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Vertical heterostructure of graphite–MoS₂ for gas sensing†

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2D materials, given their form-factor, high surface-to-volume ratio, and chemical functionality have immense use in sensor design. Engineering 2D heterostructures can result in robust combinations of desirable properties but sensor design methodologies require careful considerations about material properties and orientation to maximize sensor response. This study introduces a sensor approach that combines the excellent electrical transport and transduction properties of graphite film with chemical reactivity derived from the edge sites of semiconducting molybdenum disulfide (MoS₂) through a two-step chemical vapour deposition method. The resulting vertical heterostructure shows potential for high-performance hybrid chemiresistors for gas sensing. This architecture offers active sensing edge sites across the MoS₂ flakes. We detail the growth of vertically oriented MoS₂ over a nanoscale graphite film (NGF) cross-section, enhancing the adsorption of analytes such as NO₂, NH₃, and water vapor. Raman spectroscopy, density functional theory calculations and scanning probe methods elucidate the influence of chemical doping by distinguishing the role of MoS₂ edge sites relative to the basal plane. High-resolution imaging techniques confirm the controlled growth of highly crystalline hybrid structures. The MoS₂/NGF hybrid structure exhibits exceptional chemiresistive responses at both room and elevated temperatures compared to bare graphitic layers. Quantitative analysis reveals that the sensitivity of this hybrid sensor surpasses other 2D material hybrids, particularly in parts per billion concentrations.

New concepts

Our work presents an innovative design that comprises a unique stacking arrangement of heterolayers of 2D materials: molybdenum sulfide over graphitic layers. The vertically oriented molybdenum sulfide's (MoS₂'s) edge sites serve as receptors for gaseous analytes, ensuring selectivity, while horizontally placed graphene acts as the conductive medium for enhanced sensing performance. This architecture creates next-generation chemiresistors with superior performance achieved through synergy between hetero-atoms of 2D materials. The study further unravels atomic-scale resolution and deciphers the selectivity mechanism towards oxidizing and reducing gases in dry and humid conditions. This is the first report in the literature demonstrating the potential of vertical heterostructure as a gas sensor. These findings bridge the key knowledge gap of heterointerfaces between 2D materials, extending beyond graphitic carbon sheets. We believe our results will capture the attention of a broad audience, specifically the researchers interested in the nanoscience of 2D materials, heterostructures and gas sensors. This work will interest scientists and practitioners interested in utilizing 2D materials for selecting and targeting sensing in diverse domains, including healthcare.

Introduction

Environmental pollutants are a major concern for air quality, often associated with industrial and transport emissions, with

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significant consequences for public and individual health.¹ Nearly 4.2 million premature mortalities are reported annually due to direct results from exposure to air pollutants, including oxides of nitrogen, sulfur and carbon.² This has driven demand for gas sensing devices for domestic and public spaces, industrial and transportation environments, for the detection of toxic domestic gases and environmental monitoring.³ To this end, intensive research has fulfilled the global need for cost-effective, highly sensitive, and miniaturised gas sensors that can detect adsorbed analytes with efficient response and recovery time. Traditional approaches for gas sensing include metal oxide semiconductors, which have been implemented for commercial applications.⁴ Nevertheless, several drawbacks remain, including high operating temperature ($>200\text{ }^{\circ}\text{C}$), degradation of the sensitivity and stability of the sensor owing to variation in the microstructure of the metal oxide at high temperatures, significant baseline drift, poor selectivity, and the strong influence of humidity.⁵

2D materials have proven to be an alternative for chemiresistive gas sensors. They can be fabricated in inexpensive devices with minimal power consumption, extreme flexibility, and portability.⁶ Transition metal dichalcogenides (TMDs) have attracted huge interest in gas sensing due to their unique electronic and chemical functionalisation.^{7–11} Among the TMDs, molybdenum disulfide (MoS_2) has been widely studied as a multi-functional material due to its tuneable properties in electronics and sensors.^{12–17} MoS_2 sheets grown by chemical vapor deposition (CVD) with a controlled number of layers, sizes, and arbitrary shapes, for instance, nanoplates, nanorods, nanoflowers, *etc.* have been reported.^{8,16–21} This has led the recent development of heterostructures of MoS_2 along with other layered materials (*e.g.*, $\text{MoS}_2/\text{graphene}$,^{22,23} MoS_2/WS_2 ,²⁴ MoS_2/CdSe ,²⁵ *etc.*). Among them, heterostructures of $\text{MoS}_2/\text{graphitic carbon}$ are particularly fascinating due to outstanding properties arising from the van der Waals interfaces,^{26–29} *i.e.*, without chemical alteration at the interface that could influence the transport properties. These ultrathin MoS_2 structures with abundant active edge sites coupled to graphitic carbon support are extremely useful in enhancing the device performance in sensing, catalysis, and battery applications.^{26,30} MoS_2 flakes on carbon-based materials present a complementary architecture where individual layers contribute to each other as chemically active and electrically conductive, respectively.³¹ A wide range of applications have been demonstrated where MoS_2 is grown/transferred on conducting substrate for electronics and electrocatalysis.^{17,32,33} Nevertheless, heterostructures comprising MoS_2 and graphitic carbon as chemiresistors remain relatively unexplored, and comprehensive studies are needed in this direction.

In the present work, a heterostructure of vertically grown MoS_2 on horizontal NGF is demonstrated as a chemiresistor for gas sensing applications, where the edges of semiconducting MoS_2 promote adsorption of analyte gases (NO_2 , NH_3), while graphene as a conducting transducer. The preparation of the hybrid architecture comprises a two-step CVD process with careful control over the growth to yield centimetre-scale

heterostructure films. Surface analysis techniques have been implemented to reveal the morphology of atomically resolved edge atoms of the hybrid structure with potential lattice planes at the graphene– MoS_2 interface. Real-time gas sensing response has been demonstrated in different conditions: ambient, humid and moderately elevated temperatures ($150\text{ }^{\circ}\text{C}$). The doping mechanism is interpreted through Raman spectroscopy with the aid of Kelvin probe surface potential maps and corroborated through density functional theory simulation. This hybrid structure of MoS_2/NGF demonstrates the synergy between the two different systems for sensitivity as well as selectivity of chemical and gas detection. This design of hybrid structures will be useful for detecting a broad range of analytes in extremely low concentrations of parts per billion (ppb) and for developing sensitive and portable chemiresistors.

Results and discussion

The atomic scale structure of the hybrid architecture of MoS_2 on graphene is illustrated in Fig. 1a, where vertically aligned MoS_2 sheets are grown over NGF. The preparation of the hybrid structures through various processing steps is described in detail in Fig. S1 (ESI[†]). The MoS_2 and NGF sheets were grown on Si/SiO₂ (oxide thickness 300 nm), with the elemental composition of the $\text{MoS}_2/\text{NGF}/\text{SiO}_2/\text{Si}$ structure shown in Fig. S2 (ESI[†]) and its vertically aligned architecture over a horizontally stacked and wrinkled graphite sheet in Fig. S3 (ESI[†]). In this heterostructure, the edge atoms of semiconducting MoS_2 flakes are exposed to the local environment and connected to the conducting network of graphite layers, which leads to adsorption, enhanced charge transfer and sensitivity of the device.³⁰ The scanning electron microscopy (SEM) images in Fig. 1b of the cross-section view of the hybrid structure clearly distinguish between NGF (thickness between horizontal red color dashed line) and MoS_2 layer (thickness between horizontal yellow line). The top view of the heterostructure reveals densely packed MoS_2 grains over the basal plane and wrinkled regions that retain the NGF topographical feature (Fig. 1c). The height distribution of the MoS_2 flakes over graphene is illustrated through an AFM image and line profile (inset) in Fig. 1d, revealing size variation from 5 to 20 nm. Close observation reveals that the granular-shaped MoS_2 flakes are aligned in the vertical direction, corroborated by high-resolution SEM micrographs in Fig. S4 (ESI[†]). The width of the aligned MoS_2 granules plays an essential role in having the appropriate number of edge active sites towards the gaseous analyte and interfacial interaction between MoS_2 and carbon atoms for effective charge transfer. Recently, Mao and coworkers demonstrated that too thin or small granules of MoS_2 sheets could lead to incomplete crystallization.³⁴ Nevertheless, uncontrolled thickness could lead to continuous layers of MoS_2 with fewer dangling edge S atoms that alter its performance and efficiency.³⁰

It is crucial to monitor the crystal planes of the heterostructure to understand the arrangement of edge sulfur atoms in MoS_2 grains and their configuration at the interface with the



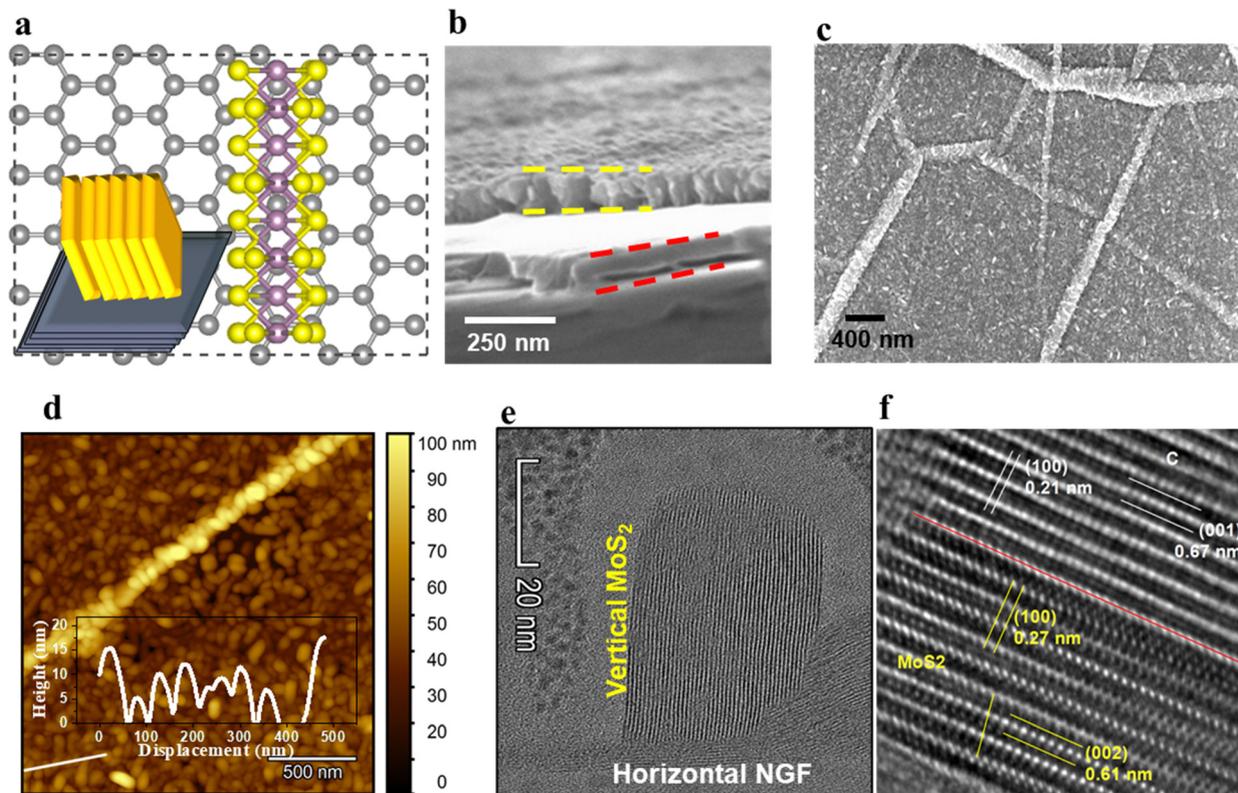


Fig. 1 Micro-to-atomic scale visualisation of the hybrid structure. (a) Schematic set-up of vertical hierarchy and simulated atomic arrangement of top view of MoS₂/NGF used as a hybrid structure. (b) SEM micrograph of a cross-section revealing the hierarchy of different sheets as MoS₂ (region marked with yellow dashed line). (c) Top view SEM micrograph of MoS₂/NGF, where grains of MoS₂ topping is distributed over different landscapes of NGF. (d) AFM image showing the height variations of MoS₂/NGF at different landscapes of edges and basal plane; inset shows the line profile of grain distribution. (e) Nanoscopic resolved TEM images of MoS₂ and NGF revealing the vertically aligned MoS₂ atoms of graphene layers. (f) The atomically resolved interface between MoS₂ and NGF shows the similar basal plane of NGF and MoS₂ as (100) close to the interface axis (red colour) and (002) MoS₂ plane far from the interface axis.

carbon atoms. Therefore, atomic scale lattice planes are analysed through high-resolution transmission electron microscopy (HR-TEM), revealing the MoS₂ (100) and (002) lattice planes at 0.27 nm (in-plane S-S separation) and 0.61 nm (out-of-plane sheet separation) respectively, and for the graphite (100) and (002) lattice planes at 0.21 nm (in-plane carbon separation) and 0.335 nm (0.67/2 nm, out-of-plane graphene sheet separation) respectively (Fig. 1e and f). The atomic arrangement at the interface in the hybrid structure reveals epitaxial symmetry between MoS₂ and NGF connected through contact planes (002) (Fig. 1f). This epitaxy can be described by the following symmetry relations as NGF [002]//MoS₂ [002], NGF [100]//MoS₂ [100], NGF *b*-axis//MoS₂ *b*-axis, where the first term represents the planar alignment, the second term is direction alignment of the interface, and the third term is the zone axis alignment. The lattice mismatch (η) between the two crystal structures is defined as the absolute difference between two lattice spacings (d_1 , d_2) along a certain crystal direction relative to the average of the lattice spacings given by relation (1):³⁵

$$\eta = \frac{|d_1 - d_2|}{\frac{1}{2}(d_1 + d_2)} \times 100 \quad (1)$$

The η values for the corresponding alignments are 58.7% for the planar equation, 24.6%, and 24.9% for the two vector equations, respectively. These findings are consistent with a semi-coherent NGF/MoS₂ interface where five NGF (002) lattice units fit four MoS₂ (002) lattice units. The attachment of sulfur atoms to central Mo atoms (*i.e.* S-end) are observed, which could play a pivotal role in detecting gases as unravelled through the density functional theory simulation.

XPS measurement of the hybrid structure reveals the elemental composition as well as the oxidation state of Mo. The NGF/SiO₂/Si configuration without MoS₂ sheet shows an asymmetric peak centred at 284.9 eV of C 1s core level (Fig. 2a), corresponding to graphitic carbon.³⁶ Nevertheless, traces of O content as a contaminant from air exposure have been observed in Fig. 2b. Interestingly, no C signal is detected after the deposition of smaller size Mo film (≈ 10 nm), indicating complete coverage of Mo films. Here, the O 1s and Mo 3d (Fig. 2b and c) confirm molybdenum oxide formation, with numerous peaks contribution. These contributions are attributed to the multiple spin-orbit doublets of MoO₂ (first doublet at 229.5 eV). The traces of Mo oxide phase is verified by a small doublet peak centered at 232.8 eV (Fig. 2c) and at O 1s core-level spectrum at 530.3 eV (Fig. 2b). The presence of MoO₃ is also confirmed by peak contributions at 233 eV. The oxidation of



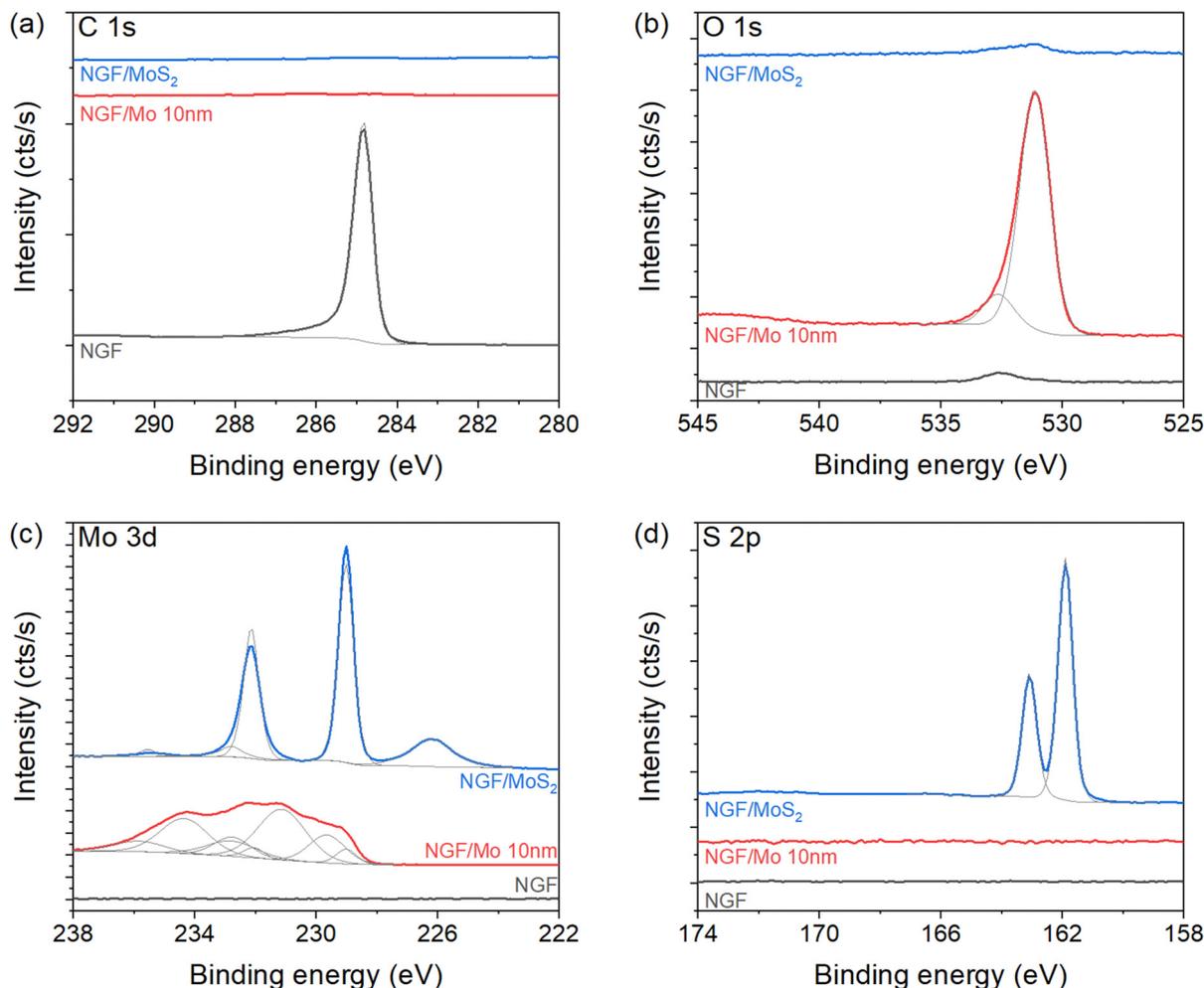


Fig. 2 Interfacial chemistry, oxidation state and charge carriers in hybrid structure. XPS analysis of the hybrid structure of different MoS₂ sizes and NGF with (a) C 1s, (b) O 1s, (c) Mo 3d, and (d) S 2p core level spectra of NGF, Mo/NGF and MoS₂/NGF.

sputtered Mo film is recurrent and confirmed on O 1s level, with a major contribution at 530.3 eV corresponding to Mo–O bonds. After sulfurisation, the formation of MoS₂ is confirmed at Mo 3d and S 2p core levels (Fig. 2c and d). The Mo signal is composed of a doublet centred at 229.1 eV (marked by a solid blue line) corresponding to Mo⁴⁺, while the contribution at 226.2 eV is attributed to S 2s signal. These results are in excellent agreement with the reference literature reporting MoS₂ formation.^{12,37–39} Additionally, the S 2p signal (Fig. 2d) is composed of a single doublet centred at 161.9 eV, further revealing the formation of MoS₂.^{12,15} The Mo:S ratio of 1:2, calculated by integrating the area of the Mo and S signals, confirms the growth of high-quality MoS₂ in good agreement with previous reports.^{39,40}

The evaluation of structural defects in graphene after Mo deposition and sulfurisation is investigated through Raman spectroscopy. The Raman modes of hybrid structure show MoS₂ peak at E_{2g}¹ ≈ 385 cm⁻¹, A_{1g} ≈ 408 cm⁻¹ and NGF (G ≈ 1582 cm⁻¹ and 2D ≈ 2756 cm⁻¹) (Fig. 3a). It is observed that the thinner regions of NGF (<10 layers) show structural

disorder validated through the D peak (≈1350 cm⁻¹), specifically after sulfurisation. A predominantly aligned MoS₂ flake over NGF leads to enhanced exposure of the edge sites, showing the higher intensity of the A_{1g} signal over E_{2g}¹ due to edge-induced Fermi level shift.^{15,16,41} The Raman shift in E_{2g}¹ (cm⁻¹) and A_{1g} (cm⁻¹) modes are quantified through Lorentzian peak fitting of Raman modes after exposure. The correlation plot of frequency shift (cm⁻¹) in A_{1g} as a function of E_{2g}¹ is presented in Fig. 3b exclusively for atomic layer investigation, carried out separately, exposing the derivative of exposed ions (i.e. nitronium and ammonium). The mechanical strain and the electrical doping in MoS₂ can be related by a linear transformation of the peak shifts of E_{2g}¹ and A_{1g} as described in our previous findings,^{42,43} and the strain and doping axis can then be constructed by taking negligible doping and strain respectively (see ESI† S5 for details). The projection of the peak position deviations (from the doped-free and strain-free intersection point) to the strain and doping axis represents the localised strain and doping of MoS₂, enabling the successful deconvolution of strain and doping in MoS₂ after the exposure of analytes.



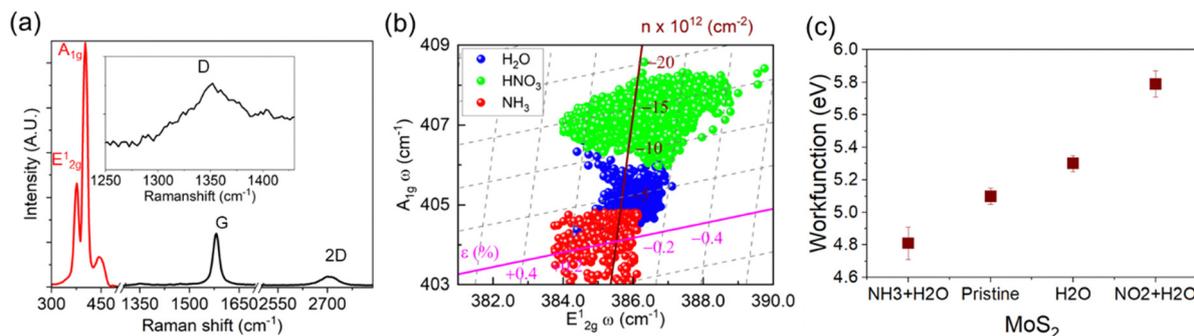


Fig. 3 Analytes induced doping in MoS₂. (a) Raman spectrum of the hybrid structure showing Raman modes of MoS₂ (E_{12g}¹ and A_{1g}) and few-layer graphene (G and 2D). (b) Correlation plot (E_{12g}¹ vs. A_{1g}) shows the shift in peak position in the MoS₂ Raman modes after exposure of solvents of dilute HNO₃ (as NO₂ + H₂O), NH₄OH (as NH₃ + H₂O) and H₂O. (c) Trend of work function (eV) of MoS₂ in different exposures of analytes.

Treatment with DI water on MoS₂ upshifts the A_{1g} peaks, showing a decrease in carrier concentration (p-type doping), which is further enhanced in the presence of dilute HNO₃. Thus, DI water acts as a milder electron acceptor (due to the presence of OH⁻)⁴⁴ than dilute HNO₃ (NO₃⁻), which will be crucial to explain the selectivity phenomenon in detecting NO₂ gases in dried and humid conditions. The Raman correlation plot results indicate that the adsorption of NO₂ gas in humid conditions would be expected to result in higher p-type doping of MoS₂ than in dry conditions. The influence of the ammonium ions redshifted the A_{1g} Raman modes of MoS₂ indicating increase in carrier concentration, nevertheless, it is not as potent as nitronium ions in the presence of water molecules.

The outcomes from Raman spectroscopy are validated through surface potential measurements by KPFM (Fig. 3c), where doping characteristics of MoS₂ in the heterostructured architecture are investigated separately in the exposed environment of dilute HNO₃ (as a mixture of H₂O and NO₂) and dilute NH₄OH through Kelvin probe force microscopy (KPFM). The surface potential map indicates an increase in electron concentration in the MoS₂, revealing n-type doping from the substrates (see ESI† Fig. S6). In the heterostructure arrangement of MoS₂ over graphite, the exposure of nitronium ions leads to p-type doping in MoS₂ as compared to the subsurface graphite. In

contrast, ammonium ions increase the electron concentration. The modulation in Fermi energy level in MoS₂ due to adsorbates leads to alteration in the work function (~5.1 eV) of MoS₂. The adsorption of NH₄⁺ salts decreases the work function (~4.81 eV), while the adsorption of nitronium salt significantly increases it to ~5.79 eV.

The charge redistribution in the graphene/MoS₂ hybrid structure for different analyte molecules (H₂O, NH₃, and NO₂) is further unravelled through density functional theory simulations. To find the best stacking configuration of MoS₂ on graphene, we calculate the binding energy ($E_{\text{graphene}} + E_{\text{MoS}_2} - E_{\text{total}}/S$) for the stacking configurations shown in Fig. S7 (ESI†), where E_{graphene} , E_{MoS_2} , and E_{total} are the total energies of graphene, MoS₂, and the graphene/MoS₂ hybrid structure, and S is the interfacial area (3.16 Å × 12.33 Å). With binding energies of 0.05 and 0.03 eV Å⁻², the AB1 and AB configurations are energetically favourable for S- and Mo-terminated MoS₂, respectively. Fig. 4a and b show that MoS₂ is a charge acceptor and accepts much more charge in the case of the S-termination (0.57 electrons) than in the case of the Mo-termination (0.10 electrons). The charge redistributions between the graphene/MoS₂ hybrid structure and analyte molecules are shown in Fig. 4c–e. In the proximity of the NO₂ molecule close to the MoS₂ edge atoms, there is a significant

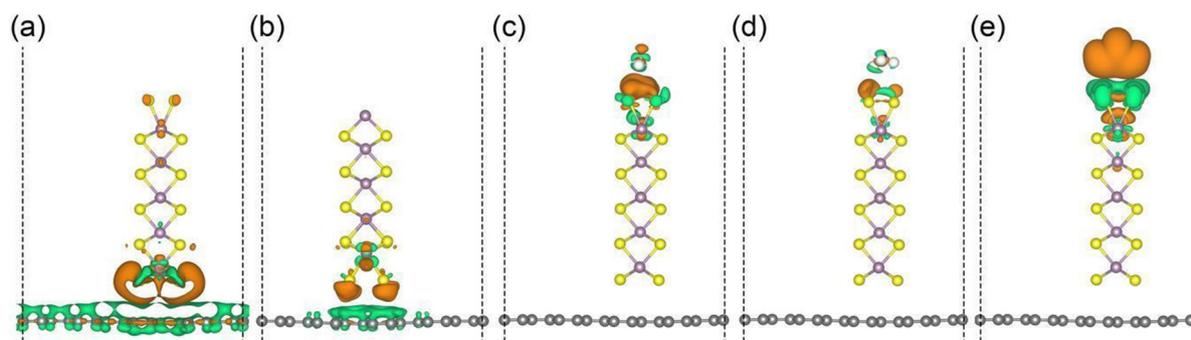


Fig. 4 Effect of atomic arrangement on charge transfer. Side views of the charge redistributions between graphene and MoS₂ with (a) S- and (b) Mo-terminations. Side views of the charge redistributions between the graphene/MoS₂ hybrid structure with S-termination (towards the analyte molecules) and (c) H₂O, (d) NH₃, and (e) NO₂. Green and orange isosurfaces (isovalue: 0.001 electrons Å⁻³) represent charge depletion and accumulation, respectively. The gray, yellow, purple, white, red, and blue spheres represent the C, S, Mo, H, O, and N atoms, respectively.



depletion of charge (green colour) around the sulphur atoms, due to interaction with the unpaired electron of the N atom, and an accumulation of charge on NO₂ (orange colour). This charge transfer is lower for the O atom of H₂O. The adsorption energy of the analyte molecule is calculated as $E_{\text{total}} - E_{\text{graphene/MoS}_2} - E_{\text{molecule}}$, where E_{total} , $E_{\text{graphene/MoS}_2}$, and E_{molecule} are the total energies of the graphene/MoS₂ hybrid structure with adsorbed molecule, the graphene/MoS₂ hybrid structure without adsorbed molecule, and the analyte molecule, respectively. We find for H₂O, NH₃, and NO₂ adsorption energies of -0.09 , -0.11 , and -0.13 eV, respectively. According to charge transfer analyses, the graphene/MoS₂ hybrid structure with S-termination is p-doped by H₂O (0.01 electrons) and NO₂ (0.21 electrons), and n-doped by NH₃ (0.01 electrons). The higher adsorption energy of MoS₂ towards the NO₂ molecule assists in more efficient charge transfer compared to the NH₃ molecule. The high susceptibility to p-doping agrees with our experimental findings by Raman spectroscopy. While the graphene/MoS₂ hybrid structure with Mo-termination shows similar tendencies towards the analyte molecules, see Fig. S8 (ESI[†]), the higher adsorption energies (-1.20 eV for H₂O, -1.55 eV for

NH₃, and -4.42 eV for NO₂) compromise the regeneration of the detector, see Table S1 (ESI[†]) for details.

The dynamic gas sensing response of the hybrid structure was investigated using NO₂ and NH₃ gases exposed to the device at controlled concentrations (Fig. 5). The investigation is conducted by cyclic exposure of the gas analyte to monitor the sensitivity, selectivity, and recovery (on each cycle) of the device. Fig. 5a and b show the typical dynamic sensing responses obtained for the hybrid structure for NO₂ and NH₃ analytes as the electron acceptor and donor molecules, respectively, demonstrating the selectivity of the device. During NO₂ (NH₃) exposure, the resistance of the hybrid-based chemiresistor decreases (increases). It reveals the different charge transport properties governed by NO₂ and NH₃ analytes that influence the overall conductivity of the device. After the adsorption of the oxidising gas NO₂, the electron concentration in MoS₂ decreases through charge transfer as shown in Fig. 4. Interaction with the unpaired electron on the N-atom shifts the Fermi level of MoS₂ into its valence band (indicated by Fig. 3a and b), extracting electron charge from MoS₂. In contrast, the lone pair of electrons of NH₃ does not have such an effect.¹⁰

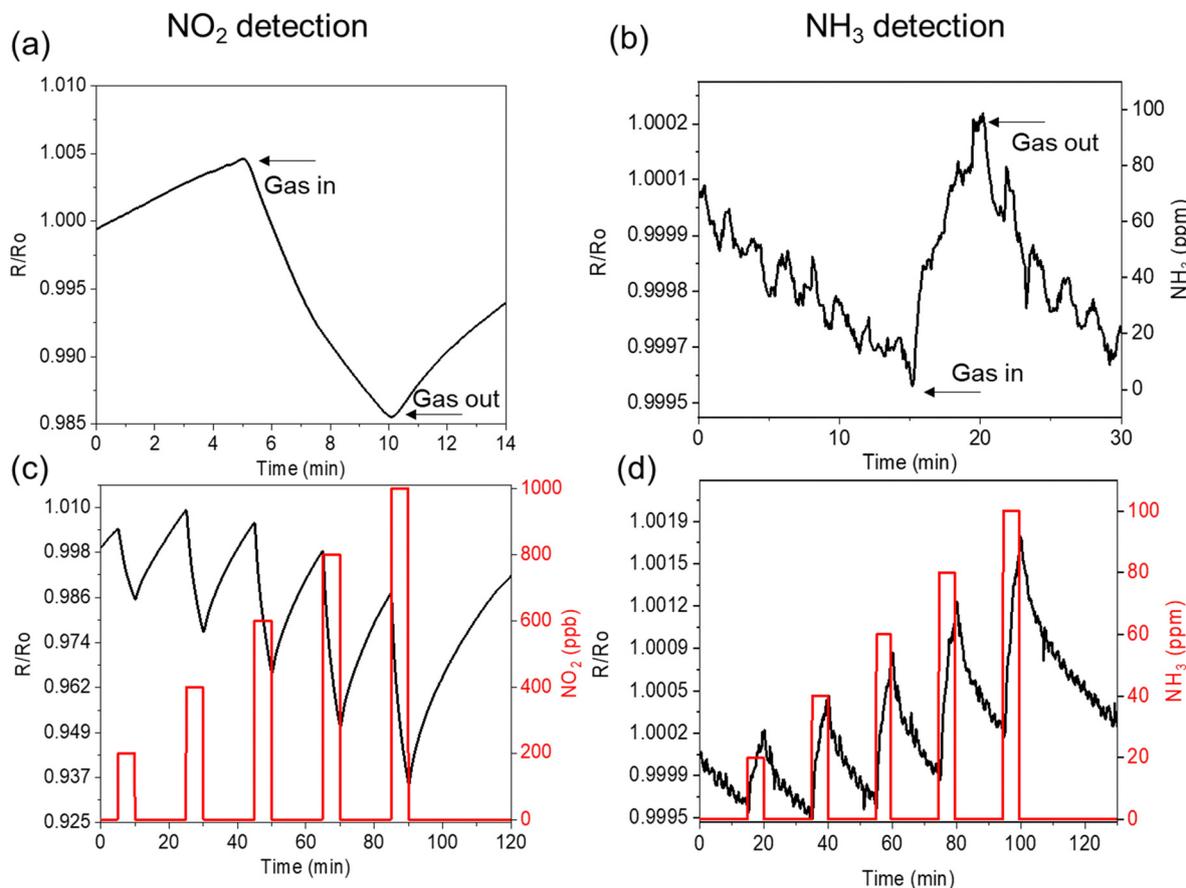


Fig. 5 Dynamic response of hybrid structure for different concentrations of NO₂ and NH₃. Resistance (R/R_0) profile of the chemiresistor against (a) NO₂ gas and (b) NH₃ gas. The arrows show the inlet and outlet of the purged gas, which influence the resistance of the chemiresistor. Resistance profile (R/R_0) at 150 °C for different concentrations (ppb) of NO₂ (c) and NH₃ (ppm) (d). The exposure of NO₂ gasses decreases resistance, while in the presence of NH₃ enhances the resistance indicating p-type behavior of the hybrid structure.



These results are in good agreement with our previous findings of p-type doping of vertically oriented MoS₂ sheets over carbon nanotubes towards NO₂ gas.¹⁷ Several research groups demonstrated and discussed similar p-type semiconducting behaviour of MoS₂ towards NO₂.^{10,30}

In the chemiresistor set-up, a slight baseline drift in the trend of the resistance profile was observed; the change in resistance during drifting is nearly 15 Ω for NO₂ and ~1 Ω for NH₃ after exposure of 90 minutes (Fig. 5c and d). This drift in the resistance profile is usually experienced in carbon-based gas sensors operating at low or moderate working temperatures, which are not enough to achieve efficient and quick desorption of the analytes during the recovery phases.⁴⁵ Several approaches can easily recover the original baseline in a potential real application. For instance, after a few hours of use, an increase in the temperature or the flow rate leads to better NO₂/NH₃ desorption, recovering the initial baseline in a few minutes. See ESI† in Fig. S9. Besides, other alternative approaches, such as the use of UV light also boost the desorption of gas molecules from the sensitive layer,^{46,48} which will be considered in future work. Unlike NO₂ resistance profile, the trend observed in NH₃ exposure shows a noisy but reproducible sub-peak (marked arrows) and slight

fluctuation at the recovery phase. It indicates the susceptibility of different basal planes ([100], [002] see TEM discussion) of the MoS₂ that could lead to a diverse range of surface and edge sites towards NH₃ molecules that differ in the selectivity with NO₂ gas analyte.

The response ($\Delta R/R_0$, %) of the individual NGF and the hybrid sensors are monitored at room temperature (RT) through purging dry air, humid conditions (relative humidity, 60%) and at elevated temperature (150 °C). The hybrid structure offers superior responses for all the testing conditions compared to its NGF counterpart. This distinction in the sensing performance is due to the abundant number of exposed edges of MoS₂ nanoplates. Therefore, analyte molecules are in higher concentration at the reactive sites in the hybrid MoS₂/NGF, thus enhancing the sensing response. There is an increment in the response with an increase in the NO₂ analyte concentration as well as for NH₃; see Fig. S10 (ESI†) for NH₃. The response is further improved in humid conditions as compared to the dried atmosphere (Fig. 6a), consistent with the findings of Sakhivel and co-workers⁴⁷ for enhanced sensitivity of graphene–MoS₂ hybrid. As discussed in the Raman spectroscopy, KPFM investigation and density functional theory simulation, the presence of NO₂ in humid conditions (*i.e.*, H₂O + NO₂)

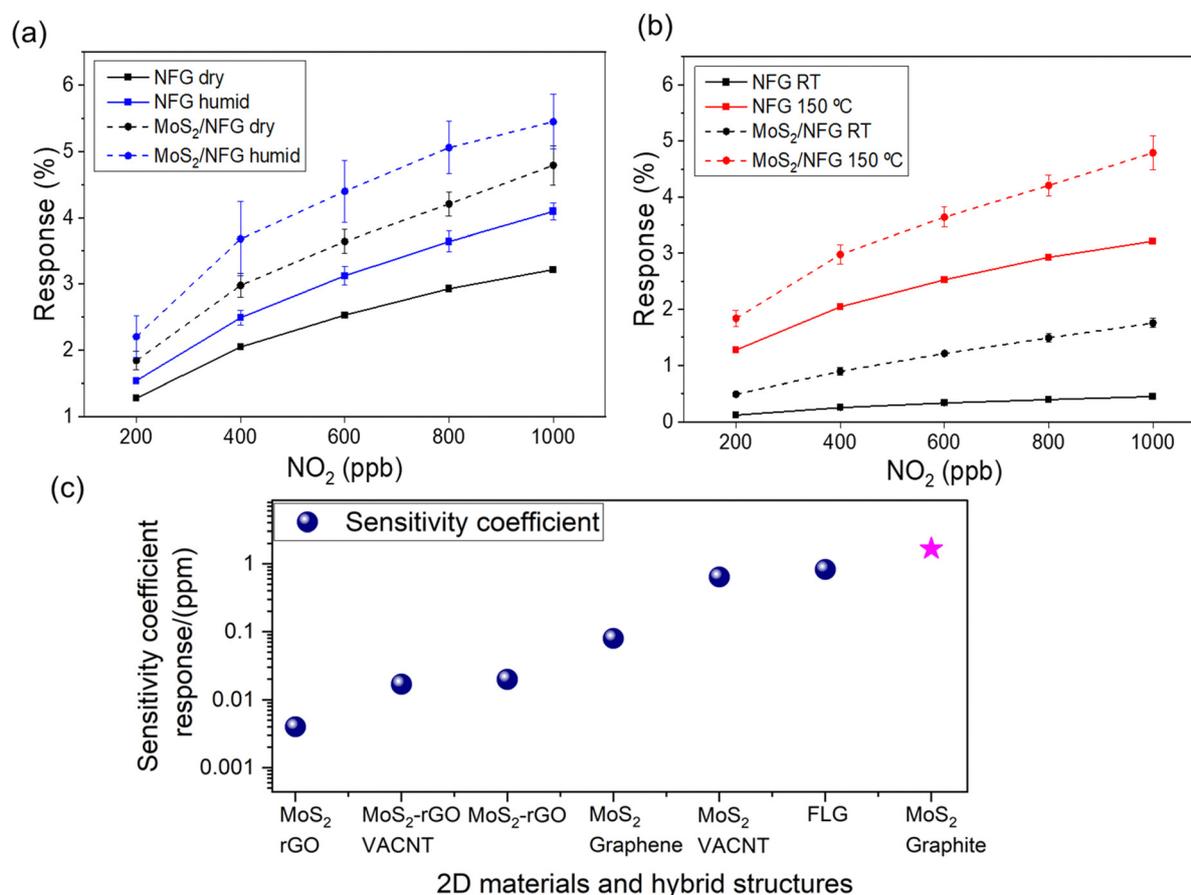


Fig. 6 Response (%) of the hybrid sensor at different conditions. (a) Effect of humidity and dried condition in the response of the graphene layers (NGF) and hybrid sensor (MoS₂/NGF). (b) Effect of temperature at 150 °C. (c) Comparison of sensing performances (Sensitivity coefficient) for NO₂ detection employing 2D materials and hybrids.



induces higher p-doped MoS₂ receptors by modulation of Fermi energy towards the valence band, consequently enhances the response. On close observation of the morphology of the MoS₂ sheet after the exposure, it is found that the edges of the MoS₂ are susceptible to salts of NO₂ while NH₄⁺ salts are distributed through the MoS₂ basal plane; see Fig. S11 (ESI†). It explains the noise in the recovery phase of the chemiresistor for ammonia gas in a humid environment owing to the crucial role of the basal plane of MoS₂ as an additional active site, which is limited for NO₂ detection.

The sensor response is usually reduced at room temperature (RT) as compared to its thermally stimulated environment, as demonstrated in ESI† (Fig. S9 and S10). This is due to lower reactivity weaker interaction between gas molecules and the 2D materials at RT.⁴⁸ The thermal activation induces better sensing kinetics, improves the recovery of the adsorbates, causes larger charge transfers for enhanced resistance changes, and influences the conduction transition.^{49–51} Therefore, moderately elevated temperatures in the range of 100–200 °C are generally applied to regulate the adsorption/desorption kinetics and improve the sensing performance in both carbon and dichalcogenide-based gas sensors.⁵¹ Here, we observed an overall enhancement in the response for both devices (*i.e.* NGF and hybrid) operated at 150 °C (Fig. 6b). Specifically, the resistance changes induced by the NO₂ exposure were increased threefold and eightfold for NGF and MoS₂/NGF samples, respectively. The sensitivity coefficient (rate of change of resistance normalised to applied concentration (sensing response/ppm) of 1.67 for NO₂ gas for the hybrid sensor outperforms other reported 2D materials (such as few-layer graphene)⁵² and hybrid heterostructures (Fig. 6c). It has been observed that several MoS₂-based devices and functionalised materials require operating temperatures (>200 °C) for sensing⁴⁸ and thus are not suited for long-term usability and stability. A brief literature survey is reported in Table S2 (ESI†); the response of pristine MoS₂ and MoS₂/carbon nanomaterial hybrids towards H₂, CO, CH₄, and H₂S has been studied, showing significantly lower responses. The present set-up operates at a moderately lower temperature (150 °C) as compared to metal oxide-based sensors attributed to the active abundant edges of MoS₂ vertical morphology on conducting substrate (NGF). This is in agreement with Cho *et al.*³⁰ results, by both experimental and theoretical studies demonstrating that edge sites of MoS₂ can adsorb five folds higher NO₂ than the basal plane MoS₂ sheets in the same amount of gas exposure time.

Conclusions

The current study presents a facile method for depositing a hybrid material composed of vertically aligned MoS₂ and conducting graphitic layers (NGF) on SiO₂/Si. The densely-packed MoS₂ flakes on NGF contribute to rationally designed edge sites for high chemiresistive sensitivity and selectivity (distinguishable p-type, n-type doping, dry-humid environment). The MoS₂/NGF hybrid, tested for NO₂, NH₃ detection in dry, humid, and

elevated temperature conditions, demonstrated superior performance compared to NGF samples. The Raman spectroscopy revealed the role of water molecules in the p-doping of MoS₂, with further improvement in the presence of nitronium ions. In addition, the Raman correlation plot distinguishes changes in total carrier concentration, NO₂ exposure indicates dedoping (*i.e.* electron transfer from MoS₂ to NO₂). Surface potential maps by KPFM technique validated this phenomenon by illustrating the modulation in the work function of MoS₂ in NH₃ and NO₂ aqueous solutions. It reveals WF increase for NO₂ exposure, indicating electron concentration decreased, consistent with p-doping of n-type MoS₂ (and *vice versa* for NH₃). Density functional theory calculations unravel the doping mechanism and structural importance (4 possible terminal ends) of hetero architecture towards different analytes and NGF. The hybrid sensor exhibited a sensitivity coefficient of 1.6 for NO₂ at 150 °C with a LOD of 50 ppb, although NH₃ measurements faced challenges due to a noisy recovery phase. Our MoS₂/NGF growth studies offer a promising strategy for designing high-performance 2D material-based heterostructures in sensing technologically relevant gases in diverse environmental applications. These findings provide fundamental insights on sensing interactions resulting from the synergy between MoS₂ and graphene, opening avenues for flexible electronics applications not achievable with traditional materials.

Materials and methods

MoS₂ synthesis

MoS₂ flakes were grown by chemical vapor deposition (CVD). 100 μL of 100 mM solution of Na₂MoO₄ (sodium molybdate: anhydrous, powder, ~100 mesh particle size, 99.9% trace metals basis; Sigma-Aldrich) was spin-coated at 3000 rpm for 1 minute onto solvent-cleaned Si/SiO₂ (n-doped Si, 280 nm thermal oxide; University Wafer) wafer chips of 1.2 × 1.2 cm size. The spin coated substrates were placed over a ceramic boat, face-up at the center of a 2 in quartz tube in a tube furnace (Lindberg Blue Tube Furnace) and 3 g of sulfur (powder, 99.98% trace metals basis; Sigma-Aldrich) was placed in a boat 15 cm upstream. The reaction was carried out under a 100 standard cubic centimetres per minute (sccm) flow of pure Ar gas. The temperature was maintained for 20 min at 25 °C to purge the tube and then ramped up to 725 °C in 15 min. The temperature was then further ramped up to 800 °C in 5 min. After reaching the 800 °C mark, the heating was shut off and the system was allowed to cool naturally to 500 °C, after which the furnace was opened to quench-cool to ambient conditions.

Graphene/MoS₂ heterostructure synthesis

NGF growth and polymer support free transfer on SiO₂/Si substrate was performed as discussed in detail in our recently published articles.^{35,36} In brief, NGF was epitaxially grown over polycrystalline Ni foils *via* the CVD method at growth temperature (900 °C), chamber pressure during annealing (10 bar) while



keeping the flow rate constant. The polymer-free transfer of grown NGF by wet-etching, where samples were kept floating overnight in 15 ml of 70% HNO₃ diluted in 600 ml of deionised (DI) water to dissolve Ni foil. The detached NGF is carefully rinsed by DI water and transferred to the target substrate (SiO₂/Si wafer).

MoS₂ nanoplates growth

To understand the effects of the growth of MoS₂ NPs on the NGF, we prepared SiO₂ (300 nm)/Si samples with NGF (half covered) transferred on it (Fig. 1). In a typical procedure, Mo films of 10 and 20 nm were deposited on both uncovered SiO₂/Si and NGF/SiO₂/Si areas of the samples. The Mo films deposited on the NGF/SiO₂/Si sample were placed in a quartz plate placed in the centre of a quartz tube. S powder was placed in two predefined zones of the tube (0.400 g were used in each zone) in order to be in the correct temperature zones when the reactor was inserted into the furnace.^{15,16} The S powder was used without any further purification. The quartz tube was then vacuumed and filled with Ar gas, and this process was repeated twice. Then the quartz tube was heated in a tube furnace at a rate of 5 °C min⁻¹ in Ar (400 sccm). When the temperature reached 850 °C, the quartz tube is inserted in the hot temperature zone so as the Mo/NGF/SiO₂/Si sample is at 850 °C while the S powder is at 250 °C. The deposition time was kept 30 min similar to previous cases of MoS₂ growth.^{15,16} After the reaction time, the tube was cooled to room temperature in Ar gas.

Characterisation techniques

Scanning electron microscopy (SEM) was performed by a Zeiss Merlin (5 kV, 50 pA) microscope. The sample surface roughness and topography was recorded *via* atomic force microscopy (AFM, Dimension Icon SPM, Bruker). The transmission electron microscopy (TEM) and high-resolution TEM (HRTEM) investigations were performed with a ThermoFisher Titan 80–300 Cubed microscope equipped with a high-brilliance field emission gun operating at 300 kV, a Wien-type FEI monochromator, and an objective lens CEOS spherical aberration corrector allowing a final spatial resolution of 0.9 Å. For both TEM and HR-TEM imaging analysis, we used a direct detection Ametek Gatan K2 IS camera. Plan-view TEM samples were prepared by scratching the specimen's surface and collecting the powder material directly onto lacey-carbon coated Cu grids. The cross-section lamella was prepared with a FEI Helios G4 focused ion beam (FIB) SEM microscope, equipped with a field emission gun and a Leica cryo-holder stage. To protect the MoS₂ NPs region of interest, we coated the sample with a 5 nm layer of Ir and 1.5 μm of Pt. Then, the milling procedure was carried out, with the cryo-stage set at -145 °C and by the use of an ion beam of 16 kV (to reduce sample damage and artifacts formation) with a current range of 10 nA to 12 pA. Once completed, the thin lamella was slowly warmed up to room temperature.

Kelvin Probe force microscopy

Kelvin Probe Force Microscopy (KPFM) is carried out at room temperature by Bruker AFM (model: ScanAsyst) using a

cantilever (model: PFQNE-AL). A single layer of the CVD MoS₂ has been used for the investigation, where it is treated with different analytes and surface potential has been measured. The KPFM measurement is carried through the two-pass technique, where the first pass is used to scan the topography of the sample, and the second pass is carried out at a specific lift height to measure the contact potential difference (CPD, mV). The reference sample (Au–Si–Al) has been used to calibrate the cantilever probe apex and measure the sample's work function (eV).

Raman spectroscopy

Raman spectroscopy (model: Renishaw) has been carried out using laser line ($\lambda = 532$ nm, grating 1800 l/mm and magnification 100×) over the samples. The spectra and mapping were carried out at a similar location for the pre and post-treated samples. More than 500 Raman spectra were carried out using different samples. All Raman shifting values were measured through Lorentzian fitting of the Raman modes and compared to its native state. Si sample was used for the calibration of the Raman spectrometer. A separate measurement is carried out using spectrometer (Alpha 300 RA, WITeC, 532 nm laser, 100×, lateral resolution 0.5) to verify the peak position and Raman shifts.

It is worth noting that the surface potential mapping and Raman spectroscopy is carried out in the presence of ammonium and nitronium ions, where salts are considered in the presence of water molecules. Our KPFM and Raman spectroscopy set-up is limited to study for gaseous analytes.

X-ray diffraction

The X-ray diffraction (XRD) analysis has been performed on NGF/SiO₂/Si before and after MoS₂ NPs growth with a powder a Bruker diffractometer equipped with D2 phaser and LYNXEYE detector and operating with a Cu radiation source.

X-ray photoelectron spectroscopy (XPS)

The chemical properties of the sample have been investigated by X-ray photoelectron spectroscopy, performed on Thermo K-Alpha spectrometer using Al K α source (1486.68 eV) and a spot size of 250 × 250 μm². Before analysis, the sample surface has been cleaned with a cluster beam (Ar₅₀₀⁺, 2 keV, 120 s, raster size of 1.25 × 1.25 mm²). The high-resolution spectra of C 1s, O 1s, S 2p, and Mo 3d levels are recorded at pass energy of 20 eV, with 20 scans. The atomic concentrations were determined from the area of each peak, considering a Shirley background.

Gas-sensing device fabrication and measurements

The MoS₂/NGF/SiO₂/Si samples were connected to 20 × 30 mm printed circuit boards (PCBs) through two platinum wire contacts made on the sensor surface using conductive silver paste from Merck KGaA (Darmstadt, Germany). This silver paste is often used to electrically contact devices to be operated in harsh conditions such as ultra-high vacuum or under heating (*e.g.* in synchrotron experiments), as it does not degases. Sample resistance was measured in a two-contact (two-probe)



configuration. Namely, the multimeter applies a constant voltage between the two contacts and measures the current. In such a simple configuration, the resistance measured comprises the sample resistance and the contact resistance. The platinum screen-printed alumina heating element was connected at the backside of the sample using a thermally conducting epoxy paste.

Both sensors NGF and MoS₂/NGF/SiO₂/Si were placed in an airtight Teflon testing chamber, with a volume of 35 cm³. This chamber was connected to a gas delivery system, which comprised two calibrated gas cylinders, one with pure dry air (Air Premier Purity: 99,995%) and the other with 1 ppm of NO₂ in a balance of dry air. Different dilutions of the target gases were performed to expose the sensors at different concentrations of NO₂, NH₃ in the range of parts per billion (ppb). Specifically, the gas sensors were stabilised under synthetic dry air for 15 minutes before every gas exposure and then exposed to a given concentration of NO₂ for 5 minutes. Besides, the effect of the ambient moisture on the sensing properties was also assessed. A controller evaporator mixer (CEM) from Bronkhorst High-Tech B.V. (Ruurlo, The Netherlands) was used to humidify the gas mixture, and the relative humidity level was monitored by using a humidity sensor from Sensirion AG (Stäfa, Switzerland) placed inside the Teflon chamber. When the sensors were operated well above room temperature, they were connected to a Keysight E3641A (Santa Rosa, CA, U.S.) power supply.

In order to operate the system under similar conditions to those needed for ambient monitoring applications, a low flow rate was applied, avoiding the use of pumps. In particular, a set of Bronkhorst High-Tech B.V. (Ruurlo, The Netherlands) mass-flows controllers were used to establish a flow rate of 100 mL min⁻¹. Finally, the changes in sensor resistance were monitored using an Agilent HP 34972A multimeter connected to the sensing chamber. The sensor responses were defined as ($\Delta R/R_0$) expressed in percentage. Where ΔR is the resistance change over the five minutes of gas exposure, while R_0 corresponds to the baseline resistance obtained under the pure air atmosphere.

Sensitivity coefficient measurement. Sensitivity is given by the slope of the calibration curve. Considering $y = mx + b$, the slope “ m ” represents the response per unit of concentration.

Thereby, the slope of our sensor at 150 °C for measuring NO₂ at the ppb level results in a 0.00167 response per ppb ($(0.00167 \text{ response/ppb}) \times (1000 \text{ ppb/1 ppm}) = 1.67 \text{ response/ppm}$)

Computational method

First-principles calculations are performed using density functional theory (Vienna Ab-initio Simulation Package) and the generalised gradient approximation of Perdew–Burke–Ernzerhof with DFT-D3 van der Waals correction.⁵³ The plane wave cutoff energy is set to 500 eV. The total energy is converged to 10⁻⁵ eV and the atomic forces are converged to 0.01 eV Å⁻¹. A Monkhorst–Pack k -sampling with 0.015 Å⁻¹ spacing is used. To achieve a small lattice mismatch (2% in the b -direction), we build our simulation cells from the supercells of graphene and MoS₂ shown in Fig. S7a and b (ESI†). Vacuum slabs of at least

15 Å thickness (in the c -direction) avoid unphysical interaction due to the periodic boundary conditions. Charge transfer analyses are performed by the Bader method.⁵⁴

Conflicts of interest

There are no conflicts to declare.

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References

- 1 I. Manisalidis, E. Stavropoulou, A. Stavropoulos and E. Bezirtzoglou, *Front. Public Health*, 2020, **8**, 14.
- 2 D. J. Buckley, N. C. Black, E. G. Castanon, C. Melios, M. Hardman and O. Kazakova, *2D Mater.*, 2020, **7**, 032002.
- 3 S. M. Majhi, A. Mirzaei, H. W. Kim, S. S. Kim and T. W. Kim, *Nano Energy*, 2021, **79**, 105369.
- 4 A. Dey, *Mater. Sci. Eng. B*, 2018, **229**, 206–217.
- 5 H. Chai, Z. Zheng, K. Liu, J. Xu, K. Wu, Y. Luo, H. Liao, M. Debliquy and C. Zhang, *IEEE Sens. J.*, 2022, **22**, 5470–5481.
- 6 B. Cho, J. Yoon, S. K. Lim, A. R. Kim, D.-H. Kim, S.-G. Park, J.-D. Kwon, Y.-J. Lee, K.-H. Lee and B. H. Lee, *ACS Appl. Mater. Interfaces*, 2015, **7**, 16775–16780.
- 7 U. Krishnan, M. Kaur, K. Singh, M. Kumar and A. Kumar, *Superlattices Microstruct.*, 2019, **128**, 274–297.
- 8 T. Zeng, Y. You, X. Wang, T. Hu and G. Tai, *Prog. Chem.*, 2016, **28**, 459.
- 9 K. F. Mak, C. Lee, J. Hone, J. Shan and T. F. Heinz, *Phys. Rev. Lett.*, 2010, **105**, 136805.
- 10 R. Kumar, W. Zheng, X. Liu, J. Zhang and M. Kumar, *Adv. Mater. Technol.*, 2020, **5**, 1901062.
- 11 F. Schedin, A. K. Geim, S. V. Morozov, E. W. Hill, P. Blake, M. I. Katsnelson and K. S. Novoselov, *Nat. Mater.*, 2007, **6**, 652–655.
- 12 G. Deokar, N. S. Rajput, J. Li, F. L. Deepak, W. Ou-Yang, N. Reckinger, C. Bittencourt, J.-F. Colomer and M. Jouiad, *Beilstein J. Nanotechnol.*, 2018, **9**, 1686–1694.
- 13 Y. Huang, J. Guo, Y. Kang, Y. Ai and C. M. Li, *Nanoscale*, 2015, **7**, 19358–19376.
- 14 A. Allain and A. Kis, *ACS Nano*, 2014, **8**, 7180–7185.



- 15 G. Deokar, N. S. Rajput, P. Vancsó, F. Ravoux, M. Jouiad, D. Vignaud, F. Cecchet and J. F. Colomer, *Nanoscale*, 2017, **9**, 277–287.
- 16 G. Deokar, D. Vignaud, R. Arenal, P. Louette and J.-F. Colomer, *Nanotechnology*, 2016, **27**, 075604.
- 17 G. Deokar, P. Vancso, R. Arenal, F. Ravoux, J. Casanova-Cháfer, E. Llobet, A. Makarova, D. Vyalikh, C. Struzzi and P. Lambin, *Adv. Mater. Interfaces*, 2017, **4**, 1700801.
- 18 Q. H. Wang, K. Kalantar-Zadeh, A. Kis, J. N. Coleman and M. S. Strano, *Nat. Nanotechnol.*, 2012, **7**, 699–712.
- 19 Y. H. Lee, X. Q. Zhang, W. Zhang, M. T. Chang, C. T. Lin, K. D. Chang, Y. C. Yu, J. T. W. Wang, C. S. Chang and L. J. Li, *Adv. Mater.*, 2012, **24**, 2320–2325.
- 20 Y.-S. Shim, K. C. Kwon, J. M. Suh, K. S. Choi, Y. G. Song, W. Sohn, S. Choi, K. Hong, J.-M. Jeon and S.-P. Hong, *ACS Appl. Mater. Interfaces*, 2018, **10**, 31594–31602.
- 21 F. Xiong, Z. Cai, L. Qu, P. Zhang, Z. Yuan, O. K. Asare, W. Xu, C. Lin and L. Mai, *ACS Appl. Mater. Interfaces*, 2015, **7**, 12625–12630.
- 22 W. Li, Y. You and J.-H. Choi, *Surf. Sci.*, 2021, **707**, 121809.
- 23 E. Lee, S. G. Lee, W. H. Lee, H. C. Lee, N. N. Nguyen, M. S. Yoo and K. Cho, *Chem. Mater.*, 2020, **32**, 4544–4552.
- 24 Y. Chen and M. Sun, *Nanoscale*, 2021, **13**, 5594–5619.
- 25 Y. Yuan, X. Zhang, H. Liu, T. Yang, W. Zheng, B. Zheng, F. Jiang, L. Li, D. Li and X. Zhu, *J. Alloys Compd.*, 2020, **815**, 152309.
- 26 X. Yu, G. Zhao, S. Gong, C. Liu, C. Wu, P. Lyu, G. Maurin and N. Zhang, *ACS Appl. Mater. Interfaces*, 2020, **12**, 24777–24785.
- 27 L. Yu, Y.-H. Lee, X. Ling, E. J. Santos, Y. C. Shin, Y. Lin, M. Dubey, E. Kaxiras, J. Kong and H. Wang, *Nano Lett.*, 2014, **14**, 3055–3063.
- 28 B. Sachs, L. Britnell, T. Wehling, A. Eckmann, R. Jalil, B. Belle, A. Lichtenstein, M. Katsnelson and K. Novoselov, *Appl. Phys. Lett.*, 2013, **103**, 251607.
- 29 W. Zan, W. Geng, H. Liu and X. Yao, *J. Alloys Compd.*, 2015, **649**, 961–967.
- 30 S.-Y. Cho, S. J. Kim, Y. Lee, J.-S. Kim, W.-B. Jung, H.-W. Yoo, J. Kim and H.-T. Jung, *ACS Nano*, 2015, **9**, 9314–9321.
- 31 G. Ciampalini, F. Fabbri, G. Menichetti, L. Buoni, S. Pace, V. Mišeikis, A. Pitanti, D. Pisignano, C. Coletti and A. Tredicucci, *ACS Nano*, 2021, **16**, 1291–1300.
- 32 Y. Shi, W. Zhou, A.-Y. Lu, W. Fang, Y.-H. Lee, A. L. Hsu, S. M. Kim, K. K. Kim, H. Y. Yang and L.-J. Li, *Nano Lett.*, 2012, **12**, 2784–2791.
- 33 M. Chatti, T. Gengenbach, R. King, L. Spiccia and A. N. Simonov, *Chem. Mater.*, 2017, **29**, 3092–3099.
- 34 Z. Mao, H. Zhu, X. Peng, J. Chen, Q. Chen, X. Chen, X. Hu and H. Chen, *Microchim. Acta*, 2022, **189**, 155.
- 35 G. Deokar, A. Genovese and P. M. Costa, *Nanotechnology*, 2020, **31**, 485605.
- 36 G. Deokar, A. Genovese, S. G. Surya, C. Long, K. N. Salama and P. M. F. J. Costa, *Sci. Rep.*, 2020, **10**, 14703.
- 37 M.-C. Chang, P.-H. Ho, M.-F. Tseng, F.-Y. Lin, C.-H. Hou, I. K. Lin, H. Wang, P.-P. Huang, C.-H. Chiang and Y.-C. Yang, *Nat. Commun.*, 2020, **11**, 3682.
- 38 J. Kong, K. T. Park, A. Miller and K. Klier, *Surf. Sci. Spectra*, 2000, **7**, 69–74.
- 39 D. Ganta, S. Sinha and R. T. Haasch, *Surf. Sci. Spectra*, 2014, **21**, 19–27.
- 40 X. Ai, H. Fan, Y. Wang, Y. Guo, X. Liu, L. Yang, H. Liu and J. Yang, *RSC Adv.*, 2018, **8**, 36280–36285.
- 41 T.-X. Huang, X. Cong, S.-S. Wu, K.-Q. Lin, X. Yao, Y.-H. He, J.-B. Wu, Y.-F. Bao, S.-C. Huang and X. Wang, *Nat. Commun.*, 2019, **10**, 5544.
- 42 F. Lee, M. Tripathi, R. Sanchez Salas, S. P. Ogilvie, A. Amorim Graf, I. Jurewicz and A. B. Dalton, *Nanoscale*, 2023, **15**, 7227–7248.
- 43 M. Tripathi, F. Lee, A. Michail, D. Anastopoulos, J. G. McHugh, S. P. Ogilvie, M. J. Large, A. A. Graf, P. J. Lynch and J. Parthenios, *ACS Nano*, 2021, **15**, 2520–2531.
- 44 J. H. Kim, J. Lee, J. H. Kim, C. Hwang, C. Lee and J. Y. Park, *Appl. Phys. Lett.*, 2015, **106**, 251606.
- 45 H. Du, G. Xie, Y. Su, H. Tai, X. Du, H. Yu and Q. Zhang, *Sensors*, 2019, **19**, 889.
- 46 A. Di Bartolomeo, A. Kumar, O. Durante, A. Sessa, E. Faella, L. Viscardi, K. Intonti, F. Giubileo, N. Martucciello, P. Romano, S. Sleziona and M. Schleberger, *Mater. Today Nano*, 2023, **24**, 100382.
- 47 R. Sakthivel, A. Geetha, B. Anandh, V. Jagadeesan, A. S. Ganesh and J. Dineshkumar, *J. Phys.: Conf. Ser.*, 2021, **2070**, 012131.
- 48 P. Srinivasan, M. Ezhilan, A. J. Kulandaisamy, K. J. Babu and J. B. B. Rayappan, *J. Mater. Sci.: Mater. Electron.*, 2019, **30**, 15825–15847.
- 49 D. Mouloua, A. Kotbi, G. Deokar, K. Kaja, M. El Marssi, M. A. El Khakani and M. Jouiad, *Materials*, 2021, **14**, 3283.
- 50 R. Kumar, X. Liu, J. Zhang and M. Kumar, *Nano-Micro Lett.*, 2020, **12**, 1–37.
- 51 H. Xie, K. Wang, Z. Zhang, X. Zhao, F. Liu and H. Mu, *RSC Adv.*, 2015, **5**, 28030–28037.
- 52 G. Deokar, J. Casanova-Chafer, N. S. Rajput, C. Aubry, E. Llobet, M. Jouiad and P. M. Costa, *Sens. Actuators, B*, 2020, **305**, 127458.
- 53 G. Kresse and D. Joubert, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1999, **59**, 1758.
- 54 W. Tang, E. Sanville and G. Henkelman, *J. Phys.: Condens. Matter*, 2009, **21**, 084204.

