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### Introduction

Ab initio-based determination of lanthanoid– radical exchange as visualised by inelastic neutron scattering<sup>†</sup>

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Magnetic exchange coupling can modulate the slow magnetic relaxation in single-molecule magnets. Despite this, elucidation of exchange coupling remains a significant challenge for the lanthanoid(III) ions, both experimentally and computationally. In this work, the crystal field splitting and  $4f-\pi$  exchange coupling in the erbium-semiquinonate complex [ErTp<sub>2</sub>dbsq] (Er-dbsq; Tp<sup>-</sup> = hydro-tris(1-pyrazolyl) borate, dbsqH<sub>2</sub> = 3,5-di-tert-butyl-1,2-semiquinone) have been determined by inelastic neutron scattering (INS), magnetometry, and CASSCF-SO ab initio calculations. A related complex with a diamagnetic ligand,  $[ErTp_2trop]$  (Er-trop; tropH = tropolone), has been used as a model for the crystal field splitting in the absence of coupling. Magnetic and INS data indicate antiferromagnetic exchange for Er-dbsq with a coupling constant of  $J_{ex} = -0.23$  meV (-1.8 cm<sup>-1</sup>) (-2 $J_{ex}$  formalism) and good agreement is found between theory and experiment, with the low energy magnetic and spectroscopic properties well modelled. Most notable is the ability of the ab initio modelling to reproduce the signature of interference between localised 4f states and delocalised  $\pi$ -radical states that is evident in the Qdependence of the exchange excitation. This work highlights the power of combining INS with EPR and magnetometry for determination of ground state properties, as well as the enhanced capability of CASSCF-SO ab initio calculations and purposely developed ab initio-based theoretical models. We deliver an unprecedentedly detailed representation of the entangled character of  $4f-\pi$  exchange states, which is obtained via an accurate image of the spin-orbital transition density between the  $4f-\pi$ exchange coupled wavefunctions.

Single-molecule magnets (SMMs) are discrete molecular species that exhibit magnetic bistability at low temperature. Determined research effort has afforded significant advances in their properties since their discovery 30 years ago.<sup>1–3</sup> Consequently, the possibility of practical applications has become realistic, for example in data-storage, qubits, or as components in spintronics devices. Compounds based on the lanthanoid (Ln) ions show higher temperature magnetic stability, utilizing the intrinsically large magnetic moments and magnetic anisotropy, especially of the latter Ln(III) ions.<sup>4</sup> Recent advances have led to the observation of open magnetic hysteresis loops above liquid nitrogen temperature.<sup>5</sup> These include a family of "dysprosocenium" complexes, where slow magnetic hysteresis has been observed up to 80 K.<sup>6-8</sup> More recently, mixed-valence binuclear  $[(Cp^{iPr5})_2Ln_2I_3]$  ( $Cp^{iPr5} =$  penta-isopropyl-cyclopentadienyl) complexes have exhibited record breaking coercive fields of greater than 14 T below 60 K for the Dy analogue, boasting bulk properties to challenge traditional magnetic media.<sup>9</sup>

For Ln(m) single-ion magnets in the absence of exchange coupling, the magnetic properties stem from crystal field (CF) splitting of the ground spin–orbit coupled *J* electronic state due to the influence of the coordinated ligands. The relative order of energy levels and the purity of their composition can be tuned by an appropriate choice of ligand scaffold.<sup>4</sup> To access slow relaxation of magnetisation, a doubly degenerate ground electronic state is required, which is fulfilled for Kramers ions such

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<sup>&</sup>lt;sup>†</sup> Electronic supplementary information (ESI) available: Experimental details, crystallography details, IR, EPR, magnetometry, INS, *ab initio* details and calculations and Q-dependence model. CCDC 2287861–2287864. Computational output files are available at <a href="https://doi.org/10.26188/25222571">https://doi.org/10.26188/25222571</a>. For ESI and crystallographic data in CIF or other electronic format see DOI: <a href="https://doi.org/10.1039/d3sc04229d">https://doi.org/10.1039/d3sc04229d</a>.

as Dy(m) and Er(m); as well as stabilisation of an easy-axis ground state, *i.e.* a ground electronic state dominated by a large  $m_{\rm I}$  projection with unidirectional magnetic moment.<sup>10</sup>

Introducing exchange coupling with another magnetic species, such as an organic radical ligand, transition metal (TM), or a second Ln(m) ion, can be used to modulate the slow magnetic relaxation in Ln(m)-SMMs.<sup>11-15</sup> A small exchange bias may suppress quantum tunnelling in zero-field in some cases, allowing for observation of slow magnetic relaxation in zero-magnetic field, and has been touted as a method of improving SMM behaviour in Ln systems.<sup>16,17</sup> In contrast, weak coupling can instead provide alternative relaxation pathways that rapidly increase relaxation rates and hinder SMM performance.<sup>18</sup> The overall magnetic moment of the complex can also be increased by coupling multiple spins into a "giant spin", which while a common approach for TM-SMMs,<sup>19</sup> has not been as widely used for Ln-SMMs, due to their typically weaker magnetic exchange coupling.<sup>13,20-23</sup>

Spectroscopic techniques to measure the lowest lying electronic energy levels in Ln(III) compounds with large spin–orbit coupling include EPR,<sup>24–26</sup> luminescence,<sup>27–31</sup> far infrared,<sup>24,32,33</sup> and inelastic neutron scattering spectroscopies.<sup>34–38</sup> Inelastic neutron scattering (INS) offers unique advantages – it requires no applied magnetic field to identify magnetic transitions and doesn't rely on the often weak luminescence of Ln(III) ions.<sup>39</sup> By combining INS measurements with *ab initio* electronic structure calculations, the CF splitting in Ln(III) systems can be accurately determined.<sup>34,35,40</sup> For example, previous literature studies combining INS and EPR spectroscopies with theoretical calculations provided the basis for a Ln–Ln exchange modelling framework, utilising experimental determination of exchange free ground state properties on which to build the interaction model.<sup>41,42</sup>

Splitting of electronic states due to a magnetic exchange interaction between a Ln(III) ion and radical ligands can also be measured by INS, such as in the Tb(III) complexes  $[Tb(hfac)_3(2pyNO)]$  (2pyNO = tert-butyl 2-pyridyl nitroxide)<sup>43</sup> and [{[(Me<sub>3</sub>Si)<sub>2</sub>N]<sub>2</sub>(THF)Tb}<sub>2</sub>(µ-N<sub>2</sub>)].<sup>44</sup> Despite a few INS studies, experimental determination of the magnetic exchange coupling in Ln(m)-radical systems is typically only achieved unambiguously by magnetometry for spin-only Gd(m), due to the large unquenched orbital angular momentum of many Ln(III).45 Examples detailing the strength of exchange and modelling method have been tabulated in Table S1.† Approaches to Ln(III)radical complex design include coordination of stable radical ligands,46 utilizing diffuse radical orbitals to offer better spatial overlap with Ln(III) electron density<sup>47-49</sup> and electron/Ln encapsulation in endohedral fullerenes.<sup>50-52</sup> Recently, a benzene dianion containing a radical bridging ligand was developed by Gould *et al.* where the binuclear Gd(III) complex  $[Ln_2(\mu-BzN_6-$ Mes)] (BzN<sub>6</sub>-Mes = 1,3,5-tris[2,6-(N-mesityl)dimethanamino-4tert-butylphenyl]benzene) holds the current record for Gdligand coupling of  $J_{\text{ex}} = -43 \text{ cm}^{-1}$ .<sup>53</sup> Comparing these systems becomes challenging, as in some strongly coupled systems it is unclear whether there is mixed valency at the metal ion, formally bonding orbitals or a localised radical interacting via exchange pathways.

Despite recent increases in the magnitude of the magnetic exchange coupling in Ln(III) systems, computational studies pose significant challenges, with no robust approach for modelling active spaces beyond the 4f orbitals.<sup>9,13,54-57</sup> To develop a computational method for elucidating the exchange coupling in Ln–radical systems, combined experimental and computational studies are required to benchmark calculations.<sup>55,58</sup>

In this study, we have returned to the [LnTp<sub>2</sub>dbsq] family,46,59,60 with the aim of elucidating the low energy CF splitting and magnetic exchange coupling in a Ln-radical compound with large spin-orbit coupling through a combination of INS spectroscopy and electronic structure calculations. We have first determined the CF splitting of the diamagnetic ligand analogue [ErTp<sub>2</sub>trop] (tropH = tropolone; Er-trop)<sup>60</sup> using ab initio electronic structure calculations, which we experimentally verify by comparison with static magnetometry and EPR spectroscopy of Er-trop. We then use a combination of INS and electronic structure calculations on [ErTp<sub>2</sub>dbsq] (Er**dbsg**) to determine the CF splitting and Er-radical magnetic exchange coupling. The dynamic magnetic susceptibility data are compared for Er-trop and Er-dbsq. These systems feature very similar coordination environments, which is critical for exchange modelling, providing experimentally observable exchange-free properties to ensure the feasibility of the final exchange coupled calculation.

### **Results and discussion**

#### Diffraction and structure description

The compounds **Y-trop**, **Er-trop**, and **Er@Y-trop** are isomorphous with the previously reported structure of the Ho(m) analogue, and consistent with the reported cell parameters for the Er(m) analogue (Table S2†).<sup>60</sup> **Er-dbsq** is isomorphous with the reported **Y-dbsq** analogue.<sup>59</sup> The rare earth centres are all eight coordinate, with an {N<sub>6</sub>O<sub>2</sub>} coordination sphere arising from coordination to the two tripodal Tp<sup>-</sup> ligands and the bidentate trop<sup>-</sup> or dbsq<sup>-</sup> ligand (Fig. 1).

The coordination geometry of the rare-earth centre is closest to a square antiprism for all analogues (Fig. 1), according to continuous shape analysis (Table 1), performed with the SHAPE 2.1 software.<sup>61,62</sup> The Er( $\Pi$ ) centre in both **Er-trop** and **Er-dbsq** have a similarly small distortion index from an ideal square antiprism of 0.501 and 0.383, respectively. The similarity between the coordination geometry at the Er( $\Pi$ ) centre in the two analogues is further demonstrated by the Er–O and Er–N distances, with average Er–O distance of 2.280(3) Å and 2.295(2) Å, and an average Er–N distance of 2.478(7) Å and 2.482(5) Å for **Er-trop** and **Er-dbsq**, respectively.

Using the software Mercury,<sup>63</sup> a comparison of the two structures has been generated (Fig. S4†) by overlaying the Er and the coordinated atoms, which shows good agreement between the coordination geometry of the two structures. Due to the similarity between the coordination spheres of the two compounds, we have analysed the spectroscopic and computational data assuming the ligand geometry and therefore CF



Fig. 1 Structural representations of Er-trop and Er-dbsq from X-ray single crystal diffraction with polyhedral representation of the Er(III) coordination sphere. Selected hydrogen atoms have been omitted for clarity. Colour code: Er (pink), O (red), N (blue), B (orange), C (grey), H (white).

effects are the same, with the only difference being the unpaired electron on the dbsq<sup>•-</sup> ligand in the Er-dbsq analogue.

#### **Electronic structure calculations**

Ab initio electronic structure calculations were performed to obtain the electronic energy levels, g-tensors of the lowest lying (pseudo-)doublets and wavefunction decompositions for Ertrop and Er-dbsq in order to simulate the spectroscopic and magnetic properties (Tables 2 and S4-S7<sup>†</sup>).

The *ab initio* calculation for Er-trop predicts a ground Kramers Doublet (KD) with majority composition  $m_{\rm I}$  =

Table 1	Selected intra-atomic bond lengths and SHAPE <sup>61</sup> parameters
for Er-tr	op and Er-dbsq

Parameter	Er-trop	Er-dbsq
Distance/Å		
Intramolecular Er–O	2.280(3)	2.2846(15), 2.3058(14)
Intramolecular	2.489(4), 2.510(4),	2.505(2), 2.4787(18),
Er-N <sub>coordinated</sub>	2.434(4)	2.4415(18), 2.5206(18), 2.4406(18), 2.5042(18)
Intermolecular Er…Er	9.3261(4)	9.7648(6)
Angle/°		
O-Er-O	68.37(17)	70.02(5)
SHAPE parameters		
SAPR <sup>a</sup>	0.501	0.383
$\operatorname{BTPR}^{b}$	2.220	1.918

<sup>a</sup> Square antiprism. <sup>b</sup> Biaugmented trigonal prism.

Table 2 Calculated electronic energy levels and experimental energies obtained from INS measurements

Compound	Calculated energy/meV	Experimental energy/meV
Er-trop	0	a
Li dop	8.09	a
	11.3	a
Er-dbsq	0/0.29	0
•	1.27/1.28	0.918(1)
	7.84/7.99	4.36(2)
	8.66/8.67	
	10.84/11.13	9.55(1)
	11.20/11.43	
<sup><i>a</i></sup> Not determi	ned in this study.	

68%  $\pm 13/2$ , with the first excited KD lying 8.09 meV ( $65.2 \text{ cm}^{-1}$ ) higher with majority composition  $m_{\rm I} = 84\% |\pm 15/2\rangle$ . The second excited KD is predicted to lie 11.3 meV (91.1 cm<sup>-1</sup>) above the ground state and is in contrast highly mixed, with the largest contributions coming from  $m_{\rm I} = 19\% |\pm 1/2\rangle$ ,  $38\% |\pm 3/2\rangle$ ,  $18\% |\pm 5/2\rangle$  and  $18\% |\pm 11/2\rangle$ .

Electronic structure calculations on Er-dbsq indicate 32 lowlying electronic states (Table S6<sup>†</sup>). These states are separated by ca. 6000  $\text{cm}^{-1}$  to the next states and clearly arise from coupling of the ground  ${}^{4}I_{15/2}$  spin–orbit term (multiplicity = 16) with the radical. The energy levels calculated in the exchange coupled spectrum are consistent with splitting of the KD states into exchange-split pairs of pseudo-doublets arising from the strong anisotropy of the Er(m) ion. The two lowest lying energy levels have a splitting of 0.29 meV  $(2.35 \text{ cm}^{-1})$  and are separated from the first excited pseudo-doublet by around 1.25 meV ( $10 \text{ cm}^{-1}$ ).

The nature of these states is represented by the effective gtensors calculated for the pseudo-doublets (Table S11<sup>†</sup>) and their comparison with those obtained for the Er-trop CF states. The ground doublet for **Er-trop** yields an effective principal  $g_z$ value of 14.7, which when coupled with the radical, yields a ground pseudo-doublet with effective  $g_z = 12.9$  (ca. 2 less than  $g_z$  of Er-trop, indicating AF coupling) and a first excited pseudodoublet with  $g_z = 16.9$ , further details in ESI.<sup>†58</sup>

Subsequent energy levels appear in a pattern of pairs of pseudo-doublets centred at energies around the CF states observed for Er-trop, indicating this exchange-free analogue is a reasonable comparison for the exchange coupled Er-dbsq molecule. The third and fourth excited pseudo-doublets also yield  $g_z$  values consistent with isolated coupling of the KD from **Er-trop** ( $g_z = 17.3$ ) with a radical to give an AF doublet at lower energy  $(g_z = 14.8)$  and F doublet at higher energy  $(g_z = 19.1)$ . For the subsequent states, the trend in  $g_z$  values is no longer observed, likely due to the highly mixed wavefunctions with more pronounced rhombic  $g_{x/y}$  components.

These calculations allow predictions of INS transitions, with four possible excitations between states  $1 \rightarrow 3, 1 \rightarrow 4, 2 \rightarrow 3$ and 2  $\rightarrow$  4, centred around 1.1 meV (9 cm<sup>-1</sup>). If the exchange coupling between radical and 4f electrons was purely Ising-type, then only two of the four excitations would be magnetic-dipole allowed and hence observable in the INS spectrum. These would correspond to a radical-spin flip, transitioning from the antiferromagnetically coupled ground pseudo-doublet to the ferromagnetically coupled excited states.

#### Magnetic properties and EPR spectroscopy

Magnetic susceptibility and magnetisation data were acquired for **Er-trop** and **Er-dbsq** (Fig. 2). The magnetic susceptibility data for both analogues agree with that reported in literature.<sup>60</sup> For the **Er-trop** analogue, the experimental room temperature magnetic susceptibility temperature product ( $\chi_M T$ ) of 11.4 cm<sup>3</sup> K mol<sup>-1</sup> is in excellent agreement with the value of 11.3 cm<sup>3</sup> K mol<sup>-1</sup> simulated for a single Er(m) ion. With decreasing temperature, there is a gradual decrease in the  $\chi_M T$  product due to depopulation of excited CF states, until 2 K.

The magnetisation of **Er-trop** is not yet saturated at 7 T, reaching a value of 4.63  $\mu_{\rm B}$  at 2.0 K, near the expected value of 4.5 for an Ising type Er(m) ion. The excellent reproduction of the experiment from calculations suggests that the predicted composition and energy of the low-lying electronic energy levels are accurate. To confirm the calculated wavefunction composition of the ground CF state for **Er-trop**, a low temperature X-band EPR spectrum was obtained on a powder sample (Fig. 3). The EPR spectrum of **Er-trop** is broad with an intense  $g_{\parallel}$  feature and a broad  $g_{\perp}$ . The broad linewidths of the two spectral features are readily attributed to dipolar interactions between neighbouring molecules. To reduce this effect the diluted **Er@Y-trop** analogue was also measured, where the Er(m) ion is diluted 5% in the diamagnetic Y(m) analogue (Fig. S8†). The



**Fig. 2** Static magnetic susceptibility data measured with  $B_{dc} = 0.1 \text{ T}$  (black dots) and profiles simulated from electronic structure calculations (red lines) for **Er-trop** (top) and **Er-dbsq** (bottom). Inset: magnetisation data at the indicated temperatures (black data points) and profiles simulated from calculations (red lines).

EPR spectrum was simulated for  $S_{\rm eff}=1/2$  with the calculated ground CF state effective *g*-values, using the Hamiltonian:

$$\hat{H}_{\rm EPR} = \mu_{\rm B} B \cdot g_{\rm eff} \cdot S_{\rm eff} + A_{\rm eff} (I \cdot S_{\rm eff}) \tag{1}$$

where *I* is the nuclear spin of the EPR-active isotope of Er (<sup>167</sup>Er, I = 7/2). Analysis of the EPR spectra of **Er-trop** affords the three components of the *g*-tensor of the ground doublet, giving  $g_x$  unobserved,  $g_y = 0.83$ , and  $g_z = 12.3$  simulated with a Gaussian line width of 50 mT, and a H-strain on  $g_y$  of 2000 MHz. As evident from the simulation, the relative intensity of the features is not well reproduced, with the simulation parameters used to best describe the anisotropic *g*-values. Attempts to include Er hyperfine coupling did not improve the simulation, so they are disregarded here. This is reasonably consistent with the degree of the calculated axiality with values of  $g_x = 0.45$ ,  $g_y = 0.59$ ,  $g_z = 14.7$  for the ground CF doublet of **Er-trop** (Table S4†).

The spectrum obtained for Er@Y-trop shows similar spectral features which are far narrower than the original Er-trop sample. Here, the g-values determined are  $g_x$  unobserved,  $g_y =$ 0.83, and  $g_z = 14.0$ , with smaller broadening evident through the obtained Lorentzian linewidth and H-strain values used (see ESI<sup>†</sup>). Through comparison of these two data sets, it is clear that the pure Er-trop system suffers from severe dipolar broadening and as such it is unclear whether the shift in g-values is a consequence of the broadening or a physical change resulting from the different matrix structure obtained via doping into the Y-trop analogue.<sup>41</sup> Despite the near axial g-tensor determined from calculations Er-trop, the relative intensity of the peaks would be even worse when employing  $g_x = g_y = 0.83$ . The **Er@Y**trop data are well fit for relative intensity with a rhombic model with  $g_x$  unobserved and given the similar coordination environment, justifies the use of a rhombic model for Er-trop.

The  $\chi_{\rm M}T$  curve for **Er-dbsq** reaches a room temperature value of 11.7 cm<sup>3</sup> K mol<sup>-1</sup>, which is consistent with an uncoupled Er(III) ion and the s = 1/2 dbsq<sup>--</sup> radical, and in agreement with the simulated value of 11.7 cm<sup>3</sup> K mol<sup>-1</sup>. Below 150 K, there is a gradual decrease in the  $\chi_{\rm M}T$  product which is mostly reproduced by calculation, likely due to depopulation of excited electronic states, as in **Er-trop**. Below 15 K, there is a sharp decrease in both the experimental and simulated profiles,



**Fig. 3** Solid state X-band EPR spectra measured at 5 K for **Er-trop** (black, top) with simulated spectrum from electronic structure calculations as described in the text (red, bottom).

consistent with AF coupling between the Er(m) and dbsq<sup>•-</sup> ligand. The experimentally observed trend towards zero is captured by calculation and is likely indicative of the singlet nature of the ground state, with a splitting of 0.3 meV calculated within the ground pseudo-doublet.

The high-field magnetisation value for **Er-dbsq** is again consistent with AF coupling between Er(m) and the dbsq<sup>--</sup> ligand, reaching a value of 3.96  $\mu_{\text{B}}$  at 2 K and 7 T, lower than that expected for an isolated Er(m) and measured for **Er-trop**. The slight deviation between experiment and theory in the magnetisation curves indicates the ground state properties are not perfectly reproduced from calculation. Nevertheless, the consistency with the susceptibility indicates a good reproduction of the overall splitting of the low-lying energy states.

To probe the effect of Er(iii)-radical exchange coupling on the slow magnetic relaxation in this system, ac magnetic susceptibility measurements were performed. Slow magnetic relaxation in Er(III) compounds is not as widely observed as for Dy(III) compounds,<sup>64</sup> and requires ligand environments that best stabilise the more prolate electron density of the higher  $m_{\rm T}$ states, such as cyclooctatetraenyl ligands.65,66 A scan of the ac magnetic susceptibility for Er-trop at 2.0 K in varying applied dc fields (Fig. S18<sup>†</sup>) shows no out-of-phase component in  $B_{dc} = 0$  T. Two applied fields of  $B_{dc} = 0.075$  T and  $B_{dc} = 0.3$  T were chosen to measure the variable temperature ac magnetic susceptibility. These data were compared with the doped sample Er@Y-trop, to identify the impact of nearest neighbour interactions. Data were interpreted using a combined relaxation model accounting for Orbach, Raman, Direct and QTM processes where appropriate as described by eqn (2).

$$\tau^{-1} = \tau_0^{-1} \cdot \exp(^{-U_{\text{eff}}}/kT) + AT + CT^n + \tau_{\text{OTM}}^{-1}$$
(2)

where  $\tau_0^{-1}$  is the pre-exponential factor,  $U_{\text{eff}}$  is the energy barrier to magnetisation reversal, k is the Boltzman factor,  $\tau_{\text{OTM}}^{-1}$  is the rate of QTM, A the direct relaxation coefficient, C the Raman coefficient, and *n* the Raman exponent. The data are shown in Fig. 4 and fitted parameters presented in Table 3. As evident from the data, the maximal measured field of 0.3 T for Er-trop reveals a slower QTM rate at low temperature compared to 0.075 T and diluting the sample in a diamagnetic host removes any observable QTM within the frequency window measured. The Er-dbsq sample was also tested but showed no out-of-phase signal even with an applied field. The ab initio calculations predict a ground singlet state as the most likely cause for lack of SMM behaviour, rather than any spin-phonon coupling changes due to the different ligand. Similar behaviour has recently been observed in a family of pseudo-tetrahedral Er(III) complexes, where the introduction of a TEMPO radical ligand leads to a loss of the SMM behaviour.67 Full details of the ac studies are available in the ESI.†

#### Inelastic neutron scattering

In order to experimentally determine the electronic energy level splitting for **Er-dbsq**, we measured variable temperature INS of the partially deuterated compound **Er-dbsq**<sup>D</sup>. Spectra of **Y-dbsq**<sup>D</sup> were collected as a pseudo-diamagnetic analogue and



**Fig. 4** Relaxation rate plot for **Er-trop** in applied fields of 0.075 T (black open circles) and 0.3 T (black closed circles), and **Er@Y-trop** in an applied field of 0.1 T. Error bars are smaller than the data points. Lines are best fits to the data as described in the text.

were subtracted from the **Er-dbsq**<sup>D</sup> spectra. The corrected variable temperature INS spectra of **Er-dbsq**<sup>D</sup> are presented in Fig. 5. Four peaks are visible at both positive and negative energy transfer. Three peaks, a, I, and II, are assigned as cold transitions on the neutron energy loss side of the INS spectrum, and the corresponding peaks a', I', and II' appear with increasing temperature on the neutron energy gain side of the spectra. An additional peak is observed, marked with an asterisk in Fig. 5. This peak can be unambiguously assigned as a phononic peak from analysis of the *Q*-dependence, which follows the  $Q^2$  dependence typical of a phonon (Fig. S15†). We therefore exclude this excitation from our further analysis.

The three remaining excitations are assigned as magnetic from analysis of their *Q*-dependence, with I and II following the magnetic form factor dependence ( $I \propto F^2(Q)$ ) expected of a CF transition of Er(m). Additional confirmation of the three magnetic excitations can be obtained by analysis of the generalised phonon density of states (ESI<sup>†</sup> discussion).

The CF peak I is a weak peak observed at E = 4.36(2) meV, most apparent on the neutron energy loss side of the variable temperature  $\lambda = 2.34$  Å spectra. Peak positions (and corresponding errors) are determined from fitting the 1.5 K data on the neutron energy loss side. Peak I has the temperature dependence expected of a magnetic excitation from the ground state – decreasing in intensity with increasing temperature as

 Table 3
 Magnetic relaxation parameters obtained from fitting of the ac magnetic susceptibility data of Er-trop and Er@Y-trop

	Er-trop		Er@Y-trop
$B_{\rm dc}$	0.075 T	0.3 T	0.1 T
$C/s^{-1} K^{-n}$	0.06(2)	0.22(5)	0.033(5)
n	7.5(1)	6.8(1)	7.9(1)
$\tau^{-1}_{\mathrm{OTM}}/\mathrm{s}^{-1}$	$2.69(9)  imes 10^4$	$6.23(8)  imes 10^{3}$	_
$A/s^{-1} K^{-1}$	_	_	30(3)



Fig. 5 Variable temperature INS spectra of Er-dbsq<sup>D</sup> corrected for the diamagnetic contribution with subtraction of Y-dbsq<sup>D</sup> data, measured with  $\lambda = 4.69$  Å and integrated over 0.28 Å<sup>-1</sup>  $\leq Q \leq 1.95$  Å<sup>-1</sup> (left), and  $\lambda = 2.34$  Å integrated over 1.30 Å<sup>-1</sup>  $\leq Q \leq 3.65$  Å<sup>-1</sup> (right). The intensities at the two wavelengths are normalised to the elastic line intensity over the corresponding *Q* range. Solid lines are a sum of Gaussian fits to the observed peaks.

the ground state depopulates, while the corresponding peak I' increases in intensity with increasing temperature. The next lowest energy CF transition II, at E = 9.55(1) meV, is observed on the neutron energy loss side of the spectrum in the  $\lambda = 2.34$  Å spectra, with a corresponding peak II' observed on the neutron energy gain side of the spectrum at elevated temperatures. As with the other CF peak, the temperature dependencies of II and II' are consistent with magnetic transitions to and from the ground state and are assigned as CF transitions to a second excited CF state.

The lowest energy peak a (and a') is observed at E = 0.918(6)meV (and E = -0.918 meV). Peak a has a temperature dependence expected of a magnetic excitation from the ground state. From inspection of the CF splitting of the Er(m) determined for Er-trop, it is unlikely that this excitation stems from a CF-like transition. Additionally, the magnitude of the magnetic exchange splitting is consistent with that observed in Er-dbsq from calculations (Table 2). The Q-dependence of an exchange transition involving two exchange-coupled metal ions is expected to follow a dependence that relies on the magnetic form factors of the two interacting spins, with an interference term, typically giving rise to a pronounced oscillatory behaviour. However, for the very common case of a radical that is coupled to a metal ion with strong axial anisotropy via a pure Ising exchange mechanism, such a transition would be dominated by a radical spin-flip, hence the Q-dependence would be dominated by the form factor of the radical spin only. The Qdependence of peak a, while not inconsistent with an exchange interaction, shows neither the typical oscillations in intensity with Q observed for dinuclear metal complexes in literature,68-70 nor the fast-decaying behaviour expected for a pure radical spin flip transition i.e. for an Ising 4f-radical exchange mechanism (Fig. S12<sup>†</sup>). The jump diffusion model used for intermolecular exchange in a Co(II)-semiquinonate compound<sup>71</sup> shows reasonable agreement when calculated with a distance of 12 Å,

however, this spacing is not consistent with the metal-radical separation nor the intermolecular separation of metal ions. As the form factor falloff for a delocalised organic radical ligand is expected to be very different to that of the radially contracted forbitals of a Ln(m) ion, the interference mechanism between the two coupled spin moieties is not simple to predict without further detailed modelling. While the transition can be tentatively assigned as a magnetic exchange transition – a transition between the AF and F states from the ground CF doublet of Er(m) split by the magnetic exchange interaction with the dbsq<sup>--</sup> radical, additional modelling of the experimental *Q*-dependence will be presented following the electronic structure calculations to better qualify this assignment.

#### Electronic structure of Er-dbsq

First, we analyse the magnetic exchange coupling in **Er-dbsq**. Upon assignment of the INS peak a, we have an exchange gap energy  $\Delta E = 0.918(6)$  meV. The data available allow for only the isotropic exchange coupling to be determined. Magnetic exchange coupling in orbitally degenerate lanthanoid ions can be treated using several methods. We begin by using the Lines model,<sup>72</sup> in which the true spin of the Ln(m) ion in the  $|L,m_{L}, S,m_{S}\rangle$  basis is used. The Lines model has been widely used for modelling exchange interactions in Ln(m) systems.<sup>73–77</sup> In the case of **Er-dbsq**, we couple the  $S_{\rm Er} = 3/2$  of Er(m) to the  $S_{\rm SQ} = 1/2$  of the dbsq<sup>--</sup> ligand with the Hamiltonian:

$$\hat{H}_{\rm ex} = -2J_{\rm ex}\hat{S}_{\rm Er}\cdot\hat{S}_{\rm SQ} \tag{3}$$

Using  $\Delta E = -4J_{ex}$ , a  $J_{ex} = -0.23$  meV is obtained. This, and the CF parameters (Table S5†) obtained from the electronic structure calculations on **Er-trop** to describe the CF splitting of the Er(III) ion are then used to simulate  $\chi_M T$  and M vs. H, shown in Fig. 6. The Lines model provides a reasonable description of the general temperature dependence of the magnetic

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susceptibility but fails to accurately describe the lowest energy magnetic properties correctly. The insufficiency of the Lines model has been discussed for several Ln(III) systems.74,75,78,79

We then analyse the magnetic exchange coupling using a pseudospin 1/2 model, where the Er(III) is treated as an effective  $S_{\text{eff}} = 1/2$  spin to couple to the dbsq<sup>•–</sup>  $S_{\text{SQ}} = 1/2$ , with  $g_{\rm eff}$  values of the ground pseudo-doublet as measured by EPR used to model the anisotropy at the Er(III) centre ( $g_x = 0.01, g_y =$ 0.8, and  $g_z = 12.3$ ). Here we consider the exchange with the ground pseudo-doublet only  $(\hat{H}_{1/2} = -2J_{1/2}\hat{S}_{\text{eff}}\cdot\hat{S}_{\text{SO}})$ . We use  $J_{1/2}$ = -0.46 meV from the observed INS transition and the experimentally determined  $g_{\text{eff}}$  values from EPR measurements. The simulations of the magnetisation curves and the low temperature  $\chi_{M}T$  are in excellent agreement with the experimental data, suggesting that the calculated isotropic exchange coupling correctly describes the exchange coupling between the ground CF doublet of Er(III) and dbsq<sup>•–</sup>.

The overall magnitude of the exchange coupling ( $J_{ex} = -0.23$ meV =  $-1.8 \text{ cm}^{-1}$ ) is similar to the Gd(m) analogue ( $J_{ex}$  = -5.7 cm<sup>-1</sup>) determined previously.<sup>46</sup> One would expect that varying the Ln(m) would change the strength of the coupling within an isostructural series, and this has indeed been observed both computationally and experimentally.49,58 Often the magnetic exchange coupling in the majority of Ln-radical systems is only determined for the Gd(m) analogue - these results highlight the need to experimentally verify the magnetic exchange coupling for each Ln(m). It is worth noting that these

current INS measurements cannot be used to unambiguously determine the sign of the magnetic exchange coupling, which instead has been determined from magnetometry. To confirm this by INS, either single crystal measurements or INS in an applied magnetic field would be required.<sup>80-83</sup>

Looking to the simulated electronic energy levels calculated using PHI,84 the exchange coupling within the Lines model results in the splitting of each KD of the Er(m) (as determined by calculation for Er-trop) into two sets of doublets. This is consistent with the assumption that the coupling is in the weak exchange limit. The Lines model and the energy level splitting from electronic structure calculations on Er-dbsq are presented in Fig. 7. From comparison of the experimentally determined splitting by INS, the origin of the INS transitions can then be assigned.

Within the Lines model, the small exchange splitting dictates that peak I (and I') must be a transition between the ground CF state, which in **Er-trop** is majority  $m_{\rm I} = |\pm 13/2\rangle$  and the next CF state, of majority composition  $m_{\rm I} = |\pm 15/2\rangle$ , consistent with INS selection rules ( $\Delta m_{\rm I} = 0, \pm 1$ ), Fig. 8. We assign the higher energy peak II (and II') as a transition between the ground CF state and the next highest CF state, originating from the 2nd lowest excited KD in Er-trop, which is a state of highly mixed wavefunction composition with an 18% contribution of  $m_{\rm I} = |\pm 11/2\rangle$ , which would be allowed by INS selection rules. Worth noting is that a transition between the first excited CF level and the second CF level is not allowed by the INS selection rules, and as such is not expected to be observed, even with thermal population of these states. Within the Lines model, one would expect the transitions between the lower energy AF state of each CF state, as well as transitions between the higher energy F states of each CF state. As the energies of



CF pseudo-doublets as calculated for Er-trop (grey); electronic energy level splitting calculated for Er-dbsq as described in the text with the Lines model (blue); experimentally determined splitting for Er-dbsg with arrows indicating observed cold INS transitions (black with dark red arrows); and singlets calculated ab initio for Er-dbsq (red). Wavefunction compositions ( $m_{\rm J}$  contributions greater than 10%) are indicated for Er-trop.



Fig. 6 Comparison of experimental static magnetic data for Er-dbsq with the ab initio calculated profiles (red dashed lines), the Lines model (blue dotted lines) and pseudospin model (green lines) as described in the text.

calc

these transitions (AF  $\rightarrow$  AF and F  $\rightarrow$  F) would be approximately the same energy in the case of isotropic exchange, this is unlikely to be observed as additional splitting of peaks in the INS spectrum.

The *ab initio* calculated splitting gives a similar overall profile to the Lines model, splitting the KDs into pairs of pseudo-doublets and we again assign the transitions as above. However, as some doublets from CASSCF show significant splitting, the observation of distinct transitions would now be apparent.<sup>34,58</sup>

To further compare the experimental INS data with the calculated properties of **Er-dbsq**, the INS spectra are simulated at 50 K and shown in Fig. 8. For the electronic structure calculations, the magnetic moment matrix element components as calculated for the spin-orbit basis from the Single\_Aniso module are used to approximate the transition probabilities. Each transition is then weighted by the Boltzmann population of the initial state at 50 K and convoluted with the instrument resolution at that energy.

The simulated spectrum at 50 K reveals transitions at 1 and 1.3 meV, and two regions of moderate intensity around 8 meV and 11 meV. The **Er-dbsq** calculation produces a spectrum with similar features to the experimental data, giving an excellent agreement with the lowest lying peak a and good agreement with peak II. However, the transitions corresponding to peak I



**Fig. 8** Experimental INS spectra of **Er-dbsq**<sup>D</sup> (top, green), spectrum simulated from electronic structure calculations (middle, red), spectrum simulated from Lines model (blue), all at 50 K. Grey curves are individual Gaussian contributions.

do not align well with the experimental data, with peak I predicted at higher energy than experimentally observed.

For the Lines model, the transition probabilities calculated in PHI<sup>84</sup> are weighted by the Boltzmann population of the initial state and convoluted with the instrument resolution. For the Lines spectrum, it is evident that this simple isotropic Lines model displays shortcomings in the simultaneous description of the transition probabilities of the exchange transitions and the CF transitions, in that it predicts larger transition probabilities for the CF transitions than for the exchange transitions, as one would naively expect if CF transitions were solely dominated by the  $13/2 \rightarrow 15/2$  components, while the exchange transitions were dominated solely by a radical spin flip process. However, as the analysis of the Q-dependence of the exchange transitions will show, the exchange states are of a more mixed 4f-radical nature, a feature that can only be recovered in the simplest approximation by relaxing the requirement of an isotropic Lines model (vide infra).

Both models predict a larger overall electronic energy level splitting than observed with similar levels of agreement with the experimental energies and intensities. As the transitions in **Er-dbsq** are at characteristically similar energies to the CF states of **Er-trop**, it implies the CF splitting energy is overestimated by CASSCF-SO calculations. This overestimation of the overall electronic energy level splitting found by calculation has been observed for other Ln(m) compounds, particularly Er systems, with the *ab initio* CF splitting sometimes adjusted with a scaling factor to accurately reproduce the experimental data.<sup>35</sup>

The final comparison necessary with computational results is to simulate a theoretical *Q*-dependence resulting from the interaction of the 4f and radical electrons. In an attempt to interpret the *Q*-dependence of the low-lying feature, we first introduce an idealised pure Ising exchange model. The major  $m_j$ contributions from the **Er-trop** calculation are used to build a Hilbert space to model only the lowest two pseudo-doublets arising from the F and AF coupling with the radical (Fig. S28; see ESI† Model 1 for full details). From this idealised model, we observe that a transition involving solely a radical flip cannot be responsible for the experimental *Q*-dependence measured, as by virtue of the delocalised character of the  $\pi$ -spin density, this would fall off too rapidly with respect to what is observed (Fig. S29†).

Extending to an *ab initio*-based model Hilbert space (see ESI<sup>†</sup> Model 2 for full details), we now consider the mixed nature of the wavefunction, where *ab initio* results for **Er-trop** showed only  $68\%|\pm13/2\rangle$  contribution to the ground state. We firstly reproduce the nature of the lowest lying four states to focus on the *Q*-dependence of those transitions. This is done by introducing anisotropic exchange parameters in the  $|m_j, m_s\rangle$  product basis in order to reproduce the energy splitting of the *ab initio* exchange manifold more accurately than with the previous Lines model. We note that this model Hamiltonian results in a qualitatively similar wavefunction decomposition of the ground and excited pseudo-doublets (Fig. S30<sup>†</sup>). Full details for the calculation of the magnetic neutron scattering form factor arising from this model are described in the ESI.<sup>†</sup> **Chemical Science** 



Fig. 9 Calculated form factor  $|F_{tot}(Q)|^2 = |\tilde{F}_{dbsg}(Q)|^2 + |\tilde{F}_{Er}(Q)|^2 +$  $2\text{Re}(\tilde{F}_{dbsg}(Q)\tilde{F}_{Fr}^{*}(Q))$  estimated using eqn S10† and Model 2 for the 1  $\rightarrow$ 3, 4 transition (black line) compared to experimental data (black circles), Er(III) only form factor (red line) and dbsg<sup>•-</sup> only form factor (green line).

Arising from the model is an interference term that originates due to contributions from both the Er(III) and radical electrons. As such, the calculated form factor for transitions between the lowest exchange states now show two qualitatively new features: (i) a much slower drop-off than would be expected for purely radical or purely Er(m) based excitations, and (ii) a maximum peak predicted around  $Q = 0.75 \text{ Å}^{-1}$ , similar to the experimental data (Fig. 9 and S33<sup>†</sup>).

It should be stressed that this resultant interference bears the fingerprint of an intrinsically non-Ising exchange interaction between the radical and 4f electrons, which results in a wavefunction where  $m_s$ -spin states on the radical are entangled with the mixed m<sub>i</sub>-superpositions describing the KDs of the Er(III) ion. Thus, even for the relatively weak exchange observed here (cf. literature Ln-radical exchange Table S1<sup>†</sup>), a far more complex mixing of states is occurring between the Er and radical orbitals, which is contrary to an expected Ising-like exchange for a 4f- $\pi$  system. The relative complex phases of the wavefunction coefficients expanded on the  $|m_i, m_s\rangle$  product basis encoded in the transition reduced density matrices  $\sigma_{m'_s m_s}^{\text{dbsq,AF} \rightarrow \text{F}}$  and  $\sigma_{m'_j m_j}^{\text{Er,AF} \rightarrow \text{F}}$  (eqn S9†), determine a non-trivial interference term  $2\text{Re}(F_{\text{dbsq}}(Q)F_{\text{Er}}^*(Q))$ , resulting in the predicted oscillatory behaviour.

### Conclusions

Herein, we report a combined spectroscopic, computational and magnetometric study of the two related compounds Er-trop and Er-dbsq. The energy level splitting of Er-trop with the diamagnetic ligand and radical ligand analogue Er-dbsq have been calculated ab initio. The simulated magnetic and spectroscopic properties from the calculations are in good agreement with experiment and allow for the determination of the lowest energy electronic levels in these compounds. Additionally, the Er(m)-semiquinonate magnetic exchange coupling was directly determined to be  $J_{\rm ex} = -0.23$  meV from INS and magnetic measurements. The agreement of this exchange coupled model with experiment, and the similarity with the CASSCF-SO calculations, indicates accurate determination of the magnetic exchange coupling in this low-symmetry compound. This result highlights the capability of ab initio computations to handle the exchange between a Ln(m) ion with significant spin-orbit coupling and an organic radical ligand.

We show slow magnetic relaxation for Er-trop is dominated by Raman relaxation and QTM, where the QTM is related to the dipolar interactions with neighbouring molecules. In the Er**dbsq** analogue the slow relaxation is quenched completely by exchange coupling with the radical ligand, due to the loss of ground state degeneracy. The moderate coupling strength here is evidently beyond the weakly coupled limit, which can improve SMM behaviour through introduction of an exchange bias.

Spectroscopic determination of lanthanoid-radical magnetic exchange coupling is rare, as is the use of *ab initio* electronic structure calculations on systems with coupling to orbitally degenerate Ln(III) ions. Comparison of the calculated magnetic and spectroscopic properties from the Lines and pseudospin models with those simulated from full CASSCF/ RASSI-SO calculations on the Er-dbsq system reinforce the need for high level calculations to accurately reproduce the low temperature properties of Ln(III)-radical systems. This work represents the first semi ab initio reproduction of the oscillatory Q-dependence behaviour for a Ln-radical molecule, which verifies the non-Ising-like, high degree of mixing between the Er 4f and radical  $\pi$  electrons. Thus, for the first time an accurate image of the transition spin and orbital current density triggered by the INS excitation has been determined for a Lnradical system.

### Data availability

Crystallographic data for all compounds have been deposited with the CCDC under deposition numbers 2287861-2287864 and can be obtained from https://www.ccdc.cam.ac.uk/. Computational output files are available at https://doi.org/ 10.26188/25222571. Other data are available from the authors upon reasonable request.

### Author contributions

M. A. D. prepared and measured the samples and wrote the draft manuscript. M. J. G. performed the calculations and spectral interpretation with support from A. S. and S. C. A. S. developed the models for Q-dependence simulation and interpretation of the INS exchange transitions. L. S. performed EPR and modelling interpretation, A. M. K.-H. prepared the deuterated ligands. R. A. M., M. A. D. and S. C. performed the INS measurements and interpreted the data. R. A. M., A. S. and C. B. supervised the project.

### Conflicts of interest

There are no conflicts to declare.

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