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# Monitoring water harvesting in metal–organic frameworks, one water molecule at a time†

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Metal–organic frameworks (MOFs) have gained prominence as potential materials for atmospheric water harvesting, a vital solution for arid regions and areas experiencing severe water shortages. However, the molecular factors influencing the performance of MOFs in capturing water from the air remain elusive. Among all MOFs,  $\text{Ni}_2\text{X}_2\text{BTDD}$  ( $\text{X} = \text{F}, \text{Cl}, \text{Br}$ ) stands out as a promising water harvester due to its ability to adsorb substantial amounts of water at low relative humidity (RH). Here, we use advanced molecular dynamics simulations carried out with the state-of-the-art MB-pol data-driven many-body potential to monitor water adsorption in the three  $\text{Ni}_2\text{X}_2\text{BTDD}$  variants as a function of RH. Our simulations reveal that the type of halide atom in the three  $\text{Ni}_2\text{X}_2\text{BTDD}$  frameworks significantly influences the corresponding molecular mechanisms of water adsorption: while water molecules form strong hydrogen bonds with the fluoride atoms in  $\text{Ni}_2\text{F}_2\text{BTDD}$ , they tend to form hydrogen bonds with the nitrogen atoms of the triazolate linkers in  $\text{Ni}_2\text{Cl}_2\text{BTDD}$  and  $\text{Ni}_2\text{Br}_2\text{BTDD}$ . Importantly, the large size of the bromide atoms reduces the void volume in the  $\text{Ni}_2\text{Br}_2\text{BTDD}$  pores, which enable water molecules to initiate an extended hydrogen-bond network at lower RH. These findings not only underscore the prospect for precisely tuning structural and chemical modifications of the frameworks to optimize their interaction with water, but also highlight the predictive power of simulations with the MB-pol data-driven many-body potential. By providing a realistic description of water under different thermodynamic conditions and environments, these simulations yield unique, molecular-level insights that can guide the design and optimization of energy-efficient water harvesting materials.

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## 1 Introduction

Lack of access to freshwater is a pressing global challenge,<sup>1</sup> with profound impacts ranging from decreased food production<sup>2</sup> to long-term public health issues.<sup>3,4</sup> With the atmosphere holding approximately  $10^{21}$  liters of water as either vapor or droplets, atmospheric water harvesting (AWH) represents an appealing solution to provide freshwater to regions grappling with water scarcity.<sup>5</sup> AWH boasts several advantages over other technologies: it is not restricted to specific geographical regions as water can be harvested from the atmosphere anywhere in the world, it eliminates the need for transportation, and it does not tap into

more limited water resources such as rivers and lakes.<sup>4,6</sup> However, for AWH to serve as a reliable source of freshwater, it necessitates sorbent materials endowed with high hydrolytic stability for enduring recycling performance, large porosity and surface area for substantial water vapor uptake, relatively mild regeneration conditions, adsorption isotherms demonstrating a sharp uptake at low relative humidity (RH), and a high deliverable capacity. Metal–organic frameworks (MOFs) hold the promise of fulfilling all these prerequisites.

Built from organic linkers and secondary building units (SBUs) composed of metal ions or clusters, MOFs are characterized by large surface areas and a wide range of structural and physical properties that can be tuned for specific applications through multivariate reticular chemistry<sup>7,8</sup> and/or post-synthetic approaches.<sup>9–11</sup> Several MOFs for potential applications in AWH technologies have been reported in the scientific literature.<sup>12–20</sup> Despite some progress, a molecular-level characterization of the water adsorption mechanisms in MOFs, which are key to the design and optimization of frameworks with high adsorption capacities at low RH, remains elusive.<sup>21–25</sup> A primary challenge stems from the intricate water–framework interface, which has so far made it difficult to link macroscopic thermodynamic properties, such as isotherms and enthalpies of adsorption, to specific interactions of individual water

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molecules with the SBUs and organic linkers of the framework as a function of RH. From a simulation perspective, these difficulties are primarily associated with intrinsic limitations of empirical force fields used to model water–framework and water–water interactions, which are unable to correctly reproduce the properties of water across different phases as a function of temperature and pressure.<sup>26</sup> Recent advances in molecular dynamics (MD) simulations of water, spurred by the advent of the MB-pol data-driven many-body potential,<sup>27–29</sup> have addressed these limitations.<sup>30</sup> This progress has paved the way for realistic computational studies, providing an unprecedented level of molecular insight into structural, thermodynamic, dynamical, and spectroscopic properties of water across the entire phase diagram, from gas-phase clusters to liquid water and ice.<sup>31–40</sup>

In this study, we employ MD simulations using the MB-pol data-driven many-body potential to characterize, at the molecular level, the individual steps associated with water adsorption in the halide-exchanged  $\text{Ni}_2\text{X}_2\text{BTDD}$  MOFs series, where X corresponds to fluoride (F), chloride (Cl), and bromide (Br) atoms. The  $\text{Ni}_2\text{X}_2\text{BTDD}$  MOFs stand out as some of the most promising candidates for AWH applications, owing to their exceptional water uptake capabilities even at low RH values.<sup>19</sup> All three  $\text{Ni}_2\text{X}_2\text{BTDD}$  MOFs share the same topology and are synthesized from the same organic linkers, specifically, bis(1*H*-1,2,3-triazolo[4,5-*b*],[4',5'-*i*']dibenzo[1,4]dioxin, or  $\text{H}_2\text{BTDD}$ . They also share the same SBUs with undercoordinated nickel ( $\text{Ni}^{2+}$ ) centers. Their distinction lies in the halide atoms (F, Cl, and Br) that bridge the  $\text{Ni}^{2+}$  centers. While the three  $\text{Ni}_2\text{X}_2\text{BTDD}$  frameworks have similar pore sizes, ranging between 22 Å and 23 Å, their different halide functionalization leads to notable variations in water adsorption capacities and framework stabilities. Specifically,  $\text{Ni}_2\text{F}_2\text{BTDD}$  and  $\text{Ni}_2\text{Cl}_2\text{BTDD}$  can adsorb over 100% of their weight in water at 32% RH, while  $\text{Ni}_2\text{Br}_2\text{BTDD}$  can only adsorb 64% of its weight in water, but this occurs below 25% RH, which is one of the lowest RH values for AWH reported for MOFs to date.<sup>19</sup> Furthermore,  $\text{Ni}_2\text{Cl}_2\text{BTDD}$  and  $\text{Ni}_2\text{Br}_2\text{BTDD}$  maintain their water uptake capacity for more than 400 adsorption/desorption cycles, while  $\text{Ni}_2\text{F}_2\text{BTDD}$  undergoes partial amorphization and displays a decline in its adsorption capacity after one cycle.<sup>19</sup> Given their structural similarities and varying water adsorption capacities, the three  $\text{Ni}_2\text{X}_2\text{BTDD}$  MOFs present an ideal platform to identify specific correlations between water–framework and water–water interactions and characterize the thermodynamics of water adsorption as a function of RH at the molecular level.

## 2 Methods

### 2.1 Force fields and potential energy functions

Each MOF structure was first optimized using periodic density functional theory (DFT) calculations carried out with the Vienna *Ab initio* Simulation Package (VASP),<sup>41–44</sup> using the PBE exchange–correlation functional<sup>45</sup> with Grimme's D3 dispersion correction<sup>46</sup> in combination with a projector-augmented wave (PAW) treatment<sup>47,48</sup> with a 700 eV kinetic energy cutoff. A  $2 \times 2 \times 4$  *k*-point grid was used, and forces were converged to

a tolerance of 0.05 eV Å<sup>−1</sup>. Following periodic optimization, a cluster surrounding the  $\text{Ni}^{2+}$  metal center was extracted for each MOF (see ESI† for details). All bonded parameters involving the  $\text{Ni}^{2+}$  metal center were fit using a genetic algorithm, and all bonded parameters for the organic linker were taken from the General Amber Force Field (GAFF).<sup>49</sup> Atomic point charges were calculated for the clusters using the charge model 5 (CM5) charge scheme.<sup>50</sup> The Lennard-Jones (LJ) parameters were taken from the Universal Force Field (UFF)<sup>51</sup> for the Ni, F, Cl, and Br atoms, and the LJ parameters for the organic linker atoms (H, C, N, O) were taken from GAFF.<sup>49</sup> Force field parameters are given in Tables S1–S15 in the ESI.†

The MB-pol data-driven many-body potential was used to simulate water.<sup>27–29</sup> MB-pol has been shown to correctly predict the properties of water from the gas to the condensed phase for structural, dynamical, and thermodynamic properties.<sup>31</sup> The many-body dipole moment (MB- $\mu$ ) was used to calculate the dipole moment of water for the IR spectra.<sup>33</sup> Local mode and local monomer (LM) calculations<sup>52–54</sup> from ref. 55 were utilized to add the Fermi resonance contribution to the IR spectra.

Nonbonded interactions between water and the MOF were represented through electrostatic and van der Waals (VDW) interactions. The Ni–OW, X–OW (X = F, Cl, Br), and X–HW interactions were fit to a Buckingham potential using a genetic algorithm (see ESI† for details). Cross interactions between water and the remaining MOF atoms were described through Lennard-Jones (LJ) parameters. The LJ parameters were derived from Lorentz–Berthelot mixing rules using the LJ parameters of the TIP4P/2005 water model,<sup>56</sup> which is the closest point-charge model to MB-pol.<sup>57</sup>

### 2.2 Molecular dynamics simulations

Classical molecular dynamics (MD) simulations were performed using an in-house code based on the DL\_POLY\_2 simulation package,<sup>58</sup> which was modified to include MB-pol water.<sup>27–29</sup> Simulations for each halide MOF were carried out for a system containing  $1 \times 1 \times 3$  unit cells under periodic boundary conditions. The initial configurations for each MD simulation of the different halide MOFs were generated from Packmol.<sup>59</sup> Each configuration started with 54 water molecules ( $1\text{H}_2\text{O}/\text{Ni}^{2+}$ ), and 54 water molecules were added to each subsequent simulation ( $+1\text{H}_2\text{O}/\text{Ni}^{2+}$ ) up to the point of pore saturation on the adsorption isotherm ( $12\text{H}_2\text{O}/\text{Ni}^{2+}$  for  $\text{Ni}_2\text{F}_2\text{BTDD}$ ,  $13\text{H}_2\text{O}/\text{Ni}^{2+}$  for  $\text{Ni}_2\text{Cl}_2\text{BTDD}$ , and  $11\text{H}_2\text{O}/\text{Ni}^{2+}$  for  $\text{Ni}_2\text{Br}_2\text{BTDD}$ ).<sup>19</sup> MD simulations were first performed utilizing a constant number of molecules, pressure, and temperature in the isobaric–isothermal ensemble ( $NPT$  = constant number of atoms, pressure, and temperature). The system was heated to 1000 K for 10 ps at 1.0 atm, cooled to 500 K for 20 ps at 1.0 atm, and cooled to 300 K for 20 ps at 1.0 atm, all in the  $NPT$  ensemble, to randomize the initial distribution of water molecules. MD simulations were then performed for 1 ns at 300 K and 1.0 atm in the  $NPT$  ensemble to equilibrate the system, calculate structural information, and obtain the average box dimensions. The average box dimensions listed in Tables S16–S18 in the ESI† were used for all subsequent simulations at each



loading. The angles of each simulation box were 90°, 90°, and 120° along the *x*, *y*, and *z* dimensions, respectively. Following the initial equilibration, the system was simulated with a constant number of molecules, volume, and temperature in the canonical ensemble ( $NVT$  = constant number of atoms, volume, and temperature) for 100 ps at 300 K followed by a constant number of molecules, volume, and energy in the microcanonical ensemble ( $NVE$  = constant number of atoms, volume, and energy) for 50 ps to obtain dynamical information, including the IR spectra, where the temperature remained constant around 300 K. Subsequent configurations were heated to 1000 K for 10 ps in the  $NVT$  ensemble to randomize the distribution of water, cooled to 300 K for 10 ps in the  $NVT$  ensemble, and then run for 50 ps in the  $NVE$  ensemble where the temperature remained constant around 300 K. 19 of these additional trajectories were run at each loading for each halide, giving a total of 20 independent MD simulations in which all dynamical properties were averaged over. The entropy was calculated as a function of RH using the two-phase thermodynamic (2PT) model.<sup>60</sup> While approximate like any other molecular model of entropy, the 2PT model provides an intuitive approach to calculating the entropy of water depending on the local environment. Using the 2PT model, MB-pol predicts a value of  $71.9 \pm 0.6 \text{ J mol}^{-1} \text{ K}^{-1}$  for the entropy of liquid water at ambient conditions, which is in good agreement with the experimental value of  $69.95 \pm 0.03 \text{ J mol}^{-1} \text{ K}^{-1}$ .

All MD simulations were performed with a time step of 0.2 fs, and the equations of motion were propagated according to the velocity-Verlet algorithm. The temperature was maintained at 300 K by a Nosé–Hoover chain of four thermostats.<sup>61</sup> The  $NPT$  ensemble was generated according to the algorithm described in ref. 62. Short-range interactions were truncated at an atom–atom distance of 9.0 Å, and the electrostatics were calculated using the Ewald sum.<sup>63</sup>

The molecular models used in the MD simulations, which were conducted using in-house software based on the DL\_POLY\_2 simulation package, are available upon request from the authors in a format compatible with the MBX interface<sup>64,65</sup> for LAMMPS.<sup>66</sup>

## 3 Results

### 3.1 Thermodynamics of water adsorption

Fig. 1 presents the experimental adsorption isotherms for the three  $\text{Ni}_2\text{X}_2\text{BTDD}$  MOFs measured as a function of RH in ref. 19,

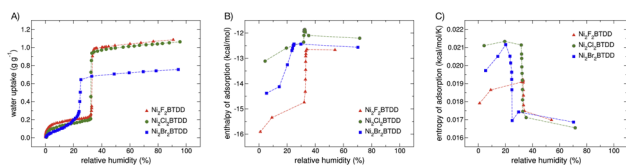


Fig. 1 (A) Experimental adsorption isotherms of the three  $\text{Ni}_2\text{X}_2\text{BTDD}$  MOFs reported in ref. 19. (B and C) Adsorption enthalpy and entropy, respectively, calculated for water in the  $\text{Ni}_2\text{X}_2\text{BTDD}$  MOFs as a function of relative humidity from MD simulations carried out with the MB-pol data-driven many-body potential of water.

along with the theoretical values of the adsorption enthalpy and entropy calculated for water in the  $\text{Ni}_2\text{X}_2\text{BTDD}$  from MD simulations carried out with the MB-pol data-driven many-body potential. The enthalpy of adsorption directly reports on the strength and interplay of water–framework and water–water interactions. The balance between these interactions is crucial for determining the onset RH values for water adsorption and pore filling, which in turn dictates the overall performance of a MOF for AWH applications.  $\text{Ni}_2\text{F}_2\text{BTDD}$  displays the most negative enthalpy of adsorption at all RH values, which directly correlates with fluoride providing the strongest interactions with water among the halide atoms.<sup>67–70</sup> The enthalpy of adsorption of  $\text{Ni}_2\text{Cl}_2\text{BTDD}$  and  $\text{Ni}_2\text{Br}_2\text{BTDD}$  does not, however, follow the order of chloride–water and bromide–water interactions, suggesting differences in the modes of adsorption between the two frameworks. The differences in the physico-chemical properties of the fluoride, chloride, and bromide ions, their interactions with water, and adsorption modes manifest in distinct pore-filling onsets (Fig. 1A). Specifically,  $\text{Ni}_2\text{Br}_2\text{BTDD}$  fills its pores at a lower RH value (24% RH), compared to  $\text{Ni}_2\text{F}_2\text{BTDD}$  and  $\text{Ni}_2\text{Cl}_2\text{BTDD}$  (32% RH), which is mirrored by steep increases calculated for the corresponding enthalpies of adsorption at the same RH values (Fig. 1B). The increase in enthalpy of adsorption predicted by the MD simulations as the RH increases indicates that water–water interactions, which become the dominant interactions at high RH values, are relatively weaker than water–framework interactions. The predicted larger enthalpy of adsorption for  $\text{Ni}_2\text{F}_2\text{BTDD}$ , indicative of stronger water–framework interactions, aligns with its observed partial amorphization and reduced water adsorption capacity in successive adsorption/desorption cycles.<sup>19</sup>

The entropy of water calculated from the MD simulations displays the opposite trend compared to the enthalpy of adsorption, with water in  $\text{Ni}_2\text{Cl}_2\text{BTDD}$  and  $\text{Ni}_2\text{F}_2\text{BTDD}$  having the highest and lowest entropy up to the pore-filling step, respectively (Fig. 1C). Since at low RH values, water in the MOF pores exhibit gas-like behavior modulated by water–framework interactions, the entropy of water adsorbed in all three  $\text{Ni}_2\text{X}_2\text{BTDD}$  MOFs remains appreciably larger than the entropy of liquid water up to the pore-filling step.<sup>23,71</sup> At the RH values corresponding to the pore-filling steps, the entropy of adsorption displays a sudden drop for all three MOFs and approaches the value of liquid water as the RH increases. The variation of the entropy of adsorption as a function of RH effectively mirrors the variation of the enthalpy of adsorption, indicating that as the pores get filled, the adsorbed water molecules acquire bulk-like behavior.

### 3.2 Early stages of water adsorption

The OH vibrations of water molecules are especially sensitive to their surrounding environment.<sup>72</sup> Monitoring the evolution of the infrared (IR) spectra of water adsorbed in the three  $\text{Ni}_2\text{X}_2\text{BTDD}$  MOFs in the OH-stretch frequency range thus provides a direct probe into the development of hydrogen-bonded networks inside the pores and offers molecular-level insights

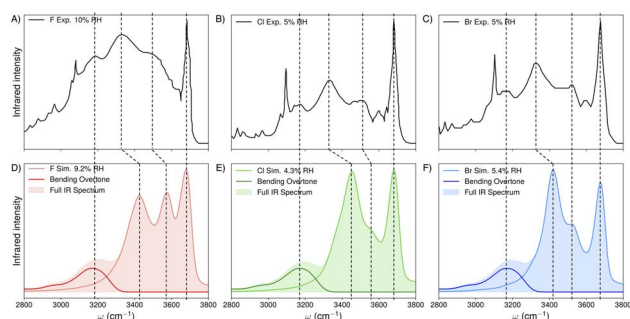


into the adsorption mechanisms. At low RH, the experimental IR spectra (Fig. 2) display four distinct features at  $\sim 3700\text{ cm}^{-1}$ ,  $\sim 3500\text{ cm}^{-1}$ ,  $\sim 3300\text{ cm}^{-1}$ , and  $\sim 3200\text{ cm}^{-1}$ . This contrasts with the broad, featureless OH-stretch band seen in the IR spectrum of liquid water.<sup>33</sup> Using the simulated IR spectra for comparison, the highest frequency peak at  $\sim 3700\text{ cm}^{-1}$  can be attributed to stretch vibrations of OH bonds not involved in hydrogen bonding. The position and width of this high-frequency peak are reminiscent of the free-OH peak in the sum-frequency generation spectrum of the air/water interface.<sup>35</sup> Our theoretical analysis suggests that the feature at  $\sim 3200\text{ cm}^{-1}$ , often linked to ice-like hydrogen-bonded structures, actually arises from Fermi resonances.<sup>73</sup> These resonances result from couplings between the OH-stretch fundamental vibrations and the HOH bending overtones.<sup>55</sup> The other two peaks indicate distinct local hydrogen-bonding environments that develop inside the MOF pores as water adsorption increases with RH. The discrepancies between the experimental and calculated IR spectra observed in Fig. 2 can be attributed to various factors. These include the omission of higher-body framework-water interactions and nuclear quantum effects in the MD simulations. Previous studies have shown that not accounting for nuclear quantum effects in MD simulations of vibrational spectra of water generally leads to spectral features that are narrower and exhibit a blueshift when compared to experimental results.<sup>74</sup>

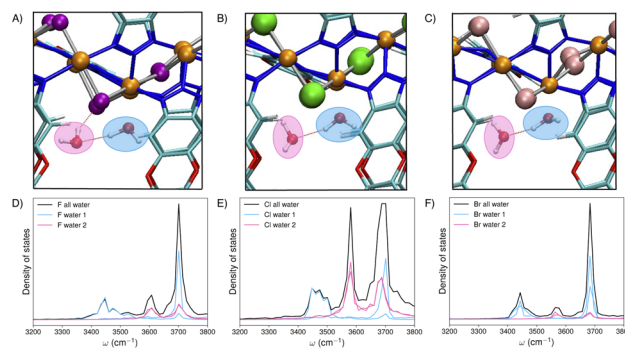
MD simulations, performed by incrementally adding water molecules, reveal that the initial water molecules entering the pores of the three  $\text{Ni}_2\text{X}_2\text{BTDD}$  MOFs bind to one of the open

$\text{Ni}^{2+}$  metal centers, as depicted by the water molecule inside the blue area in Fig. 3A–C. The spatial distributions of these Ni-bound water molecules, however, vary depending on the specific halide ion present in the framework, giving rise to characteristic peaks in the corresponding IR spectra (Fig. 3D–F).

Fig. 4A–C shows that while the Ni-bound water molecules localize in a rounded three-dimensional region perpendicular to the  $\text{Ni}^{2+}$  metal center in  $\text{Ni}_2\text{Cl}_2\text{BTDD}$  and  $\text{Ni}_2\text{Br}_2\text{BTDD}$ , they are more delocalized along the F–Ni–F axis in  $\text{Ni}_2\text{F}_2\text{BTDD}$ . These differences become more apparent from the analyses of the radial distribution functions (RDFs) describing spatial correlations between the oxygen (OW) and hydrogen (HW) atoms of the Ni-bound water molecule and the  $\text{Ni}^{2+}$ , halide, and nitrogen ( $n = n1$  and  $n2$ ) atoms of the framework (see Fig. S1, S2 and Tables S1–S15 in the ESI† for the complete lists of atom labels and force field parameters). Specifically, the OW–Ni RDFs in Fig. 4D show that the average distance of the Ni-bound water molecules from the  $\text{Ni}^{2+}$  open sites is effectively the same in all three MOFs. In contrast, the HW–X RDFs in Fig. 4E show that the average distance between the H atoms of the Ni-bound water molecules and the halide atoms of the frameworks is significantly different among the three MOFs, indicating that the Ni-bound water molecules establish hydrogen bonds with the fluoride atoms in  $\text{Ni}_2\text{F}_2\text{BTDD}$  but not with the chloride and bromide atoms in  $\text{Ni}_2\text{Cl}_2\text{BTDD}$  and  $\text{Ni}_2\text{Br}_2\text{BTDD}$ . This different spatial arrangement can be explained by considering the larger electronegativity and smaller size of fluoride compared to chloride and bromide, which lead to significantly stronger fluoride–water interactions.<sup>67</sup> The trend is reversed in the HW–n RDFs shown in Fig. 4F, with the H atoms of the Ni-bound water molecules approaching the N atoms of the triazolate linkers more closely



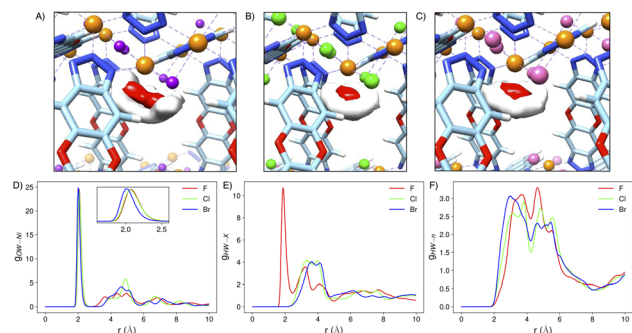
**Fig. 2** Experimental IR spectra of water from ref. 19 for  $\text{Ni}_2\text{F}_2\text{BTDD}$  at 10% RH (A),  $\text{Ni}_2\text{Cl}_2\text{BTDD}$  at 5% RH (B), and  $\text{Ni}_2\text{Br}_2\text{BTDD}$  at 5% RH (C). (D) IR spectrum of water in  $\text{Ni}_2\text{F}_2\text{BTDD}$  at 9.2% RH (light red shaded area) obtained by combining the IR spectrum calculated from the MD simulations (light red solid curve) with the Fermi resonance contributions (dark red solid curve). (E) IR spectrum of water in  $\text{Ni}_2\text{Cl}_2\text{BTDD}$  at 4.3% RH (light green shaded area) obtained by combining the IR spectrum calculated from the MD simulations (light green solid curve) with the Fermi resonance contributions (dark green solid curve). (F) IR spectrum of water in  $\text{Ni}_2\text{Br}_2\text{BTDD}$  at 5.4% RH (light blue shaded area) obtained by combining the IR spectrum calculated from the MD simulations (light blue solid curve) with the Fermi resonance contributions (dark blue solid curve). The simulated IR spectra red-shifted by  $215\text{ cm}^{-1}$  ( $\text{Ni}_2\text{F}_2\text{BTDD}$  and  $\text{Ni}_2\text{Cl}_2\text{BTDD}$ ) or  $219\text{ cm}^{-1}$  ( $\text{Ni}_2\text{Br}_2\text{BTDD}$ ) to align the position of the free-OH peak with the corresponding experimental peak, and the HOH bending overtones are blue-shifted by  $10\text{ cm}^{-1}$ , accordingly. The main features of the experimental and simulated IR spectra are indicated by dashed lines.



**Fig. 3** Snapshots of MD simulations carried out with the MB-pol data-driven potential for two water molecules in (A)  $\text{Ni}_2\text{F}_2\text{BTDD}$ , (B)  $\text{Ni}_2\text{Cl}_2\text{BTDD}$ , and (C)  $\text{Ni}_2\text{Br}_2\text{BTDD}$ . The first water molecule interacting with the  $\text{Ni}^{2+}$  metal center in each MOF is highlighted in blue, and the second water molecule H-bonded to the first is highlighted in pink. Color scheme of atoms in panels A through C: Ni = orange, F = purple, Cl = green, Br = pink, N = blue, C = cyan, O = red, H = white. Hydrogen bonds are shown as thin dashed red lines. (D–F) Vibrational densities of states calculated for two water molecules in  $\text{Ni}_2\text{F}_2\text{BTDD}$ ,  $\text{Ni}_2\text{Cl}_2\text{BTDD}$ , and  $\text{Ni}_2\text{Br}_2\text{BTDD}$ , respectively. The colors of each peak correspond to the same colors of the water molecules in panels A through C. The total density of states is shown as solid black curve in each panel.







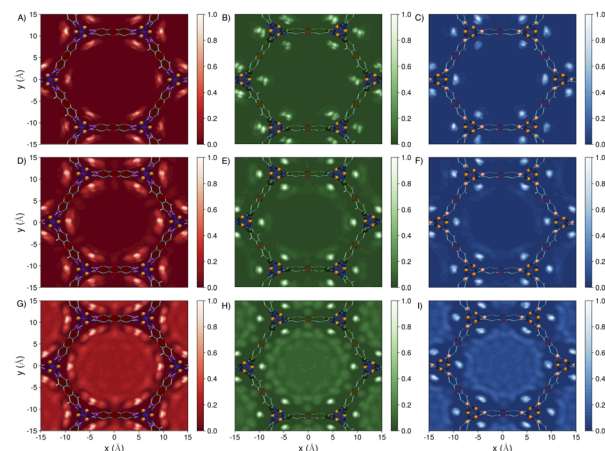
**Fig. 4** (A–C) Spatial distributions of one water molecule adsorbed in Ni<sub>2</sub>F<sub>2</sub>BTDD, Ni<sub>2</sub>Cl<sub>2</sub>BTDD, and Ni<sub>2</sub>Br<sub>2</sub>BTDD, respectively. Red iso-surfaces correspond to water oxygen atoms, and white isosurfaces correspond to water hydrogen atoms. Color scheme of atoms in panels A through C: Ni = orange, F = purple, Cl = green, Br = pink, N = blue, C = cyan, O = red, H = white. (D–F) Radial distribution functions (RDFs) describing spatial correlations for the OW–Ni (the first peak of the RDF is magnified in the inset), HW–X, and HW–n atom pairs, respectively, calculated for water in Ni<sub>2</sub>F<sub>2</sub>BTDD (red curves), Ni<sub>2</sub>Cl<sub>2</sub>BTDD (green curves), and Ni<sub>2</sub>Br<sub>2</sub>BTDD (blue curves). Refer to Fig. S1 and S2 in the ESI† for the complete list of atom labels.

in Ni<sub>2</sub>Cl<sub>2</sub>BTDD and Ni<sub>2</sub>Br<sub>2</sub>BTDD than in Ni<sub>2</sub>F<sub>2</sub>BTDD. The Ni-bound water molecules act as anchor points, initiating the formation of extended hydrogen-bonding networks by donating a hydrogen bond to a second water molecule, hereafter referred to as the H-bonded molecule (shown inside the pink area in Fig. 3A–C).

Analysis of the MD trajectories and associated vibrational density of states (Fig. 3D–F) reveals that the Ni-bound water molecules in all three MOFs exhibit a free-OH, contributing to the peak at  $\sim 3700\text{ cm}^{-1}$ . The second OH bond of the Ni-bound water molecules instead engages in hydrogen bonding with another water molecule, contributing to the peak at  $\sim 3450\text{ cm}^{-1}$ . The H-bonded water molecules also exhibit free-OH bonds that point toward the center of the pore and contribute to the peak at  $3700\text{ cm}^{-1}$ . The second OH bond of the H-bonded molecule instead interacts with the framework, forming weak hydrogen bonds with neighboring fluoride atoms in Ni<sub>2</sub>F<sub>2</sub>BTDD, and triazolate groups in Ni<sub>2</sub>Cl<sub>2</sub>BTDD and Ni<sub>2</sub>Br<sub>2</sub>BTDD, contributing to the peak at  $\sim 3550\text{--}3600\text{ cm}^{-1}$ . The IR spectra of water adsorbed in the three MOFs at low RH can thus be assigned to hydrogen-bonded water dimers that are anchored to the open Ni<sup>2+</sup> centers. As the RH increases, the four distinct peaks grow into one broad and featureless band similar to that observed in the IR spectrum of liquid water, indicating the emergence of a continuum of hydrogen-bond topologies (Fig. S4 in the ESI†).

### 3.3 Pore filling

At a loading of  $1\text{H}_2\text{O}/\text{Ni}^{2+}$ , which corresponds to  $\sim 1\text{--}5\%$  RH depending on the specific Ni<sub>2</sub>X<sub>2</sub>BTDD MOF, Fig. 5A–C show that the water molecules inside the pores of all three MOFs localize near the open Ni<sup>2+</sup> sites, following the same structural arrangements identified in Fig. 4A–C, with the water density in Ni<sub>2</sub>F<sub>2</sub>BTDD being more elongated due to relatively stronger interactions with the fluoride atoms of the



**Fig. 5** Two-dimensional density distributions of water molecules in Ni<sub>2</sub>F<sub>2</sub>BTDD (A, D, and G, red shades), Ni<sub>2</sub>Cl<sub>2</sub>BTDD (B, E, and H, green shades), and Ni<sub>2</sub>Br<sub>2</sub>BTDD (C, F, and I, blue shades). The top row shows the water distributions at  $1\text{H}_2\text{O}/\text{Ni}^{2+}$  corresponding to 0.96% RH for Ni<sub>2</sub>F<sub>2</sub>BTDD (A), 4.3% RH for Ni<sub>2</sub>Cl<sub>2</sub>BTDD (B), 5.4% RH for Ni<sub>2</sub>Br<sub>2</sub>BTDD (C). The middle row shows the water distributions at  $3\text{H}_2\text{O}/\text{Ni}^{2+}$  corresponding to 32.1% RH for Ni<sub>2</sub>F<sub>2</sub>BTDD (D), 31.7% RH for Ni<sub>2</sub>Cl<sub>2</sub>BTDD (E), 20.2% RH for Ni<sub>2</sub>Br<sub>2</sub>BTDD (F). The bottom row shows the water distributions at the highest loading simulated for each MOF:  $12\text{H}_2\text{O}/\text{Ni}^{2+}$  corresponding to 54.0% RH for Ni<sub>2</sub>F<sub>2</sub>BTDD (G),  $13\text{H}_2\text{O}/\text{Ni}^{2+}$  corresponding to 71.5% RH for Ni<sub>2</sub>Cl<sub>2</sub>BTDD (H),  $11\text{H}_2\text{O}/\text{Ni}^{2+}$  corresponding to 70.3% RH for Ni<sub>2</sub>Br<sub>2</sub>BTDD (I). Lighter shades correspond to regions of higher water density.

framework. As the RH increases to  $3\text{H}_2\text{O}/\text{Ni}^{2+}$  (Fig. 5D–F), the water molecules localized near different Ni<sup>2+</sup> sites begin to connect to each other through the formation of bridging hydrogen bonds, as revealed by the increased water density near the organic linkers, which is particularly evident in Ni<sub>2</sub>F<sub>2</sub>BTDD and Ni<sub>2</sub>Cl<sub>2</sub>BTDD. It is important to note that a loading of  $3\text{H}_2\text{O}/\text{Ni}^{2+}$  corresponds to different RH values in the three MOFs due to differences in the thermodynamics of the adsorption process, as manifested in the different trends in the corresponding isotherms, enthalpies, and entropies of adsorption. Specifically, a loading of  $3\text{H}_2\text{O}/\text{Ni}^{2+}$  is equivalent to 32.1% RH in Ni<sub>2</sub>F<sub>2</sub>BTDD and 31.7% RH in Ni<sub>2</sub>Cl<sub>2</sub>BTDD, which occur in correspondence of the pore filling step in the adsorption isotherms of both MOFs. In contrast, a loading of  $3\text{H}_2\text{O}/\text{Ni}^{2+}$  is below the onset of pore filling in Ni<sub>2</sub>Br<sub>2</sub>BTDD, corresponding to 24% RH. It is thus apparent that pore filling in the three Ni<sub>2</sub>X<sub>2</sub>BTDD MOFs occurs after the water molecules are able to form connected hydrogen-bonded chains that elongate along the pore channels and provide additional hydrogen-bonding sites. The different RH values needed for pore filling in the three MOFs is due to the different sizes of the fluoride, chloride, and bromide atoms that decorate the pore surface of the three frameworks, as originally suggested in ref. 19. In particular, due to their larger sizes, the bromide atoms of Ni<sub>2</sub>Br<sub>2</sub>BTDD protrude into the pore more than the fluoride and chloride atoms in the analogous MOFs. This implies that a relatively larger number of water molecules (*i.e.*, higher RH) is needed to encircle the bromide atoms in order to form fully connected hydrogen-bonded chains that

can then bridge with water molecules localized near the adjacent  $\text{Ni}^{2+}$  sites (Fig. 5D–F).

Fig. 5G–I show well-defined density patterns of water inside the pores at the maximum water loading in each of the three MOFs. Due to stronger fluoride–water interactions, the first layer of water molecules in  $\text{Ni}_2\text{F}_2\text{BTDD}$  is able to template the hydrogen-bonding structure of the additional water molecules down to the center of the pore. This templating effect becomes progressively weaker from  $\text{Ni}_2\text{F}_2\text{BTDD}$  to  $\text{Ni}_2\text{Br}_2\text{BTDD}$ , following the same trend as the underlying halide–water interactions.<sup>68</sup> Importantly, while water in  $\text{Ni}_2\text{Br}_2\text{BTDD}$  has the highest density near the  $\text{Ni}^{2+}$  sites, areas of little to no water density also exist right next to these regions. These areas of low water density are a direct consequence of the bromide atoms protruding further into the pores than the other two halide atoms, effectively hindering the formation of well-connected hydrogen-bonding networks that manifest in a relatively more disordered distribution of water molecules inside the pores (Fig. 5I).

Fig. 6 shows the distributions of different hydrogen-bonding topologies, classified in terms of the average number of hydrogen bonds that each water molecule donates ( $n_D$ ) and accepts ( $m_A$ ), determined from MD simulations at the maximum water loading in each of the three  $\text{Ni}_2\text{X}_2\text{BTDD}$  MOFs. In  $\text{Ni}_2\text{F}_2\text{BTDD}$ , a larger fraction of water molecules donate a hydrogen bond to the fluoride atoms of the framework, while still accepting hydrogen bonds from the neighboring water molecules. On the other hand, water molecules in  $\text{Ni}_2\text{Cl}_2\text{BTDD}$  and  $\text{Ni}_2\text{Br}_2\text{BTDD}$ , on average, donate a relatively larger number of hydrogen bonds to the N atoms of the triazolate linkers. The lifetime of the  $\text{H}_2\text{O}-\text{X}$  and  $\text{H}_2\text{O}-\text{N}$  hydrogen bonds also vary, reflecting the strength of the underlying interactions. Specifically, the hydrogen-bond correlation functions, which directly report on the probability that a hydrogen bond remains intact for a given time,<sup>75</sup> indicate that water has the longest hydrogen-bond lifetime in  $\text{Ni}_2\text{Br}_2\text{BTDD}$  when interacting with the N atom (Table 1 and Fig. S9 in the ESI†). Long hydrogen-bond lifetimes are also calculated for the hydrogen bonds established by the

**Table 1** Hydrogen-bonding timescales of water in  $\text{Ni}_2\text{X}_2\text{BTDD}$ . Life-time ( $\tau_{\text{HB}}$ ) of hydrogen bonds established by water molecules with distinct atoms of the  $\text{Ni}_2\text{X}_2\text{BTDD}$  MOFs.  $\tau_{\text{HB}}$  was calculated by fitting the corresponding hydrogen-bond time correlation function fit to the exponential  $c(t) = Ae^{-t/\tau_{\text{HB}}}$  according to the procedure described in ref. 75

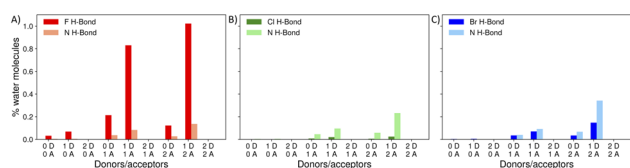
Hydrogen-bond type	$\tau_{\text{HB}}$ (ps)
$\text{H}_2\text{O}-\text{F}$ in $\text{Ni}_2\text{F}_2\text{BTDD}$	77
$\text{H}_2\text{O}-\text{N}$ in $\text{Ni}_2\text{F}_2\text{BTDD}$	20
$\text{H}_2\text{O}-\text{Cl}$ in $\text{Ni}_2\text{Cl}_2\text{BTDD}$	12
$\text{H}_2\text{O}-\text{N}$ in $\text{Ni}_2\text{Cl}_2\text{BTDD}$	78
$\text{H}_2\text{O}-\text{Br}$ in $\text{Ni}_2\text{Br}_2\text{BTDD}$	24
$\text{H}_2\text{O}-\text{N}$ in $\text{Ni}_2\text{Br}_2\text{BTDD}$	>100

water molecules with the fluoride atoms in  $\text{Ni}_2\text{F}_2\text{BTDD}$  and the N atoms in  $\text{Ni}_2\text{Cl}_2\text{BTDD}$ . The differences in hydrogen-bond lifetimes are rationalized by considering that the larger size of the bromide ions effectively push the water molecules toward the framework where they establish strong hydrogen bonds with the N atoms of the triazolate linkers. The distinct hydrogen-bonding arrangements developed by the water molecules in the pores of the three MOFs also explain the larger enthalpy of adsorption of water exhibited by  $\text{Ni}_2\text{Br}_2\text{BTDD}$  compared to  $\text{Ni}_2\text{Cl}_2\text{BTDD}$ .

## 4 Conclusions

Exchanging the halide atoms in  $\text{Ni}_2\text{X}_2\text{BTDD}$  leads to three MOFs with different properties and water adsorption capacities.  $\text{Ni}_2\text{F}_2\text{BTDD}$  and  $\text{Ni}_2\text{Cl}_2\text{BTDD}$  adsorb over 100% of their weight in water at 32% RH, while  $\text{Ni}_2\text{Br}_2\text{BTDD}$  adsorbs 64% of its weight in water below 25% RH, one of the best water uptake values reported to date.<sup>19</sup> By providing a realistic representation of water adsorption in all three MOFs, our MD simulations carried out with the MB-pol data-driven many-body potential provide unique, molecular-level insights into the underlying interactions and thermodynamic driving forces.

Theoretical analyses of the simulated IR spectra allows for monitoring the evolution of extended hydrogen-bond networks as a function of RH. At low RH values, the main spectral features are assigned to free OH stretches, water–framework hydrogen bonds, water–water hydrogen bonds, and Fermi resonances, while a continuum of hydrogen-bonding environments exist at higher RH values, resembling the hydrogen-bond network of liquid water. All three MOFs exhibit higher enthalpies of water adsorption at low RH values, indicating stronger water–framework interactions than water–water interactions. Furthermore,  $\text{Ni}_2\text{F}_2\text{BTDD}$  exhibits the largest enthalpy of adsorption at all RH values, which directly correlates with relatively stronger fluoride–water interactions. These strong interactions are likely responsible for the partial amorphization and decrease in water uptake capacity of  $\text{Ni}_2\text{F}_2\text{BTDD}$  after the first adsorption cycle. On the other hand, the distinct physico-chemical properties of the chloride and bromide atoms play a central role in modulating water adsorption in  $\text{Ni}_2\text{Cl}_2\text{BTDD}$  and  $\text{Ni}_2\text{Br}_2\text{BTDD}$ . Specifically, the water adsorption capacity is greater in



**Fig. 6** Hydrogen-bonding topologies of water molecules donating one hydrogen bond to an atom of the framework calculated from MD simulations carried out with  $12\text{H}_2\text{O}/\text{Ni}^{2+}$  for  $\text{Ni}_2\text{F}_2\text{BTDD}$  (A),  $13\text{H}_2\text{O}/\text{Ni}^{2+}$  for  $\text{Ni}_2\text{Cl}_2\text{BTDD}$  (B), and  $11\text{H}_2\text{O}/\text{Ni}^{2+}$  for  $\text{Ni}_2\text{Br}_2\text{BTDD}$  (C). A hydrogen bond is defined as having a distance  $R_{\text{D}-\text{A}} \leq 3.5$  Å between donor (D) and acceptor (A) atoms and an angle  $\vartheta_{\text{DHA}} \leq 30^\circ$ . Darker colors represent the number of donors and acceptors when water forms a hydrogen bond with a halide atom (F, Cl, or Br) of the framework. Lighter colors represent the number of donors and acceptors when water forms a hydrogen bond with a nitrogen atom (N) of the framework. The number of hydrogen-bond donors and acceptors are shown as a percent of the total number of water molecules in each MOF.



Ni<sub>2</sub>Cl<sub>2</sub>BTDD than Ni<sub>2</sub>Br<sub>2</sub>BTDD because of the smaller atomic size of Cl, resulting in a larger void space within the pores. At low RH values, the water molecules saturate the under-coordinated Ni<sup>2+</sup> centers of both Ni<sub>2</sub>Cl<sub>2</sub>BTDD and Ni<sub>2</sub>Br<sub>2</sub>BTDD frameworks before establishing hydrogen bonds with the N atoms of the triazolate linkers. Due to the large size of the bromide ions, the water molecules in Ni<sub>2</sub>Br<sub>2</sub>BTDD protrude into the pore to the greatest extent, which favors the development of extended hydrogen-bonding networks and the onset of pore filling at the lowest RH value compared to Ni<sub>2</sub>F<sub>2</sub>BTDD and Ni<sub>2</sub>Cl<sub>2</sub>BTDD. Overall, Ni<sub>2</sub>Cl<sub>2</sub>BTDD exhibits the lowest enthalpy of adsorption, which is a critical parameter for AWH applications determining the energy input required to desorb water from the MOF pores. Importantly, the entropy of adsorption calculated for water in the three MOFs exhibits the opposite trend to the enthalpy of adsorption, with higher values for water in Ni<sub>2</sub>Cl<sub>2</sub>BTDD at low RH and a progressive decrease toward the value for liquid water at high RH. The distinct trends in thermodynamic properties, spatial arrangements, and mobility of the water molecules in the three Ni<sub>2</sub>X<sub>2</sub>BTDD MOFs directly correlate with the lifetime of the underlying hydrogen bonds, which are particularly long for water hydrogen-bonded to the N atoms of the triazolate linkers in Ni<sub>2</sub>Br<sub>2</sub>BTDD.

Besides providing molecular-level insights into how structural and chemical modifications of topologically identical frameworks modulate the water adsorption capacities of MOFs, the in-depth view into the structure, thermodynamics, and dynamics of water confined in Ni<sub>2</sub>X<sub>2</sub>BTDD gained from MD simulations with the MB-pol data-driven many-body potential highlight the key role of computer simulations with realistic models in guiding the design and optimization of energy-efficient water harvesting materials.

## Data availability

All molecular dynamics trajectories and computer codes used in the analyses presented in this study are available from the authors upon request.

## Author contributions

F. P. conceived the research. K. M. H. and F. P. designed the research. K. M. H. performed the simulations. K. M. H. and F. P. analyzed and discussed the results, and wrote the paper. F. P. acquired funding and administered the project.

## Conflicts of interest

There are no conflicts to declare.

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## Notes and references

- 1 M. M. Mekonnen and A. Y. Hoekstra, *Sci. Adv.*, 2016, **2**, e1500323.
- 2 A. Dinar, A. Tieu and H. Huynh, *Global Food Secur.*, 2019, **23**, 212–226.
- 3 M. Salehi, *Environ. Int.*, 2022, **158**, 106936.
- 4 Y. Tu, R. Wang, Y. Zhang and J. Wang, *Joule*, 2018, **2**, 1452–1475.
- 5 K. Yang, T. Pan, Q. Lei, X. Dong, Q. Cheng and Y. Han, *Environ. Sci. Technol.*, 2021, **55**, 6542–6560.
- 6 X. Zhou, H. Lu, F. Zhao and G. Yu, *ACS Mater. Lett.*, 2020, **2**, 671–684.
- 7 Z. Ji, T. Li and O. M. Yaghi, *Science*, 2020, **369**, 674–680.
- 8 S. Canossa, Z. Ji, C. Gropp, Z. Rong, E. Ploetz, S. Wuttke and O. M. Yaghi, *Nat. Rev. Mater.*, 2023, **8**, 331–340.
- 9 Z. Wang and S. M. Cohen, *Chem. Soc. Rev.*, 2009, **38**, 1315–1329.
- 10 K. K. Tanabe and S. M. Cohen, *Chem. Soc. Rev.*, 2011, **40**, 498–519.
- 11 M. Kalaj and S. M. Cohen, *ACS Cent. Sci.*, 2020, **6**, 1046–1057.
- 12 H. Furukawa, F. Gándara, Y.-B. Zhang, J. Jiang, W. L. Queen, M. R. Hudson and O. M. Yaghi, *J. Am. Chem. Soc.*, 2014, **136**, 4369–4381.
- 13 A. Cadiau, J. S. Lee, D. Damasceno Borges, P. Fabry, T. Devic, M. T. Wharmby, C. Martineau, D. Foucher, F. Taulelle, C.-H. Jun, Y. K. Hwang, N. Stock, M. F. De Lange, F. Kapteijn, J. Gascon, G. Maurin, J.-S. Chang and C. Serre, *Adv. Mater.*, 2015, **27**, 4775–4780.
- 14 D. Fröhlich, E. Pantatosaki, P. D. Kolokathis, K. Markey, H. Reinsch, M. Baumgartner, M. A. van der Veen, D. E. De Vos, N. Stock, G. K. Papadopoulos, S. K. Henninger and C. Janiakothers, *J. Mater. Chem. A*, 2016, **4**, 11859–11869.
- 15 A. J. Rieth, S. Yang, E. N. Wang and M. Dincă, *ACS Cent. Sci.*, 2017, **3**, 668–672.
- 16 M. Gao, J. Wang, Z. Rong, Q. Shi and J. Dong, *RSC Adv.*, 2018, **8**, 39627–39634.
- 17 S. M. T. Abtab, D. Alezi, P. M. Bhatt, A. Shkurenko, Y. Belmabkhout, H. Aggarwal, L. J. Weseliński, N. Alsadun, U. Samin, M. N. Hedhili and M. Eddaoudi, *Chem*, 2018, **4**, 94–105.
- 18 N. Hanikel, M. S. Prévot, F. Fathieh, E. A. Kapustin, H. Lyu, H. Wang, N. J. Diercks, T. G. Glover and O. M. Yaghi, *ACS Cent. Sci.*, 2019, **5**, 1699–1706.
- 19 A. J. Rieth, A. M. Wright, G. Skorupskii, J. L. Mancuso, C. H. Hendon and M. Dincă, *J. Am. Chem. Soc.*, 2019, **141**, 13858–13866.
- 20 Z. Chen, P. Li, X. Zhang, P. Li, M. C. Wasson, T. Islamoglu, J. F. Stoddart and O. K. Farha, *J. Am. Chem. Soc.*, 2019, **141**, 2900–2905.
- 21 J. Canivet, A. Fateeva, Y. Guo, B. Coasne and D. Farrusseng, *Chem. Soc. Rev.*, 2014, **43**, 5594–5617.





- 22 N. C. Burtch, H. Jasuja and K. S. Walton, *Chem. Rev.*, 2014, **114**, 10575–10612.
- 23 A. J. Rieth, K. M. Hunter, M. Dincă and F. Paesani, *Nat. Commun.*, 2019, **10**, 4771.
- 24 X. Liu, X. Wang and F. Kapteijn, *Chem. Rev.*, 2020, **120**, 8303–8377.
- 25 N. Hanikel, X. Pei, S. Chheda, H. Lyu, W. Jeong, J. Sauer, L. Gagliardi and O. M. Yaghi, *Science*, 2021, **374**, 454–459.
- 26 G. A. Cisneros, K. T. Wikfeldt, L. Ojamäe, J. Lu, Y. Xu, H. Torabifard, A. P. Bartók, G. Csányi, V. Molinero and F. Paesani, *Chem. Rev.*, 2016, **116**, 7501–7528.
- 27 V. Babin, C. Leforestier and F. Paesani, *J. Chem. Theory Comput.*, 2013, **9**, 5395–5403.
- 28 V. Babin, G. R. Medders and F. Paesani, *J. Chem. Theory Comput.*, 2014, **10**, 1599–1607.
- 29 G. R. Medders, V. Babin and F. Paesani, *J. Chem. Theory Comput.*, 2014, **10**, 2906–2910.
- 30 F. Paesani, *Acc. Chem. Res.*, 2016, **49**, 1844–1851.
- 31 S. K. Reddy, S. C. Straight, P. Bajaj, C. Huy Pham, M. Riera, D. R. Moberg, M. A. Morales, C. Knight, A. W. Götz and F. Paesani, *J. Chem. Phys.*, 2016, **145**, 194504.
- 32 S. E. Brown, A. W. Götz, X. Cheng, R. P. Steele, V. A. Mandelshtam and F. Paesani, *J. Am. Chem. Soc.*, 2017, **139**, 7082–7088.
- 33 G. R. Medders and F. Paesani, *J. Chem. Theory Comput.*, 2015, **11**, 1145–1154.
- 34 S. K. Reddy, D. R. Moberg, S. C. Straight and F. Paesani, *J. Chem. Phys.*, 2017, **147**, 244504.
- 35 G. R. Medders and F. Paesani, *J. Am. Chem. Soc.*, 2016, **138**, 3912–3919.
- 36 D. R. Moberg, S. C. Straight and F. Paesani, *J. Phys. Chem. B*, 2018, **122**, 4356–4365.
- 37 D. R. Moberg, S. C. Straight, C. Knight and F. Paesani, *J. Phys. Chem. Lett.*, 2017, **8**, 2579–2583.
- 38 D. R. Moberg, D. Becker, C. W. Dierking, F. Zurheide, B. Bandow, U. Buck, A. Hudait, V. Molinero, F. Paesani and T. Zeuch, *Proc. Natl. Acad. Sci. U. S. A.*, 2019, **116**, 24413–24419.
- 39 T. E. Gartner III, K. M. Hunter, E. Lambros, A. Caruso, M. Riera, G. R. Medders, A. Z. Panagiotopoulos, P. G. Debenedetti and F. Paesani, *J. Phys. Chem. Lett.*, 2022, **13**, 3652–3658.
- 40 S. L. Bore and F. Paesani, *Nat. Commun.*, 2023, **14**, 3349.
- 41 G. Kresse and J. Hafner, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1993, **47**, 558.
- 42 G. Kresse and J. Hafner, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1994, **49**, 14251.
- 43 G. Kresse and J. Furthmüller, *Comput. Mater. Sci.*, 1996, **6**, 15–50.
- 44 G. Kresse and J. Furthmüller, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1996, **54**, 11169.
- 45 J. P. Perdew, K. Burke and M. Ernzerhof, *Phys. Rev. Lett.*, 1996, **77**, 3865.
- 46 S. Grimme, J. Antony, S. Ehrlich and H. Krieg, *J. Chem. Phys.*, 2010, **132**, 154104.
- 47 P. E. Blöchl, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1994, **50**, 17953.
- 48 G. Kresse and D. Joubert, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1999, **59**, 1758.
- 49 J. Wang, R. M. Wolf, J. W. Caldwell, P. A. Kollman and D. A. Case, *J. Comput. Chem.*, 2004, **25**, 1157–1174.
- 50 A. V. Marenich, S. V. Jerome, C. J. Cramer and D. G. Truhlar, *J. Chem. Theory Comput.*, 2012, **8**, 527–541.
- 51 A. K. Rappé, C. J. Casewit, K. Colwell, W. A. Goddard III and W. M. Skiff, *J. Am. Chem. Soc.*, 1992, **114**, 10024–10035.
- 52 Y. Wang and J. M. Bowman, *J. Chem. Phys.*, 2011, **134**, 154510.
- 53 X. Cheng and R. P. Steele, *J. Chem. Phys.*, 2014, **141**, 104105.
- 54 X. Cheng, J. J. Talbot and R. P. Steele, *J. Chem. Phys.*, 2016, **145**, 124112.
- 55 K. M. Hunter, F. A. Shakib and F. Paesani, *J. Phys. Chem. B*, 2018, **122**, 10754–10761.
- 56 J. L. Abascal and C. Vega, *J. Chem. Phys.*, 2005, **123**, 234505.
- 57 A. Hudait, D. R. Moberg, Y. Qiu, N. Odendahl, F. Paesani and V. Molinero, *Proc. Natl. Acad. Sci. U. S. A.*, 2018, **115**, 8266–8271.
- 58 W. Smith and T. Forester, *J. Mol. Graphics*, 1996, **14**, 136–141.
- 59 L. Martínez, R. Andrade, E. G. Birgin and J. M. Martínez, *J. Comput. Chem.*, 2009, **30**, 2157–2164.
- 60 S.-T. Lin, P. K. Maiti and W. A. Goddard III, *J. Phys. Chem. B*, 2010, **114**, 8191–8198.
- 61 M. E. Tuckerman, *Statistical Mechanics: Theory and Molecular Simulation*, Oxford University Press, 2010.
- 62 G. J. Martyna, A. Hughes and M. E. Tuckerman, *J. Chem. Phys.*, 1999, **110**, 3275–3290.
- 63 A. R. Leach, *Molecular Modelling: Principles and Applications*, Pearson Education Limited, Essex, UK, 2nd edn, 2001.
- 64 M. Riera, C. Knight, E. F. Bull-Vulpe, X. Zhu, H. Agnew, D. G. A. Smith, A. C. Simmonett and F. Paesani, *J. Chem. Phys.*, 2023, **159**, 054802.
- 65 MBX v1.0, 2021, <http://paesanigroup.ucsd.edu/software/mbx.html>.
- 66 A. P. Thompson, H. M. Aktulga, R. Berger, D. S. Bolintineanu, W. M. Brown, P. S. Crozier, P. J. Veld, A. Kohlmeyer, S. G. Moore, T. D. Nguyen, R. Shan, M. J. Stevens, J. Tranchida, C. Trott and S. J. Plimpton, *Comput. Phys. Commun.*, 2022, **271**, 108171.
- 67 P. Bajaj, A. W. Gotz and F. Paesani, *J. Chem. Theory Comput.*, 2016, **12**, 2698–2705.
- 68 F. Paesani, P. Bajaj and M. Riera, *Adv. Phys.: X*, 2019, **4**, 1631212.
- 69 P. Bajaj, X.-G. Wang, T. Carrington Jr and F. Paesani, *J. Chem. Phys.*, 2017, **148**, 102321.
- 70 P. Bajaj, M. Riera, J. K. Lin, Y. E. Mendoza Montijo, J. Gazca and F. Paesani, *J. Phys. Chem. A*, 2019, **123**, 2843–2852.
- 71 K. M. Hunter, J. C. Wagner, M. Kalaj, S. M. Cohen, W. Xiong and F. Paesani, *J. Phys. Chem. C*, 2021, **125**, 12451–12460.
- 72 H. Bakker and J. Skinner, *Chem. Rev.*, 2010, **110**, 1498–1517.
- 73 A. A. Kananenka and J. Skinner, *J. Chem. Phys.*, 2018, **148**, 244107.
- 74 M. Rossi, H. Liu, F. Paesani, J. Bowman and M. Ceriotti, *J. Chem. Phys.*, 2014, **141**, 181101.
- 75 A. Luzar and D. Chandler, *Nature*, 1996, **379**, 55–57.

