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# Elucidating the efficiency limit of silicon-based monolithic tandem cells through the combination of Auger and Shockley–Queisser limits

 Andreas Fell,<sup>ID \*ab</sup> Oliver Fischer,<sup>ID ac</sup> Martin Bivour,<sup>a</sup> Christoph Messmer,<sup>ID ac</sup> Jonas Schön,<sup>ac</sup> Martin C. Schubert<sup>a</sup> and Stefan W. Glunz<sup>ac</sup>

Accurate theoretical efficiency limits are critical for diagnosing loss mechanisms and guiding optimization in solar cell technologies. While the Shockley–Queisser (SQ) limit remains the most widely used framework for assessing tandem and multijunction devices, its assumptions—purely radiative recombination and ideal light absorption—do not account for the intrinsic limitations of silicon (Si), the dominant photovoltaic material. In particular, Si's indirect bandgap resulting in Auger recombination imposes a lower efficiency ceiling. In this work, we present a rigorous simulation approach that combines SQ-limited top cells with an Auger-limited Si bottom cell, accounting also for luminescent coupling (LC). This hybrid modeling approach yields a maximum theoretical efficiency of 43.2% for an ideal two-terminal Si-based tandem device, compared to 45.2% using the unrealistic assumption of a SQ-limited Si bottom cell. The optimal configuration features a top cell bandgap of 1.71 eV and a 300  $\mu\text{m}$ -thick Si bottom cell, with a minor efficiency penalty of only 0.1% for a more typical thickness of 120  $\mu\text{m}$ . Accounting for LC values typical for perovskite top cells reduces the optimum efficiency to 42.4%. Special emphasis is placed on the interpretation of fill factor (FF), highlighting the need for correct analytical FF limit ( $\text{FF}_0$ ) calculations using an appropriate ideality factor, which is 5/3 for silicon based tandem at the theoretical limit. To support future benchmarking, we provide lookup tables of current–voltage ( $JV$ ) parameters for a range of top cell bandgaps, bottom cell properties, multijunction stacks with up to six subcells, and perovskite-specific top cell properties. These results offer reliable efficiency limits for the evaluation of high-efficiency silicon-based tandem and multijunction solar cells.

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## Broader context

Silicon-perovskite tandem solar cells represent the most promising path to push photovoltaic efficiencies beyond the single-junction limit, and are widely regarded as the next-generation solar technology. However, theoretical benchmarks typically rely on the Shockley–Queisser (SQ) limit, which does not accurately capture the intrinsic loss mechanisms of silicon—such as Auger recombination and weak absorption close to the bandgap. This leads to overestimated efficiency targets and hampers meaningful loss analysis and device optimization. In this work, we rigorously combine the SQ-limit model for the top cell/top cells with the latest established Auger-limit model for the silicon bottom cell, enabling physically sound efficiency limits for silicon-based tandem and multijunction architectures. We provide detailed insights into the role of fill factor and luminescent coupling, and publish tabulated  $JV$  parameters across a range of configurations. By offering a more accurate and future-proof benchmark for what is likely to remain the dominant tandem architecture for years to come, this work is expected to serve as a widely referenced standard—much like the established single-junction silicon limits by (Niewelt *et al.*, *Sol. Energy Mater. Sol. Cells*, 2022, 235, 111467) and (Richter *et al.*, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2012, 86, 16, 165202). Our results will support researchers and developers in setting realistic targets and interpreting experimental performance with higher confidence.

## Introduction

Theoretical performance limits play a critical role in guiding the development of photovoltaic technologies. By defining the

maximum achievable power conversion efficiency (PCE), short-circuit current density ( $J_{sc}$ ), open-circuit voltage ( $V_{oc}$ ) and fill factor (FF) under idealized conditions, such limits provide a benchmark for identifying dominant loss mechanisms in experimental devices for directing their optimization pathways. The most widely used theoretical framework is the detailed balance limit proposed by Shockley and Queisser (SQ), which describes the efficiency ceiling for a single-junction solar cell as a function of the absorber material's band gap.<sup>1</sup> It assumes radiative recombination as the sole electrical loss mechanism together with unity absorptance of all photons

<sup>a</sup>Fraunhofer Institute for Solar Energy Systems ISE, 79110 Freiburg, Germany. E-mail: andreas.fell@ise.fraunhofer.de

<sup>b</sup>AF Simulations GmbH, Landstr. 33a, 79232 March, Germany

<sup>c</sup>Chair for Photovoltaic Energy Conversion, Department of Sustainable Systems Engineering (INATECH), University of Freiburg, Emmy-Noether-Str. 2, 79110 Freiburg, Germany



with energies above the bandgap. While the SQ model provides valuable insight, especially for direct bandgap materials, it falls short in accurately representing the performance potential of silicon (Si), the dominant material in the photovoltaic industry.

Silicon's indirect bandgap introduces additional intrinsic losses, most notably Auger recombination, and also weak absorption near the band edge, necessitating thick absorbers and efficient light management schemes. These characteristics are not captured in the idealized assumptions of the SQ model, but more elaborate electro-optical device models are required to calculate intrinsic performance limits of Si solar cells. This so-called "Auger-limit" of a single-junction silicon solar cell was most recently quantified to 29.4% by Niewelt *et al.*<sup>2</sup> While the Auger-limit models are well-established for silicon single-junction cells, this more accurate limit is to date rarely applied when evaluating multijunction solar cell architectures that incorporate silicon as a bottom cell. One reason is the non-trivial combination of the analytical SQ model for the top cell with the electro-optical device simulation model of the Si bottom cell. Furthermore, there is a lack of comprehensive, tabulated data for current-voltage ( $JV$ ) parameters besides efficiency of Si based tandem cells – even within the SQ framework.<sup>3–5</sup> Instead, the SQ framework resulting in an optimal efficiency of 45.2%, assuming a SQ-limited bottom cell with a bandgap of 1.12 eV,<sup>6</sup> is still commonly used to assess Si-based tandem cells. This leads to inconsistencies in theoretical benchmarks and hampers efforts to concisely diagnose and optimize real-world devices.

Almansouri *et al.*<sup>7</sup> first reported modeling results combining an Auger limited Si bottom cell with a SQ limited top cell, investigating a range of practical bottom cell properties. However, luminescent coupling (LC) is not considered, and no results for the specific case of a perovskite top cell are given. The study of Bowman *et al.*<sup>8</sup> focuses on analytical modeling of LC in perovskite tandem cells, and also give a maximum efficiency for an Auger limited Si bottom cell combined with an optically realistic perovskite top cell absorber as 42.0% with, and 43.0% without LC, respectively. Allen *et al.*<sup>9</sup> used a simplified equivalent-circuit model, with neglect of LC, of various Si bottom cells combined with an optically realistic perovskite top cell absorber to derive an efficiency limit of 42.45%, and a practical target of 37.8%.

Besides the various simplifications, all these previous works do not list  $JV$  parameters of the performance limits, except single values for efficiency for specific assumptions.

In this work, we rigorously simulate the performance limits of tandem and multijunction solar cells, under the assumption of SQ-limit top cells and an Auger-limit silicon bottom cell. Our model bridges detailed balance and thickness-dependent device modeling to provide accurate, physically grounded efficiency limits along with complete  $JV$  parameter sets, made available as lookup tables. Particular emphasis is placed on FF, a parameter that is easily misinterpreted in tandem analysis<sup>10</sup> due to (i) a current mismatch FF boost effect,<sup>11</sup> (ii) incorrect analytical FF limit ( $FF_0$ ) calculation using the tandem-level open-circuit voltage as per the popular formula by Green

*et al.*<sup>12</sup> with an ideality factor  $n = 1$ , and (iii) the fact that Si bottom cells can show ideality factors below unity due to Auger recombination which increases their theoretical  $FF_0$  limit.<sup>13</sup> We further quantify the role of LC in the theoretical limit, extend our analysis to multijunction architectures with Si bottom cells, and present efficiency limits for optically realistic perovskite top cells.

## Model

The top cell is modeled using the analytical detailed balance framework introduced by Shockley and Queisser.<sup>1</sup> The model assumes a step-function absorption at the bandgap energy and radiative recombination as the only loss mechanism. The resulting radiative recombination current density  $J_{r,SQ}$  as a function of device voltage  $V$  and bandgap  $E_g$  is given by eqn (4) in Rühle:<sup>14</sup>

$$J_{r,SQ}(V, E_g) = f_g q \int_{E_g}^{\infty} \frac{2\pi E^2}{h^3 c^2} \frac{1}{\exp\left(\frac{E - qV}{kT}\right) - 1} dE.$$

Here,  $k$  is Boltzmann's constant,  $T$  is the temperature,  $h$  is Planck's constant,  $c$  is the speed of light,  $f_g$  is the geometrical factor and  $q$  is the elementary charge. Neglecting the "–1" in the denominator, which is valid up to SQ-limit open-circuit voltages that remain several  $kT$  below  $E_g$ , this can be rearranged and analytically integrated to:

$$J_{r,SQ}(V, E_g) = f_g J_{01,SQ,mf}(E_g) \exp\left(\frac{qV}{kT}\right),$$

with  $J_{01,SQ,mf}$  being the dark saturation current density at the SQ limit quantifying the front-side radiative emission loss (representing a monofacial "mf" case):

$$J_{01,SQ,mf}(E_g) = \frac{2\pi q}{h^3 c^2} [kTE_g^2 + 2(kT)^2 E_g + 2(kT)^3].$$

The geometrical factor  $f_g$  quantifies the total amount of radiative emission relative to the monofacial front side emission. It equals 1 for zero rear-side emission, *i.e.* a monofacial cell with a perfect rear reflector. Most commonly in SQ calculations a symmetric bifacial cell is assumed, resulting in equal front and rear emission and thus  $f_g = 2$ . For typical monolithic tandem cells however, the top cell is inherently asymmetric, as the upper surrounding material (air) has a different refractive index than the bottom cell material. This results in different internal reflection, and means that  $f_g = 2$  is a bad assumption. As discussed *e.g.* in Bowman *et al.*,<sup>8</sup> the difference in escape cones leads to an approximately  $n_{top}^2$  higher rear emission relative to the front emission, with  $n_{top}$  being the refractive index of the top cell absorber material. Assuming negligible parasitic absorption and thus unity band-to-band absorption in the bottom cell of the rear emitted photons – a good assumption for high-efficient practical cells – the rear side emission equals the additional bottom cell generation by LC. By defining the LC efficiency  $\eta_{lc}$  as the ratio of rear emission over total radiative emission, it is linked to the geometrical factor:



$$f_g = \frac{1}{1 - \eta_{lc}}$$

A direct consequence is that higher LC efficiency means substantially higher radiative emission and recombination in the top cell, and thus lower top cell power. This loss cannot be overcompensated by the additional power generation in the bottom cell due to its lower voltage, unless for strong current-mismatch conditions (see Results section).

For our calculations, we distinguish two cases for rear emission and thus LC:

(a) “Monofacial” top cell with  $\eta_{lc} = 0$  and  $f_g = 1$ , which represents the upper theoretical limit; this is theoretically possible in a tandem configuration, by assuming an ideal dichroic mirror reflecting all radiatively emitted photons above the top cell bandgap energy, while being fully transparent to lower energy external photons; such intermediate reflectors are also implemented in experimental devices (see *e.g.* Callies *et al.*<sup>15</sup>);

(b) Top cell optically well coupled to the bottom cell for the entire spectrum, with an exemplary refractive index representative for a typical perovskite absorber of  $n_{\text{top}} = 2.58$  (Manzoor *et al.*<sup>16</sup>), resulting in  $\eta_{lc} = n_{\text{top}}^2 = 87\%$  and  $f_g = 7.7$ ; notably, this theoretical value is close to the 85% measured by Nguyen *et al.*<sup>17</sup> and also confirmed by our ray-tracing simulations (see Results section), validating it to be a useful general value for a perovskite silicon tandem cell.

Nonradiative recombination in the top cell is quantified using the external radiative efficiency (ERE).<sup>18</sup> For a monofacial SQ top cell we define the monofacial ERE  $\eta_{\text{ext,mf}}$ , which relates front side emission  $J_{01,\text{rfs}}$  and nonradiative recombination  $J_{01,\text{nr}}$  by

$$\eta_{\text{ext,mf}} = \frac{J_{01,\text{SQ,mf}}}{J_{01,\text{SQ,mf}} + J_{01,\text{nr}}}$$

We apply this definition also in the case of non-zero rear-side emission. In this situation,  $\eta_{\text{ext,mf}}$  no longer strictly quantifies the ERE as the ratio of total emission to total generation. However, the advantage is that a given  $\eta_{\text{ext,mf}}$  value corresponds to the same  $J_{01,\text{nr}}$ , independently of rear-side emission, *i.e.*, both with and without LC. In this way,  $\eta_{\text{ext,mf}}$  becomes a quantity more directly linked to the nonradiative recombination properties of the top cell material, consistent with the common intuitive interpretation of ERE.

Combined with radiative recombination  $J_{01,\text{r}}$ , the total top cell recombination  $J_{01}$  is thereby given as

$$J_{01} = J_{01,\text{r}} + J_{01,\text{nr}} = J_{01,\text{SQ,mf}} \left( \frac{1}{1 - \eta_{lc}} + \frac{1 - \eta_{\text{ext,mf}}}{\eta_{\text{ext,mf}}} \right).$$

The silicon bottom cell is modeled using the most recent and comprehensive single-junction Auger-limited efficiency model, as established by Niewelt *et al.*,<sup>2</sup> yielding a theoretical efficiency limit of 29.4%. This model includes the latest data for the silicon absorption coefficient,<sup>19</sup> combined with Lambertian

light-trapping,<sup>20</sup> the Auger parameterization of Niewelt *et al.*,<sup>2</sup> and radiative recombination by Fell *et al.*,<sup>21</sup> which incorporates Lambertian photon recycling. It further uses a consistent set of silicon semiconductor device properties—including bandgap, bandgap narrowing (BGN) and density of states (DOS)—as summarized in ref. 21. As in the single-junction Auger-limit of 29.4%, we consider intrinsic, *i.e.* undoped, Si only in this work. It was shown by Niewelt *et al.* that typical doping levels both for n-type and p-type of around 1  $\Omega$  cm result only in a moderate efficiency penalty of  $\sim 0.2\%$ <sub>abs</sub> due to increased Auger recombination. This penalty is further reduced within a tandem cell due to the lower contribution to output power, rendering it a well negligible loss.

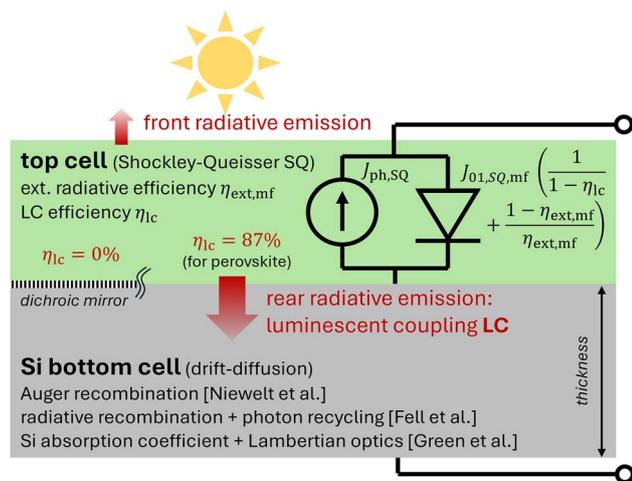
While this model incorporates a variety of physical mechanisms, it is referred to as the “Auger-limit” for simplicity, consistent with the literature, due to the dominant contribution of Auger recombination relative to the SQ-limit. In contrast to the top cell model, the Si bottom cell performance is inherently thickness-dependent, requiring numerical optimization to balance light absorption and recombination losses.

All models are implemented within the solar cell simulation software Quokka3 (ref. 22 and 23). It was used in Niewelt *et al.*<sup>2</sup> to compute Auger-limit of single-junction silicon devices, and therefore the latest Auger limit models are exactly incorporated in our calculations. Within this work, the tandem capabilities of Quokka3 have been extended to include a Shockley–Queisser-limited top cell, both in terms of optics (step-function absorption) and recombination (detailed balance radiative recombination limit with adjustable  $\eta_{lc}$  and  $\eta_{\text{ext,mf}}$ ). Notably, the bottom cell is solved by a 1D numerical drift-diffusion model instead of the 0D analytical model in previous works. This work’s model accounts for transport losses, which become noticeable for the high Si bottom cell thicknesses optimal for bottom cell limited tandems. This substantially improves accuracy of optimal bottom cell thickness prediction towards lower values. The simulation outputs include full *JV* characteristics, providing  $J_{\text{sc}}$ ,  $V_{\text{oc}}$  and FF. In addition, we compute suns- $V_{\text{oc}}$  curves of the tandem device which gives the pseudo fill factor (pFF) derived from this pseudo *JV*-curve. As no current is flowing, this pFF is free from the well-known FF boost arising from current-mismatch.<sup>11</sup> We note that in the idealized model of this work, transport, selectivity and shunt losses are negligible, which results in identical values for implied fill factor (iFF), pFF and FF, the latter however only at exact current-match conditions.

An overview of the tandem model of this work is given by the sketch in Fig. 1.

In order to model optics of a perovskite top cell with realistic absorption properties, we further employ ray-tracing simulations in Sentaurus Device, with the baseline tandem cell properties and modeling methodology described in Messmer *et al.*<sup>24</sup> In a similar approach as described in Allen *et al.*,<sup>9</sup> we scale the optical data of Manzoor *et al.*<sup>16</sup> along the photon energy axis in order to calculate wavelength-dependent absorption coefficients for varying bandgaps. Notably, this approach has been proven accurate to match experimental external quantum efficiency (EQE) measurements for a range of perovskite compositions and bandgaps,<sup>24–27</sup> including in this work (good match of





**Fig. 1** Sketch of the tandem model used in this work, which combines a detailed balance model of the top cell (“SQ-limit”) with the latest set of models used for calculating Si single junction efficiency limits (“Auger-limit”); nonradiative recombination in the top cell is considered by the external radiative efficiency assuming monofacial conditions  $\eta_{ext,mf}$ ; luminescent coupling LC is considered as both increased rear side emission (*i.e.* radiative recombination) in the top cell and additional generation current in the bottom cell.

EQE shape in Fig. 4a). It can therefore be considered a generally valid way of modeling absorption in perovskite top cells, rendering the perovskite-specific efficiency limits presented in this work well future-proof, instead of being particular to a currently best specific material composition.

We simulate the EQE of a perovskite top cell with varying thickness and bandgap, assuming negligible recombination losses at short-circuit so that optical absorption in the perovskite equals EQE. By correcting the EQE with the simulated reflection and parasitic absorption, we calculate the IQE. The IQE thereby presents an idealized top cell absorption free from reflection and parasitic absorption losses. This is consistent with the SQ assumptions, with the only distinction that a realistic optical absorption of perovskite material is used, revealing its fundamental optical limit. We note that this is more accurate than previous models using Lambert–Beer<sup>9</sup> or geometrical light-trapping models<sup>8</sup> for predicting absorption, as our approach includes transfer-matrix-method (TMM) calculations for the optically thin perovskite layer. By numerically integrating the simulated IQE using the generalized Planck-law, we can quantify associated radiative emission and recombination in a fully consistent manner as described above:

$$J_{01} = J_{01,r}(\text{IQE}) = f_g q \int_0^{\infty} \frac{2\pi E^2}{h^3 c^2} \frac{\text{IQE}}{\exp\left(\frac{E}{kT}\right)} dE$$

The equation above assumes that IQE is angle-independent, which is not obviously the case even for a textured Si bottom cell, as the conformal coating of the perovskite top cell might well produce significant angular dependence. We validate this by randomizing start directions in our ray-tracing simulations,

and find only minor difference in calculated  $J_{01}$  of <2%. Notably,  $f_g$  (and LC) are not free parameters but determined by the assumed optical properties. We quantify rear side top cell IQE by starting randomized rays within the Si wafer towards the top cell, and calculate respective  $J_{01}$ . The ratio of front and rear  $J_{01}$  gives  $\eta_{lc} = 87\%$ , accurately matching the approximate assumption described above. IQE and  $J_{01}$  are then inputted into Quokka3 for tandem performance simulations.

## Results and discussion

### Bottom cell thickness optimization

The optimal thickness of the silicon bottom cell in a tandem architecture is found to be greater than that of a single-junction device. This is due to the higher relative importance of current density in the tandem configuration, and therefore more importance of near-infrared (NIR) absorption. However, the overall efficiency sensitivity to thickness variations remains low. We find an optimal thickness of approximately 300  $\mu\text{m}$ , while typical industrial thicknesses in the range of 100–150  $\mu\text{m}$  incur a minor efficiency penalty of only  $\sim 0.1\%_{\text{abs}}$  (see Fig. 2a). The optimum thickness is lower than in previous studies,<sup>7–9</sup> caused by the increased accuracy of this work’s bottom cell model, which accounts for unavoidable transport losses becoming relevant at large thicknesses. For all subsequent simulations, we fix the silicon bottom cell thickness to 300  $\mu\text{m}$ . Although the true optimum would vary with top cell bandgap, we omit bandgap-specific thickness optimization because the associated performance improvements are marginal and the resulting optimal thicknesses in bottom-cell-limited cases would become unrealistically large.

### Top cell bandgap variation

The efficiency as a function of the top cell bandgap is shown in (Fig. 2b). The maximum efficiency is 43.2% at a bandgap of 1.71 eV when assuming  $\eta_{lc} = 0$ , *i.e.* an ideal dichroic intermediate mirror. 1.71 eV corresponds quite closely, but not exactly, to photogeneration current-matching condition (see Fig. 2c), as expected due to optimum efficiency being achieved at  $J_{mpp}$  rather than  $J_{sc}$  match. When assuming  $\eta_{lc} = 87\%$ , the optimum efficiency is reduced to 42.4%, with a bandgap at slight bottom cell limitation. The efficiency penalty arises from the increased top cell radiative emission and therefore lower  $V_{oc}$  (see Fig. 2d). As expected, LC plays a substantial role in bottom-cell-limited devices, where the additional bottom cell current increases the overall  $J_{sc}$  (see Fig. 2c). A positive effect of LC is a much lower dependence of efficiency on the bandgap for bottom cell limited devices, which is linked to a higher tolerance against current mismatch. Notably, these findings regarding LC are qualitatively and quantitatively consistent with the findings of Bowman *et al.*<sup>8</sup> Furthermore, the impact of the LC effect in providing additional bottom cell current is separated from the mentioned  $V_{oc}$  loss by the “SQ, no LC” line in Fig. 2b, for which only the additional generated current was disabled in Quokka3. While physically these two effects are fundamentally linked, the neglected bottom cell generation is often used due to lower



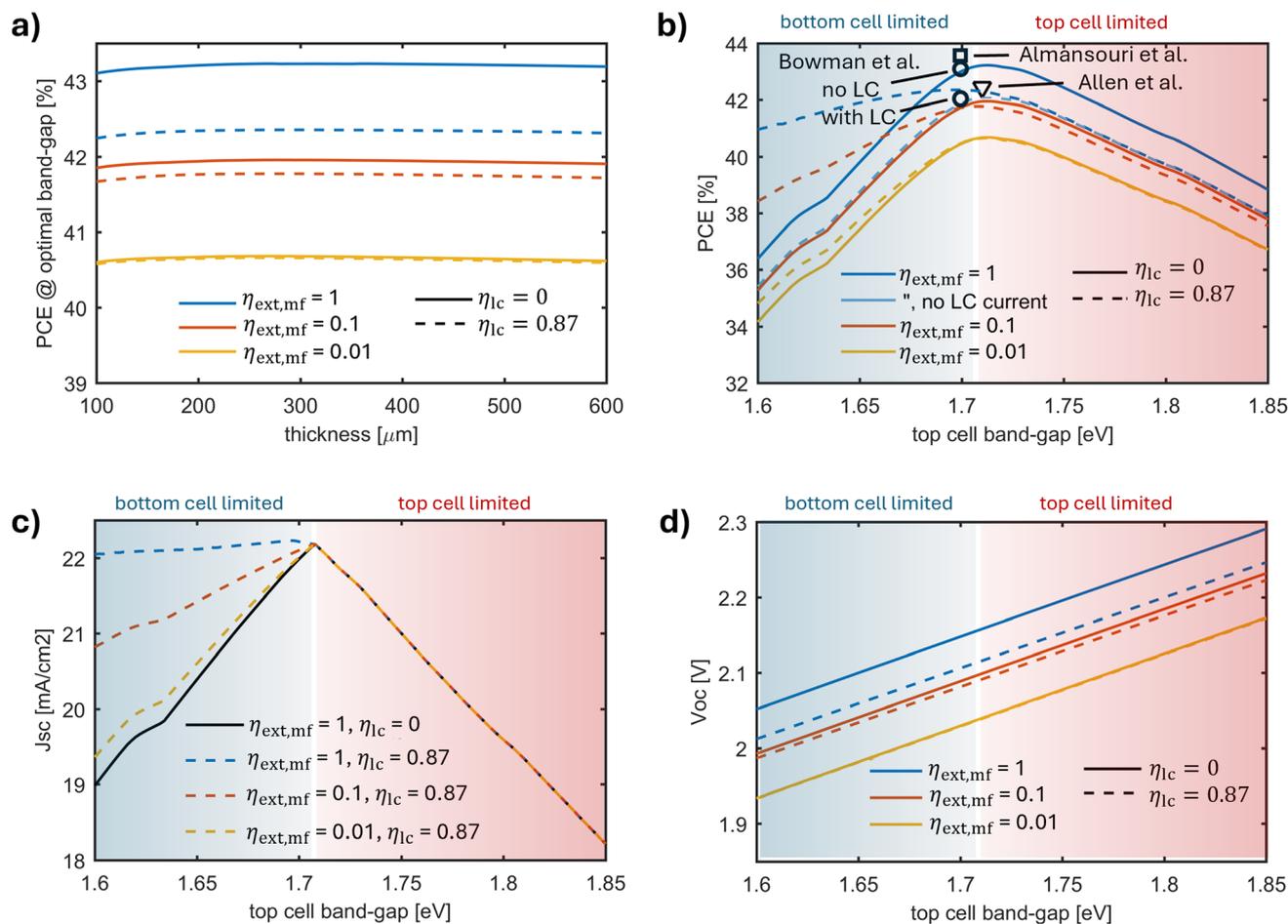


Fig. 2 Performance limit of a tandem cell combining a Auger-limit silicon bottom cell and a top cell with step-function absorption; varied is the top cell bandgap and its nonradiative recombination properties via the monofacial external radiative efficiency  $\eta_{\text{ext,mf}}$  (a value of 1 represents the SQ limit); influence of luminescent coupling (LC) is shown by assuming two different coupling efficiencies  $\eta_{\text{lc}}$ , and one simulation with switched-off bottom cell generation ( $\eta_{\text{ext,mf}} = 1$ , no LC current); (a) variation of bottom cell thickness showing minor impact of broad variation around the optimum of 300  $\mu\text{m}$ ; (b) power conversion efficiency (PCE) compared with previously published limits (Bowman *et al.* and Allen *et al.* using realistic perovskite absorption data instead of SQ limit); (c) short-circuit current density  $J_{\text{sc}}$ , showing the impact of LC and current-match; (d) open circuit voltage  $V_{\text{oc}}$ , showing substantial decrease for the high LC case due to increased top cell radiative recombination loss via rear-side emission.

model implementation effort.<sup>7,9,24</sup> This simplification is shown to be valid for top cell limited devices, and introduces only a minor error for current-match conditions, but fails for bottom cell limited devices.

Fig. 2b also illustrates the efficiency penalty resulting from non-unity (*i.e.*, below SQ limit) external radiative efficiency (ERE) of the top cell. The influence of LC decreases significantly with decreasing  $\eta_{\text{ext,mf}}$ . For  $\eta_{\text{ext,mf}}$  values below 10%, LC shows a small impact on  $V_{\text{oc}}$ , and below 1% it becomes an overall minor effect even for bottom-cell limited devices. Notably, this holds for performance and  $JV$  characteristics, and does not preclude significant impact of LC for specific measurement conditions, like *e.g.* quantum efficiency measurements.

As found earlier, the monolithic efficiency optimum is only slightly below the optimum in a four-terminal tandem configuration,<sup>3,5</sup> which we recalculate in this work to be 43.3% for an Auger-limit silicon bottom cell and a 1.73 eV top cell.

Tabular data for the optimal bottom cell thickness of 300  $\mu\text{m}$  for the two LC cases, as well as for a typical thickness of 120  $\mu\text{m}$ , are given in Tables 1–3.

### Fill factors – FF, pFF and FF<sub>0</sub>

In Fig. 3a various fill factors (FF) are shown as a function of bandgap. A well-known FF dependence on current mismatch<sup>11</sup> can clearly be seen, *i.e.* the FF rises for increasing current-mismatch and has a steep minimum at the current-match condition. A mismatch of 0.5  $\text{mA cm}^{-2}$  gives already a FF boost of around 1%<sub>abs</sub>. FF is also moderately influenced by LC. The pFF from suns- $V_{\text{oc}}$  curve on the other hand (which is identical to the implied fill factor iFF as the investigated ideal device is free from transport losses) shows a smooth linear increase with bandgap.

It is popular to compare the FF or pFF to the so-called FF<sub>0</sub> introduced by Green *et al.*<sup>12</sup> in order to identify FF losses which



**Table 1** JV parameters of tandem device with Auger limited silicon bottom cell and SQ limited top cell at the optimal bottom cell thickness of 300  $\mu\text{m}$  and  $\eta_{\text{ic}} = 0\%$

Bandgap [eV]	PCE [%]	$V_{\text{oc}}$ [V]	$J_{\text{sc}}$ [ $\text{mA cm}^{-2}$ ]	FF [%]	pFF [%]
1.58	34.7	2.03	18.3	93.6	90.0
1.6	36.4	2.05	19.0	93.4	90.0
1.62	37.9	2.07	19.6	93.2	90.1
1.64	39.0	2.09	20.1	93.1	90.2
1.66	40.6	2.11	20.7	92.8	90.2
1.68	42.0	2.13	21.4	92.2	90.3
1.7	43.0	2.15	22.0	91.1	90.4
<b>1.712</b>	<b>43.23</b>	<b>2.159</b>	<b>22.08</b>	<b>90.69</b>	<b>90.41</b>
1.72	43.2	2.17	21.9	91.1	90.4
1.74	42.8	2.19	21.3	91.8	90.5
1.76	42.1	2.21	20.7	92.2	90.6
1.78	41.4	2.22	20.1	92.5	90.6
1.8	40.7	2.24	19.6	92.7	90.7
1.82	40.0	2.26	19.1	92.8	90.8
1.84	39.2	2.28	18.5	93.0	90.8
1.86	38.4	2.30	17.9	93.1	90.9

**Table 2** JV parameters of tandem device with Auger limited silicon bottom cell and SQ limited top cell at the optimal bottom cell thickness of 300  $\mu\text{m}$  and  $\eta_{\text{ic}} = 87\%$

Bandgap [eV]	PCE [%]	$V_{\text{oc}}$ [V]	$J_{\text{sc}}$ [ $\text{mA cm}^{-2}$ ]	FF [%]	pFF [%]
1.58	40.5	1.99	22.0	92.4	89.8
1.6	40.9	2.01	22.0	92.3	89.9
1.62	41.3	2.03	22.1	92.2	90.0
1.64	41.7	2.05	22.1	92.1	90.0
1.66	42.0	2.07	22.1	91.8	90.1
1.68	42.3	2.09	22.2	91.3	90.2
1.7	<b>42.4</b>	<b>2.11</b>	<b>22.2</b>	<b>90.5</b>	<b>90.3</b>
1.72	42.1	2.13	21.9	90.8	90.3
1.74	41.7	2.14	21.3	91.3	90.4
1.76	41	2.16	20.7	91.7	90.4
1.78	40.4	2.18	20.1	92.0	90.5
1.8	39.7	2.2	19.6	92.2	90.6
1.82	39.1	2.22	19.1	92.4	90.6
1.84	38.3	2.24	18.5	92.5	90.7
1.86	37.5	2.26	17.9	92.7	90.8

are not related to ideal recombination losses corresponding to the measurable  $V_{\text{oc}}$ . The derivation of the  $\text{FF}_0$  formula assumes a single-diode model with fixed ideality factor  $n$ , which is most commonly assumed to be 1.

$$\text{FF}_0 = \frac{v_{\text{oc}} - \ln(v_{\text{oc}} + 0.72)}{v_{\text{oc}} + 1},$$

with

$$v_{\text{oc}} = \frac{qV_{\text{oc}}}{nkT}.$$

It is stressed that for multijunction devices, the  $\text{FF}_0$  formula must not directly be applied with  $n = 1$ . For a series-connection of multiple cells, the ideality factor of a single diode model representing the multijunction device instead equals the sum of

**Table 3** JV parameters of tandem device with Auger limited silicon bottom cell and SQ limited top cell at a bottom cell thickness of 120  $\mu\text{m}$  and  $\eta_{\text{ic}} = 87\%$

Bandgap [eV]	PCE [%]	$V_{\text{oc}}$ [V]	$J_{\text{sc}}$ [ $\text{mA cm}^{-2}$ ]	FF [%]	pFF [%]
1.58	40.2	2.01	21.6	92.4	89.9
1.6	40.6	2.03	21.7	92.3	90.0
1.62	41.0	2.05	21.7	92.3	90.0
1.64	41.4	2.06	21.7	92.2	90.1
1.66	41.6	2.08	21.8	91.8	90.2
1.68	41.9	2.10	21.8	91.6	90.2
1.7	42.2	2.12	21.9	91.1	90.3
<b>1.712</b>	<b>42.3</b>	<b>2.13</b>	<b>21.9</b>	<b>90.7</b>	<b>90.3</b>
1.72	42.2	2.14	21.8	90.4	90.4
1.74	41.9	2.16	21.3	91.1	90.4
1.76	41.3	2.18	20.7	91.6	90.5
1.78	40.6	2.20	20.1	91.9	90.6
1.8	40.0	2.21	19.6	92.2	90.6
1.82	39.3	2.23	19.1	92.4	90.7
1.84	38.5	2.25	18.5	92.5	90.7
1.86	37.8	2.27	17.9	92.7	90.8

the individual cells' ideality factors, as can be shown *e.g.* for the case of a tandem (subscript tnd):

$$\begin{aligned} V_{\text{tnd}} &= \frac{n_{\text{tnd}}kT}{q} \ln\left(\frac{J}{J_{0,\text{tnd}}}\right) = V_{\text{bot}} + V_{\text{top}} \\ &= \frac{n_{\text{top}}kT}{q} \ln\left(\frac{J}{J_{0,\text{top}}}\right) + \frac{n_{\text{bot}}kT}{q} \ln\left(\frac{J}{J_{0,\text{bot}}}\right) \end{aligned}$$

from which follows by basic algebra:

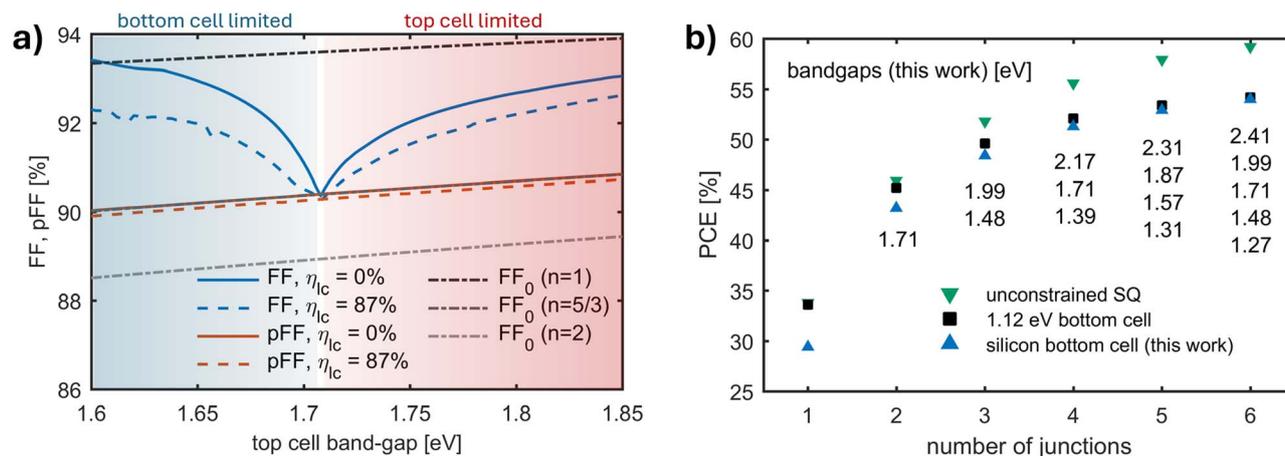
$$n_{\text{tnd}} = n_{\text{top}} + n_{\text{bot}}$$

and

$$J_{0,\text{tnd}} = J_{0,\text{top}}^{\frac{n_{\text{top}}}{n_{\text{tnd}}}} J_{0,\text{bot}}^{\frac{n_{\text{bot}}}{n_{\text{tnd}}}}.$$

Wrongly applying  $n = 1$  for a tandem device to calculate  $\text{FF}_0$  results in strong overestimation of achievable FF. Using  $n = 2$  would be appropriate for a tandem device with both sub cells following ideal  $n = 1$  recombination. For an Auger-limited Si bottom cell however, the ideality factor instead is  $2/3$ , see *e.g.* Sinton *et al.*,<sup>13</sup> and consequently  $n = 5/3$  must be assumed when calculating  $\text{FF}_0$  for a tandem cell with an Auger-limited Si bottom cell. This is evidenced by the excellent agreement of  $\text{FF}_0$  ( $n = 5/3$ ) with the simulation results in Fig. 3a. Notably, while record single-junction silicon cells indeed showing ideality factors below 1, typical experimental Si bottom cells are not yet largely Auger dominated, and so for experimental silicon based tandem cells  $n = 2$  could be more appropriate when calculating  $\text{FF}_0$ . Most accurate would be to first determine the ideality factor of the silicon bottom cell under 0.5 suns generation, by matching  $\text{FF}_0$  with the measured iFF of a single-junction or subcell-selective measurement. Subsequently, this ideality factor can be increased by 1 and then be applied to the tandem cell  $V_{\text{oc}}$  to calculate the  $\text{FF}_0$  of the experimental tandem cell. As an approximative alternative, a lookup for tandem efficiency and JV parameter limits for the case of a given silicon bottom





**Fig. 3** (a) fill factors as a function of top cell bandgap and at its radiative limit; the FF shows strongly the well-known current mismatch artifact, while pFF is artifact free and also not impacted by LC;  $FF_0$  calculated from the device'  $V_{oc}$  matches pFF only when using an ideality factor of  $n = 5/3$ ; (b) multijunction cell efficiency limit comparing the unconstrained Shockley Queisser (SQ) limit with the cases a 1.12 eV SQ-limit and an Auger-limit silicon bottom cell; adaption of Fig. 1 of Schygulla *et al.*; the bandgap values shown are the optimal bandgaps for an Auger-limit silicon bottom cell as calculated in this work.

cell with a known  $V_{oc}$  below the Auger limit is given in Table 4. Here we introduce an extrinsic recombination contribution  $J_{01,extr}$  to the silicon bottom cell with a typical thickness of 120  $\mu\text{m}$ . It is thereby assumed that the extrinsic recombination has an ideality factor of 1, which can be considered a reasonable approximation for typical high-efficiency bottom cells. It can be seen that the ideality factor, which fulfills  $FF_0 = FF$  for the bottom cell, quickly approaches 1 with increasing extrinsic recombination. Notably, with zero extrinsic recombination, *i.e.* at the Auger-limit, the ideality factor does not exactly equal  $2/3$  but 0.71. This is partly due to radiative recombination with  $n = 1$  having a non-negligible contribution.

### Multiple top cells

Furthermore, we calculate efficiency limits for multiple cells on top of a Si bottom cell, which can currently not be directly

**Table 4** *JV* parameters of 120  $\mu\text{m}$  thick Si bottom cell with additional extrinsic recombination  $J_{01,extr}$  (assuming ideality factor of 1), with a current-matched 1.72 eV bandgap SQ-limit top cell with  $\mu_{ic} = 87\%$ ; shows appropriate bottom cell ideality factors and tandem fill factor limits for devices with realistic bottom cells; ideality factor  $n_{Si}$  is determined by solving  $FF_0(V_{oc,Si}, n_{Si}) = FF_{Si}$ ; the respective tandem ideality factor for  $FF_0$  calculation is  $n_{Si} + 1$

$J_{01,extr}$ [ $\text{fA cm}^{-2}$ ]	$V_{oc,Si}$ [V]	$FF_{Si}$ [%]	$n_{Si}$ [ ]	PCE [%]	FF [%]
0	0.742	88.8	0.71	42.3	90.5
2	0.737	87.5	0.81	41.9	90.0
4	0.732	86.7	0.87	41.7	89.7
7	0.726	86.1	0.92	41.5	89.4
10	0.721	85.7	0.96	41.3	89.3
15	0.714	85.2	0.98	41.1	89.1
20	0.708	85.0	1	41.0	89.0
25	0.703	84.8	1.01	40.8	88.9
30	0.699	84.7	1.02	40.7	88.9
40	0.693	84.5	1.03	40.6	88.8

performed in Quokka3. As described by Schygulla *et al.*,<sup>6</sup> we first assume that the total absorbed photon flux is fixed regardless of the number of cells, which is  $44.37 \text{ mA cm}^{-2}$  in our tandem simulations. Note that this value is higher than in the single junction Auger limit of Niewelt *et al.*,<sup>2</sup> because of the higher optimum thickness (300  $\mu\text{m}$  vs. 100  $\mu\text{m}$ ). We then assume perfect photogeneration current-match, *i.e.*, that the current is evenly distributed between the sub cells. This defines the bandgaps of the top cells, and their  $J_{01}$  values. We then perform a single-junction Si *JV* curve simulation in Quokka3 for the bottom cell with the respective generation current density, and find the total *JV* curve by summing all sub cell voltages for each current density point calculated from the respective single diode model. This approach neglects bottom cell thickness optimization and LC, which however has been shown in the tandem case to be a minor effect at the efficiency optimum. With this data, we adapt Fig. 1 of Schygulla *et al.*<sup>6</sup> to include an Auger-limited Si bottom cell instead of a 1.12 eV SQ-limited bottom cell, providing more accurate and substantially lower theoretical limits for such devices, see results in Fig. 3b and Table 5.

**Table 5** *JV* parameters of multijunction device with Auger limited silicon bottom cell and multiple SQ limited sub cells with their bandgaps optimized for current-match and thus maximum efficiency

# of junctions	PCE [%]	$V_{oc}$ [V]	$J_{sc}$ [ $\text{mA cm}^{-2}$ ]	FF [%]
1 (Niewelt <i>et al.</i> )	29.4	0.757	43.4	89.5
2	43.2	2.16	22.2	90.3
3	48.4	3.61	14.8	90.7
4	51.3	5.09	11.1	90.9
5	52.9	6.55	8.87	91.0
6	54.0	8.02	7.39	91.0



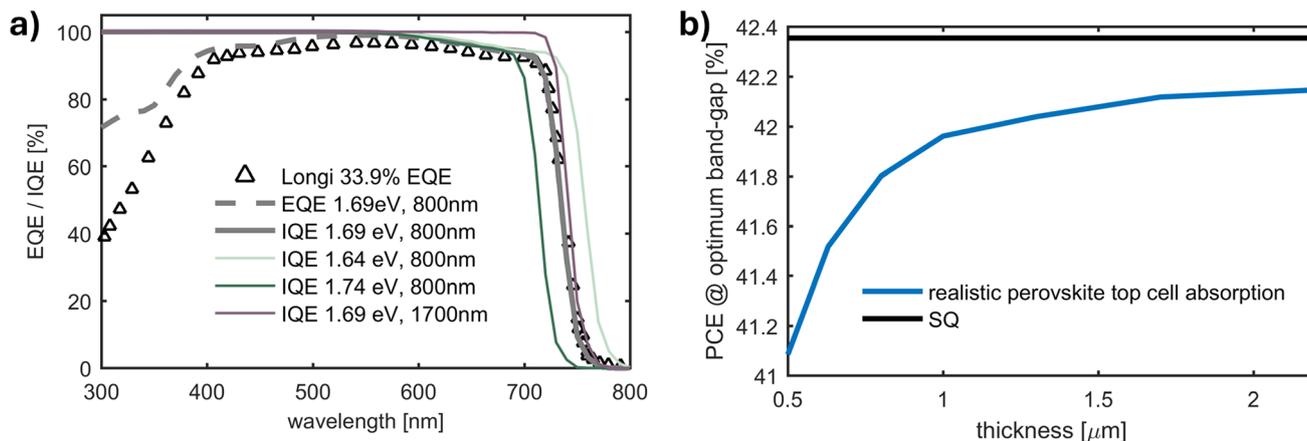


Fig. 4 (a) comparison of simulated EQEs and IQEs using ray-tracing for exemplary perovskite absorber bandgap and thickness (lines) with the measured top cell EQE of the 33.9% tandem cell by Longi; (b) efficiency at the optimal bandgap as a function of perovskite absorber thickness, compared with the SQ limit, in all cases assuming  $\eta_{\text{lc}} = 87\%$ .

### Perovskite top cell

Finally, we investigate the efficiency potential implied when considering a perovskite top cell with realistic absorption properties instead of the infinitely sharp absorption edge assumption within the SQ calculations. As described in the modeling section, our approach ensures that the only difference to the SQ limit is the intrinsic absorption capability of the perovskite absorber with the given bandgap and thickness, by using the IQE which is free from reflection and parasitic absorption losses. An overview of exemplary EQEs and IQEs is given in Fig. 4a, compared to the Longi 33.9% top cell EQE digitized from Liu *et al.*<sup>28</sup> It can be seen that using the nominal bandgap of 1.69 eV of the Longi cell, along with a realistic thickness value of 800 nm, the absorption around the bandgap, as well as the general EQE shape is well reproduced. This validates our optical modeling approach to be able to represent best-in-class experimental perovskite top cell absorbers.

In Fig. 4b the resulting efficiency at the optimum bandgap is plotted as a function of perovskite absorber thickness. Overall, the efficiency potential implied in the optical absorption properties of realistic perovskite material is close to the SQ limit. For practically achievable thicknesses around 1  $\mu\text{m}$  the fundamental efficiency penalty associated with imperfect absorption is below 0.5%, with only little headroom for larger thicknesses.

At 630 nm thickness, our simulations result in an efficiency of 41.5%. The difference to the practical limit of 39.5% published by Er-Raji *et al.*,<sup>26</sup> which used the same optical modeling approach, is explained mostly by reflection and parasitic absorption losses, and also by non-radiative recombination losses in the perovskite absorber.

## Conclusions

We have presented a rigorous and physically grounded model for evaluating the theoretical efficiency limits of tandem and multijunction solar cells incorporating a silicon bottom cell. By combining the state-of-the-art Auger-limit model for silicon

with an SQ-limit for top cells and accounting for luminescent coupling (LC), we provide an accurate picture of the performance potential of these architectures. Our results give a maximum theoretical efficiency of 43.2% for a silicon-based tandem device assuming an ideal intermediate dichroic mirror, *i.e.* free from LC—significantly lower than the 45.2% obtained when incorrectly applying the SQ limit to the bottom cell. For realistic LC conditions typical for perovskite top cells, the efficiency drops to 42.4% due to the increased radiative emission of the top cell. Assuming realistic absorption properties of a perovskite top cell with an absorber thickness of 1  $\mu\text{m}$  and full LC, the maximum efficiency is with 42.0% less than 0.5%<sub>abs</sub> below the theoretical limit, confirming the excellent tandem performance potential.

We identify an optimal silicon bottom cell thickness of 300  $\mu\text{m}$ , as opposed to the 100  $\mu\text{m}$  optimum for single-junction Auger-limited silicon, and an optimal top cell bandgap of 1.71 eV. Our analysis shows that performance is relatively insensitive to moderate variations in bottom cell thickness down to the range of typical Si cell thicknesses. The additional bottom cell current generation from LC needs to be explicitly included in the modeling only in bottom-cell-limited devices, and when the monofacial external radiative efficiency  $\eta_{\text{ext,mf}}$  of the top cell exceeds approximately 1%. Above  $\eta_{\text{ext,mf}}$  of approximately 10%, luminescent coupling significantly decreases  $V_{\text{oc}}$  due to the increased rear side radiative emission of the top cell.

In evaluating fill factor (FF), we emphasize the importance of using an appropriate ideality factor when applying the analytical FF limit ( $\text{FF}_0$ ). For an Auger-limited silicon bottom cell the appropriate ideality factor is  $n = 2/3$ , which gives a total ideality factor of  $n = 5/3$  to accurately predict the tandem  $\text{FF}_0$ , aligning with our simulations. Additionally, we emphasize that  $\text{FF}_0$  should be preferably compared to the pseudo or implied fill factor rather than the FF from  $J/V$  measurement, which is adversely impacted by current-mismatch, to avoid misleading loss interpretations.



For realistic assumptions of optical absorption for the case of a perovskite top cell, we show that the related efficiency penalty compared to the Shockley–Queisser limit is only  $\sim 0.5\%_{\text{abs}}$  for an absorber thickness of 1  $\mu\text{m}$ .

To support future modeling and device benchmarking efforts, we provide tabulated current–voltage ( $JV$ ) parameters for tandem configurations across a range of top cell properties, as well as for multijunction stacks comprising up to six subcells. These results serve as reliable limits for performance benchmarking, loss analysis, and strategic design of high-efficiency silicon-based multijunction solar cells.

## Author contributions

Andreas Fell: conceptualization, methodology, formal analysis, writing – original draft. Oliver Fischer: conceptualization, validation, writing – review and editing. Martin Bivour: methodology, project administration, writing – review and editing. Christoph Messmer: investigation, validation, writing – review and editing. Jonas Schön: methodology, validation, writing – review and editing. Martin Schubert: supervision, funding acquisition, writing – review and editing. Stefan Glunz: supervision, funding acquisition, writing – review and editing.

## Conflicts of interest

There are no conflicts to declare.

## Data availability

The most relevant result values of this article have been included as part of the main article as tables. For reproduction of the model, settings files for the solar cell simulation software Quokka3 are available on <https://www.quokka3.com/support/examples/efficiency-limit/>.

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