



Cite this: *Phys. Chem. Chem. Phys.*,  
2025, **27**, 20592

# Uncovering the photoexcited dynamics in bis(acyl)phosphine oxide photoinitiators

Marius Navickas,<sup>a</sup> Edvinas Skliutas,<sup>a</sup> Joseph Kölbel,<sup>b</sup>  
Ricardo J. Fernández-Terán,<sup>b</sup> Mangirdas Malinauskas<sup>a</sup> and Mikas Vengris<sup>a</sup>

Acylphosphine oxide photoinitiators, such as phenylbis(2,4,6-trimethylbenzoyl)phosphine oxide (**BAPO**) and diphenyl(2,4,6-trimethylbenzoyl)phosphine oxide (**TPO**), are widely used as sensitizers in UV polymerisation applications due to their favourable properties for UV lithography. However, the underlying photophysical and photochemical mechanisms that govern their performance remain poorly understood. In this study, we systematically investigate these compounds using steady-state and time-resolved fluorescence, as well as electronic and infrared transient absorption spectroscopic techniques. The results show that the lowest electronic excitation of both **BAPO** and **TPO** leads to a rapid intersystem crossing (ISC) into the hot triplet manifold—which rapidly generates radicals, further undergoing relaxation and leading to a slower rate of radical formation from the cold triplet manifold. ISC is faster in **TPO** compared to **BAPO**, which exhibits more complex dynamics. In addition, time-resolved infrared data reveal that both sensitizers generate radicals (by  $\alpha$ -cleavage) on a time scale similar to that of ISC, suggesting that the hot triplet state is very short-lived. Finally, nanosecond-to-microsecond electronic TA spectroscopy revealed multi-stage relaxation of the radicals with the formation of **TPO** photoproducts remaining detectable on microsecond time scales.

Received 28th April 2025,  
Accepted 31st August 2025

DOI: 10.1039/d5cp01612f

rsc.li/pccp

## 1. Introduction

Photopolymerisation has attracted great scientific and technological interest due to a large number of applications, ranging from microelectronics to medical areas.<sup>1–4</sup> This process enables the manufacturing of complex three-dimensional objects which cannot be cut, assembled, or carved, also at length scales that go beyond conventional machining capabilities.<sup>5,6</sup>

To maintain processing speed, reduce irradiation power and improve the control over the polymerisation process, various photoinitiators can be used.<sup>7,8</sup> This is especially important in applications where multi-focus or projection-based curing is employed for higher processing throughput. Photoinitiators are essential compounds of photoresists, which are of great interest in a variety of applications such as UV-curable inks, coatings and adhesives, lithography, table-top stereolithography 3D printing, and multiphoton 3D nanoprinting.<sup>9</sup> These organic molecules absorb light and produce free radicals that trigger the polymerisation reactions.<sup>10</sup> Since initiation is the first step in the photopolymerisation sequence, it is an important process determining the chemical and physical nature of the final product.

Organic carbonyl-based photoinitiators are often divided into two main categories, based on the mechanism of radical formation: those that undergo a photoinduced  $\alpha$ -cleavage (type I) and those that involve hydrogen abstraction by photoexcited molecules (type II).<sup>11</sup> Type I photoinitiators include acetophenone derivatives, benzoin ethers, aminoketones<sup>12,13</sup> and acylphosphine oxides.<sup>12,14</sup> In this work, we will focus on two representatives of the latter group (Fig. 1): phenylbis(2,4,6-trimethylbenzoyl)phosphine oxide (**BAPO**) and diphenyl(2,4,6-trimethylbenzoyl)phosphine oxide (**TPO**), which have received considerable recent attention due to presence of different reactive intermediates, and their absorption spectra, which makes them suitable for their use with UV-A and UV/Vis sources.<sup>15</sup> Photoexcitation of these compounds leads to  $\alpha$ -cleavage, generating mesitoyl (**R1**) and acylphosphinoyl (**R2** and **R3**) radicals.<sup>16,17</sup>

Following the high interest in their applications, Meereis *et al.* have demonstrated that **BAPO** is remarkably suitable for promoting rapid and efficient polymerisation.<sup>17</sup> **BAPO** has also become an attractive commercially-available photoinitiator for table-top stereolithography 3D printing,<sup>18</sup> especially for those employing low-cost LEDs or laser diodes emitting at 405 nm.<sup>19–21</sup> Multiphoton 3D nanoprinting is one more field, where **BAPO** is used to modify the photosensitivity of the resins.<sup>22–24</sup>

The initial radicals of **BAPO** have been extensively studied by different spectroscopic techniques to resolve their structure and decay lifetimes.<sup>16,25</sup> It was demonstrated that acylphosphinoyl

<sup>a</sup> Laser Research Centre, Vilnius University, VilniusSaulėtekio ave. 10, LT-10223, Lithuania. E-mail: marius.navickas@ff.vu.lt

<sup>b</sup> Department of Physical Chemistry, University of Geneva, CH-1205 Geneva, Switzerland. E-mail: Ricardo.FernandezTeran@unige.ch



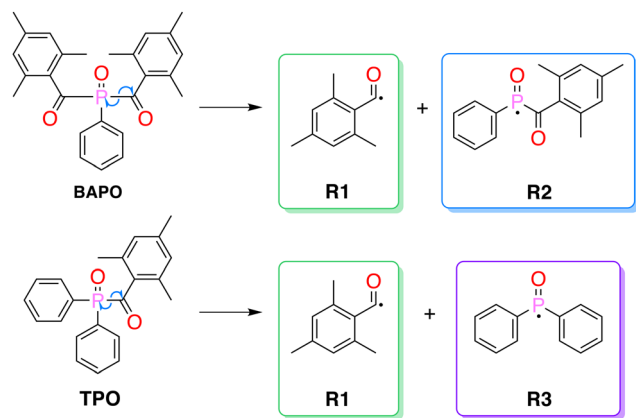


Fig. 1 Chemical structures of **BAPO** and **TPO**, and their  $\alpha$ -cleavage radical photoproducts.

radicals generated by photolysis of acyl and bis(acyl)phosphine oxides show high activity of addition to acrylate prepolymers,<sup>26–29</sup> illustrating their effectiveness as photoinitiators for free radical polymerisation.

**TPO** is another compound widely used for photopolymerisation initiation, and is compatible with epoxy acrylate systems<sup>30</sup> and natural polymers such as lignin<sup>31</sup> and vinillin.<sup>32</sup> This photoinitiator belongs to the same bis(acyl)phosphine oxides class, and has a narrow UV/Vis absorption window of 380–425 nm with the maximum peaking at *ca.* 400 nm.<sup>33</sup> Additionally, **TPO** exhibits a broader absorption range in the deeper UV region. Exposure of **TPO** to UV light promotes an  $\alpha$ -cleavage photoreaction leading to the formation of **R1** and diphenylacylphosphinoyl ( $\text{Ph}_2\text{P}^\bullet = \text{O}$ ) radicals.<sup>34</sup> An important advantage of **TPO** and **BAPO** is their ability to initiate the monomer polymerisation reaction without the assistance of co-initiators.<sup>15</sup> This can simplify the manufacturing process, and lead to a faster photopolymerisation.<sup>15</sup> The transient intermediates of **TPO** have been successfully detected by using both time-resolved and steady-state infrared spectroscopic techniques, which have been used to determine their structures and reactivity.<sup>13,35</sup>

Many studies of reactive intermediates in these photoinitiators focus on the measurements of the reactivity and determination of the polymerisation rate constants, whilst the photophysical mechanisms of initiation have received much less attention. For instance, the primary dynamics of molecule cleavage into radical fragments and the subsequent changes of the absorption spectrum during radical relaxation have been barely investigated. Moreover, the vast majority of radical dynamics studies have relied on laser flash photolysis, with ns– $\mu\text{s}$  time resolution, whilst the ultrafast events of radical formation remain poorly understood. To gain insights into the full dynamic picture covering the fs to  $\mu\text{s}$  timescales in these commonly used photoinitiators (**BAPO** and **TPO**), we employed steady-state and time-resolved fluorescence, together with ultrafast and fs to  $\mu\text{s}$  pump–probe spectroscopies in both UV/Vis and mid-IR spectral regions, supported by density functional theory (DFT) calculations.

## 2. Materials and methods

**BAPO** and **TPO** were purchased from Sigma-Aldrich, and used as received. For transient UV/Vis studies, the compounds were dissolved in 2-propanol (Sigma-Aldrich, HPLC Grade), and diluted to an optical density of *ca.* 0.2–0.3 at the lowest absorption band in a 1 mm optical pathlength quartz cuvette (Hellma) as presented in Fig. 2. For time-resolved infrared (TRIR) studies, the compounds were dissolved in dichloromethane (Sigma-Aldrich, HPLC grade). Also, the comparison UV/Vis TA experiments in dichloromethane were performed, and the obtained datasets for both **BAPO** and **TPO** are presented in the SI (Fig. S1 and S2).

Steady-state absorption and emission/excitation spectra of the 2-propanol solutions were recorded using a Shimadzu UV-3101PC UV/Vis absorption spectrometer, and a PerkinElmer LS55 fluorescence spectrometer, respectively. For steady-state fluorescence measurements, the sample solutions were diluted to an optical density below 0.1 in a  $1 \times 1$  cm Hellma quartz cells.

Time-resolved emission spectra were collected at different wavelengths using a time-correlated single-photon counting system (TCSPC), integrated in the Harpia-TF time-resolved fluorescence spectrometer (Light Conversion Ltd). In these experiments, the third harmonic (344 nm) of an Yb:KGW laser (Pharos, Light conversion, Ltd) operating at 1030 nm, and delivering 200 fs pulses at 100 kHz repetition rate. The fluorescence excitation power was set to *ca.* 65  $\mu\text{W}$  at 100 kHz pulse repetition rate.

Ultrafast transient UV/Vis absorption (TA) spectra were measured using a Harpia-TA (Light Conversion, Ltd) commercial transient absorption spectrometer. The pump pulse was generated by an optical parametric amplifier (TOPAS-800, Light Conversion) which was pumped by a commercial Ti:Sa femtosecond laser (Coherent Libra; delivering 800 nm, 50 fs, 2.5 W pulses at 1 kHz). For ultrafast pump–probe experiments, the **BAPO** and **TPO** samples were excited at the lowest absorption band. The excitation power for all performed experiments was tuned to vary by an order of 2% of the excited states, to avoid any multiphoton ionization effects. Detailed values of absorbed photons per pulse as well as concentrations of the excited states

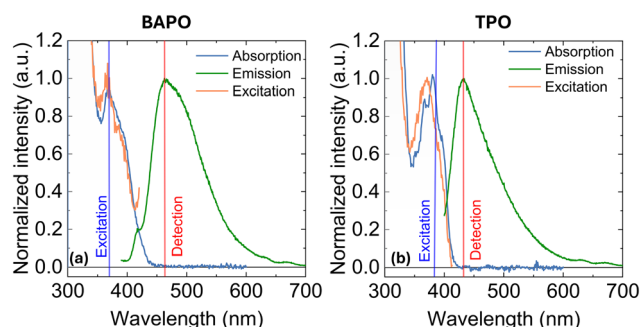


Fig. 2 Steady-state absorption, emission and excitation spectra of (a) **BAPO** and (b) **TPO** compounds. Vertical lines show the wavelengths used to excite or detect the corresponding emission/excitation spectra.



can be found in Table S1 of SI. To probe the population dynamics, we used a white light supercontinuum, generated in a mechanically-translated CaF<sub>2</sub> plate. The polarisation of the pump beam was set to magic angle (54.7°) respect to the probe beam. The pump–probe delay was controlled with an optical delay line, in the range from fs to a few ns. The TA spectra were recorded in a 1 mm pathlength Hellma quartz flow cell to ensure that each laser pulse irradiated a fresh volume of sample. At least, 10 scans with 10 bathes of 250 shots for each TA spectra were recorded and averaged over the number of performed scans that were virtually identical.

To extend the probe time window beyond the range of the optical delay line, an additional pump laser (NL640, Ekspla), which emitted 3 ns pulses at 355 nm, was used as an alternative excitation source. The samples were excited with 1 μJ pulses and 120 μm spot size beam. The operation of this laser was electronically synchronised with the Ti:Sapphire laser using the available digital outputs of the SDG-Elite synchronisation and timing unit (part of the Libra laser system). Using this arrangement, the ns laser pulse could be timed to arrive between 800 μs before to 50 ns after the femtosecond probe pulse, with a precision of 250 ps.

Time-resolved infrared spectra (TRIR) were collected at the University of Geneva. In brief, the output of a Ti:Sa amplifier (Spectra-Physics, Solstice ACE) producing *ca.* 100 fs pulses centred at 800 nm was used to pump two TOPAS-C optical parametric amplifiers (Light Conversion, Ltd). The output of the first TOPAS-C was fed into a non-collinear difference frequency generation (NDFG) stage with a AgGaS<sub>2</sub> crystal, producing mid-IR pulses centred at 1600 or 1800 cm<sup>-1</sup>, with a FWHM of *ca.* 170 cm<sup>-1</sup>. The second TOPAS-C was used to produce 640 nm pulses, which were then frequency doubled in a bismuth borate (BiBO) crystal, producing UV excitation pulses at 320 nm. The 400 nm pulses used in complementary TRIR experiments were obtained by frequency-doubling a portion of the fundamental 800 nm light in a β-barium oxide (BBO) crystal. The fundamental light was removed in both cases by the use of a Glan–Taylor polariser, after which an achromatic λ/2 waveplate was used to rotate the polarisation to magic angle (54.7°) relative to the s-polarised probe beam. The pump–probe delay was set with the help of a mechanical delay stage. The mid-IR beam was split into probe and reference beams using an uncoated CaF<sub>2</sub> wedge, and both beams were passed through two wire grid polarisers—to clean the polarisation and attenuate, respectively—both the probe and reference beams. The two mid-IR beams were focussed at the sample to a spot size of *ca.* 120 μm using off-axis parabolic mirrors, and the pump was focussed with the help of a lens (*f* = 75 cm), to a spot size of *ca.* 240 μm. All other experimental parameters for each spectroscopic technique are summarised in a Table S2 in the SI.

After the sample, both probe and reference were imaged into a spectrograph and dispersed with a 150 l mm<sup>-1</sup> grating into a 64 × 2 pixel liquid-nitrogen cooled HgCdTe detector. A multi-channel referencing scheme, as described by Ge and co-workers was used to account for fluctuations in the probe beam.<sup>36,37</sup> The resulting TRIR spectra were obtained after averaging at

least 4 scans, with 4 batches of 500 shots at each time delay. The samples were recirculated in home-made cells with 2 mm thick CaF<sub>2</sub> windows and a 400 μm PTFE spacers, and the scans were checked for scan-to-scan stability of the observed signals. Note that some spectral artifacts are observed due to the TRIR measurements being made in the absence of suitable purging, and contain the imprinted signatures of water vapour and CO<sub>2</sub>, evident as sharp features (especially in the 1550–1600 cm<sup>-1</sup> region).

Density functional theory (DFT) calculations were performed in Gaussian 16, rev. A.03,<sup>38</sup> at the B3LYP/6-311++g(d,p) level of theory,<sup>39–41</sup> including the effects of the DCM solvent using the IEFPCM model.<sup>42</sup> All structures were optimised under tight convergence criteria, and were found to be true minima in their corresponding potential energy surfaces by the absence of negative (imaginary) harmonic frequencies. Difference mid-IR spectra were obtained by convolution with a Lorentzian line shape function of 8 or 10 cm<sup>-1</sup> FWHM for the ground-state and excited-state/photoproduct spectra, respectively. The harmonic frequencies were scaled by 0.983 to better match the experimentally obtained values.

Global/target analysis was performed on both time-resolved fluorescence and pump–probe (UV/Vis and mid-IR) data, to extract meaningful time constants and species- or evolution-associated spectra,<sup>43–45</sup> using the CarpetView software (Light Conversion, Ltd).

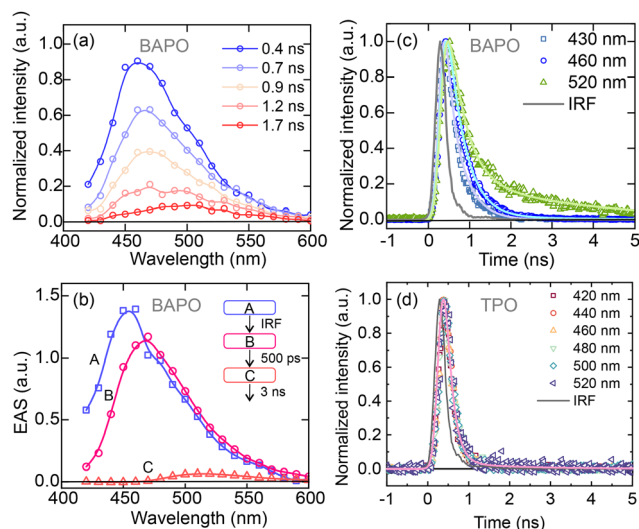
## 3. Results

### 3.1. Steady-state and time-resolved emission spectroscopy

We begin by analysing the steady-state absorption and emission spectra of **BAPO** and **TPO** in 2-propanol solutions (Fig. 2). The absorption spectra (solid blue lines) of both **BAPO** and **TPO** extends to around 420 nm, with absorption maxima at *ca.* 360 nm and 380 nm, for **BAPO** and **TPO**, respectively. Their emission spectra (solid green lines) consist of a broad, structureless band spanning the range from 420–600 nm, with emission maxima at *ca.* 470 nm and 430 nm for **BAPO** and **TPO**, respectively. Their excitation spectra (solid orange lines) closely match the corresponding absorption profiles, indeed showing that the observed emission is a result of direct excitation.

Time-resolved emission spectra (Fig. 3), obtained after pulsed excitation of **BAPO** at 344 nm, show a continuous red shift from 470 nm to 520 nm (*cf.* 400 ps *vs.* 1.7 ns spectra in Fig. 3a), accompanied with a decay that is almost complete within 1 ns. The evolution-associated spectra (EAS) obtained from global analysis reveal an IRF-limited component (the FWHM of the IRF is 300 ps), and two time constants of *ca.* 500 ps and 3 ns—the latter resulting in a significant decrease in intensity. Inspection of the emission decay traces at different detection wavelengths (Fig. 3c), reveals slightly longer decay times in the red emission wavelengths, suggesting non-trivial relaxation dynamics in the excited state. In contrast, the corresponding time-resolved emission data of **TPO** (kinetic





**Fig. 3** (a) Time-resolved emission spectra of **BAPO** in 2-propanol. (b) Evolution-associated spectra retrieved from global analysis. Panels (c) and (d) represent, respectively, the fluorescence decay traces of **BAPO** and **TPO** in 2-propanol at different detection wavelengths.

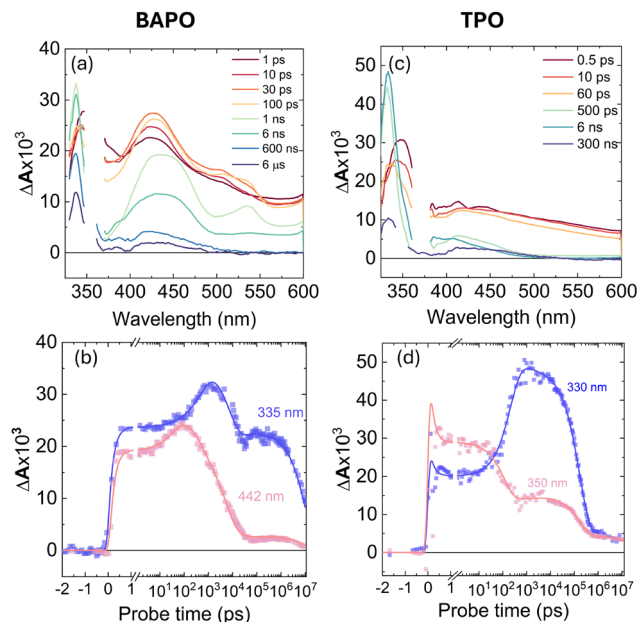
traces shown in Fig. 3d) show no emission wavelength dependence, and all kinetic traces decay with an IRF-limited time constant.

### 3.2. Transient electronic spectroscopy

To gain further insight into the excited-state photodynamics of **BAPO** and **TPO** upon UV excitation, we now turn to transient UV/Vis absorption spectroscopy (TA). In these experiments, we combined broadband fs–ns and ns– $\mu$ s pump–probe techniques to probe the entire range of timescales needed to fully address the complex photochemistry of these compounds. The resulting spectra and selected kinetic traces are shown in Fig. 4.

As shown in Fig. 4a, within several picoseconds after excitation of **BAPO**, we observe a broad and featureless excited-state absorption (ESA) background, with distinct spectral maxima around 335 nm and 442 nm which further develop over time. The peaks around 335 nm and 442 nm become more pronounced, and after approximately 30 ps, an additional peak starts to form at *ca.* 530 nm. The spectral feature around 335 nm does not change significantly after *ca.* 1 ns.

An interesting spectral transformation also takes place on a ns– $\mu$ s timescale. The decrease of ESA at 335 nm is slower compared to the ESA at 442 nm, and the IA feature at 530 nm becomes negligible after *ca.* 6 ns. Kinetic traces at 335 nm and 442 nm (Fig. 4b) reveal the temporal evolution of the **BAPO** transients. The 335 nm signal decays rapidly within picoseconds, whilst the 442 nm signal remains stable. Both traces then grow, with the 442 nm signal peaking at 200 ps and the 335 nm signal continuing to rise until *ca.* 1 ns. Subsequently, both signals decay, with the 442 nm trace showing a more gradual decline compared to the steeper slope of the 335 nm trace. Both traces plateau at *ca.* 20 ns, followed by a slow decay over *ca.* 10  $\mu$ s. A residual TA signal persists, attributed to long-lived radical species responsible for photoinitiator activity.



**Fig. 4** TA spectra (top row) and selected kinetic traces (bottom row) obtained after excitation with 360 nm pulses (fs–ns) or 355 nm (ns– $\mu$ s), of solutions of **BAPO** (a and b) and **TPO** (c and d), respectively. The spectra were merged at a delay of 6 ns.

For comparison, the photoinduced dynamics of **TPO** in 2-propanol were also studied. Excitation at 360 nm produces a broad IA band peaking at *ca.* 330 nm, and a weaker peak around *ca.* 425 nm (Fig. 4c). Kinetic traces at 330 nm and 350 nm (Fig. 4d) reveal more intricate sub-picosecond dynamics: the 335 nm signal initially decreases slightly, remains stable until *ca.* 30 ps, then grows to a maximum at *ca.* 800 ps. Meanwhile, the 350 nm signal remains stable until *ca.* 30 ps before decaying to a plateau, coinciding with the 330 nm peak. Both traces exhibit multi-stage relaxation over 1–100 ns, with weaker residual signals compared to **BAPO**, and notably, the lack of an ESA feature around 530 nm in the **TPO** spectra.

### 3.3. Quantum chemical calculations

Density functional theory (DFT) calculations reveal interesting insights into the photochemical reaction and the properties of the obtained photoproducts. Optimisation of both **BAPO** and **TPO** in the lowest singlet excited state ( $S_1$ ) yielded stable structures, whilst the optimisation of **BAPO** in the lowest triplet excited state ( $T_1$ ) resulted in cleavage of the C–P bond, which was not the case for **TPO**. The expected mesityl (**R1**) and acylphosphinoyl radicals (**R2** and **R3**) were also optimised at the same level of theory in their doublet spin configuration, and the obtained structures and spin density isosurfaces are shown in Fig. 5. The XYZ coordinates of all relevant structures are provided in the SI.

As is evidenced from the plots in Fig. 5, the spin density is largely localised in the C=O or P=O moieties for **R1** and **R2/R3**, respectively.

We used the results of these calculations, and the corresponding harmonic vibrational spectra, to guide our interpretation of the experimental time-resolved infrared spectra, discussed next.



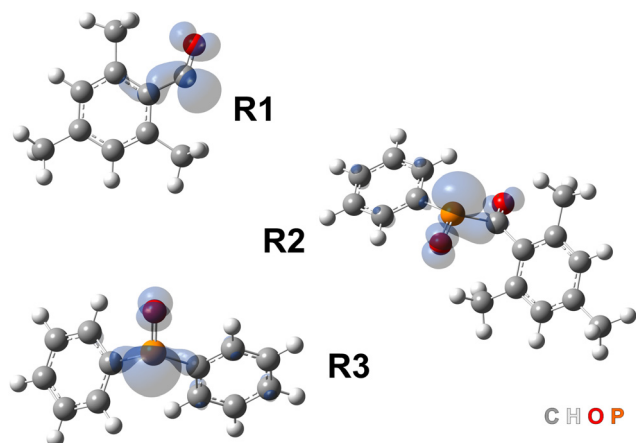


Fig. 5 Optimised structures and spin density isosurfaces of the mesityl radical (**R1**) and acylphosphinoyl (**R2** and **R3**) radicals relevant for this study. Isosurfaces plotted at 0.005 a.u. The element colour code is included in the figure (bottom right corner).

### 3.4. Time-resolved infrared spectroscopy

Transient infrared spectra (TRIR) of **BAPO** and **TPO** in the C=C and C=O stretching region ( $1550\text{--}1660\text{ cm}^{-1}$  and  $1740\text{--}1875\text{ cm}^{-1}$ ) were recorded upon excitation at both 400 nm and 320 nm. We used dichloromethane (DCM) as solvent for these experiments due to its higher IR transparency relative to 2-propanol.

Herein, we will focus on the results obtained with 400 nm excitation. The complementary datasets associated with 320 nm excitation are shown in the SI. The TRIR data for both compounds is presented in Fig. 6. The earliest time spectrum (500 fs) shows, in both cases, a negative signal at  $1609\text{ cm}^{-1}$ , corresponding to the ground-state bleach (GSB) of the C=C stretching modes localised on the mesityl moiety.

In both cases, we also observe a prompt ESA band around  $1585\text{ cm}^{-1}$ , corresponding to the C=C ring stretching modes localised on the  $\text{Ar}_2\text{P}=\text{O}$  moieties of excited **BAPO** and **TPO** in the  $S_1$  state. These signals do not show any significant spectral evolution, decaying in *ca.* 100 ps, and giving rise to a weak transient at *ca.*  $1605\text{ cm}^{-1}$  (attributed to the C=C ring stretching mode of the mesityl radical, **R1**) and a strong transient at *ca.*  $1798\text{ cm}^{-1}$  (attributed to the C=O stretching mode of **R1**). The calculated mid-IR difference spectra (Fig. 6, panels a and c) reproduce almost quantitatively the experimental data, supporting both the level of theory and the reactivity scheme used to interpret our results.

A closer inspection of kinetic traces at selected positions (Fig. 7), corresponding respectively to the bleach of the parent compound ( $1609\text{ cm}^{-1}$ , blue), the transient associated with the  $S_1$  state ( $1585\text{ cm}^{-1}$ , green), and the transient associated with the C=O stretching of the mesityl radical photoproduct ( $1798\text{ cm}^{-1}$ , red), the latter in excellent agreement with previous TRIR reports by George and co-workers.<sup>13</sup>

The  $1609\text{ cm}^{-1}$  trace exhibits a bi-exponential recovery with a fast component of 500 ps, but does not fully recover. This fast component well-matches to the emission lifetime (the detailed fitting is presented in Fig. S5). The agreement between the **BAPO** emission lifetime and the GSB recovery time constant indicates that the observed emission arises from  $S_1 \rightarrow S_0$  relaxation. Concurrently, the  $1585\text{ cm}^{-1}$  trace shows a sub-nanosecond decay, followed by the gradual rise of a new feature at  $1798\text{ cm}^{-1}$ . The decay of the  $1585\text{ cm}^{-1}$  band and the rise of the  $1798\text{ cm}^{-1}$  band occur on the same timescale. The lack of GSB recovery until the new ESA appears suggests that most of the **BAPO** excited-state population undergo decomposition, without repopulating  $S_0$ .

We obtained similar results for **TPO** under these conditions ( $\lambda_{\text{exc}} = 400\text{ nm}$ , DCM), as expected from their similarity.

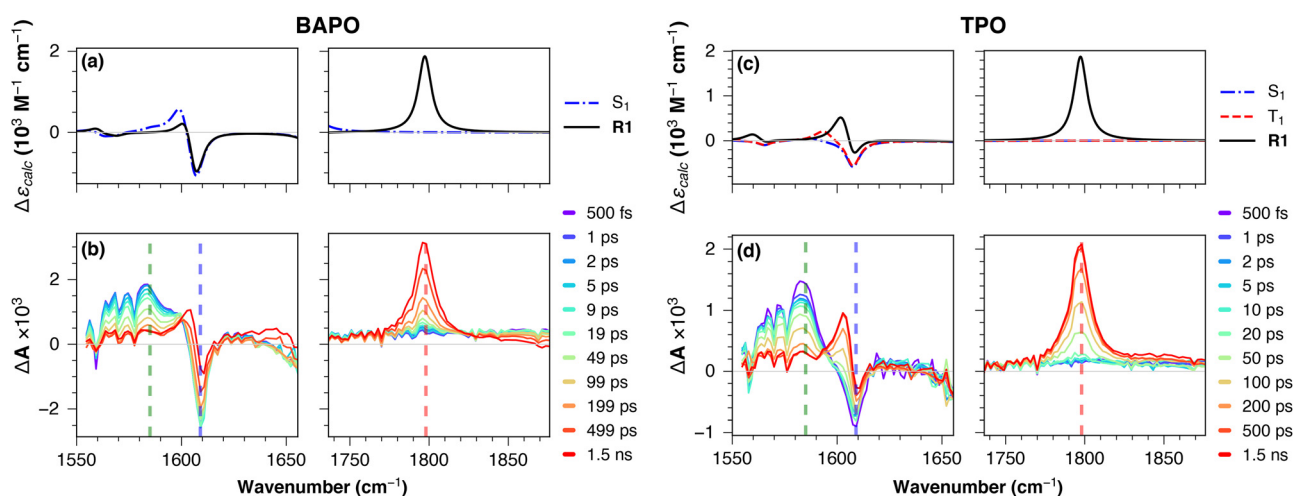


Fig. 6 Calculated mid-IR difference spectra (top row) and experimental TRIR spectra following 400 nm excitation in DCM (bottom row), of **BAPO** (a) and (b) and **TPO** (c) and (d). Vertical dashed lines indicate the probe wavelengths used for the kinetic cuts shown in Fig. 7, corresponding to the bleach of the parent compound ( $1609\text{ cm}^{-1}$ , blue), the transient associated with the  $S_1$  state ( $1585\text{ cm}^{-1}$ , green), and the transient of the photoproduct acyl radical band ( $1798\text{ cm}^{-1}$ , red). The calculated frequencies were scaled by 0.983 to better match the experiment. Optimisation of **BAPO** in the  $T_1$  state led to breaking of the C–P bond, suggesting a dissociative behaviour in this state (see text for discussion).



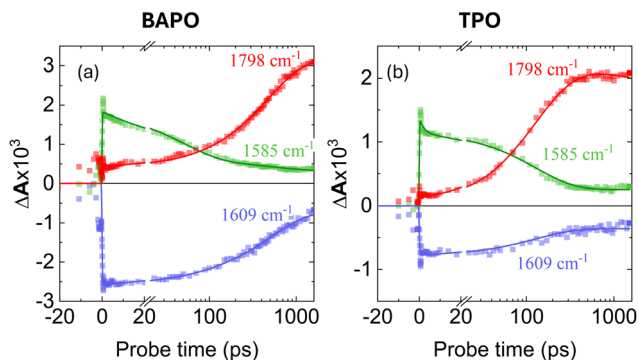


Fig. 7 Kinetic traces corresponding to the bleach of the parent compound ( $1609\text{ cm}^{-1}$ , blue), the transient associated with the  $S_1$  state ( $1585\text{ cm}^{-1}$ , green), and the transient of the photoproduct mesityl radical band ( $1798\text{ cm}^{-1}$ , red), obtained from the TRIR data of **BAPO** (a) and **TPO** (b) upon photoexcitation at  $400\text{ nm}$ . Symbols indicate measured data, and solid lines indicate the results from the global fits.

Although its TRIR dynamics have been previously discussed,<sup>46</sup> they exhibit spectral features analogous to **BAPO**, making **TPO** a suitable reference. The kinetic traces of **TPO** mirror those of **BAPO**: the ESA at  $1585\text{ cm}^{-1}$  decays, whilst the ESA at  $1798\text{ cm}^{-1}$  grows on the same timescale. Notably, the  $1609\text{ cm}^{-1}$  GSB kinetics show a bi-exponential recovery with time constants  $\tau_1 = 130(10)\text{ ps}$  (for **BAPO**, it is  $500(50)\text{ ps}$ ) and an infinite component ( $\tau_2$ ), indicating dynamics taking place beyond the *ca.*  $1.6\text{ ns}$  timescale accessible with our optical delay line, as expected from the photochemistry of these compounds and from ns- $\mu\text{s}$  TA experiments.

The short recovery time of the GSB and the short emission lifetime in **BAPO** and **TPO** indicate a quick decay of the singlet excited state. This short lifetime of the excited state may result from direct cleavage or intersystem crossing, which leads to the

population of triplet states. Altogether, the TRIR results confirm the identity of the photoproducts (**R1**, **R2** and **R3**), as well as the lack of direct  $\alpha$ -cleavage from the singlet excited state. This is in line with previous reports from Turro and co-workers,<sup>47</sup> who state that  $\alpha$ -cleavage in cyclic ketones is expected to be *ca.* 100 times faster from the triplet excited state than from the singlet excited state.

### 3.5. Naphthalene quenching experiments

To test the involvement of the triplet state, and to confirm its lifetime, we performed ultrafast UV/Vis pump-probe experiments ( $\lambda_{\text{exc}} = 350\text{ nm}$ ) aiming for triplet sensitisation, in 2-propanol solutions of **BAPO** and **TPO**, with naphthalene as triplet sensitizer. Since triplet energy transfer requires orbital overlap (*i.e.* direct contact of the molecules), we use quite high naphthalene concentrations of  $300\text{ mM}$ . The main results from these experiments are summarised in Fig. 8.

In Fig. 8a and b we present the results of triplet sensitisation of **BAPO** and **TPO**, respectively. As observed, for **BAPO** the TA spectra on the picosecond time scale are not influenced by naphthalene. Only at *ca.*  $450\text{ ps}$  some naphthalene triplet signatures (fine peaks at  $390\text{ nm}$  and  $415\text{ nm}$ )<sup>48</sup> start to appear and become more visible after *ca.*  $3\text{ ns}$ . For **TPO**, the triplet energy transfer seems to be faster, and naphthalene triplet features are better resolved.

To confirm triplet-triplet energy transfer in naphthalene solutions, we measured TA dynamics of benzophenone with  $500\text{ mM}$  naphthalene (Fig. 8e), since benzophenone has a very high triplet quantum yield. Then, if the appearance of peaks at  $390\text{ nm}$  and  $415\text{ nm}$  arises from naphthalene triplet state absorption, it should be strongly prominent in the benzophenone-naphthalene solution. Indeed, the same characteristic naphthalene triplet state absorption appears and increases over several

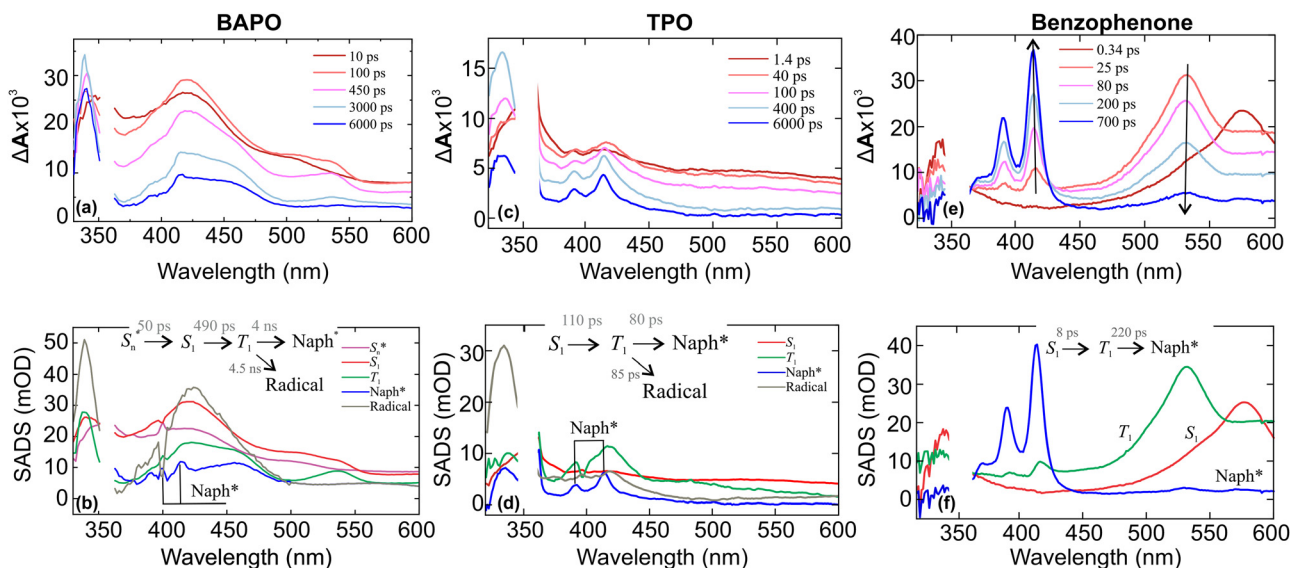


Fig. 8 (a), (c) and (e) TA spectra of **BAPO**, **TPO** and benzophenone 2-propanol solutions, containing naphthalene and (b), (d) and (f) SADS, obtained from global analysis. The empty region in spectra indicate area omitted due to the scattering of the pump pulse. The concentrations of **BAPO**, **TPO** and naphthalene were  $3\text{ mM}$ ,  $6\text{ mM}$  and  $300\text{ mM}$ , respectively.



hundred picoseconds; meanwhile, the absorption of benzophenone singlet state at 530 nm decreases. That indicates a fast triplet-triplet energy transfer between benzophenone and naphthalene.

## 4. Discussion

The current understanding of **BAPO** and **TPO** photochemistry states that triplets are usually precursors of radicals. This hypothesis is widely supported in the field of ESR spectroscopy,<sup>16,25</sup> however, the actual time constants for intersystem crossing (ISC) and subsequent  $\alpha$ -cleavage have not been clearly resolved. To determine the time scales of triplet state population in **BAPO** and **TPO**, we start with naphthalene quenching experiments. To analyze the femtosecond-to-nanosecond UV/Vis TA dataset of **BAPO** and **TPO** solutions containing naphthalene as a triplet sensitizer, we apply connectivity schemes to fit the data, which are also illustrated in the insets of Fig. 8. In the analysis of the **BAPO** TA dataset, the pathway  $S_1 \rightarrow T_1$  corresponds to an ISC occurring within 490 ps, while the transition  $T_1 \rightarrow$  Naph evolves over 4.0 ns, representing the triplet-triplet energy transfer between **BAPO** and naphthalene. Furthermore, the transition  $S_1^* \rightarrow S_1$  denotes the relaxation of the excited singlet state. The species-associated difference spectrum (SADS) of Naph reveals characteristic ESA features at 390 nm and 415 nm, which confirm the presence of triplet states in **BAPO** photodynamics. In case of **TPO**, these triplet features appear even faster, within 80 ps, and are better resolved. Although the triplets in **BAPO** and **TPO** that are surrounded by naphthalene undergo quenching, a fraction of them still undergo  $\alpha$ -cleavage, resulting in the formation of free radicals. This process is illustrated by a pathway  $T_1 \rightarrow$  radicals, with a corresponding

time constant of 4.5 ns. Despite the findings, there is a puzzling aspect regarding the timescale of triplet energy transfer for **BAPO** and **TPO**. In case of **BAPO**, the energy transfer to naphthalene takes 4.5 ns, whereas for **TPO**, it only takes 80 ps. Since the triplet sensitizer in both cases is the same and **BAPO** together with **TPO** have similar structures, we would expect the energy transfer time constants to be comparable.

The establishment of the role of the triplets in both molecules implies that fast emission decay might be related to ISC; in other words, the singlet state in both molecules is deactivated due to the population of triplets. Since the line shape of the absorption and excitation spectra are similar, we conclude that the emissive state is produced with the same probability from the absorbing state. At the same time, the fact that the Stokes shift is in the order of thousands of  $\text{cm}^{-1}$ , and the spectrum is broad, points to complex relaxation taking place in the excited state. Of particular interest here is the emission decay of **BAPO**. Here, after 500 ps, we obtain a small residual signal, spreading around 520 nm and surviving nanoseconds. Given the involvement of the triplet state from naphthalene experiments, we hypothesized that this residual signal could be phosphorescence or delayed fluorescence. This hypothesis became an argument when we measured the emission decay kinetics in oxygen-free **BAPO** solution. We found a slight prolongation of emission lifetime (see SI, Fig. S6). This effect was reversed upon bubbling the solution once more with air, confirming that the obtained changes are related to the content of oxygen in the solution. This is indeed strong evidence to attribute the red emission decaying in nanoseconds to phosphorescence of the triplets.

Assuming that triplets are precursors of radicals, they should emerge over nanoseconds for **BAPO**, but TRIR shows different. The transition B  $\rightarrow$  C (see Fig. 9a) with a lifetime

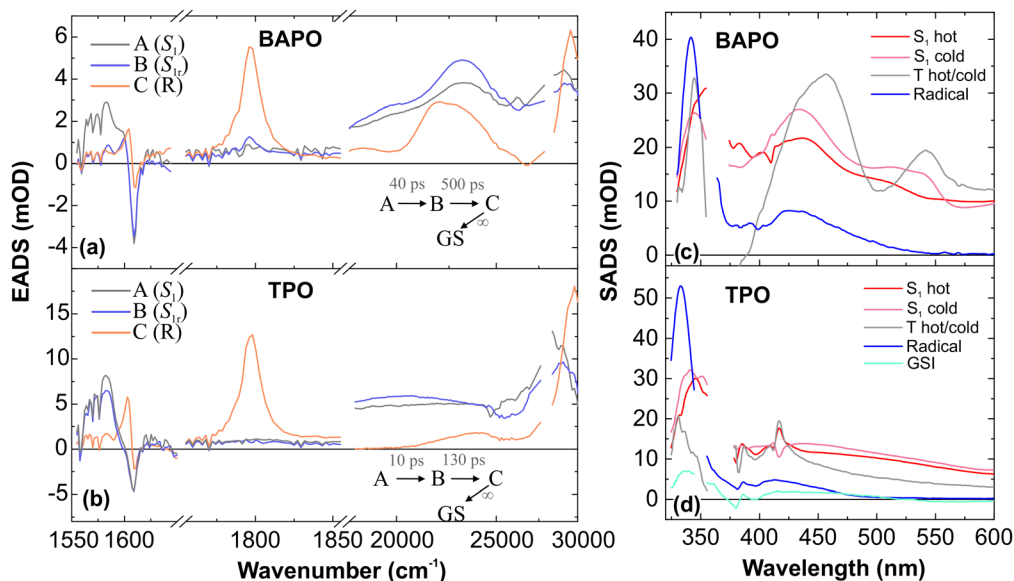


Fig. 9 Evolution-associated difference absorption spectra (EADS) and time constants obtained from a global analysis of the (a) **BAPO** and (b) **TPO** TRIR-UV/Vis data. The panels (c) and (d) show the SADS of femtosecond-microsecond UV/Vis **BAPO** and **TPO** TA data obtained by assuming the model in Fig. 10. The TRIR spectra in panel (a) and (b) were scaled to match their intensities with UV/Vis TA spectra.



close to that of ISC shows an appearance of characteristic frequencies of C=O stretching in Mes-CO radical.<sup>46</sup> Furthermore, the transition in the IR region is also accompanied by a change in UV/Vis, but the characteristic EADS C of UV/Vis TA does not resemble the spectrum of the radical; it may rather be attributed to the triplet features. This contradicts either naphthalene quenching and emission results that suggest ten times slower development of the radicals than obtained by TRIR. However, these experimental results seem to contradict each other only under the assumption that the photodynamics in **BAPO** and **TPO** are sequential, *i.e.*, singlet  $\rightarrow$  triplet  $\rightarrow$  radicals. The rapid formation of radicals and a small fraction of triplets detectable on a nanosecond timescale points to a branched reaction scheme of both molecules.

The rapid formation of radicals can occur from either singlet or triplet states. However, based on TRIR measurements and in agreement with Turro and co-workers,<sup>47</sup> we observed this process occurring on a timescale similar to ISC, therefore, it is more likely that radicals are produced from triplet states. During the radical formation stage, two types of triplet states must be involved: hot and cold triplets. Hot triplets possess much excess energy and either quickly generate free radicals or undergo internal conversion to a lower energy state, producing cold triplets, where  $\alpha$ -cleavage slows down. Additionally, molecules being in this relaxed triplet configuration can further replenish the ground state ( $S_0$ ) by emitting photons detectable in the nanosecond time scale. All considered processes, especially those that cover longer timescales, can be resolved in UV/Vis TA data.

To account for possible branching in the excited-state kinetics, we used a target model, presented in Fig. 10. The model nicely fits the kinetic traces (Fig. 4b and d) and yields SADS, which are shown in Fig. 9c and d. The first two components  $S_1(\text{hot})$  and  $S_1(\text{cold})$  denote either vibrational relaxation and solvation of the excited singlet state and both SADS are nearly similar.

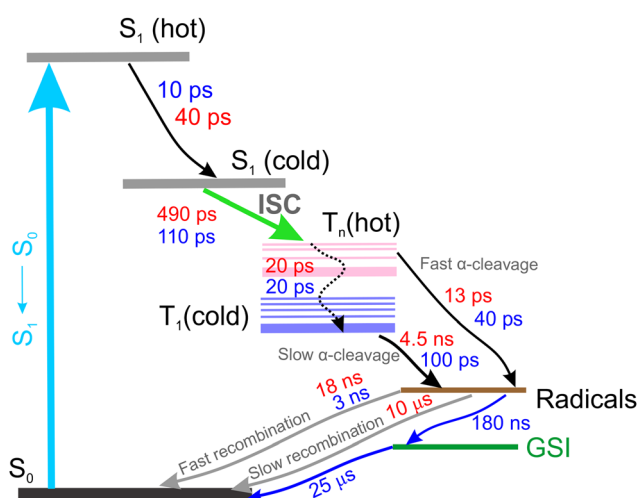


Fig. 10 Reaction scheme, summarizing the mechanisms and their time-scales in **BAPO** (red) and **TPO** (blue). The GSI component represents a small fraction of the **TPO** photoproduct, detectable on a microsecond time scale.

Furthermore, the EADS of TRIR obtained on the same timescale (EADS A and B) have the same vibrational frequencies. However, EADS B exhibits only loss of the ESA, suggesting that the observed changes belong to the same excited state. Of particular importance in this reaction scheme is ISC ( $S_1(\text{cold}) \rightarrow T_n(\text{hot})$ ), producing hot and cold triplets that further boost photofragmentation. In the TRIR data we do not observe significant changes attributed to this process. However, UV/Vis data show some spectral changes when going from singlet to triplet electronic configuration. This can be seen when comparing the UV/Vis EADS  $S_1(\text{cold})$  and  $T_n(\text{hot})$ . After an ISC, 2/3 of the hot triplet state population directly produces mesityl and acylphosphinoyl radicals,<sup>26,35,49</sup> yielding the EADS 'Radicals', which displays the features of the previously measured acylphosphinoyl radical spectrum.<sup>16</sup> The photofragmentation is accompanied by internal conversion of the hot triplet,  $T_n(\text{hot}) \rightarrow T_1(\text{cold})$ , within 20 ps into a more stable triplet configuration,  $T_1(\text{cold})$ , yielding the discussed radical pairs within 4.5 ns. The same cold triplet was detected by measuring the emission dynamics. For **TPO**, the branching ratio of internal conversion to initial radical formation is slightly different (1:2). On the other hand, the TRIR shows radical formation within 130 ps. Assuming such a branching ratio, the radical features are thus expected on a similar time scale. The plausibility of the model could also be confirmed by concentration plots of the species, presented in Fig. S7. When we compare the growth of TRIR ESA at  $1798 \text{ cm}^{-1}$  and the superposition of  $T_n(\text{hot})$  and 'Radicals' species temporal population profiles, we obtain a perfect overlap case of **BAPO**. For **TPO**, UV/Vis TA temporal profile  $c(t)$  superposition of the latter species and  $1798 \text{ cm}^{-1}$  TRIR ESA just grow at similar rates. TCSPC kinetics of **BAPO** at 520 nm can also be reconstructed in the same way as shown in Fig. S8, confirming the agreement between UV/Vis TA and TCSPC measurements. Coming back to naphthalene quenching of **BAPO** triplets, now assuming the presence of hot triplets, the **BAPO**-naphthalene quenching time constant seems to be plausible – once the triplets are formed, they are deactivated by a rapid  $\alpha$ -cleavage and only the cold triplets, where the radical splitting is slower, can undergo energy transfer to naphthalene.

The kinetic traces in Fig. 4b and d indicate that recombination of radicals in both molecules seems to be complex and occurs over several steps; therefore, to simplify this process in the global analysis, we use several parallel decays from the 'Radical' species at different time constants. Initially, the number of radicals is high, so the probability of their mutual collisions at the beginning is highest and leads to a rapid recombination. In other words, at first the radicals recombine quickly due to frequent collisions. As the concentration of radicals decreases, collisions become less frequent, and recombination slows down.<sup>50</sup>

The global analysis shows that the slow recombination of **BAPO** radicals takes about 10  $\mu\text{s}$  (Radicals  $\rightarrow S_0$ ) while the fast occurs in just over 18 ns. In **TPO**, the fast recombination is even faster – 3 ns. For **TPO**, we use an additional component (GSI) to account for slowly decaying transients. The latter component represents a slowly decaying photoproduct of **TPO** radicals,



which is formed within 180 ns and relaxes with a 25  $\mu$ s time constant. It is important to note that a quantitative description of radical decay kinetics cannot be achieved using a single-exponential model only. This is because the actual decay rate depends on the concentration of the produced radicals. Consequently, our kinetic model can only be used to qualitatively estimate the timescale of the relaxation process under our experimental conditions.

## 5. Conclusions

Ultrafast transient absorption measurements on commonly used phosphine oxide photoinitiators **BAPO** and **TPO**, with the absence and presence of naphthalene in solutions, have shown that these molecules undergo ultrafast intersystem crossing (490 ps in **BAPO** and 110 ps in **TPO**), followed by  $\alpha$ -cleavage, creating chemically active acylphosphinoyl radical species. Moreover, time-resolved IR spectroscopy revealed that radicals in both compounds start to form quickly, on a similar time scale as intersystem crossing, indicating that ISC is a primary factor influencing the activity of such compounds. In solution, these radicals were found to recombine in a multi-stage fashion, with a fraction of the radical photoproducts (**TPO**) remaining detectable on a microsecond time scale. A detailed picture of the observed photodynamics (Fig. 10) was constructed in terms of the global analysis kinetic model, providing quantitative inputs for further analysis of the dynamics of photopolymerisation reactions. Observations on **BAPO** and **TPO** clearly show the involvement of triplet states which are hypothesized to be precursors of the radicals important in photoinitiation. If the hypothesis is correct, this would imply that the development of compounds with higher intersystem crossing rates may result in more efficient photoinitiators for radical photopolymerisation.

## Conflicts of interest

There are no conflicts to declare.

## Data availability

The data that support the findings of this study are available in the supplementary material (SI) of this article. See DOI: <https://doi.org/10.1039/d5cp01612f>

## Acknowledgements

MN, MM and MV acknowledge the Universities Excellence Initiative programme by the Ministry of Education, Science and Sports of the Republic of Lithuania under the agreement with the Research Council of Lithuania (project no. S-A-UEI-23-6). RFT and JK gratefully acknowledge funding by the Swiss National Science Foundation, through an Ambizione grant (PZ00P2-216249, to RFT). RFT and MN warmly thank Prof. Dr Eric Vauthey for hosting his research group (RFT) and

for insightful discussions (RFT and MN). We also thank the University of Geneva for financial support, and for computational support through the use of the Bamboo and Yggdrasil HPC Clusters.

## References

- 1 M. Blaicher, M. R. Billah, J. Kemal, T. Hoose, P. Marin-Palomo, A. Hofmann, Y. Kutuvantavida, C. Kieninger, P. I. Dietrich, M. Lauermann, S. Wolf, U. Troppenz, M. Moehle, F. Merget, S. Skacel, J. Witzens, S. Randel, W. Freude and C. Koos, *Light: Sci. Appl.*, 2020, **9**, 1–11.
- 2 T. Abele, T. Messer, K. Jahnke, M. Hippler, M. Bastmeyer, M. Wegener, K. Göpfrich, T. Abele, K. Jahnke, K. Göpfrich Biophysical, E. Group, K. Göpfrich, T. Messer, M. Wegener, M. Hippler and M. Bastmeyer, *Adv. Mater.*, 2022, **34**, 2106709.
- 3 C. Greant, B. Van Durme, J. Van Hoorick and S. Van Vlierberghe, *Adv. Funct. Mater.*, 2023, 2212641.
- 4 B. S. Calin and I. A. Paun, *Int. J. Mol. Sci.*, 2022, **23**, 14270.
- 5 V. Hahn, P. Kiefer, T. Frenzel, J. Qu, E. Blasco, C. Barner-Kowollik and M. Wegener, *Adv. Funct. Mater.*, 2020, **30**, 1907795.
- 6 E. Skliutas, M. Lebedevaitė, E. Kabouraki, T. Baldacchini, J. Ostrauskaitė, M. Vamvakaki, M. Farsari, S. Juodkakis and M. Malinauskas, *Nanophotonics*, 2021, **10**, 1211–1242.
- 7 P. Kiefer, V. Hahn, M. Nardi, L. Yang, E. Blasco, C. Barner-Kowollik and M. Wegener, *Adv. Opt. Mater.*, 2020, **8**, 2000895.
- 8 D. Ladika, G. Noirbent, F. Dumur, D. Gigmès, A. Mourka, G. Barmparis, M. Farsari and D. Gray, *Appl. Phys. A: Mater. Sci. Process.*, 2022, **128**, 745.
- 9 H. Wang, W. Zhang, D. Ladika, H. Yu, D. Gailevičius, H. Wang, C.-F. Pan, P. N. S. Nair, Y. Ke and T. Mori, *et al.*, *Adv. Funct. Mater.*, 2023, **33**, 2214211.
- 10 C. Decker, *Polym. Int.*, 1998, **45**, 133–141.
- 11 A. Alberti, M. Benaglia, D. Macciantelli, S. Rossetti and M. Scoponi, *Eur. Polym. J.*, 2008, **44**, 3022–3027.
- 12 H. F. Gruber, *Prog. Polym. Sci.*, 1992, **17**, 953–1044.
- 13 C. S. Colley, D. C. Grills, N. A. Besley, S. Jockusch, P. Matousek, A. W. Parker, M. Towrie, N. J. Turro, P. M. Gill and M. W. George, *JACS Au*, 2002, **124**, 14952–14958.
- 14 W. Rutsch, K. Dietliker, D. Leppard, M. Köhler, L. Misev, U. Kolczak and G. Rist, *Prog. Org. Coat.*, 1996, **27**, 227–239.
- 15 A. Kowalska, J. Sokolowski and K. Bociong, *Polymers*, 2021, **13**, 470.
- 16 S. Jockusch and N. J. Turro, *J. Am. Chem. Soc.*, 1998, **120**, 11773–11777.
- 17 C. T. Meereis, F. B. Leal, G. S. Lima, R. V. de Carvalho, E. Piva and F. A. Ogliairi, *Dent. Mater.*, 2014, **30**, 945–953.
- 18 S. Miao, W. Zhu, N. J. Castro, M. Nowicki, X. Zhou, H. Cui, J. P. Fisher and L. G. Zhang, *Sci. Rep.*, 2016, **6**, 27226.
- 19 Y. Y. C. Choong, S. Maleksaedi, H. Eng, J. Wei and P. C. Su, *Mater. Des.*, 2017, **126**, 219–225.
- 20 B. Steyrer, P. Neubauer, R. Liska and J. Stampfl, *Materials*, 2017, **10**, 1445.



- 21 B. Steyrer, B. Buseti, G. Harakály, R. Liska and J. Stampfl, *Addit. Manuf.*, 2018, **21**, 209–214.
- 22 T. Zandrini, N. Liaros, L. J. Jiang, Y. F. Lu, J. T. Fourkas, R. Osellame and T. Baldacchini, *Opt. Mater. Express*, 2019, **9**, 2601.
- 23 M. Lunzer, J. S. Beckwith, F. Chalupa-Gantner, A. Rosspeintner, G. Licari, W. Steiger, C. Hametner, R. Liska, J. Fröhlich, E. Vauthey, A. Ovsianikov and B. Holzer, *Chem. Mater.*, 2022, **34**, 3042–3052.
- 24 E. Skliutas, M. Lebedevaitė, S. Rekštytė, S. Lileikis, J. Ostrauskaitė and M. Malinauskas, *Sci. Rep.*, 2020, **10**, 9758.
- 25 G. W. Sluggett, P. F. McGarry, I. V. Koptuyug and N. J. Turro, *JACS Au*, 1996, **118**, 7367–7372.
- 26 T. Sumiyoshi, W. Schnabel, A. Henne and P. Lechtken, *Polymer*, 1985, **26**, 141–146.
- 27 A. Kajiwara, Y. Konishi, Y. Morishima, W. Schnabel, K. Kuwata and M. Kamachi, *Macromolecules*, 1993, **26**, 1656–1658.
- 28 M. Kamachi, A. Kajiwara, K. Saegusa and Y. Morishima, *Macromolecules*, 1993, **26**, 7369–7371.
- 29 J. E. Baxter, R. S. Davidson, H. J. Hageman and T. Overeem, *Macromol. Chem. Phys.*, 1988, **189**, 2769–2780.
- 30 Y. Cui, J. Yang, D. Lei and J. Su, *Ind. Eng. Chem. Res.*, 2020, **59**, 11381–11388.
- 31 J. T. Sutton, K. Rajan, D. P. Harper and S. C. Chmely, *ACS Appl. Mater. Interfaces*, 2018, **10**, 36456–36463.
- 32 A. W. Bassett, A. E. Honnig, C. M. Breyta, I. C. Dunn, J. J. La Scala and J. F. Stanzione, *ACS Sustainable Chem. Eng.*, 2020, **8**, 5626–5635.
- 33 V. Miletic and A. Santini, *Dent. Mater. J.*, 2012, **31**, 717–723.
- 34 S. Jockusch, I. V. Koptuyug, P. F. McGarry, G. W. Sluggett, N. J. Turro and D. M. Watkins, *JACS Au*, 1997, **119**, 11495–11501.
- 35 G. W. Sluggett, C. Turro, M. W. George, I. V. Koptuyug and N. J. Turro, *JACS Au*, 1995, **117**, 5148–5153.
- 36 Y. Feng, I. Vinogradov and N.-H. Ge, *Opt. Express*, 2017, **25**, 26262.
- 37 Y. Feng, I. Vinogradov and N.-H. Ge, *Opt. Express*, 2019, **27**, 20323.
- 38 M. J. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria, M. A. Robb, J. R. Cheeseman, G. Scalmani, V. Barone, G. A. Petersson, H. Nakatsuji, X. Li, M. Caricato, A. V. Marenich, J. Bloino, B. G. Janesko, R. Gomperts, B. Mennucci, H. P. Hratchian, J. V. Ortiz, A. F. Izmaylov, J. L. Sonnenberg, D. Williams-Young, F. Ding, F. Lipparini, F. Egidi, J. Goings, B. Peng, A. Petrone, T. Henderson, D. Ranasinghe, V. G. Zakrzewski, J. Gao, N. Rega, G. Zheng, W. Liang, M. Hada, M. Ehara, K. Toyota, R. Fukuda, J. Hasegawa, M. Ishida, T. Nakajima, Y. Honda, O. Kitao, H. Nakai, T. Vreven, K. Throssell, J. A. Montgomery Jr., J. E. Peralta, F. Ogliaro, M. J. Bearpark, J. J. Heyd, E. N. Brothers, K. N. Kudin, V. N. Staroverov, T. A. Keith, R. Kobayashi, J. Normand, K. Raghavachari, A. P. Rendell, J. C. Burant, S. S. Iyengar, J. Tomasi, M. Cossi, J. M. Millam, M. Klene, C. Adamo, R. Cammi, J. W. Ochterski, R. L. Martin, K. Morokuma, O. Farkas, J. B. Foresman and D. J. Fox, *Gaussian 16 Revision A.03*, 2016, <https://gaussian.com/>.
- 39 A. D. Becke, *Phys. Rev. A: At., Mol., Opt. Phys.*, 1988, **38**, 3098–3100.
- 40 A. D. Becke, *J. Chem. Phys.*, 1993, **98**, 5648–5652.
- 41 C. Lee, W. Yang and R. G. Parr, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1988, **37**, 785–789.
- 42 G. Scalmani and M. J. Frisch, *J. Chem. Phys.*, 2010, **132**, 114110.
- 43 I. H. van Stokkum, D. S. Larsen and R. Van Grondelle, *Biochim. Biophys. Acta, Bioenerg.*, 2004, **1657**, 82–104.
- 44 A. R. Holzwarth, *Biophysical techniques in photosynthesis*, 1996, pp. 75–92.
- 45 R. J. Fernández-Terán, E. Sucre-Rosales, L. Echevarria and F. E. Hernández, *J. Chem. Educ.*, 2022, **99**, 2327–2337.
- 46 S. Straub, J. Lindner and P. Vöhringer, *J. Phys. Chem. A*, 2017, **121**, 4914–4922.
- 47 J. C. Dalton, D. M. Pond, D. S. Weiss, F. D. Lewis and N. J. Turro, *J. Am. Chem. Soc.*, 1970, **92**, 2564–2566.
- 48 R. McNeil, J. Richards and J. K. Thomas, *J. Phys. Chem.*, 1970, **74**, 2290–2294.
- 49 M. D. Forbes, *Carbon-centered free radicals and radical cations: structure, reactivity, and dynamics*, John Wiley & Sons, 2010.
- 50 G. Merkininkaitė, D. Gailevičius, S. Šakirzanovas and L. Jonušauskas, *Int. J. Polym. Sci.*, 2019, **2019**, 3403548.

