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Understanding the initial events of the oxidative damage and protection mechanisms of the AA9 lytic polysaccharide monooxygenase family†

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Lytic polysaccharide monooxygenase (LPMO) is a new class of oxidoreductases that boosts polysaccharide degradation employing a copper active site. This boost may facilitate the cost-efficient production of biofuels and high-value chemicals from polysaccharides such as lignocellulose. Unfortunately, self-oxidation of the active site inactivates LPMOs. Other oxidoreductases employ hole-hopping mechanisms as protection against oxidative damage, but little is generally known about the details of these mechanisms. Herein, we employ highly accurate theoretical models based on density functional theory (DFT) molecular mechanics (MM) hybrids to understand the initial steps in LPMOs' protective measures against self-oxidation; we identify several intermediates recently proposed from experiment, and quantify which are important for protective hole-hopping pathways. Investigations on two different LPMOs show consistently that a tyrosine residue close to copper is crucial for protection: this explains recent experiments, showing that LPMOs without this tyrosine are more susceptible to self-oxidation.

Introduction

The heavy fossil fuel dependence of global energy production^{1,2} represents a major challenge in transitioning to a sustainable society. Further, essential products, like fertilizers, plastics, and chemicals, are also derived from fossil fuels,³ and transitioning will thus require alternative feedstocks.

Part of the solution may come from the production of fuel and chemicals from lignocellulose which comprise an abundant and cheap resource that does not compromise food security.^{4–6} Yet, the cost-efficient production of biofuels and materials or chemicals from lignocellulose is currently hampered by its natural recalcitrance.⁷ Consequently, the discovery of a new class of enzymes that boosts the degradation of polysaccharides has attracted considerable attention.^{7–12} This enzyme class is now termed lytic polysaccharide monooxygenase (LPMO) and is currently grouped into eight families: AA9 to AA11 and AA13 to AA17.^{8,9,13–19} The LPMOs differ in substrate-specificity^{8,14–16,19–25} but they generally catalyze the oxidation of the glycosidic C–H bonds in the polysaccharide, leading ultimately to a cleavage of the glycosidic bond.^{10,11,26} This cleavage causes the boosting effect and is administered by a copper-containing active site with Cu coordinated by two histidines^{7,16–19,27} that has been coined the histidine brace.²² The

active site with the histidine brace in the Cu(II)-resting state is shown in Fig. 1.

The oxidation of inactive C–H bonds requires a potent oxidant; several theoretical studies have shown that a Cu(II)-oxyl^{11,28,29} is sufficiently potent. This Cu(II)-oxyl can be formed with a suitable co-substrate after reduction of the Cu(II)-resting state (see Fig. 1, structure 1). Whether this co-substrate is O₂ or H₂O₂ is still an ongoing debate,^{30,31} although some LPMOs have been shown to use exclusively hydrogen peroxide.^{32,33}

While a major focus in LPMO research has been the investigation of the substrate oxidation mechanism,^{10,11,27,34} the potent oxidative species can also cause oxidative damage to LPMOs in the absence of a substrate. This self-oxidation remains poorly understood but renders LPMOs inactive and thus hinders efficient exploitation of the enzymes. Proteomics techniques showed that oxidative damage is primarily confined to the two histidine residues coordinating the copper.³⁰ We recently employed quantum-mechanics/molecular mechanics (QM/MM) to show that the oxidation of the histidine brace³⁵ can be initiated by the hydrogen abstraction from either of the two histidines by a Cu(II)-oxyl (1), forming a species (3) with a histidyl radical coupled to a [CuOH]⁺ moiety (see reaction II in Fig. 1). In a complementary series of experimental and theoretical investigations, we and others have proposed that a tyrosine close to the active site effectively protects LPMOs against the self-oxidation processes by forming a tyrosyl radical (2).^{36–40} Similar oxidative stress mechanisms have been previously studied employing DFT for Cu-amyloid beta peptide model complexes.⁴¹

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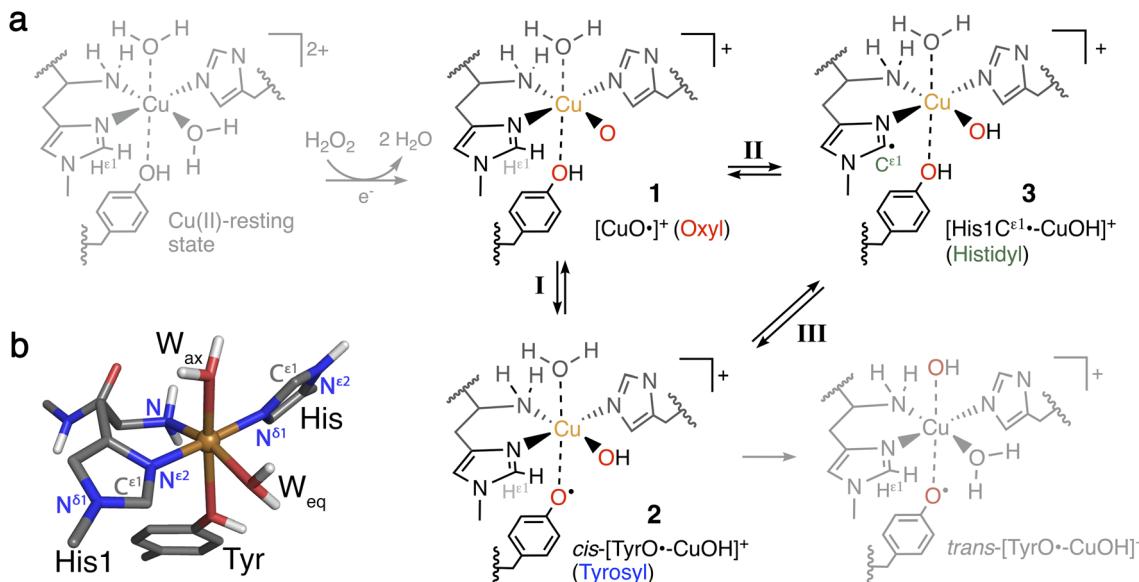


Fig. 1 (a) Formation of Cu(II)-oxyl 1 from the Cu(II)-resting state (in gray) with a reductant and H₂O₂, and generation of intermediates 2 and 3 from 1. Intermediates 1–3 (in black) are investigated here for their role as part of oxidative damage or protective pathways (species in gray are not included in this study). (b) Ball-and-stick model of the first coordination sphere of the QM/MM optimized Cu(II)-resting state of *LsAA9* employing TPSS/def2-SV(P).

The tyrosyl radical (2) can be formed by reaction I in Fig. 1 where the Cu(II)-oxyl species (1) abstracts the hydrogen from the hydroxy group of the tyrosine,⁴⁰ thus forming a [CuOH]⁺ moiety coordinated to a tyrosyl radical. Tyrosyl radicals have now been observed for a number of LPMOs and are characterized by having UV-vis spectra with strong, sharp absorption features around 400–420 nm.^{36–38,42} We note in passing that one investigation³⁹ (for *LsAA9*) showed rather different features, and based on QM/MM and time-dependent density functional theory (TD-DFT), we⁴⁰ have argued that this spectrum originated in a different species (*trans*-[TyrO[•]-CuOH]⁺ in Fig. 1).

The tyrosine residue has long mystified LPMO researchers: it is widely conserved in most LPMO families, except for the majority of AA10 LPMOs, where it is replaced by phenylalanine.^{7,8,17,27,43} It seems now that its role is to initialize a hole-hopping mechanism, known from other oxidoreductase enzymes: a radical species (“hole”) is generated close to the active site by a highly oxidizing species; this hole is then directed away (towards the surface) through electron-transfer chains comprised of aromatic/redox-active amino acids.⁴⁴ Yet, detailed investigations of these reactions have been rare for oxidoreductases;^{45,46} and has only recently begun to appear for LPMOs.^{36,37,39} To confuse matters more, a hitherto undetected histidyl species was recently claimed to be observed through transient spectroscopy in *LsAA9*.⁴² The histidyl radical was assigned based on UV-vis absorption at 360 nm and high-energy-resolution fluorescence-detected X-ray absorption spectroscopy (HERFD-XAS). It was argued that this radical was formed before the tyrosyl radical, but is otherwise part of the same protective hole-hopping mechanism. This investigation also detected a tyrosyl radical with its characteristic absorption features.

Several open questions now remain concerning the initial parts of the hole-hopping mechanism: (i) How does the tyrosyl radical formation compare energetically, relative to the histidyl radical formation? (ii) Are both tyrosyl and histidyl radicals part of the protective mechanism? (iii) Do different LPMOs display different mechanisms? The last open question arises since previous QM/MM calculations on the protective or oxidative damage mechanisms focused exclusively on a particular member of the AA9 family, namely *LsAA9*,^{35,40} whereas experimental studies focused on different LPMOs.^{36,39,42} In this paper we will address (i) and (ii) by employing QM/MM calculations for reactions I–III in Fig. 1. We will also address (iii) by, for the first time, directly comparing oxidative damage and protective mechanisms of two enzymes, namely *LsAA9* and *TaAA9* (see Fig. 2). One of our findings is that *LsAA9* and *TaAA9* (that have quite similar second coordination spheres) occasionally display significant differences in their mechanisms. Given that little is known concerning the initial events of these protective mechanisms, our findings are likely to be highly relevant for other oxidoreductases as well.

Computational details

QM/MM calculations

All QM/MM calculations were based on equilibrated and QM/MM optimized structures of *LsAA9A* from *Lentinus similis* (5ACF⁴⁷) and *TaAA9A* from *Thermoascus aurantiacus* (2YET²²) from the studies ref. 35 and 48, respectively. We denote the two enzymes *LsAA9* and *TaAA9* in the following. A detailed description of the original protein setup, choice of protonation



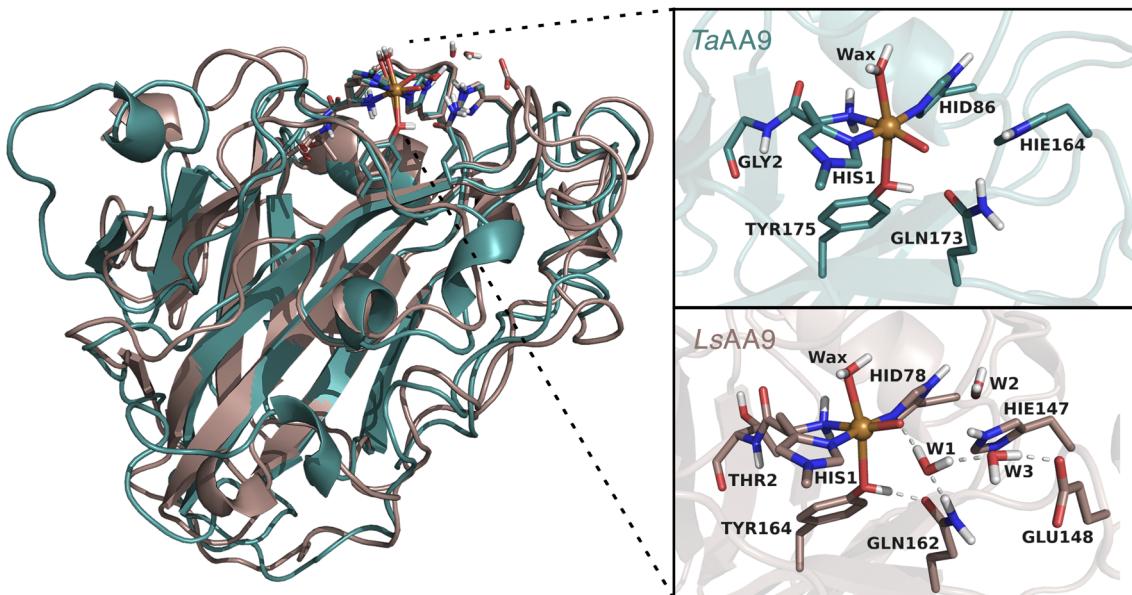


Fig. 2 Comparison of the structures and active sites of *TaAA9* and *LsAA9*. The structures are obtained from QM/MM optimizations for the oxyl $[\text{CuO}]^+$. Labels refer to PDB 2YET²² and 5ACF⁴⁷ for *TaAA9A* and *LsAA9A*, respectively. The QM region employed in this study are also shown (note that *TaAA9* has a similar hydrogen bonding network as *LsAA9* in the active site region, but this is kept in the MM region).

states and initial equilibration is given in ref. 28 and 29 for *LsAA9* and ref. 48 for *TaAA9*.

The second-sphere tyrosine is present in both enzymes (Tyr175 and Tyr164 in *TaAA9* and *LsAA9*, respectively, as shown in Fig. 2) and in both LPMOs a tyrosyl species was spectroscopically characterized.^{36,42} The protonation state of the second-coordination sphere histidine (His164) in *TaAA9* was changed from a doubly protonated form (Hip164) to a singly protonated one (Hie164), to match the protonation state present in the *LsAA9* structure.⁴⁰

The employed QM regions are shown in Fig. 2. For *LsAA9* the QM region comprised the copper ion, the oxyl/hydroxyl ligand, the complete methylated His1 residue, and the side chains of His78, Tyr164, Thr2, Gln162, Glu148 and Hie147 as well as four water molecules. This is the same QM region also employed in ref. 35 and a slightly larger QM region (extended by Glu143 and three water molecules) as employed in ref. 40. For *TaAA9*, the QM region consists of the copper ion and all the residues of the first coordination sphere, *i.e.*, the full His1 residue, the oxyl/hydroxyl ligand, the imidazole ring of His86, the phenol ring of Tyr175 and the axial water molecule. Additionally, the side chains of Gln173 and Hie164 were included, as well as parts of the backbone from the Gly2 residue. The latter was included up to the C^α atom and capped by replacing the carboxyl C with a hydrogen atom. The same applies to the corresponding residue Thr2 in *LsAA9*.^{49,50} The side chains of all other residues were capped with a hydrogen by replacing C^α . Note that *TaAA9* does not have a Glu residue corresponding to Glu148 in *LsAA9*, and we therefore did not include the residue at this position in the QM region for *TaAA9*. The QM region for *TaAA9* is larger than the one used in ref. 48 to allow better comparison to the one employed for *LsAA9*. Notably, the active site water molecules were included in *LsAA9* as part of the QM region; there are

also (non-coordinated) water molecules present in the second coordination sphere of *TaAA9*, but we decided to keep them in the MM region as done in ref. 48. The non-coordinated water molecules are therefore accounted for in both *TaAA9* and *LsAA9*, albeit at different levels of theory. Generally, the MM region only showed a small contribution to the QM/MM reaction energies ($\leq 10 \text{ kJ mol}^{-1}$) and thus we do not expect an explicit treatment of non-coordinated water molecules in the QM region of *TaAA9* to significantly change the energetics.

All structure optimizations and energy calculations employed a subtractive QM/MM approach with electrostatic embedding as implemented in the modular program ComQum.^{49,51} ComQum interfaces the QM software Turbomole⁵² and the MM software AMBER.⁵³ For all the structure optimizations the dispersion-corrected TPSS-D3 functional^{54,55} with Becke-Johnson damping⁵⁶ and a def2-SV(P) basis set^{57,58} were employed. For all calculations with TPSS the resolution of identity (RI) approximation with standard auxiliary basis sets was applied. All atoms in the MM region were kept fixed during the geometry optimizations.

The reaction and activation energies were computed as linear transit calculations without thermochemical or zero-point vibrational energy corrections as these have shown to be small for hydrogen abstractions.^{59–61} The reported energies for the reactant, product and transition state structures were obtained from single-point calculations on the QM/MM optimized structures, employing the def2-TZVPP basis set⁵⁷ (with an auxiliary basis set of the same size) and the functionals TPSS-D3 and B3LYP-D3.^{55,62–64} The single-point calculations include protein electrostatics, and an MM contribution calculated with TPSS/def2-SV(P). For brevity, we generally denote these functionals as TPSS and B3LYP throughout the paper although all calculations (for energies and geometries) always included



dispersion corrections. On several occasions we compare to previous calculations,³⁵ where generally TPSS/def2-SV(P) was used for geometry optimizations and B3LYP/def2-TZVPD for single point energies.

Our calculations took departure from the $[\text{CuO}]^+$ species (**1** in Fig. 1). This intermediate was selected as previous calculations have shown that this is the most likely species to abstract hydrogen from the substrate, and it is readily formed from the reaction between the reduced resting state and H_2O_2 .^{29,65-67} From these previous calculations, we know that this $[\text{CuO}]^+$ moiety is most accurately described as a Cu^{2+} and an O^- radical, spin-coupled to either a triplet or an open-shell singlet. Thus, the intermediates considered here are generally expected to attain either singlet and triplet spin states, and we have always attempted to obtain both of these spin states. As previous studies showed the closed-shell singlet states to be significantly higher in energy,^{35,40} the singlet states were only calculated as open-shell species in this study. The open-shell singlets were obtained in a spin-unrestricted (broken symmetry) formulation and the calculations were typically initiated from the triplet-state structures (using the triplet molecular orbital coefficients as initial guess). Similarly, for calculations with def2-TZVPP the open-shell singlet calculations were always started from the triplet molecular orbital coefficients. The convergence to an open-shell singlet was always confirmed by inspection of the Mulliken spin densities (which are reported in the ESI[†]). We compare the obtained electronic structure of **1** to previous calculations involving the $\text{Cu}(\text{II})$ -oxyl (**1**) intermediate. However, since this has been discussed frequently in the literature, this discussion is moved to the ESI.[†]

For reaction **II** of *TaAA9* we were unable to locate an open-shell singlet for **3**, although it could be obtained for a conformer of **3**, which we here denote **3'**. The calculations along reaction **II** collapse into a closed-shell singlet along the linear transit, the last open-shell singlet energy being 125 kJ mol⁻¹ (distance between O_{oxyl} and $\text{H}_{\text{Ty}}^{\text{OH}}$ restraint to 1.27 Å) for B3LYP and 84 kJ mol⁻¹ (distance restraint to 1.30 Å) for TPSS, *cf.* Table S3.[†] This issue has also been observed previously for the TPSS functional with *LsAA9*.³⁵ In these cases, we base the discussion on the triplet state.

For the calculations of reaction **III**, we observed for both *LsAA9* and *TaAA9* that the QM/MM energies for the barrier showed high MM contributions (>33 kJ mol⁻¹). We analyzed the energy contributions from individual residues and observed that this was caused by a residue close to the tyrosine (Phe43 and Pro30 in *LsAA9* and *TaAA9*, respectively). Hence, we chose to include these residues in the QM region for this reaction, reducing the MM contributions to ~13 kJ mol⁻¹ for both enzymes. Since the reaction for *TaAA9* was more favorable for the triplet state and no open-shell singlet could be located for *LsAA9*, we base the discussion on the triplet state.

UV-vis spectra

UV-vis spectra were calculated for intermediates **2** and **3/3'** considering both open-shell singlet and triplet spin states, if open-shell singlet states were obtained. We performed TD-DFT

calculations in Gaussian 16,⁶⁸ employing the CAM-B3LYP functional⁶⁹ and def2-TZVPP basis set (both as implemented in Gaussian) and including 45 states (roots). All UV-vis spectra were calculated from structures with the smaller QM regions (*i.e.*, calculations for **3'** and **3** do not include Phe43 and Pro30 for *LsAA9* and *TaAA9*, respectively). For a better comparison to the spectra calculated for *LsAA9* in ref. 40, the input structures for *LsAA9* and *TaAA9* in the TD-DFT calculations were slightly truncated compared to the QM system that was employed for the structure optimizations and energy calculations. For *LsAA9*, Thr2 was removed to the amide N, and for *TaAA9* only the amide-group of residue Gly2 was included. The cut bonds were saturated with a hydrogen atom (see Fig. S3 in the ESI[†]). Calculated oscillator strengths and energies were convoluted using a Gaussian function and a broadening factor of 0.3 eV. To probe the electrostatic effect of the environment, we in one case (for **2** of *TaAA9* in the triplet state) added point-charges from the QM/MM calculations.

Results

We start by comparing *LsAA9* and *TaAA9* for the tyrosyl radical (**2**) formation directly from the $\text{Cu}(\text{II})$ -oxyl (**1**) intermediate (*i.e.*, reaction **I** in Fig. 1). Next, we compare these two AA9 LPMOs for histidyl radical (**3**) formation (*i.e.* reaction **II** in Fig. 1). We finally discuss whether intermediates **2** and **3** can be inter-converted (reaction **III** in Fig. 1) as recently suggested.⁴²

Hydrogen abstraction from tyrosine

The calculated energy diagrams for reaction **I** of *TaAA9* and *LsAA9* are displayed in Fig. 3. We additionally show the structures and selected distances for the most stable electronic states in Fig. 4a-c (*TaAA9*) and Fig. 4d-f (*LsAA9*). A complete overview of reaction energies and barriers for both functionals is provided in Table S1.[†]

While there are some differences between the two employed functionals, both agree that the activation energies are consistently lower for *TaAA9*: we obtain a barrier of 44 kJ mol⁻¹ (28 kJ mol⁻¹ with TPSS), compared to a reaction barrier of 77 kJ mol⁻¹ for *LsAA9* (37 kJ mol⁻¹ with TPSS). Both functionals

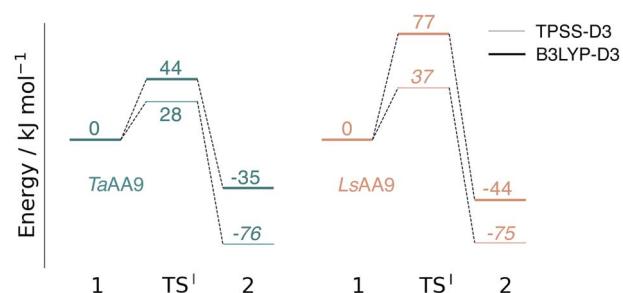


Fig. 3 Energy diagrams (in kJ mol^{-1}) of the reaction **I** for *TaAA9* (left) and *LsAA9* (right). Energies were obtained with def2-TZVPP based on structures optimized with TPSS/def2-SV(P). The reactant (**1**) in the triplet state was used as reference. Energies in bold refer to a triplet state and those in italic to open-shell singlet states.



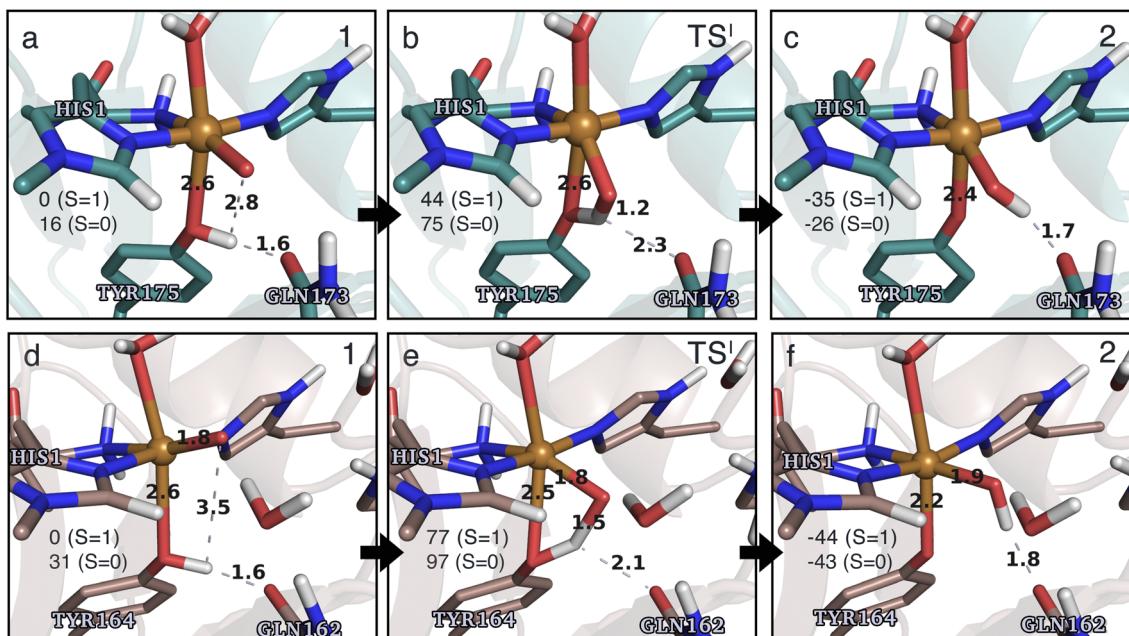


Fig. 4 The hydrogen-abstraction reaction I illustrated for *TaAA9* (a–c) and *LsAA9* (d–f). The structures were optimized using TPSS/def2-SV(P) while the energies given were obtained employing B3LYP/def2-TZVPP. Only the structures for the most stable electron configuration (either open-shell singlet or triplet) are shown. Key distances are given in Å and energies in kJ mol^{-1} with reference to **1** in the triplet state. Further bond distances for the different intermediates can be found in Table S2.†

also agree that the reaction is exothermic. In fact, the final reaction energies are quite similar between the two proteins, where we obtain an energy of -35 kJ mol^{-1} for *TaAA9* (-76 kJ mol^{-1} with TPSS) and -44 kJ mol^{-1} for *LsAA9* (-75 kJ mol^{-1} with TPSS). The result is qualitatively similar to previous calculations for *LsAA9* where a smaller QM region was employed;⁴⁰ here the reaction barrier for reaction I was predicted to be 64 kJ mol^{-1} (53 kJ mol^{-1} with TPSS). We can thus conclude that *LsAA9* and *TaAA9* are likely to form the same tyrosyl intermediate, albeit with different kinetics. We can also compare structures of **1** and **2** as well as the transition state in the two LPMOs in Fig. 4. For both enzymes, the largest structural change between **1** and **2** is the change in the Cu–O_{Tyr} bond, which shortens from 2.6 \AA (both enzymes) in **1** to 2.4 \AA (*TaAA9*) or 2.2 \AA (*LsAA9*) in **2**. Thus, the de-protonation of tyrosine leads to coordination of the tyrosyl, consistent with a previous calculation on *LsAA9*.⁴⁰ In both enzymes, the tyrosine OH-group forms a hydrogen bond to a Gln residue in **1** (the bond distance is 1.6 \AA in both *TaAA9* and *LsAA9*). This hydrogen bond is partly broken in the transition state; here *TaAA9* and *LsAA9* are different as the bond distance is 2.3 \AA in *TaAA9* and 2.1 \AA in *LsAA9*. Notably, the larger distance in *TaAA9* means that we can optimize a stable intermediate after breaking of the hydrogen bond to Gln, but this intermediate is close to degenerate with **TS**¹, and we have therefore not included it in Fig. 4 (a full figure including this intermediate is given in the ESI, see Fig. S1 and S2†). Interestingly, the hydrogen bond to Gln is re-formed in **2** (with a distance of 1.7 \AA in *TaAA9* and 1.8 \AA in *LsAA9*), where the OH group is now coordinated to Cu(II). Other distances are roughly similar (cf. Table S2†), so the energetic differences

between the reaction barrier of the two LPMOs in Fig. 3 may be traced back to the structural differences in the local hydrogen bonding of Gln to the transition state.

The spin states of **2** deserve a comment as several different results have been obtained in the literature: a somewhat noisy electron paramagnetic resonance (EPR) spectrum was recorded of **2** in *TaAA9*,³⁶ suggesting that a triplet is energetically within reach. Meanwhile Jones *et al.*³⁷ obtained an EPR silent tyrosyl (**2**) for *HjAA9*, while a very recent study proposes a triplet spin state for **2** in *LsAA9* (also based on EPR).⁴² We find that for *LsAA9*, the singlet and triplet states are essentially degenerate with both functionals. Similarly, ref. 40 reported that the splittings for *LsAA9* are small ($<5.5 \text{ kJ mol}^{-1}$) for both functionals (the triplet state was found to be slightly more stable, independent of the functional used). Thus, for this LPMO our results are commensurate with the recent EPR results in ref. 42, *i.e.*, the triplet state is energetically within reach. The results for *TaAA9* are more ambiguous: the triplet state is most stable according to B3LYP (by 9 kJ mol^{-1}), whereas TPSS predicts the open-shell singlet state to be more stable (by 20 kJ mol^{-1}). The triplet is thus still within reach, but the magnitude of the splitting is more functional dependent. Obtaining the correct ordering of close-lying spin states is a known issue with DFT.⁷⁰ However, with spin-state splittings of the size observed for **2**, even highly correlated wave function methods may be challenged. For instance, Delcey *et al.*⁷¹ obtained spin-state splittings of 24 kJ mol^{-1} with restricted active space second-order perturbation (RASPT2) calculation for another metalloenzyme (a [NiFe]-hydrogenase), but were still unable to unequivocally assign the correct ground state. Meanwhile, our previous

comparison to highly accurate multiconfigurational wave-function (CASPT2) calculations on LPMO intermediates have also shown that discrepancies for spin-state splittings can occur with DFT.⁷² In conclusion, we can only expect qualitative results for reactions where **2** is involved (this is also the reason we mostly will base our conclusions on results obtained from two different DFT functionals).

In terms of spin densities (see Tables S7 and S8 in the ESI†), both *TaAA9* and *LsAA9* have significant spin density on tyrosine in **2**. This was previously observed for *LsAA9* in ref. 40, and confirms that intermediate **2** observed in *TaAA9* and *LsAA9* have similar electronic structures. This electronic structure is a complicated coupling between the tyrosyl radical and a $[\text{CuOH}]^+$ unit, where the latter itself is best described as Cu(II) coupled to a OH radical. This complicated spin-coupling presumably also explains the observed differences between the two functionals. During reaction **I**, the spin density decreases on $\text{O}_{\text{oxyl/hydroxyl}}$ and increases on tyrosine/tyrosyl. We also note that an in-depth analysis of the spin densities over the potential energy surface of reaction **I** for *LsAA9*, reveals that for the TS^{I} state, TPSS obtains a broken-symmetry singlet with the expected spin distribution, but with the magnitude of the spin populations reduced (the issue is not seen for *TaAA9*). This issue is similar to what was described in the Computational details and has also been observed previously,³⁵ although the state obtained in Fig. 3 is not completely collapsed to a closed-shell singlet.

Hydrogen abstraction from histidine

An alternative to the reaction investigated in the previous subsection is the abstraction of a hydrogen from the histidine brace by the $[\text{CuO}]^+$ moiety in **1** (*i.e.*, reaction **II** in Fig. 1). Indeed, we recently investigated this for *LsAA9* as the first step of the oxidative self-damage reaction.³⁵ We compare the reaction profiles (barrier and reaction energy) of reaction **II** for *TaAA9* and *LsAA9* in Fig. 5 (all calculated energies associated with reaction **II** are provided in Table S3†). As can be seen from Fig. 5, *LsAA9* and *TaAA9* overall have the same energy profile for the H-abstraction from the active site histidine (His1). The reaction barrier differs only by 5 kJ mol⁻¹ (1 kJ mol⁻¹ for TPSS

and the reaction energy differs by less than 3 kJ mol⁻¹ for both functionals. Note that Fig. 5 does not include the highly exothermic OH recombination reaction investigated as part of the oxidative damage pathway in ref. 35 for *LsAA9*. However, the similarity of the reaction energetics for the first part of the reaction for *LsAA9* and *TaAA9* suggest that the oxidative damage can proceed through this pathway, independent of the LPMO.

Structures for the reactants, transition states, and products, along with selected distances, are shown in Fig. 6 (additional distances are provided in Table S4 in the ESI†). As can be seen from this figure, structural changes during reaction **II** are similar for *TaAA9* and *LsAA9*, as expected based on the similar reaction profiles: most copper-ligand bond distances undergo small changes with the sole exception being the Cu–O distance (see Table S4†), which increases by 0.09–0.10 Å, consistent with the protonation of the oxyl. We also note that for *TaAA9* we were able to obtain an isomer of **3**, here denoted as **3'** (shown in Fig. 8a). This isomer differs from **3** in Fig. 6 in that the OH-group of the tyrosine points towards the C^{ε1} in the imidazole ring of His1. In **3**, the OH group instead forms a hydrogen bond with the oxygen in Gln173. Since the energy difference between the isomers is rather small ($ΔE = 12\text{--}19$ kJ mol⁻¹ in their triplet state depending on the functional), it is likely that they both exist in solution. The structure of the **3'** conformer seems optimal for the transfer of the H-atom of the tyrosine OH-group to the histidyl, according to reaction **III** in Fig. 1 as recently suggested.⁴² We investigate this possibility in the next subsection.

Based on the spin densities, differences between the electronic changes during reaction **II** are also minor between the two LPMOs: the spin densities (Table S7†) decrease significantly on oxygen in the OH group of the $[\text{CuOH}]^+$ moiety in **3**, compared to O_{oxyl} in $[\text{CuO}]^+$ (**1**), while the spin density increases on His1. The increase mainly occurs on the de-protonated C^{ε1}, suggesting that **3** is indeed a histidyl radical, coupled to a $[\text{CuOH}]^+$ moiety. Similar to **2**, this spin-coupling in **3** is also somewhat complicated. Yet, compared to the reaction **I** involving **2**, the barrier and energetics of reaction **II** are less dependent on the employed functional, although the spin-state splitting are functional dependent, also for intermediate **3** (as will be detailed below). In the following subsection we will discuss another reaction (**III**) involving intermediate **2**, where we also see a larger influence of the employed functional. Thus, reactions involving **2** seems to be particularly functional dependent, but more detailed investigations with multi-configurational wave functions will be required to understand this difference.

In a recent experimental investigation with HERFD-XAS and UV-vis spectroscopy, a histidyl intermediate was claimed to be characterized as an open-shell spin singlet,⁴² although it is unclear how the spin state was determined. The histidyl intermediate was only characterized for *LsAA9* and we find the spin-state splitting for **3** in *LsAA9* to be small, but somewhat functional dependent: the splitting is only 5 kJ mol⁻¹ with TPSS, the open-shell singlet being most stable. With the B3LYP functional we obtain a splitting of 18 kJ mol⁻¹ with the triplet being most stable.³⁵ For *TaAA9* the open-shell singlet calculations

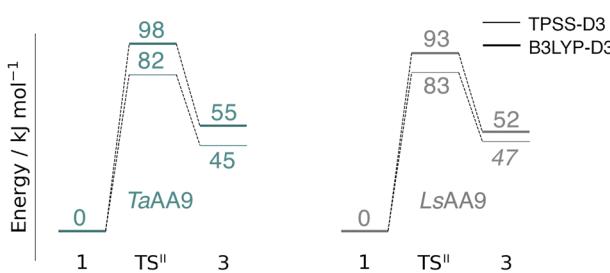


Fig. 5 Energy diagrams (in kJ mol⁻¹) of the reaction **II** for *TaAA9* (left) and *LsAA9* (right). Energies for *TaAA9* were obtained with def2-TZVPP based on structures optimized with TPSS/def2-SV(P). Energies for *LsAA9* are from ref. 35. The reactant (**1**) in the triplet state was used as reference. Energies in bold refer to a triplet state and those in italic to open-shell singlet states.



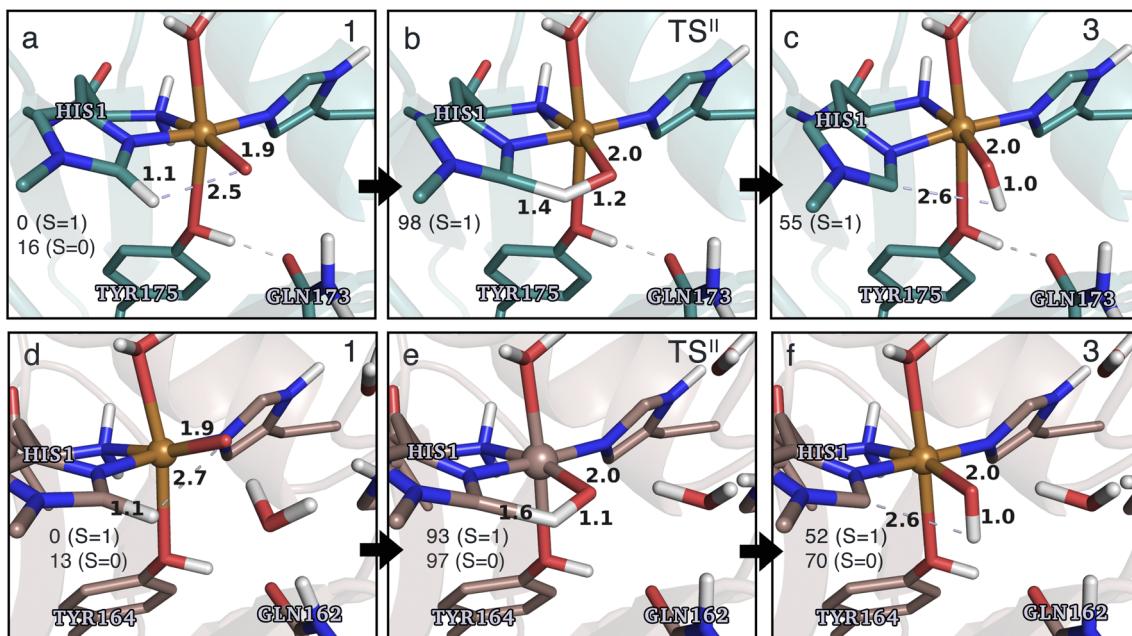


Fig. 6 The hydrogen-abstraction reaction II illustrated for *TaAA9* (a–c) and *LsAA9* (d–f). The structures for *TaAA9* were optimized using TPSS/def2-SV(P) while the energies given were obtained employing B3LYP/def2-TZVPP. Structures and energies for *LsAA9* are from ref. 35. Only the structures for the most stable electron configuration (either open-shell singlet or triplet) are shown. Key distances are given in Å and energies in kJ mol^{-1} with reference to 1 in the triplet state. Further bond distances for the different intermediates of *TaAA9* can be found in Table S4.†

converged into a closed-shell singlet for 3, while we obtain both triplet and open-shell singlet states for the conformer 3'; in this case the open-shell singlet and triplet are essentially degenerate (the triplet is only 3 kJ mol^{-1} more stable with both TPSS and B3LYP). In light of the results for the tyrosyl radical (see previous subsection), it is likely that small differences in the active-site architectures between *TaAA9* and *LsAA9* lead to small differences in spin-state splittings. Since the splittings are small, this can also lead to differences in which spin states are most stable. However, with the present accuracy of the used functionals, the splittings are generally too small to clearly differentiate the spin states.

Before investigating reaction III, we note that we have here only investigated H-abstraction from His1. This was decided based on the structure in *TaAA9*, where H-abstraction from His1 appeared more plausible compared to His86, given the notably shorter distance between $\text{H}^{\varepsilon 1}$ and O_{oxyl} (2.48 Å) in contrast to the $\text{H}^{\varepsilon 1}$ of His86 (3.39 Å). This was confirmed by a test calculation (Table S3†) for the triplet potential energy surface (PES), showing that the reaction barrier is indeed significantly higher compared to the abstraction from His1 (27–28 kJ mol^{-1} higher, depending on the functional). Moreover, the product is thermodynamically less stable by 19–21 kJ mol^{-1} . It is interesting to note that in *LsAA9*, the abstraction from His78 (equivalent to His86 in *TaAA9*) occurs with only minor changes in energy (<8 kJ mol^{-1} for the reaction barrier and energy for both functionals), compared to the H-abstraction from His1.³⁵ The difference between the two LPMOs occurs since the distances of $\text{H}^{\varepsilon 1}$ on His78 and His1 to O_{oxyl} in 1 is much closer (2.31 Å and 2.70 Å)³⁵ in *LsAA9*. This underlines that small differences in the

active site architecture can lead to mechanistic differences, and this will become more apparent in the discussion regarding reaction III in the next subsection.

Conversion of the histidyl radical to a tyrosyl radical

Spurred by the recent proposal of a histidyl radical⁴² (3) as the first intermediate in a protective hole-hopping pathway, we investigated whether the histidyl radical (3) can abstract a hydrogen from the tyrosine, thereby restoring histidine while forming a tyrosyl radical (2). This reaction is labeled III in Fig. 1. The reaction barrier and energies are shown in Fig. 7 and 8. Interestingly, the two LPMOs are remarkably different: for *TaAA9* the reaction is kinetically feasible with a barrier of 75 kJ mol^{-1} (59 kJ mol^{-1} for TPSS), and is predicted to be thermodynamically favorable with a reaction energy of

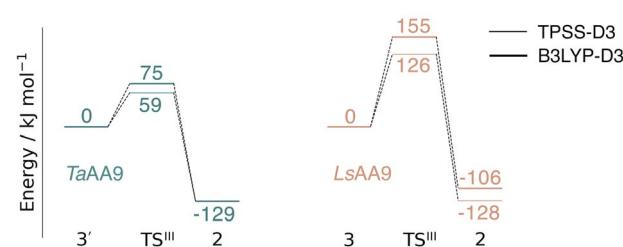


Fig. 7 Energy diagrams (in kJ mol^{-1}) for the reaction III for *TaAA9* (left) and *LsAA9* (right). The reactants 3' and 3 for *TaAA9* and *LsAA9*, respectively, were used as reference. Results were obtained for the triplet state with an extended QM region (see Computational details). Results for the smaller QM region are provided in Table S5 in the ESI.†

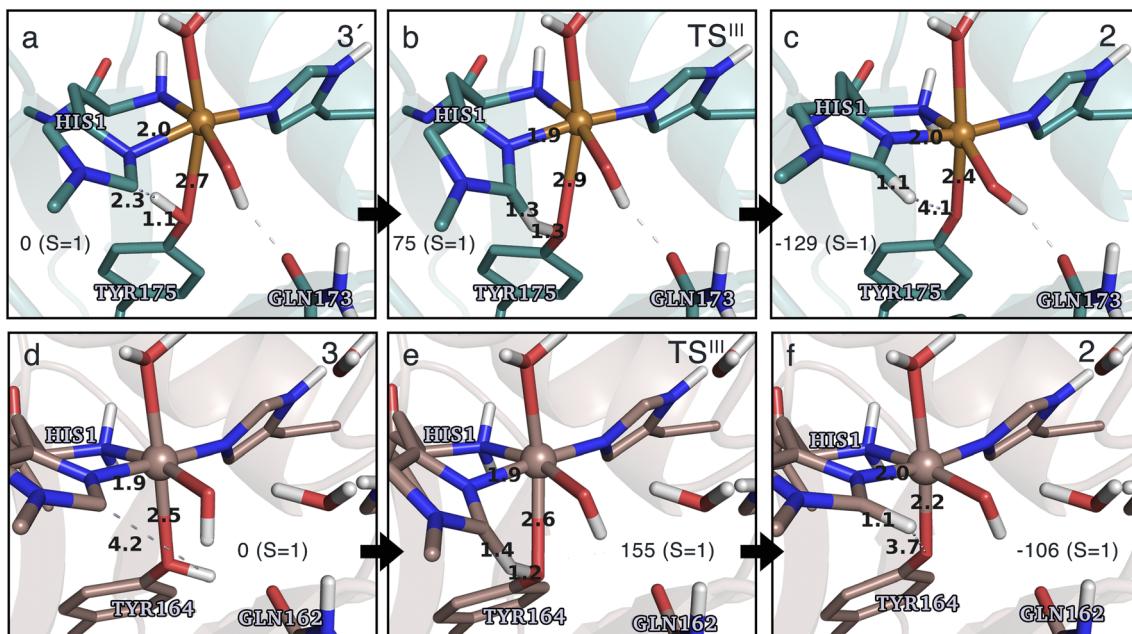


Fig. 8 The hydrogen-abstraction reaction III illustrated for *TaAA9* (a–c) and *LsAA9* (d–f). The structures were optimized using TPSS/def2-SV(P) while the energies given were obtained employing B3LYP/def2-TZVPP. Key distances are given in Å and energies in kJ mol^{-1} with reference to the reactant. Further bond distances for the different intermediates can be found in Table S6.† Note that the structures were obtained with an extended QM region including Phe43 and Pro30 for *TaAA9* and *LsAA9*, respectively.

–129 kJ mol^{-1} for both functionals. In comparison, *LsAA9* shows a high reaction barrier of 155 kJ mol^{-1} (126 kJ mol^{-1} for TPSS) with a reaction energy of –129 kJ mol^{-1} (–128 kJ mol^{-1} for TPSS). The reactant, transition state, and product structures are shown in Fig. 8 (additional distances for chosen atoms are provided in Table S6 in the ESI†). The transition states also display somewhat different bond distances in the first coordination sphere to copper; the Cu–O_{Tyr} distance varies from 2.9 Å in *TaAA9* to 2.6 Å in *LsAA9*, highlighting that the otherwise quite similar active-site architectures may lead to different reactivity. As for reaction I, the de-protonation of tyrosine's OH-group leads to a decrease in the Cu–O_{Tyr175} distance of 0.3 Å in the product (2). We finally note that we decided to use 3' for *TaAA9* in Fig. 7 and 8. This choice was made since we then have consistently small MM energies for the barrier (the reaction energy is less affected) as the QM region was enlarged for 3'. We have not attempted to extend the QM region for 3, but based on the calculations with the smaller QM region, the energy difference between 3 and 3' for *TaAA9* is sufficiently small so that none of the above conclusions change (see Table S5†).

Discussion

This work represents the first comparison of oxidative damage and protective mechanisms in two different LPMOs: we have shown that it is energetically (and kinetically) feasible for *TaAA9* and *LsAA9* to form tyrosyl (2) radicals. This commensurates with experiments for both *TaAA9* and *LsAA9*, where a number of studies detected 2 (ref. 36–38 and 42) for different LPMOs, including *TaAA9* (ref. 36) and *LsAA9*.⁴² The detection of a histidyl radical (3) has only been proposed recently for *LsAA9*.⁴² The

histidyl radical was proposed to be part of a protective mechanism, where it was converted into the tyrosyl radical. Our results show that at least the formation of this radical is energetically similar for *TaAA9* and *LsAA9*. However, if we consider the barrier for the conversion 3 → 2 alone (in Fig. 7), we find that for *LsAA9* the barrier is rather high (155 kJ mol^{-1} for B3LYP), but the overall reaction is feasible with a reaction energy of –106 kJ mol^{-1} . It is interesting to compare this to the values obtained in ref. 35 where formation of 3 (from 1) is followed by recombination of the OH group from $[\text{Cu}-\text{OH}]^+$ in 2 to the histidyl radical, forming a 2-hydroxy-histidine (see Fig. 5 of ref. 35). This reaction was proposed to be part of the oxidative damage pathway and the barrier and reaction energy were calculated to be 49 kJ mol^{-1} and –287 kJ mol^{-1} , respectively. The corresponding numbers for the overall reaction with the Cu(II)-oxyl (1) as reference were calculated to be 101 kJ mol^{-1} for the barrier and –234 kJ mol^{-1} for the reaction energy.³⁵ Clearly, this is still more favorable than the 199 kJ mol^{-1} and –44 kJ mol^{-1} in Fig. 9. Thus, the recombination reaction is more favorable than forming 3. We therefore speculate that if a histidyl radical is formed, it is more likely to lead to oxidative damage as suggested in ref. 35. We elaborate further on this in the following.

In Fig. 9 we compare different suggestions for protective mechanisms, illustrated by reactions I–III using the Cu(II)-oxyl (1) as reference. Note that we have included results involving His78 (*LsAA9*) and His86 (*TaAA9*) in the ESI (Tables S3 and S5†) but we will concentrate the discussion on His1 (this does not lead to any change of conclusions).

Reactions II and III were recently suggested as a protective hole-hopping mechanism.⁴² From Fig. 9, we can see that



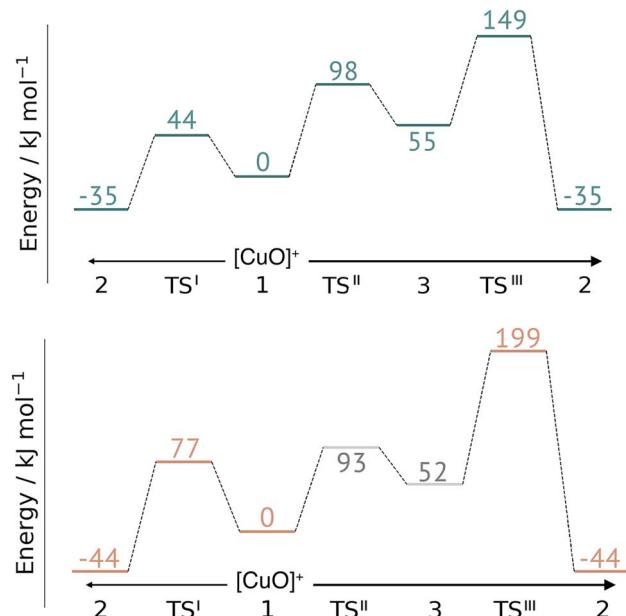


Fig. 9 Comparison of the full reaction diagrams for different suggested protective mechanisms for *TaAA9* (top) and *LsAA9* (bottom) using the Cu(II)-oxyl (1) as reference. Shown are only the most feasible energetics for B3LYP/def2-TZVPP (structures were optimized with TPSS/def2-SV(P)). Energies in gray are from ref. 35. Note that the reaction barrier TS^{III} was obtained with a bigger QM region.

reaction **I**, *i.e.*, the direct formation of the tyrosyl radical (**2**) is generally preferable to reactions **II** and **III**, where the formation **2** goes through a histidyl radical (**3**): the overall barrier for forming **2** *via* **3** is 199 kJ mol⁻¹ for *LsAA9*, using [CuO]⁺ (**1**) as a reference (see Fig. 9). The barrier is lower for *TaAA9*, but with 149 kJ mol⁻¹, it must still be considered too high to be feasible. If a histidyl radical as **3** is formed, we therefore consider it more likely that it is formed in a competing reaction with formation of a tyrosyl radical (**2**). Notably, the high overall barrier may be a result of starting from the Cu(II)-oxyl (**1**) species; we cannot exclude that another oxidative intermediate such as [Cu-OH]²⁺ or a free OH radical is responsible for initiating the oxidative damage, and we are currently investigating such alternatives.

In a broader perspective, our results provide a mechanistic explanation for that, similar to other oxidoreductases,⁴⁴ the LPMO active-site tyrosine is critical for initializing the protective hole-hopping mechanism. This tyrosine is widely conserved in most LPMO families, except for the majority of AA10 LPMOs, where it is replaced by phenylalanine.^{7,8,17,27,43} Our results can thus explain that fungal AA9 LPMOs are less prone to oxidative damage than their bacterial (AA10) counterparts as experimental studies recently have discovered.⁷³

In a previous paper⁴⁰ we compared calculated UV-vis spectra with the characteristic observed bands around 400–420 nm for **2**,^{36–38,42} but our calculations in ref. 40 were exclusively done on *LsAA9*. With our QM/MM optimized structure of **2** in *TaAA9*, we can now compare the corresponding UV-vis spectrum to the spectrum from *LsAA9* (see Fig. S4 in the ESI†). The calculated spectra are qualitatively similar displaying strong charge transfer transitions involving the tyrosyl radical close to 400 nm.

The most intense transitions in *TaAA9* are shown at 357 nm (3.48 eV) for the triplet and 386 nm (3.22 eV) for the open-shell singlet. The latter transition shows excellent correspondence with the most intense peak of the calculated open-shell singlet of *LsAA9* at 388 nm (3.20 eV).⁴⁰ Considering that we carried out the TD-DFT calculations as vacuum calculations, the intense peak from the open-shell singlet at 386 nm (3.22 eV) is in reasonable correspondence with the experimental room temperature absorption spectrum for *TaAA9* with the most intense peak at 420 nm (2.95 eV).^{36,42} The corresponding experimental value for *LsAA9* is 414 nm (2.99 eV)⁴² – other AA9 LPMOs show intense peaks in the same region (see ref. 37 and 38). For the triplet state of *TaAA9*, we also tried to calculate the spectrum of **2** including the electrostatics the enzyme as point-charges, but the effect of the point-charges on the position of the intense charge-transfer transitions from tyrosyl is minimal (see Fig. S5 in the ESI†).

Since we could qualitatively reproduce the UV-vis spectra of **2**, we additionally calculated the UV-vis spectrum of **3** for both *TaAA9* and *LsAA9* (Fig. S6 and S7 in the ESI†). Intermediate **3** has only been observed for *LsAA9* and the peak that experimentally is assigned the histidyl intermediate⁴² is obtained at 360 nm (3.44 eV). The TD-DFT calculations do predict intense transitions involving orbitals of histidyl character (further discussions are provided in the ESI†). However, these are at somewhat lower energies at 427 nm (2.91 eV) for the triplet and 406 nm (3.05 eV) for the open-shell singlet. The correspondence with the experimental values was clearly better for the tyrosyl radical (**2**). Intriguingly, the calculated spectrum for **3** in *TaAA9* is broader with two intense transitions at 342 nm (3.63 eV) and 395 nm (3.15 eV). Particularly, the former corresponds well to the observed band in *LsAA9*, but we cannot presently explain why the calculated spectrum for *TaAA9* fits better with the experimental spectrum of *LsAA9*. More benchmarks of the accuracy of TD-DFT – and preferably also investigations with multiconfigurational wave functions – will be required to confidently assign the spectrum of **3**. Along these lines, Zhao *et al.*⁴² find that 17% exact exchange is optimal to reproduce relative intensities for HERFD-XAS spectra with TD-DFT, but we have refrained from attempting to re-parameterize the functional at this point.

Conclusion

We have compared the initial steps of a protective hole-hopping mechanism of two LPMOs, namely *TaAA9* and *LsAA9*. The two investigated LPMOs have very similar active-site architectures, compared to other LPMOs (see, *e.g.*, ref. 7 and 34). We find similarities as well as remarkable differences in the protective mechanisms, highlighting that investigations on LPMOs as far as possible should consider several LPMOs. Regarding the similarities, our calculations show that the [CuO]⁺ moiety in the Cu(II)-oxyl intermediate (**1**) for both *TaAA9* and *LsAA9* is capable of oxidizing tyrosine to a tyrosyl radical. The electronic structures of the formed tyrosyl radicals are overall similar in the two LPMOs, based on their spin densities and their calculated UV-



vis spectra. The latter is commensurate with recent experimental investigations, with intense peaks around 400 nm.

We also find that the formation of a histidyl intermediate (3) has essentially the same reaction energy profile in the two LPMOs. However, it is more favorable energetically for both LPMOs to form the tyrosyl radical (2) than the histidyl radical (3). Overall, our results are commensurate with recent experiments showing that LPMOs without tyrosine are more susceptible to self-oxidation of the histidine brace.

The histidyl radical (3) has been proposed to be interconverted to the tyrosyl radical (2), making it part of the protective hole-hopping mechanism. The above comparison between barriers and reaction energies for the formation of 3 and 2 cannot be reconciled with this mechanism. In fact, we find that the formation of a tyrosyl (2) *via* a histidyl (3) intermediate is generally not feasible with a reaction starting from a Cu(II)-oxyl species (1). However, the LPMOs generally show quite different energetics regarding conversion between 3 and 2: for *TaAA9*, the conversion (the barrier between 3 and 2) is feasible, and cannot be entirely ruled out, although it is unlikely with the presently employed oxidizing species (1). Meanwhile, this conversion for *LsAA9* seems not to be possible, regardless of the oxidizing species, since conversion between 3 and 2 has a very high barrier. Another remarkable difference between the LPMOs is that the barrier for formation of the tyrosyl radical (3) is generally larger for *LsAA9*. A consequence of such differences may be that different LPMOs have different resistance toward oxidative damage. Investigation of other families (with and without a tyrosine residue close to copper) may reveal larger differences, and we are currently investigating this possibility.

Data availability

Structures of all intermediates are deposited in a zenodo repository: <https://doi.org/10.5281/zenodo.10495616>.

Author contributions

MMH: conceptualization, formal analysis, investigation (QM/MM), validation, visualization, writing – original draft, writing – review & editing; EKW: conceptualization, formal analysis, investigation (UV-vis), validation, participation in writing – original draft, writing – review & editing; EDH: conceptualization, formal analysis, funding acquisition, project administration, resources, supervision, validation, writing – review & editing.

Conflicts of interest

There are no conflicts to declare.

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