

Cite this: *Chem. Sci.*, 2024, **15**, 7144

All publication charges for this article have been paid for by the Royal Society of Chemistry

Received 2nd February 2024
Accepted 11th April 2024

DOI: 10.1039/d4sc00786g
rsc.li/chemical-science

Introduction

Since lithium-ion batteries (LIBs) were successfully commercialized by Sony, LIBs have been widely used in mobile electronic devices, electric vehicles and stationary energy storage systems because of their long cycle life, low self-discharge and environmental friendliness.^{1–3} However, the energy density of state-of-the-art LIBs cannot satisfy the increasing demand of drivers, resulting in severe range anxiety. Lithium metal batteries (LMBs) are a promising candidate for next-generation rechargeable batteries due to their high energy density obtained by using metallic lithium as anode materials.^{4–6} Unfortunately, the notorious Li dendrite growth and flammable organic liquid electrolytes are unavoidable results in severe and sporadic safety concerns.^{7–9} Therefore, unique technologies for higher energy density and safer LMBs are being urgently explored.^{10,11}

Solid-state electrolytes (SSEs) that can effectively mitigate unreliable safety have emerged as the advancement of science and technology.^{12,13} In recent years, a variety of SSEs have been explored, such as sodium superionic conductors,¹⁴ sulfides,^{15,16} perovskites¹⁷ and garnets.^{18,19} Among them, the garnet-type

Interface engineering of $\text{Li}_{6.75}\text{La}_3\text{Zr}_{1.75}\text{Ta}_{0.25}\text{O}_{12}$ via *in situ* built LiI/ZnLi_x mixed buffer layer for solid-state lithium metal batteries†

Lei Zhai,^a Jinhuan Wang,^a Xiaoyu Zhang,^{ID, a} Xunzhu Zhou,^b Fuyi Jiang,^a Lin Li,^{ID, *b} and Jianchao Sun,^{ID, *a}

Garnet-type solid-state Li metal batteries (SSLMBs) are viewed as hopeful next-generation batteries due to their high energy density and safety. However, the major obstacle to the development of garnet-type SSLMBs is the lithiophobicity of $\text{Li}_{6.75}\text{La}_3\text{Zr}_{1.75}\text{Ta}_{0.25}\text{O}_{12}$ (LLZTO), resulting in a large interfacial impedance. Herein, a LiI/ZnLi_x mixed ion/electron conductive buffer layer is constructed at the interface by an *in situ* reaction of molten Li metal with ZnI_2 film. This mixed buffer layer ensures close contact between the Li metal and garnet, significantly reducing interfacial impedance. As a result, the Li symmetrical cell with the LiI/ZnLi_x buffer layer shows an interface impedance of $10.3 \Omega \text{ cm}^2$, much lower than that of the cell with bare LLZTO ($1173.4 \Omega \text{ cm}^2$). The critical current density (CCD) is up to 2.3 mA cm^{-2} , and the symmetric cells present a long cycle life of 2000 h at 0.1 mA cm^{-2} and 800 h at 1.0 mA cm^{-2} . In addition, the full cells assembled with the LiFePO_4 cathode show a capacity of $143.9 \text{ mA h g}^{-1}$ after 200 cycles at 0.5C with a low-capacity decay of 0.021% per cycle. This work reveals a simple, feasible, and practical interface modification strategy for solid-state Li metal batteries.

$\text{Li}_{6.75}\text{La}_3\text{Zr}_{1.75}\text{Ta}_{0.25}\text{O}_{12}$ (LLZTO) attracts more attention because of its high ionic conductivity and chemical stability against Li metal.^{20–22} Unfortunately, LLZTO has poor wettability

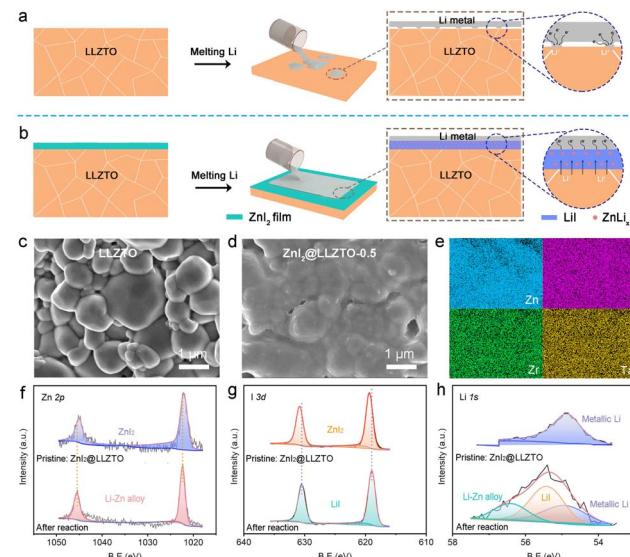


Fig. 1 Pouring molten Li metal on (a) bare LLZTO and (b) ZnI_2 @LLZTO. SEM images of (c) bare LLZTO and (d) ZnI_2 @LLZTO-0.5. (e) Corresponding Zn, I, Zr and Ta elemental mappings. High-resolution XPS analysis of ZnI_2 @LLZTO pellet surface before and after the reaction: (f) Zn 2p, (g) I 3d, and (h) Li 1s.

^aSchool of Environment and Material Engineering, Yantai University, Yantai 264005, Shandong, China. E-mail: jianchaoabc@163.com

^bInstitute for Carbon Neutralization, College of Chemistry and Materials Engineering, Wenzhou University, Wenzhou, Zhejiang 325035, China. E-mail: linli@wzu.edu.cn

† Electronic supplementary information (ESI) available. See DOI: <https://doi.org/10.1039/d4sc00786g>



on lithium metal, which can lead to point-to-point contact at the interface (Fig. 1a). The point-to-point contact not only increases interfacial resistance but also results in locally uneven Li deposition.^{23,24} Therefore, interfacial buffer layers have been proposed to improve the surface wettability of LLZTO to Li metal.^{25–27}

According to the different conductive carriers, the interfacial buffer layers can be divided into three types: electronic conductive, ionic conductive, and mixed conductive buffer layers.^{28–30} Currently, Al,³¹ Au,³² Zn,³³ Sn,³⁴ Ge³⁵ and Sb¹⁸ have been reported as electronic conductive buffer layers. Undoubtedly, the electronic conductive buffer layer improves interfacial wettability, but electrons can easily pass through the buffer layer and combine with lithium ions.^{36–38} Especially in the case of high current density, lithium ions move slowly in solid-state electrolytes, and the rapid passage of electrons through the buffer layer generates an uneven electric field, resulting in the formation of dendrites along grain boundaries.³⁹ Compared with the electron conductive buffer layer, the ionic conductive buffer layer can effectively block the attack of electrons on the garnet.⁴⁰ In general, the ionic conductive buffer layer exhibits an obvious overpotential increase during the cycle, which is caused by high interface impedance.^{41,42} In conclusion, the microstructure and wettability of the lithium/garnet interface can be improved by the construction of the lithium alloy phase. However, the high electrical conductivity of the alloy interface layer causes the lithium dendrites to penetrate the solid electrolyte. Although the use of ionic conductive layers alone can prevent the destruction of electrons during lithium deposition, the electrical insulation characteristics increase the impedance between lithium metal and the solid electrolyte, which increases the overpotential during the cycle and causes battery failure. Therefore, it is difficult for ionic or electronic conduction alone to maintain a long-term stable cycle, and a mixed ion/electron conductive layer is a promising option.^{43–48} Taking Cu₃N as an example, a mixed conductive layer Li₃N–Cu was prepared *in situ* on the LLZTO surface by the reaction of Cu₃N with Li metal.⁴⁶ Li₃N plays a role in conducting Li⁺, and Cu nanoparticles can uniformly disperse the electric field in the mixed conductive layer, thereby inhibiting the nucleation of Li dendrites. However, the mixed buffer layer is mainly prepared by magnetron sputtering, which increases the difficulty of experimental operation and the cost of production. Therefore, it is necessary to find a fast, simple and low-cost method to prepare a mixed interface buffer layer, which is helpful to realize the industrialization of solid-state lithium metal batteries.

Herein, a simple method is proposed for the construction of a mixed LiI/ZnLi_x ion/electron-conductive layer on LLZTO surface by an *in situ* reaction of ZnI₂ film and molten Li metal (Fig. 1b). The mixed LiI/ZnLi_x layer enhances the interaction between Li metal and LLZTO. As a result, the cell with the LiI/ZnLi_x buffer layer shows a lower interfacial impedance of 10.3 Ω cm² compared with the cell with bare LLZTO (1173.4 Ω cm²). The high Li⁺/e[−] conductivity of the mixed layer promotes Li⁺ transportation and Li nucleation, inhibiting the growth of Li dendrites along grain boundaries. Meanwhile, the simulation results show a much more uniform distribution of electric field

and current density at the surface of the LiI/ZnLi_x buffer layer than that of bare LLZTO. As a proof, symmetrical cells with the LiI/ZnLi_x buffer layer exhibit stable plating/stripping performance of 2000 h at a low current density/capacity of 0.1 mA cm^{−2}/0.1 mA h cm^{−2} and 800 h at a high current density/capacity of 1.0 mA cm^{−2}/1.0 mA h cm^{−2}. Moreover, the potential of the practical application of the LiI/ZnLi_x buffer layer for high-performance solid-state Li metal batteries is also demonstrated.

Results and discussion

LLZTO was prepared by the solid reaction method.⁴⁹ The X-ray diffraction results reveal that the synthesized LLZTO is a pure cubic garnet phase (Fig. S1†). As shown in Fig. 1c, the LLZTO pellet with high relative density (94.09%) is composed of generous crystalline LLZTO particles with smooth surfaces. ZnI₂@LLZTO pellets were prepared by placing the bare LLZTO pellets in ZnI₂ solution for 10 seconds (Fig. S2†). After soaking in 0.1 mol per L ZnI₂ solution, the surface of LLZTO becomes rough, but the boundary of each crystal particle is still clear (Fig. S3†). When the concentration of ZnI₂ is 0.5 mol L^{−1}, the surface of LLZTO is uniformly covered by the ZnI₂ buffer layer (Fig. 1d). Noticeably, obvious ZnI₂ particles appear on the surface of LLZTO when the concentration of ZnI₂ increases to 1.0 and 1.5 mol L^{−1}. Meanwhile, the elemental mappings of ZnI₂@LLZTO-0.5 disclose the uniform distribution of Zn, I, Zr and Ta elements (Fig. 1e). The Raman spectra of the bare LLZTO surface and the LLZTO treated with different concentrations of ZnI₂ solution were compared (Fig. S4†). The peak at 140 cm^{−1} is the typical characteristic peak of ZnI₂. The intensity of this characteristic peak increases with the increase of ZnI₂ concentration. When the concentration of ZnI₂ exceeds 1 mol L^{−1}, the characteristic peak becomes obvious. The Raman peaks belonging to LLZTO show no change in any sample. Due to the relatively low content of ZnI₂ on the LLZTO surface, the reaction between Li metal and ZnI₂ cannot be directly observed. To verify the above reaction, a large amount of ZnI₂ powder was added to the molten Li metal, and they underwent a violent reaction (Fig. S5†). XRD test was performed on the product after the reaction, and the (111) diffraction peak of LiI was clearly observed (Fig. S6†). The above experiment proves that our idea can be realized. Subsequently, the mixed conductive buffer layers were prepared *in situ* by the reaction of molten Li metal with ZnI₂ thin film on the surface of LLZTO. To determine the composition of the interface buffer layer on the LLZTO surface, the detailed Zn, I, and Li species at the interfaces were investigated by XPS. The initial peak position of Zn 2p spectra is 1045.0 and 1022.1 eV.³³ After the reaction, the peak position is shifted to 1044.2 and 1021.3 eV due to the formation of the Zn–Li bond, indicating the formation of ZnLi_x alloy (Fig. 1f). Meanwhile, analogous variation is also observed in I 3d XPS spectra (Fig. 1g). In Li 1s spectra, two new peaks appear after the reaction, corresponding to the formation of the ZnLi_x alloy at 56.4 eV and LiI at 55.4 eV, respectively (Fig. 1h). In addition, we also provided XPS spectra of oxygen element (Fig. S7†). The peak position of O 1s before and after the reaction was 531.4 eV,



without any change. This indicates that the reaction at the interface only involves ZnI_2 and Li metal.

Subsequently, we visually observed the interfacial wettability of LLZTO with or without the ZnI_2 buffer layer. A liquid Li droplet appears spherical on the bare LLZTO surface (inset in Fig. 2a). In addition, the cross-sectional SEM image shows a significant gap at the Li/LLZTO interface (Fig. 2a), which leads to high interfacial resistance and uneven interfacial electric field. In contrast, $\text{ZnI}_2@\text{LLZTO}$ has excellent lithophilic properties, and the Li metal spreads evenly on the SSE surface (inset in Fig. 2b). The SEM image clearly shows that the Li/ $\text{ZnI}_2@\text{LLZTO}$ interface is compact and seamless (Fig. 2b). To further verify the effect of the buffer layer on interface impedance, Li/SSE/Li symmetric cells were assembled. Fig. 2c shows the EIS data for the Li/SSE/Li cells at room temperature. The cell with bare LLZTO presents a large semicircle in the Nyquist plots, which means a high interfacial resistance ($1173.4 \Omega \text{ cm}^2$). In contrast, the interfacial impedance of $\text{Li}/\text{ZnI}_2@\text{LLZTO}/\text{Li}$ drops substantially to $10.3 \Omega \text{ cm}^2$. The result indicates that the ZnI_2 buffer layer can effectively reduce the interface impedance, which is beneficial for the benign electrochemical reaction at the interface. Furthermore, the interface impedances of Li/ $\text{ZnI}_2@\text{LLZTO}/\text{Li}$ cells with different concentrations of ZnI_2 were compared (Fig. 2d and e). When the concentration of ZnI_2 is 0.5 mol L^{-1} , the interface impedance is the lowest. The possible reasons for this result are as follows: (i) when the concentration of ZnI_2 is low ($<0.5 \text{ mol L}^{-1}$), the surface of LLZTO cannot be completely covered by ZnI_2 , which hardly achieves the best effect of improving the interface; (ii) the ZnI_2 buffer layer on the

surface of LLZTO is too thick to hinder the migration of Li^+ at a high concentration of ZnI_2 ($>1.0 \text{ mol L}^{-1}$).

In order to further select the optimal concentration, we assembled symmetrical cells on solid electrolytes treated with different concentrations of ZnI_2 solution for the galvanostatic charge-discharge test. As shown in Fig. S8,† when the concentration of ZnI_2 is 0.5 mol L^{-1} , the polarization voltage of the symmetric cell is the smallest (only 13 mV). The results of impedance and galvanostatic charge-discharge show that 0.5 mol L^{-1} is the best concentration. Based on the above experimental results, $\text{ZnI}_2@\text{LLZTO-0.5}$ is selected as the research object in this work. In addition, the interfacial activation energies (E_a) of Li/LLZTO and Li/ $\text{ZnI}_2@\text{LLZTO}$ were obtained through variable temperature tests (Fig. S9†). The E_a of the Li/LLZTO interface is 30.3 kJ mol^{-1} , whereas that of the Li/ $\text{ZnI}_2@\text{LLZTO}$ interface is only 19.2 kJ mol^{-1} (Fig. 2f). Therefore, LiI and ZnLi_x alloy buffer layer formed *in situ* by the reaction of ZnI_2 and Li metal can effectively reduce the impedance and promote the Li^+ migration at the interface.

The superiority of SSE with a buffer layer to suppress Li dendrites was evaluated by the critical current density (CCD). As shown in Fig. 3a, the Li/LLZTO/Li cell experiences a short circuit at a current density of 0.15 mA cm^{-2} , indicating that point contact at the interface can enhance the local electric field, leading to a “tip effect” and inducing Li dendrites to penetrate the SSE. In contrast, the CCD of the Li/ $\text{ZnI}_2@\text{LLZTO}/\text{Li}$ cell reaches 2.3 mA cm^{-2} due to the presence of the ZnI_2 buffer layer. Fig. S10† shows the CCD curve of a symmetrical cell with

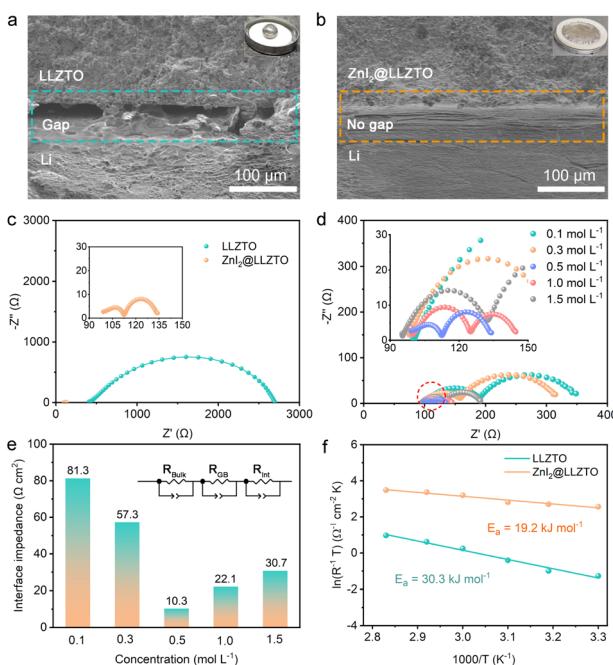


Fig. 2 Cross-sectional SEM images of interfaces: (a) Li/LLZTO and (b) Li/ $\text{ZnI}_2@\text{LLZTO}$. The insets show the wettability behavior of liquid Li on SSE. (c) Nyquist plots of Li/SSE/Li symmetric cells. (d) Nyquist plots and (e) interfacial resistance values with different concentrations of ZnI_2 . (f) Interface activation energy.

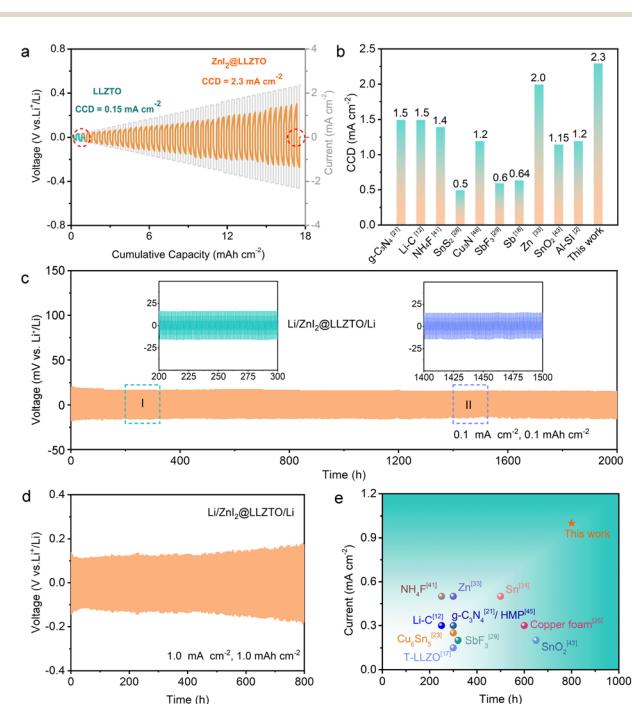


Fig. 3 (a) Critical current density test of the Li/SSE/Li cell. (b) Comparison of CCD with LLZTO modified in different ways. Charge-discharge voltage profiles of Li/ $\text{ZnI}_2@\text{LLZTO}/\text{Li}$ cells: (c) 0.1 mA cm^{-2} , 0.1 mA h cm^{-2} and (d) 1.0 mA cm^{-2} , 1.0 mA h cm^{-2} . (e) Comparison of current density and cycle life with LLZTO modified in different ways.



a ZnI_2 concentration of 1.5 mol L^{-1} . When the current density increased to 1 mA cm^{-2} , the cell experienced a soft short-circuit. This result is consistent with the impedance test, indicating that the thickness of the buffer layer has a significant impact on battery performance. Noticeably, this work is superior to the reported methods for modifying the LLZTO interface (Fig. 3b).^{2,12,18,21,28,29,33,41,43,46} This may be because the good affinity of LiI to LLZTO as well as the good lithiophilicity of the ZnLi_x alloy to Li metal is better than that of the buffer layer which only has an affinity for single LLZTO or Li metal.⁵⁰ To verify this inference, the Li/SSE/Li symmetric cell with I_2 as the buffer layer was assembled. The CCD of the Li/SSE/Li cell with the I_2 buffer layer is 0.5 mA cm^{-2} (Fig. S11†), which is lower than the 2.3 mA cm^{-2} of the mixed LiI/ZnLi_x buffer layer.

More importantly, the $\text{Li}/\text{ZnI}_2@\text{LLZTO}/\text{Li}$ symmetric cells also present excellent long-term cycling ability. As shown in Fig. 3c, the $\text{Li}/\text{ZnI}_2@\text{LLZTO}/\text{Li}$ cell delivers a stable cycle for 2000 h with an overpotential of only 13 mV at a current density of 0.1 mA cm^{-2} and an area capacity of 0.1 mA h cm^{-2} . In contrast, the $\text{Li}/\text{LLZTO}/\text{Li}$ cell exhibits a large overpotential, and the voltage rapidly decreases after 25 h, indicating a short circuit (Fig. S12†). Even under 0.5 mA cm^{-2} to 0.5 mA h cm^{-2} and 1.0 mA cm^{-2} to 1.0 mA h cm^{-2} conditions, the symmetrical cells with the ZnI_2 interface buffer layer can stably cycle for 1600 and 800 h, respectively (Fig. 3d and S13†). The electrochemical performance is better than that of most of the reported Li/SSE/Li cells (Fig. 3e, S14 and Table S1†).^{3,12,17,21,22,24,25,29,33,34,41,43,45,48} In general, most reported Li/garnet/Li cells operate at low current density with short cycle life. The above results further demonstrate that the mixed conductive buffer layers are superior to a single electronic or ionic conductive buffer layer.^{51,52}

In order to further confirm the influence of the ZnI_2 buffer layer on the interface stability, SEM and EIS were performed on

the cells after 200 h (100 cycles). As shown in Fig. 4a and b, the Li metal still maintains close contact with the $\text{ZnI}_2@\text{LLZTO}$, without gaps or lithium dendrites. Compared with the initial state, the interface impedance increases by only $3.2 \Omega \text{ cm}^2$ (Fig. 4c), which indicates that the buffer layer plays a bridging role and effectively promotes the Li^+ migration at the interface. In contrast, $\text{Li}/\text{LLZTO}/\text{Li}$ shows more obvious cracks with loose Li dendrites at the interface after cycling (Fig. 4d and e). Poor interface contact leads to large interface impedance, which is also the root cause of battery failure (Fig. 4f). It's worth noting that the interface of Li metal and LLZTO is still in close contact in the symmetrical cell with the ZnI_2 buffer layer even at high current density and capacity (1.0 mA cm^{-2} and 1.0 mA h cm^{-2}) (Fig. S15†). Moreover, the interfacial impedance of the symmetric cell increases by only $7.3 \Omega \text{ cm}^2$. The results prove that the LiI/ZnLi_x mixed layer plays a crucial role in stabilizing the interface.

Subsequently, the electric field distribution and current density of the Li/SSE interface were simulated by COMSOL Multiphysics to further investigate the effect of the ZnI_2 buffer layer on Li dendrite growth. Because Li metal and bare LLZTO are in point contact at the interface, the electric field strength and current density at the contact point are much greater than those at non-contact locations (Fig. 4g and i). This causes Li^+ in the bare LLZTO to aggregate and deposit at the contact point, which is also the reason for large polarization and dendrite formation. For the $\text{Li}/\text{ZnI}_2@\text{LLZTO}$ interface, because of the uniform distribution of Li metal on the LLZTO surface, the electric field and current density at the interface are evenly distributed (Fig. 4h and j). The transmission and diffusion of lithium ions at the interface are very smooth, resulting in small interface impedance and excellent electrochemical performance of symmetrical cells.

The $\text{Li}/\text{ZnI}_2@\text{LLZTO}/\text{LFP}$ full cell was assembled to demonstrate the potential of the practical application of the ZnI_2 buffer layer (Fig. 5a). The $\text{Li}/\text{ZnI}_2@\text{LLZTO}/\text{LFP}$ full cell shows an excellent rate performance, which hands over the specific discharge capacities of 155.8 , 147.6 , 145.1 , 143.7 , 124.4 , and $112.0 \text{ mA h g}^{-1}$ at 0.1 , 0.2 , 0.3 , 0.5 , 0.8 , and 1.0C , respectively (Fig. 5b). When the current goes back to 0.1C , the discharge

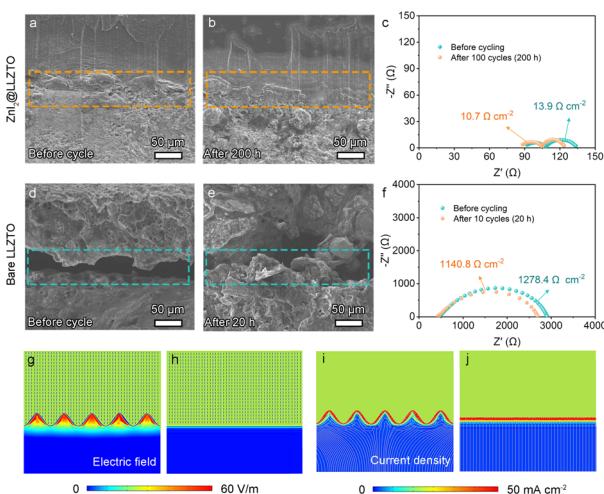


Fig. 4 The cross-sectional SEM images and Nyquist plots of Li/SSE/Li symmetric cells before and after cycles at 0.1 mA cm^{-2} , 0.1 mA h cm^{-2} : (a–c) $\text{ZnI}_2@\text{LLZTO}$; (d–f) bare LLZTO. Simulations of electric field and current density at the interface of different solid electrolytes: (g and i) bare LLZTO; (h and j) $\text{ZnI}_2@\text{LLZTO}$. Scale bar, 100 μm .

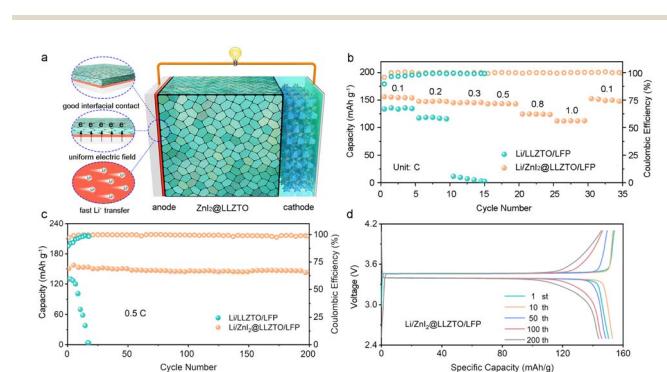


Fig. 5 (a) Schematic construction of the $\text{Li}/\text{ZnI}_2@\text{LLZTO}/\text{LFP}$ full cell. (b) Rate performance of Li/SSE/LFP full cells with $\text{ZnI}_2@\text{LLZTO}$ or LLZTO in 2.5 – 4.1 V . (c) The cycling performance of full cells at 0.5C . (d) Charge–discharge profiles of the $\text{Li}/\text{ZnI}_2@\text{LLZTO}/\text{LFP}$ full cell at 0.5C .



capacity can recover to 152.5 mA h g⁻¹, indicating significant reversibility. In contrast, as the current increases, the full cell assembled with bare LLZTO exhibits critical capacity degradation. The excellent rate performance of full cells with ZnI₂ is attributed to the smooth transfer of Li⁺ at the interface. In addition, the full cell with bare LLZTO shows poor long-term performance (Fig. 5c). By comparison, the full cell with ZnI₂@LLZTO displays a discharge capacity of 143.9 mA h g⁻¹ with high coulombic efficiency (99%) after 200 cycles. In particular, the capacity loss per cycle is only 0.021%. As shown in Fig. 5d, the smooth charge-discharge profiles suggest that the Li⁺ transmission in the full cell is rapid and uniform. More importantly, the electrochemical performance is better than that of most of the full cells reported in the literature (Fig. S16†).

Conclusions

In summary, we constructed a mixed LiI/ZnLi_x ion/electron-conductive layer on an LLZTO surface by *in situ* reaction of ZnI₂ film and molten Li metal. The mixed LiI/ZnLi_x layer enhances the interaction between Li metal and LLZTO, resulting in a low interfacial impedance of 10.3 Ω cm². The mixed conductive layer regulates the electric field at the interface, which ensures a uniform Li deposition and inhibits the formation of Li dendrites. As a result, the symmetric cells exhibit a long cycle life of 800 h at 1.0 mA cm⁻² and 1.0 mA h cm⁻². Moreover, the full cells show excellent cycling performance with a capacity of 143.9 mA h g⁻¹ after 200 cycles and the capacity loss per cycle is only 0.021%. This simple modification strategy provides a good prospect for the design of solid-state Li metal batteries based on garnets.

Data availability

All experimental data is provided in the ESI.†

Author contributions

J. S. and L. L. proposed the concept and supervised the work; L. Z. and J. W. designed the experiments and wrote the paper; X. Z. helped to analyze the data; X. Z. and F. J. helped to summarize the data and supervised the research; all authors discussed the results and revised the manuscript. All authors have approved the final version of the manuscript.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

This work was supported by the National Natural Science Foundation of China (22209140, 52202286, 52072328), Natural Science Foundation of Shandong Province (ZR2022QE059), Natural Science Foundation of Zhejiang Province (LZ21E010001), Shandong Laboratory of Advanced Materials and Green Manufacturing at Yantai (Yantai) (AMGM2023A08).

Notes and references

- D. Liu, Z. Bai, M. Li, A. Yu, D. Luo, W. Liu, L. Yang, J. Lu, K. Amine and Z. Chen, *Chem. Soc. Rev.*, 2020, **49**, 5407–5445.
- L. Zhai, K. Yang, F. Jiang, W. Liu, Z. Yan and J. Sun, *J. Energy Chem.*, 2023, **79**, 357–364.
- J. Cui, S. Yao, A. Guerfi, C. Kim, J. B. Goodenough and H. Khani, *Energy Storage Mater.*, 2022, **53**, 899–908.
- B. Liu, J. Zhang and W. Xu, *Joule*, 2018, **2**, 833–845.
- Y. Yin, J. Yang, J. Luo, G. Lu, Z. Huang, J. Wang, P. Li, F. Li, Y. Wu, T. Tian, Y. Meng, H. Mo, Y. Song, J. Yang, L. Feng, T. Ma, W. Wen, K. Gong, L. Wang, H. Ju, Y. Xiao, Z. Li, X. Tao and H. Yao, *Nature*, 2023, **616**, 77–83.
- J. Wang, R. Chen, L. Yang, M. Zan, P. Chen, Y. Li, W. Li, H. Yu, X. Yu, X. Huang, L. Chen and H. Li, *Adv. Mater.*, 2022, **34**, 2200655.
- W. Wu, Z. Song, Y. Dai, X. Zheng, G. Chai, J. Yang and W. Luo, *Adv. Energy Mater.*, 2022, **12**, 2200999.
- K. Yan, Z. Lu, H. Lee, F. Xiong, P. Hsu, Y. Li, J. Zhao, S. Chu and Y. Cui, *Nat. Energy*, 2016, **1**, 16010.
- Y. Liu, X. Tao, Y. Wang, C. Jiang, C. Ma, O. Sheng, G. Lu and X. Lou, *Science*, 2022, **375**, 739–745.
- C. Yu, F. Zhao, J. Luo, L. Zhang and X. Sun, *Nano Energy*, 2021, **83**, 105858.
- W. Luo, Y. Gong, Y. Zhu, K. K. Fu, J. Dai, S. D. Lacey, C. Wang, B. Liu, X. Han, Y. Mo, E. D. Wachsman and L. Hu, *J. Am. Chem. Soc.*, 2016, **138**, 12258–12262.
- J. Duan, W. Wu, A. M. Nolan, T. Wang, J. Wen, C. Hu, Y. Mo, W. Luo and Y. Huang, *Adv. Mater.*, 2019, **31**, 1807243.
- X. Liang, Q. Pang, I. R. Kochetkov, M. S. Sempere, H. Huang, X. Sun and L. F. Nazar, *Nat. Energy*, 2017, **2**, 17119.
- Y. Xiao, K. Jun, Y. Wang, L. J. Miara, Q. Tu and G. Ceder, *Adv. Energy Mater.*, 2021, **11**, 2101437.
- Y. Jin, K. Liu, J. Lang, X. Jiang, Z. Zheng, Q. Su, Z. Huang, Y. Long, C. Wang, H. Wu and Y. Cui, *Joule*, 2019, **3**, 1–13.
- C. Duan, Z. Cheng, W. Li, F. Li, H. Liu, J. Yang, G. Hou, P. He and H. Zhou, *Energy Environ. Sci.*, 2022, **15**, 3236–3245.
- J. Gao, J. Zhu, X. Li, J. Li, X. Guo, H. Li and W. Zhou, *Adv. Funct. Mater.*, 2021, **31**, 2001918.
- R. Dubey, J. Sastre, C. Cancellieri, F. Okur, A. Forster, L. Pompizii, A. Priebe, Y. Romanyuk, L. Jeurgens, M. Kovalenko and K. Kravchyk, *Adv. Energy Mater.*, 2021, **11**, 2102086.
- L. Fan, S. Wei, S. Li, Q. Li and Y. Lu, *Adv. Energy Mater.*, 2018, **8**, 1702657.
- T. Jiang, P. He, G. Wang, Y. Shen, C. Nan and L. Fan, *Adv. Energy Mater.*, 2020, **10**, 1903376.
- Y. Huang, B. Chen, J. Duan, F. Yang, T. Wang, Z. Wang, W. Yang, C. Hu, W. Luo and Y. Huang, *Angew. Chem., Int. Ed.*, 2020, **59**, 3699–3704.
- P. Srivastava, Y. K. Liao, K. Iputera, S. F. Hu and R. S. Liu, *ChemSusChem*, 2023, **16**, e202300504.
- W. Feng, X. Dong, Z. Lai, X. Zhang, Y. Wang, C. Wang, J. Luo and Y. Xia, *ACS Energy Lett.*, 2019, **4**, 1725–1731.
- Y. Jin, K. Liu, J. Lang, D. Zhuo, Z. Huang, C. Wang, H. Wu and Y. Cui, *Nat. Energy*, 2018, **3**, 732.



25 J. Duan, L. Huang, T. Wang, Y. Huang, H. Fu, W. Wu, W. Luo and Y. Huang, *Adv. Funct. Mater.*, 2020, **30**, 1908701.

26 H. Zheng, G. Li, R. Ouyang, Y. Han, H. Zhu, Y. Wu, X. Huang, H. Liu and H. Duan, *Adv. Funct. Mater.*, 2022, **32**, 2205778.

27 Y. Zhong, Y. Xie, S. Hwang, Q. Wang, J. J. Cha, D. Su and H. Wang, *Angew. Chem., Int. Ed.*, 2020, **59**, 14003–14008.

28 B. Zhao, W. Ma, B. Li, X. Hu, S. Lu, X. Liu, Y. Jiang and J. Zhang, *Nano Energy*, 2022, **91**, 106643.

29 A. Wang, J. Li, M. Yi, Y. Xie, S. Chang, H. Shi, L. Zhang, M. Bai, Y. Zhou, Y. Lai and Z. Zhang, *Energy Storage Mater.*, 2022, **49**, 246–254.

30 W. Huang, N. Zhao, Z. Bi, C. Shi, X. Guo, L. Fan and C. Nan, *Mater. Today Nano*, 2020, **10**, 100075.

31 K. Fu, Y. Gong, B. Liu, Y. Zhu, S. Xu, Y. Yao, W. Luo, C. Wang, S. D. Lacey, J. Dai, Y. Chen, Y. Mo, E. Wachsman and L. Hu, *Sci. Adv.*, 2017, **3**, e1601659.

32 N. Zhao, R. Fang, M. He, C. Chen, Y. Li, Z. Bi and X. Guo, *Rare Met.*, 2018, **37**, 473–479.

33 Z. Wan, K. Shi, Y. Huang, L. Yang, Q. Yun, L. Chen, F. Ren, F. Kang and Y. He, *J. Power Sources*, 2021, **505**, 230062.

34 M. He, Z. Cui, C. Chen, Y. Li and X. Guo, *J. Mater. Chem. A*, 2018, **6**, 11463–11470.

35 W. Luo, Y. Gong, Y. Zhu, Y. Li, Y. Yao, Y. Zhang, K. Fu, G. Pastel, C. Lin, Y. Mo, E. D. Wachsman and L. Hu, *Adv. Mater.*, 2017, **29**, 1606042.

36 Y. Guo, S. Wu, Y. He, F. Kang, L. Chen, H. Li and Q. Yang, *eScience*, 2022, **2**, 138–163.

37 W. Feng, X. Dong, X. Zhang, Z. Lai, P. Li, C. Wang, Y. Wang and Y. Xia, *Angew. Chem., Int. Ed.*, 2020, **132**, 5384–5387.

38 Q. Pang, X. Liang, I. R. Kochetkov, P. Hartmann and L. Nazar, *Angew. Chem., Int. Ed.*, 2018, **130**, 9943–9946.

39 Y. Lee, S. Fujiki, C. Jung, N. Suzuki, N. Yashiro, R. Omoda, D. Ko, T. Shiratsuchi, T. Sugimoto, S. Ryu, J. Ku, T. Watanabe, Y. Park, Y. Aihara, D. Im and I. Han, *Nat. Energy*, 2020, **5**, 299–308.

40 X. Ji, S. Hou, P. Wang, X. He, N. Piao, J. Chen, X. Fan and C. Wang, *Adv. Mater.*, 2020, **32**, 2002741.

41 H. Duan, W. Chen, M. Fan, W. Wang, L. Yu, S. Tan, X. Chen, Q. Zhang, S. Xin, L. Wan and Y. Guo, *Angew. Chem., Int. Ed.*, 2020, **59**, 12069–12075.

42 G. V. Alexander, C. Shi, J. O'Neill and E. Wachsman, *Nat. Mater.*, 2023, **22**, 1136.

43 Y. Chen, M. He, N. Zhao, J. Fu, H. Huo, T. Zhang, Y. Li, F. Xu and X. Guo, *J. Power Sources*, 2019, **420**, 15–21.

44 H. Huo, Y. Chen, R. Li, N. Zhao, J. Luo, J. G. Pereira da Silva, R. Mücke, P. Kaghazchi, X. Guo and X. Sun, *Energy Environ. Sci.*, 2020, **13**, 127–134.

45 Z. Qin, Y. Xie, X. Meng, D. Qian, C. Shan, D. Mao, G. He, Z. Zheng, L. Wan and Y. Huang, *Chem. Eng. J.*, 2022, **447**, 137538.

46 Q. Ma, X. Zhang, A. Wang, Y. Xia, X. Liu and J. Luo, *Adv. Funct. Mater.*, 2020, **30**, 2002824.

47 J. Fu, P. Yu, N. Zhang, G. Ren, S. Zheng, W. Huang, X. Long, H. Li and X. Liu, *Energy Environ. Sci.*, 2019, **12**, 1404–1412.

48 T. Wang, J. Duan, B. Zhang, W. Luo, X. Ji, H. Xu, Y. Huang, L. Huang, Z. Song, J. Wen, C. Wang, Y. Huang and J. B. Goodenough, *Energy Environ. Sci.*, 2022, **15**, 1325–1333.

49 Z. Zhao, Z. Wen, X. Liu, H. Yang, S. Chen, C. Li, H. Lv, F. Wu, B. Wu and D. Mu, *Chem. Eng. J.*, 2021, **405**, 127031.

50 S. Gao, P. Ju, Z. Liu, L. Zhai, W. Liu, X. Zhang, Y. Zhou, C. Dong, F. Jiang and J. Sun, *J. Energy Chem.*, 2022, **69**, 356–362.

51 Z. Lu, Z. Yang, C. Li, K. Wang, J. Han, P. Tong, G. Li, B. Vishnugopi, P. Mukherjee, C. Yang and W. Li, *Adv. Energy Mater.*, 2021, **11**, 2003811.

52 X. Liu, R. Garcia-Mendez, A. Lupini, Y. Cheng, Z. Hood, F. Han, A. Sharafi, J. Idrobo, N. Dudney, C. Wang, C. Ma, J. Sakamoto and M. Chi, *Nat. Mater.*, 2021, **20**, 1485–1490.

