


 Cite this: *RSC Adv.*, 2024, **14**, 18646

Enhancing the charge transport and luminescence properties of ethyl 4-[(*E*)-(2-hydroxy-4-methoxyphenyl)methyleneamino]benzoate through complexation: a DFT and TD-DFT study†

Dinyuy Emmanuel Kiven, ^a Fritzgerald Kogge Bine, ^b Nyiang Kennet Nkungli, ^a Aymard Dider Tamafo Fougue, ^c Stanley Numbonui Taskeh ^a and Julius Numbonui Ghogmu ^{*ad}

Organic light emitting diode (OLED) and organic solar cell (OSC) properties of ethyl 4-[(*E*)-(2-hydroxy-4-methoxyphenyl)methyleneamino]benzoate (EMAB) and its Pt^{2+} , Pd^{2+} , Ni^{2+} , Ir^{3+} , Rh^{3+} , and Zn^{2+} complexes have been theoretically studied herein. Geometry optimizations have been performed via the $r^2\text{SCAN-3c}$ composite method while single-point calculations have been carried out at the PBE0-D3(BJ)/def2-TZVP level of theory. Results have shown that complexation with selected metal ions improves hole and electron transfer rates in $\text{Pt}[\text{EMAB}]_2$ and $\text{Rh}[\text{EMAB}]_2^+$. Specifically, the hole transport rate of $\text{Pt}[\text{EMAB}]_2$, ($k_{\text{ct}}(\text{h}) = 6.15 \times 10^{14} \text{ s}^{-1}$), is found to be 44 times greater than that of [EMAB], ($k_{\text{ct}}(\text{h}) = 1.42 \times 10^{13} \text{ s}^{-1}$), whereas electron transport rate of $\text{Pt}[\text{EMAB}]_2$, ($k_{\text{ct}}(\text{e}) = 4.6 \times 10^{13} \text{ s}^{-1}$) is 4 times that of EMAB ($k_{\text{ct}}(\text{e}) = 1.1 \times 10^{13} \text{ s}^{-1}$). Charge mobility for holes and electrons are equal to $19.182 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ and $1.431 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ respectively for $\text{Pt}[\text{EMAB}]_2$, and equal to $4.11 \times 10^{-1} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ and $3.43 \times 10^{-1} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ for EMAB respectively. These results show that, charge transport in EMAB can be tuned for better performance through complexation with transition metals such as Pt^{2+} . OSC properties of the complexes were also studied by comparing their HOMO/LUMO energies with those of (6,6)-phenyl-C₆₁-butyric acid methyl ester (PCBM) and poly(3-hexylthiophene) (P3HT). It turned out that the energy gap of EMAB reduced significantly upon complexation from 2.904 eV to 0.56 eV in $[\text{Rh}(\text{EMAB})_2]^+$ and to a lesser extent in the other complexes. The energy values of the HOMOs remained higher than those of PCBM while those of the LUMOs were found to be greater than that of P3HT with the exception of $[\text{Rh}(\text{EMAB})_2]^+$. These findings show that the aforementioned species are good electron donors to PCBM. The open circuit voltage, V_{OC} , of the compounds ranged between $0.705 \times 10^{-19} \text{ V}$ and $6.617 \times 10^{-19} \text{ V}$, values that are good enough for practical usage in OSC applications. The UV-visible absorption spectra revealed absorption maxima well below 900 nm in all compounds, vital in the efficient functioning of solar cells. In general, this study has shown that platinoid complexation of EMAB can successfully modify both its OLED and OSC properties, making them better precursors in the electronic industry.

Received 24th March 2024
 Accepted 4th June 2024

DOI: 10.1039/d4ra02250e
rsc.li/rsc-advances

1 Introduction

Burgeoning world population has led to an unprecedented rise in the pursuit of better energy-saving lighting devices necessitating the search and manufacture of low-cost, environmentally

friendly and low energy consuming lighting devices or their precursors.^{1,2} Organic electroluminescent materials have attracted wide attention in both academia and industry as a result of their pre-eminence as light emitting materials in flat screen displays.^{3,4} Organic electronic materials have found extensive applications in organic light-emitting diodes (OLEDs), optoelectronics, display and lighting technologies, phototransistors, smartphones, TV applications, personal computers, photodetectors, and bio-imaging.⁴⁻⁷ Their applications in displays are due to several unique advantages such as low cost, high contrast, good colour range, flexibility, wide viewing angle, fast response time, low driving voltage, low weight, low-temperature maintenance, and absence of toxic metals.^{6,8-10} In their mode of operation, OLED devices transform electrical

^aDepartment of Chemistry, Faculty of Science, The University of Bamenda, P. O. Box 39, Bambili, Bamenda, Cameroon. E-mail: ghogsjuju@hotmail.com

^bDepartment of Fundamental and Cross-cutting Sciences, National Advanced School of Public Works, P. O. Box 510, Yaounde, Cameroon, ghogsjuju@hotmail.com

^cDepartment of Chemistry, Ecole Normale Supérieure, Université de Bertoua, P. O. Box 652, Bertoua, Cameroon

[†]Department of Chemistry, Research Unit of Noxious Chemistry and Environmental Engineering, Faculty of Science, University of Dschang, P. O. Box 67, Dschang, Cameroon



energy into light. OLED materials encompass a sequential multilayer arrangement of thin organic films (an emissive layer (EL), a hole-transport layer (HTL) and an electron transport layer (ETL)), all in the middle of two electrodes, assembled *via* a variety of high technologies.^{1,2,9}

A proficient HTL material is endowed with adequate hole injection and high hole mobility, which are primordial properties for effective transmission of charge carriers towards the emissive layer.^{11,12} Moreover, the HTL material is expected to exhibit a highest occupied molecular orbital (HOMO) energy of about 4.5 eV. This HOMO energy facilitates hole injection from the anode into the emissive layer. In addition, an appropriate lowest unoccupied molecular orbital (LUMO) energy level limits the injection of electrons from EL to HTL.¹¹ Similarly, a good ETL material is expected to possess a suitable HOMO and LUMO energy levels capable of accommodating the potential energy barrier, thus minimizing electron injection, lowering the operating voltage and making its hole-blocking ability significant.^{11,13}

Within the framework of OLED manufacture and design, *N,N*-diphenyl-*N,N*-bis(3-methyl phenyl) (1,1-biphenyl)-4,4-diamine (TPD) is the routinely used hole transport layer, whereas tris(8-hydroxyquinolato) aluminium (Alq3) is the habitually used electron transport layer prototype.^{12,13} Polymers and π -conjugated molecules are outstanding for organic-based electronics, owing to their efficient light emission and charge transport properties.^{14,15} Among organic materials used for OLED fabrication, Schiff bases have been portrayed as one of the most important photo-luminescent class of organic compounds, with promising electron transport and emission properties.⁹

Schiff bases, especially the D- π -A types have been remarkable in the manufacture of OLEDs because they are easily synthesisable and modifiable.^{16,17} Although Schiff bases are good electroluminescent materials, they are completely organic and thermally unstable. This obstructs their successful application in displays, lighting and other purposes and consequently their marketability.⁷ Again, OLEDs usually wear out through dark spot degradation, catastrophic failure and intrinsic degradation. Additionally, a noticeable band is always recorded in the luminescence spectrum of organic polymers, diminishing pure colour production ability.^{18,19} These setbacks have therefore motivated the search for more stable electroluminescent materials or their precursors for the manufacture of OLEDs.⁴ Important factors to look out for when searching materials for OLED fabrication include thermal stability, structural rearrangement resistance, and low chemical degradation. Although pure OLEDs have been commercialized as flat panel displays in the past, only single excitons have been involved because emission from spin triplet state is forbidden by spin selection rules (*i.e.* phosphorescence is not possible) at room temperature.^{9,20,21} Transition metal complexes are generally stable and heat resistant; properties highly searched for in OLED design and fabrication.^{3,17} Moreover, effective phosphorescence in late transition metal-Schiff base complexes at room temperature is associated with the heavy atom effect, which is responsible for promoting spin-orbit coupling (SOC). Maximizing the inter-

system crossing (ISC) from the singlet state (with a 25% probability of emission) to the triplet state (with a 75% probability of emission), significantly improves their quality.^{4,17,22}

In the same vein, transition metal complexes, especially platinoid complexes of iridium (Ir), platinum (Pt) and palladium (Pd) demonstrate high stability, are easily purified, and can function as both electron transporters and light emitters with excellent brightness.^{4,11,23} Schiff bases and their associated transition metal complexes have been widely explored for their applications in diverse areas such as nonlinear optics, molecular/metal ion sensing, dye-sensitized solar cells, molecular magnetism and photoluminescence.^{2,3} Also, the ease of synthesis of both Schiff base ligands and their metal complexes offers an additional advantage. Luminescent metal-organic compounds possessing functional emissive ligands represent an upcoming class of luminescent materials owing to their attributes of both organic dyes (tunable and intense emission), and transition-metal-based emitters (high photostability, large absorption-emission Stokes shifts, long emission lifetime, and tuneable excited states).²⁴⁻²⁶

To harvest both singlet and triplet excitons, which are responsible for high performance in OLEDs, heavy transition metals have generally been used in the emissive layer. Heavy metal ions including Ir(III), Rh(III), Pt(II) and Pd(II) have shown such characteristics and have been used in the manufacture of phosphorescent organic light-emitting devices possessing both singlet and triplet excitons.^{4,27-29} When both singlet and triplet excited states are involved, excitons are maximized theoretically to yield phosphorescent OLEDs with close to 100% internal quantum efficiency.^{17,30} Iridium(III) has short lifetimes of triple excitons and therefore distinguishes itself as one of the best candidates for phosphorescent OLEDs. Moreover, Iridium metal is less expensive compared to other metals routinely used for OLEDs fabrication.³¹ Transition metal complexes of Schiff bases are generally stable, owing to numerous binding sites such as nitrogen of the imine group, and oxygen of the carbonyl group which makes them more useful.³

Equally, zinc complexes of azomethine ligands exhibit good electron transport ability, high fluorescence quantum yields, and excellent thermal stability.³²⁻³⁵ Zinc has been shown to produce white OLEDs (WhOLEDs) owing to its d¹⁰ configuration, and its Schiff base complexes have been proven to possess interesting luminescent properties.^{2,36,37} Lastly, zinc(II) complexes, are cheap and universal electroluminescent materials, which have frequently yielded a higher quantum emission than free ligands.^{2,38} Since they are simple to acquire and are easily modifiable, Schiff bases and their metal complexes have been extensively developed. Bearing in mind that transition metal coordination would probably boost the mechanical, luminescent and the thermal stability properties of these molecules, the continuous design of more and more metal-organic based materials with high efficiency for OLEDs is urgent. Fortunately, theoretical studies can be employed in investigating these properties at reduced costs compared to the expensive traditional experimental approaches, notably the spectroradiometer and integrating sphere methods.



Like OLED technology, organic solar cell (OSC) technology has attracted significant attention in this era where green energy is highly solicited due to the adverse effects of fossil fuels. Organic solar cells (OSCs) are photovoltaic devices that convert sunlight into electrical energy. In contrast to conventional silicon-based solar cells, which utilize inorganic materials, OSCs employ carbon-based organic molecules or polymers to absorb and convert sunlight into electrical energy. Organic solar cells have captured both academic and industrial interest owing to their advantages such as lightweight, flexibility and roll-to-roll fabrication capabilities.³⁹ OSC devices with improved light trapping properties are particularly important for electrical energy applications.^{7,40} Although enormous progress has been made in the search and development of new materials, the low efficiency and thermal stability of small organic electronics materials remain a problem for their application.⁷ It is therefore necessary to develop new versatile organic materials with high efficiency for OLEDs and OSCs, which are more stable and capable of functioning as better light emitters, absorbers and charge transporters.⁴¹ Although a great deal of diverse OSC materials have been designed and developed, transition metal complexes have been given very little attention. Schiff base metal complexes are promising in the realm of OSCs. The coordination of transition metals with Schiff bases is readily attainable, and usually results in enhanced charge transfer, mechanical properties, and reduced energy gaps. Consequently, they are emerging as promising acceptor materials for OSCs.⁴² Pasa and colleagues explored the optoelectronic characteristics of a Schiff base ligand and its metal complexes, and observed a reduction in band gap and intriguing charge transfer properties. These features of Schiff base complexes hold significant promise for potential application in solar cell fabrication.⁴³

The aim of this study therefore has been to theoretically investigate the effects of some transition metal complexation on the OLED and OSC properties of the planar Schiff base “ethyl 4-[(E)-(2-hydroxy-4-methoxyphenyl)methyleneamino]benzoate” herein duped EMAB. The metal ions of Pt(II), Pd(II), Ir(III), Ni(II) and Zn(II) were chosen for their ability to improve charge transfer properties and thermal stability when complexed with Schiff bases.^{4,6,7} Also, some of these metal ions were chosen based on their reported ability to attain triplet excited states which leads to Inter System Crossing, a prerequisite for phosphorescence. Interestingly, EMAB and its copper complex were synthesized and characterized by Pahonlu and coworkers in 2015.⁴⁴ Moreover, the effects of metal complexation on the luminescent properties of the Schiff base under study are yet to be investigated to the best of our knowledge. The result of this investigation would provide knowledge on molecular species or precursors for the manufacture of cheaper, more durable, and environmentally friendly electronic materials. The density functional theory (DFT) method (and its time-dependent extension (TD-DFT)) have been used in this investigation because they are a good compromise between computational time and accuracy while concurrently taking into account electron correlation.^{45,46} Fig. 1 shows the proposed mode of complexation between EMAB and the selected metal ions. This mode of complexation was proposed because EMAB is a Schiff

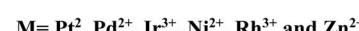
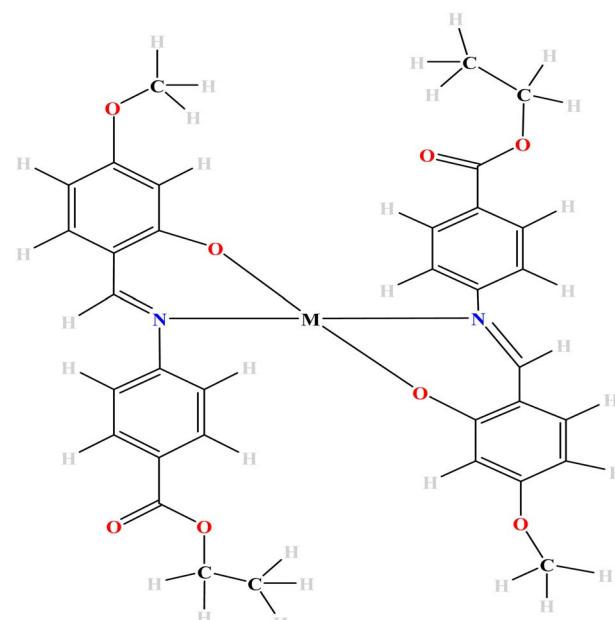


Fig. 1 The proposed mode of complexation between EMAB and selected metal ions.

base and will use only the O and N atoms close to each other for complexation. Given the bulky nature of EMAB is very bulky, it will be difficult to obtain octahedral structures due to steric hindrance reason why only complexes with coordination number of 4 have been modelled and studied herein.

2 Computational details and the theory

2.1 Effects of complexation on organic light-emitting diode (OLED) properties of EMAB

The effects of complexation on the OLED properties of the ligand and its selected complexes were studied by first carrying out geometry optimization of the species using the ORCA 5.0 program package.⁴⁷ The Density functional theory was chosen herein because it has been proven to be suitable for studying coordination compounds.^{45,46} Due to the large number of atoms in the molecules studied, pre-optimization was carried out using the Semi-empirical Quantum Mechanical (SQM) extended Tight-Binding (XTB) density functional method. The Avogadro 1.2.1 visualization program was used to prepare input files for all computations.⁴⁹ After pre-optimization and frequency calculation with the XTB method, the level of theory was augmented, and re-optimization and frequency calculation were carried out with the Swiss Army Knife composite electronic structure (r^2 SCAN-3c) method well embedded in ORCA 5 package.⁴⁸ This functional contains the mGGA functional, r^2 SCAN-3c where “3c” denotes the three corrections-(refitted)



gCP (geometrical counterpoise correction), the D4 correction, and the modification of the basis set. This method was chosen because of its advantages over other mGGAs. The mGGA r^2 SCAN describes electron-correlation effects in a more advanced way, resulting in improved conformational energies, thermochemistry, and refined covalent bond lengths. Moreover, self-interaction error, which most approximate functionals suffer from is reduced.^{50,51} This functional is also more numerically robust and precise than its predecessor (SCAN). Lastly, the r^2 SCAN-3c composite is competitive with more demanding approaches for thermochemistry and conformational energies such as B3LYP and PBE0 functionals, owing to its broader range of applications.⁵² To ascertain that geometry optimization yielded reliable and stable structures (minima on the PES), vibrational frequencies were calculated for all optimized structures at the same level of theory as that used for geometry optimization. Analysis revealed no imaginary frequencies, affirming the reliability of the structures as minima on the PES.

Transition metal ions with partially filled d-orbitals are known to demonstrate multiple spin states. As a result, spin state analysis was carried out by optimizing each complex at all possible spin states at the r^2 SCAN-3c level and determining the single point energies at r^2 SCAN-3c, PBE0/def2TZVP and M06/def2TZVP levels of theory respectively for comparison. The spin states with the lowest energies were then adopted for subsequent calculations. Based on the r^2 SCAN-3c optimized geometries, single-point energy calculations were also performed with the hybrid-GGA functional; PBE0 together with the Ahlrich's basis set def2-TZVP.^{51,53,54}

The PBE0 functional has been used herein owing to its accuracy and agreement with previous experimental results as well as theoretical results obtained at the B3LYP/6-311++g** level of theory in estimating geometrical parameters of EMAB.^{44,55} Moreover, literature indicates that the PBE0 functional, especially when combined with the Def2-TZVP basis set yields good energies for transition metals.^{9,51} Long-range dispersion corrections were done *via* the Grimme's dispersion correction along with the Becke-Johnson damping D3(BJ) scheme.⁵⁶ To determine the reorganization energies, single point calculations were carried out on the neutral, anionic and cationic structures at the RIJ-COSX-PBE0D3(BJ)/def2-TZVP level of theory. RI-J signifies the resolution of the identity approximation, and COSX refers to the chain of spheres approximation.⁵⁴ In this study, property determination was done step by step, by involving the initial screening which was carried out at tight-binding level, followed by structural refinement, and finally energy calculations with inexpensive low-level DFT, and energy re-ranking with a high-level method.⁵²

Two models have been widely used to describe charge transfer properties within entities; the coherent band model and the incoherent hopping model.⁵⁷ The band model is useful for ordered organic crystals, where there exists full delocalization of charge carriers at low temperatures.^{57,58} At elevated temperatures (above 150 K), charge carriers are localized, intermolecular forces weaken, and the band model becomes less relevant. As a result, the hopping model becomes more useful in describing charge transfer rates in organic systems.

The hopping model is described by the Marcus' theory, *via* the following equation⁵⁹

$$k_{ct} = \left(\frac{4\pi^3}{h^2 \lambda k_B T} \right)^{1/2} V_{ij}^2 \exp \left[\frac{-(\Delta G^0 + \lambda)^2}{4\lambda k_B T} \right] \quad (1)$$

where T is the temperature (298.15 K), k_B the Boltzmann constant (1.380×10^{-23} J K⁻¹), h is Planck constant (6.62 $\times 10^{-34}$ J s), ΔG^0 is the standard Gibbs free energy change of the process, λ represents the re-organization energy and V_{ij} the electronic coupling matrix element or transfer integral between two adjacent species i and j , determined by orbital overlap. In Marcus' equation, the values for transfer integrals and reorganization energies are usually converted from eV to Joules (J). Reorganization energies and importantly transfer integrals are fundamental parameters responsible for charge transfer rate.⁵

The reorganization energy, λ is the energy needed to destabilize the reactant and its solvent molecules, from their relaxed nuclear configurations to the relaxed nuclear configurations of the product and its associated solvent molecules.⁶⁰ It is clear from (1) that charge transfer rate depends directly on the square of the transfer integrals and is inversely proportional to the reorganization energy. Therefore, large transfer integrals and smaller reorganization energies are expected to boost the rate of charge transfer. However, re-organization energy is the main parameter determining charge hopping rate, which is generally sectionalized into $\lambda_{internal}$ and $\lambda_{external}$.⁵⁶ The internal reorganization energy is a measure of the structural change between the ionic and neutral states of a molecule, while the external reorganization energy depicts the effect of the surrounding medium on charge transfer.⁶¹ Essentially, $\lambda_{external}$ is not as great as $\lambda_{internal}$.⁶² Presently, attention is focused on $\lambda_{internal}$, since the species under study are isolated and are all in the gas phase. The reorganization energy for the electron λ_e and hole λ_h of the entities studied herein were determined from the following relations.^{5,6,60}

$$\lambda_e = (E_0^- - E_-^-) + (E_-^0 - E_0^0) \quad (2)$$

$$\lambda_h = (E_0^+ - E_+^+) + (E_+^0 - E_0^0) \quad (3)$$

where E_0^- and E_0^+ are the energies of the anion and cation, calculated with the optimized structure of the neutral molecule, E_-^- and E_+^+ are the energies of the anion and cation computed with the optimized anionic/cationic structures. E_-^0 and E_+^0 are the energies of the neutral molecules calculated at the anionic/cationic states and lastly, E_0^0 is the energy of the neutral molecules in the ground state.⁵

The transfer integral V_{ij} describing the ability of the holes and electrons to move from molecule i to j , is a fundamental criterion to consider in charge transfer studies, especially in semiconducting materials.⁶³ Indeed, it reveals the interaction strength between two close species involved in charge transfer.⁵⁸ Two popular methods are generally used to determine charge transfer rates; they are the direct and the indirect methods. The former determines V_{ij} from the electronic coupling element of the two wavefunctions which describe the donor and acceptor



species, while the latter otherwise termed the Energy-Splitting-in-Dimer (KT-ESD) is based on Koopman's theorem (KT). In the KT-ESD method, V_{ij} for the hole and the electron are both computed as half the energy of the HOMOs and LUMOs of two molecular dimers containing species i and j as shown in eqn (4) and (5) respectively.

$$V_e = \frac{1}{2}(E_{L+1} - E_L) \quad (4)$$

$$V_h = \frac{1}{2}(E_H - E_{H-1}) \quad (5)$$

where E_H and E_{H-1} are the energies of the HOMO and HOMO-1 energy level, while E_{L+1} and E_L are the LUMO+1 and LUMO energy levels respectively. The KT-ESD method was used in this study due to its efficiency with symmetric molecules since the molecules under study are symmetric to some extent. Moreover, this method has been instrumental herein due to its efficiency in predicting transfer integrals for organic and semi-organic molecules,⁶⁴ given that the molecules under study are predominantly organic.⁶ Transfer integrals V_{ij} were determined using single point energies of co-facial dimers. The dimers were separated by an intermolecular distance of 4.0 Å, the minimum distance of overlap of donor and acceptor molecular wave functions.⁶⁴

It is worth noting that complete overlap was observed for the ligand and partial overlap for the complexes. The energy of each dimer was obtained by carrying out single point energy calculations for each neutral complex at the PBE0-D3(BJ)/def2-TZVP level of theory, with the RIJCOSX approximation.

In order to determine the stability, which is usually desirable for OLED fabrication,⁵ Chemical hardness; a measure of a species' ability to resist reaction, was calculated using the relationship

$$\eta = \frac{1}{2} \left(\frac{\partial \mu}{\partial N} \right) = \frac{1}{2} \left(\frac{\partial^2 E}{\partial N^2} \right) = \frac{IP - EA}{2} \quad (6)$$

where IP is the ionization potential, and EA is the electron affinity.

The IP and EA, are adiabatic and are calculated as the energy difference between the cationic and neutral species and as the energy difference between the anionic and neutral species respectively.

The ionization potential and electron affinity were obtained using eqn (7) and (8)

$$IP = E(M^+) - E(M^0) \quad (7)$$

$$EA = E(M^0) - E(M^-) \quad (8)$$

where $E(M^0)$, $E(M^+)$ and $E(M^-)$ are the energies of the neutral, cationic, and anionic forms of the investigated molecules respectively. These energies were calculated at the PBE0-D3(BJ)/def2-TZVP level of theory. The stabilities of EMAB and its complexes were determined by calculating their relative hardness. Furthermore, these stabilities were predicted from their interaction energies ΔE_{int} , using the supramolecular approach as follows:

$$\Delta E_{int} = E_{\text{complex}} - (2E_L + E_{M^{n+}}) \quad (9)$$

where E_{complex} is the energy of the complex, E_L is the energy of the ligand and $E_{M^{n+}}$ is the energy of the central metal ion. The energies of the complex, ligand and central metal were calculated at the PBE0-D3(BJ)/def2-TZVP level of theory.

Lastly, the charge transfer mobility, an essential parameter for charge transfer rate studies was determined using Einstein's relation depicted in eqn (10).

$$\mu = \frac{ed^2 k_{ct}}{2k_B T} \quad (10)$$

where μ is the charge transfer mobility, d is the transport distance from the molecular centre to the centre in a dimer (in cm), k_{ct} is the charge transfer rate in s^{-1} , $k_B T$ is the Boltzmann constant and e is the electric charge in eV. The value of k_{ct} is in eV.

2.2 Effects of complexation on the organic solar cell properties of EMAB

The ability of EMAB and its designed complexes to serve as materials for OSCs was determined as follows: Firstly, the HOMO and LUMO energies, (E_{HOMO} and E_{LUMO}), the band gaps ($E_{\text{HOMO}} - E_{\text{LUMO}}$), major FMO-based performance parameters; the open circuit voltage (V_{OC}) and the energy driving force (ΔE_{L-L}) were computed, where (ΔE_{L-L}) is the difference between the LUMO energy levels of the donor (EMAB's complexes) and the standard acceptor (6,6)-phenyl-C₆₁-butyric acid methyl ester PCBM. The open circuit voltage was determined using the following relation:

$$V_{OC} = E_{\text{HOMO}}^{\text{donor}} - E_{\text{LUMO}}^{\text{acceptor}} - 0.3 \quad (11)$$

where V_{OC} is the open circuit voltage, $E_{\text{HOMO}}^{\text{donor}}$ is the HOMO energy of the donor, and $E_{\text{LUMO}}^{\text{acceptor}}$ is the LUMO energy of the acceptor, and 0.3 is an empirical parameter.^{7,58,59}

The energy driving force, ΔE_{L-L} determines the dissociation efficiency of an exciton at the donor/acceptor interface in OSCs. ΔE_{L-L} was calculated using eqn (12). We note here that the numerical value of ΔE_{L-L} must be above the threshold value of 0.3 eV; to ease dissociation.⁶⁵

$$\Delta E_{L-L} = \text{LUMO}^{\text{donor}} - \text{LUMO}^{\text{acceptor}} \quad (12)$$

where ΔE_{L-L} is the determining factor of an exciton's dissociation efficiency at the donor/acceptor interface in OSCs.

For EMAB and its complexes V_{OC} and ΔE_{L-L} were obtained by carrying out single point energy calculations with the r²SCAN-3c optimized structures at the RIJCOSX-PBE0-D3(BJ)/def2TZVP level of theory. The energy level of the HOMO indicates its ability to donate electrons, while the LUMO energy indicates its ability to act as an electron acceptor.⁶⁶ To visualize the HOMO and LUMO charge distributions, frontier molecular surfaces of EMAB and its complexes were generated using Avogadro 1.2.1 visualization program.⁴⁴ Also, the absorption spectra, and the maximum electronic absorption wavelength, λ_{max} of each complex have been computed. The TD-DFT method was used to calculate electronic absorption properties at the CAM-B3LYP/



def2-TZVP level of theory in the gas phase, based on the r^2 SCAN-3c optimized ground state geometries. Indeed, range-separated hybrid functionals like CAM-B3LYP exhibit suitability not only for conjugated systems with $n \rightarrow \pi^*$ and $\pi \rightarrow \pi^*$ electronic transitions but also for molecules characterized by significant charge transfer, such as transition metal complexes.^{67,68} It is worth noting that 10 excited states were computed for each species. Lastly, the density functional theory (DFT) method and its time-dependent extension (TD-DFT) were used in this study owing to the good compromise between computational cost and accuracy they offer, together with the considerable treatment of electron correlation.^{45,46}

3 Results and discussion

3.1 Effects of complexation on OLED properties of EMAB

3.1.1 Optimized geometries. The optimized geometries of EMAB and some of its designed transition metal complexes are shown in Fig. 2, alongside their atomic numbering schemes. Data for optimized structures of EMAB and its studied species is found in the ESI.†

It can be seen from Fig. 2 that the EMAB forms somewhat square planar complexes with the exception of $[\text{Zn}(\text{EMAB})_2]$, which has a distorted tetrahedral structure. Some geometric parameters of EMAB, computed at the PBE0(D3)BJ/def2-TZVP level of theory were compared with experimental data from Pahonțu and co-workers.⁴⁴ In addition, these parameters were also compared with theoretical parameters computed at the B3LYP/6-311++g** level of theory and the agreement between these parameters was close to 100%.⁵⁵ Moreover, some geometric parameters of the copper complex of EMAB, which was synthesized by Pahonțu and co-workers were computed and compared with those of EMAB as shown in Table 1.

Results from Table 1 have revealed excellent agreement between experimental data and data obtained from calculations studies at the B3LYP/6-311++g**.^{44,55} Moreover, the Cu–N and Cu–O bond lengths obtained theoretically fall within the range of acceptable values of 1.9–2.1 Å and 1.8–2 Å respectively.⁶⁹ Since these geometrical parameters fall within the acceptable range, we proceeded to use the PBE0-D3(BJ)/Def2-TZVP level of theory to predict the parameters of the modelled complexes. Some parameters of EMAB and its complexes were selected and presented in Table 2.

The optimized structure of the ligand exhibits some degree of planarity, rigidity and stabilization arising from the intermolecular hydrogen bond O15–H16···N18 with hydrogen bond length 1.702 Å and O–H···N bond angle 148.92°. Non-conventional hydrogen bonds of the form C–H···O are observed in nearly all the metal complexes. These non-covalent interactions contribute to the rigidity and planarity of the investigated compounds, which in turn, are likely to enhance

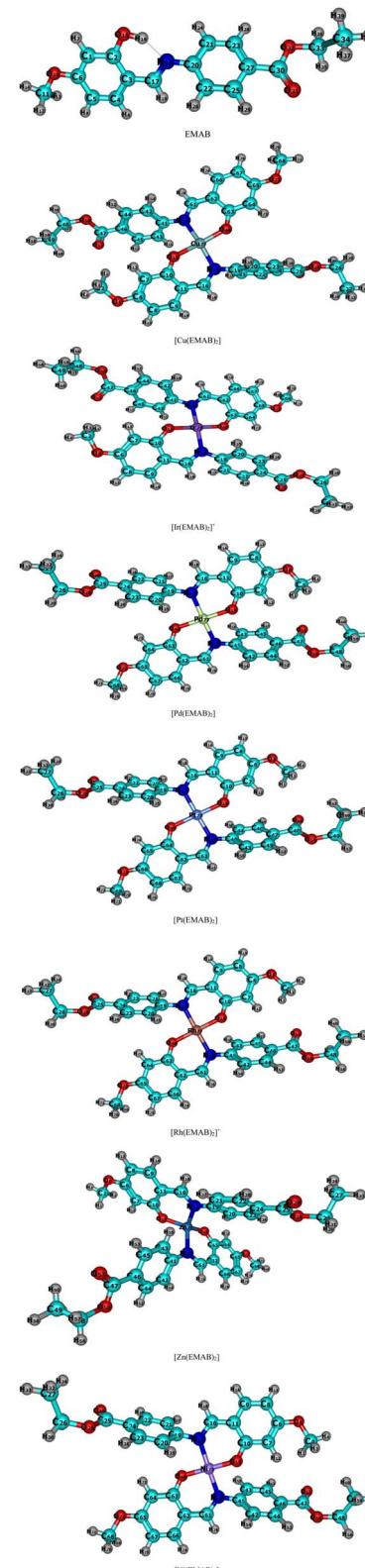


Fig. 2 Optimized structures of EMAB and its designed complexes studied at the r^2 SCAN-3c level of theory.

† Electronic supplementary information (ESI) available: The supplementary file contains the optimized geometrical coordinates of the studied compounds, data used to calculate reorganization energies and UV-visible spectra. See DOI: <https://doi.org/10.1039/d4ra02250e>

intramolecular charge transfer thereby reinforcing photochemical and electronic properties of the molecules.^{70,71}



Table 1 Some selected theoretical geometrical parameters of EMAB and its copper complex, optimized at the PBE0-D3(BJ)/def2-TZVP level of theory, alongside the theoretical parameters of EMAB calculated at the B3LYP/6-311++g** level of theory⁵⁵ and the experimental parameters of Cu[EMAB]₂ (ref. 44)

Selected structural parameters	Theoretical value/level of theory			
	B3LYP/6-311++g**		PBE0-D3(BJ)/def2-TZVP	
	EMAB	EMAB	EMAB	Cu[EMAB] ₂
Bond length (Å)				
N18–C17	1.297	1.29	1.302	1.284
C2–C1	1.394	1.388	1.401	1.386
C1–C6	1.398	1.393	1.38	1.382
N39–Cu77	—	—	2.025	—
Cu77–O40	—	—	1.9	—
Cu77–N17	—	—	2	—
Cu77–O15	—	—	1.9	—
Bond angle (°)				
C30–O32–C33	117	115.9	115.9	116.8
O32–C33–C34	111	111.2	111	107.4
O32–C30–C27	112.5	112.2	112.4	112.4
O31–C30–O32	123.3	123.6	123.5	122.7
O31–C30–C27	124.1	124.2	124.1	124.9
C21–C20–N18	118	118	119.2	116.9
C22–C20–N18	122.9	122.8	121.1	125
C17–N18–C20	121.1	120.8	118	122
N18–C17–C3	122.4	122.1	127	122.3
O15–C2–C3	121.3	121.2	123.6	121.1
O15–C2–C1	118.6	118.7	118.6	118.3
O10–C6–C1	123.8	123.8	123.3	124.7
O10–C6–C5	115.4	115.1	115.5	114.5
C6–O10–C11	119.1	118.2	118.4	117.6
N39–Cu77–O40	—	—	92.9	—
N39–Cu77–O15	—	—	90.7	—
N17–Cu77–O15	—	—	91.4	—
N17–Cu77–O40	—	—	93.3	—
Hydrogen bond parameters				
O15–H16 (Å)	1	1	—	0.9
N18–H16 (Å)	1.7	1.7	—	1.9
O15–H16–N18 (°)	147.6	148.9	—	140.7

Table 2 Some selected geometric parameters of EMAB and its transition metal complexes as predicted at the PBE0-D3(BJ)/Def2-TZVP level of theory

Geometric parameter	[Ir(EMAB) ₂] ⁺	[Pd(EMAB) ₂]	[Rh(EMAB) ₂] ⁺	[Zn(EMAB) ₂]	[Pt(EMAB) ₂]	[Ni(EMAB) ₂]
Bond length (Å)						
N17–M77	2.001	2.035	2.022	2.013	2.028	1.909
N39–M77	2.021	2.027	2.025	2.016	2.021	1.904
O40–M77	1.961	1.997	1.927	1.927	2.010	1.786
O15–M77	1.969	1.993	1.925	1.923	2.006	1.838
Bond angle (°)						
N17–M77–O15	91.1	91.0	89.6	96.4	91.7	92.3
N17–M77–O40	88.9	89.5	90.8	115.8	88.8	87.7
N39–M77–O40	91.6	91.8	90.1	96.3	92.1	92.9
N39–M77–O15	88.3	87.7	89.3	113.7	87.0	87.1
Dihedral angle (°)						
C41–N39–M77–O15	−11.0	−12.0	−9.6	−40.0	−9.5	−16.3
C41–N39–M77–N17	71.4	50.9	32.5	69.8	47.1	59.8



From the values of bond angles in Table 2, it is clear that the geometry adopted by the metal complexes is approximately square planar apart from the $[\text{Zn}(\text{EMAB})_2]$ which is nearly tetrahedral. In all cases, the geometries of the complexes slightly deviate from the 90° bond angles of perfect square planar species and 109.5° of perfect tetrahedral structures. These recorded deviations from perfect planarity or tetrahedral structures about the central metal ions may be due to electron cloud repulsion between benzene rings in each structure (ligand) and due to the Jahn Teller distortions, which causes a distortion in molecular structure and lowers the energy of the molecule.⁷² Furthermore, the C=N bond length (1.29 Å) in the ligand is smaller than that obtained in all complex structures (>1.3 Å), indicating that complexation increases the C=N bond length. Among the complexes studied, metal-ligand bond lengths in $[\text{Ni}(\text{EMAB})_2]$ are the shortest indicating higher stability compared to other complexes.

In this study, metal–oxygen and metal–nitrogen bond lengths in the designed molecules were also compared with experimental values.^{73,74} It is well known that transition metal–nitrogen bond lengths range between 1.9 and 2.03 Å, while metal–oxygen bond lengths range between 1.7 and 2.2 Å. The calculated bond lengths for these structures, at the PBE0D3(BJ)/def2TZVP level of theory are presented in Table 3, alongside the corresponding experimental values.

Results obtained from the comparison between theoretical and experimental parameters in Table 3 have revealed that the level of theory used for determining the geometrical parameters is satisfactory, since theoretical results agree with experiment.

3.1.2 Spin state analysis. It is well known that heavy transition metals exhibit multiple spin states in their complexes. Spin state analysis has therefore been carried out to determine the most stable spin state for each complex. The results for single-point energies at various levels of theory are displayed in Table 4.

Results obtained in Table 4 indicate that the most stable spin state for Pt^{2+} and Pd^{2+} is the +1 state, by virtue of their low energy minima. In the case of Ir^{3+} , results show the most stable spin state of +3, while the +1-spin state is the lowest for the Rh^{3+} .

Table 3 Comparison between predicted metal–oxygen and metal–nitrogen bond lengths at the PBE0D3(BJ)/def2TZVP level of theory and experimentally determined values

Bond lengths	Theoretical value/Å	Experimental value/Å
Pd–N	2.02	2.03
Pd–O	1.99	2.01
Ni–N	1.80	2.00
Ni–O	1.80	1.95
Rh–N	2.02	2.07
Rh–O	1.92	2.03
Zn–N	2.01	2.07
Zn–O	1.92	1.95
Cu–N	2.02	2.01
Cu–O	1.92	1.94
Pt–O	2.01	2.01
Pt–N	2.02	2.00

Table 4 Single point energies for studied EMAB complexes, calculated at $r^2\text{SCAN-3c}$, M06 (D3)/def2TZVP and PBE0-D3(BJ)/def2TZVP for the possible spin states in the gas phase

Complex	Spin multiplicity	Level of theory	Energy/kcal mol ⁻¹
Pd^{2+}	1	$r^2\text{SCAN-3c}$	-2153.8785
		PBE0	-2152.7215
		M06	-2153.6785
	3	$r^2\text{SCAN-3c}$	-2153.8210
		PBE0	-2152.6564
		M06	-2153.5728
Pt^{2+}	1	$r^2\text{SCAN-3c}$	-2145.3235
		PBE0	-2144.137
		M06	-2145.0331
	3	$r^2\text{SCAN-3c}$	-2145.2567
		PBE0	-2144.0597
		M06	-2144.9523
Ir^{3+}	1	$r^2\text{SCAN-3c}$	-2130.0956
		PBE0	-2128.9081
		M06	-2129.7920
	3	$r^2\text{SCAN-3c}$	-2130.0900
		PBE0	-2128.9497
		M06	-2129.8308
Rh^{3+}	1	$r^2\text{SCAN-3c}$	-2130.0019
		PBE0	-2128.8539
		M06	-2129.7342
	3	$r^2\text{SCAN-3c}$	-2136.3114
		PBE0	-2135.1428
		M06	-2136.0365
	5	$r^2\text{SCAN-3c}$	-2136.3059
		PBE0	-2135.1457
		M06	-2136.0307
	5	$r^2\text{SCAN-3c}$	-2136.2343
		PBE0	-2135.0616
		M06	-2135.9524

Zinc has a filled d orbital and therefore would not exhibit multiple spin states. Nickel on the other hand would not exhibit multiple spin states since is not a heavy d-block metal.

3.1.2.1 Reorganization energy and charge transport. Reorganization energy analysis is crucial in charge transport studies.⁴⁰ The calculated reorganization energies for hole/electron transport for EMAB and its metal complexes are displayed in Table 5 together with those of *N,N*-diphenyl-*N,N*-bis(3-methyl phenyl)(1,1-biphenyl)-4,4-diamine (TPD)⁶⁵ and tris(8-hydroxyquinolinato)aluminium (Alq3)^{12,60} as references for hole and electron transport materials, respectively.

Table 5 Calculated reorganization energy values for the hole (λ_h) and the electrons (λ_e) in eV for EMAB and its studied complexes at PBE0-D3(BJ)/def2TZVP in the gas phase

Species	λ_h (eV)	λ_e (eV)	λ_{total} (eV)
EMAB	0.375	0.612	0.987
$[\text{Pt}(\text{EMAB})_2]$	0.185	0.481	0.666
$[\text{Pd}(\text{EMAB})_2]$	0.136	0.489	0.625
$[\text{Zn}(\text{EMAB})_2]$	30.192	30.164	60.356
$[\text{Ir}(\text{EMAB})_2]^+$	0.653	0.925	1.578
$[\text{Rh}(\text{EMAB})_2]^+$	0.171	0.215	0.386
$[\text{Ni}(\text{EMAB})_2]$	0.147	1.12	1.267
TPD	0.290	—	—
Alq3	—	0.276	—



The information that was used to calculate the reorganization energies for the hole and the electron is found in ESI.† It is important to note that reorganization energy is easily estimated theoretically as compared to experimental procedures using spectrochemical and electrochemical methods. Equally, theoretical calculations provide flexibility in exploring different parameters that contribute to the reorganization energy such as molecular structure, external electric field, which are not easily manipulated in experiments.

Generally, molecules with high values of reorganization energies will be poor transporters of holes/electrons, while those with smaller values of reorganization energies will be good transporters. From the results in Table 5, numerical values of λ_h in EMAB complexes reveal that they are good hole transport materials except for the zinc complex. This is because all values of λ_h for EMAB's complexes are less than that of TPD, apart from $[\text{Zn}(\text{EMAB})_2]$ and $[\text{Ir}(\text{EMAB})_2]^+$. Thus the latter will serve as a mild hole transport material. This is obvious because the $[\text{Zn}(\text{EMAB})_2]$ complex has a somewhat tetrahedral structure which usually hinders charge transfer rate, unlike the other square planar complexes which enhance charge transport.⁷⁴ It should be noted also that EMAB, by virtue of its high reorganization energy value for both the hole and the electron will be a mild charge transport material. According to the ranking for λ_h , $[\text{Pd}(\text{EMAB})_2] < [\text{Ni}(\text{EMAB})_2] < [\text{Rh}(\text{EMAB})_2]^+ < [\text{Pt}(\text{EMAB})_2] < [\text{Ir}(\text{EMAB})_2]^+ < \text{EMAB} < [\text{Zn}(\text{EMAB})_2]$, it is clear that complexation improves the hole transport properties of the ligand because the λ_h values are significantly lower than those of the prototype molecule TPD. Although the ligand itself is a mild hole transport material, this property further diminishes when the ligand is complexed with zinc and iridium. Therefore, complex formation between transition metals and ligands such as EMAB can play a significant role in tuning the hole transport properties of the ligand by either significantly improving it as in Pd^{2+} , Ni^{2+} , Rh^{3+} and Pt^{2+} or diminishing it as in Zn^{2+} and Ir^{3+} . From Table 5, numerical values for reorganization energies for the electron λ_e in EMAB complexes are like or lower than those of the prototype to some extent. The $[\text{Rh}(\text{EMAB})_2]^+$ complex proved to be the best electron transport material due to its low λ_e value. Although the complexes have reorganization energies closer to those of the ligand, the zinc complex has a very high value. This suggests that the electron transfer rate in these complexes apart from the zinc complex is improved through complexation. The estimated λ_e values follow the order: $[\text{Rh}(\text{EMAB})_2]^+ < [\text{Pt}(\text{EMAB})_2] < [\text{Pd}(\text{EMAB})_2] < (\text{EMAB}) < [\text{Ir}(\text{EMAB})_2]^+ < [\text{Ni}(\text{EMAB})_2] < [\text{Zn}(\text{EMAB})_2]$. From this ordering and λ_e values in Table 5, it is glaring that transition metal complexation can tune electron transport properties of organic molecules. Our results have also indicated (from reorganization energy values) that among the complexes studied, Rh(III) complex $[\text{Rh}(\text{EMAB})_2]^+$ ($\lambda_e = 0.215$ eV) is the best electron transport material, while $[\text{Pd}(\text{EMAB})_2]$, ($\lambda_h = 0.136$) is the best hole transport material. Consequently, most of the complexes studied are potential charge transport materials for the manufacture of OLEDs. Moreover, some studies have reported that the smaller the sum of the total reorganization energy, the higher the short circuit current density which in turn improves

charge transfer rate.⁶⁶ From Table 5, the sum of reorganization energies is relatively small, except that of $[\text{Zn}(\text{EMAB})_2]$. The total reorganization energy classification is as follows: $[\text{Rh}(\text{EMAB})_2]^+ < [\text{Pd}(\text{EMAB})_2] < [\text{Pt}(\text{EMAB})_2] < [\text{Ir}(\text{EMAB})_2]^+ < ([\text{EMAB}]) < [\text{Ni}(\text{EMAB})_2] < [\text{Zn}(\text{EMAB})_2]$. Therefore, the rate of charge transfer is expected to follow this order. The hole and electron reorganization energy results obtained for the $\text{Zn}[\text{EMAB}]_2$ and $\text{Ir}([\text{EMAB}]_2)^+$ significantly deviate from the standard values (about 100 times and 3 times greater than the standard respectively).

3.1.2.2 Charge transfer integrals and charge transport. When transfer integrals are on average smaller than the reorganization energies as the case in pure organic molecules, Marcus' theory is suitable for carrying out charge transfer studies with the molecules. Transfer integrals were calculated using the Marcus' theory because the entities are predominantly organic. It is also important to note that the higher the value for charge transfer integral, the higher the molecule's ability to transfer charge. The computed charge transfer integral values for both the hole and electron transport processes of the investigated species are listed in Table 6.

Transfer integrals for the electrons range between 0.06 and 0.13 eV following the order $[\text{Rh}(\text{EMAB})_2]^+ < \text{EMAB} < [\text{Pt}(\text{EMAB})_2] = [\text{Pd}(\text{EMAB})_2] < [\text{Ni}(\text{EMAB})_2]$, while that for the hole range from 0.01–0.14 eV in the order $[\text{Rh}(\text{EMAB})_2]^+ < [\text{Pt}(\text{EMAB})_2] = [\text{Pd}(\text{EMAB})_2] < [\text{Ni}(\text{EMAB})_2] < \text{EMAB}$. Although the values for transfer integrals are considerably large for the ligand, the corresponding results for reorganization energy are larger than that of the prototypes Alq3 and TPD and therefore EMAB can be considered as a moderate electron/hole transport molecule.

From Table 6, metal chelation with EMAB increases the value of transfer integrals for the electron in most cases studied but diminishes the hole transport properties in all cases. Moreover, the ligand itself is a good hole and electron transport material by virtue of its high transfer integral values. Based on the reorganization energies, complexation between EMAB and Pd^{2+} , Ni^{2+} and Pt^{2+} increases its hole transport properties while complexation with Rh^{3+} improves its electron transport properties. Based on transfer integrals, complexing EMAB with Ni^{2+} , Pd^{2+} and Pt^{2+} increases transfer integrals for the electrons, thus making them suitable electron transport materials for the fabrication of OLED devices.

3.1.2.3 Charge transfer rate and charge transfer mobility. Marcus' theory relates charge transfer rate to both transfer integrals and reorganization energies. Therefore, studying charge transfer integrals or reorganization energies

Table 6 Calculated transfer integral values for the hole (V_h) and the electrons (V_e) in eV of EMAB and its studied complexes at PBE0-D3(BJ)/def2TZVP in the gas phase

Species	V_e (eV)	V_h (eV)
EMAB	0.11	0.14
$[\text{Pt}(\text{EMAB})_2]$	0.12	0.10
$[\text{Pd}(\text{EMAB})_2]$	0.12	0.10
$[\text{Rh}(\text{EMAB})_2]^+$	0.06	0.01
$[\text{Ni}(\text{EMAB})_2]$	0.13	0.11



independently cannot indicate the best species for the fabrication of OLED materials or their precursors. Herein therefore, the charge transfer rates for EMAB and its studied metal complexes have been computed. Charge transfer rate simultaneously accounts for the reorganization energy and charge transfer integrals of the studied species as seen in eqn (1). Also, charge mobility, another useful parameter in evaluating the performance of OLED materials has been calculated based on Einstein's relation eqn (10). The charge transfer rate and mobility calculated have been presented in Table 7.

Our results unambiguously reveal that complexation between EMAB and majority of the metal ions under investigation increases the hole transfer rate, while the electron transfer rate increases significantly only for $[\text{Pt}(\text{EMAB})_2]$ and $[\text{Rh}(\text{EMAB})_2]^+$.

Charge transfer rates for the hole range from $0.13 \times 10^{13} \text{ s}^{-1}$ to $61.5 \times 10^{13} \text{ s}^{-1}$ and follow the order $[\text{Rh}(\text{EMAB})_2]^+ < \text{EMAB} < [\text{Pd}(\text{EMAB})_2] < [\text{Ni}(\text{EMAB})_2] < [\text{Pt}(\text{EMAB})_2]$. The charge transfer rate for the electrons follows the order $[\text{Ni}(\text{EMAB})_2] < [\text{Pd}(\text{EMAB})_2] < \text{EMAB} < [\text{Rh}(\text{EMAB})_2]^+ < [\text{Pt}(\text{EMAB})_2]$. From the foregoing observations, it can be seen that $[\text{Pt}(\text{EMAB})_2]^+$ is a fascinating ambipolar material for the fabrication of OLED devices since the complex exhibits improved charge transport properties, both for the hole and the electron, compared to the ligand. As expected, charge mobility varies in the same manner as the charge transfer rate. A large value for charge mobility indicates the ease with which charges (holes/electrons) move. For the species under study, the hole mobility rate increases generally with complexation for most of the metals. Charge mobility rate follows the trend $[\text{Rh}(\text{EMAB})_2]^+ < \text{EMAB} < [\text{Pd}(\text{EMAB})_2] < [\text{Ni}(\text{EMAB})_2] < [\text{Pt}(\text{EMAB})_2]$ for both the hole and the electron. At this juncture, it can be concluded that the $[\text{Pt}(\text{EMAB})_2]$ is the best charge transporter amongst the species studied. It is therefore recommended as a good precursor for the manufacture of charge transport for OLED devices.

3.1.2.4 Effect of complexation on charge transfer rate and stability. Investigations on luminescent and/or charge transport

properties of materials requires some basic idea or knowledge of their stability. In fact, materials that are chemically and thermally stable are desirable for OLED fabrication.⁵ Chemical hardness, which is the measure of the ability of a species to resist reaction or the ability to resist exchange in electrons has been calculated through conceptual density functional theory and listed in Table 8, along with the corresponding values for ionization potential, electron affinity, and their change in interaction energies.

Table 8 reveals that the compounds $[\text{Pd}(\text{EMAB})_2]$, $[\text{Pt}(\text{EMAB})_2]$ and $[\text{Rh}(\text{EMAB})_2]^+$ have small energy gaps and are therefore expected to have the most outstanding photophysical properties.⁶⁶ In addition, metal complexation by EMAB is observed to have generally reduced the both the HOMO and LUMO energies, except those of $[\text{Rh}(\text{EMAB})_2]^+$. Furthermore, the energy gaps of the complexes in all cases are smaller than that of the ligand. This indicates that the ease of electron transfer would generally increase with complexation since a low energy gap facilitates electron jump from the HOMO to the LUMO, a primordial factor for photoluminescence. Based on energy gap, EMAB and its complexes are ranked as follows: $[\text{Rh}(\text{EMAB})_2]^+ < [\text{Ni}(\text{EMAB})_2] < [\text{Pt}(\text{EMAB})_2] < [\text{Pd}(\text{EMAB})_2] < \text{EMAB}$. Therefore, metal complexation increases the ease of charge transfer from the HOMO to LUMO and thus OLED properties. The studied complexes of EMAB have higher electron affinity than EMAB and their ionization potentials are higher or very close to that of the ligand. Ionization potentials follow the ranking $[\text{Pt}(\text{EMAB})_2] < [\text{Pd}(\text{EMAB})_2] < [\text{Ni}(\text{EMAB})_2] < \text{EMAB} < [\text{Rh}(\text{EMAB})_2]^+$, while the electron affinity classification varies thus: $[\text{Pd}(\text{EMAB})_2] < \text{EMAB} < [\text{Ni}(\text{EMAB})_2] < [\text{Pt}(\text{EMAB})_2] < [\text{Rh}(\text{EMAB})_2]^+$.

In OSC fabrication, both strong electron donors and acceptors are essential for efficient charge generation and separation. Electron donor materials typically require a low ionization potential and a high HOMO energy level, while electron acceptors should possess a high electron affinity and a low LUMO energy level. Thus, achieving a balance between the

Table 7 Charge transfer rate for the hole $k_{\text{ct}}(\text{h})$ /electron $k_{\text{ct}}(\text{e})$, and mobility rates for the hole u_{h} /electron u_{e} for EMAB and its studied complexes

Species	$k_{\text{ct}}(\text{h}) \times 10^{13} \text{ s}^{-1}$	$k_{\text{ct}}(\text{e}) \times 10^{13} \text{ s}^{-1}$	$u_{\text{h}} (\text{cm}^2 \text{ V}^{-1} \text{ s}^{-1})$	$u_{\text{e}} (\text{cm}^2 \text{ V}^{-1} \text{ s}^{-1})$
EMAB	1.412	1.1	0.441	0.343
$[\text{Pd}(\text{EMAB})_2]$	4.03	0.32	1.26	0.010
$[\text{Pt}(\text{EMAB})_2]$	61.5	4.6	19.182	1.435
$[\text{Rh}(\text{EMAB})_2]^+$	0.13	1.84	0.035	0.574
$[\text{Ni}(\text{EMAB})_2]$	12.7	0.00046	3.963	0.00014

Table 8 Energies of the LUMO, HOMO, energy gap, ionization potential, Electron affinity, chemical hardness and change in interaction energy for EMAB and some of its complexes, calculated at the PBE0-D3(BJ)/def2-TZVP level of theory in the gas phase

Species	E_{L} (eV)	E_{H} (eV)	E_{g} (eV)	IP (eV)	EA (eV)	η (eV)	ΔE_{int} (eV)
EMAB	-2.657	-5.561	2.904	7.370	1.120	3.120	0.000
$[\text{Pd}(\text{EMAB})_2]$	-2.540	-4.700	2.160	6.558	1.100	2.700	-0.977
$[\text{Pt}(\text{EMAB})_2]$	-2.478	-4.534	2.056	5.124	2.340	1.400	-26.997
$[\text{Rh}(\text{EMAB})_2]^+$	-7.621	-8.181	0.560	9.754	6.620	1.600	-25.906
$[\text{Ni}(\text{EMAB})_2]$	-2.705	-4.487	1.782	6.582	1.490	2.550	-0.990



donor and acceptor properties is critical for optimizing device performance.

From our results, $[\text{Pt}(\text{EMAB})_2]$ has the lowest ionization potential, indicating its high electron releasing ability and thus its suitability as an electron donor for OSC technology. Conversely, $[\text{Rh}(\text{EMAB})_2]^+$, with the highest ionization potential value, is the least effective electron donor and therefore unsuitable for this purpose. Ionization potential is found to generally decrease upon complexation, except in the case of $[\text{Rh}(\text{EMAB})_2]^+$, suggesting that complexation enhances electron donation. Additionally, in most cases, complexation increases electron affinity, with $[\text{Rh}(\text{EMAB})_2]^+$ exhibiting the highest potential as an electron acceptor. Consequently, the complexes investigated could serve as electron acceptors in OSC production.

Interestingly, $[\text{Pt}(\text{EMAB})_2]$ demonstrates potential as both an electron donor and acceptor, based on its ionization potential and electron affinity values. However, before final conclusions can be drawn regarding the suitability of $[\text{Pt}(\text{EMAB})_2]$ and $[\text{Rh}(\text{EMAB})_2]^+$ as electron donor and acceptor, respectively, for OSC fabrication, other relevant factors and properties must be considered.

From Table 8, it can be observed that the ligand has the highest η value. This clearly indicates that the ligand is kinetically more stable and therefore less reactive than its corresponding complexes studied. Among the complexes studied, $[\text{Pd}(\text{EMAB})_2]$ is the most kinetically stable while $[\text{Rh}(\text{EMAB})_2]^+$ is the least kinetically stable. It is equally important to note that complexation of the various metals with EMAB reduces chemical stability to some degree. Furthermore, the relative stability of EMAB's complexes was studied by determining their interaction energies ΔE_{int} . This energy was calculated using the supramolecular approach and is presented in Table 8. The more negative the ΔE_{int} value, the more stable the compound. Indeed, the negative values in Table 8 indicate that the formation process of the complex ions from the ligand and the various metal ions is energetically favoured. To sum up, the relatively high η values for $[\text{Pt}(\text{EMAB})_2]$, $[\text{Ni}(\text{EMAB})_2]$ and $[\text{Rh}(\text{EMAB})_2]^+$ indicate that they are fairly stable and can be used as precursors for the construction of OLED devices.

Generally, materials with LUMO energies below the threshold value of -2.7 eV often exhibit a limited capacity for accepting and transporting charges efficiently, which can lead to lower power conversion efficiency (PCE) in OSCs. It can be seen that apart from the LUMO energy of $[\text{Rh}(\text{EMAB})_2]^+$ which is well below -2.7 eV and that of $[\text{Ni}(\text{EMAB})_2]$ which is approximately -2.7 eV, the LUMO energies of the investigated compounds are above the threshold value. Interestingly, PCE

limitation in investigated compounds with $E_{\text{L}} \leq -2.7$ eV is significantly mitigated by their relatively low HOMO-LUMO energy gaps, which can enhance their charge transport capabilities. Other charge transport pathways such as metal-to-ligand charge transfer (MLCT) and/or ligand-to-metal charge transfer (LMCT) are likely in these metal complexes, further mitigating the PCE limitation.

3.2 Effect of complexation on photovoltaic (organic solar cell) properties

Charge transfer in organic solar cells is determined by the donor-acceptor efficiency which in turn is directly dependent on the HOMO and LUMO energy levels of the donor and acceptor. In order to have effective charge transfer within a molecule, the HOMO and LUMO energies of the molecule must be higher than that of (6,6)-phenyl-C₆₁-butyric acid methyl ester (PCBM), for which the HOMO is -6.1 eV and the LUMO is -3.75 eV and lower than that of Poly(3-hexylthiophene) (P3HT) for which the HOMO is -4.65 eV and the LUMO is -2.13 eV,^{7,66,67} which are reference electron acceptor and donor materials respectively.

Table 9 shows the energy levels of the HOMO and LUMO, the energy gap, the open circuit voltage, and the variation in LUMO energy for EMAB and its studied metal complexes in addition to the V_{OC} and $E_{\text{L-L}}$ values. It is worthy of note that, PCBM the electron acceptor prototype for OSCs has been used for comparison because of its stability, its capacity to accept electrons, and its suitable energy level.⁶⁸ From Table 9, the energy gaps for all EMAB complexes are smaller than those of the ligand, indicating that the HOMO-LUMO energy gap is reduced upon complexation. This can lead to increased effective charge transfer from the donor to the acceptor. Also, the energy gaps are observed to range between 0.56 eV and 2.904 eV and show the following trend: $[\text{Rh}(\text{EMAB})_2]^+ < [\text{Ni}(\text{EMAB})_2] < [\text{Pt}(\text{EMAB})_2] < [\text{Pd}(\text{EMAB})_2] < \text{EMAB}$. Relatively small band gaps of the complexes indicate their relative ease with which electrons can be promoted from the HOMOs to the LUMOs. Based on band gap analysis, EMAB is the least reactive and thus the most stable among the molecules studied. Also, the energy values of the HOMOs are higher than that of PCBM with the exception of $[\text{Rh}(\text{EMAB})_2]^+$, while the LUMO energy values are higher than that of P3HT besides $[\text{Rh}(\text{EMAB})_2]^+$. Consequently, these species apart from $[\text{Rh}(\text{EMAB})_2]^+$ are good electron donors to PCBM by virtue of their HOMO-LUMO energy values. As a matter of fact,

Table 9 Energy levels of the HOMO and LUMO, energy gap in eV, open circuit voltage V_{OC} in V and the change in LUMO energy ($\Delta E_{\text{L-L}}$) in eV for EMAB and its studied metal complexes, calculated at the PBE0/Def2-TZVP level of theory

Species	E_{HOMO}	E_{LUMO}	E_{gap}	V_{OC} (eV)	$V_{\text{OC}} (\times 10^{-19} \text{ V})$	$\Delta E_{\text{L-L}}$
EMAB	-5.561	-2.657	2.904	1.500	2.408	1.100
$[\text{Pd}(\text{EMAB})_2]$	-4.700	-2.540	2.160	0.650	1.041	1.210
$[\text{Pt}(\text{EMAB})_2]$	-4.534	-2.478	2.056	0.480	0.769	1.2700
$[\text{Rh}(\text{EMAB})_2]^+$	-8.181	-7.621	0.560	4.130	6.617	-3.85
$[\text{Ni}(\text{EMAB})_2]$	-4.487	-2.705	1.782	0.440	0.705	1.030
PCBM	-6.100	-3.750	—	—	—	—
P3HT	-4.650	-2.130	2.520	0.600	0.961	1.620



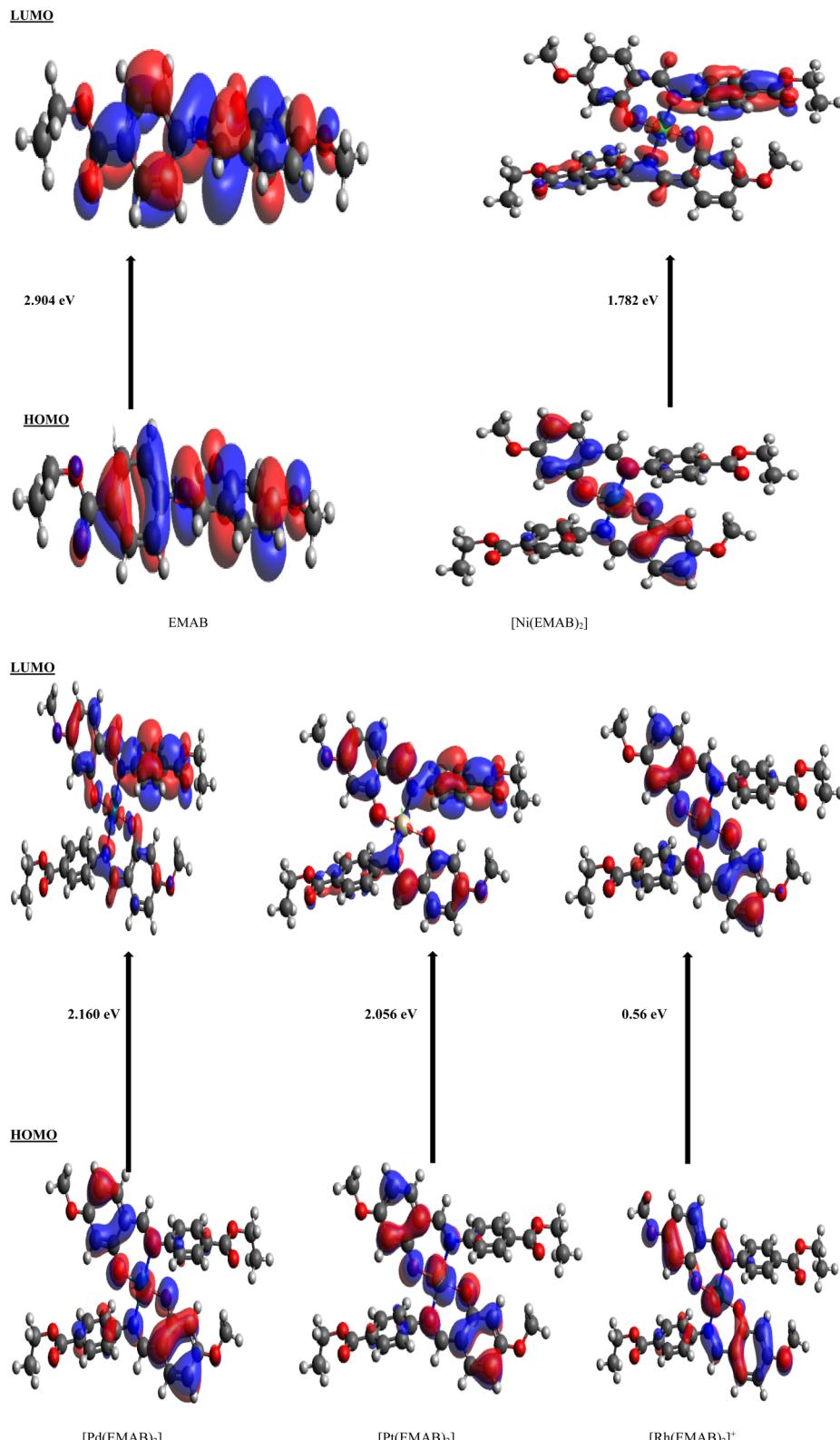


Fig. 3 Iso-surfaces of the FMOs of EMAB and its complexes computed at PBE0-D3(BJ)/def2-TZVP level of theory in gas phase.

all studied molecules stand out as potential donors, since the process of promoting electrons from the valence to the conduction band of the PCBM is feasible and thus good for

donor materials applications in organic solar cells technology. Fig. 3 shows iso-surfaces of the frontier molecular orbitals (FMOs) of EMAB and its complexes computed at PBE0-D3(BJ)/



def2-TZVP level of theory in gas phase, along with the corresponding HOMO–LUMO energy gaps.

A discernible shift in electron density is not apparent from EMAB's HOMO and LUMO distributions, unlike the case with $[\text{Pd}(\text{EMAB})_2]$, $[\text{Ni}(\text{EMAB})_2]$ and $[\text{Pt}(\text{EMAB})_2]$, where there is a clear shift in electron distribution when transitioning from the HOMO to the LUMO. This highlights the fact that metal ion complexation by EMAB can enhance intramolecular charge transfer from one part of the molecule to another. In the case of $[\text{Rh}(\text{EMAB})_2]^+$, there is no significant difference in electron distribution between the HOMO and the LUMO. This indicates that, despite a smaller HOMO–LUMO band gap, intramolecular charge transfer within the molecule is somewhat limited.

3.2.1 Open circuit voltage V_{OC} . The open circuit voltage, V_{OC} , which depends on the HOMO energy of the donor and the LUMO energy of the acceptor is primordial in determining the possibility of electron transfer from the HOMO of the donor to the LUMO of the acceptor. The V_{OC} values for the studied molecules range from 0.705×10^{-19} V to 6.617×10^{-19} V and follow the order; $[\text{Ni}(\text{EMAB})_2] < [\text{Pt}(\text{EMAB})_2] < [\text{Pd}(\text{EMAB})_2] < \text{EMAB} < [\text{Rh}(\text{EMAB})_2]^+$. Although $[\text{Rh}(\text{EMAB})_2]^+$ has a high V_{OC} , its HOMO–LUMO energy levels are low, compared to P3HT and may make it less suitable for use as a charge transport material for the fabrication of OSCs as observed in Table 9. That notwithstanding, EMAB and $[\text{Pd}(\text{EMAB})_2]$ have shown suitable open circuit voltage that are greater than 0.6 eV. On the other hand, $\Delta E_{\text{L-L}}$ values between the studied compounds and the acceptor, range from -3.85 to 1.27 eV. Apart from $[\text{Rh}(\text{EMAB})_2]^+$, all the values are higher than 0.3 eV reminiscent of the electron transfer abilities of these entities to PCBM and thus making them suitable for use in photovoltaic devices. Lastly, the $\Delta E_{\text{L-L}}$ values of the studied molecules are greater than the corresponding V_{OC} values, rendering them (especially in $[\text{Pd}(\text{EMAB})_2]$) highly recommended as materials for the fabrication of organic solar cells.

3.2.2 UV-visible absorption studies. Photovoltaic properties of materials can be greatly controlled by electronic transitions within a molecule. It is important that electronic donor materials have intense absorption properties, compared to the acceptor molecules.⁷ In order that a photovoltaic material generates an electric current, the charge must be transferred from the valence band to the conduction band, and this usually results when the energy in the form of radiation (from UV to visible) is absorbed. This has necessitated investigation of the absorption frequencies of the molecules under study. The electronic absorption frequencies of EMAB and its derivatives,

the corresponding molecular orbital transitions, as well as the oscillator strengths and excitation energies of the investigated molecules are presented in Table 10. All parameters in Table 10 pattern to the electronic transitions with the highest oscillator strength, which should correspond to the most intense band in the UV-vis spectrum. Also, the UV absorption spectra, calculated at the CAM-B3LYP/def2-TZVP level of theory for EMAB and its complexes are presented in Fig. S1, found in the ESI.†

The wavelength for a good photovoltaic material should ideally be low for substantial energy absorption.⁷⁰ For optimal absorption efficiency, an OSC should exhibit absorption onsets that align with the higher-energy regions of the solar spectrum, usually in the visible and near ultraviolet regions. From Table 10, absorption maxima for the molecules are in the order $[\text{Pd}(\text{EMAB})_2] > [\text{Ni}(\text{EMAB})_2] > [\text{Pt}(\text{EMAB})_2] > \text{EMAB} > [\text{Rh}(\text{EMAB})_2]^+$.

EMAB exhibits a maximum absorption peak in its spectra at 307.6 nm, corresponding to excitation energies of 3.785 eV with an oscillator strength (fosc) of 1.0090. This band corresponds to the molecular orbital transition $\text{H} \rightarrow \text{L}$.

The complexes $[\text{Pd}(\text{EMAB})_2]$, $[\text{Ni}(\text{EMAB})_2]$ and $[\text{Pt}(\text{EMAB})_2]$ show absorption maxima at 271.5, 273.4 and 288.8 nm in their spectra, corresponding to excitation energies 4.566 eV, 4.534 eV and 3.870 eV, with oscillator strengths f_{osc} 0.9348, 0.4075 and 0.6936, respectively. These bands correspond to the molecular orbital transitions $\text{H} \rightarrow \text{L}+2$, $\text{H} \rightarrow \text{L}$ and $\text{H}-1 \rightarrow \text{L}+1$, respectively. These species are found to have the lowest absorption frequencies, corresponding to the greatest energies.

These results are so fascinating because they show that complexation between EMAB and the chosen metal ions can significantly alter the absorption frequency. There is a significant drop in the absorption frequency in the complexes $[\text{Pt}(\text{EMAB})_2]$, $[\text{Ni}(\text{EMAB})_2]$ and $[\text{Pd}(\text{EMAB})_2]$. Thus, these complexes, based on their absorption frequencies will carry more energy per photon than EMAB and its other complexes. This will greatly increase the efficiency and therefore energy conversion with materials produced using these complexes.

On the other hand, $[\text{Rh}(\text{EMAB})_2]^+$ complex of EMAB shows absorption spectra maxima band at 538.0 nm, corresponding to the excitation energy of 2.219 eV with oscillator strength f_{osc} 0.3321. These values correspond to the molecular orbital transition $\text{H}-2 \rightarrow \text{L}$. These results show that the absorption energy is significantly reduced and therefore the amount of energy absorbed per photon will be lower.

Although this absorption frequency is higher than that of the parent molecule, it is still within the range at which visible light

Table 10 Absorption spectra data calculated via TD-DFT methods for EMAB and its studied metal complexes at CAM-B3LYP/def2-TZVP level of theory in the gas phase

Species	Transition	λ_{max} (nm)	f_{osc}	Assignment/% contribution	Energy (eV)
EMAB	$\text{S}0 \rightarrow \text{S}1$	307.6	1.0090	$\text{H} \rightarrow \text{L}$ (66%)	3.785
$[\text{Pd}(\text{EMAB})_2]$	$\text{S}0 \rightarrow \text{S}10$	271.5	0.9348	$\text{H} \rightarrow \text{L}+2$ (84%)	4.566
$[\text{Pt}(\text{EMAB})_2]$	$\text{S}0 \rightarrow \text{S}8$	288.8	0.6936	$\text{H}-1 \rightarrow \text{L}+1$ (66%)	3.870
$[\text{Rh}(\text{EMAB})_2]^+$	$\text{S}0 \rightarrow \text{S}10$	538.0	0.3321	$\text{H}-2 \rightarrow \text{L}$ (73%)	2.219
$[\text{Ni}(\text{EMAB})_2]$	$\text{S}0 \rightarrow \text{S}5$	273.4	0.4075	$\text{H} \rightarrow \text{L}$ (37%)	4.534



can be absorbed. Therefore, this complex can still be useful for the fabrication of OSC devices owing to its better mechanical properties owing to the presence of the metal centre.

These values are well below 900 nm, the maximum acceptable wavelength value for photovoltaics. It is also clear from Table 10 that the excitation energies increase as the wavelength reduces, as expected. As mentioned earlier, complexation is observed to increase absorption maxima in some cases (red shift) as in $[\text{Rh}(\text{EMAB})_2]^+$ and reduces in other cases (blue shift) as in $[\text{Pd}(\text{EMAB})_2]$, $[\text{Ni}(\text{EMAB})_2]$ and $[\text{Pt}(\text{EMAB})_2]$. This shows that complexation can conveniently be used to modify the absorption wavelength of a molecule to our interest. When the absorption wavelength is increased, absorption is more in the visible region. It must be indicated that, high-energy photovoltaics offer the advantage of extracting more energy from the sun *i.e.* by using materials that absorb more energy per photon. Photovoltaic materials containing top layers that absorb in the UV region (more energy), while the inner layers absorb at the visible and the innermost layer absorbs at the visible region, offer the advantage of absorbing energy at a range of frequencies. From the foregoing observations, complexing EMAB with metals can tune the absorption wavelength of EMAB, making it useful as a precursor for the fabrication of photovoltaic materials that absorb radiation at a range of frequencies, thus maximizing the amount of energy absorbed (especially during rainy days) from the sun ranging from the UV through visible regions. From the foregoing observations, complexation of EMAB with majority of the metal ions studied can significantly alter the charge transport properties of the molecule and make it a better precursor to the fabrication of OLED and OSC devices, especially the $[\text{Pd}(\text{EMAB})_2]$, $[\text{Ni}(\text{EMAB})_2]$ and $[\text{Pt}(\text{EMAB})_2]$ complexes.

4 Conclusion

Theoretical investigations using density functional theory (DFT) and its time-dependent variant (TD-DFT) have been performed to explore the organic light-emitting diode (OLED) and organic solar cell (OSC) properties of EMAB and its complexes with transition metals (Pt, Pd, Zn, Ir, Rh and Ni). It has been found that transition metal complexation significantly enhances the suitability of EMAB for use in OLEDs and OSCs. According to the results derived from Marcus' theory, the complex $[\text{Pt}(\text{EMAB})_2]$ emerges as the most effective transporter of both holes and electrons with a hole transfer rate ($k_{\text{ct}}(\text{h}) = 6.15 \times 10^{14} \text{ s}^{-1}$) about 44 times greater than that of EMAB ($k_{\text{ct}}(\text{h}) = 1.42 \times 10^{13} \text{ s}^{-1}$). Also, $\text{Pt}[\text{EMAB}]_2$ exhibits the best electron transport rate ($4.6 \times 10^{13} \text{ s}^{-1}$) that is about 4 times greater than that of EMAB ($1.1 \times 10^{13} \text{ s}^{-1}$). It is clear from the calculated hole and electron charge mobility for $\text{Pt}[\text{EMAB}]_2$ ($19.182 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ and $1.435 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$) and for EMAB ($4.41 \times 10^{-1} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ and $3.43 \times 10^{-1} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$, respectively) that the charge transport properties of EMAB are tuneable *via* transition metal coordination, with $\text{Pt}[\text{EMAB}]_2$ being the best charge transporter. To evaluate the applicability of the investigated compounds in OSCs, their HOMO and LUMO energies were compared with those of (6,6)-phenyl-C₆₁-butyric acid methyl ester (PCBM) and

poly(3-hexylthiophene) (P3HT). With the exception of $[\text{Rh}(\text{EMAB})_2]^+$, the HOMO energies were higher than that of PCBM and the LUMO energies were higher than that of P3HT, indicating that these molecules can be efficient electron donors to PCBM. UV-visible absorption studies revealed absorption maxima λ_{max} between 271 nm and 538 nm, in the order: $[\text{Pd}(\text{EMAB})_2] > [\text{Ni}(\text{EMAB})_2] > [\text{Pt}(\text{EMAB})_2] > \text{EMAB} > [\text{Rh}(\text{EMAB})_2]^+$. Accordingly, complexation of EMAB with some transition metal ions can shift its maximum absorption wavelength, which can be beneficial for the manufacture of photovoltaic materials.

Data availability

The data used to support the study's conclusions are contained in the ESI† file and the article.

Author contributions

Conceptualization, investigation, writing-original draft, software: D. E. K.; writing – original draft, investigation, revision, editing, software: F. K. B.; writing – original draft, revision, editing investigation, software N. K. N.; writing – original draft, investigation: A. D. T. F; writing – original draft, Investigation, software: S. N. T.; conceptualization, supervision, validation: J. N. G.

Conflicts of interest

The authors declare no conflicts of interest.

Acknowledgements

We highly acknowledge the financial support (research modernization allowance) from the Ministry of Higher Education of Cameroon.

References

1. A. F. Henwood, A. K. Bansal and D. B. Cordes, Solubilised bright blue-emitting iridium complexes for solution processed OLEDs, *J. Mater. Chem.*, 2016, **4**, 3726–3737, DOI: [10.1039/c6tc00151c](https://doi.org/10.1039/c6tc00151c).
2. S. Kagatikar and D. Sunil, Schiff bases and their complexes in organic light emitting diode application, *J. Electron. Mater.*, 2021, **50**, 1–16, DOI: [10.1007/s11664-021-09197-9](https://doi.org/10.1007/s11664-021-09197-9).
3. A. Sarwar, M. B. Shamsuddin and H. Lingtang, Synthesis, characterization and luminescence studies of metal-diimine complexes, *Mod. Chem. Appl.*, 2018, **6**, 1–7, DOI: [10.4172/2329-6798.1000262](https://doi.org/10.4172/2329-6798.1000262).
4. A. N. Gusev, M. A. Kiskin, E. V. Braga, *et al.*, Schiff base zinc (II) complexes as promising emitters for blue organic light-emitting diodes, *ACS Appl. Electron. Mater.*, 2021, **3**, 3436–3444, DOI: [10.1021/acsaelm.1c00402](https://doi.org/10.1021/acsaelm.1c00402).
5. F. Sun and R. Jin, DFT and TD-DFT study on the optical and electronic properties of derivatives of 1, 4-bis (2-substituted-



1, 3, 4-oxadiazole) benzene, *Arabian J. Chem.*, 2017, **10**, 2988–2993, DOI: [10.1016/j.arabjc.2013.11.037](https://doi.org/10.1016/j.arabjc.2013.11.037).

6 C. H. A. Alongamo, N. K. Nkungli and J. N. Ghogomu, DFT-based study of the impact of transition metal coordination on the charge transport and nonlinear optical (NLO) properties of 2-[(5-(4-nitrophenyl)-1, 3, 4-thiadiazol-2-ylimino] methyl] phenol, *Mol. Phys.*, 2019, **117**, 2577–2592, DOI: [10.1080/00268976.2019.1576932](https://doi.org/10.1080/00268976.2019.1576932).

7 C. I. L. Alongamo, S. N. Taskeh and N. K. Nkungli, Structural, electronic, and charge transport properties of new materials based on 2-(5-mercaptop-1, 3, 4-oxadiazol-2-yl) phenol for organic solar cells and light emitting diodes by DFT and TD-DFT, *J. Chem.*, 2022, **1155**, 1–15, DOI: [10.1155/2022/1802826](https://doi.org/10.1155/2022/1802826).

8 B. X. Yang, C. Yao and G. Zhou, Highly efficient phosphorescent materials based on platinum complexes and their application in organic light-emitting devices (OLEDs), *Platinum Met. Rev.*, 2013, **57**, 2–16, DOI: [10.1595/147106713X659019](https://doi.org/10.1595/147106713X659019).

9 B. Derkowska-Zielinska, M. Barwiolek and C. Cassagne, Nonlinear optical study of Schiff bases using Z-scan technique, *Opt. Laser. Technol.*, 2020, **124**, 1–7, DOI: [10.1016/j.optlastec.2019.105968](https://doi.org/10.1016/j.optlastec.2019.105968).

10 N. M. de Amorim Lima, H. J. Camargo Avila, C. F. do Nascimento Marchiori, *et al.*, Light-emitting porphyrin derivative obtained from a subproduct of the cashew nutshell liquid: A promising material for OLED applications, *Materials*, 2019, **12**, 1–16, DOI: [10.3390/ma12071063](https://doi.org/10.3390/ma12071063).

11 U. Giovanella, M. Pasini, and C. Botta, Organic Light-Emitting Diodes (OLEDs): working principles and device technology, *Applied Photochemistry: when Light Meets Molecules*, 2016, DOI: [10.1007/978-3-319-31671-0_3](https://doi.org/10.1007/978-3-319-31671-0_3).

12 M. Rahman and M. Moniruzzaman, Fundamentals of Organic Light Emitting Diode, *Paper Presented at: Proceedings of 10th Global Engineering, Science and Technology Conference*, 2015.

13 B. C. Lin, C. P. Cheng, Z.-Q. You, *et al.*, Charge transport properties of tris (8-hydroxyquinolinato) aluminum (III): Why it is an electron transporter, *J. Am. Chem. Soc.*, 2017, **127**, 66–67, DOI: [10.1021/ja045087t](https://doi.org/10.1021/ja045087t).

14 E. Mainimo, G. Ejuh and J. Ndjaka, Effect of metalation on some graphene nanoribbons for potential application as donor in organic photovoltaic cells, *J. Mater. Sci.: Mater. Electron.*, 2019, **31**, 21923–21933, DOI: [10.1007/s10854-020-04696-7](https://doi.org/10.1007/s10854-020-04696-7).

15 M. Raftani, T. Abram and W. Loued, The optoelectronic properties of π -conjugated organic molecules based on terphenyl and pyrrole for BHJ solar cells DFT/TD-DFT theoretical study, *Curr. Chem. Lett.*, 2021, **10**, 489–502, DOI: [10.5267/j.ccl.2021.4.002](https://doi.org/10.5267/j.ccl.2021.4.002).

16 R. Jin, X. Zhang and W. Xiao, Theoretical studies of photophysical properties of D- π -A- π -D-type diketopyrrolopyrrole-based molecules for organic light-emitting diodes and organic solar cells, *Molecules*, 2021, **25**, 1–13, DOI: [10.3390/molecules25030667](https://doi.org/10.3390/molecules25030667).

17 B. Blondel, F. Delarue and M. Lopes, Investigation of a sterically hindered Pt (II) complex to avoid aggregation-induced quenching: Applications in deep red electroluminescent and electrical switching devices, *Synth. Met.*, 2017, **227**, 106–116, DOI: [10.1016/j.synthmet.2017.02.021](https://doi.org/10.1016/j.synthmet.2017.02.021).

18 C. Murawski, K. Leo and M. C. Gather, Efficiency roll-off in organic light-emitting diodes, *Adv. Mater.*, 2013, **25**, 6801–6827, DOI: [10.1002/adma.201301603](https://doi.org/10.1002/adma.201301603).

19 H. Lee, J. Lee and J.-I. Lee, Improvement of colour gamut in bottom-emission organic light-emitting diodes using micro-cavity structure embedded cathodes, *Electronics*, 2018, **7**, 1–9, DOI: [10.3390/electronics7090155](https://doi.org/10.3390/electronics7090155).

20 A. Köhler and H. Bässler, Triplet states in organic semiconductors, *Mater. Sci. Eng., R*, 2009, **66**, 71–109, DOI: [10.1016/j.mser.2009.09.001](https://doi.org/10.1016/j.mser.2009.09.001).

21 J. Zhang, F. Zhao and X. Zhu, New phosphorescent platinum (II) Schiff base complexes for PHOLED applications, *J. Mater. Chem.*, 2012, **32**, 16448–16457, DOI: [10.1016/j.mser.2009.09.001](https://doi.org/10.1016/j.mser.2009.09.001).

22 Y. Yao, H.-Y. Yin and Y. Ning, Strong fluorescent lanthanide salen complexes: photophysical properties, excited-state dynamics, and bioimaging, *Inorg. Chem.*, 2019, **58**, 1806–1814, DOI: [10.1021/acs.inorgchem.8b02376](https://doi.org/10.1021/acs.inorgchem.8b02376).

23 T. Yu, W. Su, W. Li, *et al.*, Synthesis, crystal structure and electroluminescent properties of a Schiff base zinc complex, *Inorg. Chim. Acta*, 2006, **359**, 2246–2251, DOI: [10.1016/j.ica.2006.01.019](https://doi.org/10.1016/j.ica.2006.01.019).

24 H. Tanaka, S. Tokito and Y. Taga, Novel metal-chelate emitting materials based on polycyclic aromatic ligands for electroluminescent devices, *J. Mater. Chem.*, 1998, **8**(9), 1999–2003, DOI: [10.1039/a803308k](https://doi.org/10.1039/a803308k).

25 R. C. Evans, P. Douglas and C. J. Winscom, Coordination complexes exhibiting room-temperature phosphorescence: Evaluation of their suitability as triplet emitters in organic light emitting diodes, *Coord. Chem. Rev.*, 2006, **250**, 2093–2126, DOI: [10.1016/j.ccr.2006.02.007](https://doi.org/10.1016/j.ccr.2006.02.007).

26 X. Xu, X. Yang and J. Zhao, Recent advances in solution-processable dendrimers for highly efficient phosphorescent organic light-emitting diodes (PHOLEDs), *Asian J. Org. Chem.*, 2015, **4**, 394–429, DOI: [10.1002/ajoc.201402266](https://doi.org/10.1002/ajoc.201402266).

27 E. Zysman-Colman, *Iridium (III) in Optoelectronic and Photonics Applications*, John Wiley & Sons, 2017.

28 J. Kalinowski, V. Fattori and M. Cocchi, Light-emitting devices based on organometallic platinum complexes as emitters, *Coord. Chem. Rev.*, 2011, **225**, 2401–2425, DOI: [10.1016/j.ccr.2011.01.049](https://doi.org/10.1016/j.ccr.2011.01.049).

29 E. V. Puttock, J. D. Fradgley, D. S. Yufit and J. G. Williams, A family of readily synthesised phosphorescent platinum (II) complexes based on tridentate $N^+ N^+ O$ -coordinating Schiff-base ligands, *Dalton Trans.*, 2019, **48**, 15012–15028, DOI: [10.1039/c9dt03156a](https://doi.org/10.1039/c9dt03156a).

30 A. Odod, E. Nikanova and S. Y. Nikonorov, Electroluminescence of zinc complexes in various OLED structures, *Russ. Phys. J.*, 2017, **60**, 7–13, DOI: [10.1007/s11182-017-1038-2](https://doi.org/10.1007/s11182-017-1038-2).



31 M. Ibrahim-Ouali and F. Dumur, Recent advances on metal-based near-infrared and infrared emitting OLEDs, *Molecules*, 2019, **24**, 1–39, DOI: [10.3390/molecules24071412](https://doi.org/10.3390/molecules24071412).

32 A. Steinegger, O. S. Wolfbeis and S. M. Borisov, Optical sensing and imaging of pH values: Spectroscopies, materials, and applications, *Chem. Rev.*, 2020, **120**, 12357–12489, DOI: [10.1021/acs.chemrev.0c00451](https://doi.org/10.1021/acs.chemrev.0c00451).

33 A. Gusev, E. Braga and Y. Baluda, Structure and tuneable luminescence in polymeric zinc compounds based on 3-(3-pyridyl)-5-(4-pyridyl)-1, 2, 4-triazole, *Polyhedron*, 2020, **191**, 1–9, DOI: [10.1016/j.poly.2020.114768](https://doi.org/10.1016/j.poly.2020.114768).

34 H. Genç Bilgiçli, A. T. Bilgiçli and A. Günsel, Turn-on fluorescent probe for Zn²⁺ ions based on thiazolidine derivative, *Appl. Organomet. Chem.*, 2011, **3**, 3436–3444, DOI: [10.1002/aoc.5624](https://doi.org/10.1002/aoc.5624).

35 L. Kafi-Ahmadi, A. P. Marjani and M. Pakdaman-Azari, Synthesis, characterization and antibacterial properties of N, N'-Bis (4-dimethylaminobenzylidene) benzene-1, 3-diamine as new Schiff base ligand and its binuclear Zn (II), Cd (II) complexes, *S. Afr. J. Chem.*, 2018, **71**, 151–159, DOI: [10.17159/0379-4350/2018/v71a20](https://doi.org/10.17159/0379-4350/2018/v71a20).

36 B. Naskar, R. Modak, D. K. Maiti, *et al.*, A Schiff base platform: structures, sensing of Zn (II) and PPi in aqueous medium and anticancer activity, *Dalton Trans.*, 2017, **46**, 9498–9510, DOI: [10.1039/c7dt01932g](https://doi.org/10.1039/c7dt01932g).

37 Y.-Z. Xie, G.-G. Shan, P. Li, *et al.*, A novel class of Zn (II) Schiff base complexes with aggregation-induced emission enhancement (AIEE) properties: Synthesis, characterization and photophysical/electrochemical properties, *Dyes Pigm.*, 2013, **96**, 467–474, DOI: [10.1016/j.dyepig.2012.09.020](https://doi.org/10.1016/j.dyepig.2012.09.020).

38 D. J. Golding, N. Carter, D. Robinson, *et al.*, Crystallisation-Induced Emission Enhancement in Zn (II) Schiff Base Complexes with a Tuneable Emission Colour, *Sustainability*, 2020, **12**, 95–99, DOI: [10.3390/su12229599](https://doi.org/10.3390/su12229599).

39 Y. Tong, Z. Xiao, X. Du, *et al.*, Progress of the key materials for organic solar cells, *Sci. China: Chem.*, 2020, **63**, 1–8, DOI: [10.1007/s11426-020-9726-0](https://doi.org/10.1007/s11426-020-9726-0).

40 C. H. C. Wallace, K. C. Wai and Y. Yuping, Recent advances in transition metal complexes and light-management engineering in organic optoelectronic devices, *Adv. Mater.*, 2014, **26**, 5368–5398, DOI: [10.1002/adma.201306133](https://doi.org/10.1002/adma.201306133).

41 C. Duan and L. Ding, The new era for organic solar cells: small molecular donors, *Sci. Bull.*, 2020, **42**, 060502, DOI: [10.1016/j.scib.2020.05.019](https://doi.org/10.1016/j.scib.2020.05.019).

42 T. Akitsu, B. Miroslaw and S. Sudarsan, “Photo functions in Hybrid Systems of Schiff Base Metal Complexes and Metal or Semiconductor (Nano) Materials, *Int. J. Mol. Sci.*, 2022, **23**, 1005, DOI: [10.3390/ijms231710005](https://doi.org/10.3390/ijms231710005).

43 S. Pasa, Y. S. Ocak, H. Temel and T. Kilicoglu, “Synthesis, characterization, and catalytic behavior in the Suzuki reaction of Schiff base and its complexes and the optical properties of nickel complex used in the fabrication of a photodiode, *Inorg. Chim. Acta*, 2013, **405**, 493–504, DOI: [10.1016/j.ica.2013.02.038](https://doi.org/10.1016/j.ica.2013.02.038).

44 E. Pahonlu, D.-C. Ilies, S. Shova, *et al.*, Synthesis, characterization, crystal structure and antimicrobial activity of copper (II) complexes with the Schiff base derived from 2-hydroxy-4-methoxybenzaldehyde, *Molecules*, 2015, **20**, 5771–5792, DOI: [10.3390/molecules20045771](https://doi.org/10.3390/molecules20045771).

45 C. J. Cramer and D. G. Truhlar, Density functional theory for transition metals and transition metal chemistry, *Phys. Chem. Chem. Phys.*, 2009, **11**, 10757–10816, DOI: [10.1039/b907148b](https://doi.org/10.1039/b907148b).

46 E. G. Lewars, Introduction to the theory and applications of molecular and quantum mechanics, *Computational Chemistry*, 2011, DOI: [10.1007/978-90-481-3862-3](https://doi.org/10.1007/978-90-481-3862-3).

47 F. Neese, Software update: The ORCA program system—Version 5.0, *Wiley Interdiscip. Rev.: Comput. Mol. Sci.*, 2022, **12**, e1606, DOI: [10.1002/wcms.1606](https://doi.org/10.1002/wcms.1606).

48 C. Bannwarth, S. Ehlert and S. Grimme, GFN2-xTB—An accurate and broadly parametrized self-consistent tight-binding quantum chemical method with multipole electrostatics and density-dependent dispersion contributions, *J. Chem. Theory Comput.*, 2019, **15**, 1652–1671, DOI: [10.1021/acs.jctc.8b01176](https://doi.org/10.1021/acs.jctc.8b01176).

49 M. D. Hanwell, D. E. Curtis, D. C. Lonie, *et al.*, Avogadro: an advanced semantic chemical editor, visualization, and analysis platform, *J. Cheminf.*, 2012, **4**, 1–17, DOI: [10.1186/1758-2946-4-17](https://doi.org/10.1186/1758-2946-4-17).

50 P. Verma and D. G. Truhlar, Status and challenges of density functional theory, *Trends Chem.*, 2020, **2**, 302–318, DOI: [10.1016/j.trechm.2020.02.005](https://doi.org/10.1016/j.trechm.2020.02.005).

51 S. Grimme, J. G. Brandenburg, C. Bannwarth, *et al.*, Consistent structures and interactions by density functional theory with small atomic orbital basis sets, *J. Chem. Phys.*, 2015, **143**, 054107, DOI: [10.1063/1.4927476](https://doi.org/10.1063/1.4927476).

52 S. Grimme, S. Ehrlich and L. Goerigk, Effect of the damping function in dispersion corrected density functional theory, *J. Comput. Chem.*, 2011, **32**, 1456–1465, DOI: [10.1002/jcc.21759](https://doi.org/10.1002/jcc.21759).

53 F. Weigend and R. Ahlrichs, Balanced basis sets of split valence, triple zeta valence and quadruple zeta valence quality for H to Rn: Design and assessment of accuracy, *Phys. Chem. Chem. Phys.*, 2005, **7**, 3297–3305, DOI: [10.1039/b508541a](https://doi.org/10.1039/b508541a).

54 F. Neese, F. Wennmohs, A. Hansen, *et al.*, Efficient, approximate and parallel Hartree–Fock and hybrid DFT calculations. A ‘chain-of-spheres’ algorithm for the Hartree–Fock exchange, *Chem. Phys.*, 2009, **356**, 98–109, DOI: [10.1016/j.chemphys.2008.10.036](https://doi.org/10.1016/j.chemphys.2008.10.036).

55 D. E. Kiven, N. K. Nkungli and S. N. Tasheh, In silico screening of ethyl 4-[*(E*)-(2-hydroxy-4-methoxyphenyl)methyleneamino] benzoate and some of its derivatives for their NLO activities using DFT, *R. Soc. Open Sci.*, 2023, **10**, 1–24, DOI: [10.1098/rsos.220430](https://doi.org/10.1098/rsos.220430).

56 E. Caldeweyher, C. Bannwarth and S. Grimme, Extension of the D3 dispersion coefficient model, *J. Chem. Phys.*, 2017, **147**, 034112, DOI: [10.1063/1.4993215](https://doi.org/10.1063/1.4993215).

57 D. Geng, Y. Chen, Y. Chen, *et al.*, High oxygen-reduction activity and durability of nitrogen-doped graphene, *Energy Environ. Sci.*, 2011, **4**, 760–764, DOI: [10.1039/c0ee00326c](https://doi.org/10.1039/c0ee00326c).

58 J.-L. Brédas, J. P. Calbert, D. da Silva Filho, *et al.*, Organic semiconductors: A theoretical characterization of the basic parameters governing charge transport, *Proceedings of the*



National Academy of Sciences, 2002, DOI: [10.1073/pnas.092143399](https://doi.org/10.1073/pnas.092143399).

59 R. A. Marcus, Chemical and electrochemical electron-transfer theory, *Annu. Rev. Phys. Chem.*, 1964, **15**, 155–196, DOI: [10.1146/annurev.pc.15.100164.001103](https://doi.org/10.1146/annurev.pc.15.100164.001103).

60 D. P. McMahon and A. Troisi, Evaluation of the external reorganization energy of polyacenes, *J. Phys. Chem. Lett.*, 2010, **1**, 941–946, DOI: [10.1021/jz1001049](https://doi.org/10.1021/jz1001049).

61 G. R. Hutchison, M. A. Ratner and T. J. Marks, Hopping transport in conductive heterocyclic oligomers: reorganization energies and substituent effects, *J. Am. Chem. Soc.*, 2005, **127**, 2339–2350, DOI: [10.1021/ja0461421](https://doi.org/10.1021/ja0461421).

62 D. L. Cheung and A. Troisi, Theoretical study of the organic photovoltaic electron acceptor PCBM: Morphology, electronic structure, and charge localization, *J. Phys. Chem.*, 2010, **114**, 20479–20488, DOI: [10.1021/jp1049167](https://doi.org/10.1021/jp1049167).

63 V. Coropceanu, H. Li and P. Winget et al., Electronic-structure theory of organic semiconductors: charge-transport parameters and metal/organic interfaces, *Annu. Rev. Mater. Res.*, 2013, **43**, 63–87, DOI: [10.1146/annurev-matsci-071312-121630](https://doi.org/10.1146/annurev-matsci-071312-121630).

64 M. Ottonelli, M. Piccardo, D. Duce, et al., Tuning the photophysical properties of pyrene-based systems: A theoretical study, *J. Phys. Chem.*, 2012, **116**, 611–630, DOI: [10.1021/jp2084764](https://doi.org/10.1021/jp2084764).

65 N. A. Wazzan, A DFT/TDDFT investigation on the efficiency of novel dyes with ortho-fluorophenyl units (A1) and incorporating benzotriazole/benzothiadiazole/phthalimide units (A2) as organic photosensitizers with D–A2–π–A1 configuration for solar cell applications, *J. Comput. Electron.*, 2019, **18**, 375–395, DOI: [10.1007/s10825-019-01308-4](https://doi.org/10.1007/s10825-019-01308-4).

66 F. Bella, S. Galliano, G. Piana, et al., Boosting the efficiency of aqueous solar cells: A photoelectrochemical estimation on the effectiveness of TiCl₄ treatment, *Electrochim. Acta*, 2019, **302**, 31–37, DOI: [10.1016/j.electacta.2019.01.180](https://doi.org/10.1016/j.electacta.2019.01.180).

67 T. Yanai, D. P. Tew and N. C Handy, A new hybrid exchange-correlation functional using the oulomb-attenuating method (CAM-B3LYP), *Chem. Phys. Lett.*, 2004, **393**, 51–57, DOI: [10.1016/j.cplett.2004.06.011](https://doi.org/10.1016/j.cplett.2004.06.011).

68 J. Creutzberg and E. Hedegård, Investigating the influence of relativistic effects on absorption spectra for platinum complexes with light-activated activity against cancer cells, *Phys. Chem. Chem. Phys.*, 2020, **22**, 27013–27023, DOI: [10.1039/D0CP05143H](https://doi.org/10.1039/D0CP05143H).

69 O. C. Gagne and F. C. Hawthorne, Bond-length distributions for ions bonded to oxygen: Results for the transition metal and quantification of the factors underlying bond-length variation in organic solids, *IUCrJ*, 2020, **7**, 581–629, DOI: [10.1107/S2052252520005928](https://doi.org/10.1107/S2052252520005928).

70 I. Chérif, H. Raissi, K. Abiedh, B. Gassoumi, et al., Computational studies on optoelectronic and nonlinear optical properties of para-substituted nitrobenzofurazan compound, *Mater. Today Commun.*, 2023, **35**, 106133, DOI: [10.1016/j.mtcomm.2023.106133](https://doi.org/10.1016/j.mtcomm.2023.106133).

71 B. Abdelaziz, I. Chérif, B. Gassoumi, S. Patanè and S. Ayachi, Linear and nonlinear optical responses of nitrobenzofurazan-sulfide derivatives: DFT-QTAIM investigation on twisted intramolecular charge transfer, *J. Phys. Chem.*, 2023, **127**, 9895–9910, DOI: [10.1021/acs.jpca.3c04277](https://doi.org/10.1021/acs.jpca.3c04277).

72 M. A. Halcrow, Jahn-Teller distortions in transition metal compounds, and their importance in functional molecular and inorganic materials, *Chem. Soc. Rev.*, 2013, **42**(4), 1784–1795, DOI: [10.1039/c2cs35253b](https://doi.org/10.1039/c2cs35253b).

73 A. Nimmermark, L. Öhrström and J. Reedijk, Metal-ligand bond lengths and strengths: are they correlated? A detailed CSD analysis, *Z. für Kristallogr. - Cryst. Mater.*, 2013, **228**, 311–317, DOI: [10.1524/zkri.2013.1605](https://doi.org/10.1524/zkri.2013.1605).

74 G. L. Miessler and D. A. Tarr, *Inorganic Chemistry*, Pearson, Boston, 5th edn, ISBN 129202075X, 2013.

