


Cite this: *RSC Adv.*, 2024, 14, 29113

# Cholesterol inhibits oxygen permeation through biological membranes: mechanism against double-bond peroxidation†‡

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The presence of oxygen molecules ( $O_2$ ) in biological membranes promotes lipid peroxidation of phospholipids with unsaturated acyl chains. On the other hand, cholesterol is considered to be an antioxidant molecule as it has a significant barrier effect on the permeation of  $O_2$  across membranes. However, a comprehensive explanation of how cholesterol affects the distribution and diffusion of  $O_2$  within lipid bilayers is yet to be established. In this study, we investigated the interaction of oxygen molecules with polyunsaturated lipid bilayers using molecular dynamics (MD) simulations. The degree of lipid unsaturation and the concentration of cholesterol were varied to study the permeation of  $O_2$ . The free energy profile of  $O_2$  diffusing from the water phase to the lipid bilayer was calculated using biased umbrella MD simulations. The results show that  $O_2$  passively translocates into the membrane without changing the physical properties of the bilayer. Interestingly, in the unsaturated lipid bilayers the presence of cholesterol led to a significantly decreased permeation of  $O_2$  and an increase in the lipid chain order. Our results indicate that the hydroxyl groups of cholesterol strongly interact with the  $O_2$  molecules effectively inhibiting interactions between the oxygens and the double bonds in unsaturated lipid tails. In addition, a linear relationship between permeation and the ratio of membrane thickness and area per lipid was found. These insights can help our understanding of how the degree of unsaturation in a lipid tail and cholesterol affect lipid peroxidation at the molecular level.

Received 4th July 2024  
Accepted 23rd August 2024

DOI: 10.1039/d4ra04846f

rsc.li/rsc-advances

## 1 Introduction

Biological membranes separate internal and external environments of cells and organelles regulating the uptake/excretion of nutrients/waste.<sup>1,2</sup> They are made primarily of a lipid bilayer

with proteins anchored on it. Polar lipid head groups in these bilayers are exposed to internal/external environments with their acyl tails forming a non-polar core in between these environments. This creates a barrier that impedes the flow of polar molecules but allows the permeation of small non-polar molecules that are critical for life.<sup>3</sup>

Mechanisms by which molecules transport across cellular membranes can be classified into two main categories, active and passive.<sup>3</sup> Active transport involves regulatory machinery which requires energy to transport target molecules in the opposite direction of their concentration gradient. Passive transport is an entropy-driven process of transporting molecules across cellular membranes. Cross-membrane transport of hydrophilic species is typically an activated process which requires assistance of specialized protein-based transporters or ion channels.<sup>4–6</sup> Most small neutral molecules and drugs are transported across biological membranes passively.<sup>2,3,7,8</sup> This includes oxygen ( $O_2$ ) that diffuses across cell membranes with a permeability coefficient of about  $23 \text{ cm s}^{-1}$ . The passive diffusion of oxygen through the cell membrane is, however, strongly affected by lipid composition, which may help to rationalize the broad spectrum of oxygen levels observed across different cells/tissues.<sup>9</sup> Here, we focus on providing insights into the atomic mechanisms accounting for oxygen permeation

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† Videos of POPC bilayers with  $O_2$  molecules both with and without cholesterol are available at DOI: <https://doi.org/10.5281/zenodo.12570598>.

‡ Electronic supplementary information (ESI) available: The trajectories of 16 oxygen ( $O_2$ ) in lipid bilayers (Fig. S1–S6) and the density profiles of lipid and cholesterol in bilayers (Fig. S7) are provided. See DOI: <https://doi.org/10.1039/d4ra04846f>


and the role played by cholesterol and unsaturated fatty acids in this process.

Cellular oxygenation is crucial for brain function,<sup>9</sup> wound healing,<sup>10</sup> anesthesia,<sup>11</sup> adipose tissue dysfunction,<sup>12</sup> reperfusion injury,<sup>13</sup> and tumor radiation therapy.<sup>14</sup> This process is so important that oxygen concentrations below normal levels have been related to pathologies such as cancer and heart diseases. In the newly emerging cold (room temperature) plasma therapy, a partially ionized gas composed of reactive oxygen and nitrogen species (RONS) is directly delivered to living tissues, and it is being explored for treatment of skin related disorders,<sup>15,16</sup> cancer therapy,<sup>17–19</sup> and wound healing.<sup>20,21</sup> Among the RONS that cold plasma produces, is oxygen in its singlet delta state,  $O_2(^1\Delta_g)$ ,<sup>22</sup> which has been suggested to have a major role in cell apoptosis.<sup>23–25</sup> Production of singlet delta oxygen (among other reactive oxygen species) is also the basis of photodynamic therapy in which a light source and a photosensitizer are used in combination with oxygen to kill malignant cells.<sup>26,27</sup>

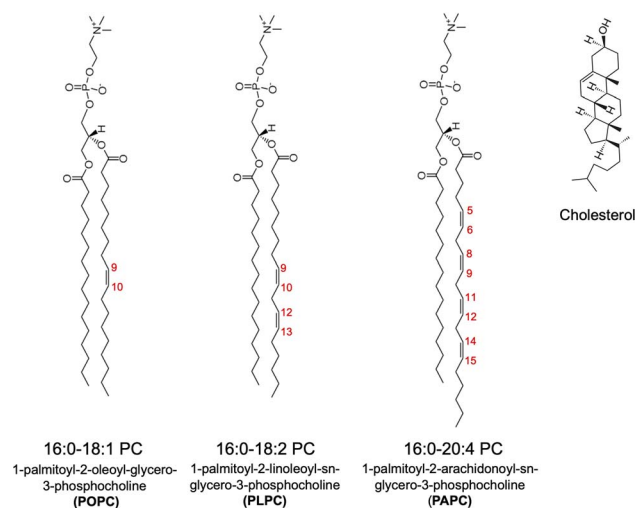
Cellular membranes are composed of a diverse set of phospholipids. They are versatile molecules with different headgroups, charge, chain length, and saturation.<sup>1</sup> Phosphatidylcholine (PC) is the main phospholipid in all mammalian cells (40–50%) and lipoprotein particles.<sup>28–30</sup> The double bonds in the unsaturated acyl chains are prone to oxidative stress caused by reactive oxygen species.<sup>31–33</sup> These reactions typically modify their chemical structure which consecutively disturbs the physiological properties of cellular membrane.<sup>34–36</sup>

In this work, we employ classical molecular dynamics (MD) simulations to study the mechanisms of permeability of  $O_2$  in six different phospholipid-based cell membrane models (details in Section 2). Although, singlet delta state of  $O_2$  is not described explicitly in our simulations, our study provides insights in the general mechanisms of its diffusion across cell membrane. We investigate how the  $O_2$  molecules interact with the double bonds of the lipid chains by exploring its location and distribution in the different lipid bilayers. As cholesterol has been suggested to protect lipid bilayers from oxidative stress,<sup>37,38</sup> we investigate the effect(s) of cholesterol on the distribution and diffusional pathways of  $O_2$  in the lipid bilayers. In addition, biased umbrella sampling MD simulations were used to calculate the free energy profiles of  $O_2$  from the water phase to the lipid bilayers.

## 2 Methods

### 2.1 Unbiased MD simulations

To study the permeability of  $O_2$  in model cell membranes, we considered phospholipid bilayers at three different saturation levels, namely 1-palmitoyl-2-oleoyl-glycero-3-phosphocholine (POPC), 1-palmitoyl-2-linoleoyl-sn-glycero-3-phosphocholine (PLPC), and 1-palmitoyl-2-arachidonoyl-sn-glycero-3-phosphocholine (PAPC). The numbers of double bonds in POPC, PLPC, and PAPC lipids are 1, 2 and 4, respectively. In addition, binary mixtures of the above lipids and cholesterol (CHOL) at 50% CHOL concentration were studied to examine the effect(s) of cholesterol on the permeability of  $O_2$ . The chemical



**Fig. 1** Chemical structures of the PC lipids (POPC, PLPC, and PAPC) and cholesterol. There are 1, 2 and 4 double bonds in the sn-2 chains of POPC, PLPC, and PAPC, respectively.

**Table 1** List of the simulated systems showing the total numbers of phospholipids, cholesterol (CHOL), water (SOL), and oxygen molecules ( $O_2$ )

System	Lipid	Cholesterol	Water	$O_2$
POPC	128	0	6400	16
PLPC	128	0	6400	16
PAPC	128	0	6400	16
POPC:CHOL	128	128	12 800	16
PLPC:CHOL	128	128	12 800	16
PAPC:CHOL	128	128	12 800	16

structures of lipids and cholesterol are shown in Fig. 1 and the details with the numbers of molecules for each simulated system are listed in Table 1. All MD simulations were performed with the GROMACS software package version 5.1.2 (ref. 39) using the CHARMM36 all-atom force field.<sup>40</sup> The TIP3P water model<sup>41,42</sup> was used in all simulations.

First, solvated lipid bilayers were constructed using the CHARMM-GUI Membrane Builder.<sup>43</sup> The pure bilayers had 128 phospholipid molecules and 6400 water molecules, and the cholesterol-containing ones had 128 lipids, 128 CHOL molecules and 12 800 water molecules. For all the systems, an equilibration simulation was run for 100 ns to optimize the bilayer structure.

After construction of the bilayers,  $O_2$  molecules were randomly inserted into the equilibrated systems. Energy minimization using the steepest descents algorithm was performed, followed by MD simulations under a constant particle number, pressure, and temperature conditions (the *NPT* ensemble). The pressure was set to 1 bar with semi-isotropic pressure coupling, time constant of 5.0 ps and compressibility of  $4.5 \times 10^{-5} \text{ bar}^{-1}$  by using the Parrinello–Rahman method.<sup>44</sup> The temperature was set to 310 K with a time constant of 1.0 ps using Nosé–Hoover thermostat technique.<sup>45–47</sup> Periodic boundary conditions were applied in all directions. A cutoff distance of 1.2 nm



was used for the van der Waals and the real space part of the electrostatic interactions. The Particle-Mesh Ewald (PME) method<sup>48,49</sup> was applied to calculate the long-range part of the electrostatic interactions. The Lennard-Jones forces were smoothly switched to zero between 1.0 and 1.2 nm. All bond lengths were constrained by the parallel linear constraint solver (P-LINCS) algorithm.<sup>50</sup> The integration time step was set to 2 fs, and the simulations were run up to 1  $\mu$ s. The Visual Molecular Dynamics (VMD) software<sup>51</sup> was used for molecular visualizations.

## 2.2 Biased umbrella MD simulations

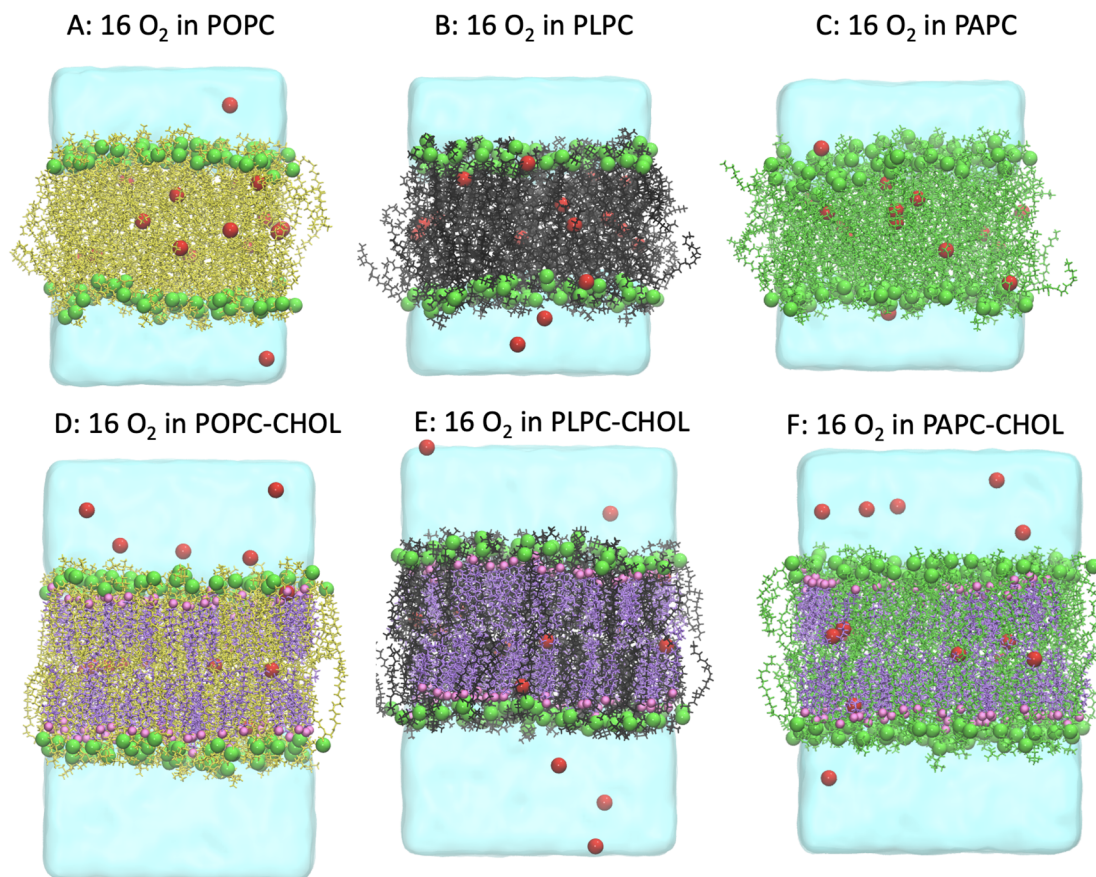
We computed the free energy profiles for O<sub>2</sub> transport in all the six bilayers. The potential of mean force (PMF) for moving an O<sub>2</sub> molecule from the water phase to the bilayer center was computed using the biased umbrella sampling MD simulations<sup>52</sup> with the Weighted Histogram Analysis Method.<sup>53,54</sup>

Initially, an O<sub>2</sub> molecule was placed in the water phase at 3.4 nm away from the lipid bilayer center. During the umbrella sampling MD simulations, the distance between the center of mass of the O<sub>2</sub> and the lipid bilayer was constrained in the *z*-direction (reaction coordinate) using a harmonic potential

function with a force constant of 1000 kJ (mol nm<sup>2</sup>)<sup>-1</sup>. In the simulations, 35 windows with varying positions of O<sub>2</sub> were performed to compute the PMF profile for each bilayer. The calculation was divided into 0.1 nm per window in the *z*-direction to pull O<sub>2</sub> from the water phase (*z* = 3.4 nm) to the bilayer center (*z* = 0 nm). All simulations were performed in the NPT ensemble at a constant of temperature at 310 K and a constant of pressure at 1 bar. Each window was run for at least 100 ns, and the last 50 ns were used to compute the PMF profile. The statistical uncertainty of the PMF profile was estimated using bootstrap analysis.<sup>54</sup>

## 3 Results

The location and distribution of O<sub>2</sub> in the different bilayers are considered to describe the potential of the O<sub>2</sub> molecules to interact with the double bonds of the lipid chains. Under oxidative stress in the cell membrane, radiation, or photo-sensitized oxidation, an O<sub>2</sub> molecules in its ground state gains energy and may jump to an excited singlet state (its lowest excited state).<sup>55</sup> This may lead to lipid peroxidation in the cell membrane which has harmful effects on the biophysical



**Fig. 2** Snapshots at the end of the 1  $\mu$ s simulation of each of the systems. The lipid bilayers without cholesterol (A–C) consist 16 oxygen molecules (O<sub>2</sub>). The light blue region is water. The red spheres are O<sub>2</sub>, and yellow, black, green, and purple lines represent POPC, PLPC, PAPC, and cholesterol molecules, respectively. The green spheres are the phosphorus atoms of the lipid headgroups. For the cholesterol-containing lipid bilayers (D–F), the pink spheres are the oxygen atoms of the hydroxyl groups in the cholesterol headgroups. Note that all O<sub>2</sub> molecules in the water phase can enter and exit the lipid bilayer.



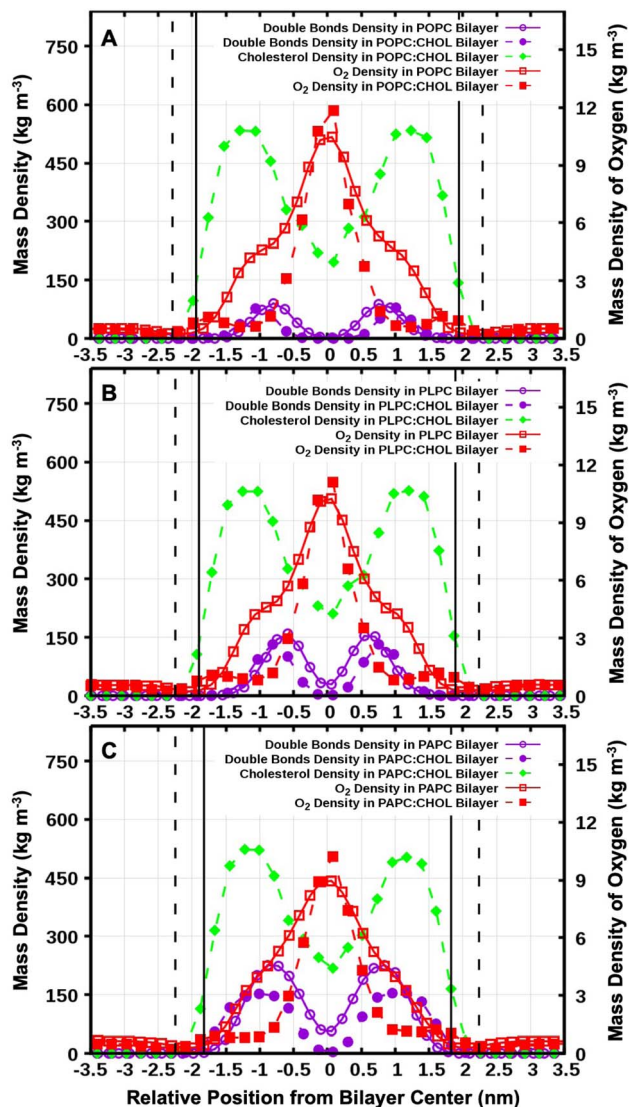


Fig. 3 Mass density profiles of the carbon atoms in the double bonds of the lipids hydrocarbon chains both with and without cholesterol. (A) POPC, (B) PLPC, and (C) PAPC bilayer. The mass density profiles of O<sub>2</sub> are also shown. The vertical solid and dashed lines refer to the average distances of the phosphorus atoms from bilayer center for the systems with and without cholesterol, respectively.

properties and functions of the cell membrane. Previous studies<sup>37</sup> have demonstrated that cholesterol help to protect lipid bilayers from oxidative attack by free radicals. Here, we

examine the effect of cholesterol on the distribution and diffusional pathway of O<sub>2</sub> in lipid bilayers. In addition, free energy calculations were performed using biased umbrella MD simulation techniques to determine the free energy of transfer of O<sub>2</sub> from the water phase to the lipid bilayer. All the simulations showed that the O<sub>2</sub> molecules passively translocate into the bilayers without causing pore formation or membrane rupture. Interestingly, the presence of cholesterol in the unsaturated lipid bilayers strongly decreases the permeability of O<sub>2</sub>.

### 3.1 Location and distribution of O<sub>2</sub> in the lipid bilayer

Fig. 2 shows snapshots of the systems after 1.0  $\mu$ s. O<sub>2</sub> is a small and non-polar molecule, and it is capable to passively and easily diffuse through the lipid bilayer. This is consistent with numerous previous experimental and computational studies.<sup>56–58</sup> We calculated the mass density profiles of the O<sub>2</sub> and cholesterol molecules, and all the carbon double bonds in the unsaturated lipid chains as a function of the distance from the bilayer center (the z-direction). The results are shown in Fig. 3. The average distance of the phosphorus atoms (of the lipid head groups) from the bilayer center was used as the operational definition for the lipid–water interface in order to determine the bilayer thickness. Bilayer thicknesses decreased as the number of double bonds in the lipid chain increased. This is due to the increase in fluidity and disorder in the presence of double bonds. The data is shown in Table 2.

Previous experimental and computational studies have reported that cholesterol increases the membrane thickness while reducing the area per lipid.<sup>59</sup> Our results are well in agreement with those studies for all lipid bilayer types studied here, see Table 2. Although the number of double bonds in the PLPC lipid chain (2 double bonds) is less than in the PAPC chain (4 double bonds), we noted no significant differences in the bilayer thicknesses between these two lipid types at 50% of cholesterol (Table 2).

In addition, we calculated the time evolution of all O<sub>2</sub> molecules moving in the system with respect to the distance from the bilayer center in the z-direction (see Fig. S1–S6†). The results demonstrate that the O<sub>2</sub> molecules passively enter and exit the lipid bilayer, the most favourable location of O<sub>2</sub> is being at the bilayer center as shown by the mass density in Fig. 3. In the presence of cholesterol, a small peak in density was observed at the position of the cholesterol hydrocarbon rings,

Table 2 Bilayer thickness, average number of O<sub>2</sub> contacts with the double bond carbons in the lipid chains (C=C) and OH in cholesterol (OH-CHOL), and the lifetimes of O<sub>2</sub> contacts with (C=C) (ps) were calculated using eqn (2)

Systems	Thickness (nm)	Average area per lipid (nm <sup>2</sup> )	Average number of O <sub>2</sub> in contact with C=C	Average number of O <sub>2</sub> in contact with OH-CHOL	Lifetime of O <sub>2</sub> in contact with C=C (ps)
POPC	3.89 $\pm$ 0.06	0.646 $\pm$ 0.014	1.87 $\pm$ 1.74	—	9.74
POPC:CHOL	4.59 $\pm$ 0.03	0.425 $\pm$ 0.004	0.60 $\pm$ 1.01	0.28 $\pm$ 0.64	7.85
PLPC	3.79 $\pm$ 0.06	0.663 $\pm$ 0.013	4.13 $\pm$ 2.65	—	7.97
PLPC:CHOL	4.49 $\pm$ 0.03	0.433 $\pm$ 0.004	1.90 $\pm$ 1.86	0.30 $\pm$ 0.65	7.92
PAPC	3.65 $\pm$ 0.06	0.711 $\pm$ 0.015	7.35 $\pm$ 3.56	—	7.03
PAPC:CHOL	4.49 $\pm$ 0.04	0.445 $\pm$ 0.005	3.31 $\pm$ 2.52	0.27 $\pm$ 0.63	6.99



and the density of O<sub>2</sub> decreased in the hydrocarbon lipid tail region especially in the double bond region. This result suggests that cholesterol decrease the possibility of O<sub>2</sub> to be in contact with the double-bond regions of lipids.

### 3.2 Number of contacts and contact lifetimes

We calculated the number of contacts between the O<sub>2</sub> molecules and the carbon atoms in the double bonds of the lipids. A contact was defined to exist when the distance was  $\leq 0.35$  nm; overlap distance criterion was taken from a carbon–oxygen hydrogen bonding.<sup>60</sup> Additionally, the average contact lifetime ( $\tau_{\text{contact}}$ ) was calculated from the autocorrelation functions  $C(t)$  of the lifetime distribution ( $P(\tau)$ ) of all O<sub>2</sub> and C=C contacts as

$$C(t) = 1 - \int_0^t P(\tau) d\tau, \quad (1)$$

$$\tau_{\text{contact}} = \int_0^\infty C(\tau) d\tau. \quad (2)$$

As the results in Table 2 show, the number of contacts increases as the number of double bonds increases. For all the lipid bilayer types, the presence of cholesterol resulted in a decrease in the number of contacts; the number of contacts was roughly halved with 50% cholesterol present.

Table 2 also shows that the lifetimes for the O<sub>2</sub> molecules in contact with the double bonds is <10 ps for all bilayers. The lifetimes were calculated using eqn (2). These results indicate that the presence of cholesterol helps to protect the lipid bilayer and to decrease the possibility of lipid oxidation by reducing molecular contacts between the O<sub>2</sub> molecules and the double-bond regions of lipids.

### 3.3 Permeability of O<sub>2</sub> in the lipid bilayer

O<sub>2</sub> is a small and relatively hydrophobic molecule, see Table 1 in Endeward *et al.*,<sup>61</sup> and it can diffuse through the lipid bilayers. We calculated the frequency of O<sub>2</sub> diffusion through a lipid bilayer. The path of oxygen translocation in a lipid bilayer was

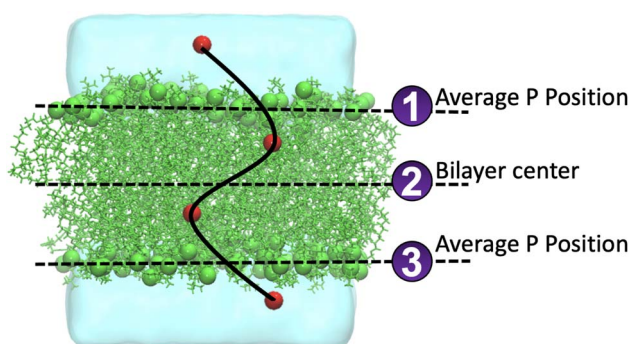


Fig. 4 Definition of the permeation pathway for the oxygen molecules to move through a lipid bilayer. The lines show the bilayer center and the average positions of the phosphorous atoms at the head groups. The latter defines the interface between the bilayer and the water solution.

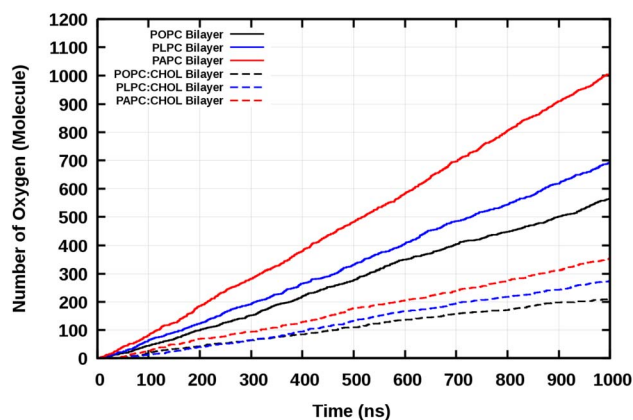


Fig. 5 The number of O<sub>2</sub> molecules passing through the bilayer as a function of time.

defined with three lines that represent the boundaries and the center of the lipid bilayer, see Fig. 4. For the upper and lower leaflets, the average positions of the phosphorous atoms are assigned to the first and the third line in Fig. 4, respectively. The center of a bilayer is shown by the second line. Translocation of O<sub>2</sub> across these three lines was counted as a successful permeation event. The durations of successful permeation events was also recorded. We monitored the systems with 16 O<sub>2</sub> molecules. The number of events of diffusion of O<sub>2</sub> molecules across the lipid bilayers during 1  $\mu$ s is shown in Fig. 5.

The total number of O<sub>2</sub> molecules permeations through the lipid bilayers per ns was determined by the slope of the total number of diffused O<sub>2</sub> molecules through the lipid bilayers and are shown in Table 3.

The results show that the more double bonds present in the lipid tails, the easier it is for the O<sub>2</sub> molecules to permeate across the lipid bilayer. This was observed both in the absence and presence of CHOL, see Table 3.

The number of O<sub>2</sub> permeation events through the bilayers is shown in Table 3. The results show that the presence of cholesterol significantly decreases the number of O<sub>2</sub> permeation across the lipid bilayer. This is consistent with the decrease in the number of contacts between O<sub>2</sub> and the double bond(s) due to the presence of cholesterol (Table 2).

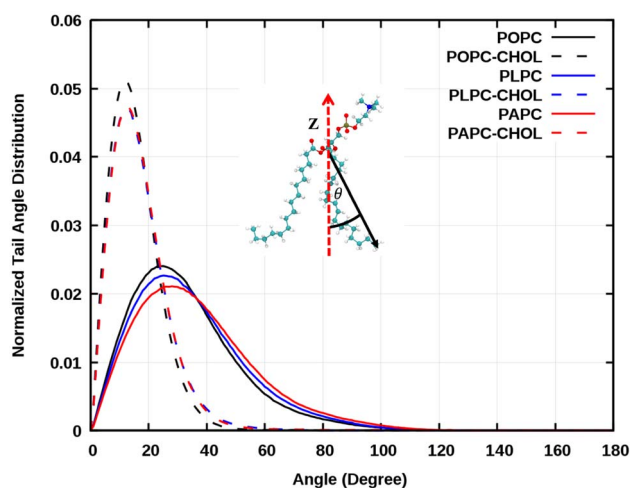
We also calculated the average times that the O<sub>2</sub> molecules reside inside the lipid bilayer, Table 3. The results show that an increase in the number of double bonds results in faster O<sub>2</sub> passage through the lipid bilayer. Interestingly, the presence of cholesterol leads to roughly doubling the residence time in all systems. This increase in residence times caused by cholesterol, which traps oxygen in the middle of the lipid bilayer, prevents the O<sub>2</sub> molecules from reaching the lipid double-bond region to exit the lipid bilayer, see Fig. 3. The reasons for this are elaborated in the next section where the free energy profiles are computed.

Interestingly, the concentration of about 50% is rather high, but there are tissues that have over 70% of cholesterol concentration. Of the lipids used in this study, POPC has the highest main phase transition temperature at around  $-2$  °C



**Table 3** The number of events that 16 O<sub>2</sub> molecules in the system diffuse across the lipid bilayer

Systems	Total number of O <sub>2</sub> passing the bilayer within 1 μs	The number of O <sub>2</sub> permeations per ns (# per ns)	Average time for O <sub>2</sub> passing through the lipid bilayer (ns)
POPC	586	0.57 ± 0.17	12.70 ± 0.50
POPC:CHOL	217	0.21 ± 0.10	24.90 ± 1.52
PLPC	695	0.72 ± 0.19	9.54 ± 0.32
PLPC:CHOL	284	0.30 ± 0.10	19.27 ± 1.03
PAPC	1072	1.07 ± 0.11	6.73 ± 0.19
PAPC:CHOL	365	0.35 ± 0.12	16.29 ± 0.82

**Fig. 6** Average tail angle (sn1 and sn2) distributions of POPC, PLPC, and PAPC lipids were calculated between the vector from the second to the last carbon of each tail and the bilayer normal (z-direction).

(about 271 K). Mixtures of POPC and cholesterol have been studied quite extensively. For example, Dotson *et al.*<sup>62</sup> performed MD simulations of POPC with different concentrations of cholesterol up to 100% cholesterol at 310 K (about 36.85 °C), the same temperature as in our simulations. They did not report any signs of a phase transition. Here, the systems were visualized, and we calculated the lipid tail's angle distributions as shown in Fig. 6. The emergence of an ordered phase was observed at 50% of cholesterol, this is also evident in the snapshots in Fig. 2. The effect of ordered phase on oxygen and water permeation to bilayer was studied by Ghysels *et al.*<sup>63</sup> They showed that the bilayer with an ordered phase is less permeable for oxygen and water compared to the one in disordered phase. This implies that the presence of an order phase may play the role on O<sub>2</sub> permeation.

### 3.4 Permeation scales with membrane thickness and area per lipid

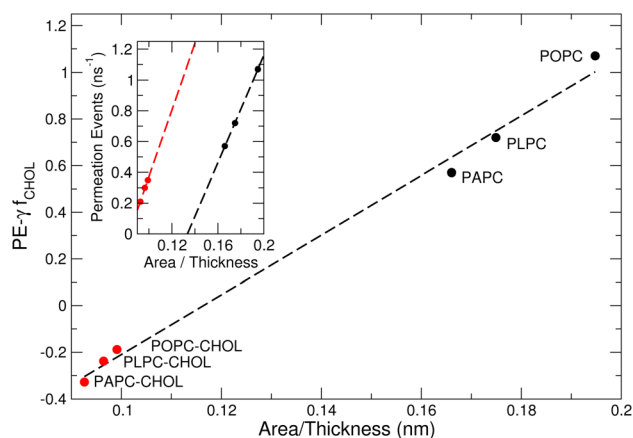
The structure of a lipid membrane is strongly affected by the degree of unsaturation of its comprising lipids and the presence of cholesterol. In particular, these quantities affect the membrane thickness and the area per lipid (Table 2) that can be related to the distance travel by small molecules during a permeation event and the space available to them within the bilayer core, respectively. Accordingly, membranes with higher

area per lipid can more easily accommodate molecules as they permeate the membrane. Conversely, the probability of a molecule to abort a permeation event increases with the membrane thickness. This suggests that the number of permeation events (PE) per unit of time, increases with the area per lipid and decrease with the membrane thickness:  $PE \propto \frac{\text{area}}{\text{thickness}}$ . The inset of Fig. 7 depicts this scaling separately for membranes without (black) and with (red) cholesterol.

To account for a scaling law that unifies membranes with and without cholesterol, one has to recognize that cholesterol occupies some of the space that is otherwise available to accommodate O<sub>2</sub> in a membrane with equal thickness/area per lipid but depleted of this sterol. Accordingly, a factor needs to be added to the proposed empirical scaling law to account for a reduction in permeation due to cholesterol:

$PE - \gamma \left( \frac{N_{\text{CHOL}}}{N_{\text{lipid}}} \right) \propto \frac{\text{area}}{\text{thickness}}$ . The constant  $\gamma$  enables superposition of best-fit-lines in the inset. For our two sets of membranes,  $\gamma$  is 0.53858.

The existence of a scaling law (Fig. 7) shows that despite the complex effects (*e.g.*, Fig. 3 and 6) induced by unsaturated lipids

**Fig. 7** The number of permeation events (PE) scales with the area per lipid and membrane thickness. The fraction  $f_{\text{CHOL}}$  of cholesterol molecules per lipids in the membrane is zero and one in simulations performed without (black) and with (red) cholesterol, respectively. The area per lipid and the membrane thickness of the six membranes studied here is given in Table 2. The number of permeation events per ns is given in Table 3. Best-fit lines to data points are shown using a dashed line.



and cholesterol on the membrane structure, the number of permeation events is dominated by the area per lipid and the membrane thickness.<sup>64</sup>

### 3.5 Free energy profiles

Fig. 8 shows the PMF profiles for O<sub>2</sub> translocation into the bilayers. In all cases, the free energy barrier is at the lipid–water interface. After the maximum at the interfacial barrier, the free energy decreases toward the bilayer interior. The minimum

value was found to be always at the bilayer center (at zero nm in Fig. 8). This represents the equilibrium position, or the most favorable location, of O<sub>2</sub> within a lipid bilayer. This minimum position directly matches the mass density profiles of O<sub>2</sub> measured from the unbiased MD simulations, Fig. 3, with the bilayer center having the highest density of O<sub>2</sub>.

The values for the free energy of O<sub>2</sub> translocation across the lipid bilayer are reported in Table 4. The minimum free energy at the bilayer center ( $G_{\text{center}}$ ) was  $-7.50 \pm 0.26$ ,  $-8.61 \pm 0.25$ , and  $-6.97 \pm 0.26$  kJ mol<sup>-1</sup> for POPC, PLPC, and PAPC lipid bilayer, respectively. In the presence of cholesterol, a deeper  $G_{\text{center}}$  was observed both in the POPC:CHOL bilayer ( $G_{\text{center}} = -10.00 \pm 0.30$ ) and the PLPC:CHOL bilayer ( $G_{\text{center}} = -9.34 \pm 0.38$ ) which may be related to the decrease in density at the bilayer center (see Fig. S7†). Except for the PAPC:CHOL bilayer,  $G_{\text{center}}$  was found to be increased to  $-5.76 \pm 0.32$  kJ mol<sup>-1</sup>. The decrease of free energy at the bilayer is related to the decrease of density at the bilayer center. It is worth noting that the local minima in the free energy profiles, located at approximately 1.75 nm, align with the positions of the cholesterol hydroxyl groups. The presence of cholesterol within the bilayer creates a dual energy barrier that significantly retards the diffusion of oxygen molecules through the central region of the bilayers. Note that the thermal energy in our simulations is 2.58 kJ mol<sup>-1</sup>, and thus thermal fluctuations play a role in aiding O<sub>2</sub> to overcome the energy barriers in Fig. 8.

In a prior computational study, Dotson *et al.*<sup>62</sup> studied mixtures of POPC and cholesterol (0–100% of cholesterol). Their free energy results for the case that can be compared, that is, POPC membrane with 50% of cholesterol are in excellent agreement. In another study, van der Paal *et al.*<sup>65</sup> investigated systems with cholesterol and 1,2-dioleoyl-*sn*-glycero-3-phosphocholine (DOPC). Similarly to our findings, they found the appearance of new free energy barriers in the presence of cholesterol.

In another study, Zuniga-Hertz and Patel<sup>66</sup> concluded that as the number of unsaturations in the lipids increases, there is an increasing number of free volume pockets in the system. Such pockets were proposed to be able to accommodate O<sub>2</sub> molecules. The current results indicate that the center of the bilayer is able to accommodate O<sub>2</sub> molecules as it is the location with the most favourable free energy. However, additional free energy minima (Fig. 8) within the bilayer may also play an important role in trapping O<sub>2</sub> molecules. The latter are more pronounced in membranes made from PLPC and PAPC lipids and in the presence of cholesterol.

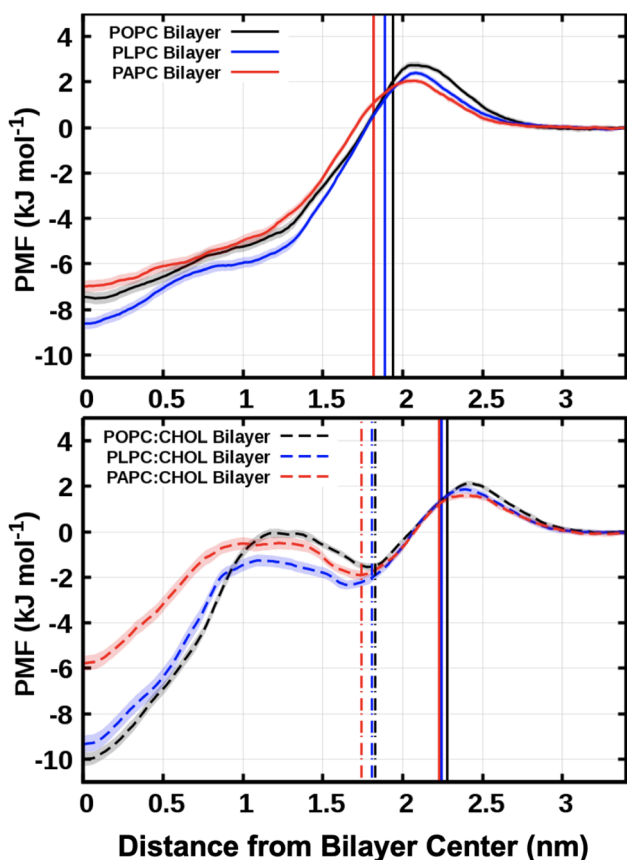


Fig. 8 The free energy profiles for moving an O<sub>2</sub> molecule from the water phase to the bilayer center. The vertical solid lines refer to the average positions of the phosphorus atoms in the systems with and without cholesterol. The vertical dot-dashed lines represent the hydroxyl groups of the cholesterol. The free energy was set to zero in the water phase ( $z = 3.4$  nm). The bootstrap method<sup>54</sup> was used to compute the error bars, shown as transparent shading.

Table 4 The free energy of O<sub>2</sub> transfer from water to the lipid bilayer (kJ mol<sup>-1</sup>)

Systems	Free energy at bilayer center	Free energy barrier at lipid–water interface	Free energy at cholesterol	Total energy
POPC	$-7.50 \pm 0.26$	$2.74 \pm 0.15$	—	$2.74 \pm 0.15$
POPC:CHOL	$-10.00 \pm 0.30$	$2.13 \pm 0.11$	$1.50 \pm 0.20$	$3.63 \pm 0.23$
PLPC	$-8.61 \pm 0.25$	$2.4 \pm 0.12$	—	$2.4 \pm 0.12$
PLPC:CHOL	$-9.34 \pm 0.38$	$1.87 \pm 0.16$	$1.07 \pm 0.31$	$2.94 \pm 0.35$
PAPC	$-6.97 \pm 0.26$	$2.05 \pm 0.11$	—	$2.05 \pm 0.11$
PAPC:CHOL	$-5.76 \pm 0.32$	$1.59 \pm 0.12$	$1.43 \pm 0.32$	$3.02 \pm 0.35$



## 4 Conclusions

In this work, we investigated the permeation of O<sub>2</sub> through the polyunsaturated lipid bilayers, with and without cholesterol, using molecular dynamics (MD) simulations. The results show that O<sub>2</sub> passively diffuses from the aqueous phase through the bilayer. O<sub>2</sub> tends to favor the central region of bilayer, in agreement with the calculated free energy transfer from the aqueous phase to the lipid bilayer. The presence of cholesterol in the unsaturated lipid bilayers notably reduces the permeation of O<sub>2</sub> through the membrane. This effect is accompanied by an increase in bilayer thickness and the emergence of a local energy minimum at the location of cholesterol hydroxyl group. Our results suggest that the presence of cholesterol serves as a protective shield for lipid bilayers by reducing the frequency (number of contacts) and the duration (lifetime) of interactions between the O<sub>2</sub> molecules and the double-bond regions of the lipids. This new insight supports the antioxidant function of cholesterol in cellular membranes.

Although beyond the current study, we would like to point out that there the matter of permeation is even more complex and there are several factors that influence it, in particular the amount of cholesterol, the phase state of the membrane, and the area per lipid and thickness, both of which depend on the phase of the membrane.

Finally, we would also like to mention that no spontaneous flip-flops were observed in simulations. We have, however, studied flip-flops in oxidized PLPC systems in a prior study.<sup>38</sup> Spontaneous flip-flops have been reported in ketosterone-containing systems,<sup>67</sup> albeit in the absence of oxidized lipids. Molecular mechanisms of flip-flop have been discussed in detail, for example, by Gurtovenko and Vattulainen.<sup>68,69</sup>

## Data availability

All data generated or analyzed during this study are included in this published article and ESI.†

## Author contributions

All authors conceived the work, wrote the manuscript and provided critical feedback on the interpretation of data, analysis and discussion. PB carried out the simulations and data analysis.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

This work was financially supported by the National Research Council of Thailand (NRCT) through the Research Grants for Talented Mid-Career Researchers with grant no. N41A640080, the National Science Research and Innovation fund (NSRF) via the Program Management Unit for Human Resources and Institutional Development, Research and Innovation [grant

number B42G670041] and Kasetsart University Research and Development Institute (KURDI) through Fundamental Fund (grant number FF(KU)33.67. MJ thanks the Development and Promotion of Science and Technology Talent Project (DPST). BB thanks the strategic alliance between TU/e, Utrecht University, and University Medical Center Utrecht for financial support. MK thanks the Natural Sciences and Engineering Research Council of Canada (NSERC) and the Canada Research Chairs Program for financial support, and the Digital Research Alliance of Canada for computational resources. CD is supported by the National Institute of General Medical Health under grant no. 1R15GM148982-01.

## References

- 1 G. Van Meer, D. R. Voelker and G. W. Feigenson, *Nat. Rev. Molec. Cell. Biol.*, 2008, **9**, 112–124.
- 2 W. K. Subczynski, J. Widomska and L. Mainali, *Adv. Exp. Med. Biol.*, 2017, **977**, 27–34.
- 3 W. Shinoda, *Biochim. Biophys. Acta*, 2016, **1858**, 2254–2265.
- 4 D. W. Deamer and J. Bramhall, *Chem. Phys. Lipids*, 1986, **40**, 167–188.
- 5 A. A. Gurtovenko and I. Vattulainen, *Biophys. J.*, 2007, **92**, 1878–1890.
- 6 J. Wong-ekkabut and M. Karttunen, *J. Biol. Phys.*, 2016, **42**, 133–146.
- 7 P. Nalakarn, P. Boonnoy, N. Nisoh, M. Karttunen and J. Wong-ekkabut, *Sci. Rep.*, 2019, **9**, 1037.
- 8 N. Nisoh, V. Jarerattanachai, M. Karttunen and J. Wong-ekkabut, *Biomolecules*, 2022, **12**, 639.
- 9 A. Carreau, B. E. Hafny-Rahbi, A. Matejuk, C. Grillon and C. Kieda, *J. Cell. Mol. Med.*, 2011, **15**, 1239–1253.
- 10 D. M. Castilla, Z.-J. Liu and O. C. Velazquez, *Adv. Wound Care*, 2012, **1**, 225–230.
- 11 J.-O. Dunn, M. Mythen and M. Grocott, *BJA Educ.*, 2016, **16**, 341–348.
- 12 N. Netzer, H. Gatterer, M. Faulhaber, M. Burtscher, S. Pramsohler and D. Pesta, *Biomolecules*, 2015, **5**, 1143–1150.
- 13 D. N. Granger and P. R. Kviety, *Redox Biol.*, 2015, **6**, 524–551.
- 14 G. Multhoff, J. Radons and P. Vaupel, *Cancers*, 2014, **6**, 813–828.
- 15 M. Wirtz, I. Stoffels, J. Dissemond, D. Schadendorf and A. Roesch, *J. Eur. Acad. Dermatol. Venereol.*, 2018, **32**, e37–e39.
- 16 J. Heinlin, G. Isbary, W. Stolz, F. Zeman, M. Landthaler, G. Morfill, T. Shimizu, J. Zimmermann and S. Karrer, *J. Eur. Acad. Dermatol. Venereol.*, 2013, **27**, 324–331.
- 17 E. A. Ratovitski, X. Cheng, D. Yan, J. H. Sherman, J. Canady, B. Trink and M. Keidar, *Plasma Processes Polym.*, 2014, **11**, 1128–1137.
- 18 R. Rutkowski, M. Schuster, J. Unger, C. Seebauer, H. Metelmann, T. v. Woedtke, K. Weltmann and G. Daeschlein, *Clin. Plasma Med.*, 2017, **7**, 52–57.
- 19 H.-R. Metelmann, C. Seebauer, V. Miller, A. Fridman, G. Bauer, D. B. Graves, J.-M. Pouvesle, R. Rutkowski, M. Schuster, S. Bekeschus, *et al.*, *Clin. Plasma Med.*, 2018, **9**, 6–13.





- 20 C. Ulrich, F. Kluschke, A. Patzelt, S. Vandersee, V. Czaika, H. Richter, A. Bob, J. v. Hutten, C. Painsi, R. Hüge, *et al.*, *J. Wound Care*, 2015, **24**, 196–203.
- 21 S. Bekeschus, A. Schmidt, M. Napp, A. Kramer, W. Kerner, T. von Woedtke, K. Wende, S. Hasse and K. Masur, *Exp. Dermatol.*, 2017, **26**, 145–147.
- 22 J. S. Sousa, K. Niemi, L. Cox, Q. T. Algwari, T. Gans and D. O'connell, *J. Appl. Phys.*, 2011, **109**, 123302.
- 23 M. Riethmüller, N. Burger and G. Bauer, *Redox Biol.*, 2015, **6**, 157–168.
- 24 G. Bauer, D. Sersenová, D. B. Graves and Z. Machala, *Sci. Rep.*, 2019, **9**, 13931.
- 25 Z. Nasri, S. Memari, S. Wenske, R. Clemen, U. Martens, M. Delcea, S. Bekeschus, K.-D. Weltmann, T. von Woedtke and K. Wende, *Chem.-Eur. J.*, 2021, **27**, 14702–14710.
- 26 B. Li, L. Lin, H. Lin and B. C. Wilson, *J. Biophotonics*, 2016, **9**, 1314–1325.
- 27 P. Agostinis, K. Berg, K. A. Cengel, T. H. Foster, A. W. Girotti, S. O. Gollnick, S. M. Hahn, M. R. Hamblin, A. Juzeniene, D. Kessel, *et al.*, *Ca-Cancer J. Clin.*, 2011, **61**, 250–281.
- 28 G. O. Fruhwirth, A. Loidl and A. Hermetter, *Biochim. Biophys. Acta, Mol. Basis Dis.*, 2007, **1772**, 718–736.
- 29 A. Makky and M. Tanaka, *J. Phys. Chem. B*, 2015, **119**, 5857–5863.
- 30 J. N. van der Veen, J. P. Kennelly, S. Wan, J. E. Vance, D. E. Vance and R. L. Jacobs, *Biochim. Biophys. Acta, Biomembr.*, 2017, **1859**, 1558–1572.
- 31 E. Frankel, *Prog. Lipid Res.*, 1984, **23**, 197–221.
- 32 M. H. Brodnitz, W. W. Nawar and I. S. Fagerson, *Lipids*, 1968, **3**, 59–64.
- 33 M. H. Brodnitz, W. W. Nawar and I. S. Fagerson, *Lipids*, 1968, **3**, 65–71.
- 34 P. Spiteller, W. Kern, J. Reiner and G. Spiteller, *Biochim. Biophys. Acta, Mol. Cell Biol. Lipids*, 2001, **1531**, 188–208.
- 35 D. A. Pratt, J. H. Mills and N. A. Porter, *J. Am. Chem. Soc.*, 2003, **125**, 5801–5810.
- 36 P. Boonnay, V. Jarerattanachai, M. Karttunen and J. Wong-Ekkabut, *J. Phys. Chem. Lett.*, 2015, **6**, 4884–4888.
- 37 T. Parasassi, A. M. Giusti, M. Raimondi, G. Ravagnan, O. Saporita and E. Gratton, *Free Radical Biol. Med.*, 1995, **19**, 511–516.
- 38 P. Boonnay, V. Jarerattanachai, M. Karttunen and J. Wong-Ekkabut, *Biophys. J.*, 2021, **120**, 4525–4535.
- 39 M. J. Abraham, T. Murtola, R. Schulz, S. Páll, J. C. Smith, B. Hess and E. Lindahl, *SoftwareX*, 2015, **1**, 19–25.
- 40 S. Lee, A. Tran, M. Allsopp, J. B. Lim, J. Hénin and J. B. Klauda, *J. Phys. Chem. B*, 2014, **118**, 547–556.
- 41 W. L. Jorgensen, J. Chandrasekhar, J. D. Madura, R. W. Impey and M. L. Klein, *J. Chem. Phys.*, 1983, **79**, 926–935.
- 42 A. D. J. MacKerell, D. Bashford, M. Bellott, R. L. J. Dunbrack, J. D. Evanseck, M. J. Field, S. Fischer, J. Gao, H. Guo, S. Ha, D. Joseph-McCarthy, L. Kuchnir, K. Kuczera, F. T. K. Lau, C. Mattos, S. Michnick, T. Ngo, D. T. Nguyen, B. Prodhom, W. E. Reiher, B. Roux, M. Schlenkrich, J. C. Smith, R. Stote, J. Straub, M. Watanabe, J. Wiórkiewicz-Kuczera, D. Yin and M. Karplus, *J. Phys. Chem. B*, 1998, **102**, 3586–3616.
- 43 S. Jo, X. Cheng, J. Lee, S. Kim, S.-J. Park, D. S. Patel, A. H. Beaven, K. I. Lee, H. Rui, S. Park, H. S. Lee, B. Roux, A. D. MacKerell, J. B. Klauda, Y. Qi and W. Im, *J. Comput. Chem.*, 2016, **38**, 1114–1124.
- 44 M. Parrinello and A. Rahman, *J. Appl. Phys.*, 1981, **52**, 7182–7190.
- 45 S. Nosé, *J. Chem. Phys.*, 1984, **81**, 511–519.
- 46 S. Nosé, *Mol. Phys.*, 1984, **52**, 255–268.
- 47 W. G. Hoover, *Phys. Rev. A: At., Mol., Opt. Phys.*, 1985, **31**, 1695–1697.
- 48 T. Darden, D. York and L. Pedersen, *J. Chem. Phys.*, 1993, **98**, 10089–10092.
- 49 U. Essmann, L. Perera, M. L. Berkowitz, T. Darden, H. Lee and L. G. Pedersen, *J. Chem. Phys.*, 1995, **103**, 8577–8593.
- 50 B. Hess, *J. Chem. Theory Comput.*, 2008, **4**, 116–122.
- 51 W. Humphrey, A. Dalke and K. Schulten, *J. Mol. Graphics*, 1996, **14**, 33–38.
- 52 G. M. Torrie and J. P. Valleau, *J. Comput. Phys.*, 1977, **23**, 187–199.
- 53 S. Kumar, J. M. Rosenberg, D. Bouzida, R. H. Swendsen and P. A. Kollman, *J. Comput. Chem.*, 1992, **13**, 1011–1021.
- 54 J. S. Hub, B. L. de Groot and D. van der Spoel, *J. Chem. Theory Comput.*, 2010, **6**, 3713–3720.
- 55 P. S. Maharjan and H. K. Bhattarai, *J. Oncol.*, 2022, **2022**, 7211485.
- 56 W. K. Subczynski, J. S. Hyde and A. Kusumi, *Proc. Natl. Acad. Sci. U.S.A.*, 1989, **86**, 4474–4478.
- 57 S. Al-Samir, F. Ite, J. Hegermann, G. Gros, G. Tsiavalariis and V. Endeward, *Cell. Mol. Life Sci.*, 2021, **78**, 7649–7662.
- 58 W. K. Subczynski, J. Widomska, M. Raguz and M. Pasenkiewicz-Gierula, *Oxygen*, 2022, **2**, 295–316.
- 59 T. Róg, M. Pasenkiewicz-Gierula, I. Vattulainen and M. Karttunen, *Biochim. Biophys. Acta, Biomembr.*, 2009, **1788**, 97–121.
- 60 S. Horowitz and R. C. Trievel, *J. Biol. Chem.*, 2012, **287**, 41576–41582.
- 61 V. Endeward, S. Al-Samir, F. Ite and G. Gros, *Front. Physiol.*, 2014, **4**, 382.
- 62 R. J. Dotson, C. R. Smith, K. Bueche, G. Angles and S. C. Pias, *Biophys. J.*, 2017, **112**, 2336–2347.
- 63 A. Ghysels, A. Krämer, R. M. Venable, W. E. Teague, E. Lyman, K. Gawrisch and R. W. Pastor, *Nat. Commun.*, 2019, **10**, 5616.
- 64 B. Andziak and R. Buffenstein, *Aging Cell*, 2006, **5**, 525–532.
- 65 J. Van der Paal, C. Verheyen, E. C. Neyts and A. Bogaerts, *Sci. Rep.*, 2017, **7**, 39526.
- 66 J. P. Zuniga-Hertz and H. H. Patel, *Front. Physiol.*, 2019, **10**, 1340.
- 67 T. Róg, L. M. Stimson, M. Pasenkiewicz-Gierula, I. Vattulainen and M. Karttunen, *J. Phys. Chem. B*, 2008, **112**, 1946–1952.
- 68 A. A. Gurtovenko and I. Vattulainen, *J. Phys. Chem. B*, 2007, **111**, 13554–13559.
- 69 A. A. Gurtovenko and I. Vattulainen, *Biomembrane Frontiers: Nanostructures, Models, and the Design of Life*, Humana Press, Totowa, NJ, 2009, pp. 121–139.

