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The impact of viscosity asymmetry on phase separating binary mixtures with suspended colloids†

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The introduction of neutrally wetting colloidal particles into coarsening binary fluids is known to arrest the dynamics of the phase separation, as the colloids tend to be captured by the growing interfaces to reduce the free energy of the system. This phenomenon has often been studied in systems with symmetric fluid viscosities. In this study, we investigate the behavior of colloidal particles introduced into asymmetric binary fluids with a viscosity contrast. Our results show that due to the broken symmetry the colloidal particles more easily escape from the interface towards the more viscous fluid, which reduces the lifetime of the jammed phase. Moreover, the presence of colloidal particles near the interfaces promotes the formation of micro-droplets with typical sizes comparable to the colloids.

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1 Introduction

Manipulation of non-equilibrium processes is a convenient tool for designing the self-assembly of novel meso-structured materials.^{1–3} Directing the assembly protocol provides controlled access to metastable configurations that often exhibit more robust mechanical properties than their equilibrium counterparts. A classical example of this process is spinodal decomposition of a binary mixture, where two immiscible fluids, initially homogeneously mixed, start phase separating and forming domains with a variety of topological features. The morphologies of the domains and their time evolution depend on the proportion of the two fluids and their physicochemical properties. Symmetric binary mixtures (composed of liquids with the same properties) have been extensively studied. If the two fluids are present in a similar amount, the system develops interpenetrating domains that coarsen in time following a universal dynamic scaling law.⁴ If one of the two liquids is a

majority phase, then droplets of the minority phase form, which subsequently grow and coalesce.⁵ The addition of colloids with similar affinities for the two liquid phases strongly affects the demixing dynamics, and the stability of the observed structures. Previous studies have established a kinetic pathway to a bicontinuous interfacially jammed emulsion gel (Bijel) composed of two interpenetrating bicontinuous fluid domains frozen by a densely jammed monolayer of colloidal particles at the fluid–fluid interface.⁶ Bijels are soft-solid materials⁷ with highly tunable mechanical properties such as elasticity.

Asymmetric binary mixtures composed of two components with different physical properties (closer to real-world situations) have been much less explored. Experiments with particle-stabilized asymmetric binary mixtures⁸ suggest that the viscosity contrast, in combination with the properties and amount of the colloidal particles, can be a good control parameter governing the complex assembly process that leads to the asymmetric formation of droplets of one fluid within the other. In this paper, we perform lattice Boltzmann simulations to computationally investigate the formation of bijels and droplets in particle-stabilized binary fluids with a viscosity contrast. We first compare the separation process of a colloid-free binary mixture with asymmetric viscosities and establish how the asymmetry in the fluid viscosity affects the coarsening route and the dynamic scaling law with effective viscosity as a parameter. Later we analyze how the introduction of colloidal particles with various packing fractions slows the dynamics of the separation process. At constant effective viscosity and colloidal packing fraction, we explore the effect of neutrally-wetting colloids on the structure and stability of the fluid

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domains, and how such particles can be used to manipulate and control the formation of arrested morphologies.

2 Methods

2.1 The model

Our system consists of colloidal particles of radius a dispersed in a binary fluid mixture undergoing spinodal decomposition. We use a lattice-Boltzmann approach to model a suspension of colloidal particles⁹ and to describe the fluid hydrodynamics that recovers the Navier–Stokes equations.¹⁰ The phase separation process is captured solving the Cahn–Hilliard equation^{11,12} through a phase-field variable, ϕ , representing the composition field of the binary fluid. The coupling between the composition field, and the hydrodynamic stress has been introduced elsewhere.¹³ We confine our study to fluids with equal density ρ , but with a viscosity contrast $\lambda = \eta_{\text{H}}/\eta_{\text{L}}$, where η_{H} (high) and η_{L} (low) are the viscosities of the two components. Within the Lattice-Boltzmann approach, different viscosities are introduced as different relaxation times for each fluid component.

The phase diagram of the binary mixture is controlled by the Landau free energy, $F[\phi] = A\phi^2/2 + B\phi^4/4 + \kappa(\nabla\phi)^2/2$, where ϕ is the local order parameter defined across the system. The parameters A and B control the stability of the mixture; the specific choice $B = -A > 0$ leads to demixing with equally stable states $\phi = \pm\phi^*$. The fluid–fluid interfacial tension is given by $\sigma = (-8\kappa/9)^{1/2}$, and the interface thickness by $\xi = (2\kappa/B)^{1/2}$. We choose a deep quench in which the fluid–fluid interface is sharp on the scale of the colloidal particle, $\xi < a$.

The system is described by a lattice, composed of fluid and solid nodes, to describe the liquid mixture and suspended colloids. The lattice-Boltzmann and the Cahn–Hilliard equations are solved on the liquid nodes. Solid nodes are identified from the positions of the geometric centers of the colloidal particles. The surface of the colloid is determined by boundary links, which connect a solid node with a fluid one. The fluid nodes connected to a solid are referred to as surface nodes. The coupling between the fluid dynamics and the motion of the colloidal particles is provided by the bounce-back on links method.^{13–15}

Fluid particles moving along a boundary link are reflected by the colloidal surface. As a result of the collision, momentum is locally exchanged between the fluid and the colloidal particles. The total hydrodynamic force on each colloid is computed by taking the sum over all boundary links defining the particle and used to compute the colloidal center of mass velocity in a self-consistent approach. At each time step, the position of the colloid is updated following an Euler forward step considering the old and updated velocities.

2.2 Simulation details

We fix the bulk free energy parameters to $B = -A = 0.002$, and $\kappa = 1.4 \times 10^{-3}$ resulting in equilibrium liquid phases at $\phi^* = \pm 1$, interfacial thickness $\xi = 1.14$, and fluid–fluid surface tension $\sigma = 1.58 \times 10^{-3}$. Colloids are neutrally wetting, and are introduced at a packing fraction $\Phi = 4N\pi(a/L)^3/3$, where

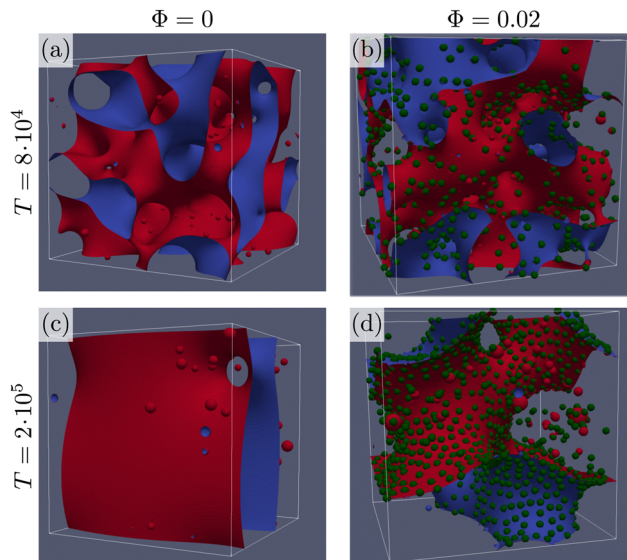


Fig. 1 Snapshots of the interface separating the two fluids with the coloring indicating the component of the fluid at each side of the interface. The fluid with low viscosity $\eta_{\text{L}} = 0.0068$ is shown in red and the fluid with high viscosity $\eta_{\text{H}} = 0.103$, with $\lambda = \eta_{\text{H}}/\eta_{\text{L}} = 15$ is shown in blue. (a) and (b) are captured at time 8×10^4 , and have colloidal packing fractions $\Phi = 0$, and $\Phi = 0.02$, respectively. Systems (c) and (d) are captured at time 2×10^5 , and have colloidal packing fractions $\Phi = 0$, and $\Phi = 0.02$, respectively.

N is the total number of colloids of radius $a = 4.3$ in lattice units.

We perform lattice Boltzmann simulations¹⁶ using Ludwig^{10,17} on a cubic lattice of $L = 256$ sites with periodic boundary conditions. All simulations start from an initial homogeneous fluid with zero mean and uniform fluctuations up to $\pm 10\%$ of the equilibrium value, ϕ^* . Colloids are initially positioned randomly avoiding overlaps. Examples of the simulations with and without colloidal particles, and their visualization, are shown in Fig. 1. The liquid–liquid interface is obtained as the surface where $\phi = 0$, as interpolated from the nearby lattice nodes where ϕ changes sign. The interfacial region corresponds to the volume of the system where $|\phi| < 1/2$. In Fig. 1, we visualize the liquid–liquid interface through two surfaces of constant values of the order parameter $\phi = -0.1$ (red), and $\phi = 0.1$ (blue). Colloidal particles are represented as green spheres. At short times, the snapshots in Fig. 1(a) and (b) show that the system is coarsening and colloidal particles are slowly captured by the interface of the growing domains. Even at a low packing fraction of colloidal particles ($\Phi = 0.02$) we observe a large difference in the behavior at longer times: from the snapshots in Fig. 1(c) and (d) it is apparent that the colloidal particles adsorbed at the interface slow down the coarsening process. Additionally, some colloids that escape the interface can be found covering droplets in the bulk of the large domains.

2.3 Measuring the domain size

In the study of the kinetics of phase separation, it is important to consistently identify characteristic domain sizes. Since the systems we study typically exhibit a small number of domains



of different shapes and sizes, the standard methods relying on Fourier transforms,^{18,19} and the geometry of interfaces^{4,20} do not resolve the characteristic length scales with sufficient precision. Here, we apply the so-called chord length method.^{21–23} We draw straight lines across the system, and define the segments belonging to the same fluid component ($\phi > 0$ or $\phi < 0$). Analyzing the statistics $P(\ell)$ of the segment lengths ℓ enables us to extract the typical length-scales present in the system. In a separating system with many growing domains, the typical domain size is well captured by the mean $\langle \ell \rangle = \int [\ell P(\ell)] d\ell$. The domain size evolution follows three major stages: (i) order parameter diffusion and initiation of domains, $T < T_{\text{init}}$, (ii) viscous domain growth regime, and (iii) saturation regime, $T > T_{\text{sat}}$, featuring macroscopic phase separation controlled by the finite system size. The initial waiting time T_{init} depends on the initial conditions and in order to study the universal properties of the domain growth dynamics, we choose a sufficiently large time to avoid the regime (i). Empirically $\langle \ell(T_{\text{init}}) \rangle = 8$ proves to be a robust choice.

3 Results

Before studying droplet formation with different fluid viscosities, it is informative to determine an effective viscosity of the system and identify a characteristic temporal scale of the coarsening process. Combining the three relevant physical quantities in a symmetric binary fluid: viscosity, surface tension, and density, the scaling of the characteristic length and time scales can be deduced. As shown in ref. 4, the universal scaling is characterized by $L_0 = \eta^2/(\rho\sigma)$, and $T_0 = \eta^3/(\rho\sigma^2)$. In asymmetric fluid mixtures, there is no *a priori* knowledge of whether there exists a universal scaling with a single pair of characteristic time and length scales, and if yes, how the different viscosities are combined to construct T_0 and L_0 . One would expect that the relevant effective viscosity would depend on the details of the model, *i.e.*, on how the viscosity of each component depends on the binary fluid order parameter ϕ . In 1:1 incompressible binary mixtures where viscosity contrast is introduced by a linear composition model $\eta(\phi) = (\phi + 1)/2\eta_L + (1 - \phi)/2\eta_H$, Henry and Tegze²⁴ compared three simple forms of effective viscosity: the arithmetic mean, $\eta_{\text{eff}}^{(a)} = (\eta_L + \eta_H)/2$, the geometric mean $\eta_{\text{eff}}^{(g)} = \sqrt{\eta_L\eta_H}$, and the harmonic mean $1/\eta_{\text{eff}}^{(h)} = (1/\eta_L + 1/\eta_H)/2$, the latter also known as Onuki's formula.²⁵ They found that $\eta_{\text{eff}}^{(g)}$ is the best among the three candidates in collapsing the growth rate onto a single curve $d\ell/dt$ for a certain range of effective viscosity values (in a large viscosity regime). An alternative to the linear model of fluid viscosity has recently been presented by Gidituri *et al.*²⁶ where an Arrhenius model $\eta(\phi) = \eta_{\text{H}}^{(1-\phi)/2}\eta_{\text{L}}^{(1+\phi)/2}$ is explored with a lattice Boltzmann environment. We follow Ledesma *et al.*²⁷ and perform the lattice Boltzmann scheme using viscosity that changes quickly at the interface: $\eta(\phi) = \eta_{\text{H}}\theta(-\phi) + \eta_{\text{L}}\theta(\phi)$, where $\theta(\phi)$ is the Heaviside step function. This choice is convenient since it prevents the appearance of interfaces with viscosities different from the bulk domains.

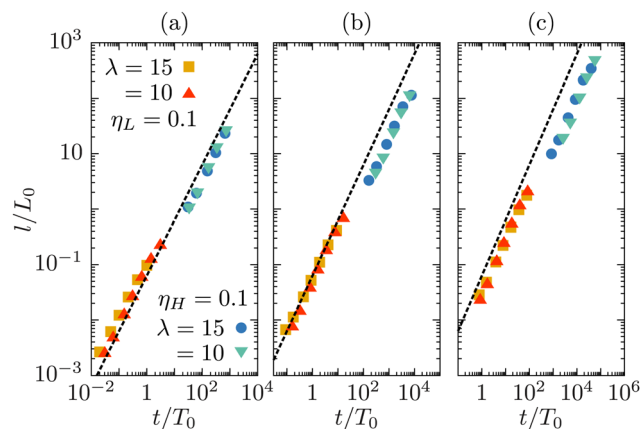


Fig. 2 Dynamic scaling of the typical domain size, at $\Phi = 0$, using the arithmetic (a), geometric (b), and harmonic (c) mean for the effective viscosity. The dashed line corresponds to the universal scaling law, $L/L_0 = 0.063t/T_0$, for a symmetric mixture.⁴ The viscosity values are displayed in the upper part of Table 1. The values of L_0 and T_0 are computed separately in each panel, using the corresponding effective viscosity. The symbols are colored by the viscosity contrast, $\lambda = \eta_{\text{H}}/\eta_{\text{L}}$.

We examine if the coarsening process in our asymmetric model follows a universal scaling law, which would enable us to determine the effective viscosity of the mixture. We follow the temporal evolution of the size of the fluid domains for a phase separating mixture starting from a homogeneous state. In Fig. 2 we show the growth of the domain length, L/L_0 with the scaled time t/T_0 . We compute $L_0 = \eta_{\text{eff}}^2/(\rho\sigma)$, and $T_0 = \eta_{\text{eff}}^3/(\rho\sigma^2)L_0$ using the three simple forms of effective viscosity introduced above. We compare the growth process to the linear universal scaling law, $L/L_0 \sim t/T_0$, as observed in symmetric binary mixtures.⁴ In our simulations we fix the surface tension, and vary η_{H} and η_{L} between 0.006 and 1.5, using the parameters listed in Table 1. The chosen range of parameter values allows us to cover five decades in dimensionless domain sizes compatible with the linear growth regime for symmetric mixtures ($L/L_0 \leq 100$).

We observe that the universal scaling characterizing symmetric mixtures⁴ is roughly recovered in all three cases. However, there are visible deviations from it. The curve obtained with the arithmetic mean hardly deviates from linear behavior in the entire range. It slightly overestimates the domain sizes for viscosity values below $\eta_{\text{H,L}} < 0.1$ and underestimates them

Table 1 Values for the viscosities simulated in Fig. 2 (upper part), and Fig. 3 and 4 (lower part). As an illustration, the viscous length L_0 , and time T_0 are shown using η_{eff}^a as the effective viscosity

η_{H}	η_{L}	λ	$\eta_{\text{eff}}^{(a)}$	L_0	T_0
1.5	0.1	15	0.8	400	2×10^5
1.0	0.1	10	0.55	189	65 000
0.1	0.01	10	0.055	1.89	65
0.1	0.0066	15	0.053	1.778	59
1.03	0.0688	15			
1	0.1	10			
0.917	0.183	5	0.55	189	65 000
0.733	0.367	2			
0.55	0.55	1			



for $\eta_{H,L} > 0.1$. The geometric mean, $\eta_{\text{eff}}^{(g)}$, in contrast, is a good match for large viscosities, but exhibits non-linear growth behavior and underestimates the domain growth rate for low viscosities. The harmonic mean seems to be the worse of the three. Comparing our observations to those of ref. 24 reveals some qualitatively similar trends, however, their evaluation was done at much larger viscosities and with a different model from ours, so quantitative comparison is not productive. We conclude that, even though the dynamic scaling results are not completely conclusive in identifying the optimal form of the effective viscosity, the scaling with the arithmetic average is reasonably close to the universal curve. We thus choose the arithmetic average as an approximation for the effective viscosity of the system. In the following, L_0 , T_0 and η_{eff} will refer to the quantities derived using η_{eff}^a as the effective viscosity.

To explore the effect of colloidal particles on the separation process, we compare the growth of the domain size, $l(t)$, with different colloidal packing fractions, Φ , and different viscosity contrasts, λ . To ensure that phase separation kinetics is dominated by viscosity, the characteristic length L_0 must be larger than the typical domain size, and the characteristic time, T_0 , must be large to keep the typical velocity $V_0 = L_0/T_0$ small enough. Accordingly, we set $\eta_{\text{eff}} = 0.55$ with $L_0 = 190$, and $T_0 = 65\,000$, $V_0 = 2.9 \times 10^{-3}$.

Colloidal particles significantly slow down the kinetics of phase separation, as is obvious from the inset in Fig. 3(a), which depicts the dynamics of the domain size, $l(t)/L_0$, for a fixed $\lambda = 10$ and varying colloidal packing fraction Φ .

The results in Fig. 3(a) for $\Phi = 0.2$ and varying λ reveal three kinetic regimes in system evolution: (i) domain formation and linear growth, $t/T_0 \lesssim 5$, (ii) slow-down regime, $5 \lesssim t/T_0 \lesssim 30$, and (iii) domain growth arrest for $t/T_0 \gtrsim 30$. The viscosity contrast has a strong effect on the arrested regime. For $\lambda = 1$ (purple symbols), the complete arrest of domain growth leads to a gel-like bicontinuous phase, called bijel.⁶ This regime has not been reported before. We managed to identify it since we have covered timescales that are an order of magnitude larger than in the previously reported studies.^{6,28} As $\lambda > 1$ the stability of the bijels is reduced, and a slow kinetic growth is observed in this third regime. The inset of Fig. 3(a) also shows that increasing Φ in asymmetric mixtures does not affect the kinetics of phase separation but decreases the domain size at which the crossover to the arrested regime takes place.

A detailed analysis of the chord length probability distribution, $P(\ell)$, for $\Phi = 0.2$ and $\lambda = 1$ (Fig. 3(b)), and $\lambda = 10$ (Fig. 3(c) and its inset), reveals the formation of small domains with sizes below $\ell = 10$. The temporal evolution of the chord length distributions is qualitatively different in the two cases. In a symmetric binary fluid the number of small domains stabilizes at long times (data not shown). The apparent domain size reduction is due to the droplet formation that alters the average domain size in two ways: by adding smaller domains in the distribution as seen by an increase in $P(\ell)$ at low values, and by splitting the larger domains in multiple segments as seen by introducing a shoulder in the distribution, and by reducing the long segments exponential decay length. In an asymmetric

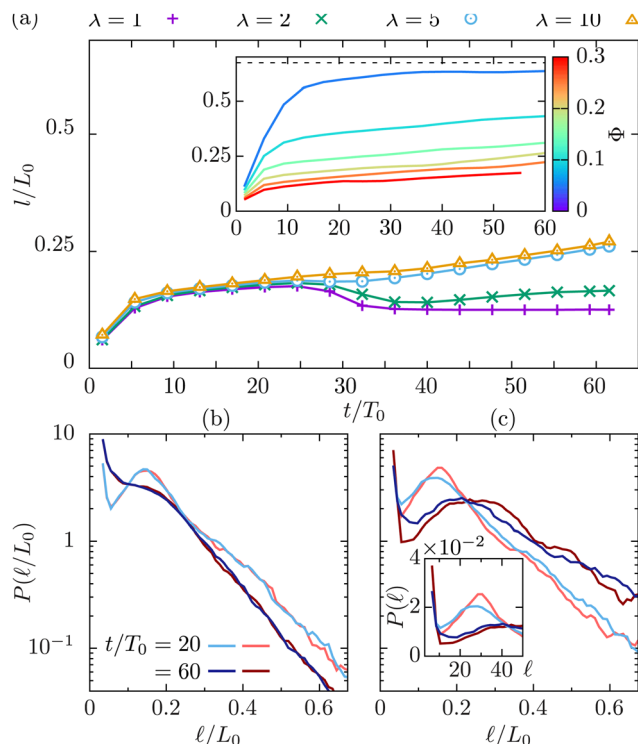


Fig. 3 (a) Domain size l/L_0 at packing fraction $\Phi = 0.2$ for different viscosity contrasts λ at fixed $\eta_{\text{eff}} = 0.55$. The inset shows the effect of varying the colloidal packing fraction Φ at constant viscosity contrast $\lambda = 10$. (b) and (c) Chord length probability density at $\Phi = 0.2$, $\lambda = 1$, and 10, respectively. Shades of red and blue indicate different measurement times $t/T_0 = 20$ for light, and $t/T_0 = 60$ for dark tones. Blue (red) colors indicate the results for the high (low) viscosity fluid. The inset in (c) shows the probability density for the unscaled chord length ℓ .

mixture, although we observe more droplets in that case, the fluid domains continue growing as shown by both the increase of the decay length of the exponential tail and the displacement of the peak of the $P(\ell)$ to larger ℓ values. Systems at large packing fractions form the bijel structure but the viscosity contrast accelerates the slow relaxation process towards the complete phase separation.

To visualize the structuring of the fluid, and colloidal particles, we present cuts of the system at constant $z = z_0$ values (see Fig. 4) at different stages of the evolution, $t/T_0 = 20, 60$. The fluid regions in the figure are colored as follows: if the lattice element is a part of a macroscopic percolating domain, it is colored in pale red (for η_L component) or pale blue (η_H); similarly, the smaller domains are colored in dark red (η_L) and dark blue (η_H). A colloidal particle is visualized if its center lies within one radius from the selected plane ($|z_i - z_0| \leq a$). The apparent radius of the displayed colloidal particle is given by $r_i = (a^2 - (z_i - z_0)^2)^{1/2}$, the colloids are shaded in grey according to the module of their instantaneous velocity $|\mathbf{V}_i|/V_0$. The instantaneous velocity of the colloids trapped at the interface is a convenient probe for the temporal evolution of the corresponding domains. At packing fraction $\Phi = 0.2$ (Fig. 4(b) and (d)), the instantaneous velocity of the particles trapped at the interface is quite different between the symmetric ($\lambda = 1$)



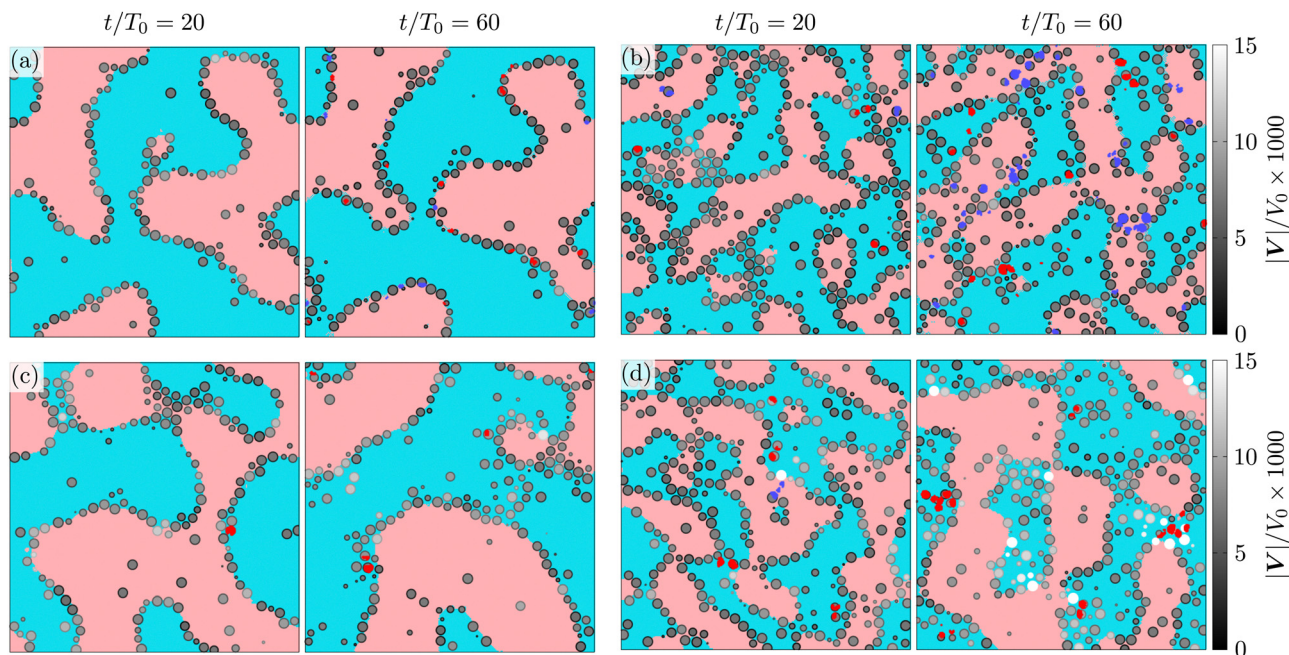


Fig. 4 Two dimensional cuts of the system at constant z_v for $\eta_{\text{eff}} = 0.55$ at $t/T_0 = 20$ or 60 as indicated by the labels. The macroscopic, and percolating clusters are identified by light red, and light blue colors, for the low and high viscosity fluids, respectively. The droplets are highlighted with saturated red and blue patches. Colloidal particles whose centers of mass $r_c = (x_c, y_c, z_c)$ lie within one radius from the plane of view, $(|z_c - z_v| \leq a)$, are visualized by circles with apparent radius proportional to their distance to the plane, and shaded in grey-scale according their speed. In (a) $\lambda = 1$, $\Phi = 0.1$, in (b) $\lambda = 1$, $\Phi = 0.2$, in (c) $\lambda = 10$, $\Phi = 0.1$, and in (d) $\lambda = 10$, $\Phi = 0.2$.

and asymmetric ($\lambda = 10$) mixtures, which indicates different domain dynamics: in asymmetric mixtures the colloids are advected two- to three-times faster than in the symmetric mixtures.

To gain further insight, we have analyzed the fluid and colloidal velocities both in the bulk (lattice nodes with $|\phi| > 0.5$), and at the interface ($|\phi| < 0.5$). Colloids are classified as interfacial if at least one of their surface nodes lies in the interfacial region. In Fig. 5 we plot the time evolution of the velocities of fluid phases and colloids – first in the bulk regions (a) and (b), and separately in the interfacial region (c) and (d). The bulk fluid velocities, Fig. 5(a) and (b), show a non-monotonous behaviour revealing three kinetic regimes in system evolution, which correlate with the observations in Fig. 3(a). In the domain growth regime, $t/T_0 \lesssim 10$, the fluid and colloidal velocities quickly decrease from their initial values, both in symmetric and asymmetric mixtures (see also the inset of Fig. 5(a)). For symmetric mixtures, the slow-down regime is characterized by a plateau with essentially constant values of fluid velocities, followed by an increase that eventually converges to a second plateau at large times. The transition between the second and third regimes is associated with a structural rearrangement of the interfacial colloids, which is reflected in the non-monotonous behaviour of the interfacial fluid velocity (a peak at $t/T_0 \sim 25$ in Fig. 5(d)). In asymmetric mixtures, the bulk fluid velocities behave qualitatively differently: they are larger in the high viscosity phase, and the plateau values are less clearly expressed than in the symmetric case. In order to understand this behaviour better, we computed the number of colloids that reside in the high viscosity

fluid, N_H , in the low viscosity fluid, N_L , and in the interface region, N_I . Clearly, $N = N_H + N_L + N_I$, and the corresponding colloidal fractions are $\zeta_k = N_k/N$, $k \in \{L, H, I\}$. In Fig. 6(a), we show their time evolution for symmetric (dashed lines) and asymmetric mixtures (solid lines). We observe that in symmetric mixtures in the entire simulation time the colloids predominantly reside at the interface. In the second and third kinetic regime, the colloids are symmetrically expelled into the bulk at a small rate, which is consistent with a gel-like behaviour of the *bijels*.⁶ In contrast, in asymmetric mixtures the colloids are expelled from the interface at a higher rate, and preferentially into the high viscosity bulk phase. This can be attributed to the asymmetric position of the colloids at the fluid interface, which is shown in Fig. 6(b), where we plot the average distance of the centers of the colloids from the fluid–fluid interface.

We calculate the dimensionless distance δ_i of a colloid i from the liquid interface by first identifying those surface nodes (fluid nodes in contact with the colloid i) that lie within the interfacial region. This set of nodes, denoted as $\{q_{i,l}\}$, can be thought of as representing an extended object in three dimensions. We use the singular value decomposition technique²⁹ to determine the principle axes of this kind of object and identify the axis that describes the normal to the interface, \hat{N}_i . By convention, the normal is pointing into the low viscosity phase. Once we know the direction of the normal to the interface at the position of the colloid X_i , we can calculate the distance $\delta_i = \langle \hat{N}_i \cdot (q_{i,l} - X_i) \rangle / a$. According to our convention, the positive values of δ_i imply colloids residing predominantly in the high viscous phase. Fig. 6(b) shows that in the asymmetric mixtures the interfacial



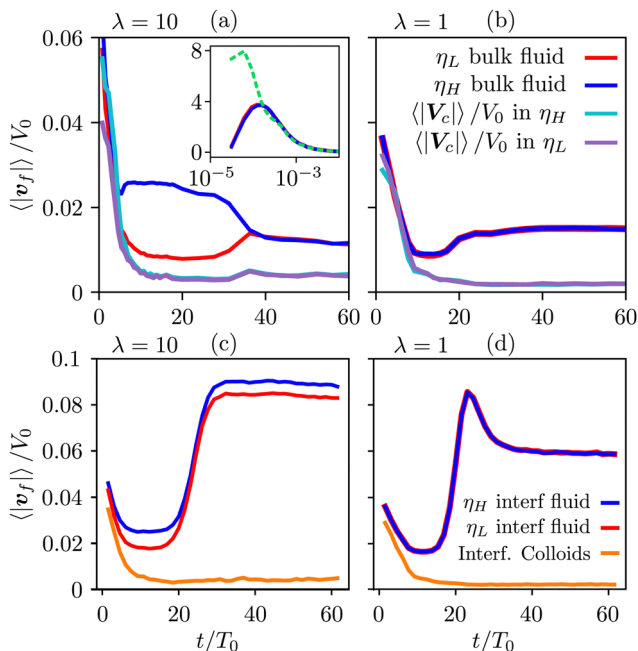


Fig. 5 Averaged velocities of both fluid phases and of colloids for $\eta_{\text{eff}} = 0.55$ and $\Phi = 0.2$. Red: low viscosity fluid, blue: high viscosity fluid, purple: colloids residing in low viscosity fluid, cyan: colloids residing in high viscosity fluid, orange: colloids trapped in the interfacial region. (a) and (b) Bulk fluid and colloids in the bulk fluid phases; (c) and (d) fluids in the interfacial region and interfacial colloids. The left column (a) and (c) is for asymmetric mixtures with viscosity contrast $\lambda = 10$, and the right column (b) and (d) for symmetric mixtures ($\lambda = 1$). The inset in (a) shows the initial stages of the phase separation process. The dashed green line is the average velocity of the colloids, which at this early stage cannot be separated into bulk and interfacial types.

colloids, before they are expelled into the bulk, are positioned asymmetrically towards the high viscous phase.

The distribution of the distances δ_i , shown in Fig. 6(c) and illustrated in the snapshots in Fig. 6(d) shows that for $\lambda = 1$, the colloidal distribution around the interface is symmetric and broadens with time. This broadening is associated with the increased tension following the structural rearrangement in the third kinetic regime.³⁰

In contrast, for $\lambda = 10$, the distribution is shifted towards the positive values of δ . The basic mechanism behind this displacement is the viscous lift force: since the fluid velocities of both phases at the interface are similar (Fig. 5(c)), and the viscosity contrast is 10, we expect a non-vanishing viscous drag force directed towards the high viscosity phase. We note that this explanation is based on a simple planar configuration with a simple flow pattern. Therefore, it can provide a qualitative rather than quantitative description of the mechanism for the onset of asymmetric interfacial stresses and the resulting asymmetric positioning of the colloids at the interface.

The asymmetric positioning, in turn, lowers the barrier for colloidal expulsion into the high viscosity bulk phase. The expulsion of a particle creates a locally enhanced fluid flow, which explains why in Fig. 5(a) and (c) the blue curves lie above the red ones. Therefore, at $\lambda = 10$, we do not observe a complete

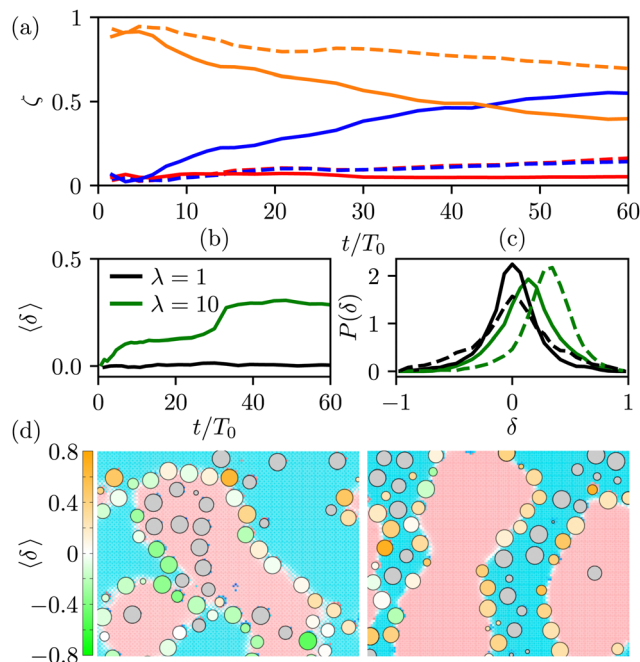


Fig. 6 Temporal evolution of the spatial distribution of the colloids with the overall volume fraction $\Phi = 0.2$, at $\eta_{\text{eff}} = 0.55$. (a) Fraction of colloids in the high viscosity (ζ_H , blue curves), and in the low viscosity (ζ_L , red curves) bulk phase. The fraction of colloids trapped in the interfacial region, ζ_i , is shown in orange. The solid lines are for asymmetric mixtures with $\lambda = 10$, and the dashed lines for symmetric mixtures with $\lambda = 1$. (b) Time evolution of the averaged distance of the colloidal center of mass from the fluid interface, $\langle \delta \rangle$. Positive values of $\langle \delta \rangle$ imply that the colloids are predominantly immersed in the high viscosity phase. The solid line is for symmetric mixtures ($\lambda = 10$), and the dashed line for $\lambda = 1$. (c) The distribution of the values of δ for symmetric (black) and asymmetric (green) mixtures. The solid lines correspond to distributions averaged over a time window $10 < t/T_0 < 20$ (second kinetic regime), and the dashed lines for $t/T_0 > 40$ (third kinetic regime). (d) 2D cuts of a small section of the system at constant z . The interfacial colloids are represented by circles whose size is scaled by the z coordinate, and coloured according to their distance from the interface δ_i ; for $\lambda = 1$ (left plot) and $\lambda = 10$ (right plot).

arrest of the domain coarsening, which indicates that the viscosity contrast destabilises the bijels.[‡]

During the expulsion from the interface into the bulk, it is likely that the colloids drag the surrounding fluid with them, which can result in the formation of droplets of one fluid phase in the bulk of the other one. Our observations of such droplets are shown in the snapshots in Fig. 4, and in the movies in the online ESI.[†] The observed droplets are typically stabilized by the colloidal particles. Fig. 7 shows the statistics of the observed micro-droplets measured in the simulations within the second and third kinetic regimes. In the symmetric fluid the amount of red and blue droplets is statistically equivalent, while for asymmetric mixtures where most free colloids are in

[‡] Unlike the fluid velocities, the average colloidal velocities in the bulk, $\langle |V_c| \rangle$, decrease monotonically and are the same for symmetric and asymmetric binary mixtures. Due to the parameters chosen in our work, the typical domain size in the stabilization phase is only about four colloidal diameters, thus the colloids cannot be considered as passive local probes of the fluid velocity field and their average speed will generally be lower than that of the bulk fluid.



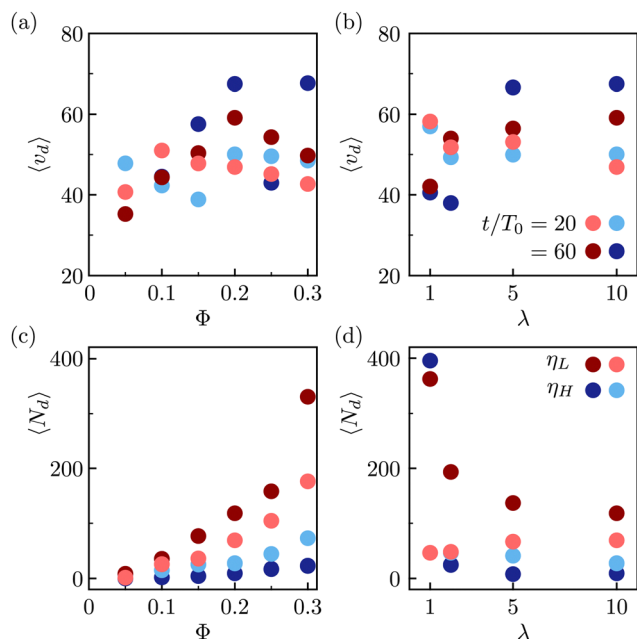


Fig. 7 The average observed droplet volume $\langle v_d \rangle$ and number $\langle N_d \rangle$ at $\eta_{\text{eff}} = 0.55$ and at two times: $t/T_0 = 20$ (lighter shades), and $t/T_0 = 60$ (darker shades). In shades of red (blue) the results for low (high) viscosity components of the fluid are presented. (a) Volume per droplet $\langle v_d \rangle$ as a function of colloidal packing fraction Φ at a constant $\lambda = 10$; (b) volume per droplet as a function of the viscosity contrast λ at constant $\Phi = 0.2$; (c) number of droplets at constant $\lambda = 10$, (d) number of droplets at constant $\Phi = 0.2$.

the high-viscosity phase, the droplets of the low viscosity fluid dominate. To identify the droplets, we performed the cluster analysis: two neighboring lattice nodes belong to the same fluid domain if they share the same fluid component (same sign of ϕ). The analysis reveals two macroscopic (percolating) domains with approximately half the lattice nodes each and on top of the macro-domains, there is a dispersion of microscopic domains with volumes ranging from tens to hundreds of lattice nodes. The quantitative analysis reveals that, in Fig. 7(a) and (b), the average volume per droplet lies around $\langle v_d \rangle = 50 \pm 10$ lattice cells, approximately a quarter of a single colloidal particle. Moreover, the number of formed droplets increases both with the packing fraction of colloidal particles and with simulation time. Whenever there is a viscosity contrast, $\lambda \neq 1$, the number of observed droplets presents a strong asymmetry with more droplets in the low viscosity fluid forming within the high-viscosity fluid domain. We speculate that the larger number of small droplets observed in the symmetric mixtures is associated to the larger free energy barrier for the colloidal escape, which in turn creates a larger distortion of the fluid interface, and as a consequence a larger number of emerging small droplets.

4 Summary

We have shown that viscosity asymmetry plays a crucial role in the survival of long-lived bijels. The asymmetry in the fluid viscosities breaks the symmetry in colloidal positions at the

fluid interface and facilitates the colloidal escape process by which bijels relax into an equilibrium structure. The escape mechanism accumulates colloids in the high viscosity fluid even in the absence of thermal fluctuations. While the time-scales associated with the domain coarsening arrest are controlled by the viscosity contrast, we observed that in the presence of colloids, the pathways leading to a coarsening process towards a bijel or complete phase separation remain unchanged. Finally, the accumulation and expulsion of colloidal particles from the interface creates complex fluid motion and favours the formation of small droplets with a characteristic size almost independent of the control parameters, *i.e.*, the packing fraction and the viscosity contrast.

Conflicts of interest

There are no conflicts to declare.

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