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# Discovery of excellent ultraviolet nonlinear optical materials in chlorates and bromates with highly stereochemically active lone pairs†

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The discovery of short-wave ultraviolet (SWUV,  $\lambda_{PM}$  < 266 nm, PM = phase-matching) and deep ultraviolet (DUV,  $\lambda_{PM}$  < 200 nm) nonlinear optical (NLO) crystals is urgently required and full of challenges. Unlike the conventional strategy of using  $\pi$ -conjugated groups (such as BO<sub>3</sub>) as core motifs for constructing UV NLO crystals, herein the long-neglected stereochemically active lone pair (SCALP) groups ClO<sub>3</sub> and BrO<sub>3</sub> are innovatively proposed to be good UV NLO functional motifs based on the group property prediction, and the NLO performance of chlorates and bromates has been investigated systematically for the first time by first-principles methods. Benefiting from the high polarizability of ClO<sub>3</sub> and BrO<sub>3</sub> and their favorable alignments, the halate crystals AClO<sub>3</sub> and ABrO<sub>3</sub> (A = NH<sub>4</sub>, K, Rb and Cs) exhibit high SHG coefficients comparable to that of classical  $\beta$ -BaB<sub>2</sub>O<sub>4</sub> (4.2–4.8 × KDP for AClO<sub>3</sub> and 6.1–7.1 × KDP for ABrO<sub>3</sub>). Meanwhile, their wide band gaps and large optical anisotropy lead to very short  $\lambda_{PM}$  deep into DUV and SWUV (185–195 nm for AClO<sub>3</sub> and 210–220 nm for ABrO<sub>3</sub>). Remarkably, ABrO<sub>3</sub> shows a rare full-wavelength phase-matching capability. Hence AClO<sub>3</sub> and ABrO<sub>3</sub> could be promising DUV and SWUV NLO candidates, respectively, and the UV NLO potential of ClO<sub>3</sub> and BrO<sub>3</sub> is further demonstrated by profound mechanism analysis. This work opens up a new avenue for the development of SWUV and even DUV NLO materials.

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## Introduction

As vital optoelectronic functional materials, nonlinear optical (NLO) crystals can extend limited and fixed laser wavelengths to the ultraviolet and infrared spectral ranges through various kinds of frequency conversion techniques, such as second harmonic generation (SHG), *etc.* Due to their wide applications in many important fields, such as ultrahigh-resolution laser lithography, biomedicine and high-precision scientific instruments, short-wave ultraviolet (SWUV,  $\lambda_{PM}$  < 266 nm, PM = phase-matching) and deep ultraviolet (DUV,  $\lambda_{PM}$  < 200 nm) NLO crystals have attracted ever-growing attention from chemists and materials scientists. <sup>1–5</sup>

The indispensable prerequisite for being an NLO crystal is the noncentrosymmetric structure constructed by NLO functional motifs.<sup>6,7</sup> Over the past decades, several important NLOactive motifs with wide UV transmittance windows, including the  $\pi$ -conjugated BO<sub>3</sub>, CO<sub>3</sub> and C(NH<sub>2</sub>)<sub>3</sub> planar triangles and the  $\sigma$ -bonded BO<sub>n</sub>F<sub>4-n</sub> and PO<sub>n</sub>F<sub>4-n</sub> tetrahedra, have been frequently involved in the exploration of UV NLO materials and some potential SWUV and DUV NLO borates, carbonates, guanidinium salts and phosphates have been developed, including the notable  $\beta$ -BaB<sub>2</sub>O<sub>4</sub> ( $\beta$ -BBO) and KBe<sub>2</sub>BO<sub>3</sub>F<sub>2</sub> (KBBF), as well as the newly discovered CsB<sub>4</sub>O<sub>6</sub>F, C(NH<sub>2</sub>)<sub>3</sub>BF<sub>4</sub>, ABCO<sub>3</sub>F, NaNH<sub>4</sub>PO<sub>3</sub>F·H<sub>2</sub>O, *etc.*<sup>8-12</sup>

The stereochemically active lone pair (SCALP) groups, represented by IO<sub>3</sub>, SeO<sub>3</sub> and TeO<sub>3</sub>, exhibit their natural advantages in the nonlinear optical field because the SCALP electrons can induce cations (I<sup>5+</sup>, Se<sup>4+</sup> and Te<sup>4+</sup>) to undergo second-order Jahn–Teller (SOJT) distortion and endow the groups with large hyperpolarizability and polarizability anisotropy. <sup>13,14</sup> In recent years, numerous excellent NLO crystals with such groups have been discovered, such as AIO<sub>3</sub> (A = Li, Rb, Cs), Li<sub>2</sub>M(IO<sub>3</sub>)<sub>6</sub> (M = Ge, Sn, Ti), BiFSeO<sub>3</sub>, Pb<sub>2</sub>Bi (SeO<sub>3</sub>)<sub>2</sub>Cl<sub>3</sub>, TlSb<sub>3</sub>Te<sub>2</sub>O<sub>12</sub>, *etc*. <sup>15–23</sup> They are mostly only used in the IR NLO field due to their relatively small band gaps. Therefore, the SCALP groups are generally considered to be the IR NLO-active motifs rather than the UV ones. However, our calculations have revealed that other less-common SCALP groups have the potential to be excellent UV NLO-active motifs.

The elements chlorine and bromine are located in the same main group as iodine in the periodic table and have the penta-

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valent ionic forms with  ${\rm Cl}^{5+}{\rm -3s}^2$  and  ${\rm Br}^{5+}{\rm -4s}^2$  SCALP electrons. Like I<sup>5+</sup> in iodates, Cl<sup>5+</sup> and Br<sup>5+</sup> are commonly found in the forms of ClO<sub>3</sub> and BrO<sub>3</sub>, and the coordinated oxygen atoms also adopt a lopsided mode in the known chlorates and bromates.<sup>24,25</sup> We have performed the quantum chemical calculations on the microscopic NLO properties of ClO<sub>3</sub> and BrO<sub>3</sub> using Gaussion09,26 and the results are shown in Fig. 1. First, the HOMO-LUMO gaps of ClO<sub>3</sub> and BrO<sub>3</sub> are calculated to be 8.23 and 7.17 eV, respectively, which fall within the UV (even deep-UV) region and are comparable to those of BO3 and CO3 (especially for ClO<sub>3</sub>); moreover, due to the SCALP effect, ClO<sub>3</sub> and BrO<sub>3</sub> possess relatively large hyperpolarizabilities of 36.08 and 39.53 a.u., respectively, which are superior to BO<sub>3</sub> and CO<sub>3</sub> though inferior to IO3. Meanwhile, their polarizability anisotropy is sufficiently large and comparable to that of IO<sub>3</sub>, enabling them to achieve shorter-wavelength phase-matching SHG output. Therefore, given the favorable properties of the microscopic groups, chlorates and bromates have the potential to be UV (even short-wave UV and deep-UV) NLO candidates. However, few exploratory studies on the NLO properties of chlorates and bromates have been reported to date.<sup>27</sup>

In this work, the potential of chlorates and bromates as UV NLO materials is systematically explored for the first time. We search the Inorganic Crystal Structure Database (ICSD) for noncentrosymmetric ternary chlorates and bromates, excluding fluorides, transition metal and rare earth halates and nonstoichiometric compounds. Cubic symmetric compounds, such as NaClO<sub>3</sub> and NaBrO<sub>3</sub> (P2<sub>1</sub>3, No. 198), are also excluded because they cannot achieve phase-matching. 28,29 In this way, 15 cases of ternary main group chlorates and bromates are finally screened out and their NLO-related properties are carefully studied by first-principles methods. The results indicate that the halate crystals AClO<sub>3</sub> and ABrO<sub>3</sub> (A = NH<sub>4</sub>, K, Rb and Cs) with well-aligned motifs exhibit excellent UV NLO performance, including very short  $\lambda_{PM}$  deep into SWUV and DUV together with high SHG coefficients comparable to β-BaB<sub>2</sub>O<sub>4</sub>,

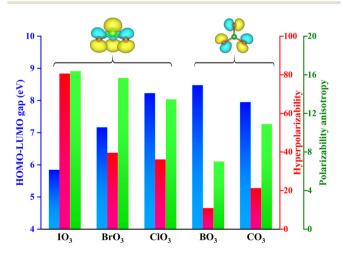


Fig. 1 Calculated HUMO-LUMO band gaps, hyperpolarizabilities and the polarizability anisotropy of ClO<sub>3</sub> and BrO<sub>3</sub> groups. The corresponding data of BO<sub>3</sub>, CO<sub>3</sub> and IO<sub>3</sub> are also calculated for comparison.

being highlighted as important candidates for solar blind and even DUV NLO materials.

## Methods

The crystal structures of the halates are fully relaxed by employing the pseudopotential plane-wave method based on density functional theory within the total energy code of CASTEP. 30-32 Based on the optimized structures, the electronic structures and optical properties are also calculated. We choose the GGA-PBE functional and norm-conserving pseudopotential, in which Cl-3s<sup>2</sup>3p<sup>5</sup>, O-2s<sup>2</sup>2p<sup>4</sup>, K-3s<sup>2</sup>3p<sup>6</sup>4s<sup>1</sup>, Rb-4s<sup>2</sup>4p<sup>6</sup>5s<sup>1</sup>, Cs-5s<sup>2</sup>5p<sup>6</sup>6s<sup>1</sup>, Sr-4s<sup>2</sup>4p<sup>6</sup>5s<sup>2</sup>, Ba-5s<sup>2</sup>5p<sup>6</sup>6s<sup>2</sup>, Tl- $5d^{10}6s^26p^1$ ,  $Pb-5s^25p^65d^{10}6s^26p^2$ ,  $Bi-5d^{10}6s^26p^3$ ,  $N-2s^22p^3$  and H-1s<sup>1</sup> are considered as valence electrons. 33,34 The number of plane waves is determined by the cutoff energies of 820 and 750 eV for Pb(ClO<sub>3</sub>)<sub>2</sub> and other compounds, respectively, while the Monkhorst–Pack k-point separation is set to 0.04  $\mathring{A}^{-1}$  to perform numerical integration of the Brillouin zone. During the optical property calculations, the number of empty bands is set to twice the number of valence bands to ensure convergence of refractive indices and SHG coefficients.

It is well known that DFT-GGA cannot accurately describe the eigenvalues of the electronic states, leading to a quantitative underestimation of the band gaps. Thus, the hybrid exchange and correlation functional of HSE06 is used to obtain more accurate band gaps.<sup>35</sup> When assessing the optical properties, the scissor operators (the gap differences between hybrid HSE06 and conventional GGA-PBE) are applied. This strategy has been demonstrated to be effective and reliable for the evaluation of optical properties of compounds without experimental band gaps.36

The NLO properties of the halates are calculated according to the full sum-over-states (SOS) formalism in the independent particle approximation, 37,38 and the static second-order NLO susceptibility can be expressed as 39,40

$$\chi^{\alpha\beta\gamma} = \chi^{\alpha\beta\gamma} (VE) + \chi^{\alpha\beta\gamma} (VH)$$

where  $\chi^{\alpha\beta\gamma}$  (VE) and  $\chi^{\alpha\beta\gamma}$  (VH) are the contributions to  $\chi^{\alpha\beta\gamma}$  from virtual-electron and virtual-hole processes, respectively, and can be given by

$$\begin{split} \chi^{(2)}_{\alpha\beta\gamma}(\mathrm{VE}) &= \frac{e^3}{2\hbar^2 m^3} \sum_{vcc'} \int \frac{d^3k}{4\pi^3} P(\alpha\beta\gamma) \mathrm{Im} \left[ p^\alpha_{cv} p^\beta_{cc'} p^\gamma_{c'v} \right] \\ &\times \left( \frac{1}{\omega^3_{cv} \omega^2_{vc'}} + \frac{2}{\omega^4_{vc} \omega_{c'v}} \right) \\ \chi^{(2)}_{\alpha\beta\gamma}(\mathrm{VH}) &= \frac{e^3}{2\hbar^2 m^3} \sum_{vv'c} \int \frac{d^3k}{4\pi^3} P(\alpha\beta\gamma) \mathrm{Im} \left[ p^\alpha_{vv} p^\beta_{v'c} p^\gamma_{cv} \right] \\ &\times \left( \frac{1}{\omega^3_{cv} \omega^2_{v'c}} + \frac{2}{\omega^4_{vc} \omega_{cv'}} \right) \end{split}$$

Based on the SOS formalism, the "band-resolved"  $\chi^{(2)}$  can be readily given, and the relationship between each band/ orbital and a specific SHG coefficient  $(d_{ij} = 1/2\chi^{(2)})$  can be conveniently established. Using the normalized SHG value of each band/orbital as the weighting coefficient, and summing all the SHG-weighted bands/orbitals in VB or CB over different k-points in space, the SHG-weighted electron density (SHGdensity) can be obtained. 41 Through the "SHG-density" plots, the crystal orbitals that make significant contributions to the optical nonlinearity can be intuitively visualized.

In addition, to reveal the structural origin of birefringence, polarizability anisotropy-weighted electron density [PAWED,  $\rho_{\Delta y}(r)$ ] has been defined.<sup>42</sup> It contains two parts of contributions (VB and CB) and can be formulated as

$$ho_{\Delta\chi}^{
m VB}(r) = \sum_i^{
m VB} \omega_i |\psi_i(r)|^2$$

$$ho_{\Delta\!\chi}^{ ext{CB}}(r) = \sum_{i}^{ ext{CB}} \omega_i |\psi_i(r)|^2$$

where  $|\psi_i(r)|^2$  is the electron density of the *i*th band/orbital; and  $\omega_i$  is the weighting factor describing the contribution ratio of the ith band/orbital to the total polarizability anisotropy  $(\Delta \chi^{(1)})$  of a crystal. In the low-frequency region,  $\omega_i$  can finally be quantified by the dielectric function  $\varepsilon$ ,

$$\omega_i = rac{\Delta \chi_i^{(1)}}{\Delta \chi^{(1)}} = rac{arepsilon_i(lphalpha) - arepsilon_i(etaeta)}{arepsilon(lphalpha) - arepsilon(etaeta)}$$

By utilizing the PAWED technique, the group/ion contributions to the birefringence (optical anisotropy) of compounds can be identified. 43,44

## Results and discussion

#### Crystal structures

Research Article

Rigorous screening yields 15 cases of simple noncentrosymmetric chlorates and bromates, including AClO<sub>3</sub> (A = NH<sub>4</sub>, K, Rb, Cs, Tl),  $M(ClO_3)_2$  (M = Sr, Ba, Pb),  $ABrO_3$  (A = K, Rb, Cs, Tl),  $M(BrO_3)_2$  (M = Sr, Ba) and  $BiO(BrO_3)^{24,25,45-50}$  We have performed the structural optimization on them, and the structural rationality is proved by the suitable bond valence sum of 4.8-5.2 for Cl<sup>5+</sup>/Br<sup>5+</sup> in ClO<sub>3</sub>/BrO<sub>3</sub> groups. They are all in the polar crystal symmetries and most of them with similar chemical formula are isostructural and the crystal structures resemble one another.

For the alkali metal/alkaline earth metal and NH4 halates, the anionic groups ClO<sub>3</sub>/BrO<sub>3</sub> are all isolated, so the geometric structures can be regarded as zero-dimensional. AClO3 and ABrO<sub>3</sub> (A = NH<sub>4</sub>, K, Rb, Cs, Tl) are crystallized in the trigonal R3m space group (No. 160) and the ClO<sub>3</sub>/BrO<sub>3</sub> themselves have C3v symmetry with three identical Cl-O/Br-O bonds and all the isolated groups show perfect alignment with the dipoles pointing towards the crystallographic axis c (Fig. 2a).  $^{24,2\overline{5},45-47}$  $M(ClO_3)_2$  and  $M(BrO_3)_2$  (M = Sr, Ba) belong to the lower-symmetry monoclinic or orthorhombic crystal systems: Sr(BrO<sub>3</sub>)<sub>2</sub> crystallizes in the Cc space group (No. 9), while Sr(ClO<sub>3</sub>)<sub>2</sub>, Ba  $(ClO_3)_2$  and Ba $(BrO_3)_2$  are in the *Fdd2* space group (No. 43). <sup>48,49</sup> They are all characterized by distorted triangular pyramidal ClO<sub>3</sub>/BrO<sub>3</sub> and ionically bonded dodecahedral MO<sub>8</sub>.

It is interesting that TI+ also adopts a lopsided coordination style with three equivalent oxygen atoms in TlClO<sub>3</sub> and TlBrO<sub>3</sub> (Fig. 2b), but the stereochemical activity of the Tl<sup>+</sup>-6s<sup>2</sup> lone pair is to be investigated. Pb(ClO<sub>3</sub>)<sub>2</sub> is isostructural to Sr and Ba chlorates, and the Pb<sup>2+</sup> cations are bisphenoidally coordinated with eight oxygen atoms from eight different ClO3, forming dodecahedral geometries (Fig. 2c). 49 The stereochemical activity of the Pb<sup>2+</sup>-6s<sup>2</sup> lone pairs is not obvious because the eight Pb-O bond lengths are relatively close, but this remains to be verified by electron localization function analysis. BiO(BrO<sub>3</sub>) crystallizes in the polar space group Pca21 (No. 29) and exhibits a layered structural topology with the BrO<sub>3</sub> neatly hanging up and down from the infinite [Bi<sub>2</sub>O<sub>2</sub>]<sub>∞</sub> layers (Fig. 2d).<sup>50</sup>

#### Nonlinear optics-related properties

For the noncentrosymmetric halates, we focus on their band gaps (UV transmittance), second-order nonlinear optical coefficients, birefringence (optical anisotropy) and phase-matching

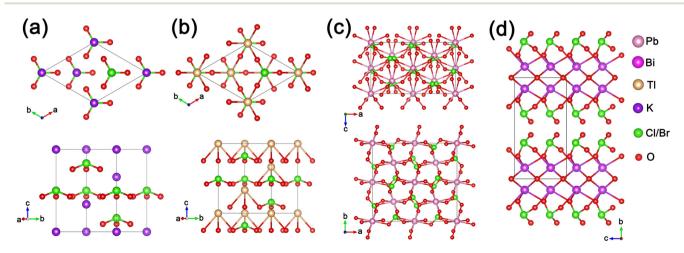


Fig. 2 Crystal structures of KClO<sub>3</sub> (a), TlClO<sub>3</sub> (b), Pb(ClO<sub>3</sub>)<sub>2</sub> (c) and BiO(BrO<sub>3</sub>) (d).

properties, all of which are closely correlated to the practical applications of NLO materials. The calculated NLO-related property data of the halates are listed in Table 1, and RbIO<sub>3</sub> and CsIO<sub>3</sub> with available experimental/theoretical data are taken as references to check our computations. It is found that our calculated optical properties of RbIO3 and CsIO3 are in good agreement with the reported results, fully demonstrating the rationality and reliability of the calculation methods.

From the band structure data obtained by HSE06, we can see that there is an absolutely clear tendency for the band gaps of the halates being  $E_g(\text{chlorate}) > E_g(\text{bromate}) >$  $E_o(\text{iodate})$  (Fig. S1,† Fig. 3a and Table 1), which is in agreement

with the results of the microscopic groups. The band gaps of alkali metal/alkaline earth metal chlorates are as wide as 6.8-7.0 eV, being comparable to or larger than those of some classical borates, such as K<sub>2</sub>Al<sub>2</sub>B<sub>2</sub>O<sub>7</sub> (6.89 eV), CsLiB<sub>6</sub>O<sub>10</sub> (6.89 eV), Cs<sub>2</sub>Al<sub>2</sub>(B<sub>3</sub>O<sub>6</sub>)<sub>2</sub>O (7.05 eV), Rb<sub>3</sub>Ba<sub>3</sub>Li<sub>2</sub>Al<sub>4</sub>B<sub>6</sub>O<sub>20</sub>F (6.26 eV),  $K_3Sr_3Li_2Al_4B_6O_{20}F$  (6.53 eV), LaBGeO<sub>5</sub> (6.40 eV), the wellknown UV NLO crystal β-BaB<sub>2</sub>O<sub>4</sub> (6.43 eV), etc.<sup>51–56</sup> Hence, they are all deep-UV transparent, which is the most important requirement for DUV NLO materials. For alkali metal/alkaline earth metal bromates, the band gaps range from 5.65 to 6.18 eV, comparable to or larger than many carbonates, such as  $ABCO_3F$  (A = alkali metal, B = alkaline earth metal,  $\sim$ 6.20 eV),

**Table 1** Space groups, calculated band gaps ( $E_a$ ) by hybrid HSE06, largest SHG coefficients ( $d_{ij}$ ), birefringence at 1064 nm ( $\Delta n$ ), UV cutoff wavelengths ( $\lambda_{cutoff}$ ), shortest phase-matching wavelengths ( $\lambda_{PM}$ ) and potential SHG applications for the halates

Material	Space group	$E_{\rm g}$ (eV)	$d_{ij}$ (pm V <sup>-1</sup> )	$\Delta n$	$\lambda_{\mathrm{cutoff}}\left(nm\right)$	$\lambda_{PM}$ (nm)	Potential SHG application
(NH <sub>4</sub> )(ClO <sub>3</sub> )	R3m	6.79	$d_{11}$ = 1.80	0.168	182.6	194.9	DUV
KClO <sub>3</sub>	R3m	6.86	$d_{11} = 1.85$	0.186	180.8	185.3	DUV
$RbClO_3$	R3m	6.81	$d_{11} = 1.88$	0.176	182.1	189.9	DUV
$CsClO_3$	R3m	6.99	$d_{11} = 1.63$	0.159	177.4	188.7	DUV
$TlClO_3$	R3m	5.23	$d_{11} = 2.37$	0.173	237.1	273.9	1064 → 532 nm
$Sr(ClO_3)_2$	Fdd2	7.03	$d_{24} = 1.23$	0.084	176.4	260.5	532 → 266 nm
$Ba(ClO_3)_2$	Fdd2	6.97	$d_{24} = 1.19$	0.089	177.9	255.7	532 → 266 nm
$Pb(ClO_3)_2$	Fdd2	5.62	$d_{24} = 2.37$	0.102	220.6	336.5	1064 → 532 nm
$KBrO_3$	R3m	5.65	$d_{11} = 2.76$	0.227	219.5	219.5	532 → 266 nm
$RbBrO_3$	R3m	5.77	$d_{11} = 2.61$	0.218	214.9	214.9	532 → 266 nm
$CsBrO_3$	R3m	5.90	$d_{11} = 2.39$	0.203	210.2	210.2	532 → 266 nm
$TlBrO_3$	R3m	5.00	$d_{11} = 3.19$	0.213	248.0	278.8	1064 → 532 nm
$Sr(BrO_3)_2$	Cc	6.14	$d_{26} = 1.64$	0.062	202.0	352.4	1064 → 532 nm
$Ba(BrO_3)_2$	Fdd2	6.18	$d_{24} = 1.80$	0.091	200.6	299.6	1064 → 532 nm
$BiO(BrO_3)^{50}$	$Pca2_1$	3.89	$d_{15} = 2.35$	0.040	318.8	815.5	NPM in UV-Vis SHG output
		$3.52^a, 3.77^b$					
RbIO <sub>3</sub> <sup>16</sup>	R3m	4.06	$d_{31} = 3.08 \ 20 \times \text{KDP}^a$	0.065	305.4	433.0	1064 → 532 nm
		$4.0^a$		$0.063^{b}$			
CsIO <sub>3</sub> <sup>17</sup>	R3m	4.38	$d_{31} = 1.99$	0.152	283.1	331.1	1064 → 532 nm
		$4.2^{a}$	$15 \times \text{KDP}^a$	$0.19^{b}$			

<sup>&</sup>lt;sup>a</sup> Experimentally measured data in references. <sup>b</sup> Theoretically calculated data in references.

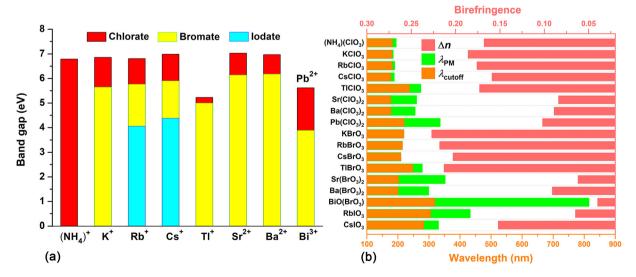


Fig. 3 (a) Distribution of band gaps of the halates and (b) summarization of the birefringence, UV cutoff wavelengths ( $\lambda_{\text{cutoff}}$ ) and the shortest phase-matching wavelengths ( $\lambda_{PM}$ ) for the halates.

 $Cs_3Ba_4(CO_3)_3F_5$  (5.90 eV),  $[C(NH_2)_3]_2Zn(CO_3)_2$  (5.90 eV), Zn (NH<sub>3</sub>)CO<sub>3</sub> (6.08 eV), NaZnCO<sub>3</sub>F (4.61 eV), Na<sub>4</sub>Zn(CO<sub>3</sub>)<sub>3</sub> (4.06 eV), etc. 9,57-59 Although the band gaps of the bromates are narrower than those of the corresponding chlorates by up to 0.79-1.21 eV, they still can be used as UV optical materials considering the good transparency in the short-wave UV region. It is worth noting that the Tl, Pb and Bi halates have much narrower band gaps than the alkali metal/alkaline earth metal halates, which is caused by the intrinsic intra-shell s-p transitions in the Tl<sup>+</sup>, Pb<sup>2+</sup> and Bi<sup>3+</sup> cations.

From Table 1, we can see that for the alkali metal and ammonium chlorates with R3m symmetry, the largest SHG tensor  $d_{11}$  values are 1.63-1.88 pm  $V^{-1}$ , being up to 4.2-4.8 times that of KDP and 0.7-0.8 times that of β-BaB<sub>2</sub>O<sub>4</sub>, which have similar DUV transmittance to these chlorates. In comparison, the alkaline earth metal chlorates with Fdd2 symmetry have relatively smaller SHG coefficients, where  $d_{24}$  = 1.23 pm  $V^{-1}$  for  $Sr(ClO_3)_2$  and 1.19 pm  $V^{-1}$  for  $Ba(ClO_3)_2$ , but they are still 3 times larger than that of KDP and 0.5 times that of β-BaB<sub>2</sub>O<sub>4</sub>. Furthermore, the largest SHG tensors of TlClO<sub>3</sub> ( $d_{11}$ ) and Pb(ClO<sub>3</sub>)<sub>2</sub> ( $d_{24}$ ) share the same value of 2.37 pm V<sup>-1</sup> and are 6.1 times larger than that of KDP.

Not surprisingly, the SHG responses of the bromates are stronger than those of the corresponding chlorates, and there is a 35%-49% increase in SHG coefficients, resulting from the greater hyperpolarizability of BrO3 compared to ClO3. Remarkably, the SHG effects of ABrO<sub>3</sub> (A = K, Rb, Cs;  $d_{11}$  = 2.39-2.76 pm  $V^{-1} \approx 6.1$ -7.1 × KDP) are greater than those of many NLO carbonates and cyanurates with similar band gaps, such as KCdCO<sub>3</sub>F (5.46 eV & 5.2 × KDP), KLi(HC<sub>3</sub>N<sub>3</sub>O<sub>3</sub>)·2H<sub>2</sub>O  $(5.23 \text{ eV } \& 5.3 \times \text{KDP}), \text{ Ln}_5(\text{C}_3\text{N}_3\text{O}_3)(\text{OH})_{12} \text{ (Ln = Y, Lu;}$ 5.28-5.51 eV & 2.5-4.2 × KDP), etc. 60-62 Meanwhile, Sr(BrO<sub>3</sub>)<sub>2</sub> in Cc and Ba(BrO<sub>3</sub>)<sub>2</sub> in Fdd2 show moderate SHG effects with the largest tensors being 1.64 and 1.80 pm V<sup>-1</sup>, respectively. In addition, TlBrO3 exhibits the highest SHG effect among the bromates ( $d_{11} = 3.19 \text{ pm V}^{-1}$ ), which originates from its wellaligned motifs and smaller band gap. However, the SHG effect of BiO(BrO<sub>3</sub>) with its very narrow band gap (3.89 eV &  $d_{15}$  = 2.35 pm V<sup>-1</sup>) is weaker than those of the alkali metal bromates, which may be because the polarity of BrO<sub>3</sub> is partially cancelled out in the compound.

The calculated refractive index results (Fig. S2<sup>†</sup> and Table 1) indicate that  $AClO_3$  and  $ABrO_3$  (A = NH<sub>4</sub>, K, Rb, Cs) in R3m possess very strong optical anisotropy and large birefringence (0.159-0.186 for  $AClO_3$  and 0.203-0.227 for  $ABrO_3$ ), enabling them to achieve phase-matching (PM) near their UV cutoff edges ( $\lambda_{\text{cutoff}}$ ; the deep-UV region for AClO<sub>3</sub> and the short-wave UV region for ABrO<sub>3</sub>). In particular, ABrO<sub>3</sub> crystals (A = K, Rb and Cs) are even full-wavelength phase-matchable (FWPM) owing to their sufficiently large birefringence and small dispersion extent of their refractive index curves. In comparison, the alkaline earth metal halates in Fdd2 and Cc show relatively small birefringence (0.062-0.091@1064 nm) and thus cannot achieve PM SHG output near their  $\lambda_{\text{cutoff}}$ . For example, although the  $\lambda_{\text{cutoff}}$  values of  $Sr(ClO_3)_2$  and  $Ba(ClO_3)_2$  are as short as 176.4 and 177.9 nm in the deep-UV region, their

shortest PM wavelengths ( $\lambda_{PM}$ ) are as long as 260.5 and 255.7 nm, implying that they can only output PM SHG light from 532 to 266 nm; so do Sr(BrO<sub>3</sub>)<sub>2</sub> and Ba(BrO<sub>3</sub>)<sub>2</sub>, which can only output PM SHG light from 1064 to 532 nm, despite their much shorter  $\lambda_{\text{cutoff}}$  values of 202.0 and 200.6 nm. It is worth noting that BiO(BrO<sub>3</sub>) cannot realize phase-matching in the  $1064 \rightarrow 532$  nm process due to its very small birefringence of 0.040@1064 nm. Interestingly, the birefringence  $\Delta n$ , the UV cutoff edges \(\lambda\_{\text{cutoff}}\) and the shortest phase-matching wavelengths  $\lambda_{PM}$  of the halates are summarized in Fig. 3b, from which we can find the intrinsic relationship among the three optical indicators, *i.e.*, the larger the  $\Delta n$ , the closer the  $\lambda_{PM}$  is to  $\lambda_{\text{cutoff}}$ , even equalling  $\lambda_{\text{cutoff}}$ , namely the full-wavelength phase-matching (FWPM).11

Summarizing the above results, it is found that AClO<sub>3</sub> (A = NH<sub>4</sub>, K, Rb, Cs) can be promising DUV NLO crystals with strong SHG responses; meanwhile,  $M(ClO_3)_2$  (M = Sr, Ba) and ABrO<sub>3</sub> (A = K, Rb, Cs), particularly the latter with higher SHG effects and FWPM capability, can be important short-wave UV NLO materials capable of producing a 266 nm laser by SHG or quadruple frequency generation techniques. In addition, TlClO<sub>3</sub>, Pb(ClO<sub>3</sub>)<sub>2</sub>, TlBrO<sub>3</sub>, Sr(BrO<sub>3</sub>)<sub>2</sub> and Ba(BrO<sub>3</sub>)<sub>2</sub> can also be potential SHG materials for application in the  $1064 \rightarrow 532$  nm range of green light output.

#### Structure-property relationship

For the screened NLO halates, the typical KClO<sub>3</sub>, TlClO<sub>3</sub> and Pb(ClO<sub>3</sub>)<sub>2</sub> are taken as representatives to study their electronic structures and analyze the origin of their strong SHG effects and large birefringence. The effects of Tl<sup>+</sup> and Pb<sup>2+</sup> on the optical properties are also to be discussed.

From the partial density of states shown in Fig. 4a-c, it can be seen that there are fully overlapping electronic states between Cl and O in the VB range of -12--5 eV and in the CB range of 5-10 eV, indicating the strong covalent interactions of Cl-O bonds in the halates. K-3p, Tl-5d and Pb-5d are highly localized in the deep energy levels (<-8 eV), so they would have little effect on the optical properties. For KClO<sub>3</sub>, the upper part of VB is dominated by the O-2p nonbonding states, mixed with a few Cl-3s3p nonbonding states, and the lower part of CB consists of the unoccupied O-2p and Cl-3p orbitals, so the wide band gap of KClO3 is determined by the ClO3 groups. The electronic structures of TlClO<sub>3</sub> and Pb(ClO<sub>3</sub>)<sub>2</sub> are more complicated: except for the usual O-2p and Cl-3s3p, Tl-6s6p and Pb-6s6p are also involved in the VB top and CB bottom, so their band gaps are determined by ClO<sub>3</sub> together with the Tl<sup>+</sup>/Pb<sup>2+</sup> cations. Considering the much narrower band gaps of TlClO<sub>3</sub> and Pb(ClO<sub>3</sub>)<sub>2</sub> compared to that of KClO<sub>3</sub>, it can be confirmed that Tl<sup>+</sup>/Pb<sup>2+</sup> has a weakening effect on the band gaps of the compounds. To further investigate the stereoactivity of the lone pair electrons (6s<sup>2</sup>) on the Tl<sup>+</sup> and Pb<sup>2+</sup> cations, the electron localization functions (ELFs) of TlClO<sub>3</sub> and Pb(ClO<sub>3</sub>)<sub>2</sub> are calculated and plotted with  $\eta = 0.9$ (Fig. S3†). It clearly shows that there are lobe-like isosurfaces on the Cl5+ ions in both compounds, confirming the highly stereoactive Cl<sup>5+</sup>-3s<sup>2</sup> lone pair, whereas the isosurfaces around

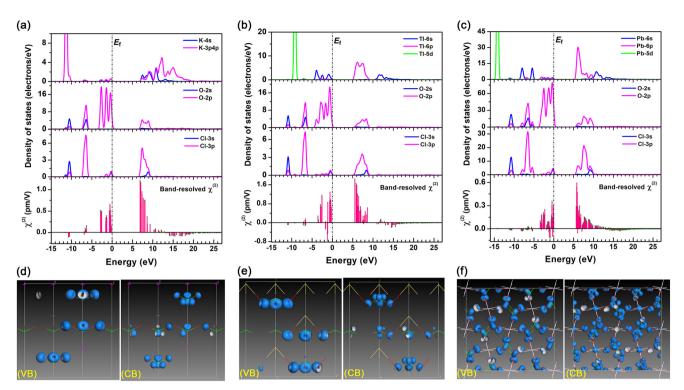


Fig. 4 Partial density of states and band-resolved NLO susceptibility  $\chi^{(2)}$  (a-c) and SHG-density plots (d-f) for KClO<sub>3</sub>, TlClO<sub>3</sub> and Pb(ClO<sub>3</sub>)<sub>2</sub>, respectively.

TI<sup>+</sup> and Pb<sup>2+</sup> (especially Pb<sup>2+</sup>) are nearly spherical, indicating that the Tl<sup>+</sup>-6s<sup>2</sup> and Pb<sup>2+</sup>-6s<sup>2</sup> lone pair electrons are much less stereoactive and almost inert.

Band-resolved  $\chi^{(2)}$  and SHG-density techniques are used to explore the origin of the strong SHG effects for these simple halates. The former can help us to know which energy levels give large contributions to SHG, and the latter can intuitively show the specific SHG-contributing orbitals and atoms in real space. From the band-resolved  $\chi^{(2)}$  in the bottom panels of the DOS graphs (Fig. 4a-c), we can see that the states distributed on both sides of their forbidden bands (-4-0 eV and 5-10 eV) give the majority of the SHG contributions, corresponding to O-2p nonbonding states and a small amount of Cl-3s3p, Tl-6s and Pb-6s6p states in VB and the empty orbitals of Cl-3p, O-2p, K-4s4p, Tl-6p and Pb-6p in CB. The SHG-density graphs (Fig. 4d-f) for the three compounds indicate that not all of the above electronic states are involved in contributing to the SHG process. In fact, only the nonbonding states of O-2p in VB and the unoccupied Cl-3p, O-2p and a few Pb-6p orbitals in CB make real contributions to the strong NLO responses of KClO<sub>3</sub>, TlClO<sub>3</sub> and Pb(ClO<sub>3</sub>)<sub>2</sub>. Furthermore, based on the SHG-density data, the contribution percentages of all groups and ions can be obtained (Table 2); apart from the core ClO<sub>3</sub> groups, which play an overwhelming role in the SHG effects, TlO3 and PbO8 groups also make some contributions, 29.37% and 26.17% respectively, while the contributions from K<sup>+</sup> are negligible.

In addition, we use the polarizability anisotropy-weighted electron density (PAWED) method to investigate the structural source of the large birefringence of these compounds

Table 2 Calculated contribution percentages of different groups to SHG coefficients and the birefringence for KClO<sub>3</sub>, TlClO<sub>3</sub> and Pb(ClO<sub>3</sub>)<sub>2</sub>

Material	Group	$d_{ij}$ (pm V <sup>-1</sup> )	$\Delta n$
KClO <sub>3</sub>		$d_{11} = 1.85$	0.186
3	$ClO_3$	98.46%	96.78%
	$\mathbf{K}^{+}$	1.53%	3.18%
TlClO <sub>3</sub>		$d_{11} = 2.37$	0.173
-	$ClO_3$	70.63%	74.00%
	$TlO_3$	29.37%	25.96%
$Pb(ClO_3)_2$		$d_{24} = 2.37$	0.102
, ,	$ClO_3$	72.31%	73.22%
	$PbO_8$	26.17%	25.40%

(Fig. S4†). In VB, the birefringence originates from the O-2p nonbonding states, and in CB, it is mainly from the Cl-3p and O-2p, mixed with some Pb-6p empty orbitals. The calculated contribution values (Table 2) are similar to those obtained for SHG effects: ClO<sub>3</sub> groups also play a predominant role in generating large birefringence, while the contributions from TlO3 and PbO<sub>8</sub> groups are not very large, which can be attributed to the almost inert stereoactivity of the 6s<sup>2</sup> lone pairs on the Tl<sup>+</sup> and Pb2+ cations.

## Conclusions

In summary, the long-neglected stereochemically active lone pair (SCALP) groups ClO<sub>3</sub> and BrO<sub>3</sub> are innovatively proposed to be good UV NLO functional motifs based on the group property prediction, and the NLO performance of chlorates and bromates is firstly investigated systematically by employing the density functional theory method within the CASTEP code. Through careful compound screening, structural optimization, property simulation and mechanism analysis, the halate crystals AClO<sub>3</sub> and ABrO<sub>3</sub> (A = NH<sub>4</sub>, K, Rb and Cs) with wellarranged SCALP motifs are highlighted as promising UV NLO materials in view of their excellent properties, including wide UV transparent ranges deep into DUV for chlorates and SWUV for bromates, high SHG effects comparable to β-BaB<sub>2</sub>O<sub>4</sub>, large birefringence and strong phase-matching ability. Concretely, AClO<sub>3</sub> crystals can be potential DUV NLO candidates (4.2-4.8 × KDP &  $\lambda_{PM}$  = 185–195 nm), and ABrO<sub>3</sub> crystals, showing very strong SHG responses (6.1-7.1 × KDP) and rare full-wavelength phase-matching capability ( $\lambda_{PM} = \lambda_{cutoff} = 210-220$  nm), can be important short-wave UV NLO materials. Hence the UV NLO potential of ClO<sub>3</sub> and BrO<sub>3</sub> is fully demonstrated, and this work provides novel functional motifs and opens up a new avenue for the development of SWUV and even DUV NLO materials.

## Author contributions

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Chun-Li Hu performed the theoretical calculations, data analyses, and paper writing. Qian-Qian Chen and Fang Kong performed crystal structure screening. Jiang-Gao Mao provided data analyses and major revisions of the manuscript. All the authors discussed the results and commented on the manuscript.

## Conflicts of interest

There are no conflicts to declare.

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