

## PAPER

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Cite this: *Nanoscale*, 2024, **16**, 657

# Novel hollow MoS<sub>2</sub>@C@Cu<sub>2</sub>S heterostructures for high zinc storage performance†

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Heterostructured materials have great potential as cathodes for zinc-ion batteries (ZIBs) because of their fast Zn<sup>2+</sup> transport channels. Herein, hollow MoS<sub>2</sub>@C@Cu<sub>2</sub>S heterostructures are innovatively constructed using a template-engaged method. The carbon layer improves the electrical conductivity, provides a high *in situ* growth area, and effectively restricts volume expansion during the recycling process. MoS<sub>2</sub> nanosheets are grown on the surfaces of hollow C@Cu<sub>2</sub>S nanocubes using the *in situ* template method, further expanding the specific surface area and exposing more active sites to enhance the electrical conductivity. As expected, an admirable reversible capacity of 197.2 mA h g<sup>-1</sup> can be maintained after 1000 cycles with a coulombic efficiency of 91.1%. Therefore, we firmly believe that this work points the way forward for high-performance materials design and energy storage systems.

 Received 17th October 2023,  
 Accepted 17th November 2023

DOI: 10.1039/d3nr05231a

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 † Electronic supplementary information (ESI) available. See DOI: <https://doi.org/10.1039/d3nr05231a>


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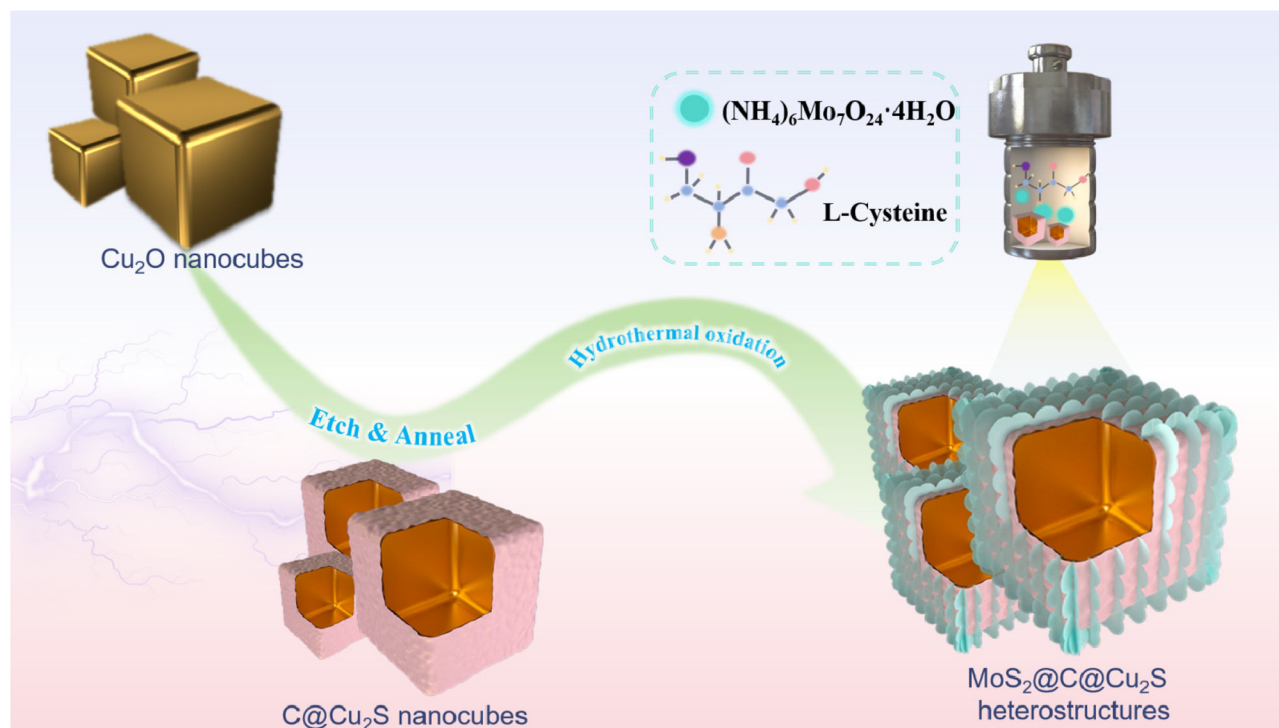
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## Introduction

With the development of advanced energy storage technologies, lithium-ion batteries (LIBs) have exhibited some inevitable defects, such as high cost and potential safety hazards.<sup>1–3</sup> Currently, aqueous system zinc ion batteries (ZIBs) are expected to be alternatives to LIBs due to their low cost, abundant resources, and high security.<sup>4,5</sup> Although there have been many advances in the research on ZIBs, the design of high-performance ZIBs still faces great challenges, such as slow diffusion of zinc ions and low reversibility, which seriously hinder their development and commercialization.<sup>6,7</sup>

As a new type of 2D structure, many transition metal dichalcogenides (TMDs) are attracting attention as novel energy storage materials, because of their efficient energy conversion.<sup>8–11</sup> For example, VS<sub>4</sub>@rGO nanosheets were explored using a hydrothermal process for super-performance ZIBs;<sup>12</sup> as a novel electrode, WS<sub>2</sub> nanosheets exhibited high specific capacity and were applied in sodium-ion batteries;<sup>13</sup> SnS<sub>2</sub>/Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub> nanosheets were fabricated using an *in situ* hybridization method for high-performance potassium-ion batteries.<sup>14</sup> More interestingly, CuS nanocubes exhibited high energy density when used as cathode materials,<sup>15</sup> and MoS<sub>2</sub> nanosheets are conducive to ion intercalation/deintercalation in lithium/sodium-ion batteries.<sup>16,17</sup> However, there are some challenges in the application of MoS<sub>2</sub> nanosheets in ZIBs, such as poor conductivity, higher hydrophilicity, agglomeration, and volume expansion.<sup>18–20</sup> Therefore, it is important to explore novel MoS<sub>2</sub> structures with high conductivity and cycling stability for energy storage.



**Scheme 1** Synthetic process of the hollow  $\text{MoS}_2\text{@C@Cu}_2\text{S}$  heterostructures.

Herein, we fabricated novel  $\text{MoS}_2\text{@C@Cu}_2\text{S}$  heterostructures using a template-engaged process for high-capacity  $\text{Zn}^{2+}$  storage (Scheme 1). These heterostructures are composed of  $\text{MoS}_2$  nanosheets and hollow  $\text{C@Cu}_2\text{S}$  nanocubes, and exhibit high specific surface area, exposing many active sites and enhancing the electrical conductivity in ZIBs. The hollow cavity facilitates rapid  $\text{Zn}^{2+}$  transport to reduce the energy barrier of  $\text{Zn}^{2+}$  absorption and transport, and reinforces the volume elasticity to accommodate volume change during  $\text{Zn}^{2+}$  de-/intercalation. At current densities of  $0.1 \text{ A g}^{-1}$  and  $1.0 \text{ A g}^{-1}$ , the as-designed ZIBs exhibit stable specific capacities of  $679.1 \text{ mA h g}^{-1}$  and  $192.9 \text{ mA h g}^{-1}$ , respectively.

## Results and discussion

The hollow  $\text{MoS}_2\text{@C@Cu}_2\text{S}$  heterostructures (size of  $\sim 600 \text{ nm}$ ) were prepared using an etching process followed by hydrothermal reduction (Scheme 1 and Fig. 1a and b). First, hollow  $\text{CuS}$  nanocubes with uniform sizes (size of  $\sim 500 \text{ nm}$ , Fig. S1b<sup>†</sup>) were prepared by the reaction of  $\text{Cu}_2\text{O}$  nanocubes (size of  $\sim 550 \text{ nm}$ , Fig. S1a<sup>†</sup>),  $\text{Na}_2\text{S}$ , and  $\text{HCl}$ . After carbonization,  $\text{PDA@CuS}$  nanocubes converted into hollow  $\text{C@Cu}_2\text{S}$  nanocubes (size of  $\sim 550 \text{ nm}$ , Fig. S1c<sup>†</sup>). Finally, hollow heterostructures were prepared after folding  $\text{MoS}_2$  nanosheets were grown on the surface of hollow  $\text{C@Cu}_2\text{S}$  nanocubes (Fig. 1c).

The X-ray powder diffraction (XRD) pattern shows that all peaks of the crystal phases are indexed to  $\text{MoS}_2$  (JCPDS: 37-1492) and  $\text{Cu}_2\text{S}$  (JCPDS: 72-2276) (Fig. S2<sup>†</sup>). As exemplified in elemental distribution scanning (Fig. 1e), the distribution of C can be seen in the intermediate layer (Fig. S3<sup>†</sup>). These hollow  $\text{CuS}$  nanocubes with folding  $\text{MoS}_2$  nanosheets and the C layer are  $\text{MoS}_2\text{@C@Cu}_2\text{S}$  heterostructures. As further evidence, the lattice diffraction fringe at  $0.606 \text{ nm}$  corresponds to the  $(0\ 0\ 2)$  crystal plane of  $\text{MoS}_2$ , and the fringe at  $0.260 \text{ nm}$  corresponds to the  $(1\ 0\ 3)$  crystal plane of  $\text{Cu}_2\text{S}$  in the high-resolution TEM image (Fig. 1d). X-ray photoelectron spectroscopy (XPS) further confirmed that the hollow heterostructure is composed of S, Mo, C, and Cu (Fig. S3<sup>†</sup>). The  $\text{E}_{1g}^*$  and  $\text{A}_{1g}$  peaks from  $\text{MoS}_2$ , the Cu-S vibration, and the D-band (defect vibration) and the G-band (graphitic carbon vibration) from C are observed in Raman spectroscopy (Fig. S4<sup>†</sup>), demonstrating the presence of  $\text{MoS}_2$ ,  $\text{CuS}$  and the carbon layer in the heterostructures. In addition, the heterostructures show typical *meso*-shell structures with a large specific surface area of  $34.10 \text{ m}^2 \text{ g}^{-1}$  (Fig. S5<sup>†</sup>).

The electrochemical properties of the  $\text{MoS}_2\text{@C@Cu}_2\text{S}$  heterostructures were investigated as shown in Fig. 2. As shown in Fig. 2a, the irreversible cathodic peaks at  $0.32 \text{ V}$  result from multiple electrochemical reactions, which involve the insertion of  $\text{Zn}^{2+}$  to form  $\text{Zn}_x\text{Cu}_2\text{S}/\text{Zn}_x\text{MoS}_2$  and to generate  $\text{ZnS}$  and  $\text{Mo}$ . In the subsequent positive scan, the peak at  $0.71 \text{ V}$  corresponds to the dezincing of  $\text{ZnS}$  to re-generate  $\text{Cu}_2\text{S}$  and  $\text{MoS}_2$ . Due to the formation of a solid electrolyte layer (SEI) film during the initial activation process, the loss of irre-

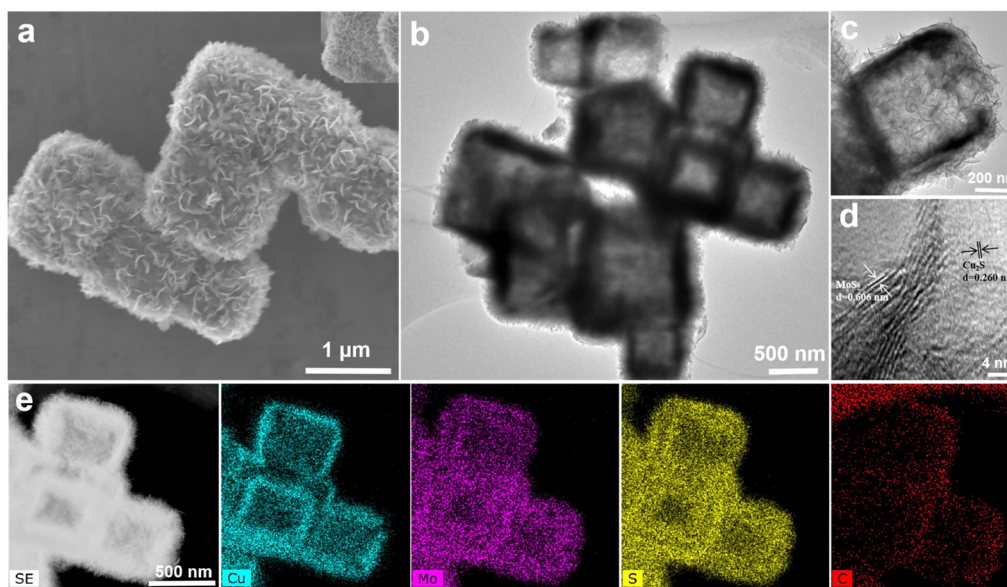


Fig. 1 Characterization of  $\text{MoS}_2@\text{C}@\text{Cu}_2\text{S}$  heterostructures: (a) SEM image; (b–d) TEM images; (e) elemental mapping images.

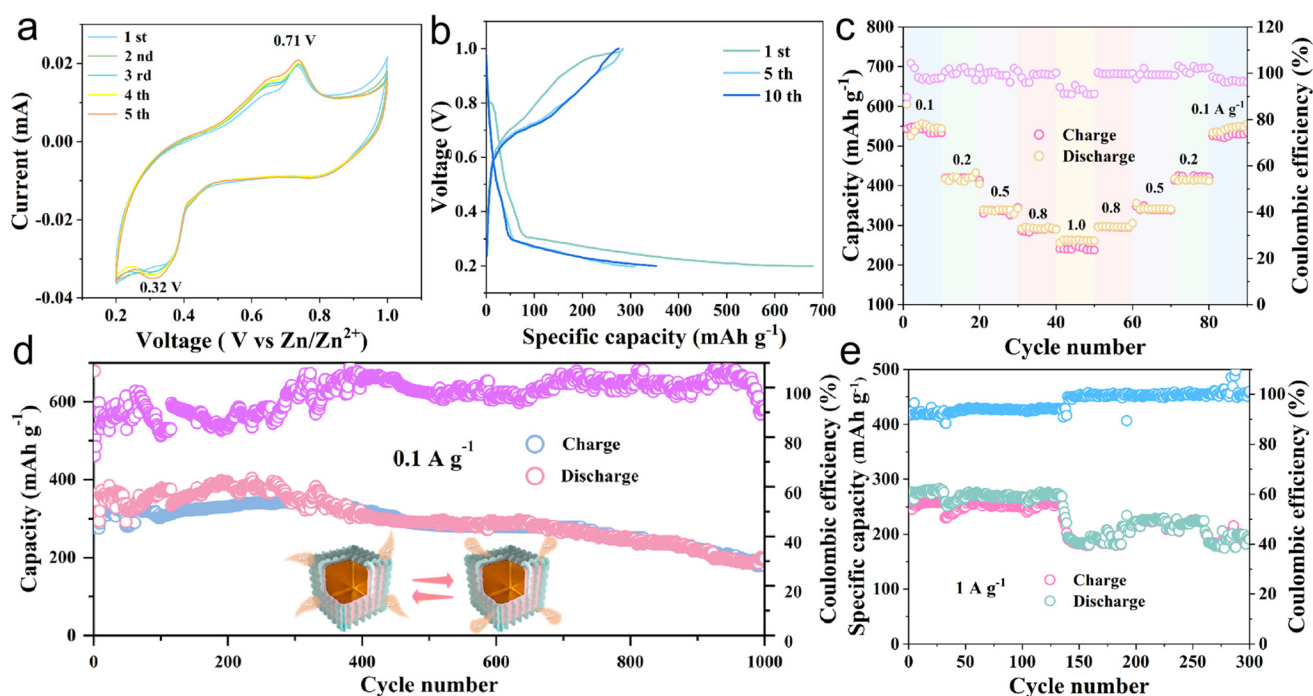


Fig. 2 The electrochemical performance tests of  $\text{MoS}_2@\text{C}@\text{Cu}_2\text{S}$  heterostructures: (a) CV profiles at  $0.1 \text{ mV s}^{-1}$ ; (b) charge–discharge curves of the first, fifth and tenth cycles; (c) rate performance; (d) long-term cycling performance; (e) high current density performance.

versible capacity results in a larger area in the initial cycle than that in the following cycles. From the second cycle to the fifth cycle, the curves show good overlap, indicating good cycling stability.

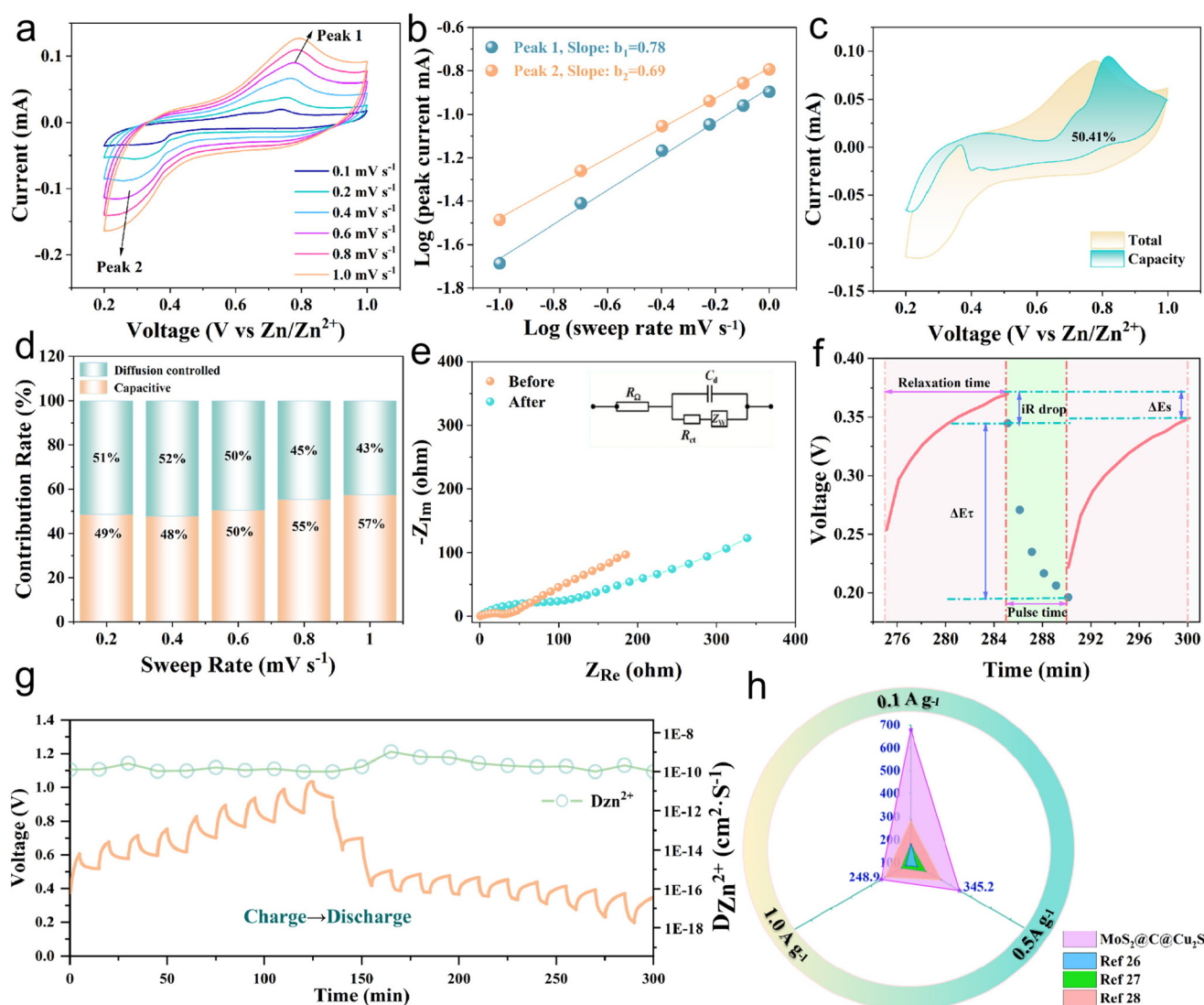
In Fig. 2b, the galvanostatic charge–discharge (GCD) test exhibits the charge–discharge curves of the heterostructures at a current density of  $0.1 \text{ A g}^{-1}$ . The charge and discharge curves

of the first, fifth, and tenth cycles show a high degree of similarity, indicating good cycling stability. For the insertion of  $\text{Zn}^{2+}$ , the voltage platforms are formed in the discharge process at  $0.32 \text{ V}$ . At  $0.71 \text{ V}$ , the platform is in the charging state, corresponding to the dezincing process. These results indicate good agreement between the GCD test and the CV curve (Fig. 2a and b). At different current densities of  $0.1, 0.2,$

0.5, 0.8, 1.0, 0.8, 0.5, 0.2, and 0.1 A g<sup>-1</sup>, the corresponding specific discharge capacities are 548.6, 420.2, 345.2, 292.2, 248.9, 304.8, 349.1, 424.3, and 534.1 mA h g<sup>-1</sup>, respectively (Fig. 2c).

The specific capacities of the heterostructures regularly decrease with increasing current density. In contrast, when the current density decreases, the specific capacity increases, proving the excellent rate performance. Fig. 2d shows the cycling performance and the corresponding coulombic efficiency of the heterostructures after 1000 cycles at 0.1 A g<sup>-1</sup>. The initial specific capacity is 679.1 mA h g<sup>-1</sup>, and the specific capacities during the first few cycles are gradually enhanced, which is related to the activation of the heterostructures.<sup>22</sup> After 1000 cycles, the specific capacity is 197.2 mA h g<sup>-1</sup> and the coulombic efficiency is 91.1%. Furthermore, at a high

current density, *i.e.*, 1 A g<sup>-1</sup>, the specific capacity is 192.9 mA h g<sup>-1</sup> after 300 cycles. These results confirm that the MoS<sub>2</sub>@C@Cu<sub>2</sub>S heterostructures have excellent cycling stability for ZIBs at high current density. In order to investigate the reaction mechanism of Zn<sup>2+</sup>, the electrochemical characteristics of MoS<sub>2</sub>@C@Cu<sub>2</sub>S heterostructures in an aqueous electrolyte were studied using CV in the potential range of 0.2–1.0 V (vs. Zn<sup>2+</sup>/Zn). In Fig. 3a, a pair of redox peaks at 0.25/0.75 V is observed, representing the reversible de-intercalation process of Zn<sup>2+</sup> in MoS<sub>2</sub> nanosheets. There is no change in the shape of the CV curves when the scanning rate increases from 0.1 to 1.0 mV s<sup>-1</sup>, which indicates that the heterostructures possess good electrochemical reversibility. With an increased scan rate, the anodic peak moves to a high potential (*i.e.*, 0.76 V) from 0.72 V, and the cathodic peak shifts to a low potential



**Fig. 3** (a) CV profiles at different scan rates; (b) plot of log (sweep rate) versus log (peak current) for anodic and cathodic peaks; (c) the capacitive contribution to the total capacity at 0.6 mV s<sup>-1</sup>; (d) capacity contribution at different scan rates; (e) EIS plots of MoS<sub>2</sub>@C@Cu<sub>2</sub>S heterostructures before and after 100 cycles; (f) GITT test conditions for the heterostructures during the discharge process; (g) GITT profiles and the calculated  $D_{\text{Zn}^{2+}}$  values; (h) comparison of the rate performance with those of previously reported MoS<sub>2</sub>-based electrode materials applied in ZIBs.

(i.e., 0.25 V) from 0.38 V. These potential changes are related to the enhanced diffusion resistance.

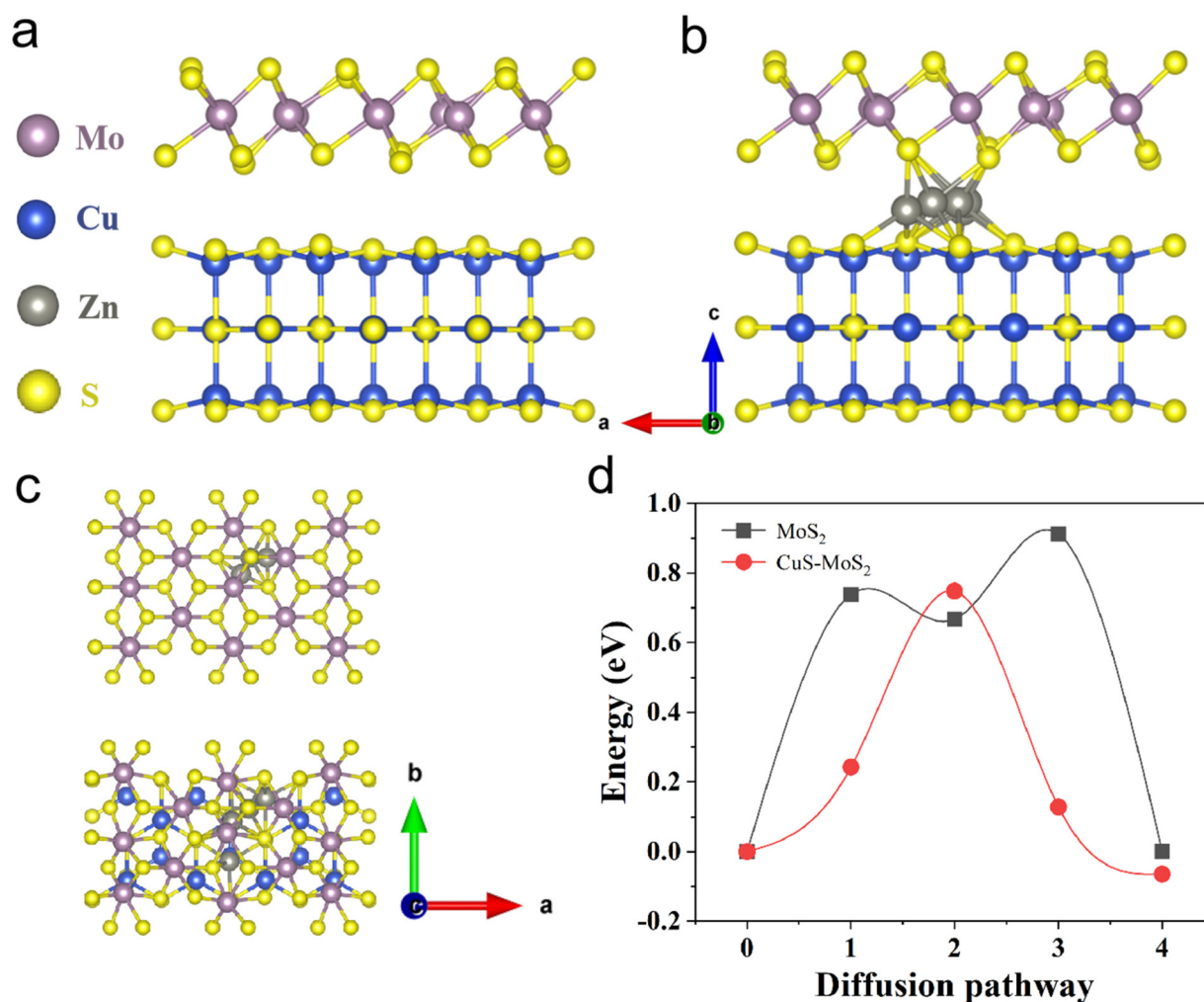
The electrochemical kinetics of the reduction of  $\text{Zn}^{2+}$  are presented in Fig. 3b. The power formula shows that  $I = a\nu^b$ ,  $\log(i) = \log(a) + b \log(\nu)$ , where  $I$  is the current,  $\nu$  is the scanning rate, and  $a$  and  $b$  are adjustable parameters. The coefficient  $b$  represents the type of electrochemical charge storage. When  $b$  is in the range of 0.5 to 1, the capacitance controls the process of charge storage.<sup>21</sup> Here,  $b$  was calculated to be 0.78 and 0.69 for the oxidation peak and the reduction peak, respectively, indicating that the  $\text{Zn}^{2+}$  storage process is controlled by a synergistic mechanism between the capacitive contribution and the diffusion process. The contributions of the capacitive effect ( $k_1\nu$ ) and diffusion effect ( $k_2\nu^{1/2}$ ) were further calculated according to the formula:  $i = k_1\nu + k_2\nu^{1/2}$ . The contribution ratio of the capacitance to the diffusion capacity of the heterostructures was calculated to be ~50.41% at a scan rate of 0.6  $\text{mV s}^{-1}$  (Fig. 3c). In Fig. 3d, with the increasing scan rate, i.e., 0.2 to 1.0  $\text{mV s}^{-1}$ , the capacitance contribution

is enhanced from 49% to 57%, indicating that a capacitive-controlled process dominates the storage behavior to promote the rate performance and capacity performance of ZIBs.

Next, electrochemical impedance spectroscopy (EIS) and galvanostatic intermittent titration technique (GITT) tests were executed to evaluate the kinetics of the electrode reaction in detail (Fig. 3e and f). There is a small semicircle diameter in the EIS curves of  $\text{MoS}_2\text{@C@Cu}_2\text{S}$  heterostructures before and after the cycle (Fig. 3e), indicating that the heterostructures possess good electrical conductivity and rapid  $\text{Zn}^{2+}$  kinetic diffusion. Based on Fick's second law, the difference in  $\text{Zn}^{2+}$  diffusion kinetics is explored as follows:

$$D_{\text{Zn}^{2+}} = \frac{4}{\tau\pi} \left( \frac{m_B V_M}{S M_B} \right)^2 \left( \frac{\Delta E_S}{\Delta E_\tau} \right)^2$$

where  $\tau$ ,  $m_B$ ,  $M_B$ ,  $V_M$ , and  $S$  stand for the current pulse time (s), the mass, the molar mass, the molar volume of active



**Fig. 4** (a) Optimized structure of  $\text{MoS}_2\text{@C@Cu}_2\text{S}$ ; (b) optimized structure of  $\text{Zn}^{2+}$  insertion in  $\text{MoS}_2\text{@C@Cu}_2\text{S}$ ; (c) top view of the optimized structure of  $\text{Zn}^{2+}$  insertion in pure  $\text{MoS}_2$  and  $\text{MoS}_2\text{@C@Cu}_2\text{S}$ ; (d) energy barriers of diffusion behavior of  $\text{Zn}^{2+}$  in pure  $\text{MoS}_2$  and  $\text{MoS}_2\text{@C@Cu}_2\text{S}$ .

materials, and the area of the electrode, respectively;<sup>23</sup> the calculation method of  $\Delta E_s$  and  $\Delta E_t$  is also displayed in Fig. 3f. It is clearly seen that the pulse time is 5 min, the relaxation time is 10 min, and the quasi-open circuit voltage is reached, and by calculation the  $D_{\text{Zn}^{2+}}$  was obtained in the range from  $10^{-8}$  to  $10^{-10}$  cm<sup>2</sup> s<sup>-1</sup> (Fig. 3g), indicating that the electrode possesses good reaction kinetics.<sup>24,25</sup> Compared with previous reports on MoS<sub>2</sub>-based zinc-ion batteries (Fig. 3h),<sup>26–28</sup> MoS<sub>2</sub>@C@Cu<sub>2</sub>S heterostructures exhibit high performance for ZIBs at different current densities. The excellent electrochemical characteristics of the heterostructures are closely related to the following factors: (1) ultra-thin MoS<sub>2</sub> nanosheets cover the surface of hollow C@Cu<sub>2</sub>S nanocubes, which would promote fast ion transport and effectively alleviate the volume expansion;<sup>29</sup> (2) the hollow cavity, the high specific surface area (*i.e.*, 34.10 m<sup>2</sup> g<sup>-1</sup>), and a large number of electrochemical active sites would shorten the ion diffusion path;<sup>19,25,30</sup> (3) the intermediate C layers maintain the stability of the heterostructures.<sup>26</sup> All three factors enable MoS<sub>2</sub>@C@Cu<sub>2</sub>S heterostructures to achieve superior reversible cycling performance. Density functional theory (DFT) calculations were used to analyze the diffusion of Zn<sup>2+</sup> inserted in the lattice of MoS<sub>2</sub>@C@Cu<sub>2</sub>S heterostructures at different energy barriers. Fig. 4a exhibits the theoretical laminated structure, and Fig. 4b shows the heterostructure of the interlayer model of Zn<sup>2+</sup> insertion. Moreover, compared with the Zn<sup>2+</sup> diffusion barriers of pure MoS<sub>2</sub> (Fig. 4c and d), the interlayer diffusion energy barrier of Zn<sup>2+</sup> decreases in MoS<sub>2</sub>@C@Cu<sub>2</sub>S heterostructures after MoS<sub>2</sub> nanosheets grow on the surface of hollow C@Cu<sub>2</sub>S nanocubes, which further indicates that the heterostructures are conducive to Zn<sup>2+</sup> diffusion.

## Conclusions

In summary, novel hollow MoS<sub>2</sub>@C@Cu<sub>2</sub>S heterostructures were fabricated using a template-engaged process for Zn<sup>2+</sup> storage. These heterostructures with sizes of ~600 nm are packed by folding MoS<sub>2</sub> nanosheets and hollow C@Cu<sub>2</sub>S nanocubes. Because of the high surface area, *i.e.*, 34.10 m<sup>2</sup> g<sup>-1</sup>, these heterostructures exhibit a high initial specific capacity of 679.1 mA h g<sup>-1</sup> and a reversible capacity of 197.2 mA h g<sup>-1</sup> with a coulombic efficiency of 91.1% after 1000 cycles at 0.1 A g<sup>-1</sup>. When the capacity was tested at 1.0 A g<sup>-1</sup>, the structure delivered 192.9 mA h g<sup>-1</sup> after 300 cycles. Furthermore, the hollow heterostructures exhibit excellent rate performance. Therefore, this work provides an effective design of heterogeneous materials for applications in power grid energy storage systems using ZIBs in the future.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

The authors acknowledge the financial support from the Talent Introduction Start-Up Foundation of Guangxi Minzu University (2021KJQD08), the Natural Science Foundation of China (No. 22074130), the Key Scientific Research Projects of Henan Province (232102320040), the National Natural Science Foundation of China (22004032) and the Hunan Provincial Natural Science Foundation of China (2021JJ20020). The authors also acknowledge the great support from the Analysis Testing Center of Xinyang Normal University for materials characterization.

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