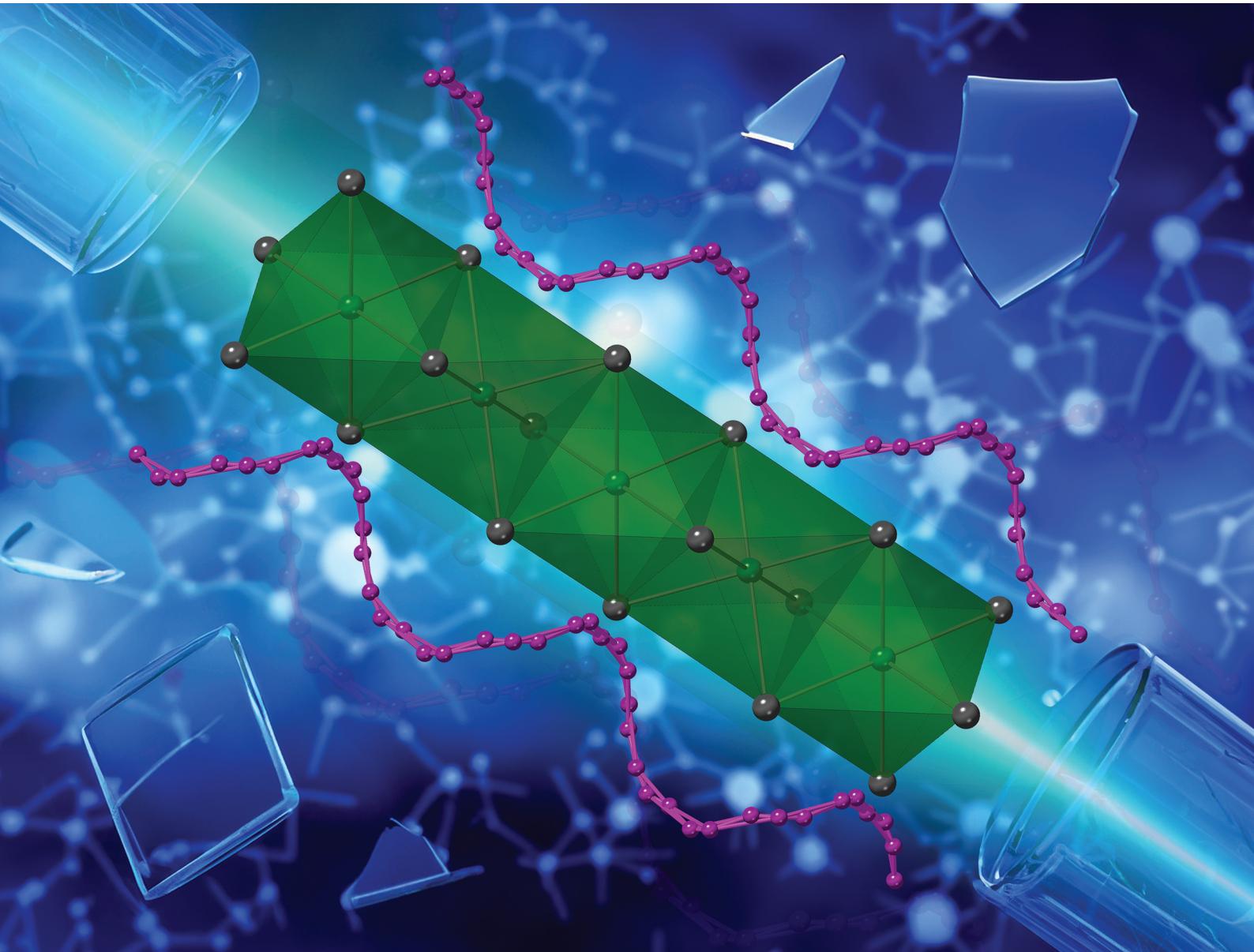


Nanoscale

rsc.li/nanoscale



ISSN 2040-3372

PAPER

Sviatoslav Baranets *et al.*
Novel ternary Zintl phosphide halides $\text{Ba}_3\text{P}_5\text{X}$ ($\text{X} = \text{Cl}, \text{Br}$)
with 1D helical phosphorus chains: synthesis, crystal and
electronic structure



Cite this: *Nanoscale*, 2024, **16**, 7916

Novel ternary Zintl phosphide halides $\text{Ba}_3\text{P}_5\text{X}$ ($\text{X} = \text{Cl}, \text{Br}$) with 1D helical phosphorus chains: synthesis, crystal and electronic structure[†]

Mohd Ishtiyak, S. M. Gayomi K. Samarakoon, Thimira Kandabadage Don, Spencer R. Watts and Sviatoslav Baranets *

Black single crystals of two novel ternary phosphide halides, $\text{Ba}_3\text{P}_5\text{Cl}$ and $\text{Ba}_3\text{P}_5\text{Br}$, were grown using molten metal Pb-flux high-temperature reactions. These compounds were structurally characterized with the aid of the single-crystal X-ray diffraction (SCXRD) method at 100(2) K. The SCXRD shows that both compounds are isostructural and adopt a new structure type (space group $\text{R}\bar{3}c$, No. 167, $Z = 6$) with unit cell parameters $a = 14.9481(16)$ Å, $c = 7.3954(11)$ Å and $a = 15.045(4)$ Å, $c = 7.537(3)$ Å for $\text{Ba}_3\text{P}_5\text{Cl}$ and $\text{Ba}_3\text{P}_5\text{Br}$, respectively. Cl^- and Br^- anions are octahedrally coordinated by Ba^{2+} cations, thus composing a face-sharing 1D infinite chain ${}_{\infty}^1[\text{XBa}_3]^{5+}$ running along the [001] direction. Moreover, the crystal structures feature peculiar one-dimensional disordered infinite helical chains of ${}_{\infty}^1\text{P}^-$, composed of partially occupied phosphorous atoms, each being a superposition of three symmetrical copies of the ordered phosphorus chain, with continuity along the c -axis. $\text{Ba}_3\text{P}_5\text{X}$ ($\text{X} = \text{Cl}, \text{Br}$) compounds are charge-balanced heteroanionic Zintl phases according to the charge-partitioning scheme $(\text{Ba}^{2+})_5[\text{P}^-]_5\text{X}^-$. The presumed semiconducting behavior of both compounds corroborates well with the results of the electronic structure calculations performed with the aid of the TB-LMTO-ASA code.

Received 20th December 2023,
Accepted 6th March 2024

DOI: 10.1039/d3nr06492a
rsc.li/nanoscale

Department of Chemistry, Louisiana State University, Baton Rouge, Louisiana, 70803, USA. E-mail: sbaranets@lsu.edu

† Electronic supplementary information (ESI) available: Ba–P–X compositional diagrams, EDS and XPS spectra, PXRD patterns, structural and computational details of the electronic structure calculations. CCDC 2313247–2313248. For ESI and crystallographic data in CIF or other electronic format see DOI: <https://doi.org/10.1039/d3nr06492a>



Sviatoslav Baranets

Dr. Sviatoslav Baranets is an Assistant Professor of Chemistry at Louisiana State University. He obtained his Ph.D. in Inorganic Chemistry from the Institute of General and Inorganic Chemistry of the National Academy of Sciences of Ukraine in 2015. Later, he was a DAAD scholarship holder and worked as a visiting researcher at the Technische Universität Dresden, Germany. He was also a postdoctoral researcher at the University

of Delaware. His research area primarily focuses on the design, synthesis, and characterization of emerging heteroanionic inorganic functional materials for energy applications.

Introduction

Heteroanionic quantum materials have recently received increased attention due to unusual structural characteristics and exceptional properties.^{1,2} These compositionally and structurally complex phases that are marked by an anionic sublattice comprising of multiple anion types remain relatively underdeveloped compared to the families of homoanionic compounds.¹ However, they have already demonstrated a broad range of physical properties and have recently emerged as promising platforms for nanotechnology,^{3,4} energy-related applications,^{5–11} catalysis,^{12,13} and developing topological quantum materials.¹⁴

A particularly intriguing subset of heteroanionic compounds is composed of pnictogens (*i.e.*, electronegative elements of group 15, Pn = P, As, Sb, Bi) combined with halogens (*i.e.*, nonmetallic elements of group 17, X = F, Cl, Br, I). These heteroanionic ternary and multinary solid-state phases contain negatively charged pnictogen and halogen anions and are termed pnictide halides (*e.g.* $[\text{Ba}^{2+}]_2[\text{P}_7]^{3-}[\text{Cl}^-]$). They are distinct from perovskite-like pnictohalides materials, where pnictogens exhibit positive oxidation states (*e.g.* $[\text{Ag}^+][\text{Bi}^{3+}]_2[\text{I}^-]$).¹⁵ These compounds are known for their rich structural chemistry due to their potential to stabilize various types of homoatomic finite and infinite pnictogen-based units

or metal cluster compounds.^{16–20} Recently discovered pnictide halides and pnictohalides, such as A_2As_7X ($A = Ba, Sr, Eu$; $X = Cl, Br, I$) or W_6PCl_{17} , are considered unique low-dimensional semiconductors with promising tunability that may lead to the discovery of promising electronic and spintronic materials.^{20,21}

Our research group is interested in advancing novel hetero-anionic compounds that are characterized by distinct charge transfer, aiming to develop narrow-gap semiconducting materials suitable for energy applications. Within this domain, polar intermetallic Zintl phases are known for their abundant chemical compositions, variability of crystal structures, and wide range of physical properties.²² The bonding in these charge-balanced compounds is a blend of ionic and covalent interactions. Typically, the electropositive elements from groups 1 and 2 donate electrons to the more electronegative elements, which then use these electrons to form covalent bonds and create complex anionic substructures.²³

Heteroanionic pnictide halides, such as A_2Pn_7X ,^{16,20,24} A_5Pn_3X ,²⁵ or A_2PnX ($A = Ca, Sr, Ba, Eu$; $Pn = P–Bi$; $X = Cl–I$)^{26–28} exemplify charge-balanced double Zintl compounds, showing a range of band gaps and an excellent tunability of crystal structures and electronic properties. Several A_2Pn_7X phases crystallize in the chiral space group $P2_13$ – a desirable prerequisite for the materials requiring chiral control and enantioselectivity.²⁰ As these phosphide and arsenide halides do not have inversion symmetry, they can be used for application in second harmonic generation and ferroelectricity.²⁹

Among pnictogens, phosphorous exhibits remarkable homoatomic bonding formation capability, which results in the constant discovery of novel phosphorus allotropes, like recently reported P_{11} ,³⁰ structurally diverse polyanions, such as infinite $^{1\infty}P^-$ chains in BaP_2 ,³¹ oligomeric $[P_9]$ units in Ba_5P_9 , or $0D$ $[P_7]^{3-}$ clusters in A_2Pn_7X .^{16,32} Being motivated by the rich structural chemistry of Zintl phosphides and developing frontiers for the synthesis of heteroanionic materials that crystallize in the non-centrosymmetric and chiral space groups, we discovered two novel isostructural phosphide halides, Ba_3P_5Cl and Ba_3P_5Br . These compounds adopt a novel structure type featuring a peculiar infinite helical phosphorus chain. We report their synthesis, structural characterization, and chemical bonding study.

Results and discussion

Synthesis

The $Ba–P–X$ ($X = Cl, Br$) ternary compositional diagrams are relatively underexplored, with a limited number of compounds currently known. These include $Ba_2P–X$ of Ba_2P_7Cl structure type,¹⁶ Ba_2PCl of $\alpha\text{-NaFeO}_2$ structure type,²⁸ and the solid solution $Ba_{16}P_{10.14}Br_{1.59}$ of Th_3P_4 structure,²⁸ which is rather believed to be a defect binary $Ba_4P_{2.67}$ phase (Fig. S1†).³¹

The initial discovery of Ba_3P_5Cl was made during an experiment targeting a quaternary $Ba–In–Cl–P$ composition. This process, mirroring the synthetic procedure described in the Experimental section, included $InCl_3$ in the reaction mixture

instead of $BaCl_2$. However, the resulting product contained no known In-bearing compounds, yielding only crystals of the new Ba_3P_5Cl phase. EDS studies also confirmed the absence of foreign elements in the obtained crystals. Subsequently, we attempted the synthesis of the title phase without $InCl_3$ to elucidate its role. Initially, it was hypothesized that the partial decomposition of $InCl_3$ at high temperatures might provide a source of chlorine to the reaction mixture. However, our findings revealed that indium halides are not essential for synthesizing Ba_3P_5Cl or Ba_3P_5Br . We found that the presence of anhydrous barium halides alone is sufficient, serving as an equally excellent source of halogens and yielding nearly high-purity samples of the title phases.

We also attempted to explore periodic trends and synthesize the isostructural Ba_3P_5I phase, but several synthetic attempts were unsuccessful, and the crystals of Ba_2P_7I were mainly found as a major phase.¹⁶ Experimental efforts to synthesize this hitherto unknown iodide using InI_3 also yielded the same ternary 2–7–1 phase.

The single crystals of Ba_3P_5Cl and Ba_3P_5Br synthesized in our study exhibit extreme sensitivity to air and moisture, posing significant challenges in the study of their properties. These crystals rapidly decompose upon brief exposure to ambient conditions and undergo slow deterioration even with the presence of trace amounts of moisture.

Crystal structure description of Ba_3P_5X ($X = Cl, Br$)

The detailed crystallographic analysis of the Ba_3P_5X phases uncovers a great degree of complexity, as depicted in Fig. 1. Both phosphide halides are isostructural and crystallize in the rhombohedral centrosymmetric space group $R\bar{3}c$ (No. 167) within the trigonal crystal system with six formula units per

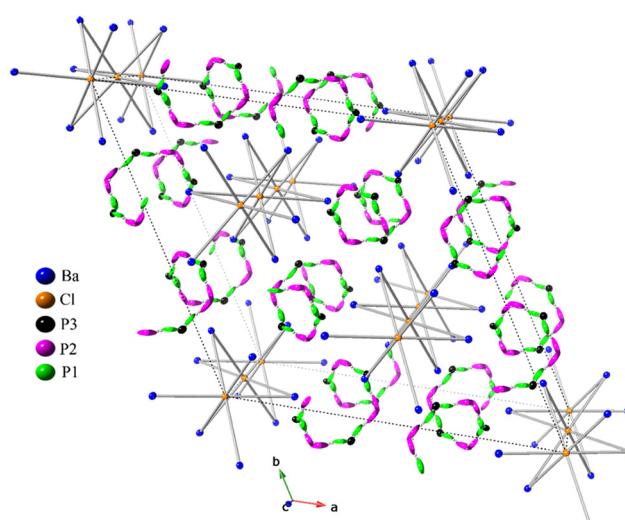


Fig. 1 A unit cell view of the crystal structure of Ba_3P_5X ($X = Cl$ and Br) approximately along the $[001]$ direction. Barium atoms are blue, halogen atoms are orange, whereas three different phosphorus atoms are depicted as green, pink, and black. Ellipsoids for phosphorus atoms are plotted with the 50% probability level. The unit cell is dashed.

unit cell. The crystalline arrangement of both compounds, being unprecedented in existing literature, represents a novel structure type. As can be observed from Table 1, the unit cell parameters for the bromide analog are substantially larger, aligning with the larger atomic radius of bromine compared to chlorine.

The asymmetric unit of $\text{Ba}_3\text{P}_5\text{X}$ ($\text{X} = \text{Cl}, \text{Br}$) comprises a total of 5 crystallographically independent sites with respective site symmetries: $\text{Ba}1$ (2), $\text{X}1$ (3), $\text{P}1$ (1), $\text{P}2$ (1), and $\text{P}3$ (2). The atomic arrangement in these 3–5–1 phases demonstrates that the structures of both compounds can be considered as pseudo-one-dimensional (Fig. 1). The halogen atoms are octahedrally coordinated by six equidistant Ba atoms, with $\text{Ba}-\text{X}$ distances of *ca.* 3.32 Å and 3.41 Å for the chloride and the bromide phases, respectively. These values are slightly longer than the sum of Ba^{2+} and X^- ionic radii³³ or those observed in the $\text{Ba}_2\text{P}_5\text{Cl}$ phase,¹⁶ indicating a high degree of ionicity for these contacts, *vide infra*. $[\text{XBa}_6]$ octahedra are interconnected by sharing two opposite faces, forming a 1D chain $^1\infty[\text{XBa}_3]^{5+}$ along the [001] direction. Compared to the edge- and vertex-sharing polyhedral arrangements, the face-sharing of polyhedral units is often regarded as the least stable arrangement,³⁴ primarily because of the closest positioning of two ions (X^-) to each other.

However, the most interesting component of the reported crystal structures is the arrangement of phosphorus atoms. Their connectivity forms an anionic substructure identified as 1D infinite disordered chains also oriented along [001] direction. These anionic $^1\infty\text{P}^-$ helical chains are a hallmark of the $\text{Ba}_3\text{P}_5\text{X}$ crystal structure and exemplify the structural abundance of phosphorus Zintl chemistry.

To better understand the observed structural disorder in these anionic helical chains, we provide the reader with the details of the structure determination flow for $\text{Ba}_3\text{P}_5\text{Cl}$. The initial structural model, derived from the first iteration of the

structure solution, indicated four crystallographically independent sites. Based on the geometrical parameters and coordination environments, these atomic positions were assigned as one Ba, one Cl, and two P atoms. In this model, each Ba atom can be viewed as 6-coordinated by two chlorine atoms and four P–P pairs. However, we observed immediate problems with this disorder-free structure with the refined composition of $\text{Ba}_3\text{P}_6\text{Cl}$. Among them were (i) anomalously large atomic displacement parameters for both P atoms, (ii) a residual electron density peak of $4.5 \text{ e}^- \text{ \AA}^{-3}$ and a hole of $-4.9 \text{ e}^- \text{ \AA}^{-3}$ near one phosphorus atom, and (iii) a relatively short interatomic P–P distance of *ca.* 1.84 Å. These observations clearly pointed to the disorder of the P atoms. Allowing free occupancies for these atoms indicated a clear underoccupancy of one phosphorus atom (P1) and a refined formula approaching the $\text{Ba}_3\text{P}_5\text{Cl}$ composition. However, the second phosphorus atom (P2) still exhibited large ADP alongside a significant residual electron density peak, located *ca.* 0.9 Å away. By assigning this peak as another partially occupied P atom (P3) and splitting the P1 atom, we achieved a nearly featureless Fourier difference map with reasonable ADP values without introducing additional constraints. This structural model was successfully applied to the bromide analog, yielding an identical structural solution with heavy disorder on phosphorus sites.

Refined occupancies of phosphorus atomic sites for both chloride and bromide phases were found to be close to 1/3, yielding nearly stoichiometric refined compositions of $\text{Ba}_3\text{P}_{4.8(3)}\text{Cl}$ and $\text{Ba}_3\text{P}_{4.9(5)}\text{Br}$, respectively. Given the large calculated ESD values, we considered the “off-stoichiometry” as an artifact arising from differences in U_{eq} values for all three phosphorus atomic sites. Constraining their occupancies to 1/3 did not affect the quality of the structural refinement and resulted in minor changes in ADP values. On the contrary, such a structural model represents a stoichiometric $\text{Ba}_3\text{P}_5\text{X}$ ($\text{X} = \text{Cl}, \text{Br}$) composition, which corroborates with the idea of charge balance, as discussed further below.

The rationalization of the disorder within the $^1\infty\text{P}$ chain aligns well with the model containing constrained occupancies. The observed minimal P–P interatomic distances are too short to be classified as covalent bonds, so the simultaneous existence of all phosphorus atomic sites can only be achieved through their partial SOFs. At first glance, the chain looks like a continuous flow of phosphorus atoms within the helix with almost no breaches, leading to nearly uniformly distributed or “smeared” electron density, as can be seen from the alignment of refined ellipsoids in Fig. 2a. However, the partial occupancies of phosphorus atoms allows for various structural scenarios with specific atomic sequence. Given the covalent radius of 1.07 Å for P,³⁵ one should expect an interatomic P–P distance of *ca.* 2.15 Å. In the presented structural model, we observe a wide range of similar distances for P1–P1 (2.20–2.22 Å), P1–P2 (2.18–2.22 Å), and P2–P3 (2.15–2.19 Å) contacts so that each disordered helix can be viewed as a superposition of three symmetry copies of the same ordered chain sequenced to the {P1–P1–P2–P3–P2} fragment of partially occupied phosphorus atoms (Fig. S6†), each with SOF of

Table 1 Selected crystallographic data and structure refinement details for $\text{Ba}_3\text{P}_5\text{Cl}$ and $\text{Ba}_3\text{P}_5\text{Br}$ ($T = 100(2)$ K, Mo $\text{K}\alpha$, $\lambda = 0.71073 \text{ \AA}$)

Refined formula	$\text{Ba}_3\text{P}_5\text{Cl}$	$\text{Ba}_3\text{P}_5\text{Br}$
$fw/\text{g mol}^{-1}$	602.32	646.78
Space group	$R\bar{3}c$	
$a/(\text{\AA})$	14.9481(16)	15.045(4)
$c/(\text{\AA})$	7.3954(11)	7.537(3)
$V/(\text{\AA}^3)$	1431.1(4)	1477.6(10)
Z	6	
$\rho_{\text{cal}}/\text{g cm}^{-3}$	4.193	4.361
$\mu(\text{Mo K}\alpha)/\text{cm}^{-1}$	132.8	166.5
Collected/independent reflections	11 569/493	5390/471
$R_1 (I > 2\sigma_{(I)})^a$	0.025	0.048
$wR_2 (I > 2\sigma_{(I)})^a$	0.040	0.106
R_1 (all data) ^a	0.030	0.066
wR_2 (all data) ^a	0.041	0.111
$\Delta\rho_{\text{max,min}}/\text{e}^{-\text{\AA}^{-3}}$	0.83, -2.09	1.82, -2.45
CCDC	2313247	2313248

^a $R_1 = \sum ||F_o| - |F_c|| / \sum |F_o|$. $wR_2 = \{\sum [w(F_o^2 - F_c^2)^2] / \sum wF_o^4\}^{1/2}$, $w = 1 / [\sigma^2(F_o^2) + (AP)^2 + (BP)]$, where $P = (F_o^2 + 2F_c^2)/3$; A and B are weight coefficients.

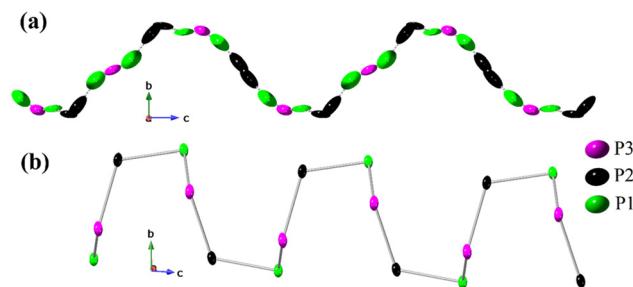


Fig. 2 Propagation of 1D $1\infty P1^-$ helical chains found in (a) the crystal structures of Ba_3P_5X and (b) AP_2 ($A = Sr, Ba$). The P–P bonding interactions have been represented with dotted lines in Ba_3P_5X , while solid-bond has been used in the case of AP_2 . Ellipsoids are plotted with a 50% probability level.

1/3. The random shift of adjacent chains in the trigonal *ab* plane with respect to each other justifies the observed disorder. Lastly, the described sequence of five phosphorus atoms exemplifies one complete turn of the helix, with a pitch equal to the *c* lattice parameter.

Described charged helical phosphorus chains are stabilized by the $1\infty[ClBa_3]^{5+}$ matrix, yet their fundamental existence arises from the ability of pnictogen elements to form homoatomic Pn–Pn contacts (Pn = P, As, Sb, Bi) and to construct electron-rich polyanionic networks.³⁶ Such structural fragments are commonly observed in the realm of Zintl pnictides or heteroanionic Zintl-like salts.^{23,37} In addition, phosphorus is known for forming a variety of allotropes of different dimensionalities,^{30,38–42} including some exotic predicted forms like helical coil phosphorus allotropes.⁴³

The disordered helical phosphorus chains described in this work (Fig. 2a) have never been previously reported, although disorder-free phosphorus helices are known for binary phosphides AP_2 ($A = Sr, Ba$) (Fig. 2b)³¹ and ternary heteroanionic compounds like $SnPX$ (Fig. S6†)^{44,45} and Cd_2P_3X ($X = Cl, Br$, and I).⁴⁶ The crystal structure of the $SnPI$ compound reported by Pfister *et al.* in 2016⁴⁴ can be described as a double helical $SnIP$ strand composed of inner $1\infty[P^-]$ and outer $1\infty[SnI]^+$ chains oriented along the [100] direction. Reported $1\infty P$ chains in $SnIP$ (Fig. S6†) are structurally similar, yet not identical to those presented in this work. Bond angles within the {P1–P1–P2–P3–P2} fragment in Ba_3P_5Cl range from 127.2° to 130.2° , whereas observed angles within the phosphorus chain in $SnPI$ are much smaller, being within 95.5 – 96.7° range. The zigzag phosphorus chains were also found running along the *c*-axis in the structures of valence-precise compound Cd_2P_3X ($X = Cl, Br$, and I) and monoclinic- NiP_2 .⁴⁷ 1D helical chains of $1\infty P^-$ were also reported in α - CdP_2 and β - CdP_2 phases.⁴⁸ These compounds exhibit P–P interatomic distances within the range of 2.10 to 2.35 Å, similar to those in the title Ba_3P_5X phases close to the sum of their covalent radii.³⁵ It would be worth noting that one-dimensional heteroanionic semiconducting material, $SnPI$, exhibited anisotropic mechanical properties like steel, which strongly motivate the further exploration of novel helix-containing compounds.⁴⁴

Although heavy disorder significantly complicates the bonding analysis, we observe that the shortest and longest Ba–P bond distances in Ba_3P_5X range from *ca.* 3.13(1) to 3.29(1) Å, comparable with those in BaP_2 .³¹ Similar Ba–P distances are also observed in Ba_2P_7X ($X = Br$ and I).¹⁶ As Table 2 shows, Ba–X bond distances are longer in Ba_3P_5Br compared to Ba_3P_5Cl , which aligns with the larger ionic radius of Br, while Ba–P distances do not show significant variation in Ba_3P_5X .

The described structural model of Ba_3P_5X ($X = Cl, Br$) with constrained occupancies of phosphorus atoms, as detailed in Table 3, leads to the charge-balanced composition, where disordered phosphorus helix fully accounts for a -5 charge necessary for electron balancing the space-filler between the 1D chains of $1\infty[ClBa_3]^{5+}$. A charge-balanced composition can be represented as $(Ba^{2+})_3(P^-)_5X^-$ ($X = Cl, Br$), satisfying the Zintl count and pointing to the potential semiconducting behavior.

XPS studies

To further elucidate the validity of the proposed charge-partitioning and understand the chemical states in Ba_3P_5X , we conducted X-ray Photoelectron Spectroscopy (XPS) analysis on the polycrystalline samples. The results of these studies are presented in the ESI (Fig. S4 and S5).† The survey spectra (Fig. S4a and S5a†) of Ba_3P_5X indicated the presence of peaks corresponding to the assigned elements, corroborating well with the SEM-EDS analysis, *vide supra*. The characteristic binding energy peaks of $Ba 3d_{5/2}$ (780.26 eV for Ba_3P_5Cl and 780.41 eV for Ba_3P_5Br) confirmed that Ba is in the +2 oxidation state, in agreement with previous reports on Ba^{2+} ions.⁴⁹ The high-resolution spectrum of Cl exhibited $2p_{3/2}$ and $2p_{1/2}$ peaks at 198.66 eV and 200.43 eV, respectively, while the Br $3d_{5/2}$ and $3d_{3/2}$ peaks were observed at 69.11 eV and 70.17 eV, corresponding to the binding energies of Cl^- and Br^- , respectively.⁴⁹

The high-resolution P spectra for both compounds reveal two broad asymmetric peaks. The deconvolution of these peaks suggested the presence of binding energies indicative of several phosphorus states, categorized into higher (BE > 132.5 eV), mid-range (BE ~ 130 eV), and lower binding (BE ~ 129 eV) energy regions. Such spectral behavior has already been previously reported for transition metal phosphides, indicating the presence of oxidized, unreacted, and reduced P chemical states.⁵⁰

As previously mentioned, both title compounds were found to be air-sensitive, hence, a certain degree of surface oxidation

Table 2 Selected interatomic distances (Å) in Ba_3P_5Cl and Ba_3P_5Br

Ba_3P_5Cl	Distances (Å)	Ba_3P_5Br	Distances (Å)
Ba1–Cl1	3.3207(4)	Ba1–Br1	3.4110(11)
Ba1…P1	3.229(6)	Ba1…P1	3.254(13)
Ba1…P2	3.161(9)	Ba1…P2	3.176(14)
Ba1…P3	3.236(4)	Ba1…P3	3.257(6)
P1…P1	2.204(12)	P1…P1	2.22(3)
P1…P2	2.178(10)	P1…P2	2.219(19)
P2…P3	2.144(7)	P2…P3	2.194(10)

Table 3 Fractional atomic coordinates and U_{eq} values for $\text{Ba}_3\text{P}_5\text{Cl}$ and $\text{Ba}_3\text{P}_5\text{Br}$

Atoms	Site	<i>x</i>	<i>y</i>	<i>z</i>	U_{eq}^a
$\text{Ba}_3\text{P}_5\text{Cl}$					
Ba1	18e	0.8154(1)	0	$\frac{1}{4}$	0.014(1)
P1	36f	0.0254(5)	0.2674(5)	0.1776(8)	0.027(2)
P2	36f	0.0861(4)	0.3150(7)	0.1047(18)	0.032(2)
P3	18e	0.2383(8)	0	$\frac{1}{4}$	0.021(2)
Cl1	6b	0	0	0	0.012(1)
$\text{Ba}_3\text{P}_5\text{Br}$					
Ba1	18e	0.8110(1)	0	$\frac{1}{4}$	0.015(1)
P1	36f	0.0256(11)	0.2664(12)	0.1746(15)	0.023(3)
P2	36f	0.0857(9)	0.3142(10)	0.105(2)	0.023(3)
P3	18e	0.2362(11)	0	$\frac{1}{4}$	0.016(3)
Br1	6b	0	0	0	0.013(1)

^a U_{eq} is the one-third value of the trace of orthogonalized U_{ij} tensor. Atom site occupancies of P1, P2, and P3 were constrained as 1/3.

was anticipated. This was further evidenced by the presence of an O 1s peak (\sim 531 eV) and the higher binding energy (P 2p) peaks of 132.40 eV and 132.57 eV found in the XPS spectra of $\text{Ba}_3\text{P}_5\text{Cl}$ and $\text{Ba}_3\text{P}_5\text{Br}$, respectively. These bands are characteristic of oxidized species, such as phosphates, although this phenomenon was also observed during the XPS studies of metal phosphides, such as $\text{Ba}_2\text{Si}_3\text{P}_6$ and NiP_2 , indicating similar surface oxidation.^{47,49–51}

The phosphorus peaks observed in the lower binding energy regions of the recorded XPS spectra are particularly noteworthy, as they indicate reduced chemical states. This feature is commonly seen in the XPS spectra of metal phosphides, such as $\text{Ni}^{2+}(\text{P}^{1-})_2$, which typically exhibit characteristic P 2p_{3/2} peaks in this region. For instance, the bands at 128.7 eV and 129.3 eV correspond to P^{1-} in NiP_2 .⁴⁷ Similarly, the XPS spectra of $\text{Ba}_3\text{P}_5\text{Cl}$ and $\text{Ba}_3\text{P}_5\text{Br}$ displayed major P 2p_{3/2} peaks at 128.98 eV and 129.71 eV, respectively. These binding energies suggest that phosphorous predominantly exists in a -1 oxidation state in $\text{Ba}_3\text{P}_5\text{X}$, aligning well with the proposed charge-partitioning scheme. Distinct peaks corresponding to the phosphorous 2p_{3/2} and 2p_{1/2} were not observed, likely due to the high sensitivity of the compounds to air and moisture. Additionally, the use of high pass energy during the XPS measurement and the presence of oxidized species impurities may have contributed to the broadening of the P 2p peaks for both compounds.⁵⁰

Electronic structure and chemical bonding

While the Zintl-Klemm formalism, considerations of electronegativity differences, and experimental studies of oxidation states have provided insights into an empirical understanding of the bonding situation in the title $\text{Ba}_3\text{P}_5\text{X}$ ($\text{X} = \text{Cl}, \text{Br}$) phases, a more comprehensive understanding of the electronic structure and stability of these compounds was gained from our electronic structure calculations. Due to extensive structural disorder in both title phases, calculations using experimentally obtained atomic coordinates were not feasible. Consequently, we considered a structurally similar model that keeps structural features and chemical composition of the title phases.

As we discussed above, the structural disorder likely originates from the different packing of the phosphorus chains in the trigonal ab plane. With the assistance of the ISODISTORT tool,^{52,53} we transformed the trigonal structure with the space group $R\bar{3}c$ to a monoclinic one with the space group $C2/c$, in which the original three phosphorus sites split into eight with lower multiplicities.

We can remove some of the P sites, leaving three that are symmetrically equivalent to those in the starting structure. The resulted structural model (Fig S7a and Tables S1, S2†) retained the helical $_{\infty}\text{P}^-$ chains while maintaining full occupancy of all atoms, refined 3–5–1 composition, and physically meaningful P–P distances. The periodicity of the helixes remained identical to those in the refined structures. We could reduce the symmetry even further by transforming the structure to the non-centrosymmetric space group Cc while yielding a reasonable structural model (Fig. S7b and Tables S1, S4†). The results of the electronic structure calculations for $\text{Ba}_3\text{P}_5\text{Cl}$ and $\text{Ba}_3\text{P}_5\text{Br}$ compositions using the $C2/c$ structural model are presented in Fig. 3 and 4, respectively, whereas the Total/PDOS calculations for the Cc structural model are included and displayed in Fig. S8 of the ESI.† Crystal Orbital Hamilton Population curves are presented in Fig. 5.

The analysis of the electronic structure indicates an almost identical band gap opening of *ca.* 0.73 eV for both compounds, even though it is expected to observe smaller band gaps for compounds with heavier elements within the same family of isostructural phases.⁵⁴ On the contrary, the calculated band gap for the $\text{Ba}_3\text{P}_5\text{Cl}$ model in space group Cc is significantly larger (*ca.* 1.2 eV, Fig. S8†). Here we should note, that none of the discussed ordered models were geometrically optimized, which is believed to be a major contributing factor for the observed inconsistency. However, details of the electronic structure can provide insightful results for the understanding of the bonding of the title compounds, primarily due to the similar crystal structures of $\text{Ba}_3\text{P}_5\text{X}$ in $C2/c$ and $R\bar{3}c$ models.

Our calculations indicate that the electronic structures of the chloride and bromide compounds are nearly identical; therefore, we will focus our discussion primarily on the characterization of the chloride phase. The total (TDOS) and partial (PDOS) density of states diagrams demonstrate that the conduction band is predominantly populated by the Ba-d and P-p orbitals, with minor contributions from chlorine. Conversely, the states in the valence band are contributed mainly by P-p and Cl-p orbitals. Cl-p states are located in a relatively narrow energy window between -0.5 eV and -4.5 eV, with Ba-s, Ba-p, and Ba-d orbitals also contributing within this range. Such orbital mixing may be indicative of Ba–Cl bonding interactions, as this region exhibits domains of bonding character in the computed COHP curves (Fig. 5).

States above -5.7 eV are significantly contributed by P-p orbitals, while in the lower energy regions (not shown in Fig. 3), P-s orbital contributions emerge, indicating a small degree of hybridization in the phosphorus atoms. Empty Ba-d orbitals also contribute between the VBM and -5.7 eV, where the significant contributions from P-p orbitals are observed,

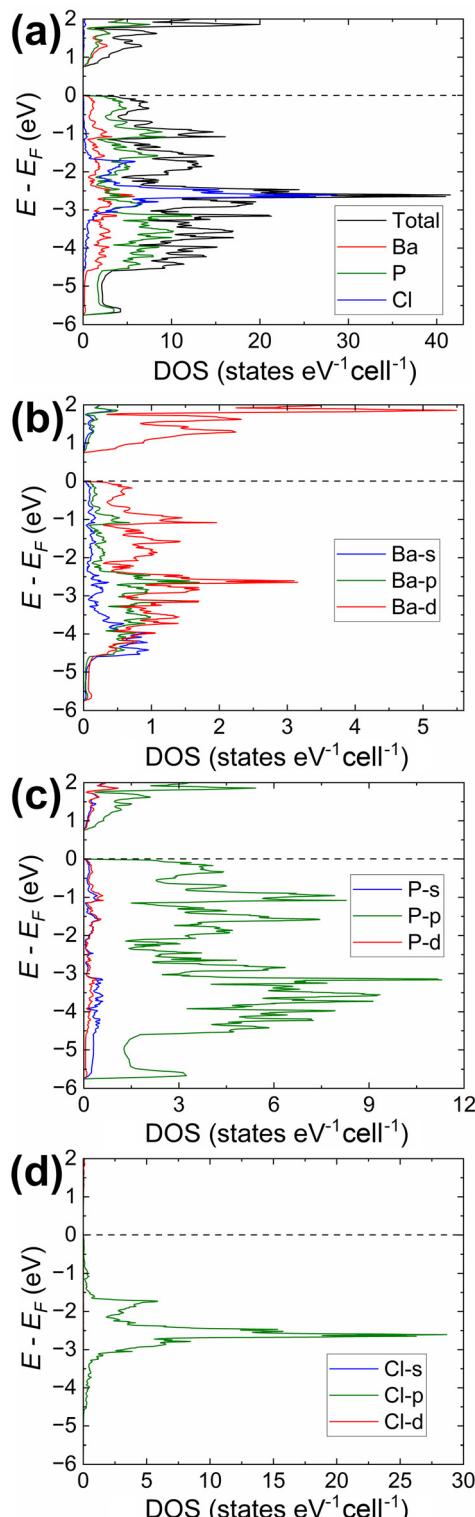


Fig. 3 Electronic structure in $\text{Ba}_3\text{P}_5\text{Cl}$ ($\text{C}2/\text{c}$ structural model; Fig. S7a and Tables S1, S2†): (a) total DOS together with projected densities of states for constituting elements; (b)–(d) partial DOS calculated for the different orbitals of Ba, P, and Cl, respectively.

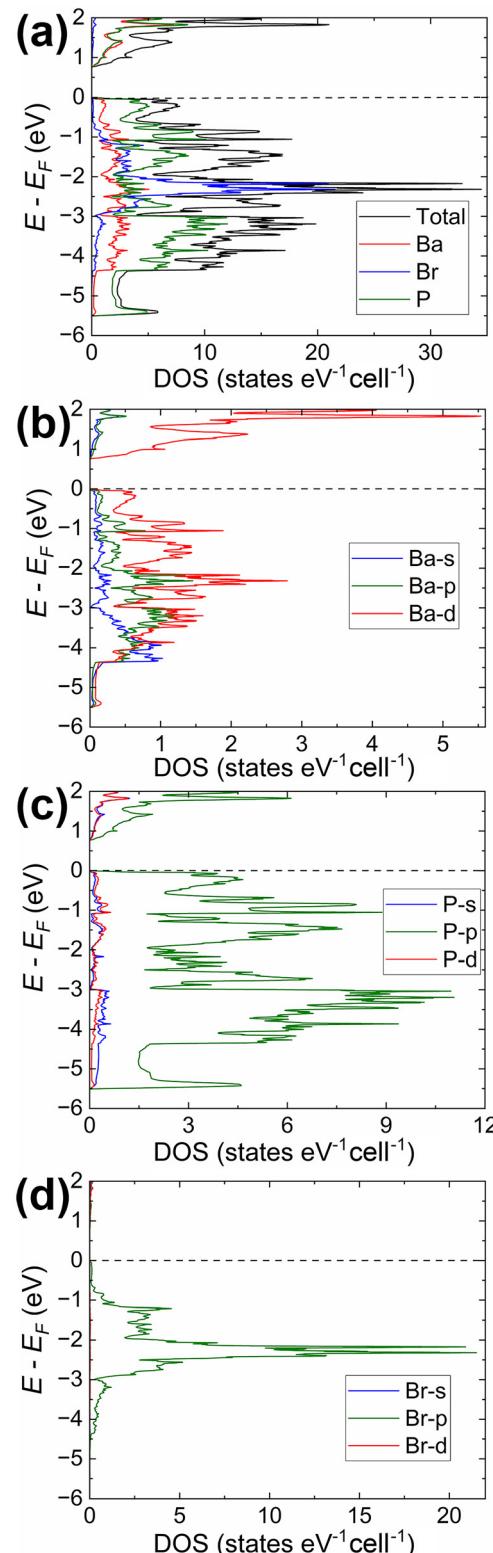


Fig. 4 Electronic structure in $\text{Ba}_3\text{P}_5\text{Br}$ ($\text{C}2/\text{c}$ structural model; Fig. S7a and Tables S1, S3†): (a) total DOS together with projected densities of states for constituting elements; (b)–(d) partial DOS calculated for the different orbitals of Ba, P, and Br, respectively.

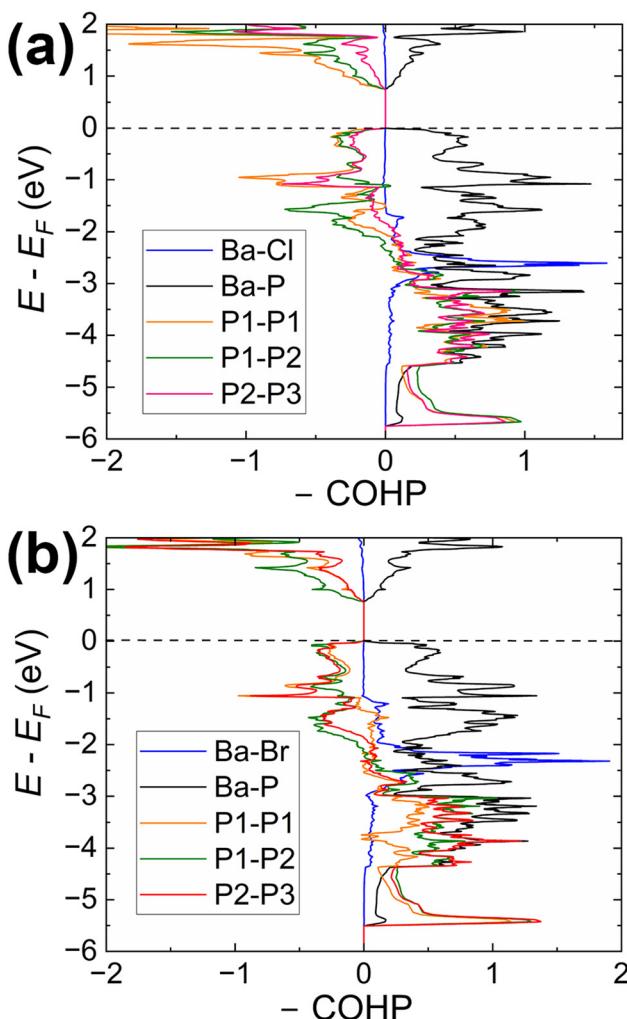


Fig. 5 (a) COHP curves for selected averaged Ba–X and Ba–P interactions and three specific interactions (P1–P1, P1–P2, and P2–P3) within the helical α -P chain in $\text{Ba}_3\text{P}_5\text{Cl}$ (a) and $\text{Ba}_3\text{P}_5\text{Br}$ (b) calculated for the C2/c structural model.

pointing to the orbital mixing between atoms and, therefore, Ba–P interactions. The involvement and mixing of Ba-d orbitals with P-p and Cl-p orbitals clearly illustrates the departure from the fully ionic approximation of the Zintl–Klemm concept, which posits the role of metal cations as electron donors and space-filling entities.

The Crystal Orbital Hamilton Population curves (Fig. 5) provide insights into the bonding situation in the title compounds. COHP analysis indicates optimized Ba–P and Ba–X contacts, while P–P contacts are underoptimized, as evidenced by domains of antibonding character in both the valence and conduction bands. This pattern is characteristic for homoatomic electron-rich pnictogen–pnictogen bonds and was previously observed in several Zintl pnictides, such as $\text{Ba}_2\text{Cd}_2\text{P}_3$ and $\text{Ba}_3\text{Cd}_2\text{P}_4$ with homoatomic P–P contacts.⁵⁵

The average integrated COHP (–ICOHP) values for Ba–X and Ba–P contacts are almost identical, falling within the ranges of

0.20–0.23 eV per bond and 0.28–0.38 eV per bond, respectively. These values indicate a significant charge transfer from electron-donating Ba to electron-accepting P and halogen atoms. –ICOHP values for P–P bonds in $\text{Ba}_3\text{P}_5\text{Cl}$ range between 2.96–3.58 eV per bond, pointing to the large degree of covalency for these interactions.

Experimental

Synthesis

The synthesis of barium phosphide halides, $\text{Ba}_3\text{P}_5\text{Cl}$ and $\text{Ba}_3\text{P}_5\text{Br}$, was accomplished utilizing the following supplied chemicals: Ba rod (99.9%, Thermo Scientific), BaCl_2 powder (99.9%, J.T. Baker), BaBr_2 powder (99.9%, Thermo Scientific), P powder (99.9%, Beantown Chemical) and Pb powder (99.99%, Strem Chemicals). All synthetic and post-synthetic procedures were performed inside an argon-filled glovebox with oxygen and moisture levels below 1 ppm. An oxide layer was removed from the barium rod, whereas other reagents were used as received. Black-colored, irregular-shaped single crystals of the phosphide halide phases can be readily synthesized utilizing a metal-flux reaction in alumina crucibles. $\text{Ba}_3\text{P}_5\text{Cl}$ can be produced in a reaction with stoichiometric amounts of barium chloride (86.4 mg, 0.415 mmol), phosphorous (128.5 mg, 4.149 mmol), barium (285.0 mg, 2.075 mmol), and lead flux (2500 mg, 12.065 mmol). The reaction mixture was loaded into an alumina crucible in the argon-filled glove box. Subsequently, the crucible was placed inside a fused silica tube, which was then sealed under vacuum. Quartz wool was used to cover the crucible, serving as a filter during the centrifugation step. After that, the tube was placed inside a muffle furnace for the heat treatment using the following temperature profile. The furnace temperature was initially ramped up from 373 K to 1273 K at a rate of 100 K h^{–1}, followed by a sustained hold at 1273 K for 20 h. Subsequently, the temperature was gradually decreased to 923 K at a cooling rate of 5 K h^{–1}. Upon reaching the target temperature, the tube was removed from the furnace, and the sample was centrifuged, enabling the separation of the lead flux from the reaction product. Once cooled, the ampoule was transferred to the glovebox and crack-opened. The crystalline reaction product was retrieved from the alumina crucible and stored under the protective atmosphere of the argon-filled glovebox.

After the successful synthesis of single crystals of $\text{Ba}_3\text{P}_5\text{Cl}$, we succeeded in producing the isostructural bromide analog by loading stoichiometric amounts of barium bromide (114.9 mg, 0.529 mmol), phosphorous (119.7 mg, 3.865 mmol), barium (264.4 mg, 1.925 mmol), and the excess of lead (2500 mg, 12.065 mmol) using the same reaction procedure and temperature profile that were used for the synthesis of single crystals of $\text{Ba}_3\text{P}_5\text{Cl}$. It was noted that the crystals of the bromine-containing compound were found to be more stable than the crystals of the chlorine-containing compound. However, neither phase's crystals are sufficiently stable

in the presence of air to allow for the comprehensive analysis of physical properties.

X-ray diffraction methods

Powder X-ray diffraction (PXRD) patterns were recorded in the reflection mode at room temperature using a Rigaku Miniflex 600 diffractometer (Cu K α radiation, $\lambda = 1.5406 \text{ \AA}$, Ge monochromator). The PXRD data were acquired using a θ - θ step-scan mode with 0.05° increments and a 2 s collection period per step. Due to the air sensitivity of the samples, they were covered by a thin amorphous film that prevented immediate oxidation and allowed for the data collection. Collected PXRD patterns indicate the presence of the title phases, yet we cannot assert the phase purity due to the high degree of phase amorphization (Fig. S3 \dagger).

Structural studies were performed with the aid of the single-crystal X-ray diffraction (SCXRD) method. Low-temperature data were collected at 100(2) K on a Bruker APEX II CCD diffractometer equipped with a graphite-monochromized Mo K α radiation ($\lambda = 0.71073 \text{ \AA}$). Suitable-sized single crystals of $\text{Ba}_3\text{P}_5\text{Cl}$ and $\text{Ba}_3\text{P}_5\text{Br}$ were selected, cut to the desired dimensions under the Paratone-N oil, and mounted on a low-background MiTeGen plastic loop attached to the goniometer head. A constant flow of chilled nitrogen gas stream was applied throughout the data collection process to solidify the oil and protect the crystals.

Single crystal data for both compounds were acquired with a scan width of 0.5° and a scan time of 20 s per frame, maintaining a crystal-to-detector distance of 50 mm. The intensity data acquisition, data reduction, cell refinement processes, and absorption correction were conducted using the APEX3 software package.^{56,57} The single-crystal data of both compounds were integrated with a Rhombohedral R setting. The centrosymmetric $R\bar{3}c$ space group was used to solve the crystal structures. The determination of atomic positions was successfully located by the direct method utilizing the SHELXT program, whereas the structure refinement was performed with full-matrix least-squares methods on F^2 with SHELXL.^{58,59} Olex2 program was used as a graphical interface.⁶⁰ Atomic coordinates were standardized using the STRUCTURE TIDY Program.⁶¹ The crystal structures were visualized with the help of CrystalMaker software.⁶²

Selected details of the data collection and relevant crystallographic parameters are given in Tables 1–3. CSD 2313247 and 2313248 contain the complete supplementary crystallographic data for the compounds discussed in this paper. \dagger

Elemental microanalysis

Irregularly shaped black crystals of $\text{Ba}_3\text{P}_5\text{Cl}$ and $\text{Ba}_3\text{P}_5\text{Br}$ were carefully selected and mounted on carbon tape affixed to an aluminum stub. All sample preparation procedures were carried out inside the glovebox due to the extreme air sensitivity of the title compounds.

Energy-dispersive X-ray spectroscopy (EDS) analysis was performed using a Thermo Scientific (TFS Helios G5 PFIB CXe) scanning electron microscope equipped with an OXFORD

Instruments Ultim Max Detector spectrometer. For both sets of crystals, EDS data were acquired at an operating acceleration voltage of 20 kV. Semiquantitative analyses were executed on various crystals at multiple sites to verify the uniformity of elemental compositions throughout the crystals. The EDS studies detected the presence of elements barium, phosphorous, and chlorine/bromine in the crystals, with approximate elemental ratios of $\text{Ba}:\text{P}:X \approx 3:5:1$, as shown in Fig. S2. \dagger These compositions are consistent with the ratios obtained from the SCXRD structure refinements.

X-ray photoelectron spectroscopy (XPS) studies

The XPS measurements were performed to investigate the formal oxidation states of elements in both compounds ($\text{Ba}_3\text{P}_5\text{Cl}$ and $\text{Ba}_3\text{P}_5\text{Br}$). The polycrystalline samples were first ground into fine powder and then pressed into pellets using a hydraulic press. These pellets were placed on the surface of conductive adhesive carbon tape, which was affixed to an aluminum sheet. All these procedures were performed inside the argon-filled glove box to minimize the chances of oxidation, as both compounds are susceptible to air and moisture. The prepared samples were kept inside an argon-filled sealed container for transportation to the XPS instrument and loaded instantly inside the analysis chamber. The pressure in the analysis chamber was maintained between 2.7×10^{-9} and 6×10^{-9} mbar during the measurements. The data were collected on Scienta Omicron ESCA 2SR X-ray Photoelectron Spectroscopy instrument equipped with a high-performance monochromatic Al-K α X-ray radiation source. The operating working voltage and current were set to 15 kV and 20 mA throughout the data measurement. Due to the semiconducting nature of samples, no charge correction was performed. Before plotting the spectra, the observed binding energies for both samples were calibrated with respect to carbon ($\text{C}_{1s} = 284.8 \text{ eV}$). The spectra of polycrystalline samples were analyzed using CasaXPS software.⁶³ Survey spectra of both compounds were collected in the binding energy range of 0 to 1200 eV.

Electronic structure calculations

Electronic structure calculations were performed for the $\text{Ba}_3\text{P}_5\text{Cl}$ and $\text{Ba}_3\text{P}_5\text{Br}$ compositions within the framework of the TB-LMTO-ASA code.⁶⁴ Two disorder-free structural models with lowered symmetries in space groups $C2/c$ and Cc were considered for calculations. Unit cell parameters and fractional atomic coordinates used for calculations are presented in Tables S1–S4. \dagger The visualization of the crystal structures used for calculations is presented in Fig. S7. \dagger The von-Barth-Hedin exchange-correlation functional was utilized, and empty spheres were introduced to satisfy the atomic sphere approximation (ASA) requirements.⁶⁵ Integration over the k -space was achieved using the tetrahedron method.⁶⁶ After confirming convergence, a $10 \times 10 \times 10$ k -point grid was employed to sample the Brillouin zone. The Fermi level, denoted as $E_F = 0 \text{ eV}$, was set as the reference energy level. We carried out density of states (DOS) and atom-projected density of states (PDOS) calculations to elucidate the electronic struc-

ture of these novel phosphide halide compounds. The calculations employed a basic set of orbitals including Ba [5s, (6p), 5d], P [3s, 3p, (3d)], Cl [(4s) 3p, (3d)], Br [(5s) 4p, (4d)]; the downfolded orbitals are in parentheses. Furthermore, Crystal Orbital Hamilton population (COHP) calculations were performed to provide deeper insights into the chemical bonding within these compounds.⁶⁷ All calculations were facilitated with the assistance of the modules available in the LMTO program.

Conclusions

Two new ternary phosphide halides, $\text{Ba}_3\text{P}_5\text{Cl}$ and $\text{Ba}_3\text{P}_5\text{Br}$, have been synthesized using the Pb flux technique. Both compounds represent a novel structure type and crystallize in the rhombohedral space group $R\bar{3}c$. These double Zintl salts are charge-balanced and are presumed to function as narrow band gap semiconductors. $\text{Ba}_3\text{P}_5\text{Cl}$ and $\text{Ba}_3\text{P}_5\text{Br}$ phases exhibit peculiar one-dimensional infinite helical ${}^1\text{P}^-$ chains of partially occupied phosphorus atoms. The discovery of these previously unknown compounds with unique structural unit opens a possibility for further exploration and development of novel functional inorganic materials and nanomaterials. The $\text{Ba}_3\text{P}_5\text{X}$ compounds may act as precursors for synthesizing hitherto unknown non-centrosymmetric solids by utilizing pre-existing helical phosphorus chains.⁶⁸ Further exploratory work involving different synthetic approaches to sequester these phosphorus chains will be worthwhile pursuing, as it could lead to the stabilization of this novel phosphorus form at the nanoscale or within a more suitable matrix than ${}^1\text{XBa}_3$ ⁵⁺, potentially enhancing the ambient stability of compounds featuring this unique building block.

Author contributions

Mohd Ishtiyak: writing – original draft, methodology, formal analysis, investigation, data curation, and validation. S. M. Gayomi K. Samarakoon, Thimira Kandabadage Don, Spencer R. Watts: formal analysis, validation. Sviatoslav Baranets: conceptualization, methodology, supervision, validation, project administration, writing – review and editing.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

This work was financially supported by the College of Science and Department of Chemistry at Louisiana State University (start-up funding). SB also acknowledges the LSU Provost's Fund for Innovation in Research – Council on Research Summer Stipend Program for the summer support. We thank

Dr Dongmei Cao for helping with XPS measurements and data processing. The XPS was performed at the Shared Instrumentation Facility (SIF) at Louisiana State University.

References

- 1 H. Kageyama, K. Hayashi, K. Maeda, J. P. Attfield, Z. Hiroi, J. M. Rondinelli and K. R. Poeppelmeier, *Nat. Commun.*, 2018, **9**, 772.
- 2 J. K. Harada, N. Charles, K. R. Poeppelmeier and J. M. Rondinelli, *Adv. Mater.*, 2019, **31**, 1805295.
- 3 Y. Nalawade, J. Pepper, A. Harvey, A. Griffin, D. Caffrey, A. G. Kelly and J. N. Coleman, *ACS Appl. Electron. Mater.*, 2020, **2**, 3233–3241.
- 4 B. Chitara, T. B. Limbu, J. D. Orlando, Y. Tang and F. Yan, *Nanoscale*, 2020, **12**, 16285–16291.
- 5 K. Xhaxhiu, C. Kvarnström, P. Damlin and K. Bente, *Mater. Res. Bull.*, 2014, **60**, 88–96.
- 6 L.-D. Zhao, J. He, D. Berardan, Y. Lin, J.-F. Li, C.-W. Nan and N. Dragoe, *Energy Environ. Sci.*, 2014, **7**, 2900–2924.
- 7 P. Vaqueiro, G. Guélou, M. Stec, E. Guilmeau and A. V. Powell, *J. Mater. Chem. A*, 2012, **1**, 520–523.
- 8 Y. Kamihara, T. Watanabe, M. Hirano and H. Hosono, *J. Am. Chem. Soc.*, 2008, **130**, 3296–3297.
- 9 Y.-F. Shi, W.-B. Wei, X.-T. Wu, H. Lin and Q.-L. Zhu, *Dalton Trans.*, 2021, **50**, 4112–4118.
- 10 B. M. Oxley, J. B. Cho, A. K. Iyer, M. J. Waters, J. He, N. C. Smith, C. Wolverton, V. Gopalan, J. M. Rondinelli, J. I. Jang and M. G. Kanatzidis, *J. Am. Chem. Soc.*, 2022, **144**, 13903–13912.
- 11 M. Yang, W. Liu and S.-P. Guo, *Inorg. Chem.*, 2022, **61**, 14517–14522.
- 12 Y. Subramanian, A. Dhanasekaran, L. A. Omeiza, M. R. Somalu and A. K. Azad, *Catalysts*, 2023, **13**, 173.
- 13 K. Chatterjee and S. E. Skrabalak, *ACS Appl. Mater. Interfaces*, 2021, **13**, 36670–36678.
- 14 B. Khamari and B. R. K. Nanda, *Mater. Res. Express*, 2019, **6**, 066309.
- 15 A. Chakraborty, N. Pai, J. Zhao, B. R. Tuttle, A. N. Simonov and V. Pecunia, *Adv. Funct. Mater.*, 2022, **32**, 2203300.
- 16 J.-A. Dolyniuk and K. Kovnir, *Crystals*, 2013, **3**, 431–442.
- 17 M. Ströbele, O. Oeckler, M. Thelen, R. F. Fink, A. Krishnamurthy, S. Kroeker and H.-J. Meyer, *Inorg. Chem.*, 2022, **61**, 17599–17608.
- 18 M. Ströbele, K. Eichele and H.-J. Meyer, *Eur. J. Inorg. Chem.*, 2011, **2011**, 4063–4068.
- 19 J.-A. Dolyniuk, N. Tran, K. Lee and K. Kovnir, *Z. Anorg. Allg. Chem.*, 2015, **641**, 1422–1427.
- 20 J.-A. Dolyniuk, S. Lee, N. Tran, J. Wang, L.-L. Wang and K. Kovnir, *J. Solid State Chem.*, 2018, **263**, 195–202.
- 21 Y. Qiao and H. Yin, *Nanoscale*, 2023, **15**, 9835–9842.
- 22 R. Nesper, *Z. Anorg. Allg. Chem.*, 2014, **640**, 2639–2648.
- 23 S. Baranets, G. M. Darone and S. Bobev, *Z. Kristallogr. - Cryst. Mater.*, 2022, **237**, 1–26.

24 E. R. Franke and H. Schäfer, *Z. Naturforsch., B: Anorg. Chem., Org. Chem., Biochem., Biophys., Biol.*, 1972, **27**, 1308–1315.

25 W. M. Hurng and J. D. Corbett, *Chem. Mater.*, 1989, **1**, 311–319.

26 C. Hadenfeldt and W. Held, *J. Less-Common Met.*, 1986, **123**, 25–35.

27 C. Hadenfeldt and H. Herdejürgen, *Z. Anorg. Allg. Chem.*, 1988, **558**, 35–40.

28 C. Hadenfeldt, *Z. Anorg. Allg. Chem.*, 1977, **436**, 113–121.

29 M. D. Donakowski, R. Gautier, J. Yeon, D. T. Moore, J. C. Nino, P. S. Halasyamani and K. R. Poeppelmeier, *J. Am. Chem. Soc.*, 2012, **134**, 7679–7689.

30 G. Cicirello, M. Wang, Q. P. Sam, J. L. Hart, N. L. Williams, H. Yin, J. J. Cha and J. Wang, *J. Am. Chem. Soc.*, 2023, **145**, 8218–8230.

31 J.-A. Dolyniuk, H. He, A. S. Ivanov, A. I. Boldyrev, S. Bobev and K. Kovnir, *Inorg. Chem.*, 2015, **54**, 8608–8616.

32 B. Eisenmann and U. Rößler, *Z. Anorg. Allg. Chem.*, 2003, **629**, 459–462.

33 R. D. Shannon, *Acta Crystallogr., Sect. A: Cryst. Phys., Diff., Theor. Gen. Crystallogr.*, 1976, **32**, 751–767.

34 M. Ishtiyak, G. Panigrahi, S. Jana, J. Prakash, A. Mesbah, C. D. Malliakas, S. Lebègue and J. A. Ibers, *Inorg. Chem.*, 2020, **59**, 2434–2442.

35 B. Cordero, V. Gómez, A. E. Platero-Prats, M. Revés, J. Echeverría, E. Cremades, F. Barragán and S. Alvarez, *Dalton Trans.*, 2008, 2832–2838.

36 G. A. Papoian and R. Hoffmann, *Angew. Chem., Int. Ed.*, 2000, **39**, 2408–2448.

37 S. Baranets, A. Ovchinnikov and S. Bobev, in *Handbook on the Physics and Chemistry of Rare Earths*, eds. J.-C. G. Bünzli and V. K. Pecharsky, Elsevier, 2021, vol. 60, pp. 227–324.

38 M. Aykol, J. W. Doak and C. Wolverton, *Phys. Rev. B*, 2017, **95**, 214115.

39 J. Liu, Y. Guo, S. Zhang, Q. Wang, Y. Kawazoe and P. Jena, *J. Phys. Chem. C*, 2015, **119**, 24674–24680.

40 A. Pfitzner, *Angew. Chem., Int. Ed.*, 2006, **45**, 699–700.

41 A. Pfitzner, M. F. Bräu, J. Zweck, G. Brunklaus and H. Eckert, *Angew. Chem., Int. Ed.*, 2004, **43**, 4228–4231.

42 Z. Sun, B. Zhang, Y. Zhao, M. Khurram and Q. Yan, *Chem. Mater.*, 2021, **33**, 6240–6248.

43 D. Liu, J. Guan, J. Jiang and D. Tománek, *Nano Lett.*, 2016, **16**, 7865–7869.

44 D. Pfister, K. Schäfer, C. Ott, B. Gerke, R. Pöttgen, O. Janka, M. Baumgartner, A. Efimova, A. Hohmann, P. Schmidt, S. Venkatachalam, L. van Wüllen, U. Schürmann, L. Kienle, V. Doppel, E. Parzinger, B. Miller, J. Becker, A. Holleitner, R. Weihrich and T. Nilges, *Adv. Mater.*, 2016, **28**, 9783–9791.

45 F. Reiter, M. Pielmeier, A. Vogel, C. Jandl, M. Plodinec, C. Rohner, T. Lunkenbein, K. Nisi, A. Holleitner and T. Nilges, *Z. Anorg. Allg. Chem.*, 2022, **648**, e202100347.

46 P. C. Donohue, *J. Solid State Chem.*, 1972, **5**, 71–74.

47 B. Owens-Baird, J. Xu, D. Y. Petrovykh, O. Bondarchuk, Y. Ziouani, N. González-Ballesteros, P. Yox, F. M. Sapountzi, H. Niemantsverdriet, Y. V. Kolen'ko and K. Kovnir, *Chem. Mater.*, 2019, **31**, 3407–3418.

48 J. Chen, C. Lin, F. Xu, S. Yang, Y. Sun, X. Zhao, X. Jiang, B. Li, T. Yan and N. Ye, *Chem. Mater.*, 2020, **32**, 10246–10253.

49 J. F. Moulder, W. F. Stickle, P. E. Sobol and K. D. Bomben, *Handbook of X-ray Photoelectron Spectroscopy; Physical Electronics, Inc.*: Eden Prairie, MN, 1995.

50 P. E. R. Blanchard, A. P. Grosvenor, R. G. Cavell and A. Mar, *Chem. Mater.*, 2008, **20**, 7081–7088.

51 J. Mark, J. Wang, K. Wu, J. G. Lo, S. Lee and K. Kovnir, *J. Am. Chem. Soc.*, 2019, **141**, 11976–11983.

52 H. T. Stokes, D. M. Hatch and B. J. Campbell, *ISODISTORT, ISOTROPY Software Suite*, <https://iso.byu.edu>.

53 B. J. Campbell, H. T. Stokes, D. E. Tanner and D. M. Hatch, *J. Appl. Crystallogr.*, 2006, **39**, 607–614.

54 S. Baranets, A. Balvanz, G. M. Darone and S. Bobev, *Chem. Mater.*, 2022, **34**, 4172–4185.

55 A. Balvanz, S. Baranets and S. Bobev, *J. Solid State Chem.*, 2020, **289**, 121476.

56 SAINT, BrukerAXS Inc., 2014.

57 SADABS, BrukerAXS Inc., 2014.

58 G. M. Sheldrick, *Acta Crystallogr., Sect. A: Found. Adv.*, 2015, **71**, 3–8.

59 G. M. Sheldrick, *Acta Crystallogr., Sect. C: Struct. Chem.*, 2015, **71**, 3–8.

60 O. V. Dolomanov, L. J. Bourhis, R. J. Gildea, J. K. Howard and H. Puschmann, *J. Appl. Crystallogr.*, 2009, **42**, 339–341.

61 L. M. Gelato and E. Parth'e, STRUCTURE TIDY - a computer program to standardize and crystal structure data, *J. Appl. Crystallogr.*, 1987, **20**, 139–143.

62 D. C. Palmer, *CrystalMaker*, CrystalMaker Software Ltd, Begbroke, Oxfordshire, England, 2014. <https://journals.iucr.org/m//services/stdswrefs.htm>.

63 N. Fairley, *CasaXPS, version 2.2.19*, Casa Software, Ltd, Teignmouth, Devon, U.K., 2003, <https://www.casaxps.com>.

64 O. Jepsen and O. K. Andersen, *The Stuttgart TB-LMTO program*, 1999.

65 U. von Barth and L. Hedin, *J. Phys. C: Solid State Phys.*, 1972, **5**, 1629.

66 P. E. Blöchl, O. Jepsen and O. K. Andersen, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1994, **49**, 16223–16233.

67 S. Steinberg and R. Dronskowski, *Crystals*, 2018, **8**, 225.

68 P. A. Maggard, C. L. Stern and K. R. Poeppelmeier, *J. Am. Chem. Soc.*, 2001, **123**, 7742–7743.