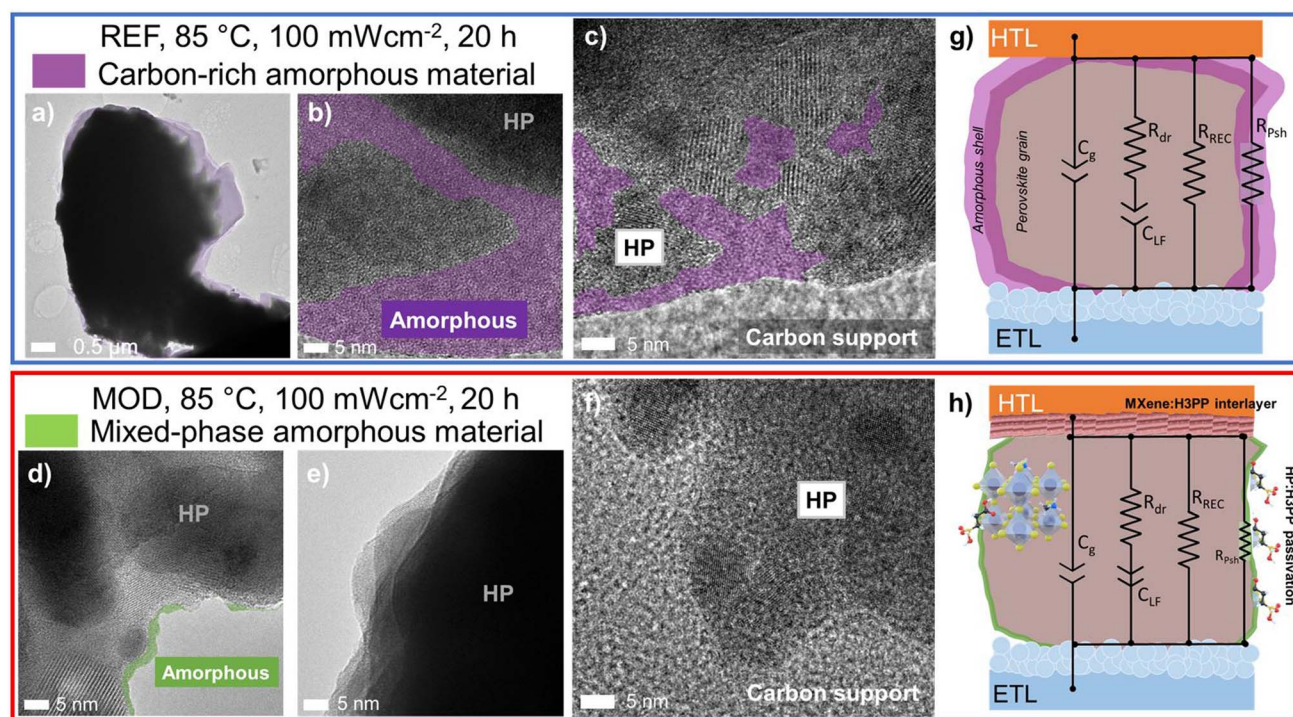


Fig. 6 Selection of TEM images of HP from PSC treated at $\sim 10^5 \text{ mW cm}^{-2}$ and 85 °C 1 week before TEM experiments. (a)–(c) HP from the treated REF solar cell half, with a carbon-rich amorphous shell layer shaded purple. (d)–(f) HP from the treated MOD solar cell, with scarce mixed-phase amorphous material shaded green. (g) REF PSC schematic showing how an amorphous material has been created around HP, leading to additional shunt paths, while in (h) MOD HP, most areas do not show any new material.



HP grain in the inside (as confirmed by EDX, Fig. S11), while the outside or the “shell” is characterized by a transparent and amorphous surface material surrounding the HP grains, that would also be undetectable by XRD (colored in purple in Fig. 6a–c, and S10). In addition, the treated MOD HP sample is observed to be more isolated with more easily resolved grains, despite a very scarce presence of another mixed-phase material (colored in green in Fig. 6d–f, and S10). The maintained intensity of peaks observed in XRD for the MOD HP sample after treatment, together with the low abundance of the mixed-phase compared to the amorphous phase observed in the REF sample, indicates that, in MOD, the employment of the organic molecule H3PP circumvents and delays the formation of the amorphous material on the surface of the grains. This mitigation is probably due to the strong interaction between the H3PP molecule and the HP that could immobilize ions,²¹ maintaining the integrity of the HP surface. Further investigation with EDX (Fig. S11) showed that the “shell” material had a high carbon content. We deduce that the large A-site organic cation, being the only available source of carbon in the HP material, should be highly involved in the creation of the decomposition product. Even if formation of a new material or defects (Fig. S12) was not completely prevented in MOD, the morphology and make-up of the two materials were vastly different. New phase regions in MOD had particles imbedded in the dimmer amorphous area as seen in TEM (Fig. S10) and STEM (Fig. S11). EDX of the bright particles revealed that they contained strong Ti and O signals together with components of the HP, indicating a likely mixture of TiO₂/MXene and HP. The MOD HP samples also showed a stronger titanium and oxygen signal, likely from the added content of titanium carbide from the MXene layer, indicating a close contact between the MXene layer and the perovskite.

Extended discussion

We have observed a completely different response in REF and MOD PSCs at 85 °C. We observed that reversible processes in the HP material can be related to the PSC device's performance. Creation of amorphous-like regions in devices occurs before the irreversible step of PbI₂ formation. Our observations of an amorphous “shell” layer around HP are in accordance with previous studies made by TEM, that agree with the layer by layer decomposition from the surface of the perovskite grains at 85 °C previously reported.^{68–71} It has been reported that the presence of intermediate, amorphous like, areas of MAPbI₃ with what looks like ionic channels facilitate the reverse intercalation of MAI and PbI₂, and that the appearance of PbI₂ is frequently related to ion migration, specifically of the large cation.^{71–75} In our experiments we do not observe any significant PbI₂ formation (Fig. S7a), instead, we observe that the amorphous material found after operation at 85 °C has a very high carbon content (Fig. 6g), which implicates the involvement of the large organic cations. In addition, it is possible that this material can facilitate both fast electron and ion diffusion.^{76,77} Faster electron and ion transport in this new amorphous phase is coherent with changes observed on the photovoltaic performance of the solar

cells (Fig. 2, and 5) and the EIS signal changes after stress (Fig. 3) that indicate reversible formation of shunts; conductive channels (R_{psh}), and increased ion mobility (decreased τ_{LF}). The findings of the formation of this core–shell material observed after treatment at 85 °C, are in good agreement with the observed conductivity and mobility increases measured by the EIS analysis (Fig. 3) and seen on *in situ* *J*–*V* curves (Fig. 5e),⁷⁸ which are smaller in the presence of H3PP and MXene (in MOD devices). This effect is represented by the electric circuits shown in Fig. 6g and h that accentuate the shunt paths created by the amorphous phase.

We propose that the carbon-rich amorphous material formed around the perovskite grains supposes highly conductive channels leading to both lower barriers for ionic migration (lower τ_{LF}) and electrical shunt-losses (lower R_{psh}). *Ab initio* calculations have shown that surface defects dominate the paths for ion migration in lead-iodide perovskites due to surface-assisted formation of migrating defects.⁷⁹ We show that the presence of these defects relates to the reversible formation of a surface layer in the HP, that leads to 3-fold increased ion conductivity, likely due to increased ion mobility. In contrast, in the MOD PSCs, the HP grain surface is passivated by organic additive H3PP, the migration in the shallow defects in the surface can be partially prevented, resulting in enhanced stability. The *ab initio* calculations support that the prevention of access to shallow trap states seen by passivation with H3PP,²¹ can be synonymous with the prevention of both ion migration and shallow trap state propagation. At higher temperatures, this new “shell material” surrounding the HP grains seems to lead to further voltage drops and shunting-like features similar to what was seen on EIS and during operational testing. The recovery of XRD intensity upon removal of the heating explains the reversibility in efficiency as the shell converts back into crystalline perovskite. Therefore, the new “shell” formed around the HP grains is likely responsible for the reversible component seen by EIS. Incorporation of the H3PP molecule in the precursor of HP prevents its amorphization as the intensity of XRD at higher temperature is not gradually lost, yet it seems like if any material is formed its properties could be different, leading to a different loss-in-performance behavior.

Although the general voltage drop in MOD devices is concerning, it is a common trend upon operation of PV at higher temperatures.⁵ Even if additional voltage drops might still occur, we show how further improvements in grain boundary passivation and interfacial dynamics at varying temperature and illumination to account for changes could reduce the problems faced in these conditions. All factors considered, we could relate the morphological and crystalline changes occurring to electronic effects observed. Our study shows the value of simultaneous data-acquisition to arrive at certainty of the effects of observed changes both electronically and in physical and chemical properties.

Conclusions

In conclusion, we carried out stability analysis of PSCs at 85 °C and white LED light ($\sim 100 \text{ mW cm}^{-2}$) illumination under



operational testing with MPP tracking. Analyses were carried out under *in operando* XRD testing conditions for n-i-p devices with and without additive engineering (passivation with H3PP additive) and functionalized with MXene interlayers between the perovskite and the HTM (MOD and REF). Our findings reveal that perovskite lattice expansion occurs under conductive heating of the device and, additionally, from illumination-induced heat, independent of the surface temperature level. The expansion observed in the bulk of the perovskite was non-differentiable between MOD and REF, and unchanging with time. Instead, the difference in degradation seen during MPPT was due to a gradual transformation of the HP without additive into an amorphous phase around the crystal grains. In post-mortem analysis, this phase appeared as a carbon-rich material in REF samples, and as a mixed-composition material in MOD samples. The formation of this carbon-rich phase was directly related to an increase in ionic conductivity, shallow trap state propagation and apparent shunting of devices, all partially reversible. Detailed EIS analysis revealed increased ionic conductivity and lowered shunt resistances, related to the creation of ionic and electronic conductive paths through the device, a phenomenon delayed by the phosphonate-group based additive at grain boundaries. The direct observation of the surface-induced decomposition and performance decay together when employing *in situ* XRD, provides proof that paths are located at the surfaces of grains. In addition, we demonstrate that the functionalized-MXene interlayer protects the interface during thermal expansion to partially prevent voltage loss. The results provide a framework for future efforts with the aim of stabilizing perovskite solar cells under high temperature operation. Additionally, interfacial effects between the (expanding) perovskite and HTM layer created an additional sharp, temporary, voltage drop during changing temperature, and MXene interlayers can aid in compensating for problems caused by interfacial strain, with a result of higher solar cell stability.

Author contributions

F. B. performed the stability testing, additional measurements and experiments, and prepared the manuscript under supervision by S. R. and M. L. C. J. P. P. operated the XRD and J. P. P. and J. S. helped with XRD experimental design and analysis. S.R. created the EIS protocol, operated by F. B. F. B. designed data collection, collected data and wrote analysis code. J. M. C. R. fabricated the XRD stage with partial support from F. B. N. V. performed the CS-SEM acquisition, B. B. performed the TEM acquisition. M. K. fabricated the solar cell devices with assistance from Z. H. M. L. and S. R. secured the funding for the project and advised on the objectives. All authors contributed in discussions about results, and made comments on the manuscript.

Conflicts of interest

There are no conflicts to declare.

Data availability

Supplementary information: supporting figures, fittings, calculations, drift diffusion EIS simulations, TEM images, EDX and more detailed depictions of analysis have been included as part of the supplementary information (SI). See DOI: <https://doi.org/10.1039/d5el00104h>. Data for this article, including the raw data for all figures in the article are available at CORA – Research Data Repository (RDR) at <https://doi.org/10.34810/data2409>.

Acknowledgements

The authors acknowledge the Spanish Ministry of Science and Innovation for the predoctoral contract to F. B. with Reference No. PRE2020-092669 of the Project No. SEV-2017-0706-20-3 and the contract of N. V. with ref. PRE2022-103001 of the project CEX2021-001214-S-20-6. This work constitutes part of the Materials Science PhD Program for F. B., N. V. and Z. T. at the Universitat Autònoma de Barcelona (UAB, Spain). This project was funded by the project PerMXSol (ref. PID2022-143344OB-I00), project TAMPOPO, (ref. PID2021-122349OAI00) and project SPOT-IT (ref. PCI2023-146003-2) funded by MCIN/AEI/10.13039/501100011033 and “ERDF A way of making Europe”, co-financed by the European Union. ICN2 is funded by the CERCA Programme/Generalitat de Catalunya. Thanks to Agencia de Gestio d’Ajuts Universitaris i de Recerca for Grant No. SGR 01617 and the Xarxa d’ R + D + I Energy for Society. The ICN2 is supported by the Severo Ochoa Centres of Excellence programme, Grant CEX2021-001214-S, funded by MCIN/AEI/10.13039.501100011033. We would like to acknowledge the team from Syre (syre.ai), especially Dr Brian Carlsen, for lending us their software prototype facilitating steps in the automatic processing of results from experiments. We would like to thank Juan A. Anta and colleagues for their help in setting up the drift diffusion simulations. We would like to thank Damaso Torres for his aid with graphic design, and Rafael León Lavin for modifications of the XRD holder and LITOS customized sample stages.

Notes and references

- 1 S.-P. Feng, Y. Cheng, H.-L. Yip, Y. Zhong, P. W. K. Fong, G. Li, A. Ng, C. Chen, L. A. Castriotta, F. Matteocci, L. Vesce, D. Saranin, A. D. Carlo, P. Wang, J. Wei Ho, Y. Hou, F. Lin, A. G. Aberle, Z. Song, Y. Yan, X. Chen, Y. Yang, A. A. Syed, I. Ahmad, T. Leung, Y. Wang, J. Lin, A. M. C. Ng, Y. Li, F. Ebadi, W. Tress, G. Richardson, C. Ge, H. Hu, M. Karimipour, F. Baumann, K. Tabah, C. Pereyra, S. R. Raga, H. Xie, M. Lira-Cantu, M. V. Khenkin, I. Visoly-Fisher, E. A. Katz, Y. Vaynzof, R. Vidal, G. Yu, H. Lin, S. Weng, S. Wang and A. B. Djurišić, *J. Phys.: Mater.*, 2023, **6**, 032501.
- 2 S. He, L. Qiu, L. K. Ono and Y. Qi, *Mater. Sci. Eng., R*, 2020, **140**, 100545.
- 3 T. J. Jacobsson, A. Hultqvist, A. Garcia-Fernandez, A. Anand, A. Al-Ashouri, A. Hagfeldt, A. Crovetto, A. Abate,



- A. G. Ricciardulli, A. Vijayan, A. Kulkarni, A. Y. Anderson, B. P. Darwich, B. W. Yang, B. L. Coles, C. A. R. Perini, C. Rehmann, D. Ramirez, D. Fairen-Jimenez, D. Di Girolamo, D. L. Jia, E. Avila, E. J. Juarez-Perez, F. Baumann, F. Mathies, G. S. A. Gonzalez, G. Boschloo, G. Nasti, G. Paramasivam, G. Martinez-Denegri, H. Nasstrom, H. Michaels, H. Kobler, H. Wu, I. Benesperi, M. I. Dar, I. B. Pehlivan, I. E. Gould, J. N. Vagott, J. Dagar, J. Kettle, J. Yang, J. Z. Li, J. A. Smith, J. Pascual, J. J. Jeronimo-Rendon, J. F. Montoya, J. P. Correa-Baena, J. M. Qiu, J. X. Wang, K. Sveinbjornsson, K. Hirslandt, K. Dey, K. Frohna, L. Mathies, L. A. Castriotta, M. H. Aldamasy, M. Vasquez-Montoya, M. A. Ruiz-Preciado, M. A. Flatken, M. V. Khenkin, M. Grischek, M. Kedia, M. Saliba, M. Anaya, M. Veldhoen, N. Arora, O. Shargaieva, O. Maus, O. S. Game, O. Yudilevich, P. Fassl, Q. S. Zhou, R. Betancur, R. Munir, R. Patidar, S. D. Stranks, S. Alam, S. Kar, T. Unold, T. Abzieher, T. Edvinsson, T. W. David, U. W. Paetzold, W. Zia, W. F. Fu, W. W. Zuo, V. R. F. Schroder, W. Tress, X. L. Zhang, Y. H. Chiang, Z. Iqbal, Z. Q. Xie and E. Unger, *Nat. Energy*, 2022, 7, 107–115.
- 4 Y. Zhao, J. Zhang, Z. Xu, S. Sun, S. Langner, N. T. P. Hartono, T. Heumueller, Y. Hou, J. Elia, N. Li, G. J. Matt, X. Du, W. Meng, A. Osvet, K. Zhang, T. Stubhan, Y. Feng, J. Hauch, E. H. Sargent, T. Buonassisi and C. J. Brabec, *Nat. Commun.*, 2021, 12, 2191.
- 5 M. Jošt, B. Lipovšek, B. Glažar, A. Al-Ashouri, K. Brecl, G. Matič, A. Magomedov, V. Getautis, M. Topič and S. Albrecht, *Adv. Energy Mater.*, 2020, 10, 2000454.
- 6 M. Koehl, M. Heck, S. Wiesmeier and J. Wirth, *Sol. Energy Mater. Sol. Cells*, 2011, 95, 1638–1646.
- 7 G. A. Landis, Review of solar cell temperature coefficients for space, <https://ntrs.nasa.gov/citations/19950014125>.
- 8 M. V. Khenkin, E. A. Katz, A. Abate, G. Bardizza, J. J. Berry, C. Brabec, F. Brunetti, V. Bulović, Q. Burlingame, A. Di Carlo, R. Cheacharoen, Y.-B. Cheng, A. Colsmann, S. Cros, K. Domanski, M. Dusza, C. J. Fell, S. R. Forrest, Y. Galagan, D. Di Girolamo, M. Grätzel, A. Hagfeldt, E. von Hauff, H. Hoppe, J. Kettle, H. Köbler, M. S. Leite, S. Liu, Y.-L. Loo, J. M. Luther, C.-Q. Ma, M. Madsen, M. Manceau, M. Matheron, M. McGehee, R. Meitzner, M. K. Nazeeruddin, A. F. Nogueira, Ç. Odabaşı, A. Osherov, N.-G. Park, M. O. Reese, F. De Rossi, M. Saliba, U. S. Schubert, H. J. Snaith, S. D. Stranks, W. Tress, P. A. Troshin, V. Turkovic, S. Veenstra, I. Visoly-Fisher, A. Walsh, T. Watson, H. Xie, R. Yildirim, S. M. Zakeeruddin, K. Zhu and M. Lira-Cantu, *Nat. Energy*, 2020, 5, 35–49.
- 9 K. T. Tanko, Z. Tian, S. Raga, H. Xie, E. A. Katz and M. Lira-Cantu, *MRS Bull.*, 2025, 50, 512–525.
- 10 P. Holzhey and M. Saliba, *J. Mater. Chem. A*, 2018, 6, 21794–21808.
- 11 Q. Jiang, R. Tirawat, R. A. Kerner, E. A. Gaulding, Y. Xian, X. Wang, J. M. Newkirk, Y. Yan, J. J. Berry and K. Zhu, *Nature*, 2023, 623, 313–318.
- 12 R. K. Gupta, D. K. Kumar, V. Sudhakar, J. M. Beckedahl, A. Abate, E. A. Katz and I. Visoly-Fisher, *Adv. Energy Mater.*, 2025, 15, 2403844.
- 13 S. Su, T. Ahn and Y. Yang, *Adv. Funct. Mater.*, 2025, 35, 2408480.
- 14 G. Y. Kim, A. Senocrate, T. Yang, G. Gregori, M. Grätzel and J. Maier, *Nat. Mater.*, 2018, 17, 445.
- 15 Y.-C. Li, C. Ge, H.-Y. Song, P. Wang, X.-R. Ma and S.-B. Liu, *J. Appl. Phys.*, 2022, 132, 013102.
- 16 L. E. Mundt and L. T. Schelhas, *Adv. Energy Mater.*, 2020, 10, 1903074.
- 17 A. Zohar, M. Kulbak, S. H. Turren-Cruz, P. K. Nayak, A. Kama, A. Hagfeldt, H. J. Snaith, G. Hodes and D. Cahen, *ACS Appl. Mater. Interfaces*, 2022, 14, 34171–34179.
- 18 P. Lopez-Varo, M. Amara, S. Cacovich, A. Julien, A. Yaïche, M. Jouhari, J. Rousset, P. Schulz, J.-F. Guillemoles and J.-B. Puel, *Sustainable Energy Fuels*, 2021, 5, 5523–5534.
- 19 S. Tan, I. Yavuz, N. De Marco, T. Huang, S.-J. Lee, C. S. Choi, M. Wang, S. Nuryyeva, R. Wang, Y. Zhao, H.-C. Wang, T.-H. Han, B. Dunn, Y. Huang, J.-W. Lee and Y. Yang, *Adv. Mater.*, 2020, 32, 1906995.
- 20 Y. Wu, F. Xie, H. Chen, X. Yang, H. Su, M. Cai, Z. Zhou, T. Noda and L. Han, *Adv. Mater.*, 2017, 29, 1701073.
- 21 H. B. Xie, Z. W. Wang, Z. H. Chen, C. Pereyra, M. Pols, K. Galkowski, M. Anaya, S. Fu, X. Y. Jia, P. Y. Tang, D. J. Kubicki, A. Agarwalla, H. S. Kim, D. Prochowicz, X. Borriese, M. Bonn, C. X. Bao, X. X. Sun, S. M. Zakeeruddin, L. Emsley, J. Arbiol, F. Gao, F. Fu, H. I. Wang, K. J. Tielrooij, S. D. Stranks, S. X. Tao, M. Gratzel, A. Hagfeldt and M. Lira-Cantu, *Joule*, 2021, 5, 1246–1266.
- 22 C. Ma, F. T. Eickemeyer, S.-H. Lee, D.-H. Kang, S. J. Kwon, M. Grätzel and N.-G. Park, *Science*, 2023, 379, 173–178.
- 23 F. Baumann, M. Karimipour, J. Padilla-Pantoja, E. Chávez-Angel, J. M. Caicedo Roque, R. Pouteaux, A. Alcalá Ibarra, S. R. Raga, J. Santiso and M. Lira-Cantu, *ACS Energy Lett.*, 2025, 10, 476–483.
- 24 C. Cheng, Y. Yao, L. Li, Q. Zhao, C. Zhang, X. Zhong, Q. Zhang, Y. Gao and K. Wang, *Nano Lett.*, 2023, 23, 8850–8859.
- 25 W.-H. Zhang, L. Chen, Z.-P. Zou, Z.-A. Nan, J.-L. Shi, Q.-P. Luo, Y. Hui, K.-X. Li, Y.-J. Wang, J.-Z. Zhou, J.-W. Yan and B.-W. Mao, *ACS Appl. Mater. Interfaces*, 2022, 14, 31911–31919.
- 26 M. Karimipour, A. Paingott Parambil, K. Tabah Tanko, T. Zhang, F. Gao and M. Lira-Cantu, *Adv. Energy Mater.*, 2023, 13, 2301959.
- 27 A. Bou, A. Pockett, D. Raptis, T. Watson, M. J. Carnie and J. Bisquert, *J. Phys. Chem. Lett.*, 2020, 11, 8654–8659.
- 28 F. Ebadi, M. Aryanpour, R. Mohammadpour and N. Taghavinia, *Sci. Rep.*, 2019, 9, 11962.
- 29 K. T. Tanko, S. R. Raga, N. Vahedigharehchopogh, F. Baumann, M. Karimipour, R. A. Miranda-Gamboa and M. Lira-Cantu, *Sol. RRL*, 2025, 9, 2500162.
- 30 L. Contreras-Bernal, S. Ramos-Terrón, A. Riquelme, P. P. Boix, J. Idígoras, I. Mora-Seró and J. A. Anta, *J. Mater. Chem. A*, 2019, 7, 12191–12200.



- 31 O. Almora, P. López-Varo, R. Escalante, J. Mohanraj, L. F. Marsal, S. Olthof and J. A. Anta, *J. Appl. Phys.*, 2024, **136**, 094502.
- 32 O. Almora, K. T. Cho, S. Aghazada, I. Zimmermann, G. J. Matt, C. J. Brabec, M. K. Nazeeruddin and G. Garcia-Belmonte, *Nano Energy*, 2018, **48**, 63–72.
- 33 D. A. Jacobs, H. Shen, F. Pfeffer, J. Peng, T. P. White, F. J. Beck and K. R. Catchpole, *J. Appl. Phys.*, 2018, **124**, 225702.
- 34 M. T. Neukom, A. Schiller, S. Züfle, E. Knapp, J. Ávila, D. Pérez-del-Rey, C. Dreessen, K. P. S. Zanoni, M. Sessolo, H. J. Bolink and B. Ruhstaller, *ACS Appl. Mater. Interfaces*, 2019, **11**, 23320–23328.
- 35 S. Lammar, R. Escalante, A. J. Riquelme, S. Jenatsch, B. Ruhstaller, G. Oskam, T. Aernouts and J. A. Anta, *J. Mater. Chem. A*, 2022, **10**, 18782–18791.
- 36 A. Riquelme, L. J. Bennet, N. E. Courtier, M. J. Wolf, L. Contreras-Bernal, A. B. Walker, G. Richardson and J. A. Anta, *Nanoscale*, 2020, **12**, 17385–17398.
- 37 D. Moia, I. Gelmetti, P. Calado, W. Fisher, M. Stringer, O. Game, Y. Hu, P. Docampo, D. Lidzey, E. Palomares, J. Nelson and P. R. F. Barnes, *Energy Environ. Sci.*, 2019, **12**, 1296–1308.
- 38 C. Aranda, J. Bisquert and A. Guerrero, *J. Chem. Phys.*, 2019, **151**, 124201.
- 39 A. Guerrero, J. Bisquert and G. Garcia-Belmonte, *Chem. Rev.*, 2021, **121**, 14430–14484.
- 40 E. von Hauff and D. Klotz, *J. Mater. Chem. C*, 2022, **10**, 742–761.
- 41 A. Dualeh, T. Moehl, N. Tetreault, J. Teuscher, P. Gao, M. K. Nazeeruddin and M. Gratzel, *ACS Nano*, 2014, **8**, 362–373.
- 42 J. Torres, I. Zarazua, D. Esparza, J. M. Rivas, M. Saliba, I. Mora-Seró, S.-H. Turren-Cruz and A. Abate, *ACS Appl. Energy Mater.*, 2022, **5**, 12545–12552.
- 43 S. D. Stranks, R. L. Z. Hoyer, D. W. Di, R. H. Friend and F. Deschler, *Adv. Mater.*, 2019, **31**, 1803336.
- 44 M. Ledinsky, T. Schönfeldová, J. Holovský, E. Aydin, Z. Hájková, L. Landová, N. Neyková, A. Fejfar and S. De Wolf, *J. Phys. Chem. Lett.*, 2019, **10**, 1368–1373.
- 45 J. Xu, S. Yu, X. Shang and X. Chen, *Adv. Photonics Res.*, 2023, **4**, 2200193.
- 46 M. Wang, Z. Yang, Y. Chen, L. Shi, K. Cao, Y. Liang, G. Cao, X. Li and L. Qin, *Sol. Energy Mater. Sol. Cells*, 2025, **282**, 113393.
- 47 T. Kirchartz, J. A. Márquez, M. Stolterfoht and T. Unold, *Adv. Energy Mater.*, 2020, **10**, 1904134.
- 48 M. Dailey, Y. Li and A. D. Printz, *ACS Omega*, 2021, **6**, 30214–30223.
- 49 C. Zhu, X. Niu, Y. Fu, N. Li, C. Hu, Y. Chen, X. He, G. Na, P. Liu, H. Zai, Y. Ge, Y. Lu, X. Ke, Y. Bai, S. Yang, P. Chen, Y. Li, M. Sui, L. Zhang, H. Zhou and Q. Chen, *Nat. Commun.*, 2019, **10**, 815.
- 50 S. Ma, J. Tang, G. Yuan, Y. Zhang, Y. Wang, Y. Wu, C. Zhu, Y. Wang, S. Wu, Y. Lu, S. Chi, T. Song, H. Zhou, M. Sui, Y. Li and Q. Chen, *Energy Environ. Mater.*, 2024, **7**, e12739.
- 51 G. G. Njema, J. K. Kibet and S. M. Ngari, *Meas.: Energy*, 2024, **2**, 100005.
- 52 S. N. Habisreutinger, N. K. Noel and H. J. Snaith, *ACS Energy Lett.*, 2018, **3**, 2472–2476.
- 53 R. García-Rodríguez, A. J. Riquelme, M. Cowley, K. Valadez-Villalobos, G. Oskam, L. J. Bennett, M. J. Wolf, L. Contreras-Bernal, P. J. Cameron, A. B. Walker and J. A. Anta, *Energy Technol.*, 2022, **10**, 2200507.
- 54 W. Clarke, M. V. Cowley, M. J. Wolf, P. Cameron, A. Walker and G. Richardson, *J. Appl. Phys.*, 2023, **133**, 095001.
- 55 F. Wu, R. Pathak, K. Chen, G. Wang, B. Bahrami, W.-H. Zhang and Q. Qiao, *ACS Energy Lett.*, 2018, **3**, 2457–2460.
- 56 D. A. Shohonov, D. B. Migas, A. B. Filonov, V. E. Borisenko, R. Takabe and T. Suemasu, *Thin Solid Films*, 2019, **686**, 137436.
- 57 S. Kim, S. Bae, S. W. Lee, K. Cho, K. D. Lee, H. Kim, S. Park, G. Kwon, S. W. Ahn and H. M. Lee, *Sci. Rep.*, 2017, **7**, 1200.
- 58 S. Reichert, Q. An, Y. W. Woo, A. Walsh, Y. Vaynzof and C. Deibel, *Nat. Commun.*, 2020, **11**, 6098.
- 59 S. Kim, J. H. Jang, Z. Wu, M. J. Lee, H. Y. Woo and I. Hwang, *Small*, 2021, **17**, 2101839.
- 60 C. Eames, J. M. Frost, P. R. F. Barnes, B. C. O'Regan, A. Walsh and M. S. Islam, *Nat. Commun.*, 2015, **6**, 7497.
- 61 M. H. Futscher, M. K. Gangishetty, D. N. Congreve and B. Ehrler, *J. Chem. Phys.*, 2020, **152**, 044202.
- 62 P. Srivastava, R. Kumar and M. Bag, *Phys. Chem. Chem. Phys.*, 2021, **23**, 10936–10945.
- 63 A. D. Sheikh, R. Munir, M. A. Haque, A. Bera, W. Hu, P. Shaikh, A. Amassian and T. Wu, *ACS Appl. Mater. Interfaces*, 2017, **9**, 35018–35029.
- 64 A. M. Soufiani, Z. Hameiri, S. Meyer, S. Lim, M. J. Y. Tayebjee, J. S. Yun, A. Ho-Baillie, G. J. Conibeer, L. Spiccia and M. A. Green, *Adv. Energy Mater.*, 2017, **7**, 1602111.
- 65 A. D. Bui, N. Mozaffari, T. N. Truong, T. Duong, K. J. Weber, T. P. White, K. R. Catchpole, D. Macdonald and H. T. Nguyen, *Prog. Photovoltaics*, 2022, **30**, 1038–1044.
- 66 A. A. B. Baloch, F. H. Alharbi, G. Grancini, M. I. Hossain, M. K. Nazeeruddin and N. Tabet, *J. Phys. Chem. C*, 2018, **122**, 26805–26815.
- 67 M. Lira-Cantú, *Nat. Energy*, 2017, **2**, 17115.
- 68 Z. Fan, H. Xiao, Y. Wang, Z. Zhao, Z. Lin, H.-C. Cheng, S.-J. Lee, G. Wang, Z. Feng, W. A. Goddard, Y. Huang and X. Duan, *Joule*, 2017, **1**, 548–562.
- 69 E. J. Juarez-Perez, L. K. Ono, M. Maeda, Y. Jiang, Z. Hawash and Y. Qi, *J. Mater. Chem. A*, 2018, **6**, 9604–9612.
- 70 A. Senocrate, G. Y. Kim, M. Grätzel and J. Maier, *ACS Energy Lett.*, 2019, **4**, 2859–2870.
- 71 T. W. Kim, N. Shibayama, L. Cojocar, S. Uchida, T. Kondo and H. Segawa, *Adv. Funct. Mater.*, 2018, **28**, 1804039.
- 72 D. B. Khadka, Y. Shirai, M. Yanagida and K. Miyano, *ACS Appl. Energy Mater.*, 2021, **4**, 11121–11132.
- 73 I. S. Zhidkov, D. W. Boukhalov, A. F. Akbulatov, L. A. Frolova, L. D. Finkelstein, A. I. Kukharenko, S. O. Cholakh, C.-C. Chueh, P. A. Troshin and E. Z. Kurmaev, *Nano Energy*, 2021, **79**, 105421.



- 74 S. R. Raga, L. K. Ono and Y. Qi, *J. Mater. Chem. A*, 2016, **4**, 2494–2500.
- 75 B. Yang, O. Dyck, W. Ming, M.-H. Du, S. Das, C. M. Rouleau, G. Duscher, D. B. Geohegan and K. Xiao, *ACS Appl. Mater. Interfaces*, 2016, **8**, 32333–32340.
- 76 T. Guo, P. Hu, L. Li, Z. Wang and L. Guo, *Chem*, 2023, **9**, 1080–1093.
- 77 C. Bommier, D. Mitlin and X. Ji, *Prog. Mater. Sci.*, 2018, **97**, 170–203.
- 78 D. A. Jacobs, Y. Wu, H. Shen, C. Barugkin, F. J. Beck, T. P. White, K. Weber and K. R. Catchpole, *Phys. Chem. Chem. Phys.*, 2017, **19**, 3094–3103.
- 79 D. Meggiolaro, E. Mosconi and F. De Angelis, *ACS Energy Lett.*, 2019, **4**, 779–785.

