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$[(C_5H_6N_2)_2H](Sb_4F_{13})$: a polyfluoroantimonite with a strong second harmonic generation effect†

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It is of great difficulty to create a new antimonite with second-harmonic-generation (SHG) intensity larger than 6 times that of KDP. In this study, a polyfluoroantimonite strategy has been proposed to explore fluoroantimonites with large nonlinear optical (NLO) coefficients. Under the cooperation of chemical (highly asymmetric π -conjugated organic amine) and physical (viscous reaction medium ethylene glycol) methods, two novel polyfluoroantimonites, namely, $(3PC)_2(Sb_4F_{14})$ and $(3AP)_2(Sb_4F_{13})$, have been achieved. Interestingly, these two structures contain two new polyfluoroantimonite groups respectively, an isolated $(Sb_4F_{14})^{2-}$ four-member polyhedral ring and an infinite $[Sb_4F_{13}]_{\infty}^{-}$ helical chain. More importantly, the polar $(3AP)_2(Sb_4F_{13})$ displays a strong SHG intensity of $8.1 \times$ KDP, a large birefringence of 0.258@546 nm and a high laser-induced damage threshold (LIDT) value of 149.7 MW cm^{-2} . Theoretical calculations indicated that its strong SHG effect stems from the synergistic effect of the helical $[Sb_4F_{13}]_{\infty}^{-}$ polyfluoroantimonite chain and π -conjugated $3AP^+$ cation, with a contribution ratio of 48.93% and 50.77% respectively. This work provides a new approach for the design and synthesis of high-performance fluoroantimonites.

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Introduction

Nonlinear optical crystals capable of widening the wavelength range of output radiations have attracted widespread attention. SHG crystals, as the most representative NLO materials, have been applied in fields such as optical communication, laser technology, biomedicine, photonics, and sensor technology.^{1–9} Fluoroantimonites are considered as a potential candidate for excellent SHG materials. On the one hand, the Sb^{3+} cation with stereochemically active-lone-pair (SCALP) electrons can undergo second-order Jahn–Teller (SOJT) distortion, resulting in high polarizability.¹⁰ On the other hand, fluorine, being the most electronegative, can enlarge the band gaps of target compounds due to the deeper F-2p orbit than O-2p.^{11–14} Recently, scientists have developed various design strategies to create new SHG fluoroantimonites, mainly including the following: (a) introduction of tetrahedral groups (SO_4^{2-} and PO_4^{3-}), which can widen the band gaps of fluoroantimonites,

such as $CsSbF_2SO_4$ ($3.0 \times$ KDP, 4.76 eV) and $Rb_2Sb(P_2O_7)F$ ($5.1 \times$ KDP, 4.74 eV);^{15,16} (b) combination of Sb^{3+} with delocalized π -conjugated groups such as oxalate and nitrate as in $RbSb_2(C_2O_4)F_5$ ($1.3 \times$ KDP) and $Rb_2SbF_3(NO_3)_2$ ($2.7 \times$ KDP);^{17,18} (c) introduction of SOJT distorted d^0 transition metal cations (Mo^{6+} , W^{6+} , Nb^{5+} , etc.) such as $Cs_6Sb_4Mo_3O_5F_{26}$ ($0.7 \times$ KDP).¹⁹ However, the development of new fluoroantimonites with high SHG intensity is still a challenge. The SHG intensities of most reported fluoroantimonites are less than $6 \times$ KDP. Consequently, it is necessary to develop a new strategy to enhance the SHG effects of fluoroantimonites further.

It is believed that the condensation of functional building blocks can enhance the density of SHG active groups, thereby impacting the structure and SHG capability of target compound.^{20–25} This approach has facilitated the development of numerous high-performance polyiodates, such as GdI_5O_{14} ($15 \times$ KDP) and $[o-C_5H_4NHOH]_2[I_7O_{18}(OH)] \cdot 3H_2O$ ($8.5 \times$ KDP).^{26,27} If fluoroantimonate groups can be condensed into polyfluoroantimonite groups, it is expected to obtain new NLO materials with strong SHG effects in polyfluoroantimonite system. However, research on polyfluoroantimonites is very rare, especially for the NLO material, which can be attributed to the following two reasons: (i) the synthetic conditions for polyfluoroantimonites are rather complicated and still unclear. (ii) The bond energy of the Sb–F bond is lower than that of the Sb–O bond, which makes the polymerization of fluoroantimonite groups difficult compared to oxysalts.^{28,29}

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In order to explore new fluoroantimonites with strong SHG effects, we have dedicated our efforts to polyfluoroantimonite system, and found a universal synthesis method for polyfluoroantimonites, which involves using organic amine groups (chemical method) and enhancing the viscosity of the reaction media (physical method). Highly asymmetric organic amine groups may not only facilitate the polymerization of fluoroantimonite groups during their protonation process, but also trigger the formation of noncentrosymmetric (NCS) structures. Furthermore, the viscous reaction medium can reduce the migration rate of fluoroantimonite groups, which can also promote the combination of fluoroantimonate groups to form new polyfluoroantimonite anions. Two organic amine cations, namely, 3-pyridinecarboxamide [$3PC^+$, $(C_6H_7N_2O)^+$] and 3-aminopyridine [$3AP^+$, $(C_5H_7N_2)^+$], were identified as the highly asymmetric NLO active groups through theoretical calculations.^{30–32} The $3PC^+$ and $3AP^+$ cations feature significant polarizability anisotropy (δ) and hyperpolarizability (β), surpassing most of the inorganic or organic π -conjugated groups, such as $B_3O_6^{3-}$, $(C_3N_3O_3)^{3-}$, $(C_5H_6NO)^+$, and $(C_4H_6N_3)^+$ (Fig. 1). The selection of viscous solvents is also very important. For instance, phosphoric acid exhibits a propensity to selectively bind with organic ligands, leading to the formation of by-products during the reaction process, and highly viscous glycerol is not conducive to crystallization.^{33–35} Finally, ethylene glycol, a viscous and neutral organic solvent, was selected as the reaction medium.

Guided by the above method, an appropriate amount of hydrofluoric acid was added dropwise into a mixture of organic amines (3PC and 3AP) and SbF_3 to protonate the amines, and ethylene glycol was used to enhance the viscosity of the reaction medium. Two novel organic–inorganic hybrid polyfluoroantimonites, namely, $(3PC)_2(Sb_4F_{14})$ and $(3AP)_2(Sb_4F_{13})$,

have been successfully synthesized by a facile solvent evaporation method. $(3AP)_2(Sb_4F_{13})$ crystallizes in a NCS space group because the $3AP^+$ cation features a larger hyperpolarizability than the $3PC^+$ cation. Interestingly, $(3AP)_2(Sb_4F_{13})$ displays a unique $[Sb_4F_{13}]_\infty^-$ helical chain, a strong SHG effect ($8.1 \times$ KDP) under 1064 nm laser radiation, a large birefringence ($0.258 @ 546$ nm) and a high LIDT value (149.7 MW cm^{-2}). Herein, we report the syntheses, structural analyses, optical properties and first-principles studies of $(3PC)_2(Sb_4F_{14})$ and $(3AP)_2(Sb_4F_{13})$.

Results and discussion

Single crystals of $(3PC)_2(Sb_4F_{14})$ and $(3AP)_2(Sb_4F_{13})$ were synthesized using a facile solvent evaporation method (Fig. S1†). The detailed crystallographic data for both structures are given in Table S1.† The atomic coordinates and bond lengths are listed in Tables S2 and S3.† The presence of Sb, F, C and N elements was confirmed by field-emission scanning electron microscopy (FESEM) analyses (Fig. S2†). Powder XRD analyses confirmed the purity of the products (Fig. S3†).

$(3PC)_2(Sb_4F_{14})$ crystallizes in the triclinic crystal system with the CS space group $P\bar{1}$ (No. 2). There are one $(Sb_2F_7)^-$ anion and one $3PC^+$ cation in its asymmetric unit (Table S2†). Sb(1) and Sb(2) are coordinated to four and five fluorine atoms, respectively, forming the SbF_4 seesaw and SbF_5 square pyramid, with lone-pair electrons located in the open sites. The Sb–F bond lengths range from 1.927(3) to 2.521(3) Å (Table S3†). Two Sb(1) F_4 and two Sb(2) F_5 groups further linked to a $(Sb_4F_{14})^{2-}$ four-member polyhedral ring (4-MPR) *via* corner-sharing of F(1) and F(4) atoms (Fig. 2a). The bond-valence-sum results show that the bond valences of Sb(1) and Sb(2) are 3.093 and 3.181, respectively, suggesting their oxidation states of +3. As shown in Fig. 2b, the $(Sb_4F_{14})^{2-}$ anionic clusters, separated by the $3PC^+$ cations, were stacking linearly along the *c*-axis to form the structure of $(3PC)_2(Sb_4F_{14})$. The $3PC^+$ cations were connected with $(Sb_4F_{14})^{2-}$ anion clusters through hydrogen bonds with hydrogen bond lengths in the range of 3.108(6)–3.229(5) Å (Table S4†). The adjacent $3PC^+$ cationic pyrimidine rings are parallel to each other, with an inter-ring distance of 3.323 Å (Fig. S4a†).

When the $3PC^+$ cation was replaced by the $3AP^+$ cation, a NCS polyfluoroantimonite, namely, $(3AP)_2(Sb_4F_{13})$, was obtained. $(3AP)_2(Sb_4F_{13})$ crystallizes in a chiral triclinic space group $P1$ (No. 1) and its asymmetric unit is composed of one $3AP^+$ cation, one 3AP molecule and one $(Sb_4F_{13})^-$ anion (Table S2†). Sb(1), Sb(3) and Sb(4) adopt the SbF_4 seesaw configuration, whereas Sb(2) is in the SbF_5 square pyramid configuration. The Sb–F bond distances are in the range of 1.915(5)–2.522(6) Å (Table S3†). The calculated total bond valences for Sb(1), Sb(2), Sb(3) and Sb(4) are 3.045, 3.104, 2.857 and 2.818, respectively, indicating that their oxidation states are all +3. Its structure features a unique $[Sb_4F_{13}]_\infty^-$ helical polyfluoroantimonite chain along the *b* axis with $3AP^+$ cations and 3AP molecules located at the inter-chain space to keep charge balance (Fig. 3c). Sb(1) F_4 , Sb(2) F_5 , Sb(3) F_4 and Sb(4) F_4 groups are interconnected by corner-sharing [F(4), F(8) and F(11)] to form a $(Sb_4F_{14})^{2-}$ tetramer (Fig. 3a).

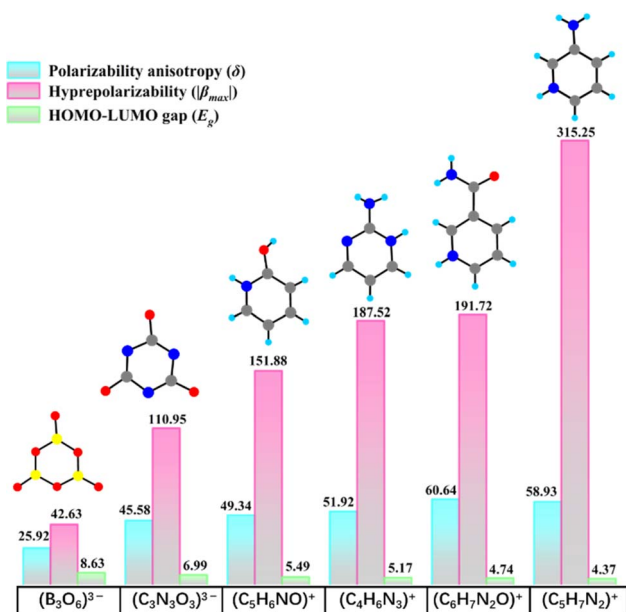


Fig. 1 Geometry and key NLO parameters of the conventional functional groups, the $3PC^+$ cation and the $3AP^+$ cation.



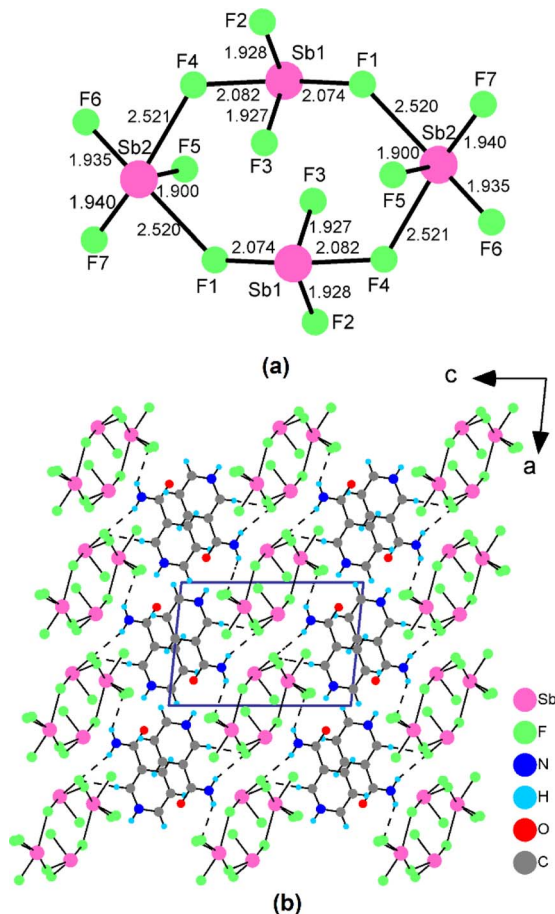


Fig. 2 The configuration of the $(\text{Sb}_4\text{F}_{14})^{2-}$ 4-MPR and the Sb–F bond lengths (Å) (a) and the crystal structure of $(3\text{PC})_2(\text{Sb}_4\text{F}_{14})$ (b).

Neighbouring $(\text{Sb}_4\text{F}_{14})^{2-}$ tetramers are further interconnected into a unique $[\text{Sb}_4\text{F}_{13}]_{\infty}^{-}$ anionic chain along the *b* axis *via* Sb(1)–F(10)–Sb(4) bridges (Fig. 3b). The 3AP^+ cations are linked with the $[\text{Sb}_4\text{F}_{13}]_{\infty}^{-}$ helical polyfluoroantimonite chains through hydrogen bonds with the hydrogen bond lengths in the range of 2.697(12)–3.080(15) Å (Table S4†). There is also π – π interaction between neighbouring pyridine rings with a dihedral angle of 2.257° and inter-ring distances ranging from 3.837 to 4.001 Å (Fig. S4b†). The dipole moment of the Sb_4F_{14} tetramer in $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ was calculated. The local dipole moments have the values of 8.988–19.160 *D* for SbF_4 units, and 8.719 *D* for SbF_5 units. The local dipole moment of the Sb_4F_{14} tetramer is 18.558 *D* (Table S5†), implying its polar nature.

Benefitting from the viscous reaction media and the organic amine groups, both $(3\text{PC})_2(\text{Sb}_4\text{F}_{14})$ and $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ exhibit novel polyfluoroantimonite structures formed by the condensation of the SbF_4 seesaw and SbF_5 square pyramid (Fig. S5†). These two types of polyfluoroantimonite groups have not been reported previously. For $(3\text{PC})_2(\text{Sb}_4\text{F}_{14})$ and $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$, the Sb–F bonds of the Sb–F–Sb bridges are 2.074(3)–2.521(3) and 1.940(6)–2.522(6) Å, respectively, which are longer than their terminal Sb–F bonds [1.900(3)–1.940(3) and 1.915(5)–1.964(6) Å]. These values are comparable with those of polyfluoroantimonite crystals reported before.²⁹

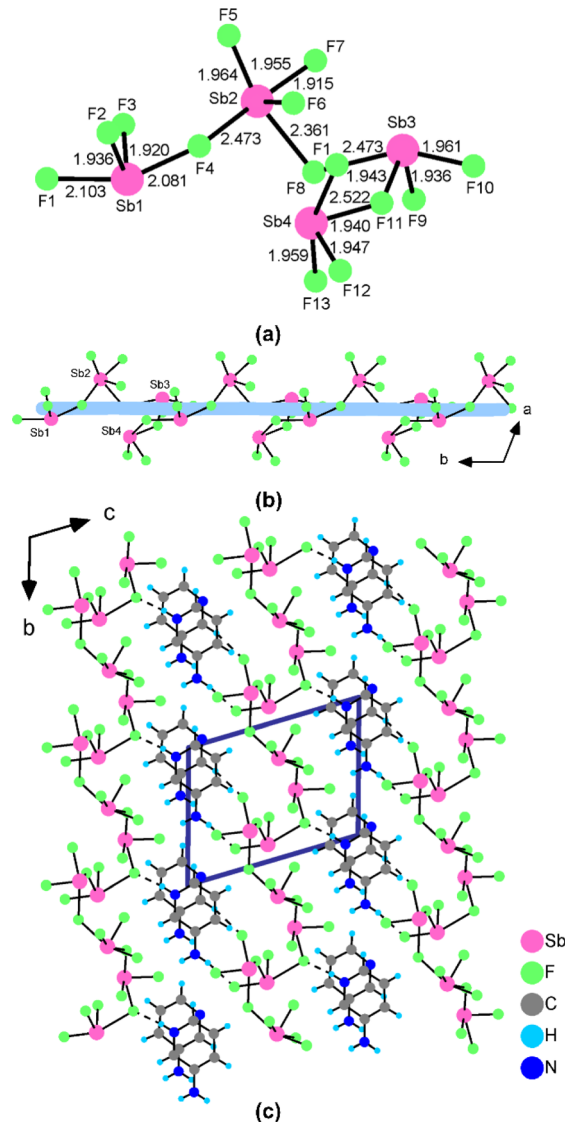


Fig. 3 The configuration of the $(\text{Sb}_4\text{F}_{14})^{2-}$ tetramer and the Sb–F bond lengths (Å) (a), the $[\text{Sb}_4\text{F}_{13}]_{\infty}^{-}$ helical polyfluoroantimonite chain (b) and the crystal structure of $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ (c).

The crystal structures of $(3\text{PC})_2(\text{Sb}_4\text{F}_{14})$ and $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ undergo a transition from CS to NCS arrangements. To comprehend the reasons behind this transition, a detailed structural comparison has been conducted between the two compounds. Specifically, the basic building units in these two compounds are similar, including the condensed fluoroantimonite groups and highly asymmetric π -conjugated organic amine cations. The difference lies in the arrangement of organic amine cations and the polymerization modes of fluoroantimonite groups. In comparison to the antiparallel arrangement of 3PC^+ cations, the 3AP^+ cations and 3AP molecules exhibit a parallel pairing pattern (Fig. 4a and b), favouring the formation of NCS structures. It is worth noting that this parallel pairing pattern is exceptionally rare in organic amines. Furthermore, there have been significant changes in the polymerization modes of fluoroantimonite groups. The $(\text{Sb}_4\text{F}_{14})^{2-}$

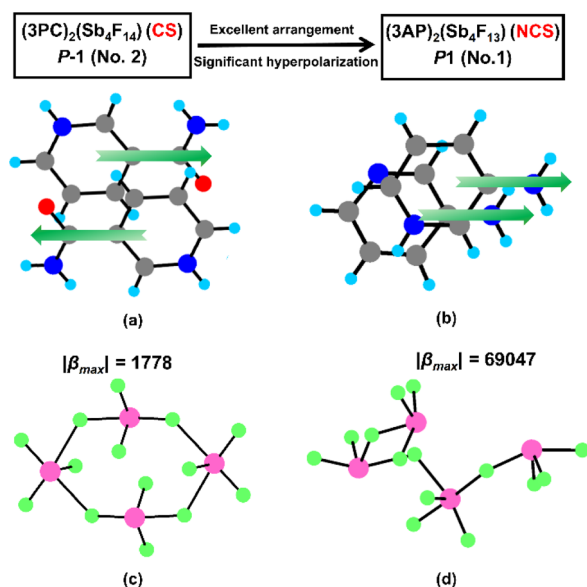


Fig. 4 The arrangement modes of $3PC^+$ (a) and $3AP^+$ (b) cations, and the Sb_4F_{14} tetramers in $(3PC)_2(Sb_4F_{14})$ (c) and $(3AP)_2(Sb_4F_{13})$ (d).

tetramer in $(3PC)_2(Sb_4F_{14})$ appears as a 4-MPR configuration, while the tetramer in $(3AP)_2(Sb_4F_{13})$ features an open semi-circular configuration (Fig. 4c and d). We calculated the hyperpolarizability of these two condensed Sb_4F_{14} groups. The hyperpolarizability of the semi-circular Sb_4F_{14} tetramer ($\beta = 69\,047$) is 38 times more than that of the Sb_4F_{14} 4-MPR ($\beta = 1778$).³⁶ Therefore, the polymerization mode of the fluoroantimonite groups is also crucial. The hyperpolarizabilities of closed cyclic clusters are often small due to the opposite arrangements of the functional groups within the rings.

Thermogravimetric analyses of $(3PC)_2(Sb_4F_{14})$ and $(3AP)_2(Sb_4F_{13})$ were performed in the temperature range of 20–800 °C under a nitrogen atmosphere. $(3PC)_2(Sb_4F_{14})$ and $(3AP)_2(Sb_4F_{13})$ can stabilize to 134 and 153 °C, respectively (Fig. S6†). The NCS polyfluoroantimonite, $(3AP)_2(Sb_4F_{13})$, underwent complete decomposition in the temperature range of 153 to 600 °C (Fig. S6†). Compared to other inorganic–organic hybrid NLO crystals, such as α - and β -($C_4H_5N_2O$)(IO_3)· HIO_3 (120 °C), $(C_9H_{14}N)SbCl_4$ (122 °C) and $KLi(HC_3N_3O_3) \cdot 2H_2O$ (125 °C), $(3AP)_2(Sb_4F_{13})$ exhibits better thermal stability.^{37–39}

The infrared (IR) spectra of $(3PC)_2(Sb_4F_{14})$ and $(3AP)_2(Sb_4F_{13})$ were analysed at room temperature in the wavelength range of 4000–400 cm^{-1} . Fig. S7† displays the IR spectra of the two compounds, and the corresponding assignments of the IR absorption peaks can be found in Table S6.† The absorption peaks around 3460 and 3250 cm^{-1} can be attributed to the N–H vibrations. The peaks observed at 3070–2850 cm^{-1} and 1050–650 cm^{-1} correspond to the C–H vibrations. The bands at 1690–1380 cm^{-1} and 1350–1110 cm^{-1} are assigned to C–C vibrations and C–N vibrations, respectively. The sharp absorption peaks in the range of 600–450 cm^{-1} can be attributed to the Sb–F vibrations. The above assignments are consistent with previously reported compounds.^{40–43}

The UV-vis-NIR diffuse reflectance spectra reveal that the cutoff edge of $(3PC)_2(Sb_4F_{14})$ and $(3AP)_2(Sb_4F_{13})$ is 272 and 313 nm, respectively. Their experimental optical band gaps were estimated to be 3.71 and 3.05 eV (Fig. S8†). Notably, the NCS $(3AP)_2(Sb_4F_{13})$ possesses a larger bandgap in comparison with the 1.97 eV of the 3AP molecule (Fig. 5). This phenomenon can be attributed to the highly electronegative fluorine anions. Compared with C-2p or N-2p states, the F-2p orbitals are usually located at deeper energy levels, which induces the high level valence band to shift toward lower energy, resulting in a wide band gap.^{44,45} Under the excitation at 330 nm, $(3AP)_2(Sb_4F_{13})$ exhibits a blue emission band with a peak at 400 nm (Fig. S9†), but its photoluminescence quantum yield is less than 1%. A large red shift is observed compared to the neutral 3AP, which is likely because of the protonation of the organic molecule and the influence of $[Sb_4F_{13}]_{\infty}^-$ on its highest occupied molecular orbital (HOMO) and lowest unoccupied molecular orbital (LUMO).^{46,47}

The birefringence of $(3PC)_2(Sb_4F_{14})$ and $(3AP)_2(Sb_4F_{13})$ was determined at $\lambda = 546$ nm using a polarized microscope (ZEISS AxioScope. A1).⁴⁸ Both crystals can achieve complete extinction under ortho-polarized light (Fig. S10 and S11†). The optical path differences of the two samples were measured to be 9.591 and 7.323 μm with the thicknesses of 64.88 and 28.39 μm (Fig. S10c and S11c†), respectively. According to the formula $R = |n_e - n_o| \times d = \Delta n \times d$, the experimental birefringences of $(3PC)_2(Sb_4F_{14})$ and $(3AP)_2(Sb_4F_{13})$ at 546 nm were measured to be 0.148 and 0.258, respectively. Due to the SCALP electrons of Sb^{3+} and the large anisotropy of organic amine cations, they both exhibit high birefringence. In particular, the birefringence of $(3AP)_2(Sb_4F_{13})$ exceeds those of all inorganic fluoroantimonites, such as $K_2SbMoO_2F_7$ (0.220@550 nm), $(NH_4)_3SbF_3(NO_3)_3$ (0.164@546 nm) and $K_2SbP_2O_7F$ (0.157@546 nm).^{49–51}

Under 1064 nm laser irradiation, powder SHG measurement was performed on $(3AP)_2(Sb_4F_{13})$ using KDP (particle size 210–300 μm) as a reference. As shown in Fig. 6a, $(3AP)_2(Sb_4F_{13})$ exhibits a strong SHG intensity of $8.1 \times$ KDP. Furthermore, it can realize phase matching at 1064 nm (Fig. 6b). In terms of structure–property relationships, the excellent SHG property of

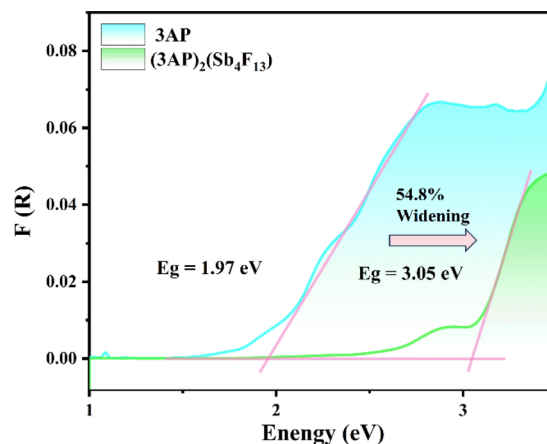


Fig. 5 The experimental bandgap spectra of 3AP and $(3AP)_2(Sb_4F_{13})$.

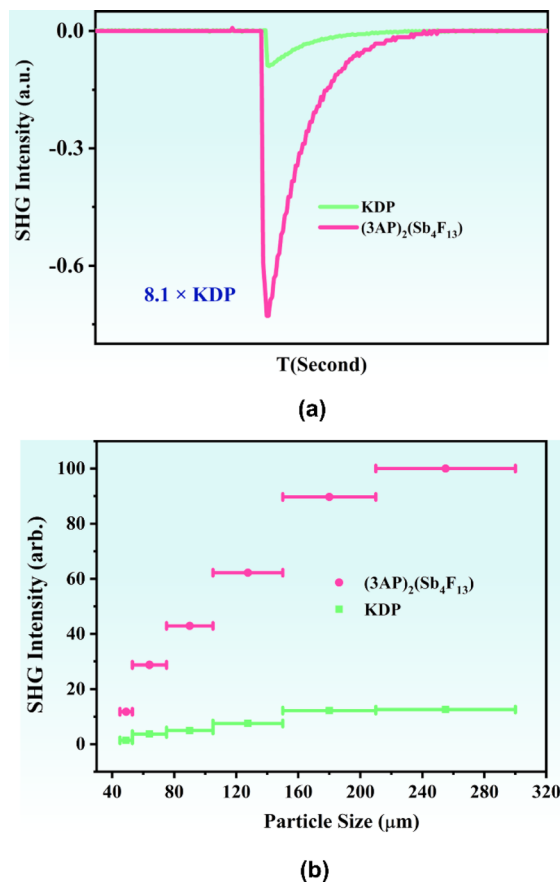


Fig. 6 Oscilloscope traces of the SHG signals for $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ (150–210 μm) (a) and the SHG intensity vs. particle size of $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ under 1064 nm laser irradiation (b).

$(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ can be ascribed to its polar structure, characterized by the unique $[\text{Sb}_4\text{F}_{13}]_{\infty}^-$ polyfluoroantimonite chain and the highly asymmetric organic amine groups 3AP. Note that the SHG intensity of $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ is larger than those of most of the reported fluoroantimonites (Table S7†). To evaluate the comprehensive performance of $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$, we further compared $(3\text{AP})_2\text{Sb}_4\text{F}_{13}$ with other inorganic–organic hybrid Sb(III)-based compounds in terms of SHG intensity, band gap, and thermal stability (Table S8†). As shown in Fig. 7, compared with representative compounds such as $(\text{C}_9\text{H}_{14}\text{N})\text{SbCl}_4$, $(\text{C}_9\text{H}_{26}\text{N}_3)\text{SbCl}_6$ and $\text{L-H}_2\text{his} \cdot \text{SbBr}_5 \cdot \text{H}_2\text{O}$, $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ shows great comprehensive performance.^{38,52,53}

Based on the reported powder method, we conducted a preliminary estimation of the LIDT value for $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ using sieved crystals.⁵⁴ The powder LIDT of $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ and AGS was assessed with a Q-switched pulse laser featuring a flat-top laser beam distribution, yielding values of 149.7 and 4.0 MW cm^{-2} , respectively. It suggests that the LIDT of $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ is approximately 37.5 times that of AGS. The powder LIDT of $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ is comparable with that of $\text{NaLu}(\text{SeO}_3)_2$, despite the substantial difference in their band gaps (3.05 vs. 5.30 eV).⁵⁵

To gain further insights into the structure–property relationships of $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ and $(3\text{PC})_2(\text{Sb}_4\text{F}_{14})$, theoretical

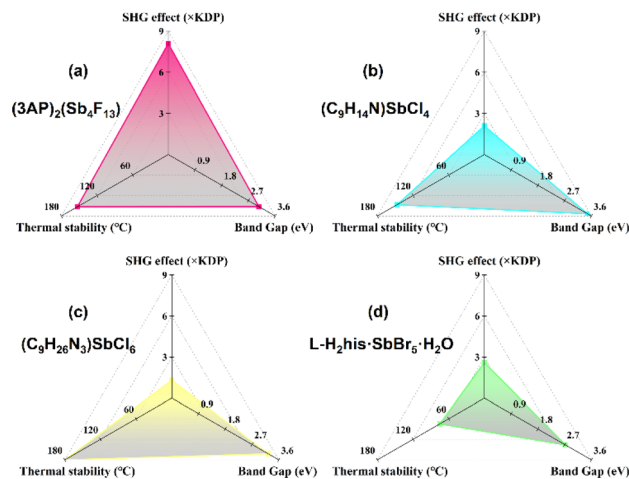


Fig. 7 The comparison of the SHG effect, band gap and thermal stability between $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ (a) and other representative inorganic–organic hybrid Sb(III)-based compounds: (b) $(\text{C}_9\text{H}_{14}\text{N})\text{SbCl}_4$, (c) $(\text{C}_9\text{H}_{26}\text{N}_3)\text{SbCl}_6$ and (d) $\text{L-H}_2\text{his} \cdot \text{SbBr}_5 \cdot \text{H}_2\text{O}$.

calculations were carried out using density functional theory (DFT). The calculation results along the high symmetric points in the first Brillouin region revealed that $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ is a direct bandgap compound with a bandgap of 2.28 eV (Fig. 8a), whereas $(3\text{PC})_2(\text{Sb}_4\text{F}_{14})$ is an indirect bandgap compound with a bandgap of 2.47 eV (Fig. S12a†). The calculated band gaps of $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ and $(3\text{PC})_2(\text{Sb}_4\text{F}_{14})$ are smaller than the experimental values (3.05 and 3.71 eV, respectively), which can be attributed to the limitation of the GGA method. Accordingly, in order to accurately analyse the optical properties of $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ and $(3\text{PC})_2(\text{Sb}_4\text{F}_{14})$, the scissor operators of 0.77 and 1.24 eV were adopted, respectively.

The total and partial density of states of $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ and $(3\text{PC})_2(\text{Sb}_4\text{F}_{14})$ are shown in Fig. 8b and S12b.† Since the optical properties of the compounds were mainly determined by the states near the Fermi level, we analyzed the top of the valence band (VB) and the bottom of the conduction band (CB). The VB near the Fermi level of $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ is primarily composed of C-2p, and N-2p states, while the CB is mainly attributed to C-2p and N-2p states (Fig. 8b). Therefore, the band gap of $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ is predominantly determined by C and N atoms. For $(3\text{PC})_2(\text{Sb}_4\text{F}_{14})$, the top of the VB is mainly contributed by O-2p orbitals, whereas the bottom of the CB is occupied by C-2p and N-2p orbitals (Fig. S12b†), indicating that the bandgap of $(3\text{PC})_2(\text{Sb}_4\text{F}_{14})$ is primarily dominated by O, C, and N atoms.

The refractive index dispersion curves depicted in Fig. 8c and S12c† demonstrate that $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ and $(3\text{PC})_2(\text{Sb}_4\text{F}_{14})$ exhibit strong anisotropy. For $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$, the trend of optical anisotropy is $n_{010} > n_{001} > n_{100}$, and the derived birefringence (Δn_{cal}) is 0.267@546 nm (Fig. 8b), in good agreement with the experimental value (0.258@546 nm). The calculated birefringence of $(3\text{PC})_2(\text{Sb}_4\text{F}_{14})$ at 546 nm is 0.154, which also closely matches the experimental value (0.149@546 nm). The electron differential density (EDD) maps of $(3\text{AP})_2(\text{Sb}_4\text{F}_{13})$ and $(3\text{PC})_2(\text{Sb}_4\text{F}_{14})$ indicate the strong anisotropy of π -conjugated organic amine cations and the highly stereochemically active

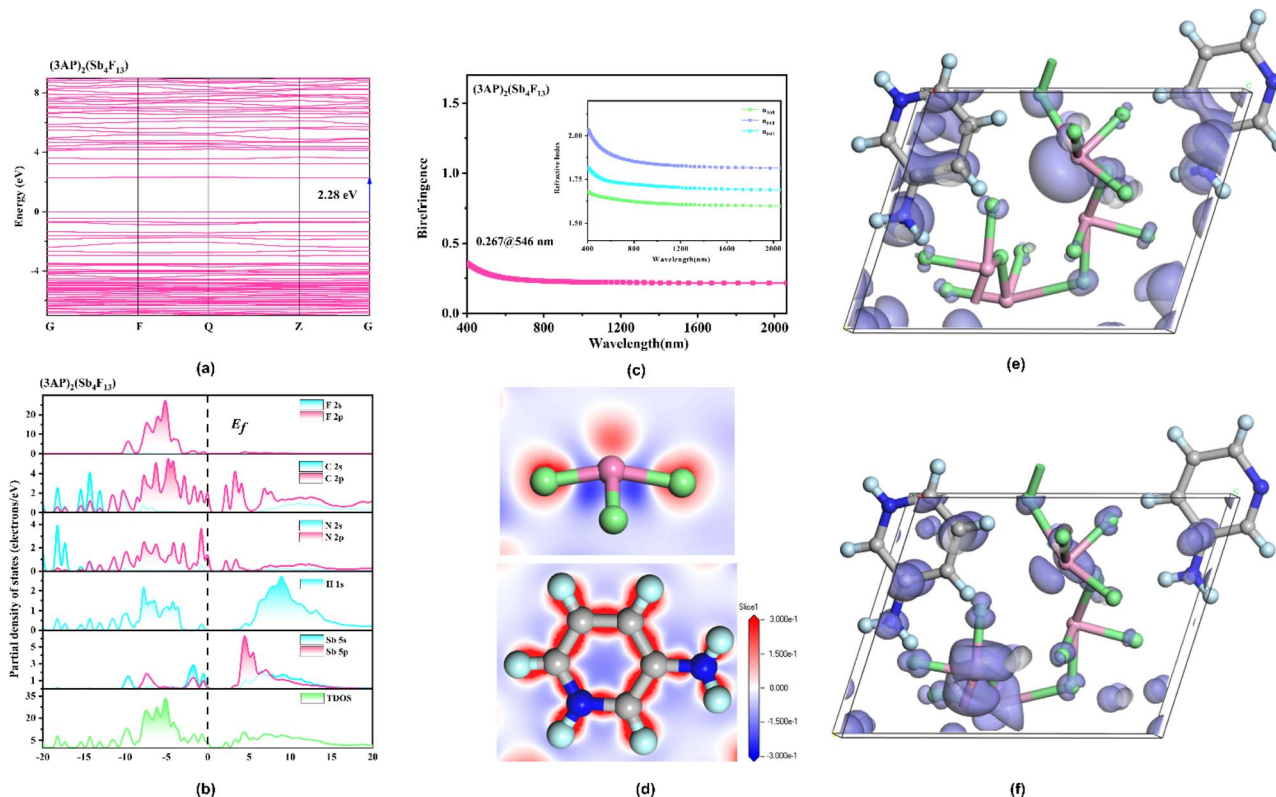


Fig. 8 Calculated band structure (a), total and partial densities of states (b), birefringence (c), electron density difference maps (d) and SHG density of d_{22} in the VB (e) and the CB (f) for $(3AP)_2(Sb_4F_{13})$.

Sb^{3+} (Fig. 8d and S12d[†]), which is beneficial for generating the large birefringence. In addition, the π - π stacking of organic amine cations in these two compounds also facilitates the enhancement of their birefringence.

The second-order nonlinear optical susceptibilities of $(3AP)_2(Sb_4F_{13})$ were also evaluated. With the constraints of Kleinman's symmetry and the space group ($P1$), its SHG tensor d_{22} gives the largest value (3.69 pm V^{-1}), which agrees well with the experimental value. Furthermore, to precisely reveal the nonlinear origin of $(3AP)_2(Sb_4F_{13})$, the SHG-weighted electron density (SHG density for short) of d_{22} was analyzed (Fig. 8e and f). Specifically, the SHG effect in the VB originates primarily from Sb-5s5p orbitals, the nonbonding N-2p orbitals, π_p bonding orbitals between C-C bonds and a small number of F-2p orbitals, and in the CB, the dominating SHG sources are the antibonding orbitals of π_p^* on $3AP^+$ groups, the unoccupied Sb-5p orbitals and some unoccupied F-2p orbitals. Numerically speaking, the SHG contribution percentages of helical $[Sb_4F_{13}]_{\infty}^-$ polyfluoroantimonite chain and π -conjugated $3AP^+$ cations were calculated to be 48.93% and 50.77%, respectively, indicating that the strong SHG effect of $(3AP)_2(Sb_4F_{13})$ originates from the synergistic effect of the $[Sb_4F_{13}]_{\infty}^-$ polyfluoroantimonite chain and π -conjugated $3AP^+$ cations.

Conclusions

In summary, a universal synthesis method for polyfluoroantimonites has been developed for the first time,

affording two novel polyfluoroantimonites, namely, $(3PC)_2(Sb_4F_{14})$ and $(3AP)_2(Sb_4F_{13})$. They exhibit totally different polymerization modes, which can be regulated by the organic amines. The $3AP^+$ cation with stronger hyperpolarizability results in a $[Sb_4F_{13}]_{\infty}^-$ helical polyfluoroantimonite chain, while the $3PC^+$ cation brings about a unique $(Sb_4F_{14})^{2-}$ 4-MPR. The NCS and polar $(3AP)_2(Sb_4F_{13})$ features a strong SHG effect of $8.1 \times \text{KDP}$, a large birefringence of $0.258@546 \text{ nm}$ and a high LIDT value of 149.7 MW cm^{-2} . Detailed structural analysis and theoretical calculations reveal that its strong SHG response originates from the synergistic effect of the helical $[Sb_4F_{13}]_{\infty}^-$ polyfluoroantimonite chain and π -conjugated $3AP^+$ cations. Our work demonstrates that the polyfluoroantimonite strategy is an effective approach for designing and synthesizing high-performance SHG materials. Further studies on NCS polyfluoroantimonites are in progress.

Data availability

The data that support the findings of this study are available in the ESI[†] of this article.

Author contributions

Jia-Hang Wu: conceptualization, investigation, data curation, and writing – original draft; Chun-Li Hu: software and formal analysis; Ya-Feng Li: data curation and visualization; Jiang-Gao



Mao: project administration and funding acquisition; Fang Kong: methodology, writing – review & editing, and supervision.

Conflicts of interest

There are no conflicts to declare.

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