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# From molecular salt to layered network: cation-driven tuning of band gap, structure, and charge transport in $A_3Bi_2I_9$ ( $A = Cs, Rb$ ) perovskites†

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The increasing demand for eco-friendly and stable optoelectronic materials has led to interest in all-inorganic lead-free halide perovskites. This study reports the synthesis of  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ) perovskites *via* a solvothermal technique. The materials crystallize in hexagonal and monoclinic structures, with micrometer-sized particles. Optical investigations reveal direct band-gaps of 2.03 eV for  $Cs_3Bi_2I_9$  and 1.90 eV for  $Rb_3Bi_2I_9$ . Raman spectroscopy highlights distinct vibrational modes, influenced by their structural differences. Space charge limited current (SCLC) measurements indicate varying threshold voltages and trap densities. Impedance spectroscopy and Jonscher's power law analysis reveal different polaron tunneling mechanisms in each compound. Ultrafast transient absorption spectroscopy shows the formation of self-trapped states upon photoexcitation, linked to lattice distortion and the formation of small polarons, which affect electrical conductivity.

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## 1. Introduction

The burgeoning demand for environmentally sustainable and stable optoelectronic materials has prompted intensive exploration into lead-free alternatives, particularly in the realm of halide perovskites.<sup>1–10</sup> Among these, all-inorganic lead-free halide perovskites have emerged as promising candidates owing to their potential for high performance and reduced environmental impact.<sup>11–13</sup> These materials have shown versatility beyond traditional applications, including their potential utility in advanced X-ray detection technologies. For instance, white emission metal halides derived from all-inorganic perovskites offer advantages such as flexibility and transparency, making them promising candidates for next-generation X-ray scintillators.<sup>14</sup> Additionally, one-dimensional  $CsCu_2I_3$  single-crystal X-ray detectors have shown remarkable

sensitivity and efficiency, highlighting the diverse applications of inorganic perovskites in radiation detection.<sup>15</sup>

In this study, we delve into the influence of monovalent cations on the conduction mechanism in lead-free  $A_3Bi_2I_9$  perovskites ( $A = Cs, Rb$ ).<sup>16,17</sup> The choice of Cs and Rb as the monovalent cations stems from their distinct properties, which can profoundly affect the structural and electrical characteristics of the perovskite materials.<sup>18,19</sup> Understanding these effects is pivotal for tailoring the properties of lead-free perovskites to suit specific applications. The synthesis of  $A_3Bi_2I_9$  perovskites was accomplished *via* the solvothermal technique, employing a meticulously controlled slow cooling process.<sup>20</sup> Characterization techniques, including powder X-ray diffraction (PXRD) and scanning electron microscopy (SEM), provided insights into the structural and morphological features of the synthesized materials. Optical investigations, such as UV-Vis spectroscopy, unveiled direct band-gap values, while Raman spectroscopy shed light on vibrational modes characteristic of the crystal structures. Exploration of the charge transport mechanisms in  $A_3Bi_2I_9$  perovskites encompassed SCLC measurements and impedance spectroscopy. These techniques facilitated the elucidation of carrier mobility, trap densities, and contributions from grains and grain boundaries to the total impedance. Jonscher's power law analysis was instrumental in identifying the presence of overlapping large polaron tunneling mechanisms in  $Cs_3Bi_2I_9$  and non-overlapping small polaron tunneling mechanisms in  $Rb_3Bi_2I_9$ , thereby elucidating the charge transportation mechanisms in these lead-free halide perovskites. Furthermore, a comparative analysis of the optical and electrical properties of  $Cs_3Bi_2I_9$  and  $Rb_3Bi_2I_9$  highlighted the

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structural disparities between the two compounds.  $\text{Cs}_3\text{Bi}_2\text{I}_9$ , with its zero-dimensional (0D) molecular salt crystal structure, exhibited distinct optical properties compared to  $\text{Rb}_3\text{Bi}_2\text{I}_9$ , which featured 2D layered derivatives of the perovskite structure. These structural variances significantly influenced the bandgap and electronic confinement, consequently impacting the electrical conductivity of the materials.

In conclusion, this study contributes to the understanding of the intricate interplay between monovalent cations, crystal structure, and functional properties in lead-free  $\text{A}_3\text{Bi}_2\text{I}_9$  perovskites. The elucidated insights offer valuable guidance for the design and optimization of lead-free perovskite materials for diverse optoelectronic applications, underscoring the significance of environmentally sustainable materials in advancing modern technology.

## 2. Experimental section

### 2.1. Materials

Cesium iodide (purity: 99.999%, Sigma-Aldrich), rubidium iodide (purity: 99.9%, Sigma-Aldrich), bismuth(m) iodide (purity:  $\geq 99.998\%$ , Sigma-Aldrich), and hydroiodic acid (purity: 99.99%, contains no stabilizer, distilled, 57 wt% in  $\text{H}_2\text{O}$ , Sigma-Aldrich) were used in their as-received state for the experiments.

### 2.2. Synthesis

$\text{A}_3\text{Bi}_2\text{I}_9$  (where  $\text{A} = \text{Cs, Rb}$ ) single crystals were grown *via* the solvothermal method, implementing a carefully controlled slow cooling procedure as depicted in Fig. 1. Instead of relying on natural cooling, the temperature of the precursor solution was gradually decreased from 180 °C to room temperature at a controlled rate of 1 °C per hour, following established protocols.<sup>4</sup>

### 2.3. Characterization

$\text{A}_3\text{Bi}_2\text{I}_9$  ( $\text{A} = \text{Cs, Rb}$ ) powders were characterized using  $\text{Cu K}\alpha$  radiation ( $\lambda = 1.54 \text{ \AA}$ ) *via* PXRD at room temperature on a Rigaku MiniFlex 600 benchtop diffractometer. The analysis covered a  $2\theta$  range of  $5^\circ$ – $40^\circ$  with a step size of  $0.02^\circ$ .

The morphology of  $\text{A}_3\text{Bi}_2\text{I}_9$  ( $\text{A} = \text{Cs, Rb}$ ) powders was investigated using SEM. The analysis was performed with a Philips XL30 microscope operated at an accelerating voltage of 20 kV.

Thermogravimetric analysis (TGA) was employed to investigate the thermal decomposition behavior of powders derived from  $\text{A}_3\text{Bi}_2\text{I}_9$  ( $\text{A} = \text{Cs, Rb}$ ) single crystals. The analysis was conducted using a PerkinElmer Pyris 6 instrument. The samples were heated at various heating rates, ranging from 0 to 800 °C, within a closed, perforated aluminum pan.

The diffuse reflectance of  $\text{A}_3\text{Bi}_2\text{I}_9$  ( $\text{A} = \text{Cs, Rb}$ ) single crystals was measured in the ultraviolet-visible (UV-Vis) region using a PerkinElmer Lambda 950 spectrophotometer. The analysis was conducted at room temperature. The thicknesses of the single crystals of  $\text{Cs}_3\text{Bi}_2\text{I}_9$  and  $\text{Rb}_3\text{Bi}_2\text{I}_9$  are 0.93 mm and 0.87 mm, respectively.

Raman spectra of  $\text{A}_3\text{Bi}_2\text{I}_9$  ( $\text{A} = \text{Cs, Rb}$ ) single crystals were collected at room temperature using a Horiba-Jobin-Yvon T64000 multi-channel spectrometer. The excitation source for the analysis was a laser operating at a wavelength of 785 nm.

Current–voltage ( $I$ – $V$ ) characteristics of  $\text{A}_3\text{Bi}_2\text{I}_9$  ( $\text{A} = \text{Cs, Rb}$ ) were measured using a Keithley 2400 Source Meter.

The AC conductivity of  $\text{A}_3\text{Bi}_2\text{I}_9$  ( $\text{A} = \text{Cs, Rb}$ ) pellets was investigated using a TH2828A impedance analyzer. Measurements were conducted at various frequencies ranging from 10 to  $10^6$  Hz with an applied AC voltage of 0.5 V. To explore the temperature dependence of conductivity, the samples were subjected to a heating ramp from 300 to 400 K at a rate of 4 K  $\text{min}^{-1}$  controlled by a TP94 temperature controller (Linkam, Surrey, UK).

This work employed the Helios pump-probe ultrafast spectroscopy system for conducting time-resolved measurements known as transient absorption (TA) spectroscopy. The Helios system comprises several key components:

- Laser source: A Ti:sapphire-based laser oscillator (Coherent Micra) acts as the initial light source.
- Amplification: The laser pulse energy is then amplified using a Coherent Legend amplifier.
- Pulse shaping: An optical parametric amplifier (OPA) equipped with a barium borate (BBO) crystal allows for shaping the laser pulse characteristics for optimal TA measurements.

The amplifier within the Helios system generates a near-Gaussian shaped laser pulse with a duration of around 35



Fig. 1 Unveiling the solvothermal route to  $\text{A}_3\text{Bi}_2\text{I}_9$  ( $\text{A} = \text{Cs, Rb}$ ) single crystals: a schematic depiction.



femtoseconds (fs). This ultrashort pulse is centered at a wavelength of 700 nanometers (nm) and delivers an average power of 3.5 Watts (W) at a repetition rate of 1 kHz. Following amplification, the laser beam undergoes a precise splitting process. Approximately 70% of the beam's power is directed towards the optical parametric amplifier (OPA) for further manipulation. The remaining 30% of the beam is directed towards the spectrometer for potential reference or analysis. A critical component of the setup is a delay stage, capable of introducing a precise time delay within a range of 0 to 8 nanoseconds (ns) between the two split portions of the amplified laser beam. This time delay allows the 30% portion to act as a "probe beam" for subsequent spectroscopic analysis. To generate a suitable probe beam, the split beam is directed through a sapphire plate. This process creates a broad spectrum of light, known as a white-light continuum (WLC), encompassing wavelengths ranging from 375 to 750 nanometers (nm). This WLC serves as the ideal probe for the spectroscopic analysis. The highly stable output from the optical parametric amplifier (OPA) provided the pump beam for the experiment. This beam operated at a wavelength near 350 nanometers (nm) and delivered a fluence of 0.5 milliwatts (mW). Careful spatial alignment was achieved by focusing both the pump and probe beams onto the sample.

For subsequent data analysis and fitting, the researchers utilized the Surface Xplorer software. To ensure accurate spectral resolution, the system was calibrated using a well-characterized dye, ZnTPP, in conjunction with an in-house, pre-calibrated ocean optics spectrometer.

## 3. Results and discussion

### 3.1. Structural analysis

PXRD was employed to acquire the crystal structure of  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ) at room temperature (Fig. 2(a) and (b)). The Celref3 software facilitated the indexing of all reflection peaks using the powder data through the least square method.<sup>21,22</sup> PXRD analysis revealed that  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ) exists in two distinct crystal structures.  $Cs_3Bi_2I_9$  crystallizes in the hexagonal  $P6_3/mmc$  (194) space group with unit cell parameters:  $a = b = 8.4122(5)$  Å and  $c = 21.1798(3)$  Å. On the other hand  $Rb_3Bi_2I_9$  adopts a monoclinic  $P2_1/n$  space group with unit cell parameters:  $a = 14.6435(7)$  Å,  $b = 8.1801(4)$  Å,  $c = 20.8848(6)$  Å, and  $\beta = 90.418^\circ$ . These results suggest that  $Cs_3Bi_2I_9$  and  $Rb_3Bi_2I_9$  exhibit distinct crystallographic arrangements. The unit cell parameters obtained from the PXRD analysis are consistent with previously reported values in the literature.<sup>23,24</sup> This validates the structural models identified for  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ). Furthermore, Fig. 2(a) and (b) present insets that depict the crystallographic arrangements of both compounds. These visualizations were generated using the Vesta software.<sup>25,26</sup>

The PXRD analysis revealed that  $Cs_3Bi_2I_9$  crystallizes in the  $P6_3/mmc$  space group, exhibiting a 0D molecular salt structure with hexagonal symmetry.<sup>23</sup> This is evident from the presence of hexagonal channels occupied by Cs cations (the inset of Fig. 2(a)). Notably, unlike perovskite-related structures,  $Cs_3Bi_2I_9$  lacks corner-sharing  $BiI_6$  octahedra. Instead, these octahedra share faces to form discrete  $[Bi_2I_9]^{3-}$  anions, contributing to the

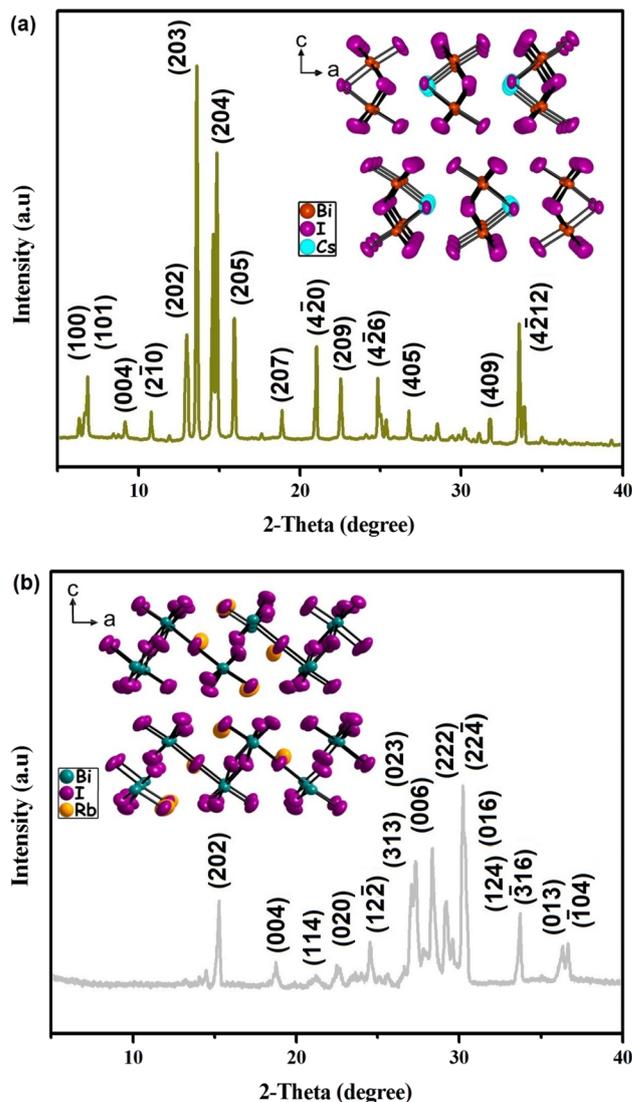


Fig. 2 Crystallographic characterization of (a)  $Cs_3Bi_2I_9$  and (b)  $Rb_3Bi_2I_9$  at room temperature: unveiling lattice structure through powder XRD (inset included).

distinct 0D arrangement. On the other hand, the PXRD analysis suggests that  $Rb_3Bi_2I_9$  adopts a 2D layered perovskite-derivative crystal structure within the monoclinic  $P2_1/n$  space group. Unlike  $Cs_3Bi_2I_9$ , this structure retains the characteristic corner-sharing  $BiI_6$  octahedra observed in classic perovskites. However, these octahedra exhibit a distorted rhombohedral elongation, deviating from the ideal cubic shape. Rb cations occupy the channels between the distorted  $BiI_6$  octahedra, but these channels follow irregular pathways within the 2D layers.<sup>27</sup>

The morphology and particle shapes of  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ) powders were investigated using SEM imaging. Fig. 3(a) and (b) present the corresponding SEM micrographs. ImageJ software was employed to analyze the particle size distribution of  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ) powders. Multiple particle sizes were measured from the SEM images (Fig. 3(a) and (b)) to generate histograms. The average particle size for  $Cs_3Bi_2I_9$  was determined to be 0.620



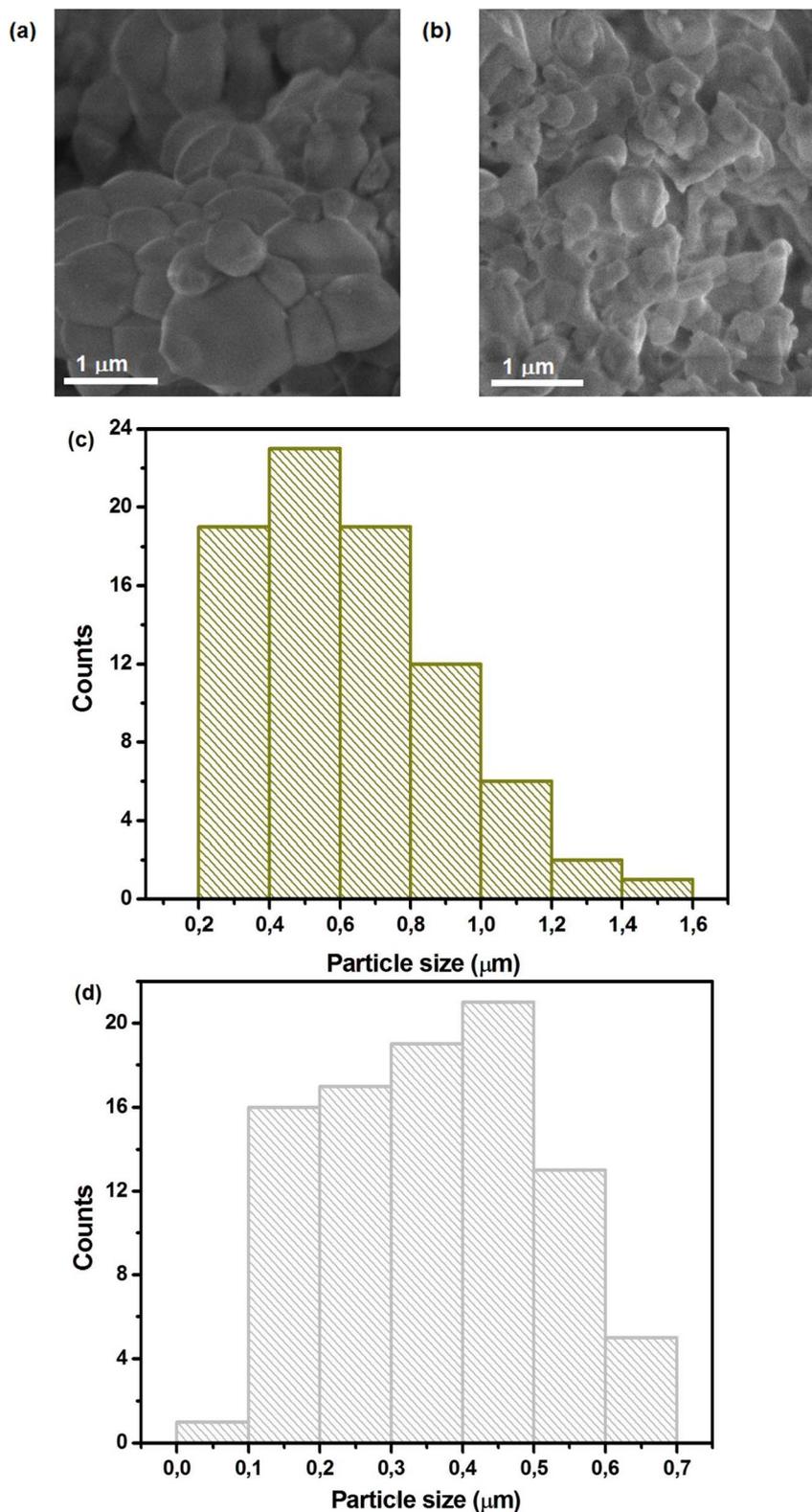


Fig. 3 A comparative look at the morphology of A<sub>3</sub>Bi<sub>2</sub>I<sub>9</sub> (A = Cs, Rb): (a) and (b) SEM images and (c) and (d) particle size analysis.

μm ± 0.269 μm, while Rb<sub>3</sub>Bi<sub>2</sub>I<sub>9</sub> exhibited an average particle size of 0.359 μm ± 0.147 μm.

To assess the thermal stability of A<sub>3</sub>Bi<sub>2</sub>I<sub>9</sub> (A = Cs, Rb), we conducted TGA measurements, as depicted in Fig. 4. Initially,

we noted a slight decrease in weight around 315 and 265 °C for Cs<sub>3</sub>Bi<sub>2</sub>I<sub>9</sub> and Rb<sub>3</sub>Bi<sub>2</sub>I<sub>9</sub>, respectively. This decrease can be attributed to the breakdown of organic capping ligands.<sup>24</sup> Subsequently, a secondary weight reduction commenced



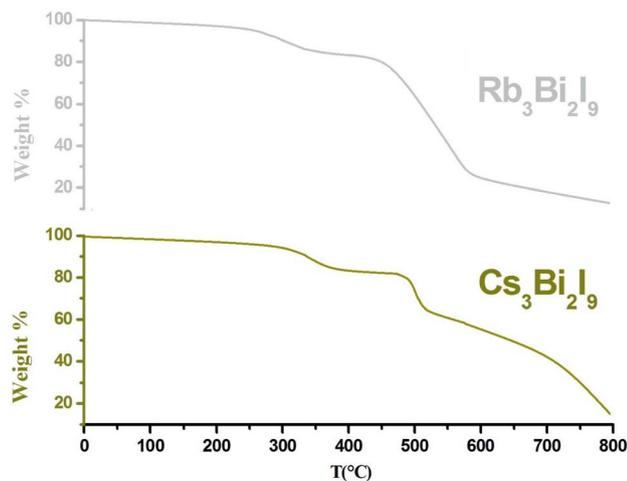


Fig. 4 Investigating thermal stability of  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ): a thermogravimetric analysis (0–800 °C).

around 485 and 445 °C for  $Cs_3Bi_2I_9$  and  $Rb_3Bi_2I_9$ , respectively. This decline signifies the onset of  $A_3Bi_2I_9$  decomposition into  $AI$  ( $A = Cs, Rb$ ) and  $BiI_3$ .<sup>24</sup>

Fig. 5(a) and (b) present the UV-Vis absorption spectra of  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ) measured across a wavelength range of 300 nm to 800 nm. For a more detailed analysis, the Tauc method<sup>28–30</sup> was applied to the absorption spectrum obtained from a single crystal of  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ) (the insets of Fig. 5(a) and (b)). This analysis revealed a direct bandgap ( $E_g$ ) value of 2.03 eV for  $Cs_3Bi_2I_9$  and 1.90 eV for  $Rb_3Bi_2I_9$ .

The Urbach energy ( $E_U$ ), symbolized by  $E_U$ , is a material property associated with the exponential tail observed in absorption spectra at the band edge. This tail arises from electronic transitions between extended states in the valence band and localized states within the conduction band.<sup>29</sup> The Urbach energy is an empirically derived parameter that provides valuable information about the defect density within a material.<sup>31</sup> Urbach's law establishes an empirical relationship between the absorption coefficient ( $\alpha$ ) and photon energy ( $h\nu$ ) as expressed in eqn (S1).<sup>†</sup> This equation allows researchers to estimate the degree of structural disorder within a material by analyzing the slope of the linear portion of the  $\alpha$  vs.  $\ln(h\nu)$  plot derived from the absorption spectrum. Analysis of the absorption spectra using Urbach's law yielded Urbach energy ( $E_U$ ) values of 0.086 eV and 0.044 eV for  $Cs_3Bi_2I_9$  and  $Rb_3Bi_2I_9$ , respectively.

Raman spectroscopy was employed to investigate the vibrational modes of  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ) single crystals at room temperature using a 785 nm laser excitation source (Fig. 6). The Raman spectra of both  $Cs_3Bi_2I_9$  and  $Rb_3Bi_2I_9$  were acquired across a wavenumber range of 2000  $cm^{-1}$ . Interestingly, no significant Raman peaks were detected above 200  $cm^{-1}$  for either material. The Raman spectra of  $Rb_3Bi_2I_9$  and  $Cs_3Bi_2I_9$  reveal distinct vibrational modes.  $Rb_3Bi_2I_9$  exhibits two prominent peaks located above 125  $cm^{-1}$ . Additionally, several overlapping peaks are observed below 75  $cm^{-1}$ . However,  $Cs_3Bi_2I_9$  shows a distinctive feature – many strong vibrational

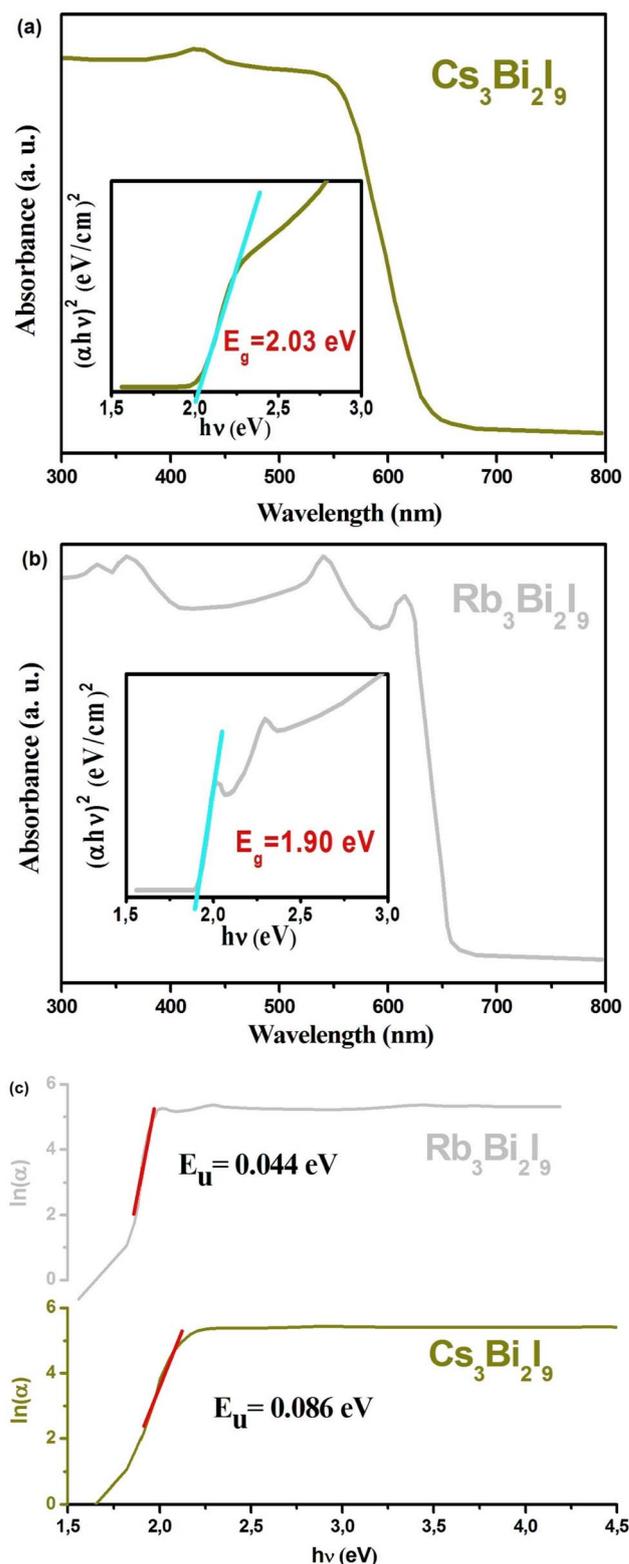


Fig. 5 Optical band gap of (a)  $Cs_3Bi_2I_9$  and (b)  $Rb_3Bi_2I_9$  single crystals: a combined approach utilizing room-temperature absorption measurements, Tauc plot and (c) Urbach analysis.

modes present between 120 and 80  $cm^{-1}$  that are absent in  $Rb_3Bi_2I_9$ . The observed disparity in the Raman spectra can be attributed to the distinct crystal structures of  $Cs_3Bi_2I_9$  and



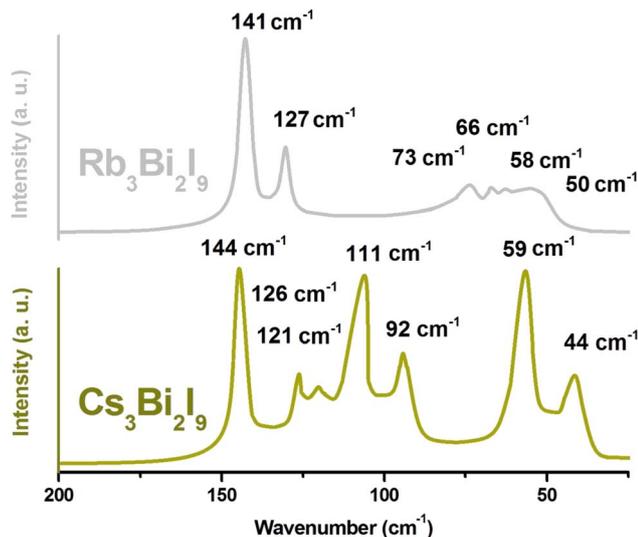


Fig. 6 Comparative analysis of  $A_3\text{Bi}_2\text{I}_9$  perovskites ( $A = \text{Cs}, \text{Rb}$ ) by room-temperature Raman spectroscopy (785 nm excitation).

$\text{Rb}_3\text{Bi}_2\text{I}_9$ ,  $\text{Cs}_3\text{Bi}_2\text{I}_9$  adopts a molecular salt structure, featuring discrete  $[\text{Bi}_2\text{I}_9]^{3-}$  anions. These complex anions allow for additional vibrational modes involving Bi–I bonds, leading to the presence of several strong peaks between 120 and  $80\text{ cm}^{-1}$  that are absent in  $\text{Rb}_3\text{Bi}_2\text{I}_9$ . In contrast,  $\text{Rb}_3\text{Bi}_2\text{I}_9$  possesses a more uniform structure with  $\text{BiI}_6$  octahedra, resulting in the two prominent peaks observed above  $125\text{ cm}^{-1}$  for this material. Table S1† summarizes the prominent peaks of  $A_3\text{Bi}_2\text{I}_9$  ( $A = \text{Cs}, \text{Rb}$ ). A closer examination of the vibrational modes observed in the Raman spectrum of  $\text{Cs}_3\text{Bi}_2\text{I}_9$  (space group  $P6_3/mmc$ ) suggests their primary source lies within the  $[\text{Bi}_2\text{I}_9]^{3-}$  anion.<sup>32</sup> The crystal structure of  $\text{Cs}_3\text{Bi}_2\text{I}_9$  plays a crucial role in understanding the observed Raman peaks. This structure features discrete, covalently bonded  $[\text{Bi}_2\text{I}_9]^{3-}$  anions. Each anion is surrounded by and electrostatically bound to three  $\text{Cs}^+$  cations. This arrangement leads to two distinct types of vibrational modes:

- **Dominant vibrations:** These originate from within the strongly bound  $[\text{Bi}_2\text{I}_9]^{3-}$  units due to the covalent Bi–I bonds. These are expected to be the most prominent peaks in the Raman spectrum.
- **Weaker modes:** These arise from the weaker ionic interactions between the  $[\text{Bi}_2\text{I}_9]^{3-}$  anions and the bridging  $\text{Cs}^+$  cations. These modes are likely to be less intense compared to the dominant vibrations.

The dominant vibrational modes observed in the Raman spectrum of  $\text{Cs}_3\text{Bi}_2\text{I}_9$  are expected to originate from within the  $[\text{Bi}_2\text{I}_9]^{3-}$  anion itself. These modes primarily involve stretching vibrations of the Bi–I bonds. There are two key types of Bi–I bonds within the anion:

- **Terminal Bi–I stretching:** These involve the six iodine atoms located at the periphery of the  $[\text{Bi}_2\text{I}_9]^{3-}$  unit.
- **Bridging Bi–I stretching:** These involve the three iodine atoms that bridge the two central bismuth atoms within the anion, linking them together. The stretching vibrations of both

terminal and bridging Bi–I bonds are expected to contribute significantly to the observed Raman peaks for  $\text{Cs}_3\text{Bi}_2\text{I}_9$ .

Following the established Nakamoto notation for vibrational modes,<sup>33</sup> the Raman spectrum of  $\text{Cs}_3\text{Bi}_2\text{I}_9$  is expected to exhibit specific peaks corresponding to stretching vibrations within the  $[\text{Bi}_2\text{I}_9]^{3-}$  anion.<sup>34</sup> Each set of Bi–I bonds (terminal and bridging) can potentially give rise to three Raman-active stretching modes:

- $A'_1$  symmetric stretch: This mode involves the simultaneous contraction and expansion of both bonds in a symmetrical manner.
- $E''$  antisymmetric stretch (1): This mode involves the bonds stretching out of phase with each other (one contracts while the other expands).
- $E'$  antisymmetric stretch (2): This mode represents another type of out-of-phase stretching motion, potentially different from the first  $E''$  mode.

Therefore, considering both terminal and bridging Bi–I bonds within the  $[\text{Bi}_2\text{I}_9]^{3-}$  anion, a total of six distinct stretchings modes could be observed in the Raman spectrum. The observed Raman spectrum of  $\text{Cs}_3\text{Bi}_2\text{I}_9$  likely reveals a correlation between the bond strength and the energy (and therefore, position) of the corresponding Raman peaks for Bi–I stretching modes. This can be explained by the differences in the Bi–I bonding environment within the  $[\text{Bi}_2\text{I}_9]^{3-}$  anion. Terminal Bi–I bonds involve iodine atoms located at the periphery of the anion and participate in only one Bi–I bond. In contrast, bridging Bi–I bonds involve iodine atoms that connect the two central bismuth atoms and participate in bonds with both Bi atoms. As a result, terminal Bi–I bonds are expected to be stronger due to the absence of bond sharing. Since the energy of a vibrational mode is often related to the strength of the involved bonds, the Raman peaks corresponding to the higher energy stretching modes in the spectrum are likely attributed to the stronger terminal Bi–I bonds. The prominent peak observed at  $144\text{ cm}^{-1}$  is assigned to the most likely candidate – the symmetric terminal Bi–I stretch denoted as  $A'_1$  ( $\nu_1$ ). Its antisymmetric counterparts, the  $E''$  ( $\nu_{15}$ ) and  $E'$  ( $\nu_{10}$ ) modes, are tentatively assigned to the peaks at 126 and  $121\text{ cm}^{-1}$ , respectively. However, it is important to acknowledge that further investigations might be necessary to confirm these assignments definitively. The Raman spectrum of  $\text{Cs}_3\text{Bi}_2\text{I}_9$  reveals a distinct peak at  $111\text{ cm}^{-1}$  which can be tentatively assigned to the symmetric stretching mode of the bridging Bi–I bonds within the  $[\text{Bi}_2\text{I}_9]^{3-}$  anion. This mode corresponds to the  $A'_1$  ( $\nu_2$ ) notation in the Nakamoto convention.<sup>33</sup> Its antisymmetric counterpart, the  $E''$  ( $\nu_7$ ) mode, is likely responsible for the peak observed at  $92\text{ cm}^{-1}$ . These assigned peak positions are consistent with values reported in previous studies.<sup>35,36</sup> Interestingly, the spectrum appears to lack a distinct peak corresponding to the  $E'$  ( $\nu_{11}$ ) mode, which is expected for the antisymmetric stretching of bridging Bi–I bonds. However, the observed width of the  $E''$  ( $\nu_7$ ) peak is broader compared to the  $E''$  and  $E'$  peaks assigned to terminal Bi–I bonds. This broader width might suggest the possible overlap of the  $E''$  ( $\nu_7$ ) peak with a weaker  $E'$  ( $\nu_{11}$ ) peak, which could be obscured within the broader signal. The broad Raman peaks in both spectra can be



attributed to overlapping vibrational modes and potential structural disorder within the crystals. The Raman spectrum of  $\text{Cs}_3\text{Bi}_2\text{I}_9$  exhibits lower energy peaks compared to the stretching modes discussed previously. These lower energy peaks likely correspond to various bending vibrations of the bonds within the  $[\text{Bi}_2\text{I}_9]^{3-}$  anion, as well as weaker vibrational modes arising from the interaction between these anions and the surrounding  $\text{Cs}^+$  cations. However, due to the complexity of these modes, assigning them to specific bonds or interactions definitively requires additional techniques like polarization experiments and theoretical calculations. One interesting observation is the unusually strong peak centered at  $58\text{ cm}^{-1}$  compared to the higher energy stretching modes. This can be attributed to the unique crystal structure of  $\text{Cs}_3\text{Bi}_2\text{I}_9$ , classified as a molecular salt. In this structure, the  $[\text{Bi}_2\text{I}_9]^{3-}$  anions are not rigidly bound in specific orientations by the  $\text{Cs}^+$  cations. This allows the anions a certain degree of rotational and twisting freedom within the lattice. This freedom of movement within the lattice enables the bending modes of the Bi–I bonds to occur as frequently as the stretching modes. This behavior contrasts sharply with compounds that possess a covalently bonded network, where such bending motions are restricted. The observed strong peak at  $58\text{ cm}^{-1}$  is likely a consequence of these bending modes. Additionally, the rotational symmetry of the  $[\text{Bi}_2\text{I}_9]^{3-}$  anion and the strong Bi–I coupling within it contribute to the situation where the rotational bending of iodine atoms around the central Bi–Bi axis becomes energetically similar to the symmetric stretching motions.

The Raman spectrum of  $\text{Rb}_3\text{Bi}_2\text{I}_9$  presents a distinct contrast compared to  $\text{Cs}_3\text{Bi}_2\text{I}_9$ . Unlike the broader range of vibrational modes observed in  $\text{Cs}_3\text{Bi}_2\text{I}_9$ ,  $\text{Rb}_3\text{Bi}_2\text{I}_9$  exhibits a simpler spectrum. This difference can be attributed to the varying crystal structures of these materials. The key feature of the  $\text{Rb}_3\text{Bi}_2\text{I}_9$  spectrum is the presence of two dominant peaks positioned above  $125\text{ cm}^{-1}$ . These peaks likely correspond to the same Bi–I stretching modes observed in  $\text{Cs}_3\text{Bi}_2\text{I}_9$ , but with slightly lower wavenumbers. Based on the Nakamoto notation, the higher energy peak at  $141\text{ cm}^{-1}$  can be tentatively assigned to the symmetric  $A_{1g}(\nu_1)$  Bi–I stretching mode. The lower energy peak at  $127\text{ cm}^{-1}$  is then likely its counterpart, the  $E_g(\nu_2)$  antisymmetric Bi–I stretching mode.<sup>33</sup>

The  $\text{Rb}_3\text{Bi}_2\text{I}_9$  Raman spectrum also exhibits several overlapping peaks at lower wavenumbers compared to the dominant high-energy features. These lower energy peaks are likely a consequence of various bending vibrations within the molecule. Due to the complex nature of these bending modes, assigning them to specific bonds definitively requires further investigation. The relatively low intensity of these bending modes can be attributed to the constraints imposed by the covalent network in  $\text{Rb}_3\text{Bi}_2\text{I}_9$ . This behavior is similar to what is observed in  $\text{Cs}_3\text{Bi}_2\text{I}_9$ , where the covalently bonded  $\text{BiI}_6$  octahedra restricts the bending motions of the bonds.<sup>18</sup>

In examining the electrical characteristics of  $\text{A}_3\text{Bi}_2\text{I}_9$  ( $\text{A} = \text{Cs}, \text{Rb}$ ), we conducted an analysis of SCLC on the  $\text{Au}/\text{A}_3\text{Bi}_2\text{I}_9/\text{Au}$  device, as depicted in Fig. 7. This analysis was carried out based on eqn (S2).<sup>†</sup> Consequently, the determined trap density ( $n_{\text{traps}}$ ) for  $\text{Cs}_3\text{Bi}_2\text{I}_9$  stands at  $7.76 \times 10^9\text{ cm}^{-3}$ , exhibiting a lower value

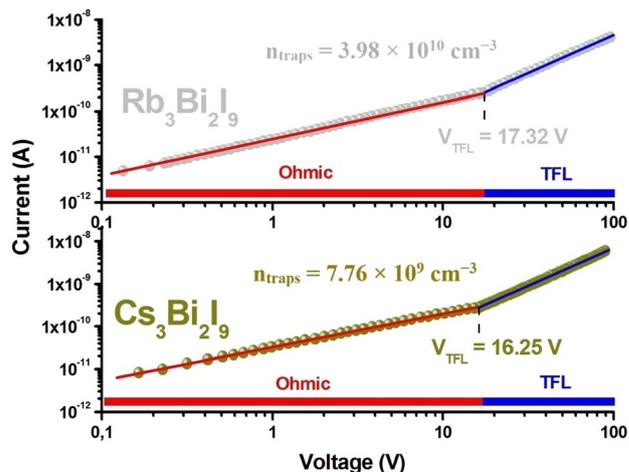


Fig. 7 Electrical characterization of  $\text{A}_3\text{Bi}_2\text{I}_9$  ( $\text{A} = \text{Cs}, \text{Rb}$ ) perovskites at room temperature: unveiling current–voltage behavior.

compared to  $\text{Rb}_3\text{Bi}_2\text{I}_9$ , which registers at  $3.98 \times 10^{10}\text{ cm}^{-3}$ . This suggests  $\text{Cs}_3\text{Bi}_2\text{I}_9$  exhibits comparatively fewer charge trapping defects than  $\text{Rb}_3\text{Bi}_2\text{I}_9$ , indicating potentially superior electrical properties.<sup>37</sup>

Complex impedance spectroscopy (CIS) emerged as a valuable tool for characterizing the electrical behavior of  $\text{A}_3\text{Bi}_2\text{I}_9$  ( $\text{A} = \text{Cs}, \text{Rb}$ ). This technique allows for the simultaneous investigation of two key electrical properties: ion mobility and electrical conductivity. By analyzing the impedance response of the material, CIS provides insights into both the movement of ions within the lattice and the overall ability of the material to conduct electricity.<sup>38–40</sup> CIS stands out as a powerful, non-destructive technique for probing the electrical properties of materials. Unlike other methods, CIS avoids sample damage while offering crucial insights into key parameters like electrical conductivity, relaxation time, and ion mobility within  $\text{A}_3\text{Bi}_2\text{I}_9$  perovskites. Furthermore, CIS possesses a unique capability – the ability to differentiate contributions from various sources within the material. By employing a time-reversed electric field, CIS can separate the electrical response of the bulk material, grain boundaries, and electrodes.<sup>38–40</sup> This ability to deconvolute contributions provides a much more comprehensive understanding of the electrical behavior within  $\text{A}_3\text{Bi}_2\text{I}_9$ . In CIS, the application of an alternating electric field plays a crucial role in investigating charge transport within  $\text{A}_3\text{Bi}_2\text{I}_9$  ( $\text{A} = \text{Cs}, \text{Rb}$ ). This alternating field interacts with the material, activating charge carriers (ions and electrons). This activation leads to a series of events:

- Formation of space charge: The movement of charges creates a localized region of excess or deficit charge within the material.
- Charge displacement: The alternating field induces the movement of these charges, contributing to the overall current flow.
- Dipole orientation: The field can also influence the orientation of polar molecules or defect dipoles present within the material.



These combined effects allow CIS to distinguish the contributions of various components to the overall impedance of the material. By analyzing the frequency dependence of the measured impedance (as described by eqn (S3)†), it becomes possible to differentiate the contributions of electrodes, grain boundaries, and the bulk material (grains) to the transport of charge carriers within  $A_3Bi_2I_9$ .

The electrical properties of  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ) exhibit a temperature dependence, as revealed by the analysis of the real part ( $Z'$ ) of the complex impedance (Fig. 8(a)). Notably, increasing temperature and frequency both lead to a decrease in  $Z'$ . This observed trend suggests a negative temperature coefficient of resistance (NTCR) behavior in these materials.<sup>41</sup> The observed decrease in  $Z'$  with increasing temperature and frequency suggests two potential mechanisms influencing the electrical behavior of  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ):

- **Enhanced AC conductivity:** The decrease in  $Z'$  can be attributed to an increase in the alternating current (AC) conductivity within the material. This could be due to a reduction in the density of trapped charges. As temperature rises, trapped charges may gain enough energy to escape their localized positions and contribute to overall conductivity. Additionally, increased frequency can limit the time available for charges to become trapped, further enhancing conductivity.<sup>39–41</sup>

- **Space charge influence:** At high frequencies, the observed temperature-independent behavior of  $Z'$  points towards the presence of space charge within the material for both  $Cs_3Bi_2I_9$  and  $Rb_3Bi_2I_9$ . Space charge refers to a localized region of excess or deficit charge within the material. At high frequencies, the rapid fluctuations of the alternating field may not effectively penetrate the entire material, resulting in a dominant contribution from the space charge layer to the impedance, leading to a less pronounced temperature dependence of  $Z'$ .<sup>39–41</sup>

The variations in the loss factor ( $Z''$ ) depicted in Fig. 8(b) for  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ) across different frequencies and temperatures offer valuable insights into the relaxation processes and dominant charge transport mechanisms within these materials. The observed changes in  $Z''$  can be attributed to the relaxation of mobile charges after an applied electric field is removed. At higher temperatures, the relaxation process is likely dominated by the movement of defects or vacancies within the crystal structure. These defects or vacancies can act as temporary traps for charge carriers, hindering their movement. However, at higher temperatures, these trapped charges gain enough energy to overcome the trapping potential and contribute to the overall conductivity. This results in a more efficient relaxation process, reflected by a lower peak in  $Z''$  at higher temperatures.<sup>42</sup> Conversely, at lower temperatures, the relaxation process is expected to be governed by the movement of immobile species or electrons. These species or electrons may be tightly bound within the crystal lattice or possess limited mobility due to the lower thermal energy. Consequently, the relaxation process becomes sluggish, leading to a broader and potentially higher peak in  $Z''$  observed at lower temperatures. Analysis of the loss factor ( $Z''$ ) spectra for  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ) (as shown in Fig. 8(b)) reveals the presence of a distinct peak, often referred to as the “relaxation frequency.” This peak signifies a characteristic

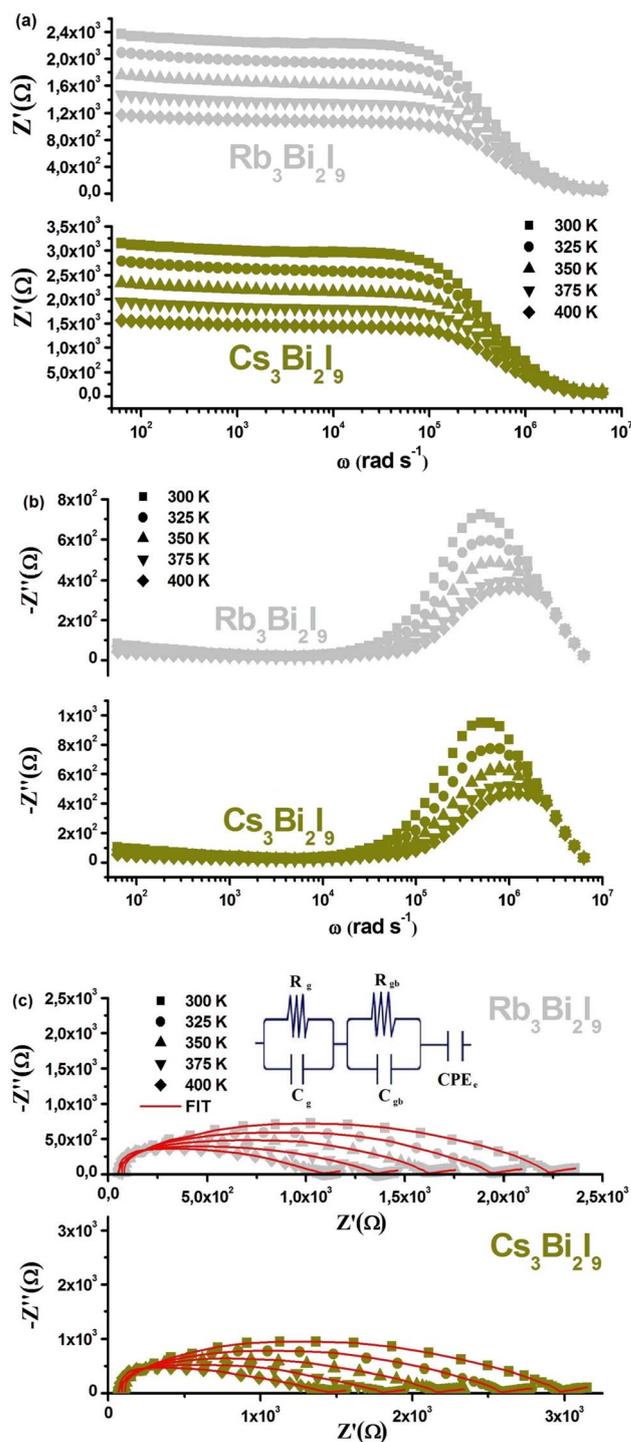


Fig. 8 Deciphering the impedance behavior of  $A_3Bi_2I_9$  perovskites ( $A = Cs, Rb$ ): (a)–(c) a combined approach using frequency and temperature dependence with equivalent circuit modeling.

frequency at which the relaxation process of mobile charges within the material becomes most prominent.<sup>39</sup> Notably, the broadness of the peak observed in these materials suggests a deviation from the ideal Debye relaxation model, which assumes a single, well-defined relaxation time for all charge carriers. The spectra offer insights into the dominant modes of



charge transport across different frequency ranges. Below the relaxation peak frequency, the movement of charge carriers is likely dominated by a long-range transport mechanism. This could involve the migration of charges across relatively large distances within the crystal lattice.<sup>41</sup> In contrast, at frequencies exceeding the relaxation peak frequency, the spectra suggest a shift towards the movement of more localized charge carriers. These localized carriers may be temporarily trapped within the lattice or possess limited mobility due to interactions with defects or other localized states. As temperature increases, the relaxation peak exhibits a distinct shift towards higher frequencies. This observation suggests a thermally activated relaxation process within the material.<sup>39,40</sup> In simpler terms, the relaxation of mobile charges within the crystal lattice becomes faster at elevated temperatures. This could be attributed to factors like increased thermal energy, allowing charge carriers to overcome potential barriers or defects more efficiently. Interestingly, the peak height also exhibits a temperature dependence, decreasing with increasing temperature. This trend points towards a decline in the overall resistance experienced by the charge carriers. This potentially indicates a decrease in both grain boundary resistance and resistance within the bulk grains (grains) of the material. As temperature rises, the mobility of charge carriers might improve, leading to a reduction in the resistance encountered at these interfaces and throughout the material. This aligns with observations reported in previous studies.<sup>38–41</sup>

The Cole–Cole plots presented in Fig. 8(c) for  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ) offer valuable insights into the electrical relaxation processes within the materials across a temperature range of 300 to 400 K. Notably, both materials exhibit depressed semicircles instead of perfect semicircles in the Cole–Cole plots. This deviation from the ideal Debye model indicates the presence of non-Debye relaxation behavior. In simpler terms, the relaxation of mobile charges within the material does not occur with a single, well-defined relaxation time but rather involves a distribution of relaxation times. Furthermore, the Cole–Cole plots display a trend where the radii of the semicircles decrease with increasing temperature. This observation aligns with the negative temperature coefficient of resistance (NTCR) behavior observed in these materials. As temperature rises, the overall resistance to charge transport within the material decreases, which is reflected by the shrinking radii of the semicircles in the Cole–Cole plots. This finding aligns with previous reports on NTCR behavior in similar materials.<sup>43</sup> The electrical properties of  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ), as revealed by the Nyquist plots (Fig. 8(c)), suggest characteristics typical of semiconductors. This implies that the dominant conduction mechanism within these materials is thermally activated. In simpler terms, the conductivity increases with rising temperature as more charge carriers gain enough thermal energy to participate in conduction processes. The ideal Nyquist plot for this type of material would typically showcase three distinct semicircles. Each semicircle represents the impedance contribution from a specific region within the material.<sup>43</sup>

• **High-frequency semicircle:** This semicircle observed at higher frequencies corresponds to the contribution of the bulk

material (grains). It reflects the intrinsic resistance and capacitance associated with the movement of charge carriers within the crystal lattice.

• **Low-frequency semicircle:** The semicircle observed at lower frequencies represents the contribution of the grain boundaries. Grain boundaries can act as regions of higher resistance compared to bulk material due to defects or limited mobility of charge carriers at these interfaces.

However, interpreting the contributions from different regions of the material can be challenging. While the presence of a single, depressed semicircle often reflects the combined influence of grains and grain boundaries due to their overlapping frequency ranges, the Nyquist plot can also reveal a separate contribution from the electrode–material interface at ultra-low frequencies. This contribution from the electrode–material interface arises due to interfacial phenomena such as charge transfer processes or the formation of a double electric layer at the contact points between the electrodes and the material. It typically manifests as a smaller semicircle at very low frequencies in the Nyquist plot. To fully disentangle the contributions from these various regions, researchers often employ the Maxwell–Wagner equivalent circuit model.<sup>40</sup> This model represents the electrical behavior using a parallel combination of resistors ( $R$ ) and capacitors ( $C$ ), where each element corresponds to a specific region or process: even when they appear as a combined response in the Nyquist plot. This approach provides a more comprehensive understanding of the electrical properties within the material. To gain a deeper understanding of the electrical behavior of  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ), an equivalent circuit model has been employed. This model, fitted using ZView software for the impedance spectra across various temperatures, consisted of two parallel RC (resistor–capacitor) elements connected in series. Each element represents the electrical properties of a specific region within the material:

• **Bulk material (grains):** Represented by the first RC element, it reflects the intrinsic resistance ( $R_g$ ) to charge transport and the associated capacitance ( $C_g$ ) within the crystal lattice.

• **Grain boundaries:** Modeled by the second RC element ( $R_{gb}$ ,  $C_{gb}$ ), it accounts for the higher resistance encountered at grain boundaries due to factors like defects or limited charge carrier mobility.

The Nyquist plots (refer to Fig. 8(c)) also reveal a straight line following the semicircle at lower frequencies. This behavior is attributed to the presence of a constant phase element (CPE), which often arises due to inhomogeneities in the electrode surface and the formation of a double-layer capacitance. The inset figures provide a closer look at this low-frequency region. A significant observation from the Nyquist plots is the decrease in the diameter of the semicircle and its intersection with the real axis as temperature increases. This trend suggests an enhancement in the DC conductivity of  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ) at higher temperatures.<sup>39</sup> Furthermore, analysis of the fitted parameters (Table S2†) reveals that the resistance associated with grain boundaries ( $R_{gb}$ ) consistently exceeds the resistance within the grains ( $R_g$ ) across all temperatures. This difference in resistance values contributes to the accumulation of space



charges at the grain boundaries due to a phenomenon known as the Maxwell–Wagner effect.<sup>40</sup> As mentioned previously,<sup>39</sup> grain boundaries are generally more resistive than the bulk material (grains) in perovskite materials.<sup>38</sup>

Fig. 9(a) depicts the fascinating interplay between temperature and frequency on the AC conductivity of  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ) crystals. The AC conductivity, calculated using eqn (S4),<sup>†</sup> dominates the overall conductivity at high frequencies. In contrast, the low-frequency range reflects the temperature-dependent and frequency-independent DC conductivity. This explains the plateau observed in the curves at lower frequencies.<sup>44</sup> As frequency increases, the AC conductivity contributions become significant, leading to a pronounced rise and frequency dispersion (dependence on frequency) in the conductivity curves. This transition from DC to AC-dominated behavior occurs at a specific frequency, often referred to as the “hopping frequency.” Notably, the hopping frequency (represented by a symbol, which can be removed) shifts towards higher frequencies with increasing temperature. Funke's jump relaxation model (JRM) offers a compelling explanation for the observed trends.<sup>45</sup> At low frequencies, the model suggests that ions within the material can successfully hop to neighboring sites. These successful hops contribute to long-range, translational motion of the ions, ultimately leading to a sustained DC conductivity observed at low frequencies. However, at higher frequencies, the hopping process becomes more dynamic. According to the JRM, these hops can either be successful (resulting in ion movement and conductivity) or unsuccessful (where the ion returns to its original position and doesn't contribute to conductivity). This interplay between successful and unsuccessful hops at high frequencies explains the frequency dispersion observed in the AC conductivity.<sup>45</sup>

The JRM model provides a framework for understanding the observed frequency dependence of AC conductivity. According to this model, the hopping behavior of ions within the material plays a crucial role. At low frequencies, the hopping process is dominated by successful events. Here, an “excited” ion successfully jumps to a neighboring site, contributing to long-range translational motion within the material. These successful hops ultimately lead to the sustained DC conductivity observed at lower frequencies. However, as the frequency increases, the hopping dynamics become more complex. JRM suggests that hops can now be both successful and unsuccessful. In an unsuccessful hop, the excited ion may oscillate back and forth before returning to its initial position. This unsuccessful hopping doesn't contribute to the overall movement of ions and therefore doesn't contribute to conductivity. With increasing frequency, the proportion of unsuccessful hops increases compared to successful ones. This imbalance leads to the observed dispersion in the conductivity curves at high frequencies (dispersive conductivity). The measured electrical conductivity can be described by Jonscher's power law (eqn (S5)<sup>†</sup>), which mathematically captures this frequency dependence.<sup>45</sup>

The experimental data for AC conductivity variation ( $\sigma_{AC}(\omega)$ ) presented in Fig. 9(a) was employed to determine the DC conductivity ( $\sigma_{DC}$ ) values for  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ) across all investigated temperatures. This analysis utilized the Universal Dielectric Response (UDR) approach. Following the UDR analysis, Fig. 9(b) depicts a plot of the natural logarithm of DC conductivity ( $\ln(\sigma_{DC})$ ) versus the inverse of temperature ( $1000/T$ ). This plot allows for the calculation of the average activation energies for both materials using the Arrhenius equation (eqn (S6)<sup>†</sup>). The Arrhenius equation relates the temperature dependence of conductivity to the activation energy, a parameter that

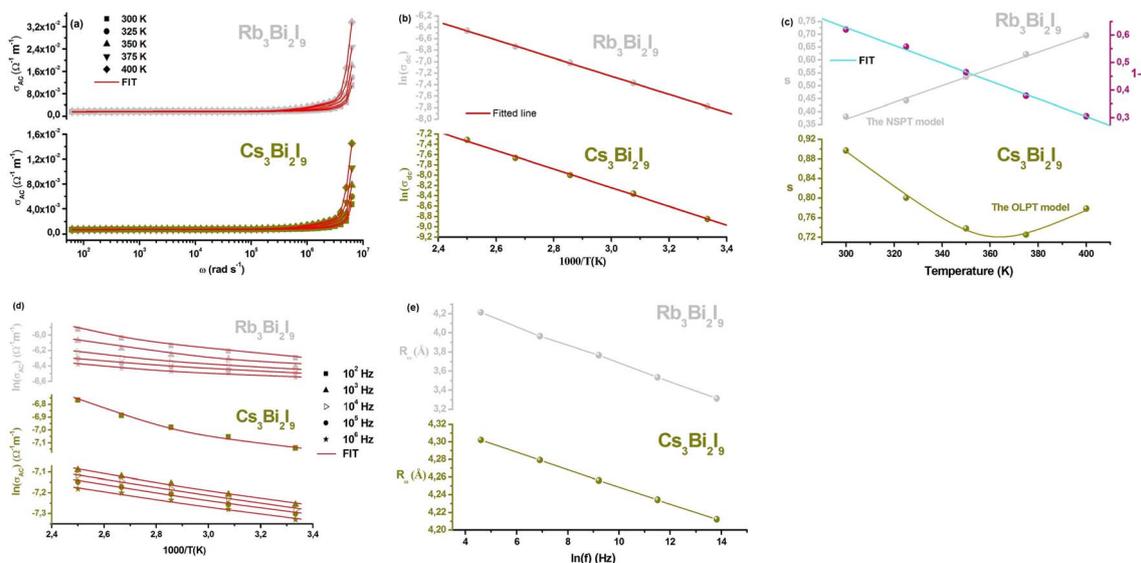


Fig. 9 (a) Unveiling the temperature–frequency landscape of AC conductivity in  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ). (b) Extracting activation energy for electrical conduction using an Arrhenius-like approach. (c) Unveiling the temperature dependence of charge transport mechanisms in  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ). (d) Investigating charge transport dynamics in  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ): unveiling the role of relaxation processes through  $R_\omega$  analysis. (e) Frequency-dependent of  $R_\omega$ .



reflects the energy barrier for charge carriers to participate in conduction processes. The activation energy values for  $\text{Cs}_3\text{Bi}_2\text{I}_9$  and  $\text{Rb}_3\text{Bi}_2\text{I}_9$  stand at 0.155 eV and 0.135 eV, respectively.

Understanding the charge transport mechanisms within  $\text{A}_3\text{Bi}_2\text{I}_9$  ( $\text{A} = \text{Cs}, \text{Rb}$ ) crystals is crucial. By analyzing the temperature dependence of the exponent 's' (refer to Fig. 9(c)), researchers can gain valuable insights into the dominant processes. Several theoretical models have been proposed in the literature to explain the behavior of 's'.<sup>38–41</sup> These models often involve a combination of quantum mechanical tunneling and/or hopping of charge carriers (electrons/polarons or even atoms) across potential barriers. Some prominent examples include overlapping large polaron tunneling (OLPT), non-overlapping small-polaron tunneling (NSPT), correlated barrier hopping (CBH), and pure quantum mechanical tunneling (QMT). Insights from 's' for  $\text{Cs}_3\text{Bi}_2\text{I}_9$  and  $\text{Rb}_3\text{Bi}_2\text{I}_9$ :

- $\text{Cs}_3\text{Bi}_2\text{I}_9$ : The observed decrease in 's' with increasing temperature, followed by a slight increase, suggests that the OLPT model might be most suitable for  $\text{Cs}_3\text{Bi}_2\text{I}_9$ . This behavior aligns with the characteristics of overlapping large polarons, where the tunneling probability increases at lower temperatures due to a larger overlap between adjacent polarons.

- $\text{Rb}_3\text{Bi}_2\text{I}_9$ : In contrast,  $\text{Rb}_3\text{Bi}_2\text{I}_9$  exhibits an increase in 's' with temperature. This trend points towards the NSPT model as a better fit. NSPT describes the hopping of smaller polarons with minimal overlap between them. As temperature rises, the thermal energy likely aids these hops, leading to a higher value of 's'.

For a deeper understanding of the fitting process, detailed descriptions of the OLPT and NSPT models are provided in the ESI.† Fig. 9(d) presents a crucial piece of evidence supporting the chosen models. It depicts the variation of the natural logarithm of AC conductivity ( $\ln(\sigma_{\text{AC}})$ ) with temperature for both  $\text{A}_3\text{Bi}_2\text{I}_9$  ( $\text{A} = \text{Cs}, \text{Rb}$ ) materials. The key takeaway from this figure is the excellent agreement between the fitting curves based on the OLPT and NSPT models and the actual experimental data points. This agreement strengthens the validity of our model selection. This fitting process served another important purpose: extracting the key parameters associated with both the OLPT and NSPT models. These extracted parameters hold valuable information about the charge transport mechanisms within the materials. Tables S3–S5† present the complete set of parameters obtained for  $\text{Cs}_3\text{Bi}_2\text{I}_9$  and  $\text{Rb}_3\text{Bi}_2\text{I}_9$ .

The analysis of the extracted parameters for  $\text{Cs}_3\text{Bi}_2\text{I}_9$  provides compelling evidence supporting the selection of the overlapping large polaron tunneling (OLPT) model. Notably, the value of the parameter  $\alpha$  falls within the range reported by Murawski *et al.*<sup>46</sup> for similar systems involving localized states, which aligns with the characteristics of large polarons in the OLPT model. Fig. 9(e) further strengthens this connection. It depicts the frequency dependence of the parameter  $R_{\omega}$  for  $\text{Cs}_3\text{Bi}_2\text{I}_9$ . Interestingly, the calculated tunneling distances ( $R_{\omega}$ ) range from 4.212 to 4.302 Å. This range is remarkably close to the interatomic distances between Cs and I ions (Cs–I) within the  $\text{Cs}_3\text{Bi}_2\text{I}_9$  crystal lattice, which is reported to be between 4.204 and 4.318 Å.<sup>47</sup> This close correspondence between the tunneling distances and interatomic spacings suggests that the

mobility of  $\text{Cs}^+$  cations might play a crucial role in facilitating AC conductivity. The large polarons associated with the OLPT model can leverage the movement of these cations to enhance conductivity through the material.

Similar to  $\text{Cs}_3\text{Bi}_2\text{I}_9$ , the extracted parameters for  $\text{Rb}_3\text{Bi}_2\text{I}_9$  provide valuable insights into the chosen model (NSPT). The observed values for the relevant parameters fall within the expected range for localized states, consistent with the characteristics of small polarons in the NSPT model.<sup>38</sup> Table S5† depicts the frequency dependence of two key parameters,  $N(E_{\text{F}})$  and  $\alpha$ , for  $\text{Rb}_3\text{Bi}_2\text{I}_9$ . Interestingly, both parameters exhibit a decreasing trend with increasing frequency. This observed behavior aligns well with previous reports on similar systems.<sup>39</sup> Fig. 9(e) sheds further light on the hopping behavior within  $\text{Rb}_3\text{Bi}_2\text{I}_9$ . It displays the frequency dependence of  $R_{\omega}$ , revealing a decrease in the hopping distance as frequency increases. This trend suggests a shift from long-distance to short-distance tunneling by the charge carriers at higher frequencies. This shift might be linked to the observed increase in  $N(E_{\text{F}})$  (carrier density) with frequency, as reported elsewhere.<sup>39</sup> Another noteworthy observation is the close match between the calculated tunneling distances ( $R_{\omega}$ , ranging from 3.312 to 4.214 Å) and the interatomic distances between Rb and I ions (Rb–I) within the  $\text{Rb}_3\text{Bi}_2\text{I}_9$  crystal lattice (reported to be between 3.273 and 4.224 Å).<sup>48</sup> This similarity suggests that the mobility of  $\text{Rb}^+$  cations might play a significant role. The NSPT model proposes that small polaron motion can leverage the movement of these cations to facilitate AC conductivity within the material.

In this section, we aim to compare the optical and electrical properties of the compounds  $\text{A}_3\text{Bi}_2\text{I}_9$  ( $\text{A} = \text{Cs}, \text{Rb}$ ). The structural disparity between  $\text{Cs}_3\text{Bi}_2\text{I}_9$  and  $\text{Rb}_3\text{Bi}_2\text{I}_9$ , with  $\text{Cs}_3\text{Bi}_2\text{I}_9$  adopting a 0D molecular salt crystal structure and  $\text{Rb}_3\text{Bi}_2\text{I}_9$  featuring 2D layered derivatives of the perovskite structure, likely influences their optical properties.  $\text{Cs}_3\text{Bi}_2\text{I}_9$ 's hexagonal symmetry and disparate  $[\text{Bi}_2\text{I}_9]^{3-}$  anions might lead to a wider bandgap ( $E_{\text{g}}$ ) of 2.03 eV compared to  $\text{Rb}_3\text{Bi}_2\text{I}_9$ 's 1.90 eV, suggesting stronger electronic confinement. Moreover, the higher Urbach energy in  $\text{Cs}_3\text{Bi}_2\text{I}_9$  (0.086 eV) compared to  $\text{Rb}_3\text{Bi}_2\text{I}_9$  (0.044 eV) may signify a broader distribution of localized states, possibly due to its more complex crystal structure, impacting its optical absorption properties. Fig. 10 presents a combined analysis of the electrical conductivity behavior in  $\text{A}_3\text{Bi}_2\text{I}_9$  ( $\text{A} = \text{Cs}, \text{Rb}$ ) at room temperature (300 K). This analysis includes both Nyquist plots and the frequency dependence of AC conductivity ( $\sigma_{\text{AC}}$ ). A key takeaway from this figure is the noticeably larger conductivity observed in  $\text{Rb}_3\text{Bi}_2\text{I}_9$  compared to  $\text{Cs}_3\text{Bi}_2\text{I}_9$ . The specific features of the Nyquist plots and the trend in the  $\sigma_{\text{AC}}$  curves with frequency both support this conclusion. The observed higher conductivity in  $\text{Rb}_3\text{Bi}_2\text{I}_9$  compared to  $\text{Cs}_3\text{Bi}_2\text{I}_9$  may be attributed to the mass difference of the  $\text{A}^+$  cation ( $\text{Cs} = 132.9 \text{ u}/\text{Rb} = 85.4 \text{ u}$ ). With Rb having a smaller mass relative to Cs, it suggests that the delocalization of the A cation in  $\text{Rb}_3\text{Bi}_2\text{I}_9$  is more facile compared to  $\text{Cs}_3\text{Bi}_2\text{I}_9$ . This facilitates easier movement of charge carriers, leading to enhanced AC conductivity, as evidenced by the decrease in hopping distances with increasing frequency in  $\text{Rb}_3\text{Bi}_2\text{I}_9$ . The proximity of the tunneling distances to the interatomic



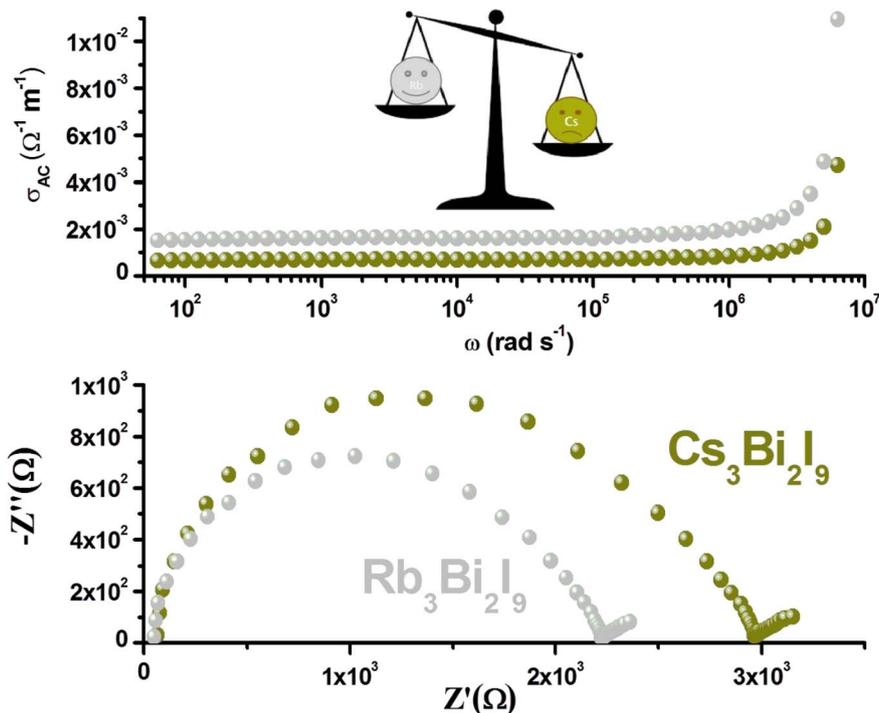


Fig. 10 Unveiling electrical properties of  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ) perovskites: a combined approach using Nyquist plots and AC conductivity (300 K).

distances within the crystal lattice further supports this interpretation, indicating efficient charge transport facilitated by small polaron motion in  $Rb_3Bi_2I_9$ . In addition to the mass difference of the  $A^+$  cation, the disparity in crystal structure between  $Cs_3Bi_2I_9$  (0D) and  $Rb_3Bi_2I_9$  (2D) likely contributes to the variation in conductivity. The 0D molecular salt crystal structure of  $Cs_3Bi_2I_9$ , with its hexagonal symmetry and disparate  $[Bi_2I_9]^{3-}$  anions, may restrict charge carrier mobility compared to the 2D layered derivatives of the perovskite structure in  $Rb_3Bi_2I_9$ . The 2D layered structure of  $Rb_3Bi_2I_9$  facilitates charge transport along the layers, potentially leading to enhanced AC conductivity due to reduced hindrance for charge carriers compared to the more constrained 0D structure of  $Cs_3Bi_2I_9$ .

To further investigate the conduction mechanism and gain insights into the nature of polarons within  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ), transient absorption spectroscopy was employed. This technique involved exciting the samples with a laser pulse centered at a wavelength of 350 nm and monitoring the resulting changes in light absorption over time. By analyzing these transient absorption signals, researchers aimed to elucidate the dominant charge transport mechanisms and characterize the properties of the polarons, which are quasiparticles that influence conductivity in these materials.

Fig. 11(a) and (b) display the measured absorption spectrum after light excitation. A prominent photoinduced absorption (PIA) band is observed within the 400–700 nm wavelength range. This broad band signifies the formation of self-trapped states (STS) just below the material's band edge. The presence of a broad PIA band aligns with established observations during STS formation in other materials.<sup>49,50</sup>

The intensity of the PIA signal exhibits a rise until around 800 femtoseconds (fs), followed by a rapid decay that reduces the signal significantly within approximately 30 ps. Notably, the entire development of PIA kinetics across the probed wavelength range occurs within a timeframe of roughly 400–800 fs. This rapid rise time suggests the absence of a potential barrier separating free excitons (excited electron–hole pairs) and trapped carriers within the self-trapped states.

To further understand the dynamics of the observed photo-induced absorption (PIA), researchers employed a biexponential decay model. This model assumes the decay can be described by the combined effect of two distinct processes, each characterized by its own time constant. The analysis focused on the data obtained at a wavelength of 475 nm, which coincides with the absorbance band edge of  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ). The results revealed that the PIA decay could be effectively fitted using two time constants, denoted as  $\tau_1$  and  $\tau_2$ . Fig. 11(c) and (d) illustrate these findings. For  $Cs_3Bi_2I_9$ , the decay was attributed to  $\tau_1$  with a value of 1.3 ps, accounting for 74.8% of the total decay, and  $\tau_2$  with a value of 6.7 ps, contributing the remaining 25.2%. Similarly, for  $Rb_3Bi_2I_9$ , the decay was described by  $\tau_1$  (1.7 ps, 71.9%) and  $\tau_2$  (6.6 ps, 18.1%).

The biexponential decay model analysis was extended to a broader range, encompassing wavelengths from 450 nm to 650 nm. The focus then shifted to the decay kinetics observed at 600 nm, as shown in Fig. 11(c) and (d) for  $Cs_3Bi_2I_9$  and  $Rb_3Bi_2I_9$ , respectively. These decay profiles could be effectively described by two time constants, denoted as  $\tau_1$  and  $\tau_2$ . For  $Cs_3Bi_2I_9$ , the values were  $\tau_1 = 0.4$  ps and  $\tau_2 = 4.9$  ps, while for  $Rb_3Bi_2I_9$ , they were  $\tau_1 = 0.2$  ps and  $\tau_2 = 3.7$  ps.



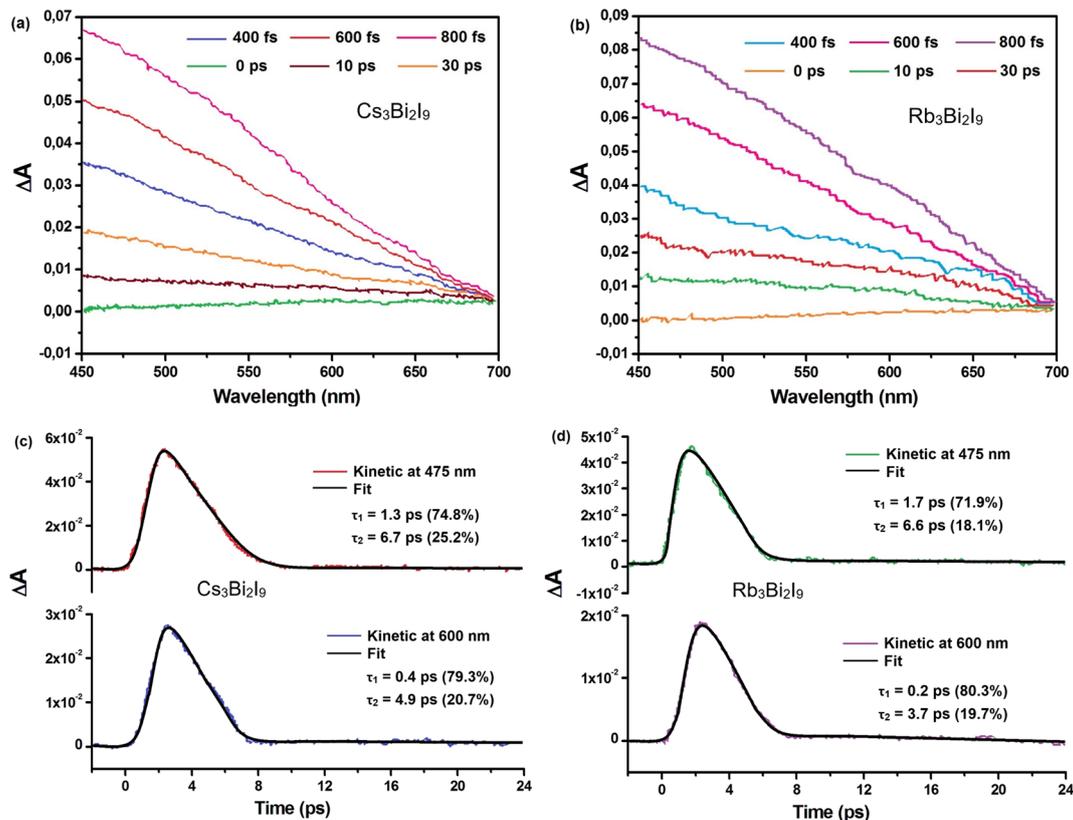


Fig. 11 (a) and (b) Femtosecond-scale transient absorption spectroscopy of  $A_3Bi_2I_9$  (where  $A = Cs$  or  $Rb$ ) following 350 nm optical excitation and probing within the 450–700 nm wavelength range. (c) and (d) Exponential fitting of experimental data indicated by continuous red lines at wavelengths 475 nm and 600 nm.

The ultrafast component,  $\tau_1$ , observed at 475 nm (1.3 ps for  $Cs_3Bi_2I_9$  and 1.7 ps for  $Rb_3Bi_2I_9$ ) is attributed to carrier relaxation *via* coupling with longitudinal optical (LO) phonons. This phenomenon is associated with the formation of Fröhlich polarons, as documented in previous studies.<sup>51,52</sup> Interestingly, the time constant for Fröhlich polaron formation in lead-halide perovskites is typically around 0.2 ps, which closely aligns with the faster time constants observed at 600 nm ( $\tau_1 = 0.4$  ps and 0.2 ps for  $Cs_3Bi_2I_9$  and  $Rb_3Bi_2I_9$ , respectively).

The significant distortion within the soft lattice caused by the Fröhlich polaron can activate various acoustic phonon modes. The interaction between these acoustic phonons and the electrons, along with the ultrafast self-trapping of carriers by these modes, is believed to be responsible for the second decay component observed at both 475 nm and 600 nm. This component is represented by  $\tau_2$  with values of 6.7 ps and 6.6 ps for  $Cs_3Bi_2I_9$  and  $Rb_3Bi_2I_9$  at 475 nm, respectively, and 4.9 ps and 3.7 ps at 600 nm.

The impedance study and absorption analysis together provide a comprehensive understanding of charge transport and polaron dynamics in  $A_3Bi_2I_9$  ( $A = Cs, Rb$ ) compounds. The impedance study offers detailed insights into the suitability of different theoretical models and the impact of structural differences on conductivity. Meanwhile, the absorption analysis complements this by revealing the ultrafast dynamics and

mechanisms of polaron formation and relaxation. By integrating findings from both analyses, a more complete picture emerges, showing how structural and dynamic properties work together to influence the overall behavior of these materials. This complementarity underscores the importance of using multiple analytical approaches to fully characterize complex material systems.

## 4. Conclusions

In sum, this study elucidates the influence of monovalent cations on the conduction mechanism in lead-free  $A_3Bi_2I_9$  perovskites ( $A = Cs, Rb$ ). The synthesis and characterization reveal distinct structural and optical properties between  $Cs_3Bi_2I_9$  and  $Rb_3Bi_2I_9$ , with  $Cs_3Bi_2I_9$  exhibiting a wider bandgap and higher Urbach energy due to its 0D molecular salt crystal structure. Raman spectroscopy revealed distinct vibrational modes in  $Cs_3Bi_2I_9$  and  $Rb_3Bi_2I_9$ , reflecting their differing crystal structures.  $Cs_3Bi_2I_9$  (molecular salt) exhibited a broader range of peaks compared to  $Rb_3Bi_2I_9$  (octahedral  $BiI_6$  units), indicating a more complex vibrational environment within the  $Cs_3Bi_2I_9$  lattice. Furthermore, the observed higher AC conductivity in  $Rb_3Bi_2I_9$  is attributed to the facile delocalization of the smaller mass  $Rb^+$  cation, coupled with its 2D layered structure, enabling more efficient charge transport compared to  $Cs_3Bi_2I_9$ .



This study investigated the carrier dynamics in  $A_3Bi_2I_9$  ( $A = Cs$  or  $Rb$ ) using transient absorption spectroscopy. The results suggest that polarons play a key role in trapping excited charge carriers within the material. Following light excitation, the carriers lose energy through interaction with longitudinal optical (LO) phonons, forming Fröhlich polarons. These polarons further interact with various acoustic phonon modes, leading to their stabilization. The combined effect creates self-trapped states that confine the charge carriers within a limited region of the crystal lattice. This mechanism provides valuable insights into the fundamental processes influencing carrier transport in these materials. These findings underscore the significance of both cation mass and crystal structure in dictating the electrical properties of lead-free perovskite materials, offering insights for their potential applications in optoelectronic devices.

## Data availability

The data that support the findings of this study are available on request from the corresponding author. The data are not publicly available due to [restrictions *e.g.* their containing information that could compromise the privacy of research participants].

## Conflicts of interest

There are no conflicts to declare.

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