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# Precise control of TiO<sub>2</sub> overlayer on hematite nanorod arrays by ALD for the photoelectrochemical water splitting†

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The short lifetime of electron–hole pairs and high electron–hole recombination rate at surface states significantly limit the practical applications of hematite ( $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>) photoanodes in photoelectrochemical (PEC) water splitting. Surface modification with a TiO<sub>2</sub> overlayer has been demonstrated to be an efficient way to improve the PEC performance. However, a fine control of the TiO<sub>2</sub> overlayer and a deep understanding of the impact of the TiO<sub>2</sub> overlayer with variable thickness on the PEC performance, to the best of our knowledge, has yet to be done. Here, a conformal ultrathin TiO<sub>2</sub> overlayer is successfully deposited on hydrothermal grown one-dimensional hematite nanorod arrays by atomic layer deposition. The morphology and thickness of the TiO<sub>2</sub> overlayer can be precisely controlled. The effect of the thickness of the TiO<sub>2</sub> overlayer on the overall water splitting efficiency of hematite photoanodes under visible and UV light has been systematically investigated. The charge excitation and transfer mechanism at the semiconductor–electrolyte interface has also been studied.

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## 1. Introduction

To solve the energy crisis, it has become urgent to develop novel technologies that can efficiently utilize solar energy and convert sunlight into chemical fuels.<sup>1,2</sup> Photoelectrochemical (PEC) water splitting using semiconductors has been extensively studied for solar hydrogen generation.<sup>3–5</sup> Hematite ( $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>) is a promising n-type photoanode material because of its favorable optical band gap (2.0–2.2 eV), chemical stability, earth abundance and low cost.<sup>6,7</sup> One-dimensional (1D) hematite nanostructures such as nanorods, nanowires, nanosheets and nanotubes can enhance charge collection efficiency in PEC water splitting due to their unique morphology and large surface to volume ratio.<sup>8–11</sup> Nevertheless, the water splitting efficiency of hematite is severely limited by the very short

lifetime of electron–hole pairs, short hole diffusion length (2–4 nm), poor charge carrier mobility and the high charge recombination rate at surface states.<sup>12,13</sup>

Numerous efforts have been made to overcome these issues and improve the PEC efficiency of hematite, such as element doping, constructing heterojunction and surface modification.<sup>14–17</sup> Among these strategies, surface modification of semiconductor photoanodes has proven to be an effective strategy to improve the charge separation and enhance the efficiency of water oxidation of hematite photoanodes. TiO<sub>2</sub> has attracted considerable interest as a photoanode for water splitting due to its good chemical stability, outstanding optoelectronic properties and nontoxicity.<sup>18–20</sup> Recently, many studies have focused on the modification of TiO<sub>2</sub> overlayer on hematite for PEC water splitting applications. For instance, Yang *et al.* showed that an ultrathin TiO<sub>2</sub> layer grown on hematite could lead to better charge separation and a 100 mV cathodic shift of the turn on potential because of the enhanced charge separation.<sup>12</sup> Feng *et al.* displayed an effective strategy in promoting the charge transfer by decorating TiO<sub>2</sub> at the grain boundaries of hematite.<sup>21</sup> However, these reports mainly demonstrated an improved PEC performance of hematite photoanodes by surface treatment with TiO<sub>2</sub>. The precise control over the morphology and thickness of the TiO<sub>2</sub> overlayer remains challenging. Furthermore, the impact of the thickness of the TiO<sub>2</sub> overlayer on the PEC performance and the charge transfer mechanisms at the semiconductor–electrolyte interface, to the best of our knowledge, have yet to be investigated.

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A TiO<sub>2</sub> overlayer with optimized thickness can increase electronic conductivity and suppress surface charge carrier recombination.<sup>18,19,22</sup> So far, a variety of techniques have been used to deposit a TiO<sub>2</sub> overlayer, including spin coating, atomic layer deposition (ALD), sputtering, chemical bath deposition and dip coating.<sup>21,23–26</sup> ALD is a thin film deposition technique based on self-terminating surface reactions, which allows the conformal and homogeneous coating on high aspect ratio substrates, with precise control of thin film thickness at the Ångström scale.<sup>27–29</sup> Progress has already been made in the formation of homogeneous TiO<sub>2</sub> overlayer on hematite by ALD.<sup>27,30,31</sup> Nevertheless, achieving a conformal, pinhole-free and high quality ultrathin TiO<sub>2</sub> film is still challenging due to limited understanding of the TiO<sub>2</sub> growth at early stages.

In this work, TiO<sub>2</sub> continuous ultrathin films with variable thickness are successfully deposited onto 1D hematite nanorod arrays by ALD. The morphology and the thickness of the TiO<sub>2</sub> overlayer can be accurately controlled. Fe<sub>2</sub>O<sub>3</sub> can be excited to generate electron–hole charge carriers under both visible and UV light owing to its low band gap, while TiO<sub>2</sub> can only be excited under UV light because of its large band gap (3.2 eV for anatase).<sup>32,33</sup> To understand the role of a TiO<sub>2</sub> overlayer and the underlying charge transfer mechanism in PEC water splitting, the PEC performances of as-synthesized TiO<sub>2</sub>/Fe<sub>2</sub>O<sub>3</sub> under visible and UV light as a function of the TiO<sub>2</sub> thickness were systematically investigated. The morphology and thickness of the resulting films were characterized by high-resolution transmission electron microscopy (HRTEM). The energy level alignment was analyzed by X-ray photoelectron spectroscopy (XPS). This work provides a promising strategy to elucidate the role of a TiO<sub>2</sub> overlayer and the charge transfer mechanisms in PEC water splitting of hematite, which can be further applied to other nanostructured hematite photoanodes.

## 2. Experimental section

### 2.1 Materials

Titanium tetrachloride (TiCl<sub>4</sub>, 99%) was purchased from ABCR GmbH. Ozone (O<sub>3</sub>) was provided by a BMT803N ozone delivery system using pure oxygen at a pressure of 0.5 bar, and nitrogen (99.99%) was purchased from Air Liquide. Iron(III) chloride hexahydrate (FeCl<sub>3</sub>·6H<sub>2</sub>O) and urea (CH<sub>4</sub>N<sub>2</sub>O) were purchased from Carl Roth and Sigma-Aldrich, respectively. VWR GmbH supplied absolute ethanol and acetone for synthesis. Water with a resistivity of 18.2 MΩ cm<sup>-1</sup> was used in all syntheses. 1 M KOH standard aqueous solution (pH = 13.6) was supplied by Carl Roth. All other chemicals and reagents were of analytical grades. SnO<sub>2</sub>:F transparent conducting glass (FTO, TEC8, thickness 2.2 mm, resistance 6.70 ± 0.27 Ω square<sup>-1</sup>) was purchased from Ossila. Silicon wafers (B014002) were provided by Siegert wafer GmbH and used after cleaning with piranha solution.

### 2.2 Preparation of Fe<sub>2</sub>O<sub>3</sub> nanorods

Hematite (α-Fe<sub>2</sub>O<sub>3</sub>) nanorod arrays were grown on FTO substrates through an easy hydrothermal method followed by

an annealing procedure in air flow. Prior to the synthesis, an accurate cleaning of the FTO-glass is required, including ultrasonication with absolute acetone, absolute ethanol, and ultrapure water (MilliQ) for 10 minutes, respectively. 1.8 mmol FeCl<sub>3</sub>·6H<sub>2</sub>O and 4.2 mmol CH<sub>4</sub>N<sub>2</sub>O were dissolved in 15 mL of water and stirred at room temperature for 20 min. Subsequently, the solution was transferred into a Teflon-lined stainless-steel autoclave with the FTO substrates placed with the conductive side facing down and then heated at 100 °C in the oven for 12 h. The product of β-FeOOH forming a uniform yellow film on the FTO substrates was thoroughly washed with absolute ethanol and water and dried with nitrogen gas. After a two-step annealing treatment of β-FeOOH in a muffle furnace (250 °C for 30 min and 500 °C for 30 min with a heating rate of 10 °C min<sup>-1</sup>), α-Fe<sub>2</sub>O<sub>3</sub>, hematite samples were obtained.

### 2.3 Preparation of TiO<sub>2</sub>/Fe<sub>2</sub>O<sub>3</sub> heterostructures

TiO<sub>2</sub> was deposited on hematite and Si-wafers using a PLAY 2018-01 (CTECHnano) thermal ALD system. Before starting the ALD process, the ALD system was evacuated, and the temperature of the ALD chamber was well stabilized. The baseline pressure was maintained at 7.3 × 10<sup>-1</sup> mbar with a 40 sccm of nitrogen flow. TiCl<sub>4</sub> and H<sub>2</sub>O were used as the metal precursor and the oxygen source, respectively. Both the precursors were kept in stainless steel canisters at room temperature. The temperatures of the manifold and the ALD chamber were maintained at 100 °C and 120 °C, respectively. TiCl<sub>4</sub> and H<sub>2</sub>O were introduced into the ALD chamber in a sequence using nitrogen as purging and carrier gas. Typically, pulse time, exposure time and purge time are 0.5 s, 50 s and 30 s for TiCl<sub>4</sub> and 0.15 s, 40 s and 30 s for H<sub>2</sub>O, respectively. The samples were deposited with 10, 20, 40, 80 and 150 ALD cycles and named as TiO<sub>2</sub>-10/Fe<sub>2</sub>O<sub>3</sub>, TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub>, TiO<sub>2</sub>-40/Fe<sub>2</sub>O<sub>3</sub>, TiO<sub>2</sub>-80/Fe<sub>2</sub>O<sub>3</sub>, and TiO<sub>2</sub>-150/Fe<sub>2</sub>O<sub>3</sub>, respectively. After TiO<sub>2</sub> coating, the hematite samples were annealed at 450 °C for 2 h in air.

### 2.4 Photoelectrochemical measurements

PEC measurements were carried out on a Bio-Logic VMP3 potentiostat in a homemade three-electrode electrochemical system with a 0.385 cm<sup>2</sup> working area. Platinum sheet (1 × 1 cm<sup>2</sup>) and Hg/HgO (1.0 mol L<sup>-1</sup> NaOH) were employed as counter and reference electrodes, respectively. 1 M KOH (pH = 13.6) was used as electrolyte. The electrolyte solution was degassed with Ar flow for 20 min before the PEC measurement to remove the dissolved oxygen. LEDs nominal wavelengths (λ = 430 nm, M430L4 Spectrum) and (λ = 365 nm M365L4 Spectrum) were used as the light source. The light power density for LED illumination under 430 nm and 365 nm were 57 mW cm<sup>-2</sup> and 26 mW cm<sup>-2</sup>, respectively. The LEDs power was measured with a Compact Power and Energy Meter (PM100D, ThorLabs). Linear sweep voltammetry (LSV) curves were measured in the potential window -0.1 V to 0.6 V bias *versus* Hg/HgO with a scan rate of 10 mVs<sup>-1</sup>. Electrochemical impedance spectroscopy (EIS) measurements were carried out in a frequency range from 100 kHz to 0.1 Hz with an amplitude of 5 mV.



## 2.5 Characterizations

Spectroscopic ellipsometry (SENpro Sentech) was utilized to measure the relative thickness of the TiO<sub>2</sub> thin film on Si-wafers with a wavelength range from 370 to 1000 nm under a 70° incidence angle. Field-emission scanning electron microscopy (FESEM) was performed with Zeiss SUPRA 40 instruments. All micrographs were acquired at a 5 kV accelerating voltage. Transmission electron microscopy (TEM) was performed on a Philips CM 200 equipped with a LaB<sub>6</sub> filament. High-resolution transmission electron microscopy (HR-TEM), high-angle annular dark field scanning transmission electron microscopy and energy dispersive X-ray analysis elemental mapping were performed on a FEI Talos F200S scanning/transmission electron microscope operated at 200 kV. Raman spectrometer (Thermo Scientific) equipped with an ANDOR CCD camera and excited by a solid-state laser operating at 2.33 eV (532 nm) was used to measure Raman scattering. The spectra were recorded in the range 100–800 cm<sup>-1</sup> with an acquisition of 16 s. Before measurement, the Raman spectra were calibrated using the 520 cm<sup>-1</sup> phonon mode of Si. UV-vis-NIR diffuse reflectance spectra were measured with a PerkinElmer Lambda 1050+ UV-vis-NIR spectrophotometer in the range of 200–800 nm. XPS measurements were performed using a JEOL JPS-9030 setup with a base pressure of 2 × 10<sup>-9</sup> mbar, employing the K<sub>α</sub>-radiation of a non-monochromated Al X-ray source ( $h\nu = 1486.6$  eV) for excitation and a hemispherical analyzer to detect the kinetic energy of the emitted electrons. A pass energy of 50 eV was used for the survey scans and a pass energy of 10 eV for the narrow scans. The binding energy scale was calibrated by measuring a sputter-cleaned polycrystalline gold foil and setting the Au 4f<sub>7/2</sub> peak to 84.00 eV. The core level spectra were fitted with CasaXPS, using Voigt-functions.<sup>34</sup>

## 3. Results and discussions

### 3.1 Structural characterizations

Fig. 1a represents the synthesis procedure of ALD-grown TiO<sub>2</sub> film on hematite nanorods (TiO<sub>2</sub>/α-Fe<sub>2</sub>O<sub>3</sub>) with FTO glass as a substrate. In detail, vertically aligned β-FeOOH nanorods were prepared by hydrothermal method and subsequently converted to hematite nanorods *via* post annealing treatment. Afterward, the TiO<sub>2</sub> film with various thicknesses was deposited on hematite nanorods by ALD.

The phase composition and structure of β-FeOOH and α-Fe<sub>2</sub>O<sub>3</sub> were characterized by XRD (Fig. S1, ESI†). The diffraction peak at 35.2° of 2θ can be assigned to the (211) reflection of tetragonal FeOOH (JCPDS 34-1266, red line). The diffraction peaks located at 35.9° and 64.5° can be ascribed to the (110) and (300) reflections of α-Fe<sub>2</sub>O<sub>3</sub>, respectively (COD 96-901-5504, black line). In addition, the reflections located at 26.6°, 33.9°, 38.0°, 51.8°, 61.9° and 66.0° and marked with asterisks correspond to the (110), (101), (200), (211), (130) and (301) reflections of the cassiterite phase of SnO<sub>2</sub>, respectively. Both samples show the characteristic signals of SnO<sub>2</sub>, which come from the FTO substrate.

The morphology and microstructure of β-FeOOH and hematite nanorods were thoroughly investigated by TEM (Fig. S2, ESI†). The typical solid β-FeOOH nanorods are uniform, with diameters ranging from 40 to 80 nm (Fig. S2a and b, ESI†). The main rings in the selected area electron diffraction (SAED) pattern correspond to the (200), (310) and (211) planes for tetragonal FeOOH (Fig. S2c, ESI†). Fig. S2d and e of the ESI† illustrate TEM images of pristine hematite nanorods. The concentric rings in the SAED corresponding to the (102), (104), (110) and (113) planes can be indexed to hematite, indicating the polycrystalline structure of Fe<sub>2</sub>O<sub>3</sub> (Fig. S2f, ESI†).

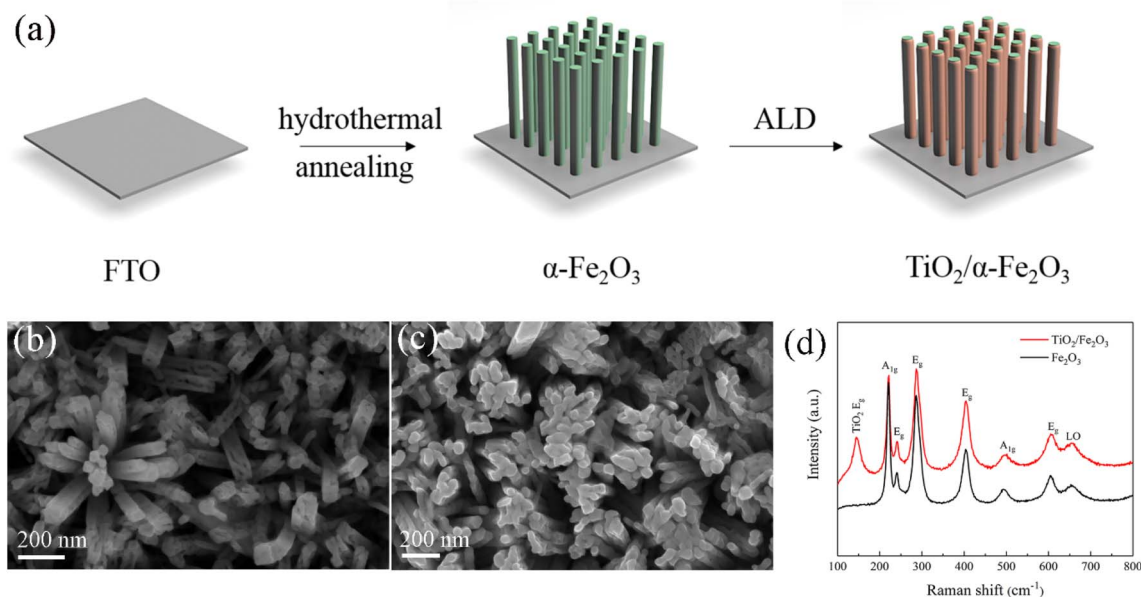


Fig. 1 (a) Schematic view of the synthesis process for TiO<sub>2</sub>/α-Fe<sub>2</sub>O<sub>3</sub> hierarchical nanostructure. (b) SEM images of pristine hematite nanorods, (c) SEM images of 40ALD cycles of TiO<sub>2</sub> deposited on pristine hematite, (d) Raman-active modes of pristine hematite and TiO<sub>2</sub> deposited on hematite.





The morphology and homogeneity of the pristine and TiO<sub>2</sub>-coated hematite samples were examined by FESEM. Fig. 1b displays that the pristine hematite nanorods are vertically aligned on the FTO. The diameter of the nanorods is around 60 to 80 nm. The SEM images of TiO<sub>2</sub>-40/Fe<sub>2</sub>O<sub>3</sub> demonstrate a well-preserved hematite morphology (Fig. 1c). However, it is difficult to distinguish the TiO<sub>2</sub> film because of its low thickness. The composition of the hematite before and after ALD-TiO<sub>2</sub> coating was confirmed by Raman spectra (Fig. 1d). The Raman peaks at 220, 241, 286, 404, 494 and 606 cm<sup>-1</sup> are assigned to the hematite structure. The LO peak center at 658 cm<sup>-1</sup> is detected, which is attributed to the disorder-induced breaking of Raman symmetry properties.<sup>14,35</sup> After TiO<sub>2</sub> deposition, an intense Raman peak located at 144 cm<sup>-1</sup> is observed, which can be indexed to the E<sub>g</sub> mode of anatase TiO<sub>2</sub>.<sup>19</sup> The Raman peaks of the Fe<sub>2</sub>O<sub>3</sub> nanorods are not affected by ALD of TiO<sub>2</sub>.

The HRTEM images of TiO<sub>2</sub>-coated Fe<sub>2</sub>O<sub>3</sub> nanorods with a different number of ALD cycles (from 10 to 150 cycles) is

shown in Fig. S3, ESI.† At 10 ALD cycles, a continuous TiO<sub>2</sub> film coverage cannot be clearly observed due to the small amount of TiO<sub>2</sub> coating (Fig. S3a, ESI†). After 20 ALD cycles, Fe<sub>2</sub>O<sub>3</sub> nanorods are homogeneously and uniformly coated with an ultrathin TiO<sub>2</sub> film (Fig. S3b–e, ESI,† red dotted lines). The thicknesses of the TiO<sub>2</sub> on Fe<sub>2</sub>O<sub>3</sub> nanorods increase with the further increase of the number of ALD cycles. Fig. S3f of ESI† shows the lattice fringes of TiO<sub>2</sub>-150/Fe<sub>2</sub>O<sub>3</sub> match well with the interplanar distance of the TiO<sub>2</sub> anatase structure. The calculated lattice spacing of 0.35 nm can be assigned to the interplanar distance of the (011) plane of TiO<sub>2</sub>.

The HAADF-STEM images and the corresponding elemental mappings for all the TiO<sub>2</sub>-coated hematite samples are shown in Fig. 2. Fig. 2a shows small TiO<sub>2</sub> nuclei randomly deposited on hematite nanorods after 10 ALD cycles. As the deposition proceeds, small nuclei grow up, and small nanocrystallites coalesce together to form a very thin, continuous TiO<sub>2</sub> layer. The HAADF-STEM images of TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub>, TiO<sub>2</sub>-40/Fe<sub>2</sub>O<sub>3</sub>, TiO<sub>2</sub>-80/

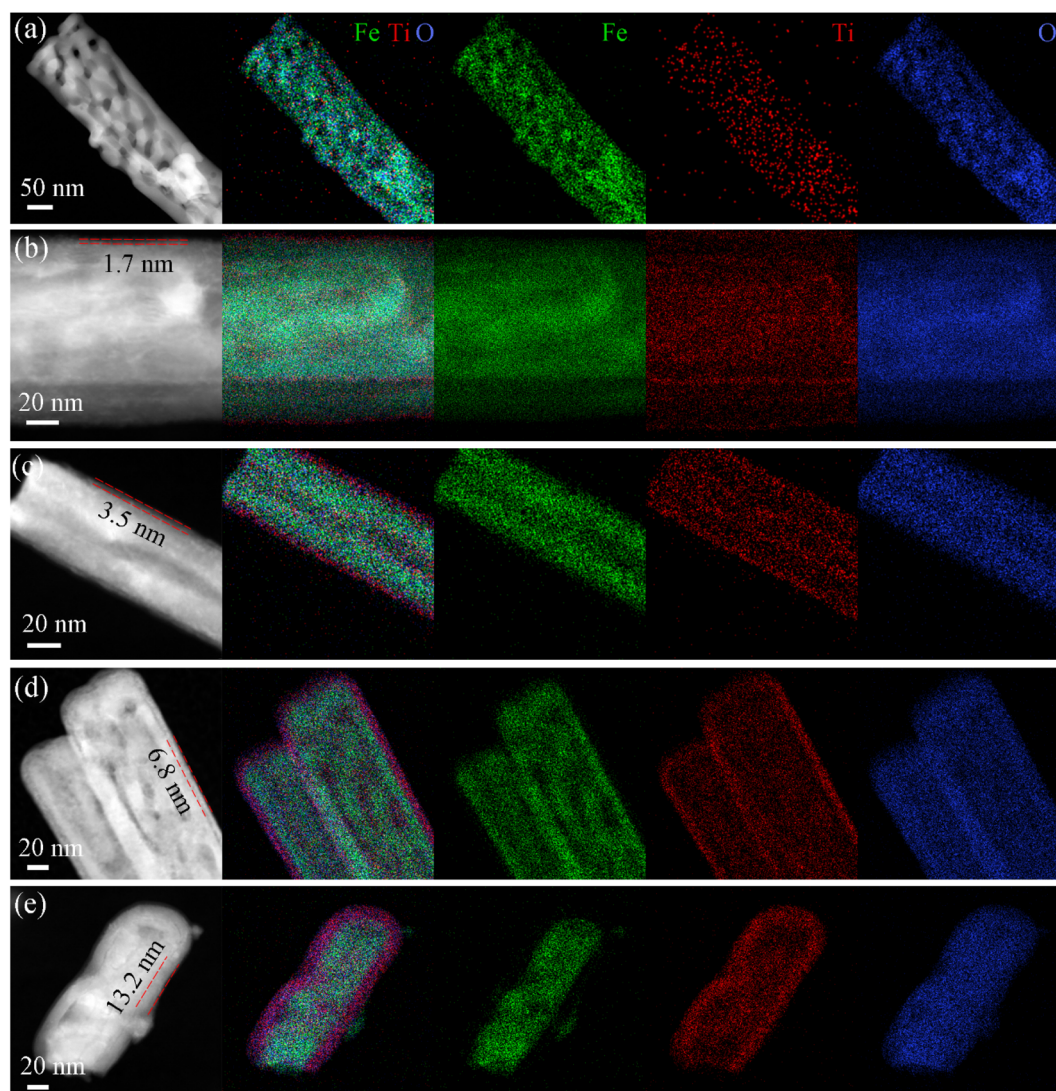


Fig. 2 HAADF-STEM and corresponding EDX elemental maps for (a) TiO<sub>2</sub>-10/Fe<sub>2</sub>O<sub>3</sub>, (b) TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub>, (c) TiO<sub>2</sub>-40/Fe<sub>2</sub>O<sub>3</sub>, (d) TiO<sub>2</sub>-80/Fe<sub>2</sub>O<sub>3</sub> and (e) TiO<sub>2</sub>-150/Fe<sub>2</sub>O<sub>3</sub>.



Fe<sub>2</sub>O<sub>3</sub> and TiO<sub>2</sub>-150/Fe<sub>2</sub>O<sub>3</sub> show a conformal and homogeneous TiO<sub>2</sub> deposition on the Fe<sub>2</sub>O<sub>3</sub> nanorods, where the phases of TiO<sub>2</sub> can be easily identified because of their different Z-contrast. The average thickness of the TiO<sub>2</sub> film is calculated as 1.7, 3.5, 6.8, and 13.2 nm for 20, 40, 80 and 150 ALD cycles, respectively (Fig. 2b–e). This indicates that the shell thickness of the TiO<sub>2</sub> coatings *versus* the number of ALD cycles shows good linearity with a slope corresponding to a growth per cycle (GPC) of 0.88 Å. The thickness of the TiO<sub>2</sub> overlayer from TEM images is well-controlled and shows the same trend as the ellipsometry data on Si wafers (Table S1, ESI†). The corresponding elemental maps further elucidate a homogenous distribution of Ti and O elements on the Fe<sub>2</sub>O<sub>3</sub> nanorods. The EDX spectra for the pristine Fe<sub>2</sub>O<sub>3</sub> and TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub>, further indicates the presence of both Fe and Ti elements (Fig. S4, ESI†). The atomic/weight percentage of all the samples with different ALD cycles are summarized in Table S2 of ESI†.

The optical properties of the pristine TiO<sub>2</sub> on FTO substrate, pristine hematite and TiO<sub>2</sub>-coated hematite were measured by UV-vis-NIR diffuse reflectance spectroscopy (Fig. S5, ESI†). As expected, TiO<sub>2</sub> shows a high absorbance in the UV region, while Fe<sub>2</sub>O<sub>3</sub> and TiO<sub>2</sub>/Fe<sub>2</sub>O<sub>3</sub> show a high absorbance in the visible region (Fig. S5a, ESI†). Fig. S5b–d in the ESI† shows the estimated optical band gap, which can be measured from the plot of  $(F(R) \cdot h\nu)^{1/2}$  *versus*  $h\nu$ .<sup>36</sup> The band gaps obtained from Tauc-plots for pristine TiO<sub>2</sub>, pristine Fe<sub>2</sub>O<sub>3</sub> and TiO<sub>2</sub>/Fe<sub>2</sub>O<sub>3</sub> are around 3.25, 2.07 and 2.02 eV, respectively, that match well with the reported values.<sup>4,10</sup> The slight change of band gap between Fe<sub>2</sub>O<sub>3</sub> and TiO<sub>2</sub>/Fe<sub>2</sub>O<sub>3</sub> indicates that the doping of Ti in hematite is negligible.

The surface composition and chemical states of the samples were investigated by high-resolution XPS (Fig. S6, ESI†). The survey spectra of pristine TiO<sub>2</sub>, pristine hematite and hematite

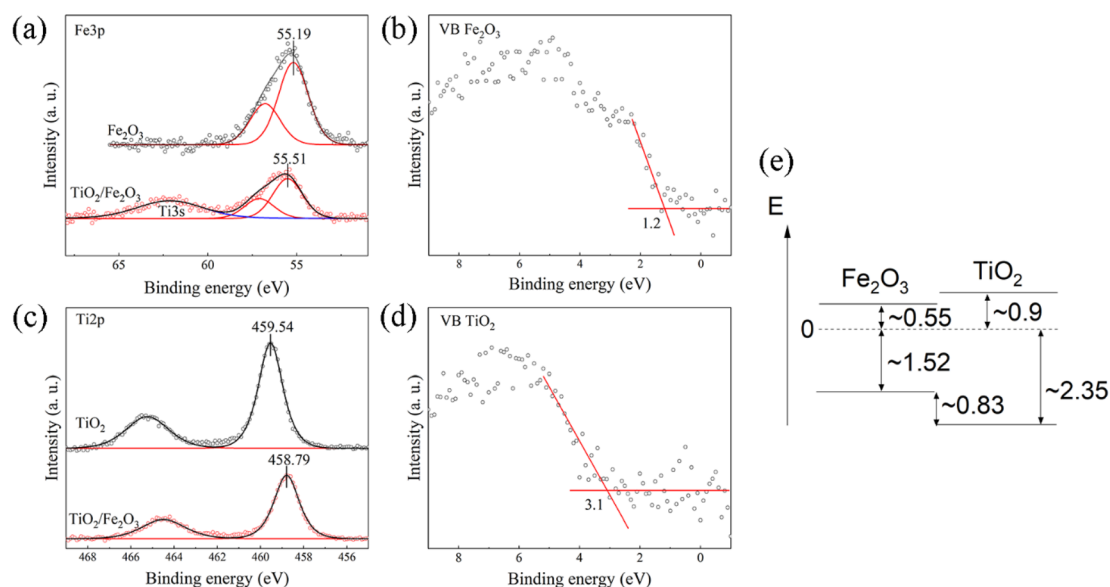
coated with 20 TiO<sub>2</sub> ALD cycles are shown in Fig. S6a, ESI†, confirming the presence of all expected elements. Fig. S6b of ESI† shows the Fe 2p peaks of the pristine Fe<sub>2</sub>O<sub>3</sub> and TiO<sub>2</sub>-coated Fe<sub>2</sub>O<sub>3</sub>. The iron peaks are still well visible for the TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub> sample, indicating that the TiO<sub>2</sub> layer is sufficiently thin, which is in good agreement with TEM measurements. The two peaks located at binding energies of 711.9 and 725.1 eV can be assigned to Fe 2p<sub>3/2</sub> and Fe 2p<sub>1/2</sub> of hematite, respectively, indicating the oxidation state III of iron.<sup>37,38</sup> XPS spectrum of O 1s peak for hematite samples is centered at 530.2 eV, which is consistent with literature values (Fig. S6c, ESI†).<sup>21,39</sup>

To determine the energy level offset at the  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>/TiO<sub>2</sub> heterostructures, the Kraut's method was used, as a direct measurement of the valence band offsets is typically difficult due to overlapping valence band signals.<sup>40,41</sup> Assuming the energy difference between the valence band and the core levels is the same for the pristine materials and the heterostructure, the valence band offset in the heterostructure  $\Delta E_V^{\text{Fe}_2\text{O}_3/\text{TiO}_2}$  can be calculated with the following equation:

$$\Delta E_V^{\text{Fe}_2\text{O}_3/\text{TiO}_2} = \left( E_{\text{Ti } 2p_{3/2}}^{\text{TiO}_2} - E_V^{\text{TiO}_2} \right) - \left( E_{\text{Fe } 3p_{3/2}}^{\text{Fe}_2\text{O}_3} - E_V^{\text{Fe}_2\text{O}_3} \right) - \left( E_{\text{Ti } 2p_{3/2}}^{\text{Fe}_2\text{O}_3/\text{TiO}_2} - E_{\text{Fe } 3p_{3/2}}^{\text{Fe}_2\text{O}_3/\text{TiO}_2} \right) \quad (1)$$

The corresponding core level and valence spectra are shown in Fig. 3a–d.

While the binding energy of the core levels was determined from the peak position after fitting, the valence band onsets were determined from a linear extrapolation of the leading edges. This results in a valence band offset of 0.83 eV, with the valence band of  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> being closer to the Fermi-level. Considering the band gaps determined by absorption spectroscopy (Fig. S5, ESI†), we find that the  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>/TiO<sub>2</sub> interface



**Fig. 3** (a) Fe 3p core level spectra of Fe<sub>2</sub>O<sub>3</sub> (black) and TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub> (red). Shirley backgrounds have been removed. (b) Valence band spectrum of Fe<sub>2</sub>O<sub>3</sub> with a linear extrapolation of the valence band onset. (c) Ti 2p core level spectra of TiO<sub>2</sub> (black) and TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub> (red). Shirley backgrounds have been removed. (d) Valence band spectrum of TiO<sub>2</sub> with a linear extrapolation of the valence band onset. (e) Schematic energy level diagram of TiO<sub>2</sub>/Fe<sub>2</sub>O<sub>3</sub>.



forms a straddling gap (type I heterostructure) with a conduction band offset of 0.35 eV (Fig. 3e).

### 3.2 Photoelectrochemical activities

The PEC performance was measured on pristine hematite and hematite modified with various thicknesses of the TiO<sub>2</sub> overlayer (Fig. 4). Fig. 4a shows the photocurrent–potential ( $J$ – $V$ ) curves of hematite photoanodes under visible light. Compared to pristine hematite, the photoanodes of TiO<sub>2</sub>-coated hematite show an obvious lower overpotential for water oxidation reaction and a cathodic shift of the onset potential. Furthermore, the photocurrent density shows a significant enhancement by deposition of the TiO<sub>2</sub> overlayer.

To further study the impact of the TiO<sub>2</sub> overlayer on hematite photoanodes on the PEC performance, the measurement of the photocurrent response under visible light at 0.5 V bias vs. Hg/HgO (1.4 V vs. RHE) was carried out on all the samples (Fig. 4b). As mentioned above, Fe<sub>2</sub>O<sub>3</sub> can be excited to produce photogenerated electron–hole pairs under visible light due to the low band gap, while TiO<sub>2</sub> shows almost no photocurrent response because of its large band gap. Therefore, the photocurrent response is attributed only to Fe<sub>2</sub>O<sub>3</sub> under visible light. The photocurrent density of pristine hematite is 0.01 mA cm<sup>-2</sup>. The low photocurrent response is due to the high

recombination rate of photogenerated charge carriers. By the deposition of the TiO<sub>2</sub> overlayer, all the composite samples display improved photocurrent densities, which can be attributed to the passivation effect of TiO<sub>2</sub>.<sup>18,42</sup> It is well known that the passivation layer is normally very thin (<100 nm), which can prevent parasitic light absorption and charge-transfer inhibition.<sup>43</sup> TiO<sub>2</sub> passivates the surface states, reduces the recombination rate of the photogenerated electron–hole pairs, facilitates photogenerated hole transfer and improves the photocurrent response.<sup>22</sup> The photocurrent densities of TiO<sub>2</sub> coated hematite with 10, 20, 40, 80 and 150 ALD cycles are 0.06, 0.27, 0.13, 0.06 and 0.02 mA cm<sup>-2</sup>, respectively. The highest photocurrent response is recorded for TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub> (~1.7 nm), showing a photocurrent density *ca.* 27 times higher than that obtained for the pristine hematite. At this optimized thickness, the TiO<sub>2</sub> overlayer passivates surface recombination states of hematite and the photogenerated holes can be transferred from hematite to electrolyte. With a further increase of the TiO<sub>2</sub> thickness, the TiO<sub>2</sub> overlayer can hinder the hole transfer efficiency and block light from entering the inner hematite layer, leading to a dramatic decrease in photocurrent response.

During visible light illumination, photogenerated holes in pristine hematite travel to the semiconductor/liquid junction (SCLJ) and accumulate there, owing to the existence of surface

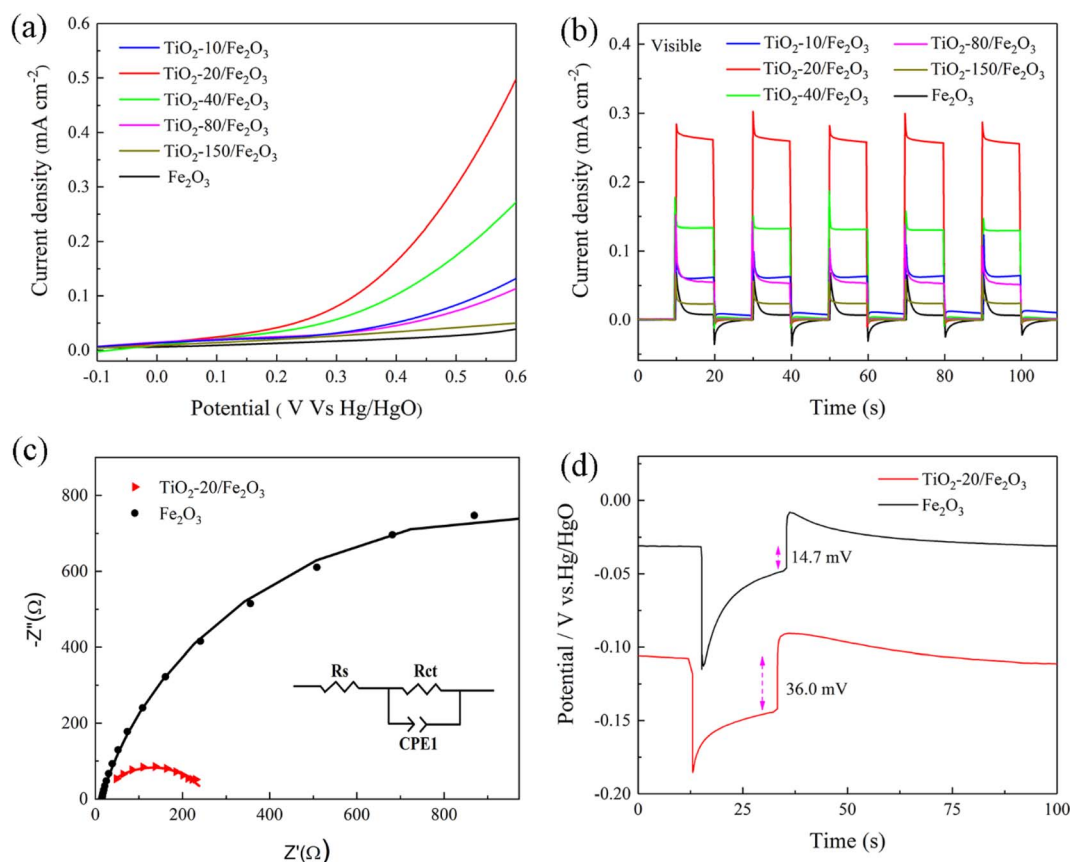


Fig. 4 (a) Photocurrent versus potential plots of a different number of ALD cycles for TiO<sub>2</sub> coated hematite photoanodes under visible light irradiation. (b) Time-based photocurrent density at 0.5 V bias vs. Hg/HgO with the light off and on for 10 s. (c) Nyquist plots of pristine Fe<sub>2</sub>O<sub>3</sub> and TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub> photoanodes measured under light illumination. (d) OCV curves of pristine Fe<sub>2</sub>O<sub>3</sub> and TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub> photoanodes with the light on for 20 s.





trap states. This leads to the appearance of large anodic transient peaks, which decay until they finally reach an equilibrium between charge carrier transfer and recombination at the interface.<sup>44</sup> On the other hand, cathodic transient peaks overshoot and decay back to zero under dark conditions, which are attributed to the electrons diffusing from the external circuit and recombining with holes accumulated at the SCLJ.<sup>12,22</sup> However, the transient photocurrent spike decreases obviously after the deposition of TiO<sub>2</sub>. Moreover, the feature was highly repeatable for numerous ALD cycles. This result further demonstrates the passivation effect of the TiO<sub>2</sub> overlayer. Therefore, Fe<sub>2</sub>O<sub>3</sub> essentially acts as a light absorption layer, absorbing the visible light and producing photogenerated electron–hole pairs.

EIS was measured under light illumination to study the interfacial kinetics of the charge transfer process. Fig. 4c displays the Nyquist plots of pristine Fe<sub>2</sub>O<sub>3</sub> and TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub> that are simulated using the equivalent circuit (EC) shown in Fig. 4c inset. The equivalent circuit elements include the series resistance in the PEC cell ( $R_s$ ), the charge transfer resistance ( $R_{ct}$ ) at the Fe<sub>2</sub>O<sub>3</sub>/electrolyte interface and a constant phase element (CPE1).<sup>39,45</sup> The fitting parameters for the equivalent circuit elements are shown in Table 1. Compared to pristine hematite, TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub> has a lower charge transfer resistance than pristine hematite, which indicates that the charge transfer across hematite/electrolyte interface is significantly facilitated due to the passivation effect of TiO<sub>2</sub>. The EIS result confirms that TiO<sub>2</sub>/Fe<sub>2</sub>O<sub>3</sub> heterostructure can significantly enhance the charge transfer and separation and improve the PEC water splitting performance.

The open circuit voltage (OCV) of pristine hematite and hematite coated with 20 TiO<sub>2</sub> ALD cycles were measured in the dark and illumination conditions to study the behavior of the photogenerated carriers (Fig. 4d). Under illumination, the value of pristine Fe<sub>2</sub>O<sub>3</sub> moves to a negative potential, while the potential of TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub> shifts more negative than pristine Fe<sub>2</sub>O<sub>3</sub>. This suggests that both samples are n-type semiconductors, in line with the XPS results.<sup>46–48</sup> Under illumination, both samples exhibit a similar trend and create a similar number of photogenerated electrons, which can accumulate on the surface of photoanodes, leading to a potential negative shift and a sudden increase in the OCV. Meanwhile, the recombination of photogenerated electrons and holes occurs, which leads to a decrease of the OCV until a relatively steady state is reached. It can be seen that the potential change for Fe<sub>2</sub>O<sub>3</sub> and TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub> is comparable. However, compared to pristine Fe<sub>2</sub>O<sub>3</sub>, the photoanode of TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub> is more stable due to the passivation effect of TiO<sub>2</sub> and the low photogenerated carrier recombination dynamics. After irradiation for 20

seconds, the light is turned off, and the OCV continues to decline due to the high recombination of photoelectron–hole pairs. The photovoltages of the pristine Fe<sub>2</sub>O<sub>3</sub> and TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub> photoanodes are 14.7 mV and 36.0 mV, respectively, as obtained from the difference between the OCV under dark and light conditions. Compared to pristine Fe<sub>2</sub>O<sub>3</sub>, the photovoltage obtained with TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub> is much higher due to the fact that more photogenerated electrons are collected and a more efficient separation of the photogenerated electron–hole pairs.

To study the photo-response properties further, the incident photon to converted electron (IPCE) and stability test of pristine and TiO<sub>2</sub>-coated hematite photoanodes were measured under visible light. The IPCE values have been calculated according to the equation:

$$\text{IPCE (\%)} = \frac{1240 \times J_{\text{photo}}}{\lambda \times P_{\text{light}}} \times 100 \quad (2)$$

where  $J_{\text{photo}}$  is the photocurrent density (mA cm<sup>-2</sup>) under visible light,  $\lambda$  is the wavelength of the incident light (nm), and  $P_{\text{light}}$  is the light power density (mW cm<sup>-2</sup>).<sup>21,49</sup>

A maximum IPCE of 1.35% at 430 nm is achieved by TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub>, which is much higher than the corresponding value of 0.05% for pristine Fe<sub>2</sub>O<sub>3</sub> photoanodes (Fig. S7a, ESI†). The significant improvement of photoconversion efficiency indicates higher utilization efficiency of photoinduced carriers for the TiO<sub>2</sub>-coated hematite photoanodes, which is consistent with the  $J$ - $V$  measurements (Fig. 4a).

Applied bias photon-to-current efficiency (ABPE) was calculated from the photocurrent–potential plots according to the following equation:

$$\text{ABPE} = \frac{I \times (1.23 - V_{\text{bias}})}{P_{\text{light}}} \times 100\% \quad (3)$$

where  $I$  (mA cm<sup>-2</sup>) is the measured photocurrent density,  $V_{\text{bias}}$  ( $V$  vs. RHE) is the applied potential, and  $P_{\text{light}}$  is the incident illumination power density.

The ABPE was plotted as a function of the applied potential. The TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub> photoanode achieves the maximum ABPE value of 0.01% at 1.02 V, which is significantly higher than that of 0.004% at 0.98 V for pristine hematite electrode (Fig. S7b, ESI†). Clearly, the improvement of PEC water splitting performance can be attributed to the enhanced charge separation caused by the passivation effect of TiO<sub>2</sub>.

Fig. S8 in ESI† shows the current retention as a function of the time of pristine hematite and hematite coated with 20 TiO<sub>2</sub> ALD cycles under continuously visible light irradiation for 60 min. A significant improvement of the photocurrents of pristine hematite electrode upon TiO<sub>2</sub> deposition as well as a good stability can be obtained. The morphology and the composition of the synthesized photocatalyst after stability test are well maintained, which are shown in Fig. S9 and S10 in ESI.†

Fig. S11 in ESI† shows the Mott–Schottky plots of the pristine Fe<sub>2</sub>O<sub>3</sub> and TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub>. The slopes determined from the Mott–Schottky plots are used to estimate the carrier densities. The positive slopes display the n-type semiconductor nature of the pristine Fe<sub>2</sub>O<sub>3</sub> and TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub>. Moreover, the plot of TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub> depicts a suppressed slope compared to pristine

**Table 1** Fitted parameters for the Nyquist plots of pristine Fe<sub>2</sub>O<sub>3</sub> and TiO<sub>2</sub>-20/Fe<sub>2</sub>O<sub>3</sub>

	$R_s$ ( $\Omega$ cm <sup>2</sup> )	$R_{ct}$ ( $\Omega$ cm <sup>2</sup> )
Pristine-Fe <sub>2</sub> O <sub>3</sub>	7.5	247.3
TiO <sub>2</sub> -20/Fe <sub>2</sub> O <sub>3</sub>	13.9	1807.0



$\text{Fe}_2\text{O}_3$ , suggesting a higher carrier concentration in the heterostructure than in the pristine  $\text{Fe}_2\text{O}_3$ . The as-calculated carrier densities for pristine  $\text{Fe}_2\text{O}_3$  and  $\text{TiO}_2$ -20/ $\text{Fe}_2\text{O}_3$  are  $1.6 \times 10^{19} \text{ cm}^{-3}$  and  $4.4 \times 10^{19} \text{ cm}^{-3}$ , respectively.

Contrary to the visible light system, both  $\text{Fe}_2\text{O}_3$  and  $\text{TiO}_2$  participate in the photocurrent response under UV light irradiation due to the suitable band gap of  $\text{Fe}_2\text{O}_3$  and  $\text{TiO}_2$ . As a result, the photocurrent response is not only attributed to  $\text{Fe}_2\text{O}_3$  but is also related to the  $\text{TiO}_2$  overlayer. To verify the role of  $\text{TiO}_2$  overlayer in the PEC performance of hematite, the photocurrent response for hematite and  $\text{TiO}_2$ -coated hematite samples were evaluated under UV light (Fig. S12, ESI†).

The experimental results show that the photocurrent response varies significantly with the thickness of the  $\text{TiO}_2$  overlayer. Compared to pristine hematite, the  $\text{TiO}_2$  overlayer results in a remarkable enhancement of photocurrent response. Before 80 ALD cycles, the variation trend of the photocurrent response of hematite samples with  $\text{TiO}_2$  thickness under UV irradiation was similar to that under visible light irradiation. The photocurrent densities of pristine hematite and  $\text{TiO}_2$  coated hematite with 10, 20, 40 and 80 ALD cycles are 0.03, 0.07, 0.41, 0.22 and 0.09  $\text{mA cm}^{-2}$ , respectively. The optimized  $\text{TiO}_2$ -20/ $\text{Fe}_2\text{O}_3$  (ca. 1.7 nm) shows a photocurrent response of 0.41  $\text{mA cm}^{-2}$ , which is 14 times higher than the photocurrent response of the pristine hematite.  $\text{TiO}_2$  thicknesses up to 1.7–6.8 nm (20 ALD cycles to 80 ALD cycles) lead to a decrease in photocurrent response. At this stage, the  $\text{TiO}_2$  overlayer mainly acts as a passivation layer and  $\text{Fe}_2\text{O}_3$  acts as a light absorption layer. However, the photocurrent response increases with a further increase of the thickness of the  $\text{TiO}_2$  overlayer (*i.e.*, from 6.8 to 13.2 nm, 80 ALD to 150 ALD), in contrast to what was observed under visible light. The photocurrent density of  $\text{TiO}_2$ -150/ $\text{Fe}_2\text{O}_3$  is 0.14  $\text{mA cm}^{-2}$ . This change of photocurrent response at a higher thickness suggests that the photocurrent response is not dominated by  $\text{Fe}_2\text{O}_3$  anymore, but is restricted to the  $\text{TiO}_2$  overlayer.

Fig. 5 compares the photocurrent response of hematite samples coated with different thicknesses of the  $\text{TiO}_2$  overlayer under different illumination. The experimental results reveal a substantial variation of the photocurrent response as a function of the thickness of the  $\text{TiO}_2$  overlayer. The highest photocurrent response exhibited by the  $\text{TiO}_2$ -20/ $\text{Fe}_2\text{O}_3$  photoanodes is recorded under both visible and UV light. After 20 ALD cycles, the photocurrent response decreases with a further increase of the  $\text{TiO}_2$  thickness under visible light, while it first decreases and then increases under UV light.

Under visible light, the  $\text{TiO}_2$  overlayer acts as a passivation layer. The photogenerated holes from  $\text{Fe}_2\text{O}_3$  can cross the passivation layer primarily by tunneling, which gradually vanishes within a few nanometers.<sup>43,50,51</sup> Because of this mechanism, the  $\text{TiO}_2$  overlayer must be of the optimal thickness. It should not be too thin to passivate the surface recombination states enough, and not too thick to allow the photogenerated holes transfer between  $\text{Fe}_2\text{O}_3$  and electrolyte to participate in the water oxidation. Under UV light, when the  $\text{TiO}_2$  thickness is  $\leq 6.8$  nm, it does not absorb enough UV light, and therefore only few photogenerated charge pairs are produced in the  $\text{TiO}_2$ . In

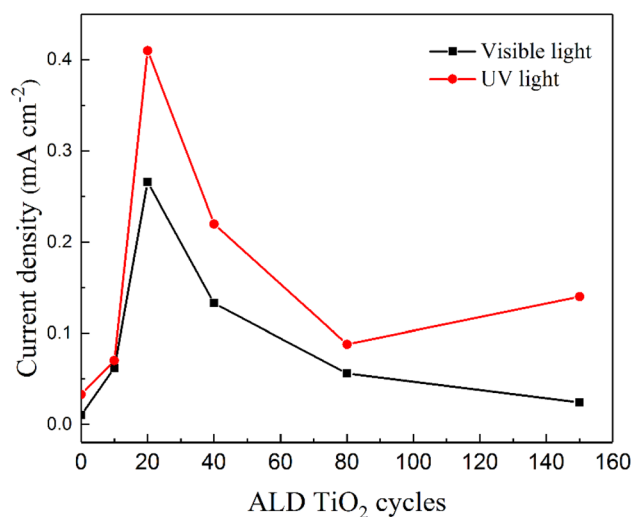


Fig. 5 Current densities of hematite photoanodes with different ALD  $\text{TiO}_2$  cycles at 0.5 V bias vs.  $\text{Hg}/\text{HgO}$ , 1 M  $\text{KOH}$ , under visible light and UV light irradiation, respectively.

this case,  $\text{TiO}_2$  mainly acts as a passivation layer, and the photocurrent response is dominated by  $\text{Fe}_2\text{O}_3$ . Therefore, the photocurrent response under UV light follows the same trend as under visible light. However, when the thickness of  $\text{TiO}_2$  is larger than the optimal 1.7 nm (20 ALD cycles), it hinders the transfer of the photogenerated holes. Therefore, the photogenerated holes in  $\text{Fe}_2\text{O}_3$  cannot efficiently cross the  $\text{TiO}_2$  overlayer. Conversely, when the  $\text{TiO}_2$  overlayer is thick enough, it absorbs more UV light and produces more photogenerated charge pairs that participate in water oxidation. The contribution of the  $\text{TiO}_2$  overlayer for photocurrent response becomes dominant at 150 ALD cycles (13.2 nm). At this stage,  $\text{TiO}_2$  mainly behaves as the active layer, and the contribution of  $\text{Fe}_2\text{O}_3$  becomes negligible. This demonstrates that the photoresponse under UV illumination is largely related to the thickness of the  $\text{TiO}_2$  overlayer.

### 3.3 Charge transfer mechanism

Fig. 6 displays the schematic diagram of the charge separation and transport mechanism under visible and UV light illumination. Under visible light,  $\text{Fe}_2\text{O}_3$  with a suitable band gap can produce photogenerated electrons in the conduction band and photogenerated holes in the valence band. The photogenerated electrons driven by the bias voltage will transfer through the external circuit (Fig. 6a). The photogenerated holes from  $\text{Fe}_2\text{O}_3$  cannot be transferred to the valence band of  $\text{TiO}_2$  owing to an unfavorable valence band edge alignment of  $\text{Fe}_2\text{O}_3$  and  $\text{TiO}_2$ . Therefore,  $\text{TiO}_2$  behaves as a passivation layer and  $\text{Fe}_2\text{O}_3$  acts as a light absorption layer. In the case of UV light illumination, both  $\text{Fe}_2\text{O}_3$  and  $\text{TiO}_2$  have suitable band gaps that can be excited to produce electron-hole pairs. The photogenerated electrons of  $\text{TiO}_2$  migrate to the conduction band of  $\text{Fe}_2\text{O}_3$ , which lies at a more positive potential than that of  $\text{TiO}_2$ . Subsequently, the electrons in  $\text{Fe}_2\text{O}_3$  flow into the external circuit. The photoinduced holes in  $\text{Fe}_2\text{O}_3$  can pass through the  $\text{TiO}_2$  thin layer by





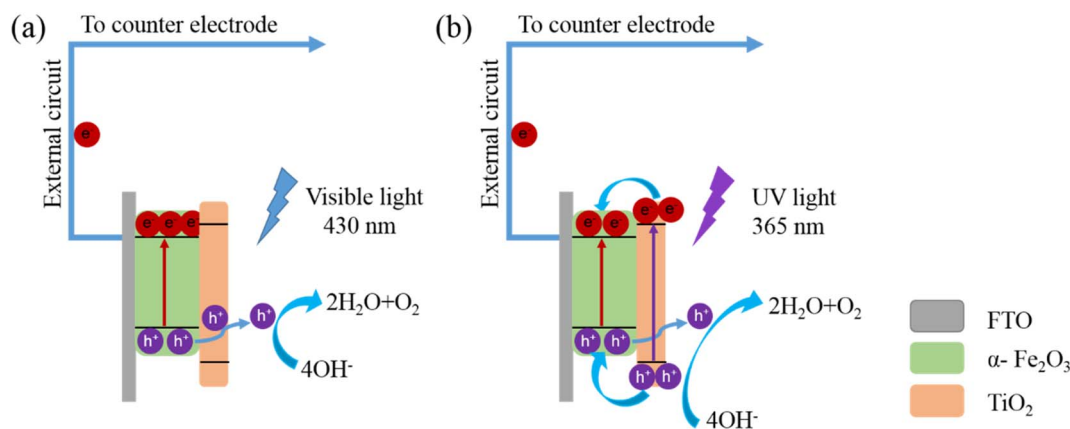


Fig. 6 Schematic illustration of the charge-transfer pathways in the  $\text{TiO}_2/\text{Fe}_2\text{O}_3$  system under visible light and UV light.

tunneling effect for the oxidation reaction. The photogenerated holes from  $\text{TiO}_2$  can react directly with the electrolyte. It can also transfer to the valence band of  $\text{Fe}_2\text{O}_3$  and then cross the  $\text{TiO}_2$  thin layer to participate in the water oxidation reaction (Fig. 6b). This indicates that the photocurrent response is not always governed by the  $\text{Fe}_2\text{O}_3$ . The  $\text{TiO}_2$  overlayer also plays an essential role in producing photocurrent response under UV light. As we have discussed above, the impact of the  $\text{TiO}_2$  overlayer is largely dependent on the film thickness.

## 4. Conclusion

$\text{TiO}_2/\text{Fe}_2\text{O}_3$  heterostructures have been synthesized by a hydrothermal method and atomic layer deposition. The as-synthesized samples show a homogeneous and conformal coating of  $\text{TiO}_2$  onto the hematite nanorods. The morphology and thickness of a  $\text{TiO}_2$  overlayer can be precisely controlled. The pristine hematite and hematite coated with various thicknesses of  $\text{TiO}_2$  are studied for their photocurrent response under visible and UV light. Compared to pristine hematite, the  $\text{TiO}_2$  coating leads to a significant improvement in photocurrent response. The photocurrent response of the  $\text{TiO}_2/\text{Fe}_2\text{O}_3$  heterostructures is strictly related to the thickness of the  $\text{TiO}_2$  overlayer. The best photoanode response is observed with a coating of 1.7 nm ( $\text{TiO}_2\text{-20}/\text{Fe}_2\text{O}_3$ ) under both visible and UV light. Our results show that the photocurrent density of  $\text{TiO}_2\text{-20}/\text{Fe}_2\text{O}_3$  reaches  $0.27 \text{ mA cm}^{-2}$  at  $0.5 \text{ V vs. Hg/HgO}$  under visible light, which is 27 times higher than that of pristine  $\alpha\text{-Fe}_2\text{O}_3$ . Under visible light, the  $\text{TiO}_2$  overlayer acts as a passivation layer facilitating the photogenerated holes transfer and reducing the surface recombination rate of electron-hole pairs.  $\text{Fe}_2\text{O}_3$  is the major contributor to the photogeneration of electron-hole pairs. Under UV light, the  $\text{TiO}_2$  overlayer behaves mainly as a passivation layer at lower thicknesses ( $\leq 6.8 \text{ nm}$ ). At higher thickness,  $\text{TiO}_2$  acts as a light absorption layer, while the contribution of  $\text{Fe}_2\text{O}_3$  is reduced. This work not only provides deep insights into the impact of a  $\text{TiO}_2$  overlayer on the PEC water splitting performance of hematite, but can be also valuable and promising for the design and development of high-performance photoanodes.

## Data availability

The data supporting this article have been included as part of the ESI.†

## Conflicts of interest

There are no conflicts to declare.

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